

Conformally and Disformally Coupled Vector field Models of Dark Energy

L. Gabriel Gómez Díaz
Fondecyt Postdoc Fellow

Departamento de Física, Universidad de Santiago de Chile



UNIVERSIDAD
DE SANTIAGO
DE CHILE

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- Different representations of the same theory:

Jordan frame: fields and curvature couple directly, but matter only involves the metric.

Einstein frame: the gravitational Lagrangian has the Einstein-Hilbert form but the matter sector is affected by fields, which mediates an additional force.

-Coupled Multi-Proca (Gómez et. al. PDU (2020)):

$$\mathcal{S} = \int d^4x \sqrt{-g} \left(\frac{M_p^2}{2} R - \frac{1}{4} F_{\mu\nu}^a F_a^{\mu\nu} - V(\tilde{X}) + f(X) \tilde{\mathcal{L}}_m \right),$$

$$\bar{g}_{\mu\nu} = f(X) g_{\mu\nu}.$$

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-Anisotropic Scalar Field Dark Energy with a Disformally Coupled Yang-Mills Field (Gómez et. al. arXiv:2103.11826 [gr-qc]):

$$S = \int d^4x \sqrt{-g} \left[\frac{M_p^2}{2} R + p(\phi, X) + \mathcal{L}_M(g_{\mu\nu}, \Psi_i) \right] - \int d^4x \sqrt{-\hat{g}} \left[\frac{1}{4} \hat{g}^{\alpha\beta} \hat{g}^{\gamma\delta} F_{\alpha\gamma}^a F_{\beta\delta a} \right].$$

$$\hat{g}_{\mu\nu} = C(\phi) g_{\mu\nu} + D(\phi) \phi_{,\mu} \phi_{,\nu}.$$

-Vector-tensor sector coupling to dark matter:

$$\mathcal{S} = \int d^4x \left[\sqrt{-g} \left(\frac{M_p^2}{2} R + \mathcal{L}_A(X, Y) \right) + \sqrt{-\bar{g}} \bar{\mathcal{L}}_c[\bar{g}_{\mu\nu}, \psi_c] \right],$$

$$X = -\frac{1}{2} g^{\mu\nu} A_\mu A_\nu, Y = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu},$$

$$F_{\mu\nu} \equiv \nabla_\mu A_\nu - \nabla_\nu A_\mu.$$

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$$F_{\mu\nu} \equiv \nabla_\mu A_\nu - \nabla_\nu A_\mu.$$

-Vector disformal transformation:

$$\bar{g}_{\mu\nu} = C(X) g_{\mu\nu} + B(X) A_\mu A_\nu,$$

$$\bar{g}^{\mu\nu} = \frac{1}{C} \left(g^{\mu\nu} - \frac{B}{C-2BX} A^\mu A^\nu \right).$$

-Energy momentum tensor:

$$T_{(A)}^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta(\sqrt{-g}\mathcal{L}_A)}{\delta g_{\mu\nu}}, T_{(c)}^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta(\sqrt{-g}\bar{\mathcal{L}}_c)}{\delta g_{\mu\nu}}, T_{(i)}^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta(\sqrt{-g}\mathcal{L}_i)}{\delta g_{\mu\nu}}.$$

$$\bar{T}_{(c)}^{\alpha\beta} = \frac{2}{\sqrt{-\bar{g}}} \frac{\delta(\sqrt{-\bar{g}}\bar{\mathcal{L}}_c)}{\delta \bar{g}_{\alpha\beta}}.$$

-Cold dark matter sector:

$$T_{(c)}^{\mu\nu} = \sqrt{\frac{\bar{g}}{g}} \frac{\partial \bar{g}_{\alpha\beta}}{\partial g_{\mu\nu}} \bar{T}_{(c)}^{\alpha\beta},$$

$$T_{(c)}^{\mu\nu} = \sqrt{\frac{\bar{g}}{g}} \left[C \bar{T}^{\mu\nu} + \frac{1}{2} A^\mu A^\nu (C_{,X} g_{\alpha\beta} + B_{,X} A_\alpha A_\beta) \bar{T}_{(c)}^{\alpha\beta} \right].$$

-Vector-tensor sector:

$$\mathcal{L}_{A,Y} \nabla_\beta F^{\alpha\beta} + \mathcal{L}_{A,X} A^\alpha + \mathcal{M}_\beta F^{\alpha\beta} = Q^\alpha,$$

$$\mathcal{M}_\beta = \mathcal{L}_{A,XX} A^\nu \nabla_\beta A_\nu + \mathcal{L}_{A,YY} F^{\rho\nu} \nabla_\beta \nabla_\nu A_\rho.$$

-Source term:

$$Q^\alpha = -\frac{1}{\sqrt{-g}} \frac{\delta(\sqrt{-g}\bar{\mathcal{L}}_c)}{\delta A_\alpha} = -\frac{1}{\sqrt{-g}} \left(\frac{\partial(\sqrt{-g}\bar{\mathcal{L}}_c)}{\partial A_\alpha} - \nabla_\mu \frac{\partial(\sqrt{-g}\bar{\mathcal{L}}_c)}{\partial(\nabla_\mu A_\alpha)} \right).$$

-Continuity equations:

$$\begin{aligned} \nabla_\mu T_{(A)\nu}^\mu &= -\nabla_\mu T_{(c)\nu}^\mu = Q^\mu \nabla_\nu A_\mu - \nabla_\mu (Q^\mu A_\nu) \\ &= Q^\mu F_{\nu\mu} - A_\nu \nabla_\mu Q^\mu. \end{aligned}$$

$$\nabla^\mu T_{\mu\nu}^{(A)} + \nabla^\mu T_{\mu\nu}^{(c)} = 0.$$

-Proca theory with a potential: $\mathcal{L}_A = Y + m^2 X - V(X)$.

$$\nabla_\mu F^{\mu\nu} + (V_{,X} - m^2) A^\nu = -\frac{B}{C} T_{(c)}^{\nu\mu} A_\mu + \frac{D}{2C} (C - 2BX) \left(C_{,X} T_{(c)} + B_{,X} T_{(c)}^{\alpha\beta} A_\alpha A_\beta \right) A^\nu$$

$$D \equiv \frac{1}{C - C_{,X} X + 2B_{,X} X^2}.$$

Background Cosmology:

-FRW Metric:

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j,$$

-Temporal ansatz:

$$A_\mu \equiv (A(t), 0, 0, 0).$$

-Field equations:

$$3M_p^2 H^2 = \left(\frac{m^2}{2} - V_X\right) A^2 + V + \rho_c + \rho_r,$$

$$M_p^2(3H^2 + 2\dot{H}) = V - \frac{1}{2}m^2 A^2 - \frac{\rho_r}{3},$$

$$(m^2 - V_X)A = \frac{A\rho_c(C_X - B_X A^2 - 2B)}{B_X A^4 - C_X A^2 + 2C}.$$

-Effective equation of state:

$$w_{\text{eff}} \equiv \frac{p_T}{\rho_T} = - \left(1 + \frac{2\dot{H}}{3H^2}\right).$$

-Continuity equations:

$$\begin{aligned}\dot{\rho}_A + 3H(\rho_A + p_A) &= -\tilde{Q}, \\ \dot{\rho}_c + 3H\rho_c &= \tilde{Q},\end{aligned}$$

with

$$\rho_A = \left(\frac{m^2}{2} - V_X\right) A^2 + V, \quad p_A = \frac{m^2 A^2}{2} - V,$$

Interacting term: $\tilde{Q} = \rho_c \frac{\dot{A}}{2A} \gamma = \rho_c \frac{\dot{A}}{2A} (\gamma_C + \gamma_B),$

$$\gamma_C = \frac{-2\frac{C_X}{C} A^2 + A^4 \left(\frac{C_X^2}{C^2} - 2\frac{C_{XX}}{C}\right)}{\left(\frac{C_X}{C} A^2 - 2\right) \left(\frac{C_X}{C} A^2 - 1\right)}, \quad \gamma_B = \frac{A^2 (10B_X A^2 + 4B + A^6 (B_X^2 - 2B_{XX} B) + 2A^4 (B_{XX} - 3B_X B))}{(2 + B_X A^4)(1 + B_X A^4 + B A^2)}.$$

Dynamical system analysis:

$$x \equiv \sqrt{\frac{-V_X A^2}{3M_p^2 H^2}}; \quad y \equiv \sqrt{\frac{V}{3M_p^2 H^2}}; \quad z \equiv \sqrt{\frac{\rho_m}{3M_p^2 H^2}}; \quad r \equiv \sqrt{\frac{\rho_r}{3M_p^2 H^2}}; \quad u \equiv \sqrt{\frac{m^2 A/2}{3M_p^2 H^2}}; \quad v \equiv \frac{A}{M_p}.$$

$$x^2 + y^2 + z^2 + u^2 = 1. \quad V(X) = V_0 e^{-2\lambda X/M_p^2}.$$

-Autonomous system:

$$x' = x \left(-\epsilon_H + \frac{v'}{v} (1 - \lambda v^2) \right),$$

$$y' = -y (\epsilon_H + \lambda v' v),$$

$$z' = \frac{z}{2} \left(-2\epsilon_H - 3 + \frac{\gamma v'}{2v} \right),$$

$$r' = -r (2 + \epsilon_H),$$

$$\frac{v'}{v} = \frac{6u^2 + 3x^2}{4\lambda^2 v^4 y^2 - x^2 - 2u^2 - \frac{\gamma}{2} z^2}.$$

Conformal case:

-Model 1: $C(X) = C_0 \left(\frac{X}{M_p^2} \right)^q \rightarrow \gamma_C = \frac{2q}{1-2q}.$

-Model 2: $C(X) = \tilde{C}_0 e^{\frac{4\alpha X}{M_p^2}} \rightarrow \gamma_C = -\frac{4\alpha v^2(1+2v^2\alpha)}{1-6\alpha v^2+8\alpha^2 v^4}.$

Point	r_c	y_c	z_c	u_c	v_c	Ω_A	w_A	w_{eff}	Existence	Acceleration
(A $_{\pm}$)	± 1	0	0	0	0	0	-	1/3	$\forall q(\alpha), \lambda$	No
(B $_{\pm}$)	0	0	± 1	0	0	0	-	0	$\forall q(\alpha), \lambda$	No
($\tilde{B}_{1,2}$)	0	0	± 1	0	$\mp \frac{1}{\sqrt{2\alpha}}$	0	-	0	$\alpha \neq 0, \forall \lambda$	No
($\tilde{B}_{3,4}$)	0	0	± 1	0	$\pm \frac{1}{\sqrt{2\alpha}}$	0	-	0	$\alpha \neq 0, \forall \lambda$	No
(C $_{1,2}$)	0	0	$\pm \sqrt{\frac{-2+6q}{-2+5q}}$	$\mp \sqrt{\frac{q}{2-5q}}$	0	$\frac{q}{2-5q}$	1	$\frac{q}{2-5q}$	$q \neq 2/5, 0 < q < 1/3, \forall \lambda$	$2/5 < q < 1$
(C $_{3,4}$)	0	0	$\pm \sqrt{\frac{-2+6q}{-2+5q}}$	$\pm \sqrt{\frac{q}{2-5q}}$	0	$\frac{q}{2-5q}$	1	$\frac{q}{2-5q}$	$q \neq 2/5, 0 < q < 1/3, \forall \lambda$	$2/5 < q < 1$
($\tilde{D}_{1,2}$)	0	1	$\pm \sqrt{-\frac{\lambda}{\alpha}}$	0	$\mp \frac{1}{\sqrt{2\alpha}}$	$1 + \frac{\lambda}{2\alpha}$	$-\frac{1}{1+\frac{\lambda}{2\alpha}}$	-1	$\alpha \neq 0, \lambda < 0 \alpha < 0$	Yes
($\tilde{D}_{3,4}$)	0	1	$\pm \sqrt{-\frac{\lambda}{\alpha}}$	0	$\pm \frac{1}{\sqrt{2\alpha}}$	$1 + \frac{\lambda}{2\alpha}$	$-\frac{1}{1+\frac{\lambda}{2\alpha}}$	-1	$\alpha \neq 0, \lambda < 0 \alpha < 0$	Yes
($\tilde{F}_{1,2}$)	0	1	$\pm \sqrt{-\frac{\lambda}{2\alpha}}$	0	$\mp \frac{1}{2\sqrt{\alpha}}$	$1 + \frac{\lambda}{2\alpha}$	$-\frac{1}{1+\frac{\lambda}{2\alpha}}$	-1	$\alpha \neq 0, \lambda < 0 \alpha < 0$	Yes
($\tilde{F}_{3,4}$)	0	1	$\pm \sqrt{-\frac{\lambda}{2\alpha}}$	0	$\pm \frac{1}{2\sqrt{\alpha}}$	$1 + \frac{\lambda}{2\alpha}$	$-\frac{1}{1+\frac{\lambda}{2\alpha}}$	-1	$\alpha \neq 0, \lambda < 0 \alpha < 0$	Yes
(E $_{\pm}$)	0	0	0	± 1	0	1	1	1	$\forall q(\alpha), \lambda$	No
(D $_{1,2}$)	0	$\pm \frac{1}{\sqrt{1+v^2\lambda}}$	0	$\mp v \sqrt{\frac{\lambda}{-1-v^2\lambda}}$	-	1	-1	-1	$\forall q(\alpha), \text{eqn. (36)}$	$\forall q(\alpha), \lambda$
(D $_{3,4}$)	0	$\pm \frac{1}{\sqrt{1+v^2\lambda}}$	0	$\pm v \sqrt{\frac{\lambda}{-1-v^2\lambda}}$	-	1	-1	-1	$\forall q(\alpha), \text{eqn. (36)}$	$\forall q(\alpha), \lambda$
(S)	0	1	0	0	0	1	-1	-1	$\forall q(\alpha), \lambda$	$\forall q(\alpha), \lambda$

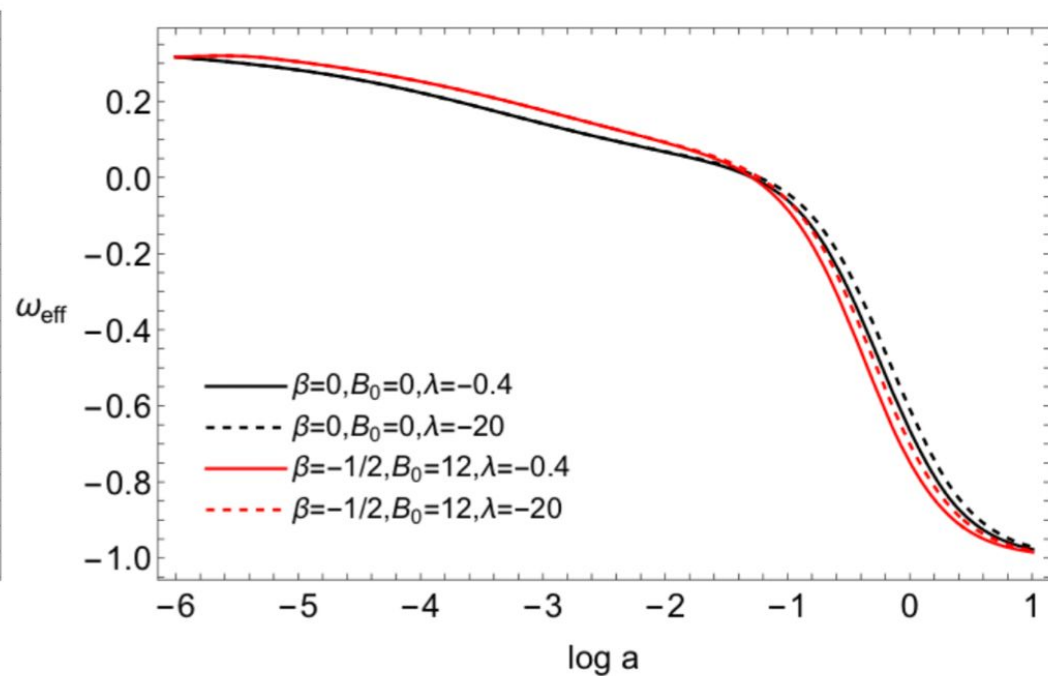
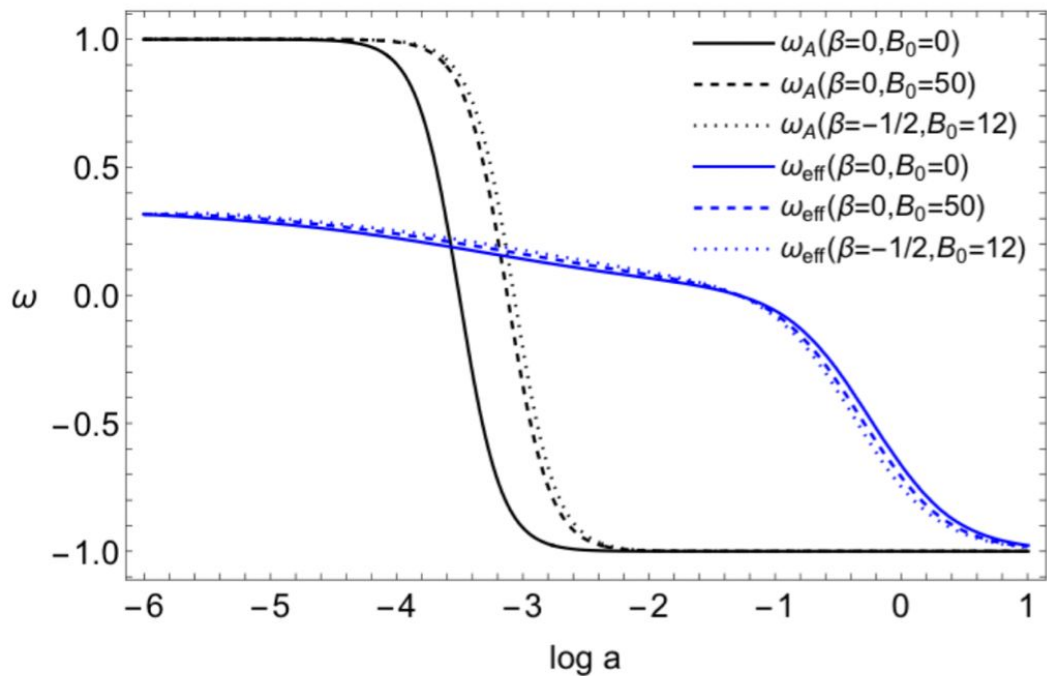
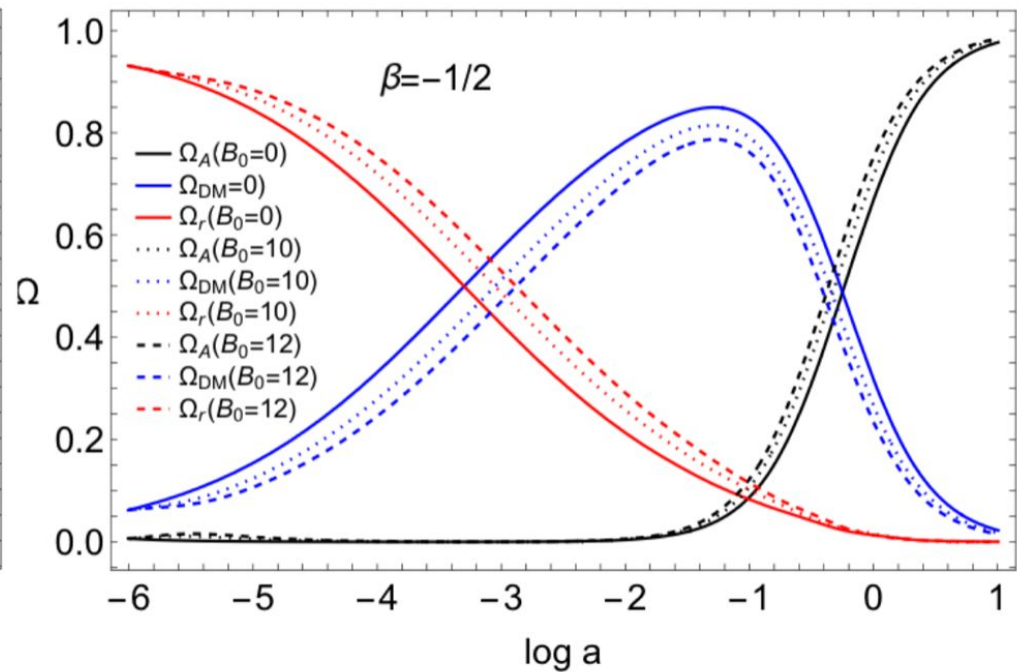
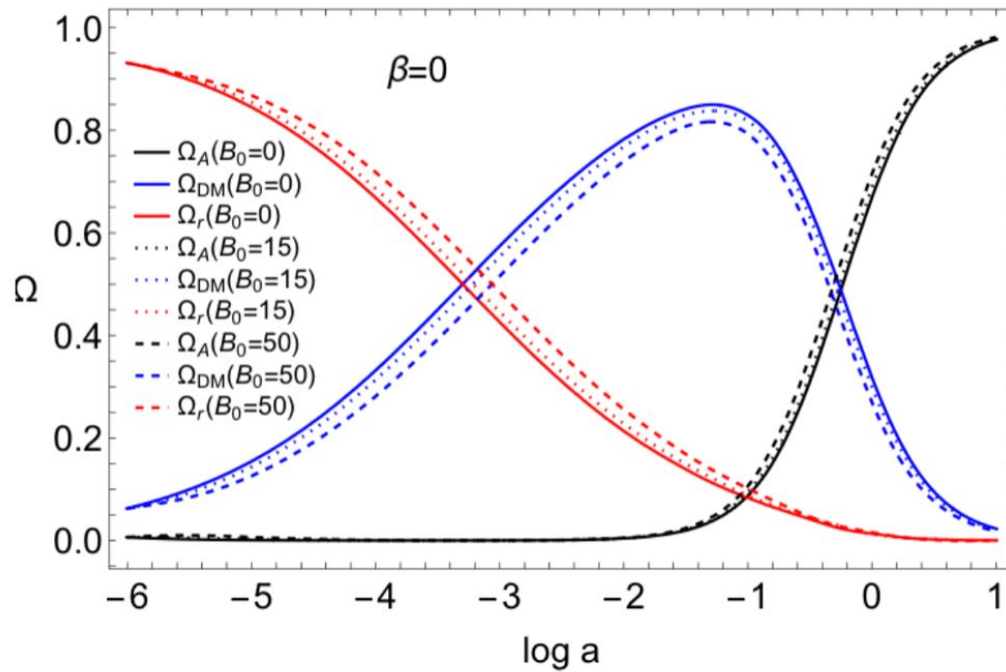
Disformal case:

-Model: $C(X) = 1$, $B(X) = B_0 \frac{2^\beta X^\beta}{M_p^{2+2\beta}}$

$$\rightarrow \gamma_B = -B_0 \frac{2v^{2+2\beta}(1+\beta)(-1+\beta(-2+B_0v^{2+2\beta}))}{(1+\beta B_0v^{2+2\beta})(1+v^{2+2\beta}(1+2\beta)B_0)}.$$

Point	r_c	y_c	z_c	u_c	v_c	Ω_A	w_A	w_{eff}	Existence	Acceleration
$(H_{1,2})$	0	± 1	$\pm 2 \frac{\sqrt{-2\lambda}}{B_0}$	0	$\mp \frac{2}{B_0}$	$1 + \frac{8\lambda}{B_0^2}$	$-\frac{1}{1 + \frac{8\lambda}{B_0^2}}$	-1	$B_0 \neq 0, \lambda < 0$	Yes
$(H_{3,4})$	0	± 1	$\pm 2 \frac{\sqrt{-2\lambda}}{B_0}$	0	$\mp \frac{2}{B_0}$	$1 + \frac{8\lambda}{B_0^2}$	$-\frac{1}{1 + \frac{8\lambda}{B_0^2}}$	-1	$B_0 \neq 0, \lambda < 0$	Yes
$(G_{1,2})$	0	0	± 1	0	$\pm \frac{2}{B_0}$	0	-	0	$B_0 \neq 0$	No
$(G_{3,4})$	0	0	± 1	0	$\mp \frac{2}{B_0}$	0	-	0	$B_0 \neq 0$	No
$(\tilde{H}_{1,2})$	0	± 1	$\pm \sqrt{2 - 2(1 + cB_0\lambda)}$	0	$\mp \sqrt{cB_0}$	$1 + 2cB_0\lambda$	$-\frac{1}{1 + 2cB_0\lambda}$	-1	$B_0 > 0, \lambda < 0$	Yes
$(\tilde{H}_{3,4})$	0	± 1	$\pm \sqrt{2 - 2(1 + cB_0\lambda)}$	0	$\pm \sqrt{cB_0}$	$1 + 2cB_0\lambda$	$-\frac{1}{1 + 2cB_0\lambda}$	-1	$B_0 > 0, \lambda < 0$	Yes
$(\tilde{G}_{1,2})$	0	0	± 1	0	$\pm \sqrt{cB_0}$	0	-	0	$B_0 > 0$	No
$(\tilde{G}_{3,4})$	0	0	± 1	0	$\mp \sqrt{cB_0}$	0	-	0	$B_0 > 0$	No

Cosmological background evolution:



Conclusions and perspectives:

- We have proposed a novel coupling between vector fields and CDM at the level of the action through a general disformal transformation.
- Several critical points arise in comparison to the uncoupled case, leading to interesting cosmological solutions: **stable attractor solutions, scaling solutions.**
- The effects of the disformal coupling can affect both the early and late-time universe, leaving some imprints that can be contrasted with observational data at different redshifts.
- Study disformal coupling for spatial vector fields: **Muti-fields, non-Abelian vector fields.**
- Involve the field strength in the disformal transformation or its dual in order to make the structure of the group explicit.
- Constrain the background dynamics, matter density perturbations: spherical collapse, number countst, the growth rate and the redshift-space distortion.



Thank you!