

RD51 Collaboration Meeting

An update on the modelling of signal formation in detectors with resistive elements.

Djunes Janssens

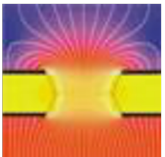
djunes.janssens@cern.ch

CERN Supervisor: Dr. Heinrich Schindler

University Supervisor: Prof. Dr. Freya Blekman

Co-Supervisors: Dr. Eraldo Oliveri, Dr. Werner Riegler and Dr. Rob Veenhof

June 17th, 2021



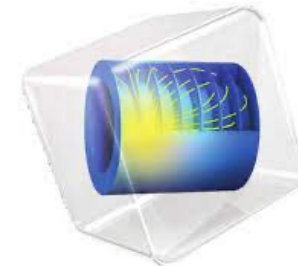
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Introduction

We want to use Garfield++ and COMSOL to model **the signal formation in detectors with resistive elements.**

Outline:

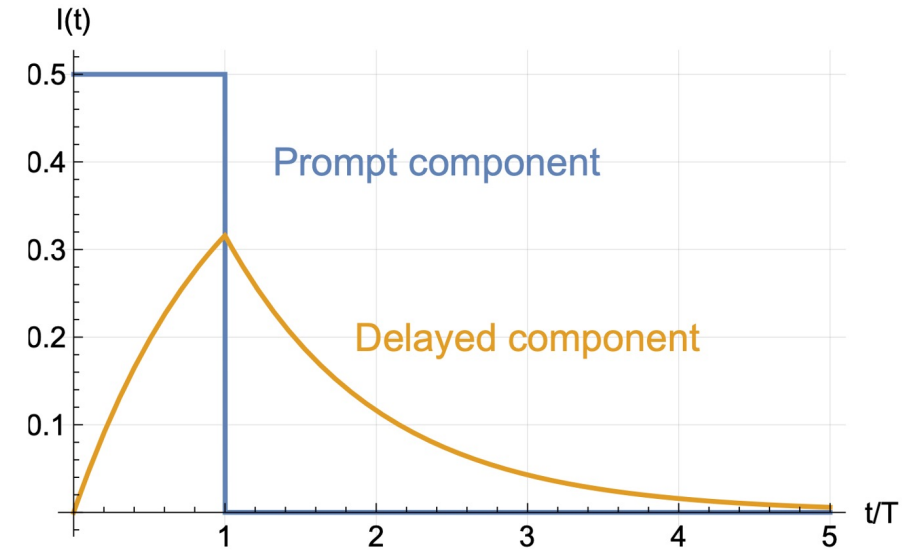
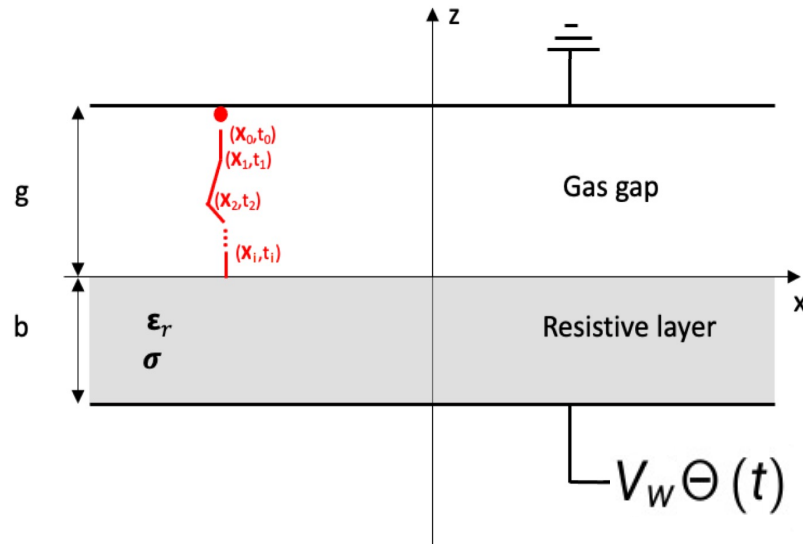
- Components of the dynamical weighting potential
- Simulating a MRPC
- Signal formation in an AC-coupled LGAD
- Signal formation in a MicroCAT readout structure
- Summary



Components of the dynamical weighting potential

The time-dependent weighting potential is a superposition of a static **prompt** and a dynamic **delayed** component.

The prompt component of your signal is the instantaneous induction of current on the electrode source by the movement of the charge carrier.

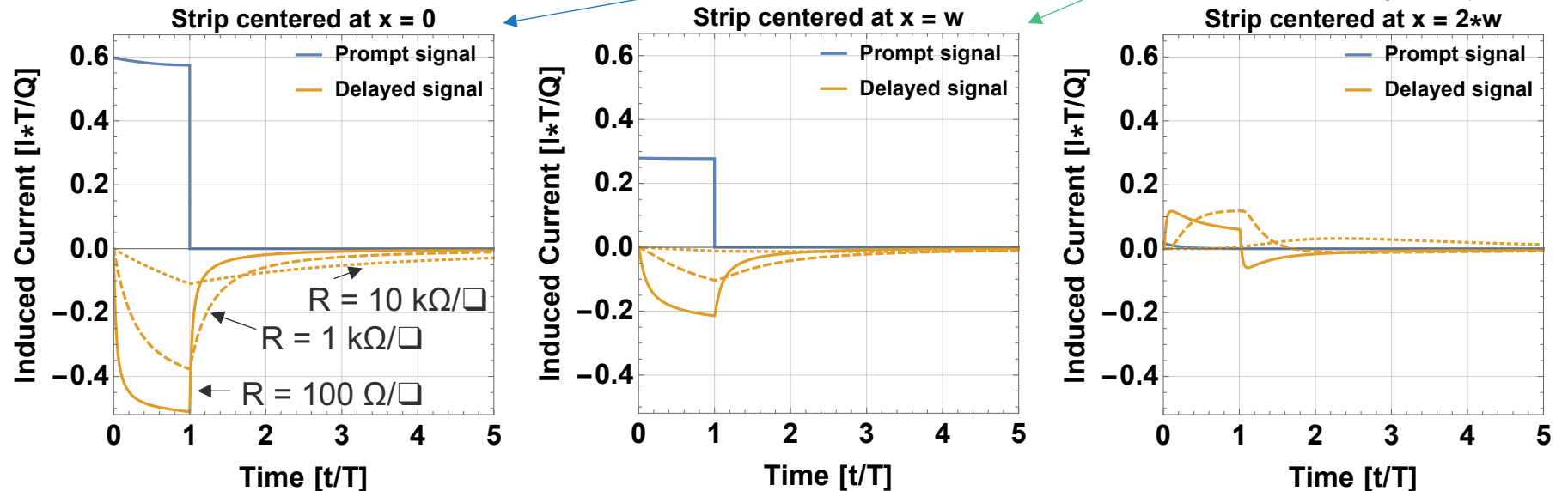
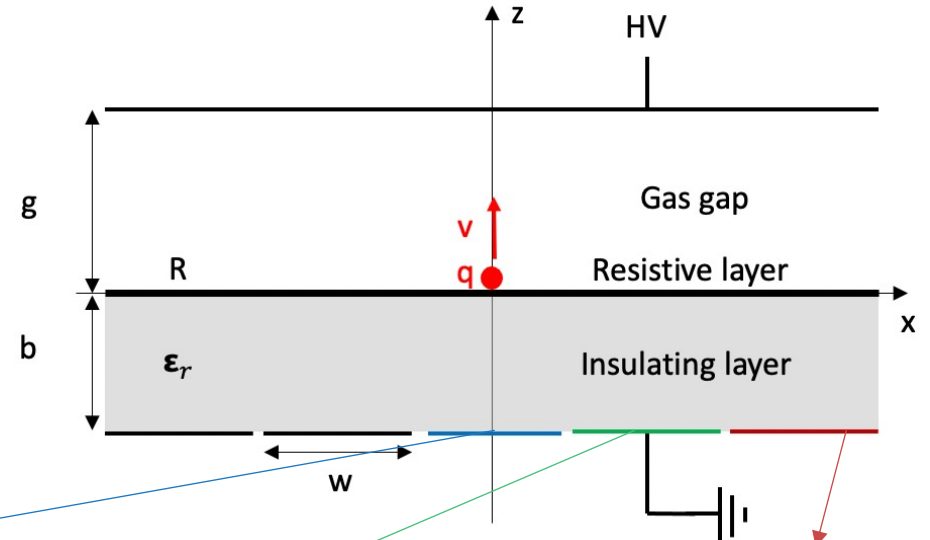


Components of the dynamical weighting potential

The **delayed component** of your signal encodes the responds of the resistive material to the field lines of this charge carrier.

Let us take a simple geometry to get an intuition:

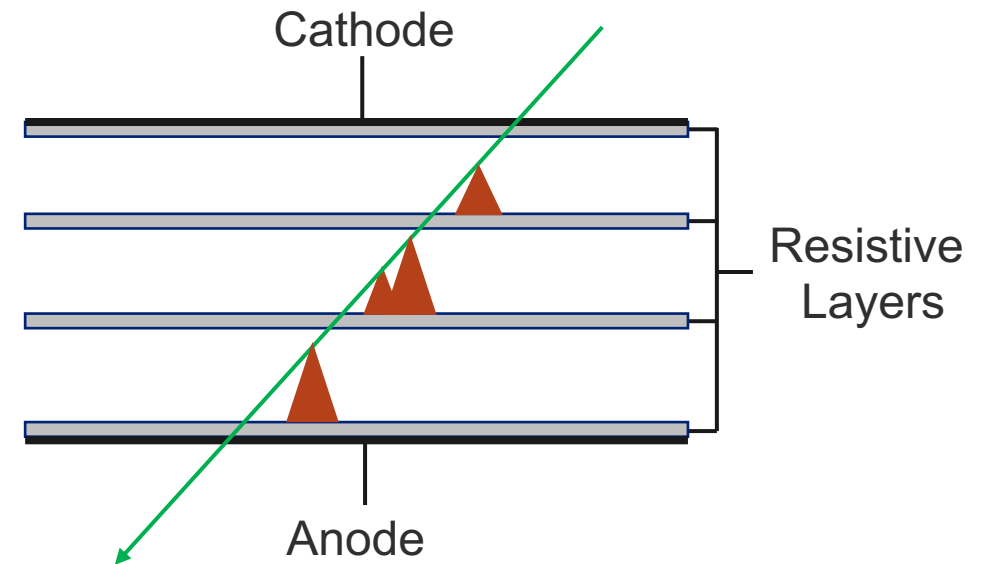
- $g = 50 \mu\text{m}$
- $b = 5 \mu\text{m}$
- $w = 200 \mu\text{m}$
- $T = g/v = 0.4 \text{ ns}$



Simulating a MRPC

In order to relatively efficiently simulate a MRPC two separate questions ought to be addressed:

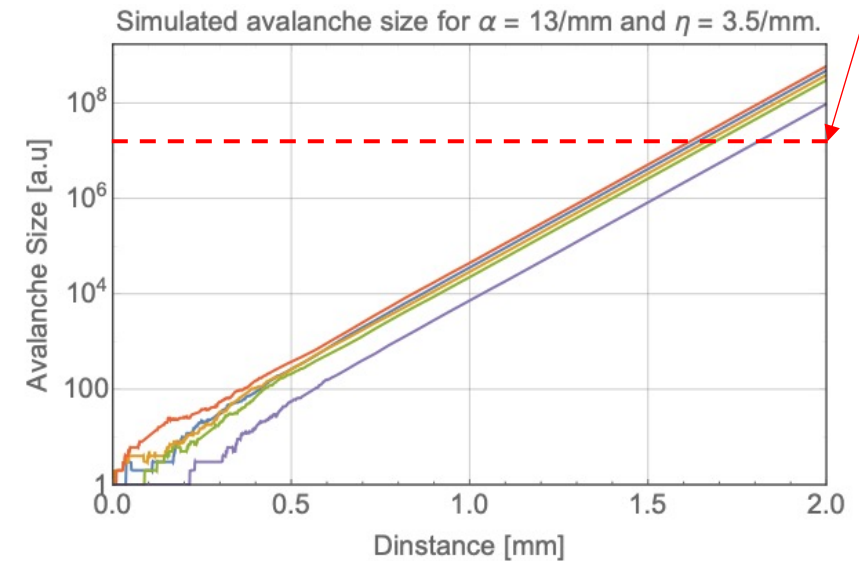
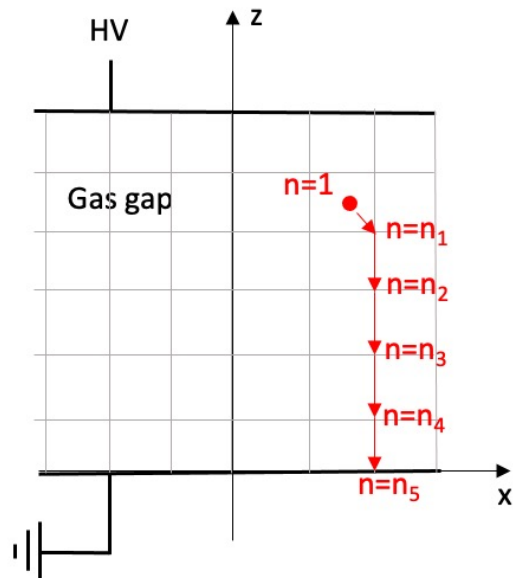
- A grid-based Monte Carlo simulation for the avalanche dynamics
- The weighting potential of the readout electrodes.



Grid-based avalanche dynamics calculations

When propagating electrons through nodes of a lattice, the growth in its resulting effective Townsend avalanche can be done using a simple Monte Carlo method.

Due to space-charge effects there is a suppression of avalanche growth. → Saturation at $1.6 \cdot 10^7$ electrons.



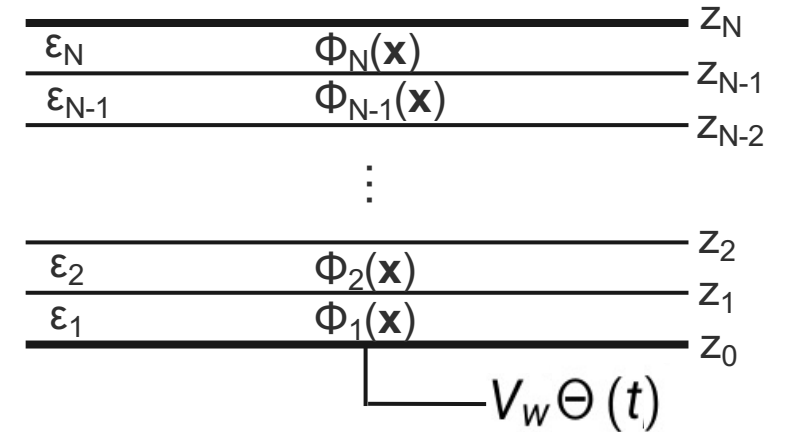
Weighting potentials in a MRPC

On the question of the weighting potential: for a strip of width w_x centered at the origin when $z_{m-1} < z < z_m$ it is:

$$\phi_m(x, z) = \frac{2V_w}{\pi} \int_0^\infty \cos(kx) \sin\left(\frac{kw_x}{2}\right) \frac{h_m(k, z)}{k} dk .$$

In the case of a readout plane the solution is:

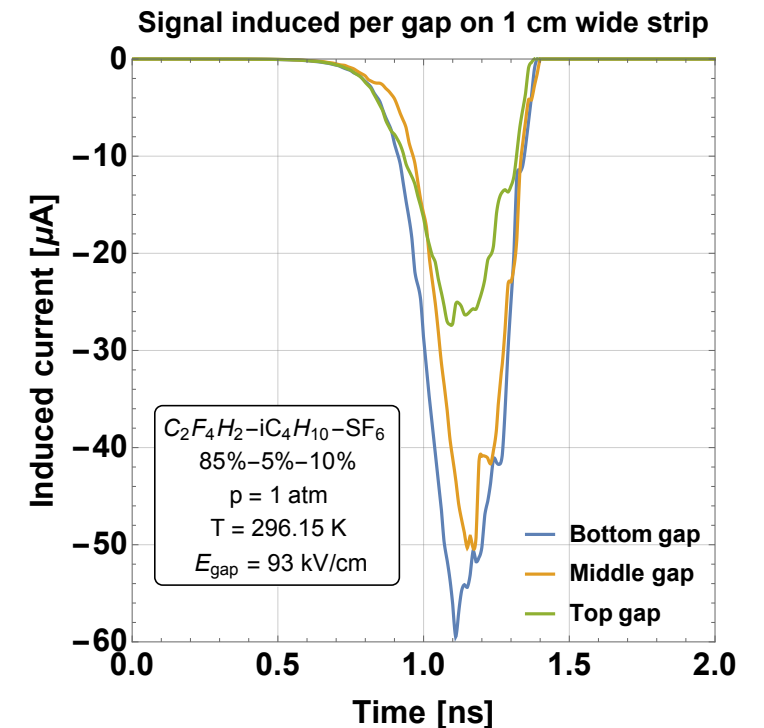
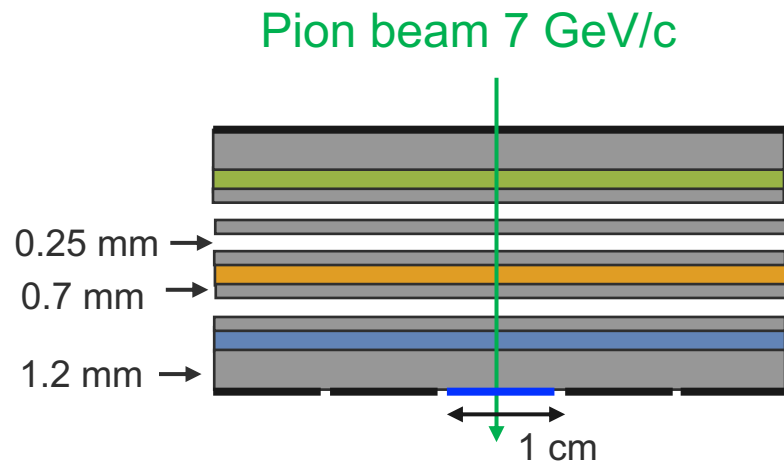
$$\phi_m(z) = V_w - \frac{V_w}{\sum_{l=1}^N \frac{z_l - z_{l-1}}{\epsilon_l}} \left(\frac{z - z_{m-1}}{\epsilon_m} + \sum_{n=1}^{m-1} \frac{z_n - z_{n-1}}{\epsilon_n} \right) .$$



Strip width and the contributions of layers

In general, a readout strip's signal will not be comprised equally from that induced by each layer.

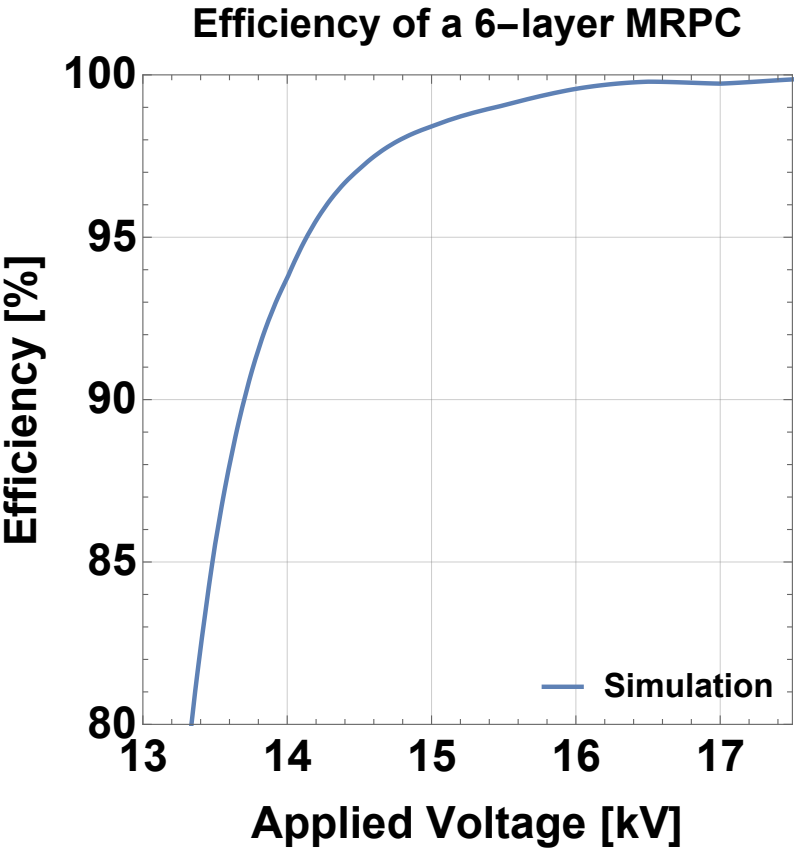
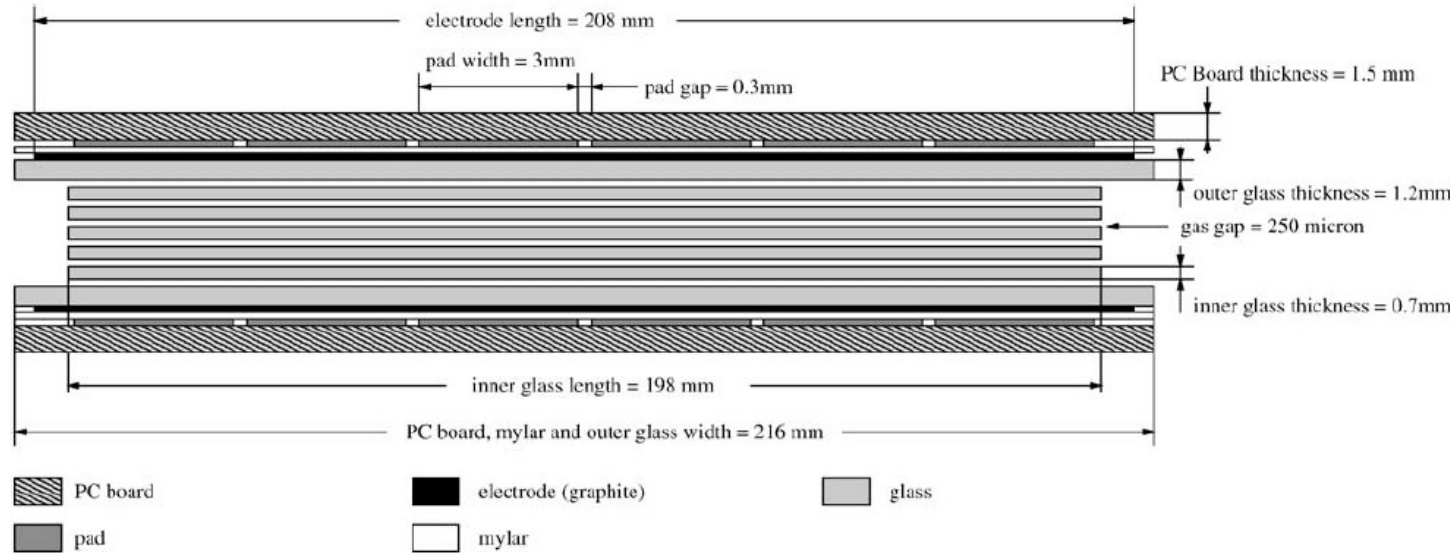
This imbalance will decrease when the strip's width increases.



Next step: benchmark the simulations

We will do a first comparison of the efficiency and time to measurements taken with a MRPC

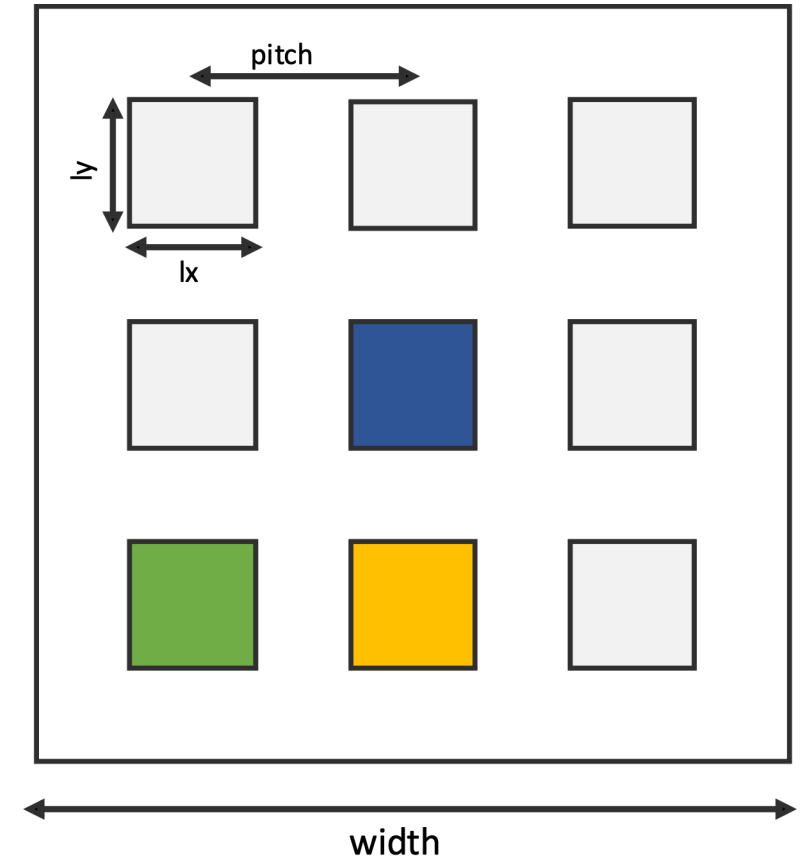
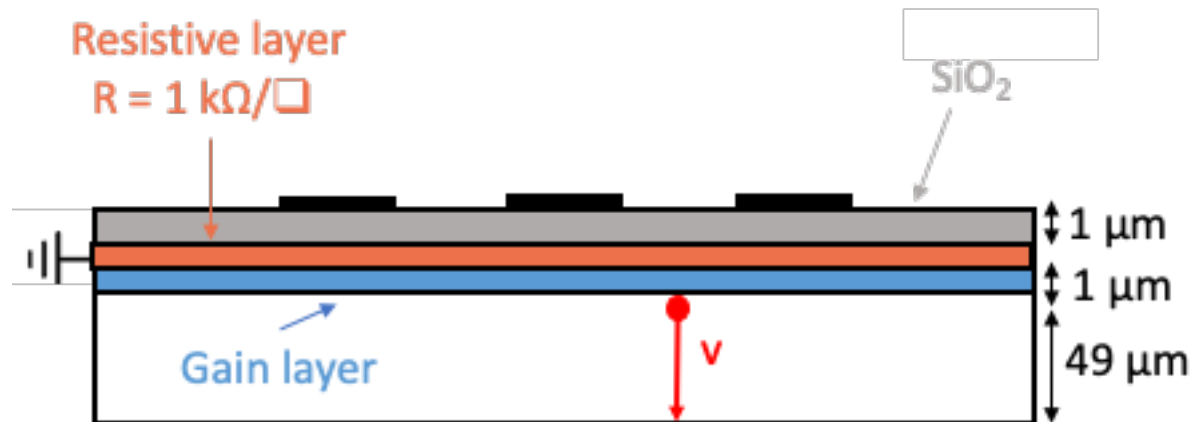
Here a discriminator threshold of 30 fC is assumed.



Weighting potential for AC-coupled LGAD

Using COMSOL the dynamic weighting potential can be calculated for the AC-coupled LGAD's readout system.

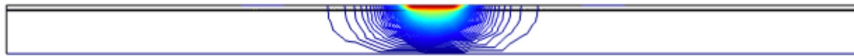
This has been done for five different geometries.



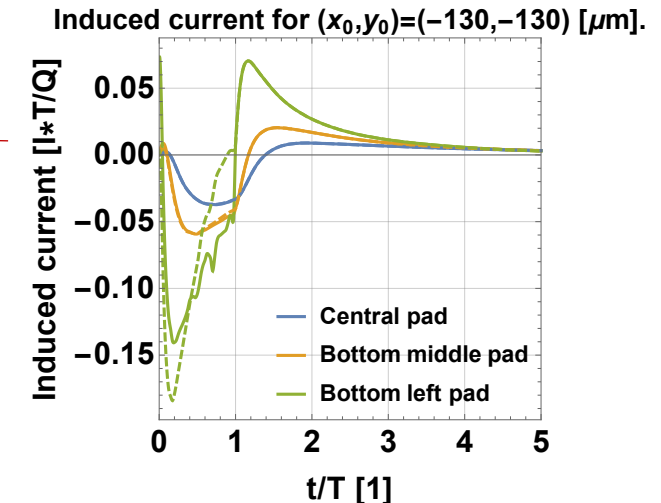
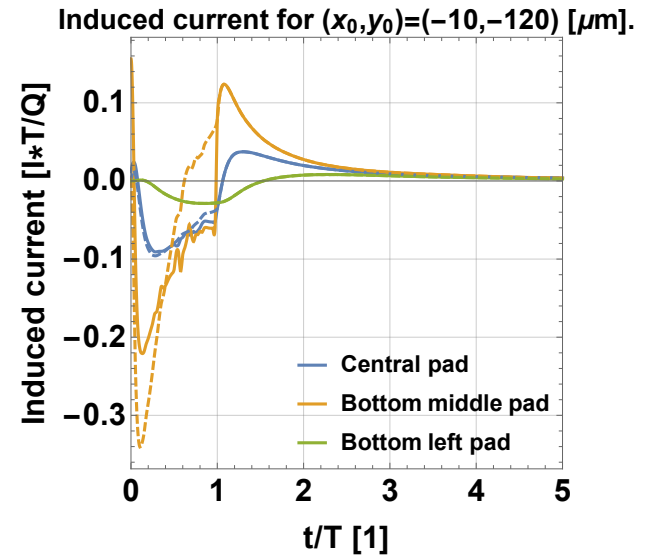
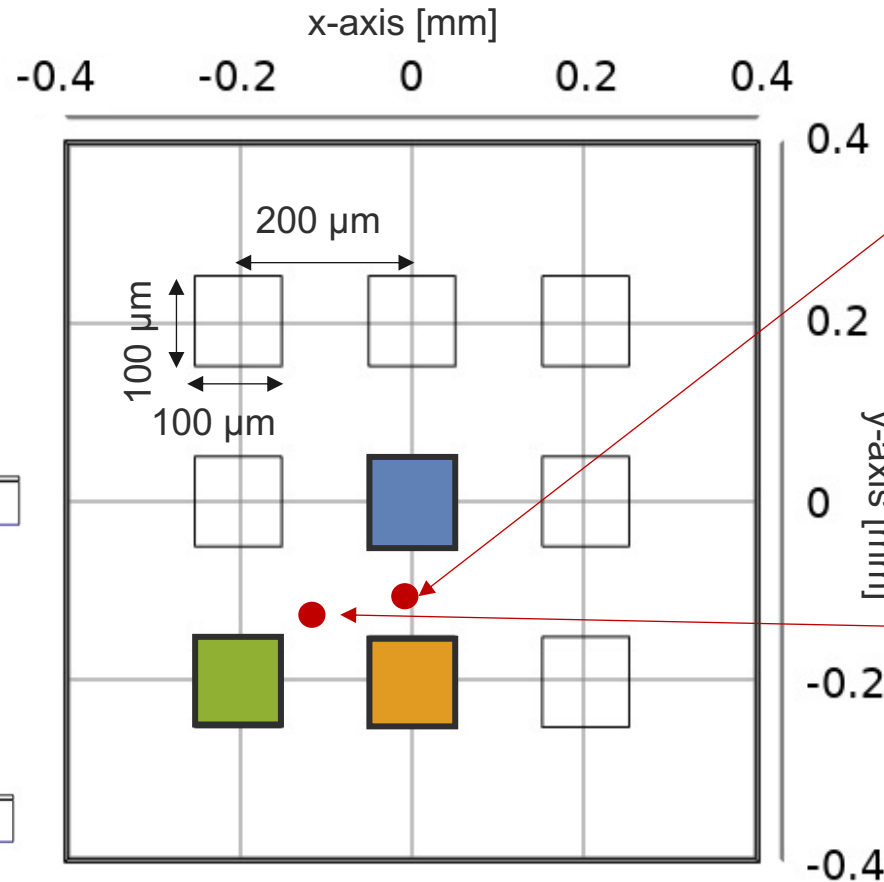
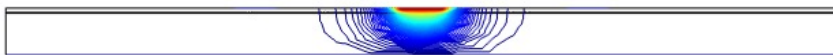
Evaluating signals for a 100-200 geometry

To understand the roll that the prompt component plays in this signal one can look at **the central pad on the bottom row**.

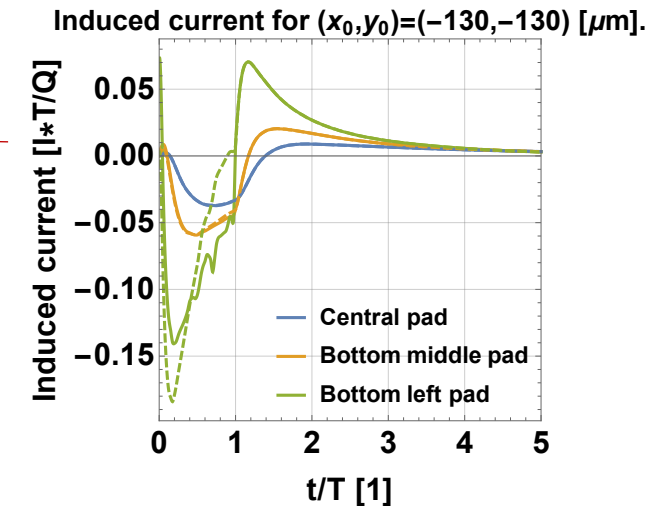
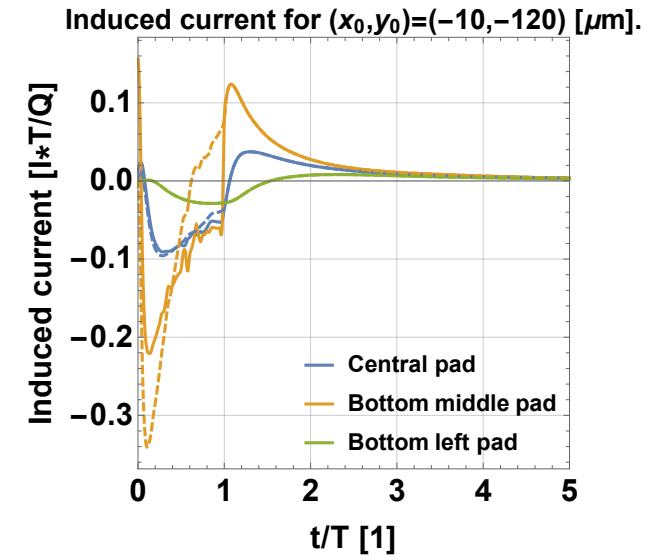
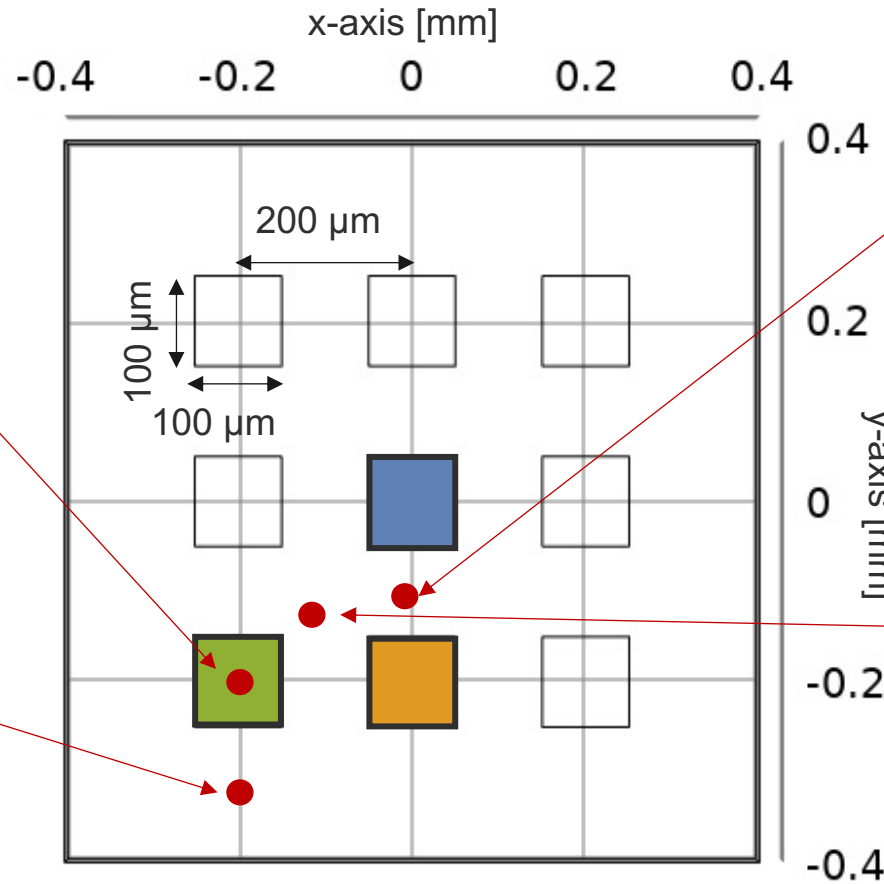
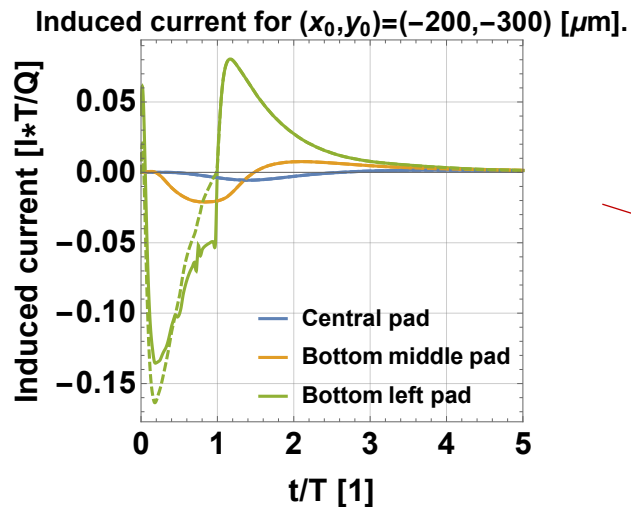
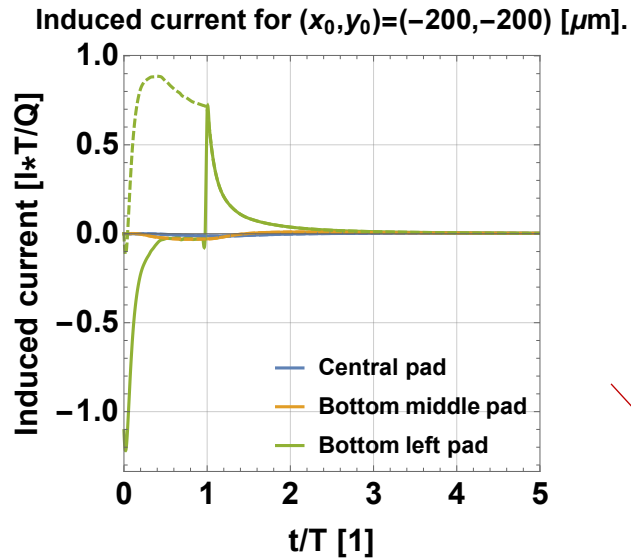
Prompt weighting potential:



Full weighting potential:

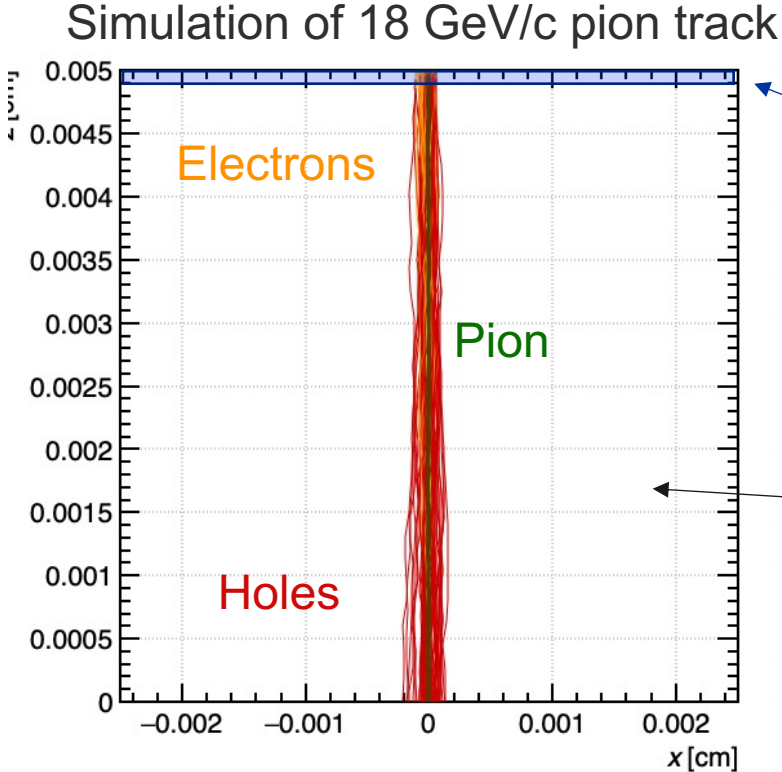


Evaluating signals for a 100-200 geometry



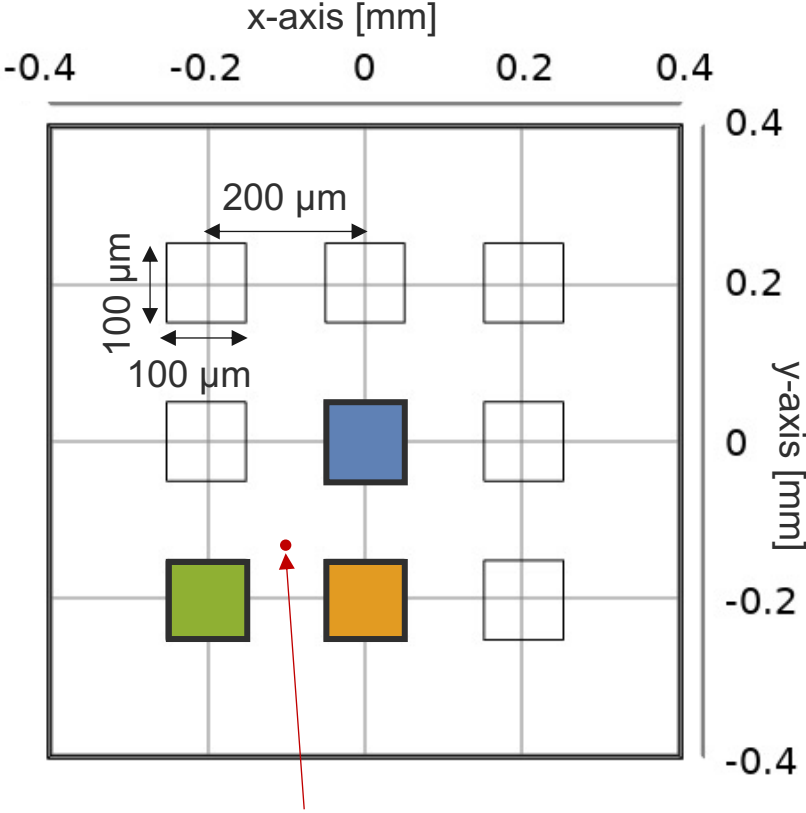
Induced current for an 18 GeV/c pion track

Given a pion tracking through the sensor the resulting signal can be calculated.



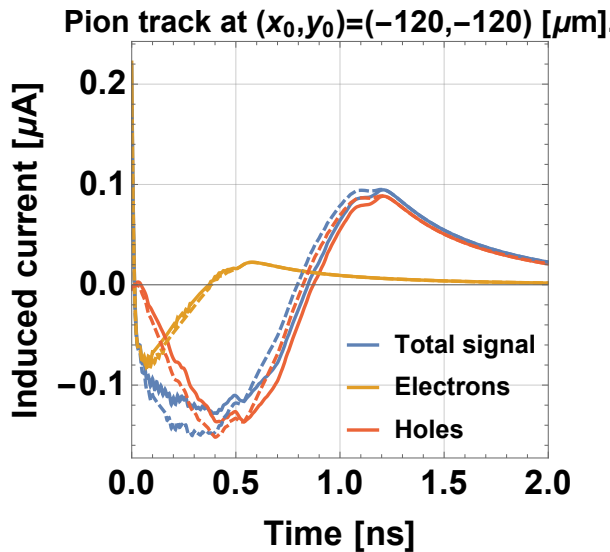
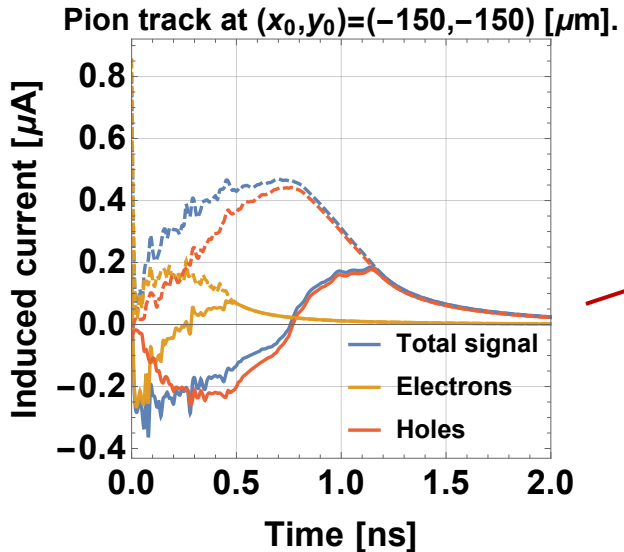
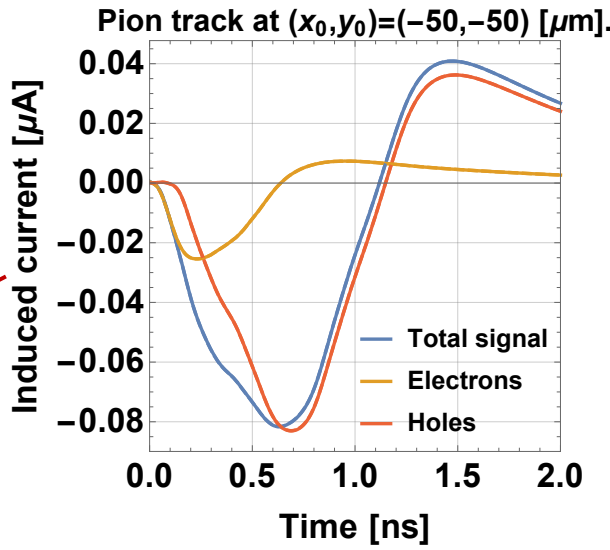
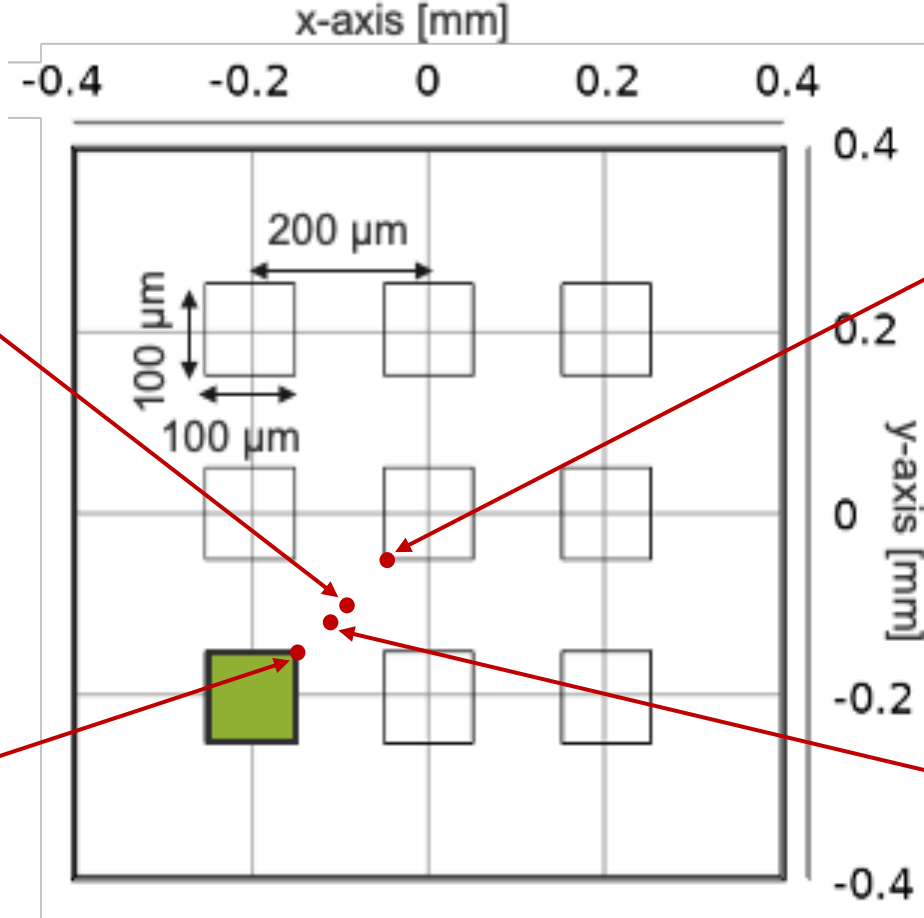
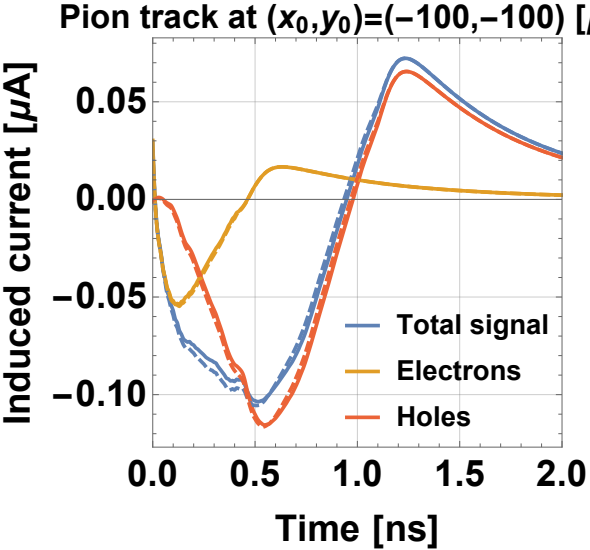
Gain layer of 1 μm thick with $E = 300 \text{ kV/cm}$

Bulk with $E = 38 \text{ kV/cm}$



An example of the position of the pion track

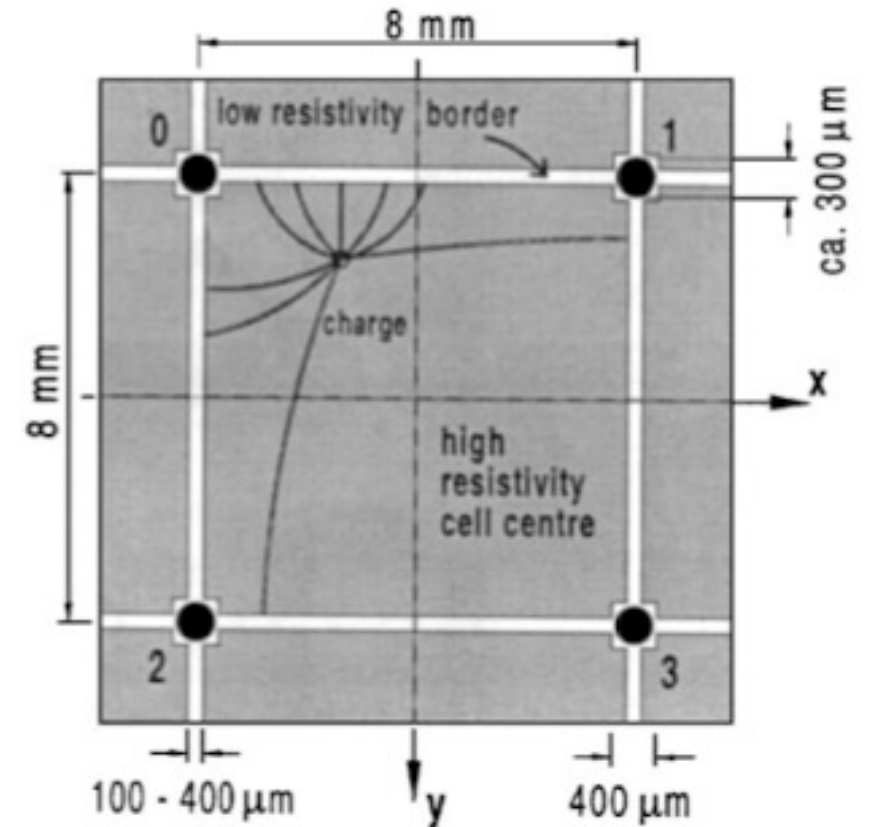
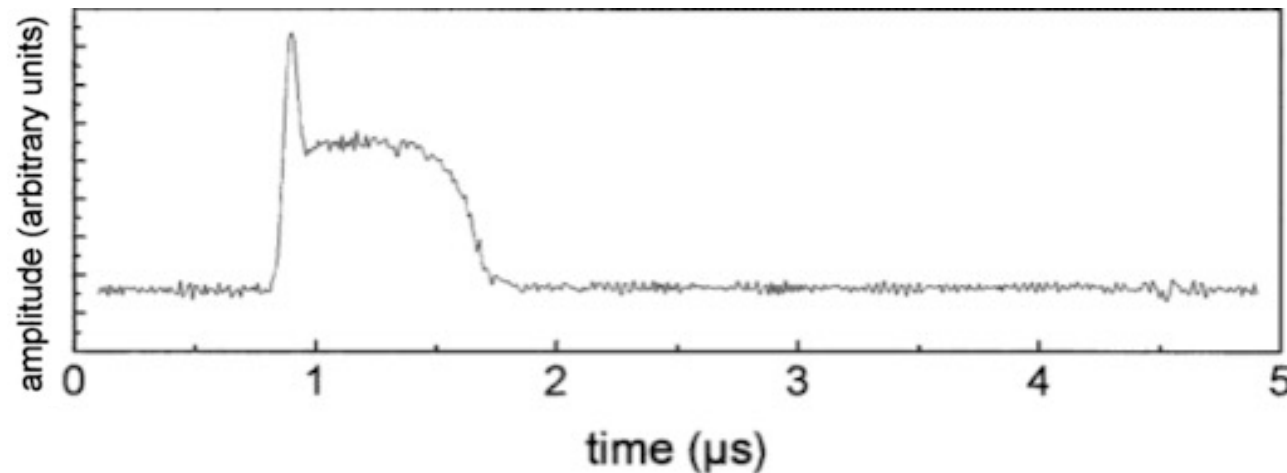
Induced current for an 18 GeV/c pion track



Signal formation in a MicroCAT detector

As an further example we take the MicroCAT's two-dimensional interpolating readout structure.

A reduced number of electronic readout channels nevertheless reaching a spatial nevertheless reaching good spatial resolution by using the delayed component of the signal.

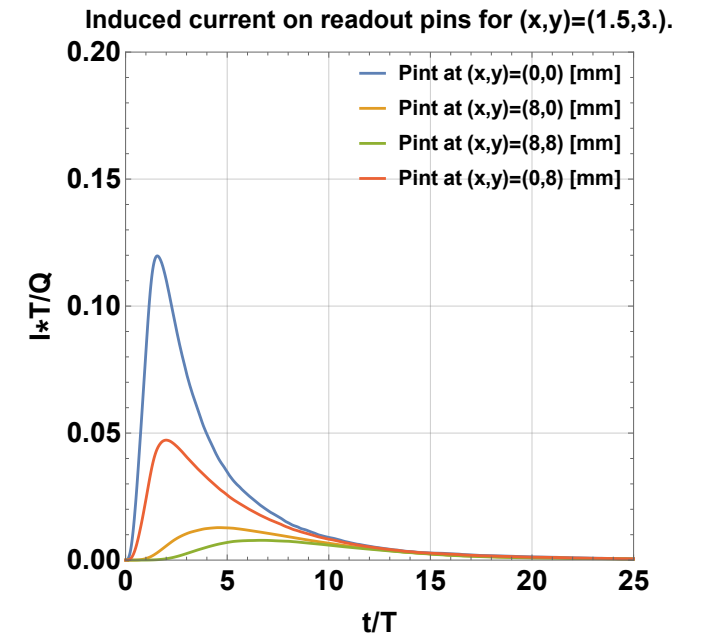
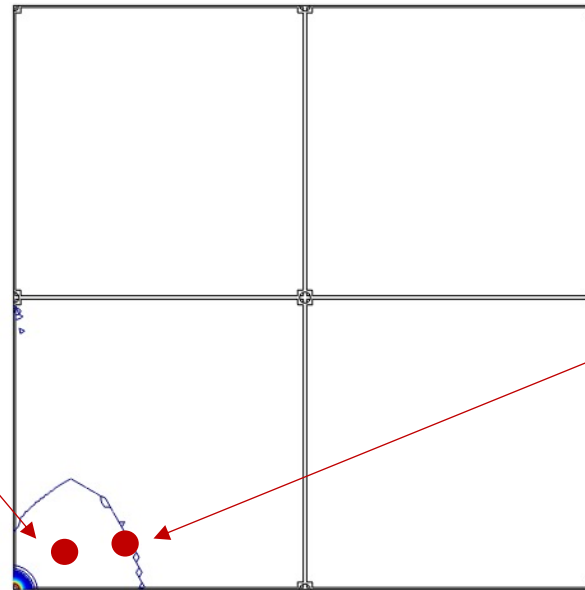
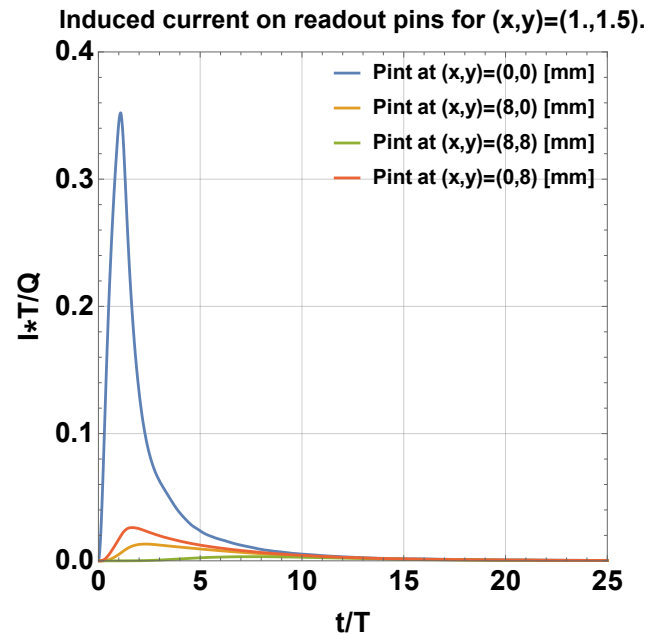


H. Wagner et al. Nucl.Instrum.Meth. A 482 (2002) 334–346

Signal formation in a MicroCAT detector

The induced current on the readout pins on one pad has been calculated for a charge carrier traversing an induction gap of 200 μm downwards in $T = 4$ ns. The surface resistivity is:

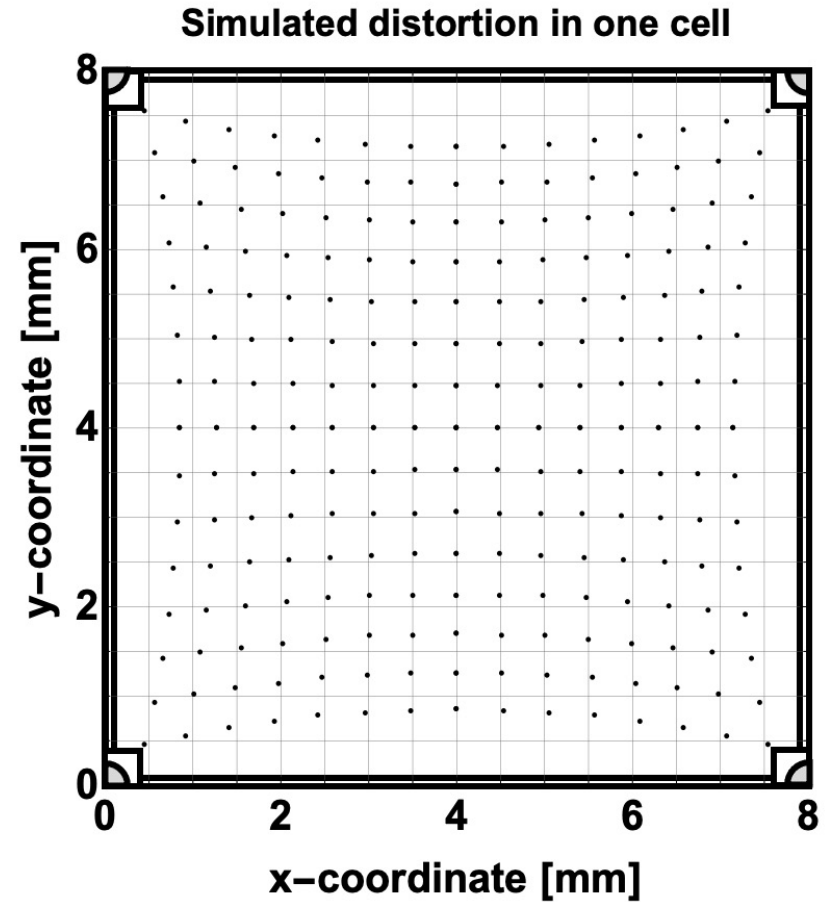
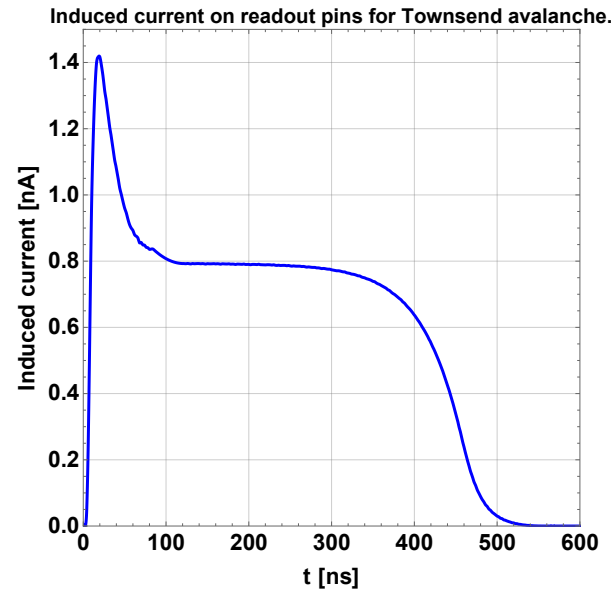
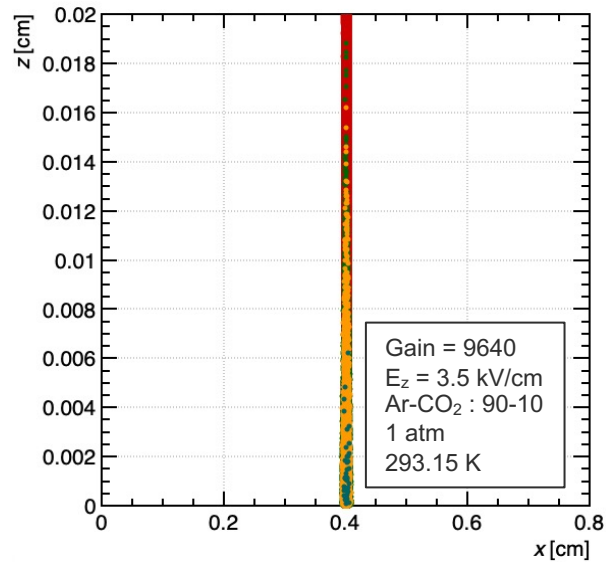
- $R_{\text{pad}} = 100 \text{ k}\Omega/\square$
- $R_{\text{strip}} = 1 \text{ k}\Omega/\square$



Signal formation in a MicroCAT detector

The same can be done for a Townsend avalanche.

Here a single electron has been placed in a uniform electric field, starting on the top of the induction gap.



Summary

We want to use Garfield++ and COMSOL to model **the signal formation in detectors with resistive elements**.

- Efficient multigap resistive plate chamber simulations will soon be possible in Garfield++.
- The tools developed during the project allow us to make a full description of the signal of:
 - The MicroCAT readout,
 - The AC-coupled LGAD.

Outlook:

- Benchmarking MRPC simulations against measurements.
- Using the developed tools to look at more detectors with resistive elements.

Thank you for your attention!