J/ψ and ψ(2S) production in small systems with PHENIX

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Introduction

Recent quarkonium analyses in PHENIX have focused on small systems results for J/ ψ and ψ (2S).

p+Al, p+Au and ³He+Au data from the 2014 and 2015 RHIC runs have been added to our d+Au data from 2008.

This allows us to address questions about J/ ψ and ψ (2S) production:

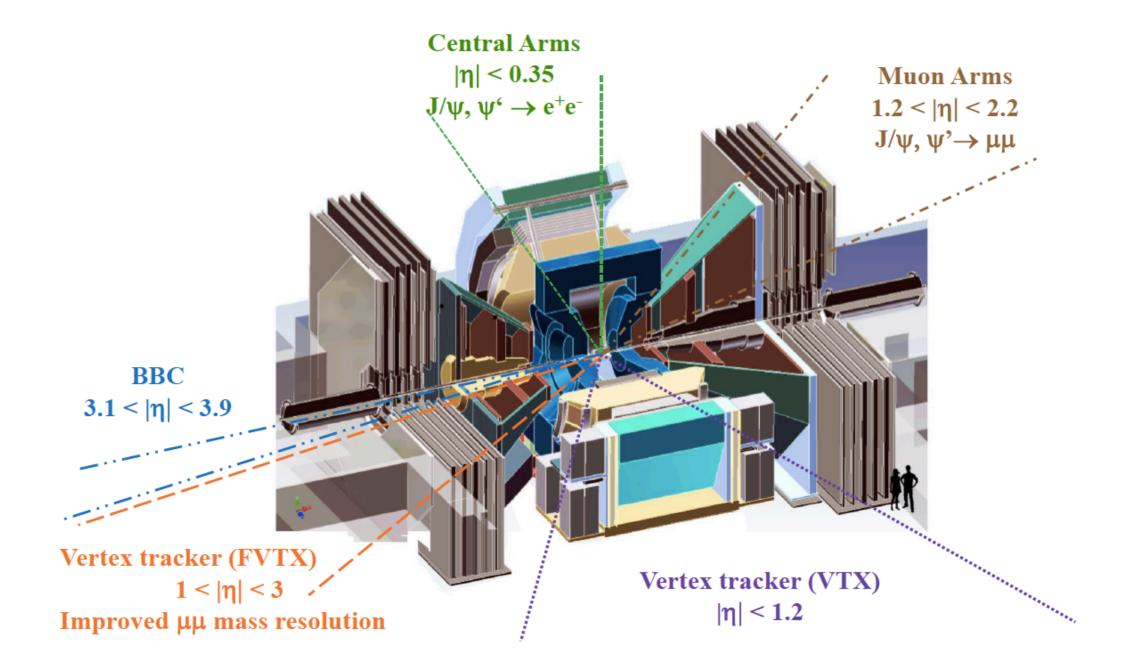
- Do we see evidence of final state effects in light systems?
- How well do we understand quarkonia modification in light systems?

The results shown here are from:

Measurement of $\psi(2S)$ nuclear modification at backward and forward rapidity in p+p, p+AI, and p+Au collisions at $\sqrt{sNN} = 200$ GeV, PHENIX Collaboration (arXiv:2202.03863)

Measurement of J/ ψ at forward and backward rapidity in p + p, p + Al, p + Au, and 3He+Au collisions at $\sqrt{sNN} = 200$ GeV, PHENIX collaboration, *Phys.Rev.C* 102 (2020) 1, 014902

Quarkonia - PHENIX

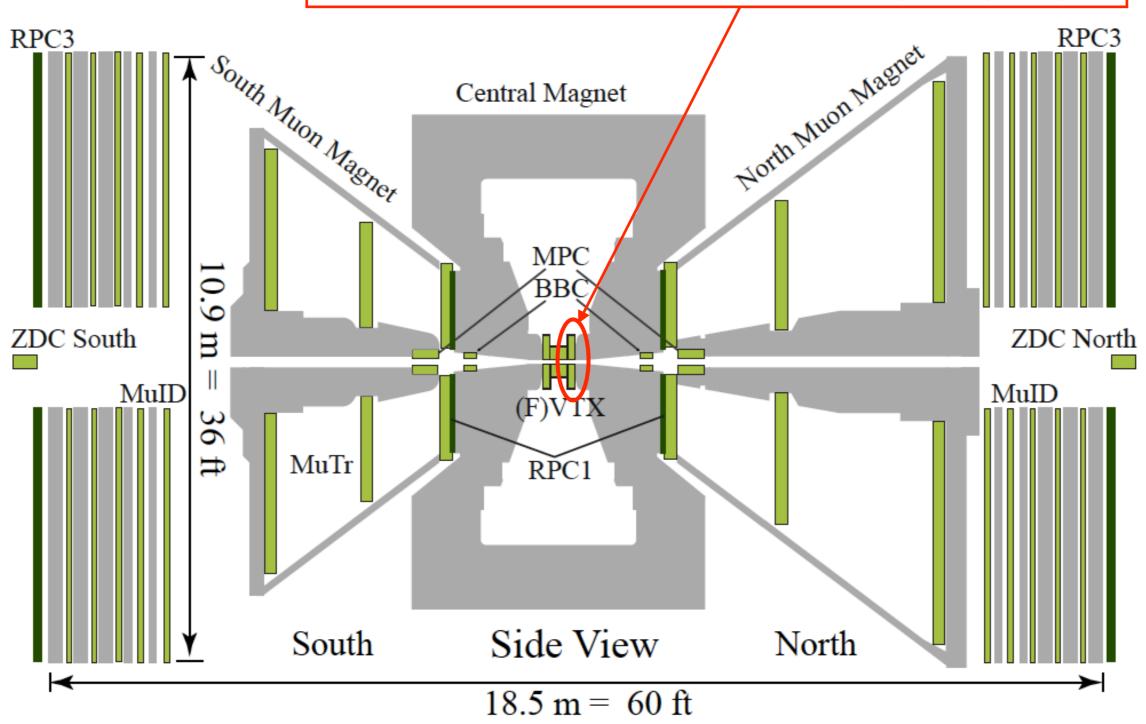


Di-leptons measured in y = (-2.2, -1.2) (-0.35, 0.35) (1.2, 2.2)

- No triggering required in Au+Au
- Efficient triggers in p+p, (p,d,³He)+A

The muon arms

FVTX provides precise space point before passing through magnet



Quarkonia production in a nucleus

Processes that modify the quarkonia yield in a nuclear target - called cold nuclear matter (CNM) processes. $R_G^{\rm Pb}$

Gluon shadowing - parton distributions are modified in a nucleus

Absorption - breakup of the precursor quarkonium by collision with a target nucleon

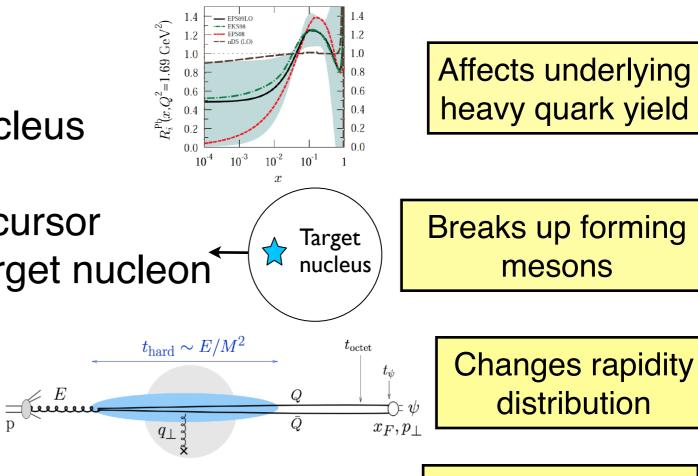
Initial state energy loss of a parton in cold nuclear matter

Cronin effect - multiple elastic scattering of partons

There is also a possibility that quarkonium states may be broken up in the final state by interactions with particles produced in the collision.

Modifies the p_T

distribution



Shadowing

Recent shadowing parameterizations

- EPPS16 (Eskola et. al., Eur. Phys. J. C 77, 163 (2017))
- nCTEQ15 (Kovarik et. al., Phys. Rev. D 93, 085037 (2016))

Bayesian re-weighting of EPPS16 and nCTEQ15 gluon nPDF's

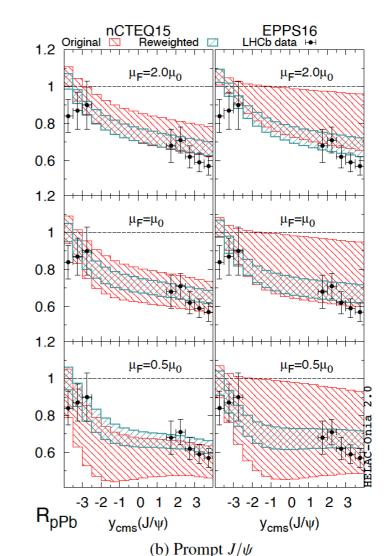
- (Kusina et. al., Phys. Rev. Lett. 121, 052004 (2018))
- Adds LHC pPb data gluon dominated processes

• D₀, J/ ψ , B \rightarrow J/ ψ , and Y(1S) mesons See also Eskola et. al. arXiv:1906.03943, and *Nucl.Phys.A* 1005 (2021) 121944.

- Considerably narrows uncertainty band
- Reduces R_g at forward rapidity
- "Absorbs" initial state energy loss into nPDF?

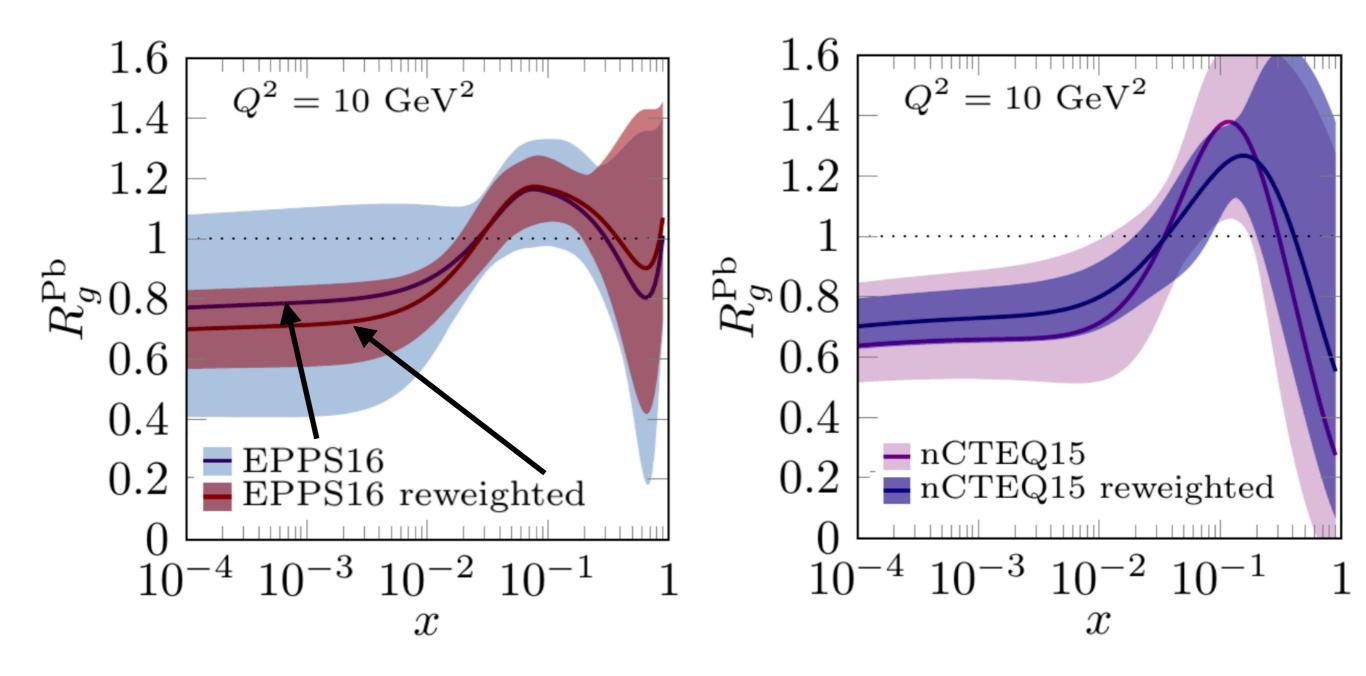
Fitted to centrality integrated data only

- Has no information about centrality dependence
- Centrality dependence has to be invented



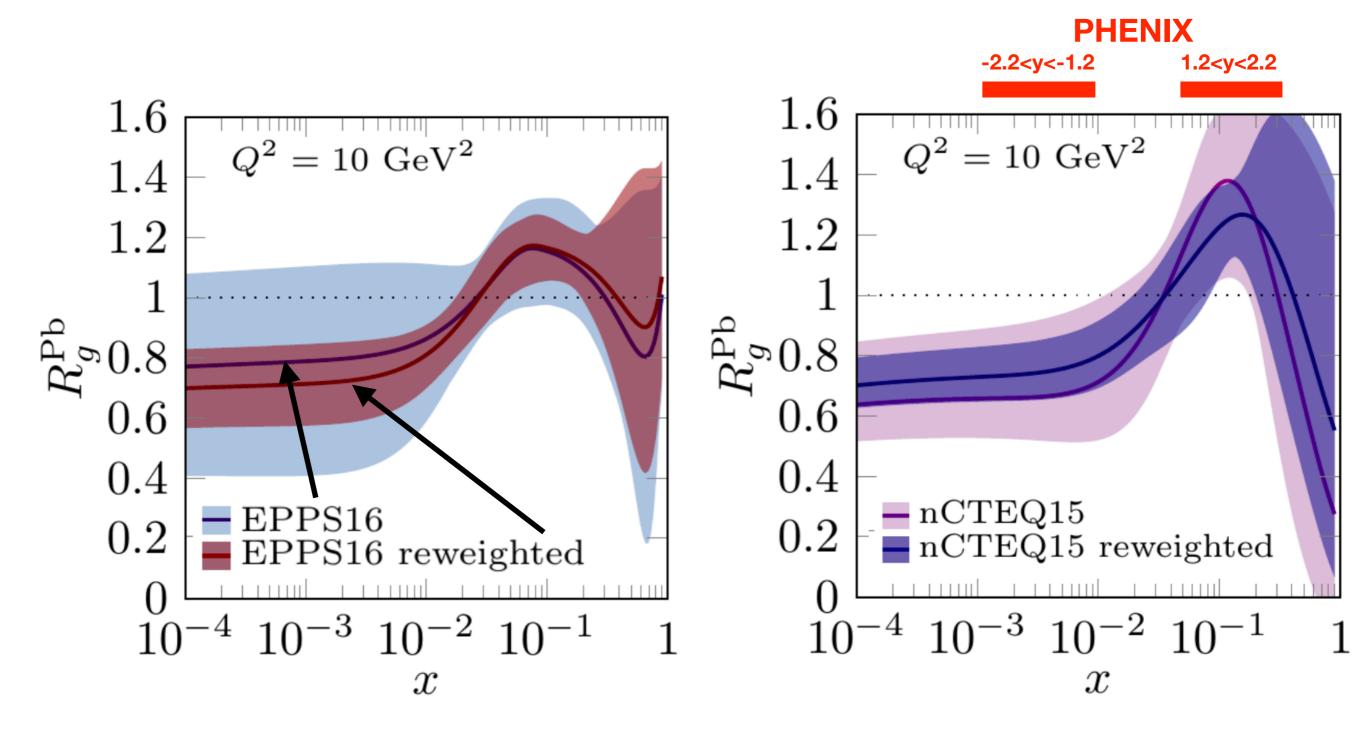
Impact of re-weighting of gluon nPDF's

Eskola et. al., arXiv:1906.03943, Nucl. Phys. A 1005 (2021) 121944
Hesssian re-weighting of nPDF's using LHCB D⁰ data.



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Initial state energy loss

Incoming parton loses energy due to gluon radiation associated with p_T broadening.

Examples:

- Arleo et. al. JHEP 05 (2013) 155.
- Sharma and Vitev, PRC 87 (2013) 044905.
- Kopeliovich et al., Phys.Rev. C95 (2017) 065203.

The Bayesian re-weighted shadowing seems to explain p+A data reasonably well without additional effects from initial state energy loss.

 Absorbs initial state energy loss effects into the shadowing parameterization?

J/ψ absorption

Backward rapidity J/ψ in PHENIX experience a significant "absorption" cross section - in addition to substantial anti-shadowing.

Parameterized using model of expansion of color neutral charmonium precursor as it crosses the target (Arleo et. al., PRC 61, 054906 (2000)).

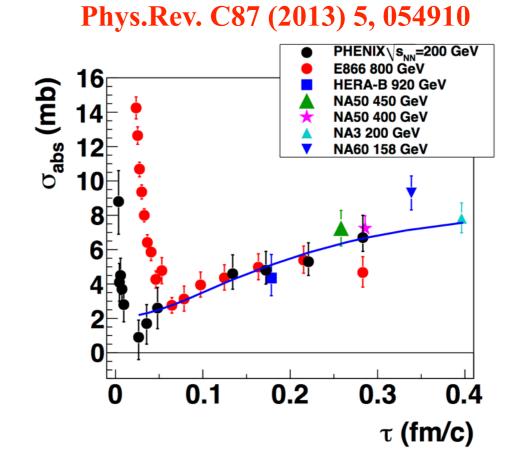
- Applied in Glauber model of collisions integrated over relevant y range.
- Fitted to world's σ_{abs} data for nuclear crossing time $\tau > 0.05$ fm/c
- All data corrected for shadowing with EKS98 or EPS09

Provides good description of $\tau > 0.05$ fm/c data from $\sqrt{s_{NN}} = 17$ to 200 GeV

Anti-shadowing parameterizations have remained stable over several generations.
Well constrained by DIS data.
Not the case for low x shadowing (low τ).

Strong absorption is not expected at LHC

Nuclear crossing times very short at all y



Transport model

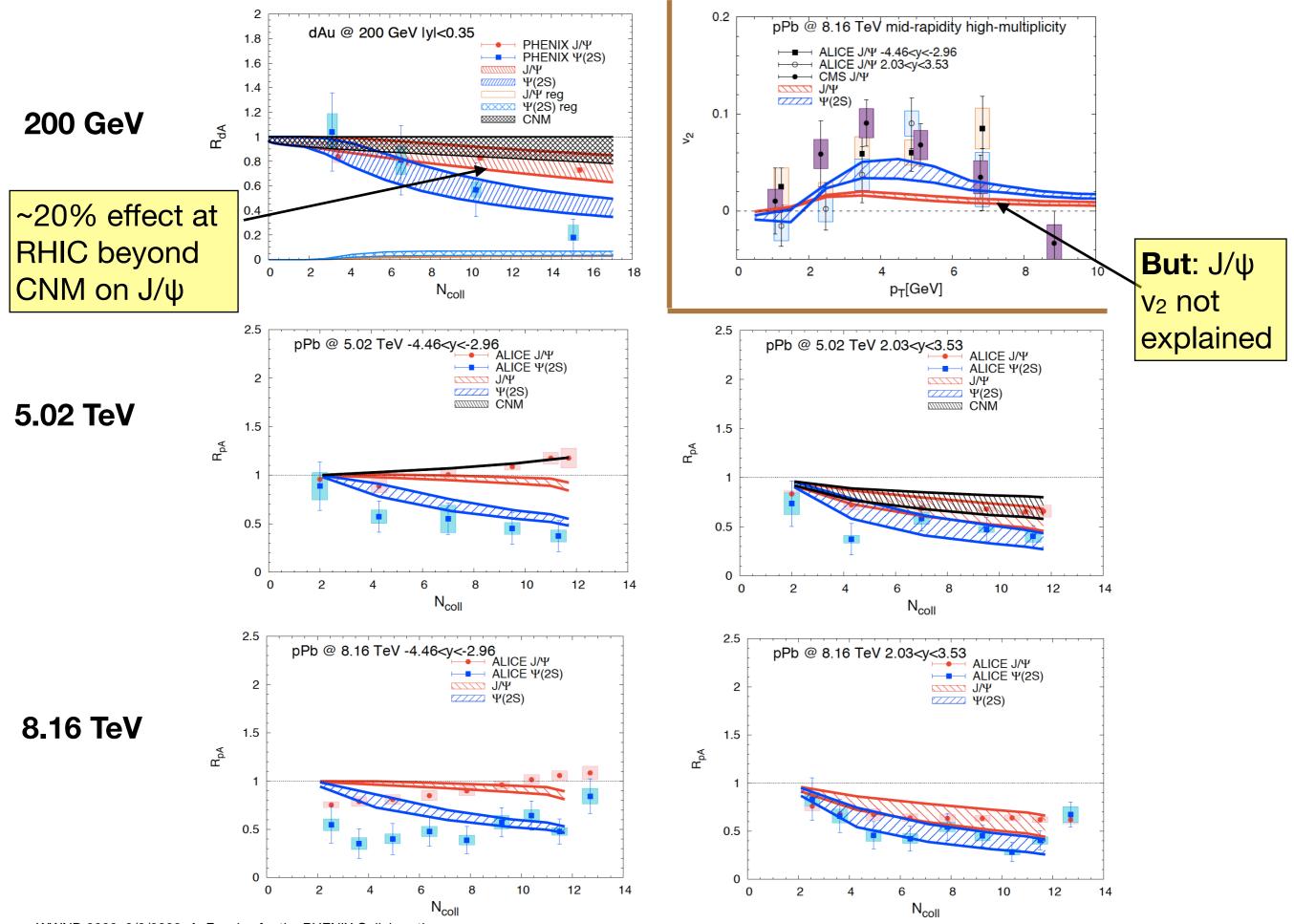
Du and Rapp (JHEP 1903 (2019) 015) adapted their transport model, used to describe heavy ion collisions, for use in small systems. They tried to describe available charmonium J/ ψ and ψ (2S)

- RHIC: PHENIX J/ ψ and ψ (2S) data at midrapidity.
- LHC: ALICE J/ ψ and ψ (2S) data at forward/backward rapidity, including the J/ ψ v₂.

The transport model uses

- A rate equation approach within a fireball model
- Initial geometry of the fireball from a Monte-Carlo event generator
- Initial anisotropies are caused by fluctuations
- Includes corrections for CNM effects
 - EPS09 shadowing with assumed linear centrality dependence
 - Assumes constant nuclear absorption at backward rapidity

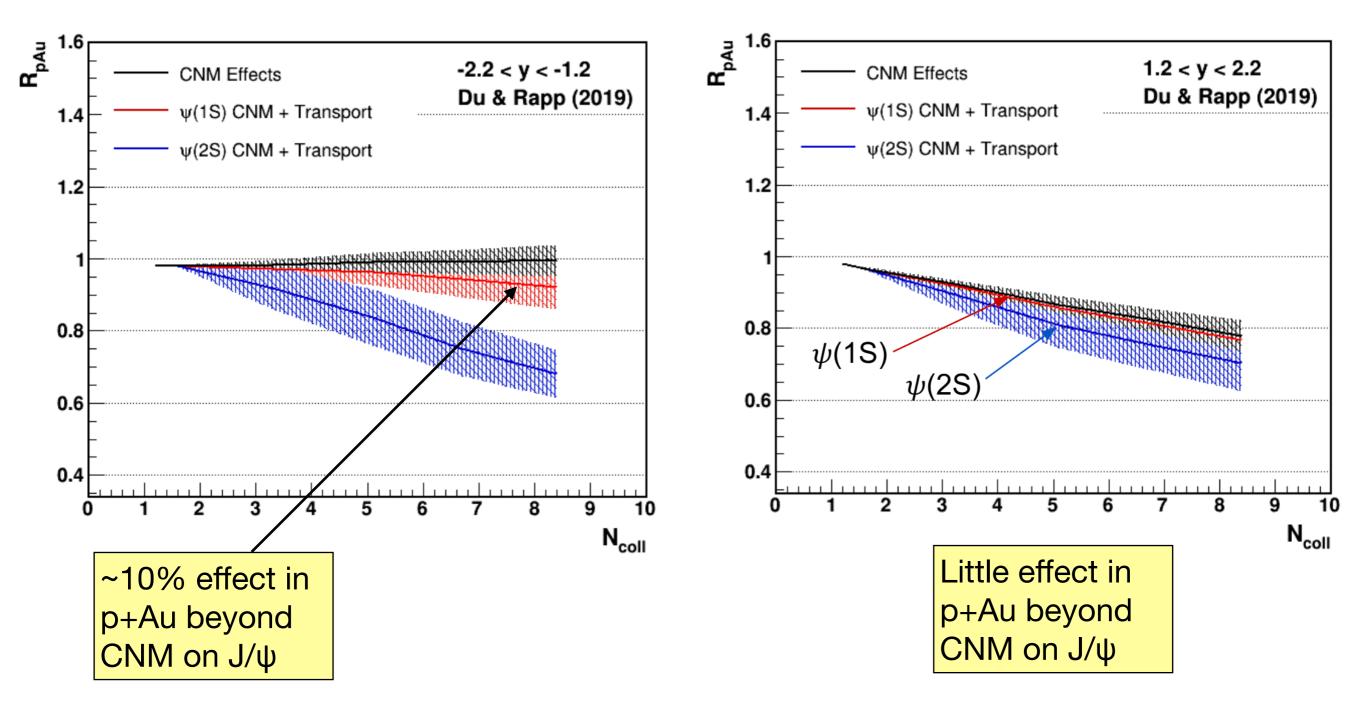
Some comparisons with data from the paper on the next slide.



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Prediction for p+Au at 200 GeV

Du and Rapp transport model prediction (really!) for 200 GeV p+Au collisions at forward and backward rapidity.

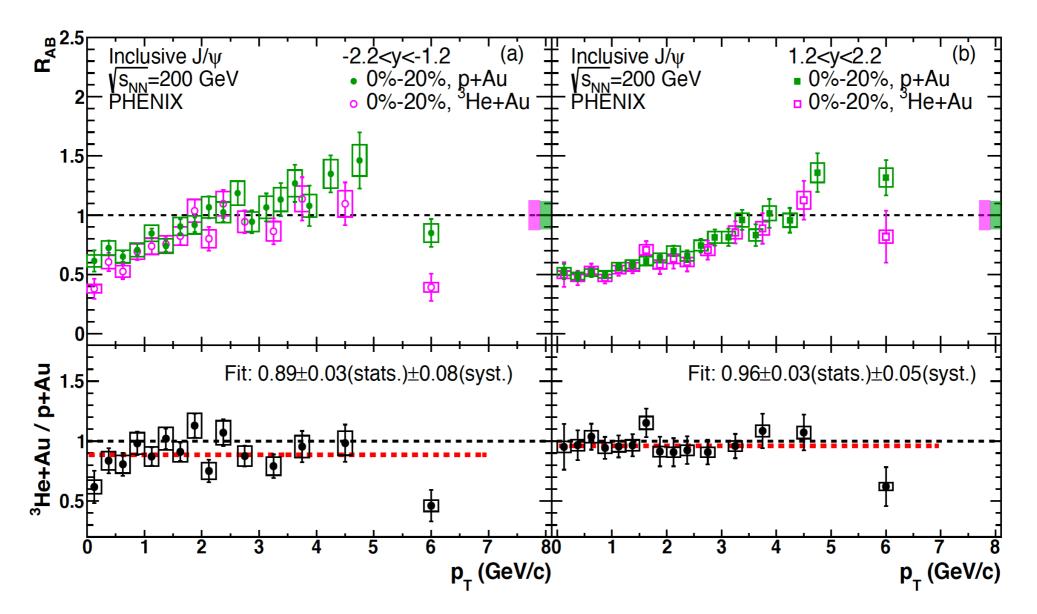


Review of J/ ψ results at forward/backward rapidity

³He+Au to p+Au ratio (0-20%centrality)

Backward rapidity ratio $0.89 \pm 0.03 \pm 0.08$

- Consistent with some additional suppression (90% probability).
- But not far outside the systematic uncertainty.
- Forward rapidity ratio $0.96 \pm 0.03 \pm 0.05$
- Consistent with 1

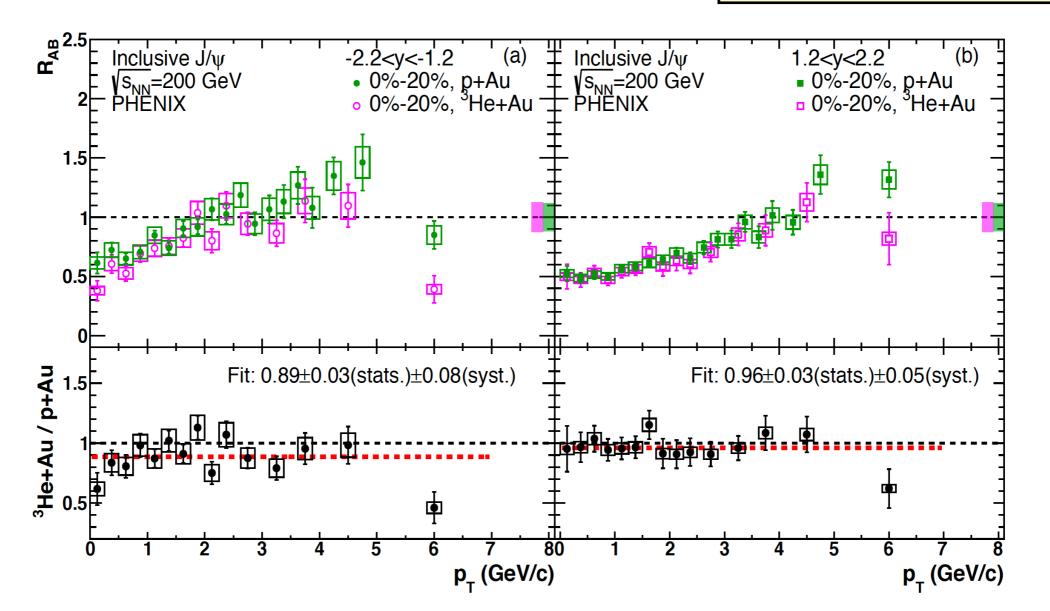


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Little evidence for strong suppression of J/ψ in final state

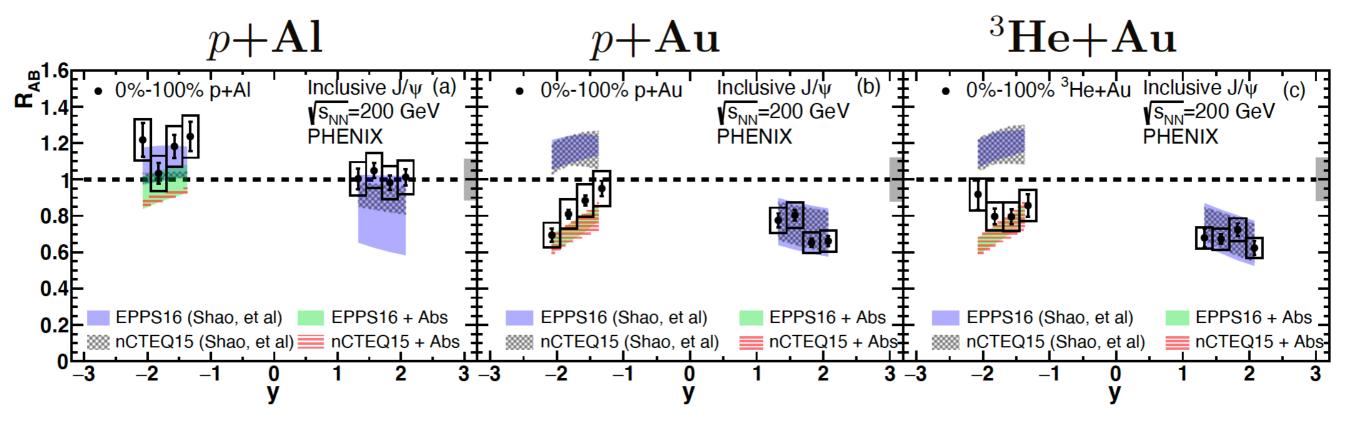


Rapidity dependence, p_T and centrality integrated

0-100% centrality, p_T integrated.

Add EPPS16 and nCTEQ15 shadowing, with Bayesian re-weighting.

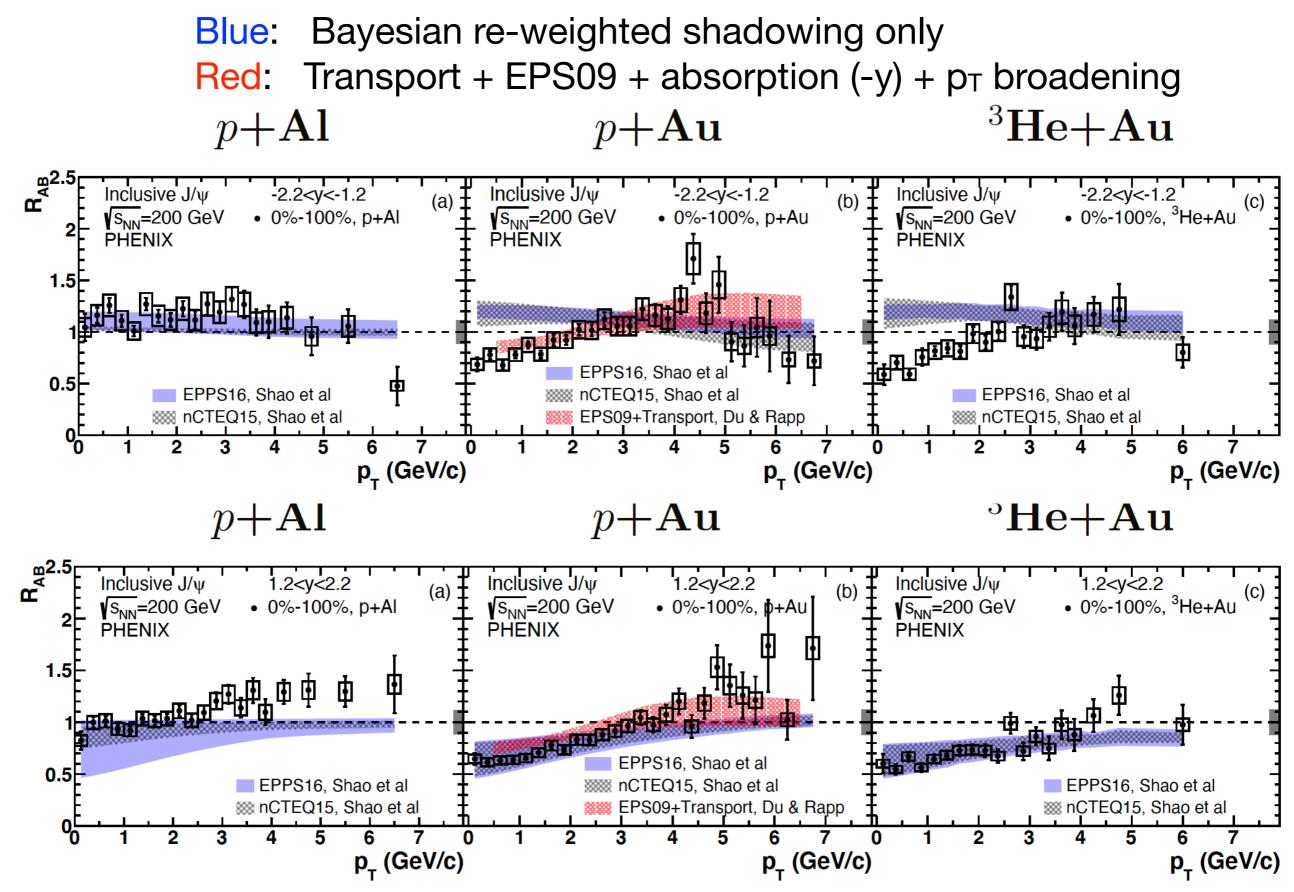
• Fold in absorption prediction with shadowing at backward rapidity.



Not so bad!

How about the p_T dependence?

p_T dependence, 0-100% centrality

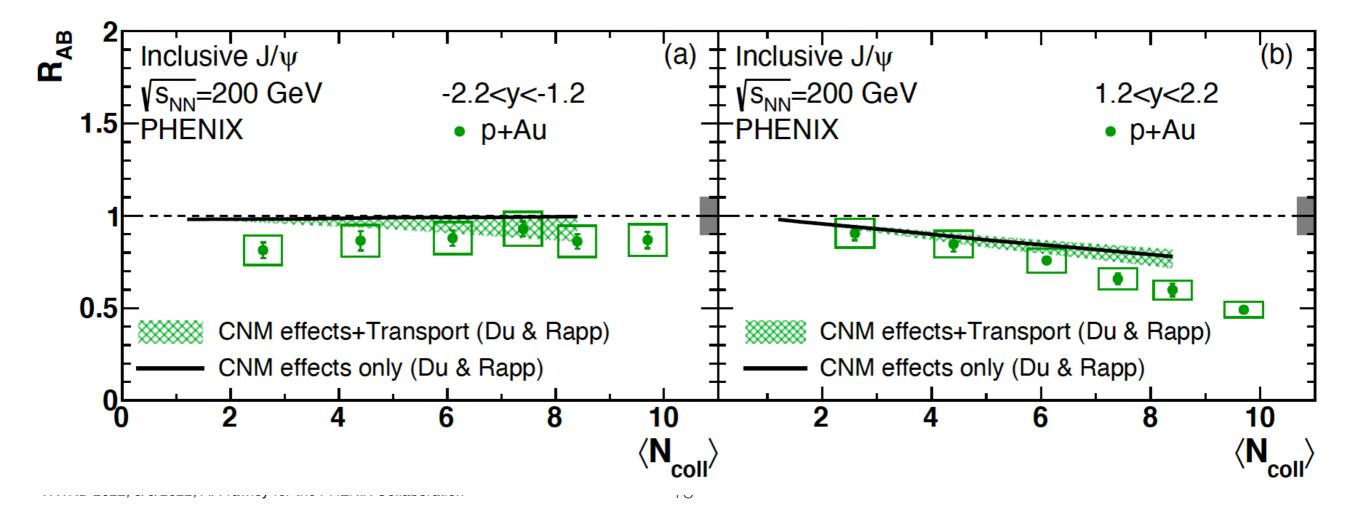


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$p+Au N_{coll}$ dependence for J/ψ

Compare transport calculation with N_{coll} dependence of p_T integrated data.

- At backward rapidity anti-shadowing + absorption + small final state effect.
- At forward rapidity suppression is dominated by EPS09 shadowing.
 - Centrality dependence is assumed to be linear.
 - EPS09 under-predicts suppression considerably!
 - But later parameterizations have stronger low x shadowing.

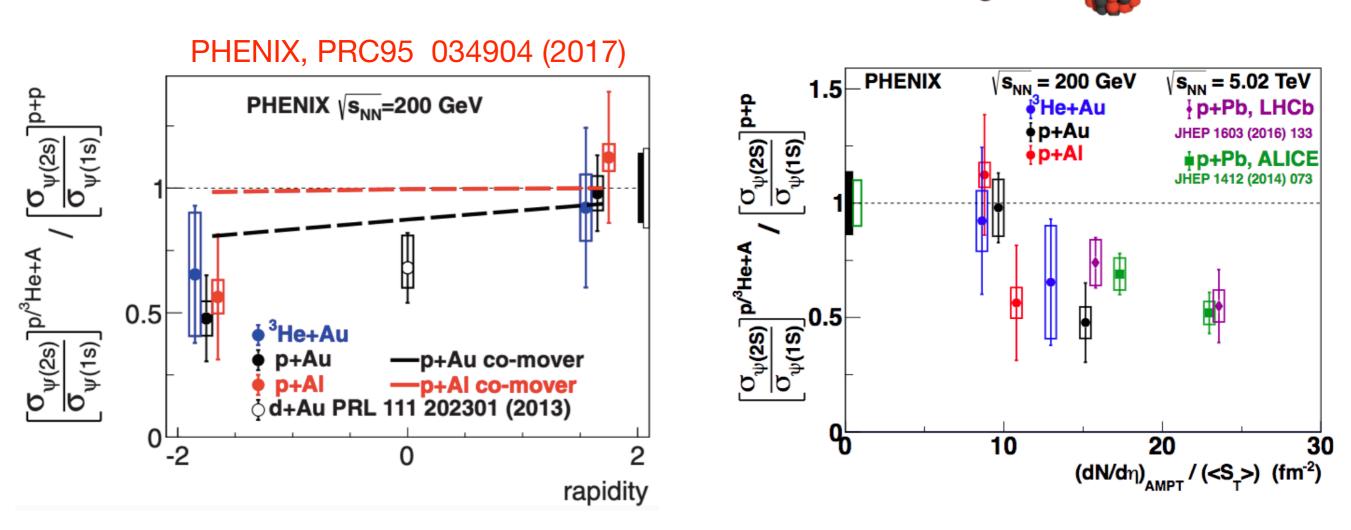


ψ(2S) results at forward/backward rapidity

Earlier $\psi(2S)$ results - centrality integrated

Strong suppression at mid & backward rapidity relative to J/ψ

Not explained by CNM effects (CNM very similar for both states)



The differential suppression is correlated with particle multiplicity

- Suggestive of a final state effect
- Would like more yield, and centrality dependence.

New - ψ(2S) production vs centrality

Very statistically challenged measurement!

• Feasible only for the p+Au data set.

Issues:

- Pair acceptance of FVTX/muon arm drops with distance from Z = 0.
- $\psi(2S)$ yield is low when both tracks must pass through FVTX.
- Mass resolution inadequate when neither track passes through FVTX.

Solution:

- Use all muon pairs with at least one track passing through FVTX.
- Use pairs with two FVTX associations for systematic estimates

Provides best compromise between mass resolution and yield.

Mass spectrum fitting

Crystal Ball line-shapes used for quarkonia.

- J/ ψ mass and width fixed from MB data.
- $\psi(2S)$ mass & width **ratio** to J/ψ from simulations.
- Tail parameters from simulations.

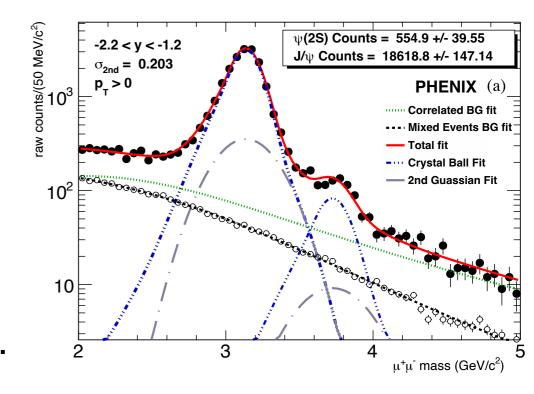
 $\psi(2S)$ yield is very sensitive to high mass tail of J/ ψ .

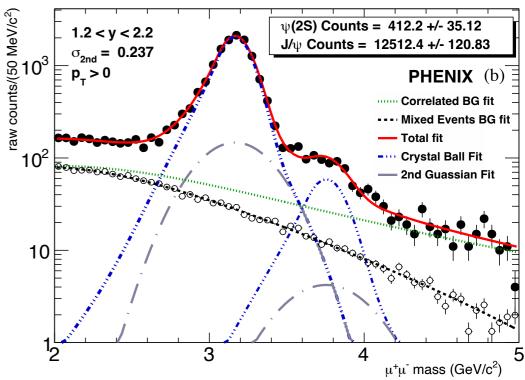
- Caused by mis-association of MuTr and FVTX hits.
- Included in fit using second gaussian.
 - Parameters determined from simulation.

Combinatorial background estimated from event mixing.

Correlated background.

- Open HF, Drell Yan, charged hadron muon decays.
- Poorly constrained by the data.
- Use a modified Hagedorn function, constrained by detailed simulations of all components.

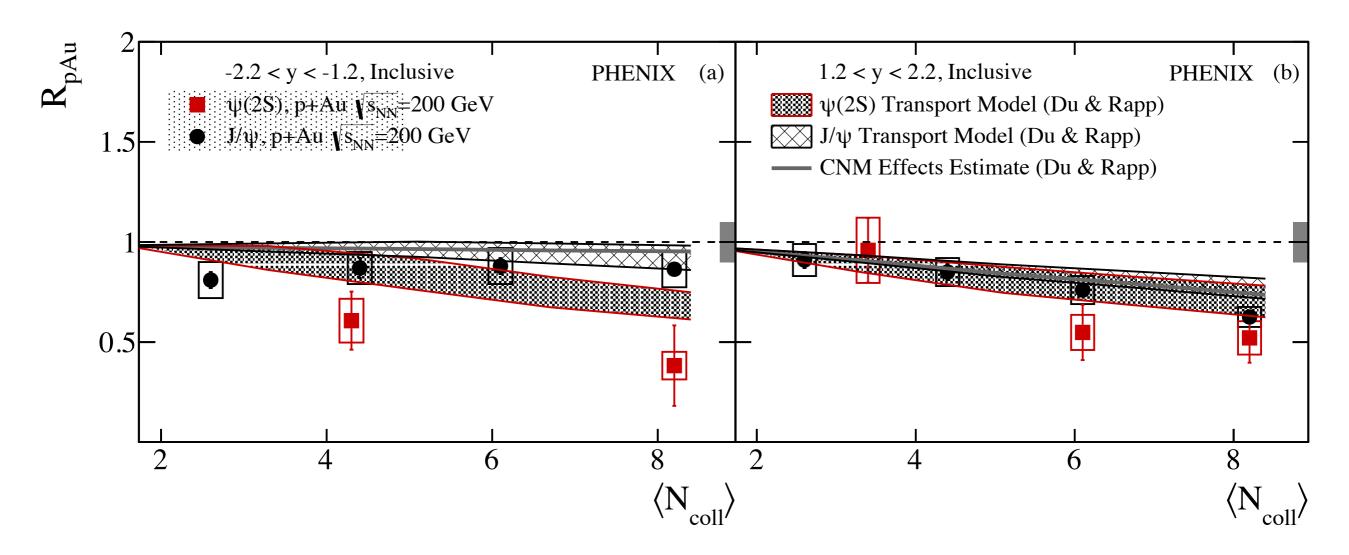




ψ(2S) R_{pAu} - centrality dependence

Nuclear modification in p+Au collisions for J/ ψ and ψ (2S) as a function of $<N_{coll}>$.

Du and Rapp transport model somewhat under-predicts the suppression, but gets the suppression **ratios** about right.



ψ(2S) R_{pAu} centrality dependence - compare with shadowing alone

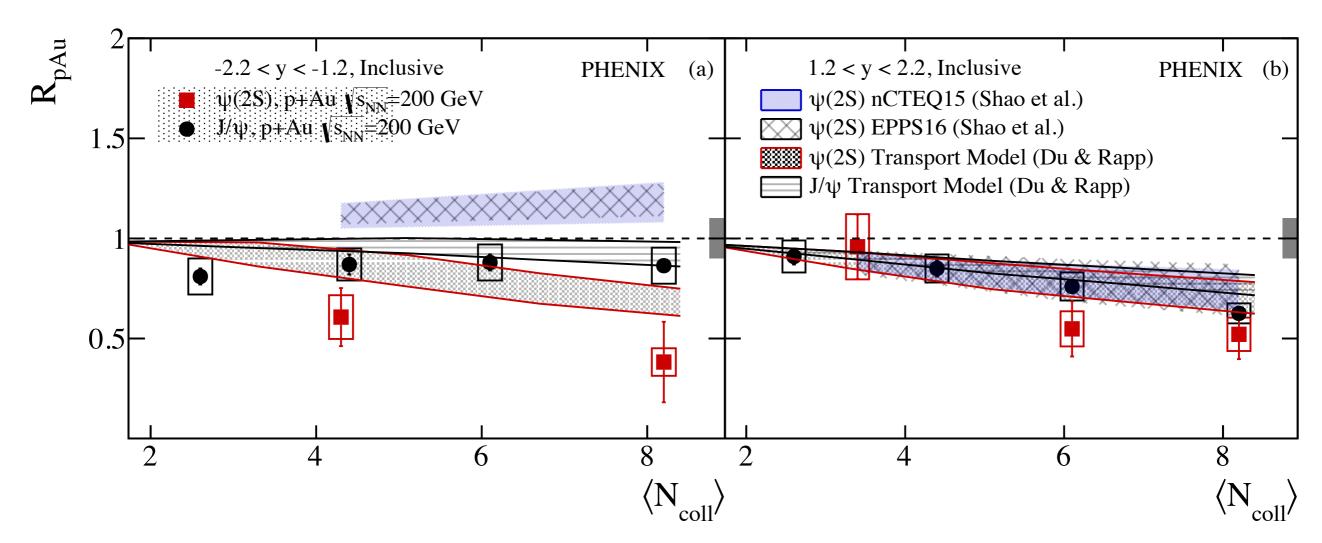
Add re-weighted shadowing comparison to plot.

Forward rapidity:

Modification consistent with shadowing alone.

Backward rapidity:

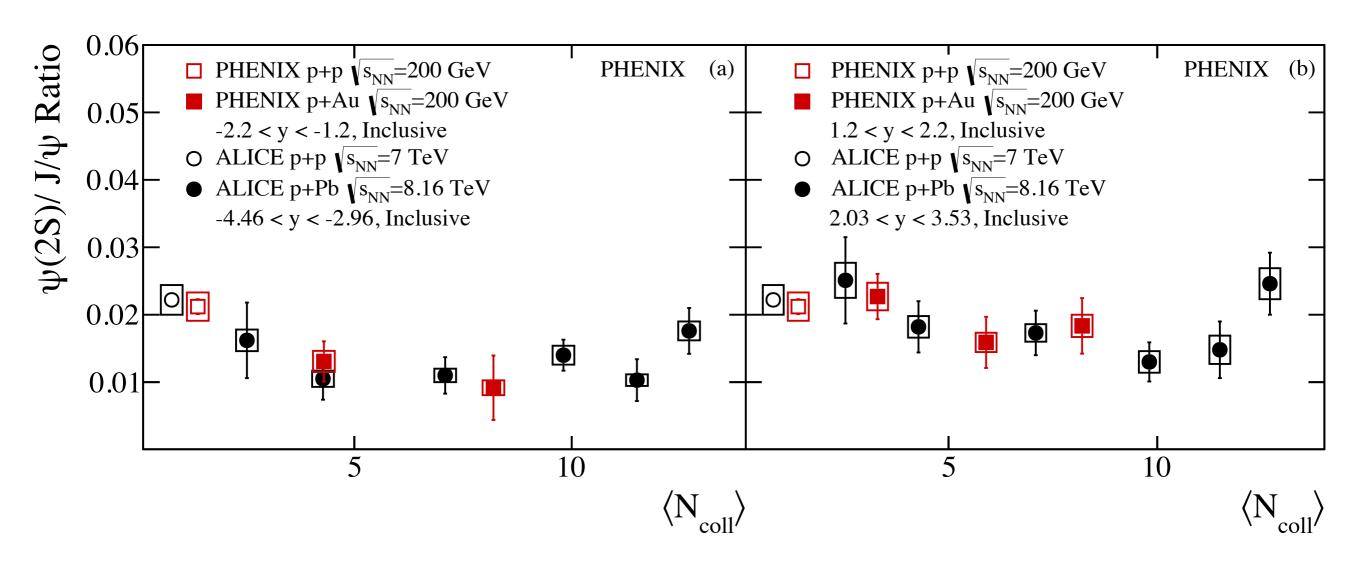
Require addition of strong absorption + differential $\psi(2S)$ suppression.



$\psi(2S)/J/\psi$ ratio vs N_{coll} - PHENIX/ALICE

PHENIX and ALICE $\psi(2S)$ to J/ ψ ratio plotted together.

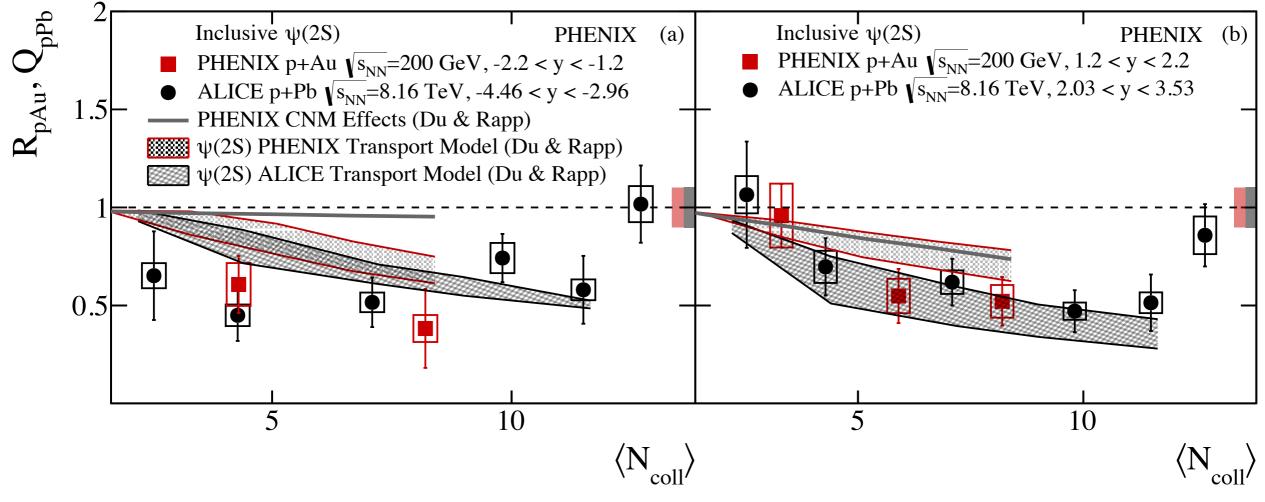
- Behavior is very similar at the two energies.
- The ratio is considerably smaller at backward rapidity.



ψ(2S) R_{pAu} vs N_{coll} - PHENIX/ALICE

Simultaneous comparison of PHENIX and ALICE $\psi(2S)$ modification data with Du & Rapp transport model.

- Similar suppression at backward rapidity
 - Combination of anti-shadowing, absorption, final state effects.
- The different model suppression at forward rapidity is due to differences in shadowing.

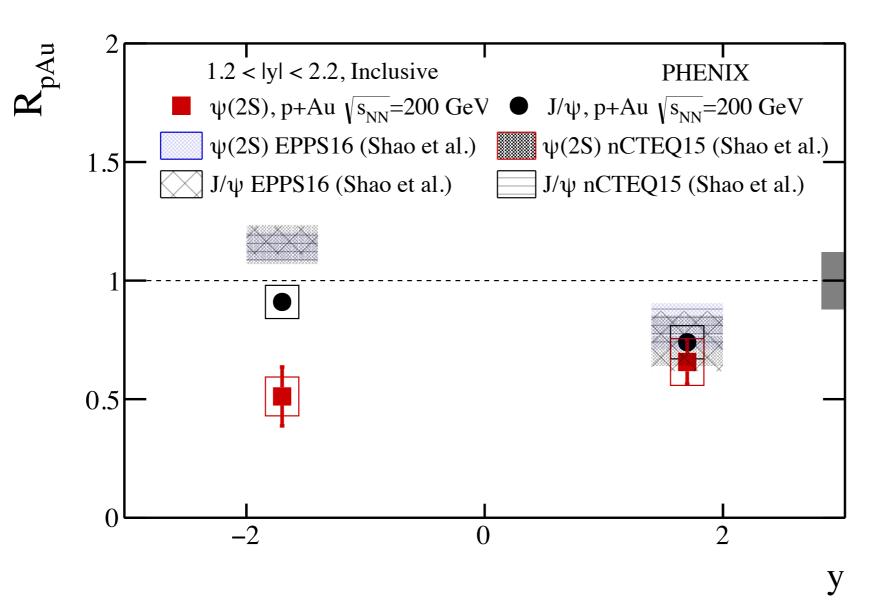


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ψ(2S) R_{pAu} vs rapidity - compare shadowing only

Centrality integrated modification vs rapidity.

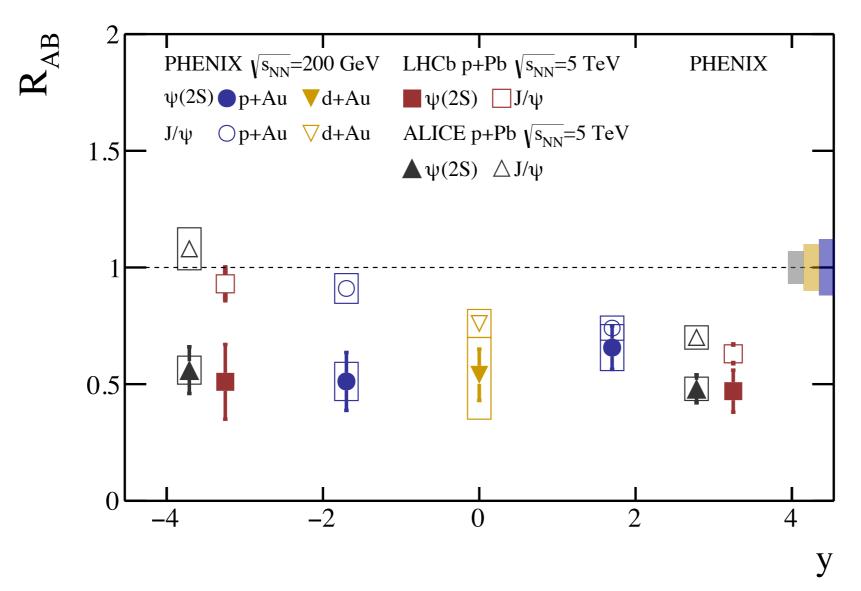
- Forward rapidity: good agreement with data for shadowing alone.
- Backward rapidity: Requires strong absorption + differential $\psi(2S)$ suppression to achieve the measured modification.



ψ(2S) R_{pAu} vs rapidity trend in world data

PHENIX, ALICE and LHCb modification for J/ ψ and ψ (2S) vs rapidity.

Clear trend of increasing differential suppression from forward to backward rapidity.



Conclusions

J/ψ modification in p+Au:

At forward rapidity:

• Described reasonably well by shadowing alone.

At backward rapidity:

- Described reasonably well by anti-shadowing + absorption.
 - There is room for a small (~10%) contribution from final state effects.
 - But there is not strong evidence for it.

ψ(2S) modification in p+Au:

At forward rapidity:

• Described reasonably well by shadowing alone.

At backward rapidity:

- Requires substantial additional suppression from final state effects.
 - Transport model accounts reasonably well for the differential $\psi(2S)$ suppression.

Backup

