

# Calculation of 3-pion Coulomb Scattering using Scalar QED

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Winter Workshop on Nuclear Dynamics 2022

# Outline

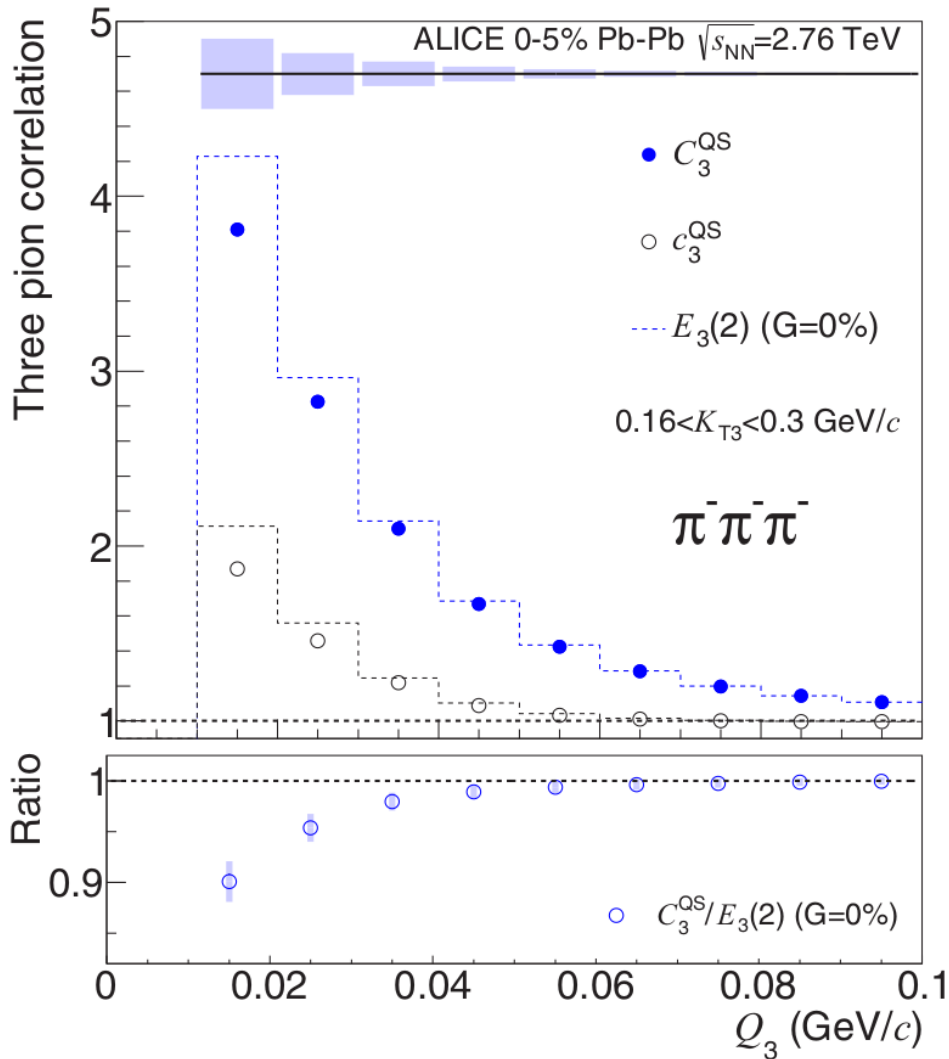
## **Central question**

How do the true 3-pion Coulomb interactions compare to the commonly used asymptotic approximations in high-energy hadronic collisions?

- 1) Why are 3-pion Coulomb interactions important?
- 2) Feynman diagram approach to calculating Coulomb interactions.
- 3) Scalar QED Feynman rules.
- 4) Benchmark studies using 2-pion QED calculations at NNLO.
- 5) 3-pion QED calculations at NNLO.

# Motivation:

## Why are 3-pion Coulomb interactions important?



PRC 93 054908 (2016)

ALICE measurements of 3-pion Bose-Einstein correlations revealed a suppression wrt expectations (dashed lines).

3-body Coulomb interactions were taken into account using an asymptotic ansatz that is valid at sufficiently “large” triplet relative momentum  $Q_3$ .

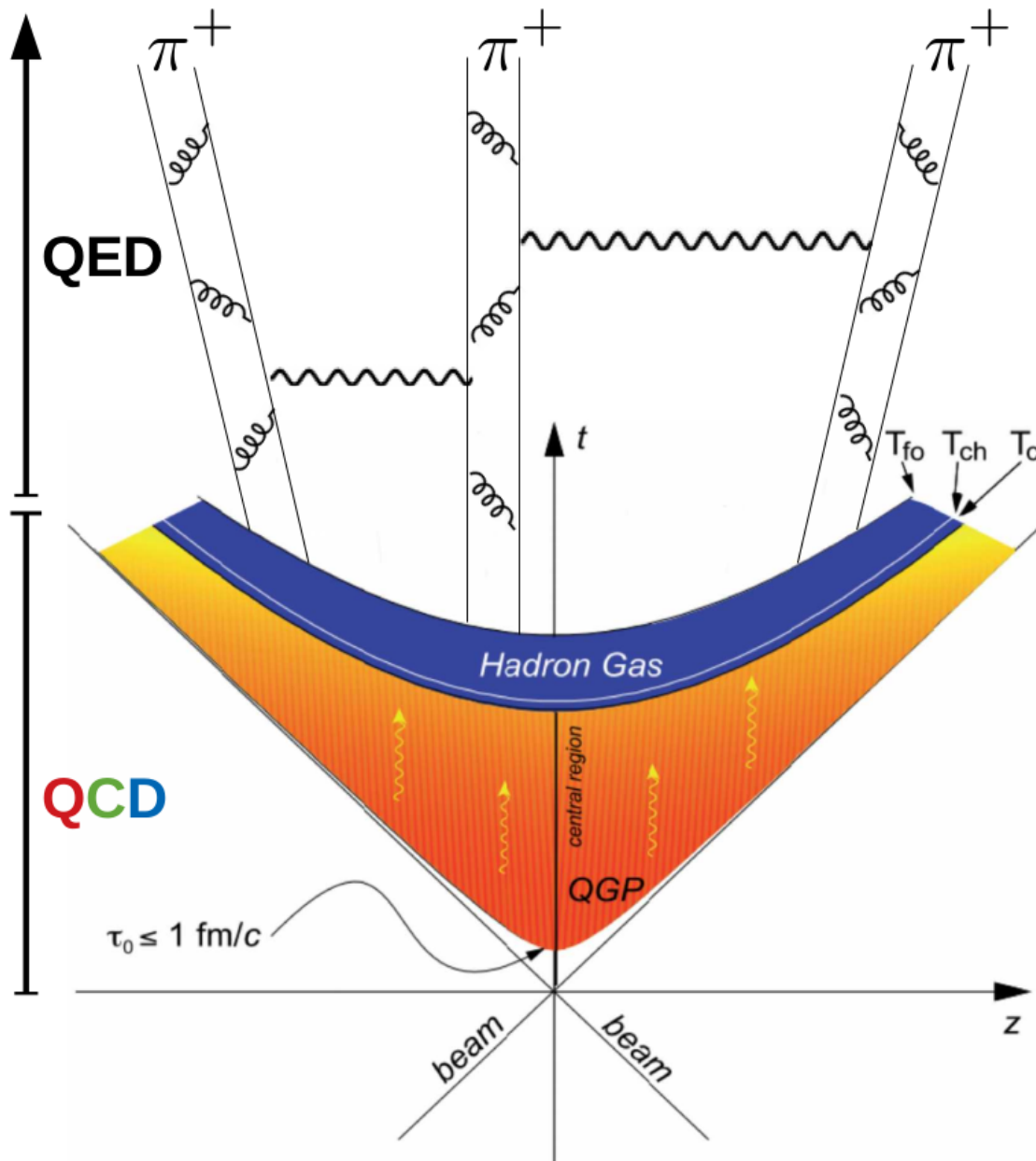
Ansatz Coulomb factor used:

$$K_3 = K_{12} K_{13} K_{23}$$

A calculation of genuine 3-body Coulomb interactions is needed to be sure about the origin of this suppression.

$$Q_3 = \sqrt{q_{12}^2 + q_{13}^2 + q_{23}^2}$$

# Feynman diagram approach to calculating Coulomb interactions



Before freeze out,  
**QCD** processes dominate.

After freeze out,  
**QED** dominates the interaction  
 Between charged pions.

The production amplitude of a pair/triplet  
 at kinetic freeze out is referred to as  $M_0$

For simplicity,  
 $M_0$  is treated as momentum independent  
 (point-source Gamow approximation).

The calculation presented here pertains  
 to the QED interactions after freeze out.

The calculation is needed to help  
 interpret the suppression of 3-pion  
 Bose-Einstein correlations.

**Quantum coherence at freeze out or  
 “extra” Coulomb repulsion?**

# Scalar QED Feynman rules

Lagrangian:

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^2 + (D_\mu\phi)^*(D^\nu\phi) - m^2\phi^*\phi$$

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Propagators:

$$\text{---} \underset{k}{\text{---}} \text{---} = \frac{i}{k^2 - m^2 + i\varepsilon}$$

$$\text{~~~~~} \underset{k}{\text{~~~~~}} \text{~~~~~} = \frac{-ig^{\mu\nu}}{k^2 + i\varepsilon}$$

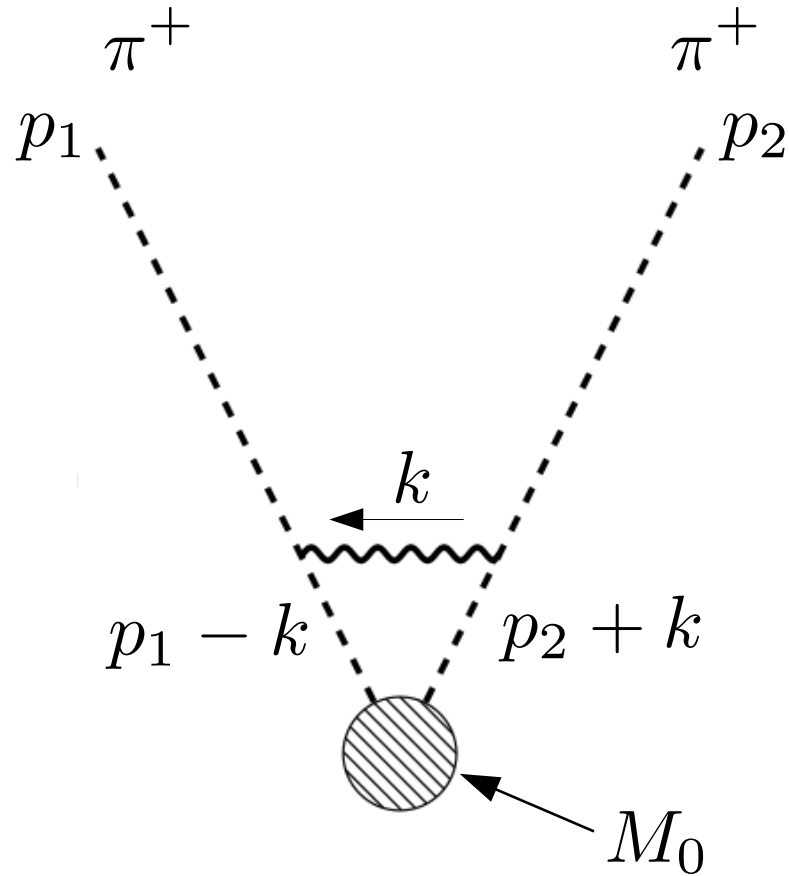
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Vertices:

$$\begin{array}{c} \text{---} \\ \text{---} \end{array} \underset{k'}{\text{---}} \text{~~~~~} \mu = -ie(k + k')^\mu$$

$$\begin{array}{c} \text{~~~~~} \\ \text{~~~~~} \end{array} \underset{\mu}{\text{~~~~~}} \underset{\nu}{\text{~~~~~}} = 2ie^2 g^{\mu\nu}$$

# Example of 1-loop amplitude



Dimensionally regulated integral  
 $d = 4 - 2\epsilon$

**QCD** amplitude factor  
(momentum independent)

$$I = M_0 \underbrace{\mu^{2\epsilon} e^{\gamma\epsilon}}_{\text{green arrow}} \int_{-\infty}^{\infty} \frac{d^d k}{(2\pi)^4} \frac{i}{(p_1 - k)^2 - m^2 + i\epsilon} [-ie(2p_1 - \cancel{k})^\mu] \frac{-ig^{\mu\nu}}{k^2 + i\epsilon} [-ie(2p_2 + \cancel{k})_\nu] \frac{i}{(p_2 + k)^2 - m^2 + i\epsilon}$$

*(Red arrows point to the crossed-out terms in the numerator)*

$\mu = m_{\text{pion}}$

$\gamma = \text{Euler constant} = 0.577$

Loop momenta can be dropped from the numerator.  
This obviates renormalization.

# Non-relativistic solution to 2-pion problem: Gamow factor

Some non-relativistic simplifications to the propagator are justified after a scale transformation:

$$\mathbf{k} \rightarrow \mathbf{p}\mathbf{k}, \quad k^0 \rightarrow \mathbf{p}^2/m k^0$$

Baier & Fadin  
Sov. Phys. JETP 30  
127 (1970)

$$\frac{1}{(p - k)^2 - m^2} = \frac{1}{k^2 - 2pk} \rightarrow \left[ \frac{1}{\mathbf{p}^2} \right] \frac{1}{-\mathbf{k}^2 - 2k^0 + 2\hat{\mathbf{p}}\mathbf{k}}$$

After this, Baier and Fadin showed how to resum the entire perturbative series analytically to obtain the well-known QM result of 2-body Coulomb from George Gamow

$$|\psi|^2 \equiv C_2 = \frac{\frac{4\pi}{qa}}{e^{\frac{4\pi}{qa}} - 1}$$

**Gamow** 2-body Coulomb factor  
Valid for point-source emission

$$q = \sqrt{-(p_1^\mu - p_2^\mu)^2}$$

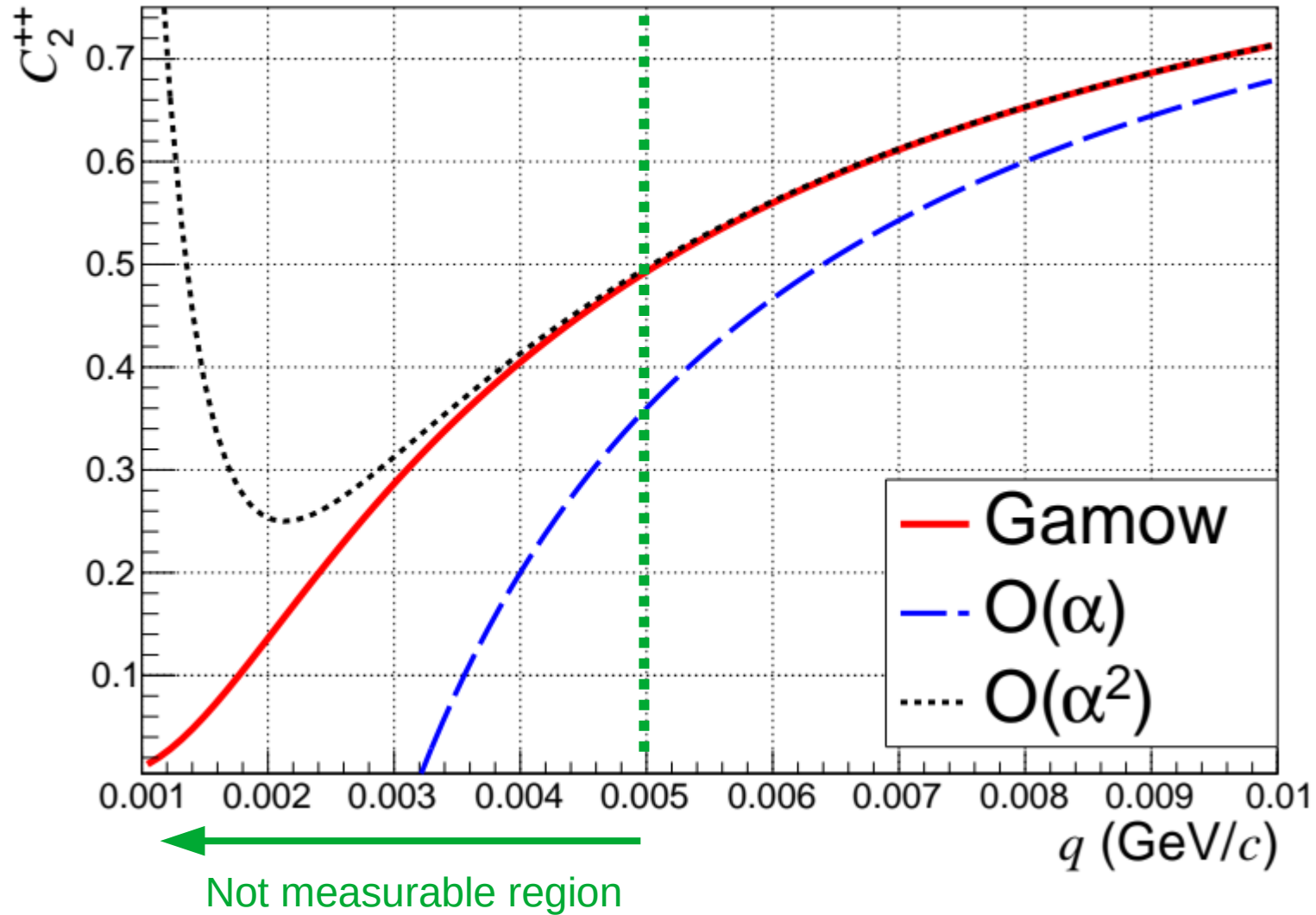
Lorentz invariant relative momentum

$a$

Bohr radius of the pair: 388 fm for pions

# 2-pion Gamow compared to $O(\alpha)$ and $O(\alpha^2)$ expansion

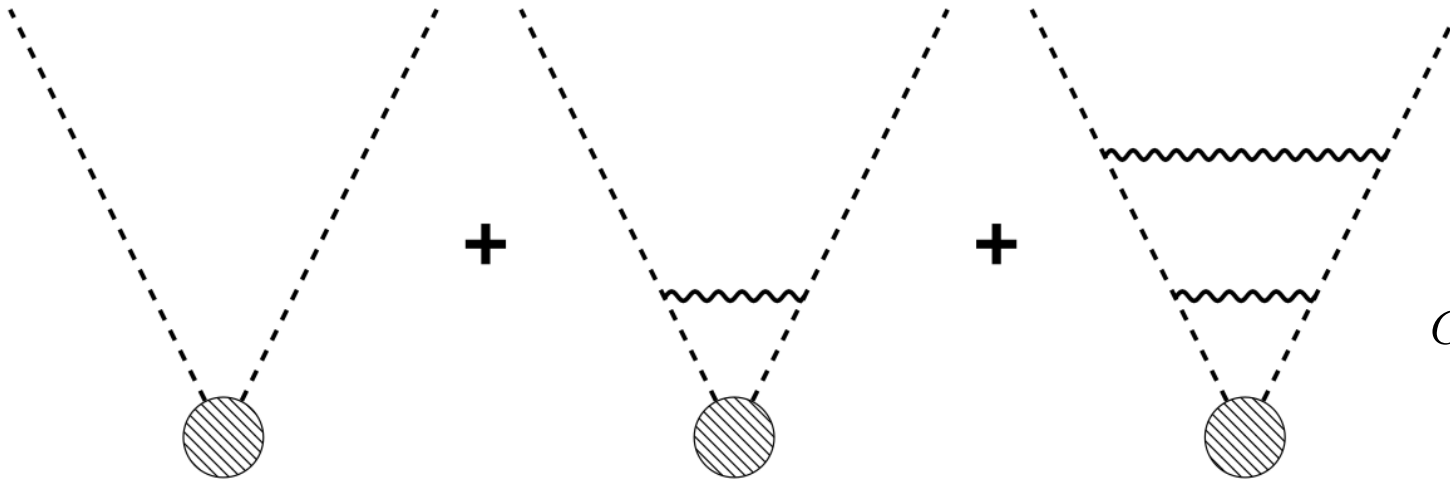
$\pi^+ \pi^+$



Just 2 terms in the series is very accurate in the measurable region!!



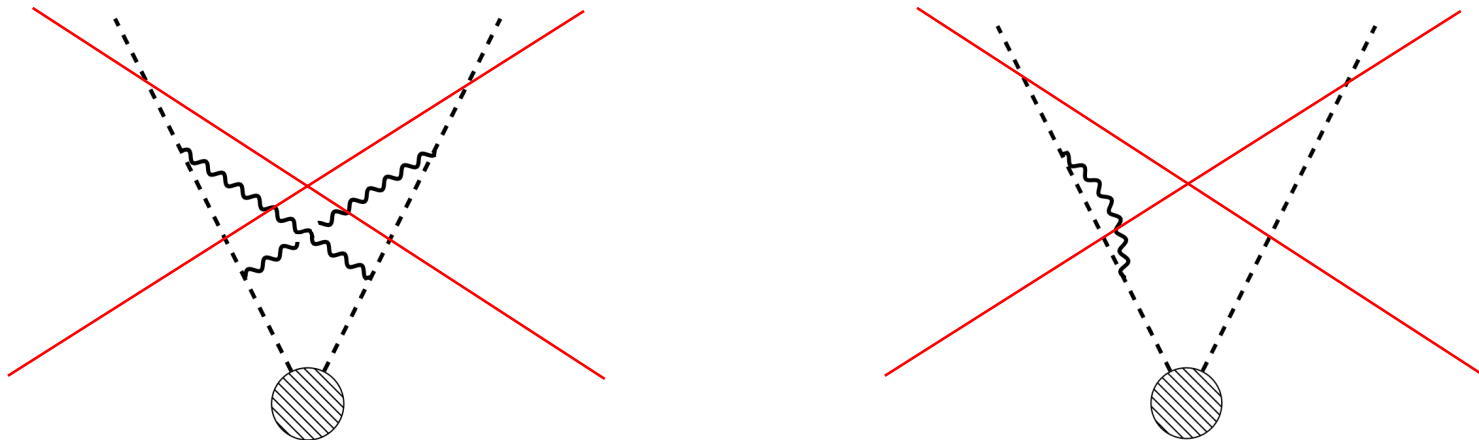
# 2-pion diagrams to be calculated using scalar QED



**2-pion correlation function**

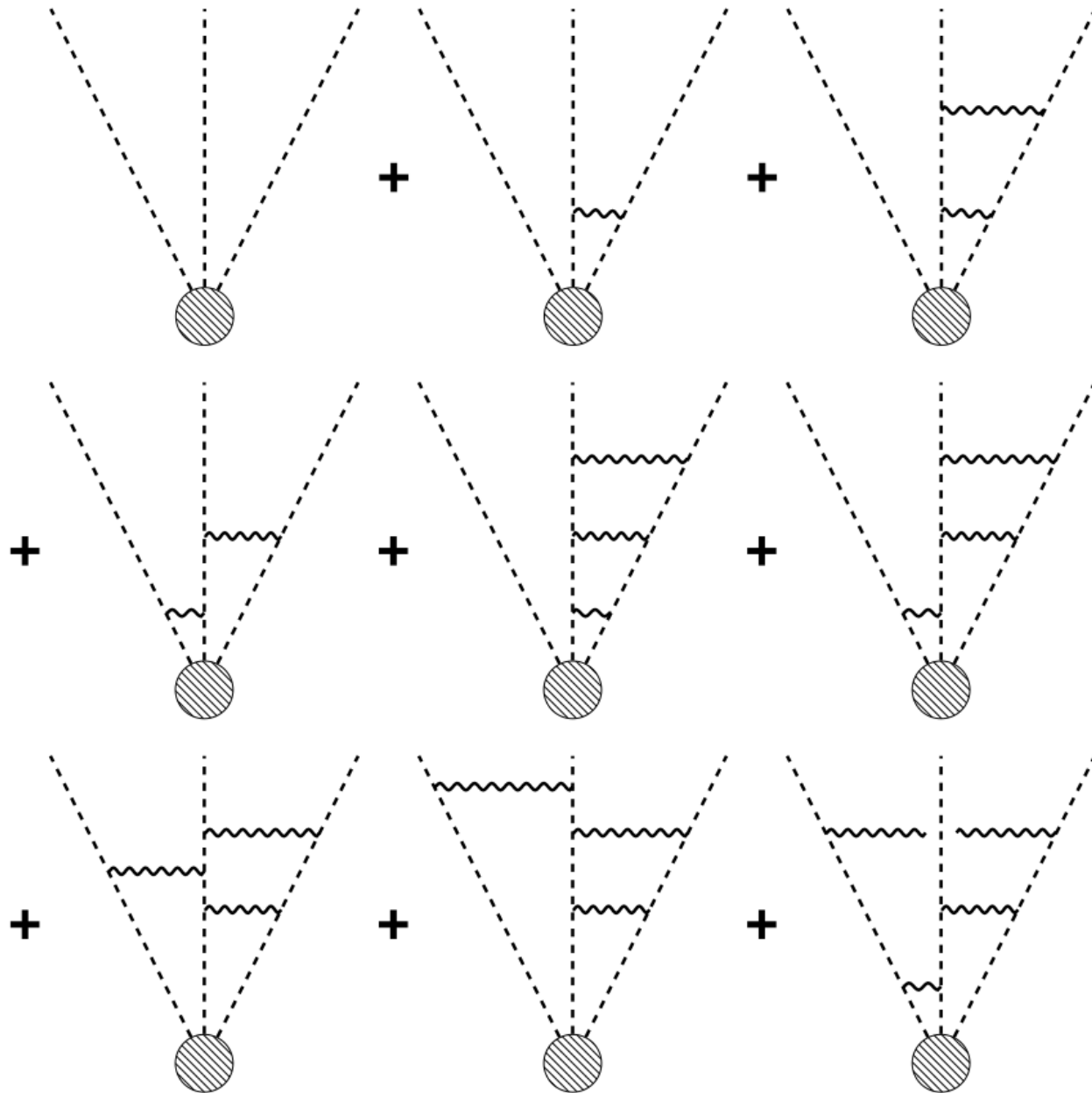
$$C_2 = \frac{|\text{sum of all terms}|^2}{|\text{no photon term}|^2}$$

Only straight “ladder” diagrams are important at low relative momentum



Baier & Fadin  
Sov. Phys. JETP 30  
127 (1970)

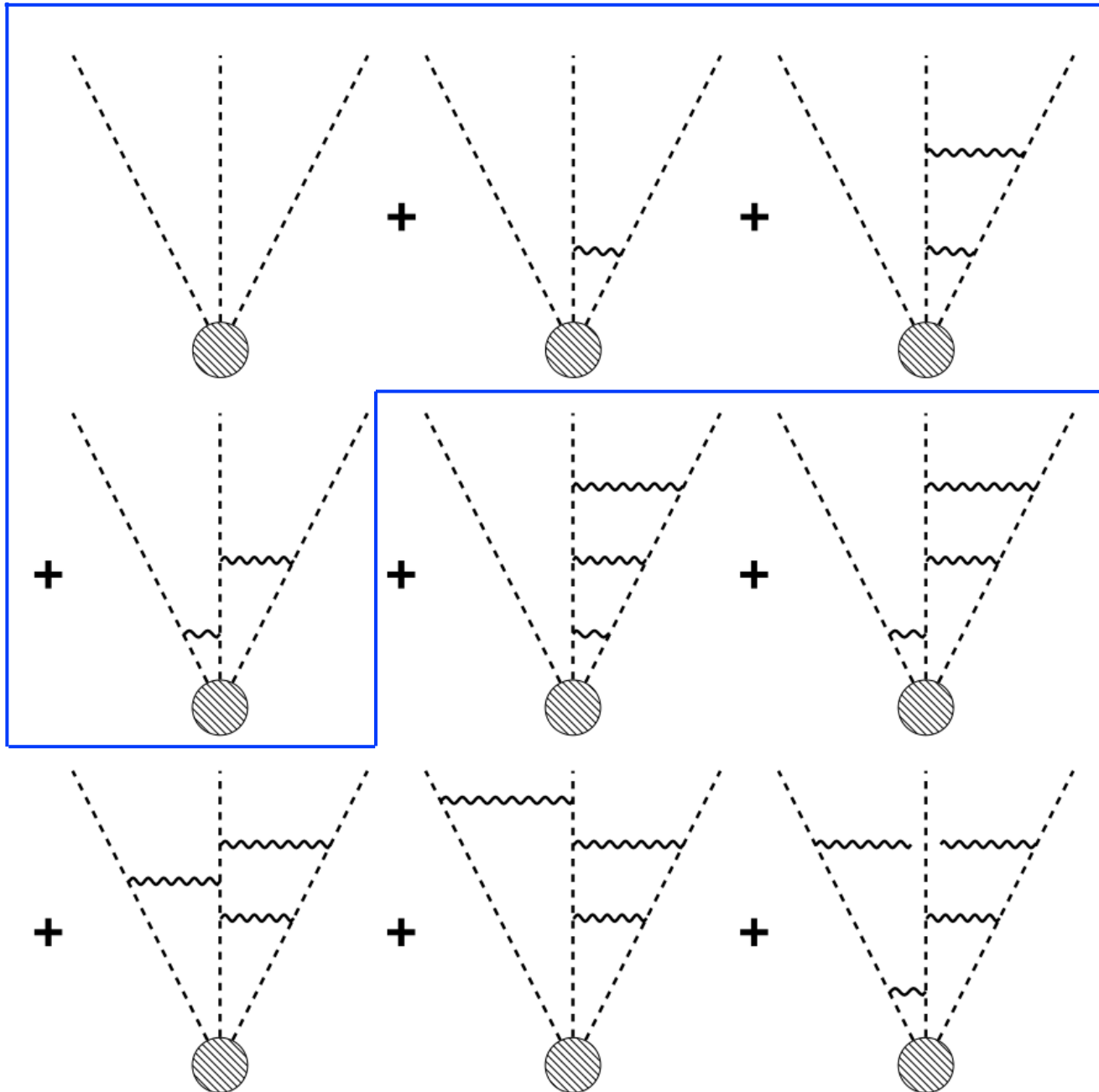
# 3-pion diagrams to be calculated using scalar QED



**3-pion correlation function**

$$C_3 = \frac{|\text{sum of all terms}|^2}{|\text{no photon term}|^2}$$

# 3-pion diagrams to be calculated using scalar QED



## 3-pion correlation function

$$C_3 = \frac{|\text{sum of all terms}|^2}{|\text{no photon term}|^2}$$

Only these terms have been calculated at the moment

The rest need special treatment on a GPU farm

# Sector decomposition and dimensional regularization

Regularize by shifting exponents  
by a small parameter  $\epsilon$

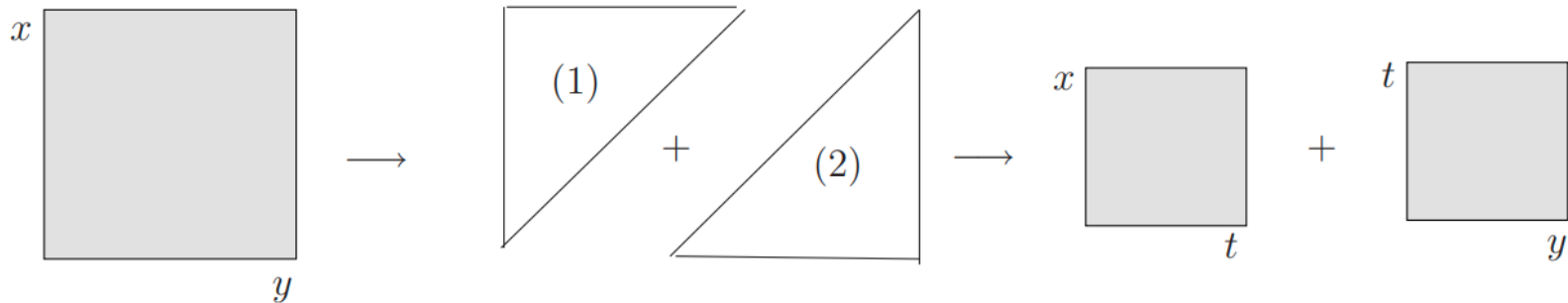
Decompose integral into 2 sectors  
to factor out the divergent part

$$\int_0^1 \frac{dx dy}{(x+y)^2} \rightarrow \int_0^1 \frac{dx dy}{(x+y)^{2-\epsilon}} \rightarrow \int_0^1 \frac{dx dy}{(x+y)^{2-\epsilon}} \left[ \overset{(1)}{\Theta(x-y)} + \overset{(2)}{\Theta(y-x)} \right]$$

Change variables

$$y \rightarrow xt \qquad x \rightarrow yt$$

$$0 \leq t \leq 1 \qquad 0 \leq t \leq 1$$



G. Heinrich  
arXiv:0803.4177

# Sector decomposition and dimensional regularization

$$\int_0^1 \frac{dx dy}{(x+y)^{2-\epsilon}} [\Theta(x-y) + \Theta(y-x)] = \int_0^1 \frac{x dx dt}{(x(1+t))^{2-\epsilon}} + \int_0^1 \frac{y dy dt}{(y(1+t))^{2-\epsilon}}$$

Change variables  
 $y \rightarrow xt \quad x \rightarrow yt$   
 $0 \leq t \leq 1 \quad 0 \leq t \leq 1$

$$= 2 \left[ \int_0^1 \frac{dx}{x^{1-\epsilon}} \right] \left[ \int_0^1 \frac{dt}{(1+t)^{2-\epsilon}} \right] = 2 \left[ \frac{1}{\epsilon} \right] \left[ \frac{1}{2} (1 + \epsilon(1 - \ln 2)) \right]$$

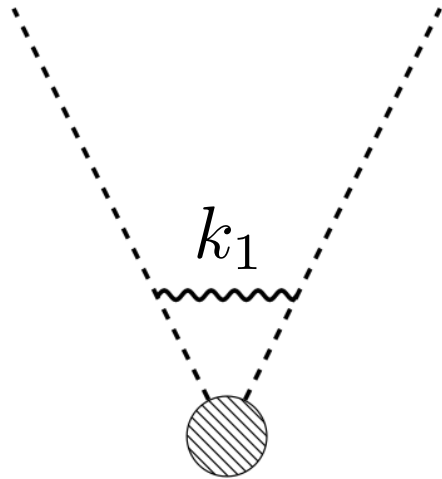
$$= \frac{1}{\epsilon} + (1 - \ln 2)$$



Divergent part

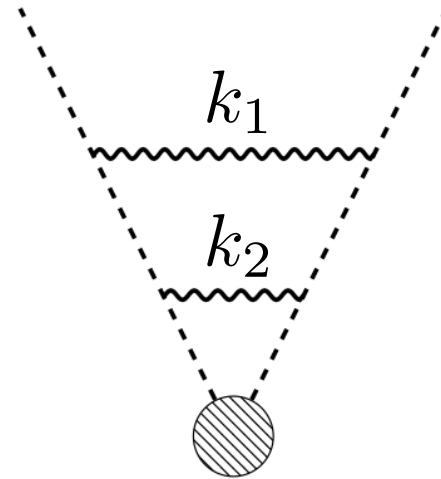
Finite part

# General form of Feynman integrals with dimensional regularization



$$I = \frac{A}{\epsilon} + B + C\epsilon + \mathcal{O}(\epsilon^2)$$

Divergence  
when  $k_1 \rightarrow 0$



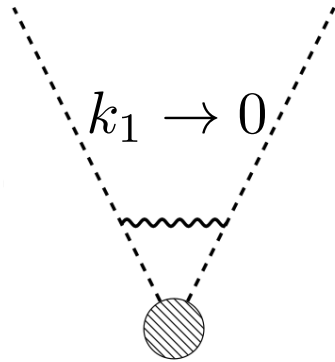
$$I = \frac{A}{\epsilon^2} + \frac{B}{\epsilon} + C + D\epsilon + \mathcal{O}(\epsilon^2)$$

Divergence  
when  $k_1$  &  $k_2 \rightarrow 0$

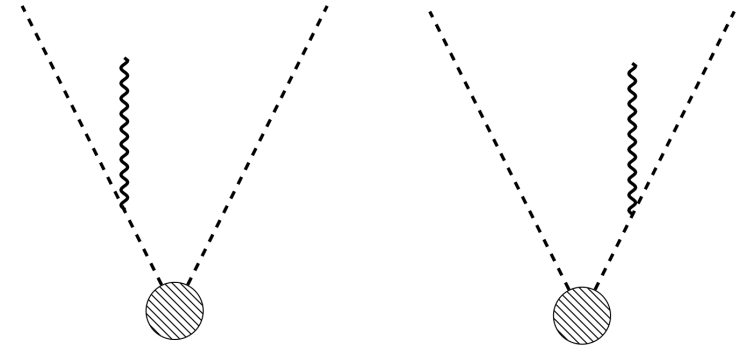
Divergence  
when  $k_1$  or  $k_2 \rightarrow 0$

# Infrared and Ultraviolet divergencies

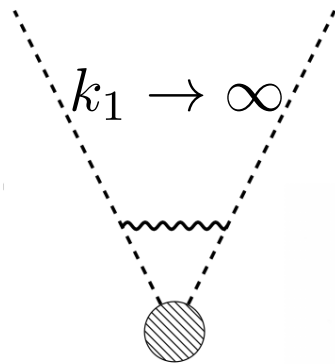
## Infrared divergencies



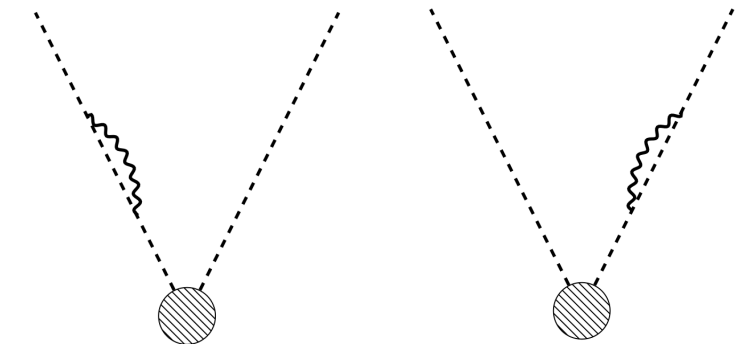
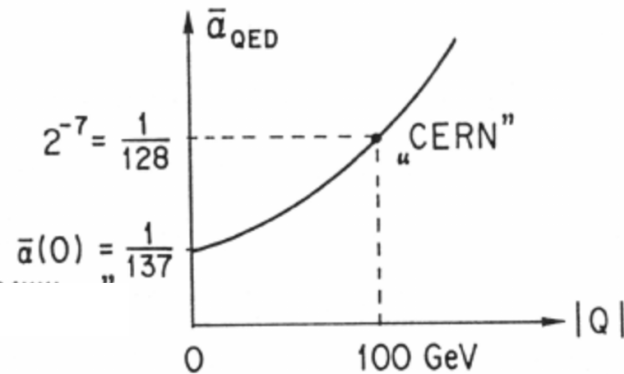
**Canceled** by contributions with a soft real photon emission



## Ultraviolet divergencies



**Absorbed** into a modified coupling and pion mass (Renormalization)



Also need these diagrams to renormalize

Baier & Fadin  
Sov. Phys. JETP 30  
127 (1970)

**Important simplification:** For low momentum pions, real photon emission is highly suppressed and also the "standard" values of  $\alpha_{\text{QED}}$  and  $m_{\text{pion}}$  can be used. That means that the divergent terms can simply be discarded.

# GPGPU-based computation

The Feynman integrals are evaluated numerically using the software package:

## pySecDec

Gudrun Heinrich et al.  
arXiv: 1703.09692

- General-Purpose GPUs (GPGPU) are used to speed up the calculation.
- 1- and 2-loop diagrams can be calculated on a single GPGPU card
- 3-loop diagrams need a farm of GPGPUs

Compute times with a **single** high-end GPGPU

	1-loop	2-loop	3-loop
Compute time	~1 min	~1 day	> 1 month

My GPGPU

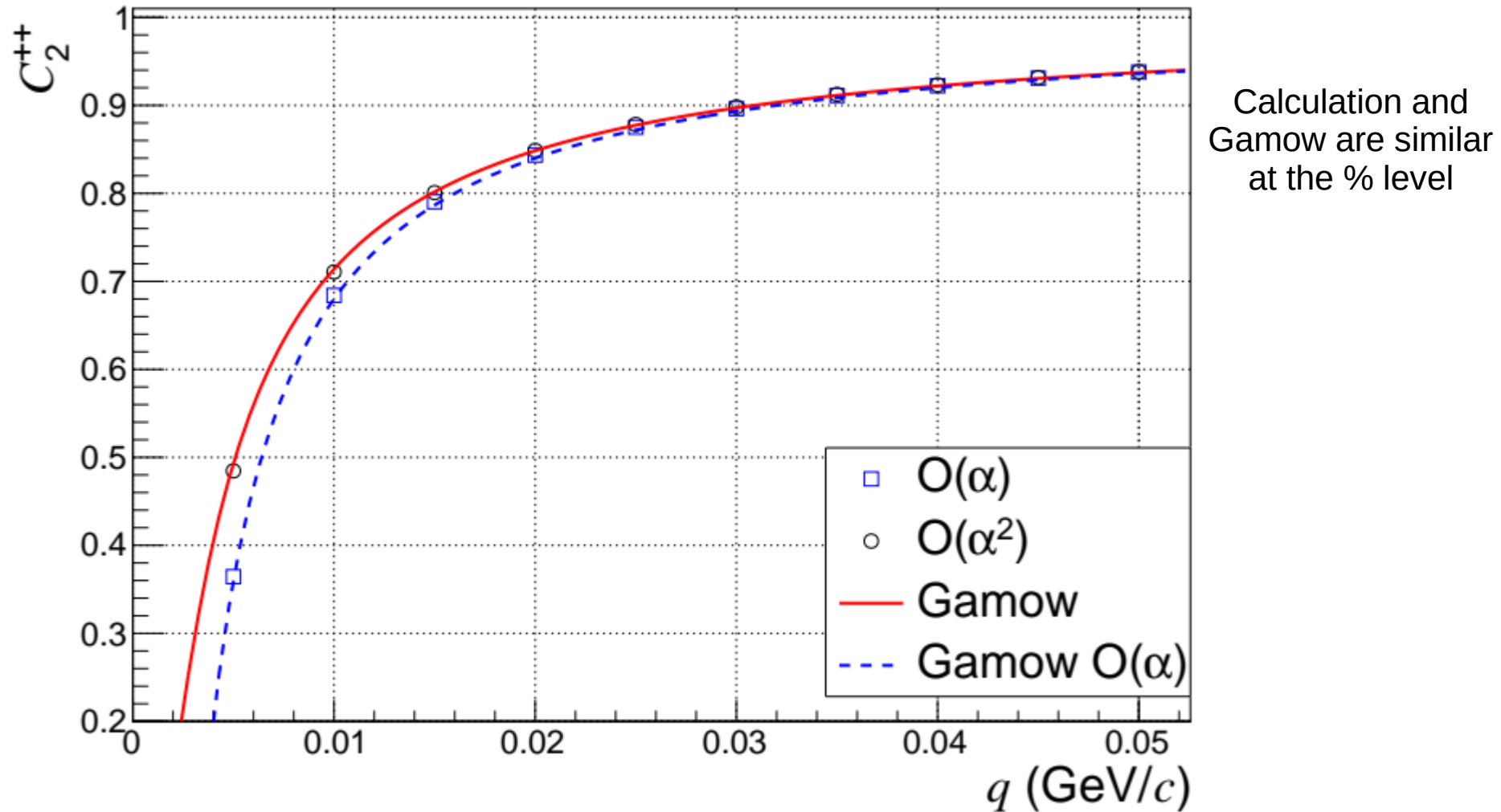


Claudia Ratti's GPGPU farm at UH





# Benchmark results: 2-pion calculation compared to non-relativistic expectation



The calculation should **not** be identical to Gamow but they should be close:

Quantum Field Theory  $\rightarrow$  Quantum Mechanics as  $q \rightarrow 0$

# 3-pion kinematics

Amplitudes are projected against the usual triplet Lorentz invariant relative momentum

$$Q_3 = \sqrt{q_{12}^2 + q_{13}^2 + q_{23}^2}$$

For a single value of  $Q_3$ , there is a continuum of pair configurations.

For a point-source, the momentum spectrum is flat, which yields the phase-space weight:

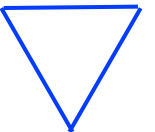
$$W \propto q_{12} q_{13} q_{23}$$

Two schemes of 3-pion kinematics are chosen:

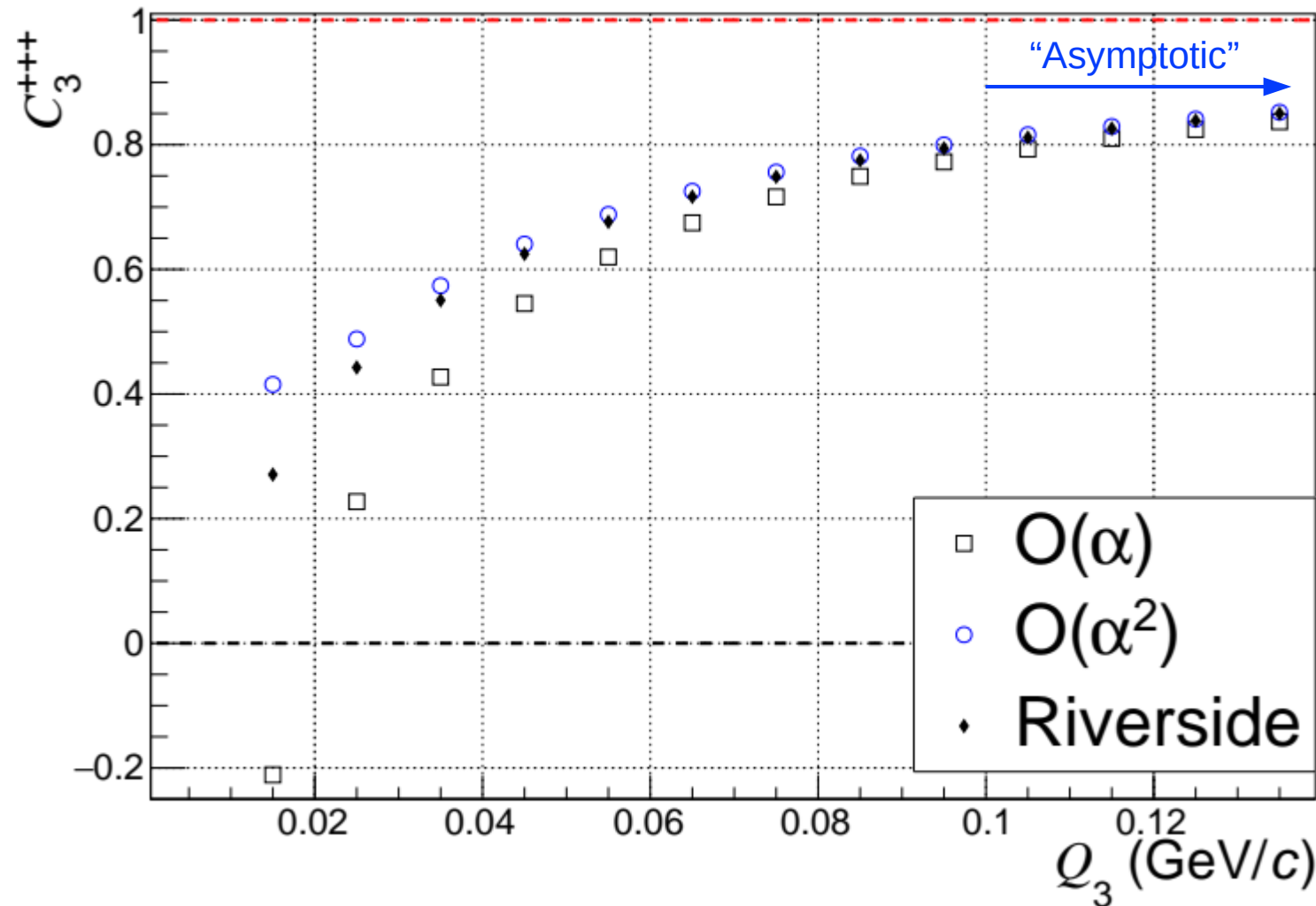
1) Weighted pair configurations (4 chosen): ranges from the most asymmetric to symmetric pair configurations. For each, 6 pair permuted diagrams need to be computed: **4\*6 = 24 diagrams per  $Q_3$  value.**

2) Symmetric configuration:  $q_{12} = q_{13} = q_{23}$ . **1 diagram per  $Q_3$  value.**

This scheme might be the only way forward for 3-loop diagrams.



# 3-pion Result: Weighted pair configurations

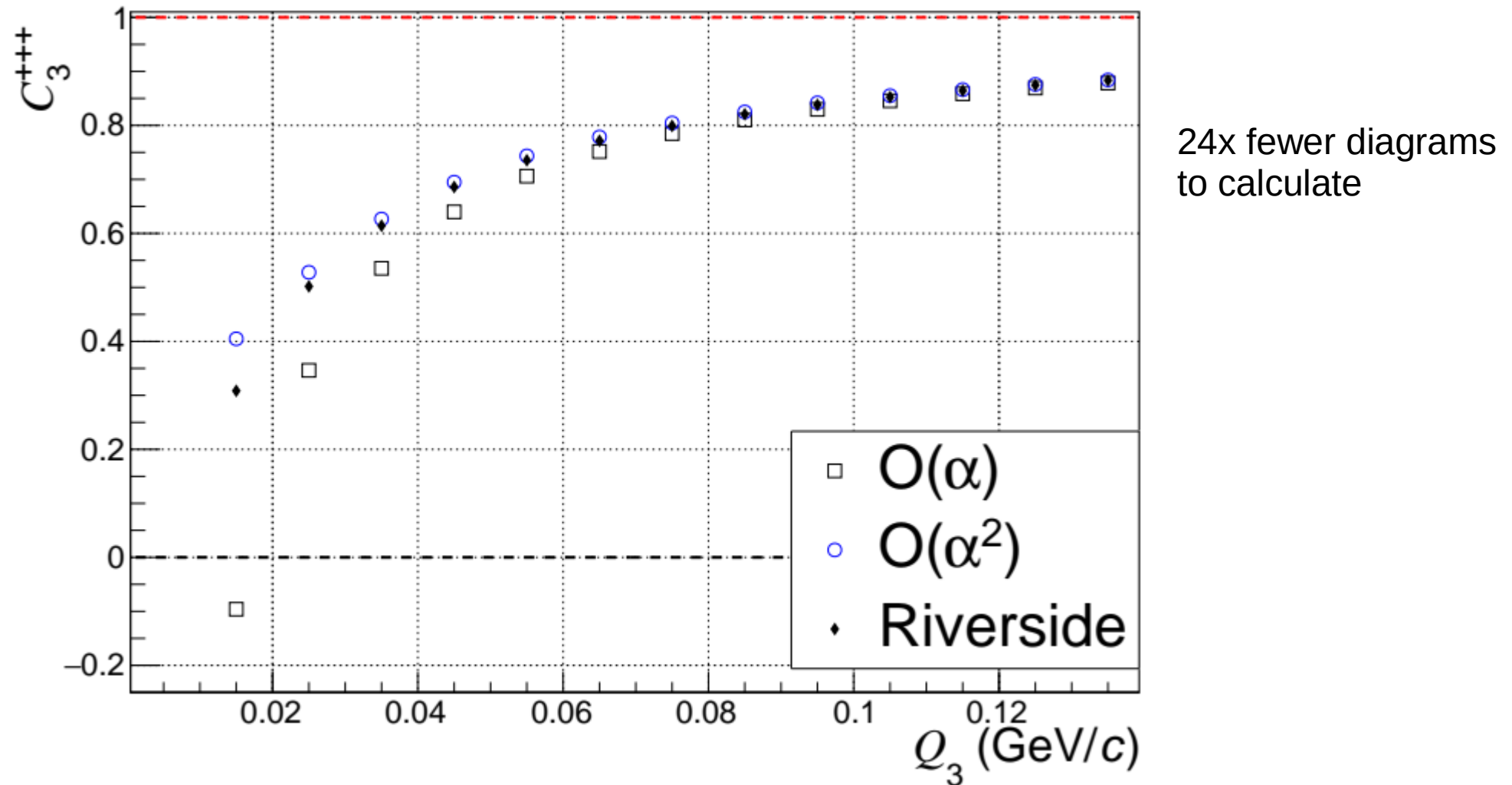


## Central Question:

Is the genuine 3-body  
Coulomb correlation  
greater or less than  
Riverside??

- $Q_3 > 0.1$  GeV,  $O(\alpha^2)$  is very similar to Riverside (asymptotic solution).
- Each order appears to contribute with sign =  $(-1)^N$ :  $O(\alpha)$  negative,  $O(\alpha^2)$  positive.
  - Need one more order to get an answer to the central question.

# 3-pion Result: Symmetric pair configuration



- ToDo: calculate 3-loop contributions,  $O(\alpha^3)$ , at lowest and highest  $Q_3$  points only.
- That should be decisive in answering the central question.

# Summary

- Coulomb interactions between 2- and 3-pions are perturbatively calculable with scalar QED in the measurable kinematics of virtually all high-energy experiments.
- Sector decomposition is a powerful tool to numerically calculate the Feynman integrals.
- At low relative momentum (non-relativistic case), IR and UV divergencies can simply be discarded from the dimensionally regularized integrals.
- For 2-pions, 2 terms in the series is sufficient.
- For 3-pions, one may need 3 terms for  $Q_3 < \sim 100$  MeV/c (series decreases less rapidly).

## ToDo

- Calculate  $O(\alpha^3)$  diagrams. Should be decisive in answering the central question:  
Is the asymptotic 3-body Coulomb solution justified in high-energy hadronic collisions??