



# Reinterpretation of novel CMS analysis in low-mass LLP models

Andrea Mitridate

based on 21xx.xxxx with Michele Papucci and Christina Wang



Caltech



# INTRODUCTION

- current ATLAS/CMS searches for light LLPs **limited** by
  - **SM backgrounds** (calorimeter not thick enough to veto SM)
  - LLPs decay to few charged particles (**# tracks** controlled by  $m_{\text{LLP}}$ )
- recent CMS search looking for clusters in the endcap muon detector system ameliorates these problems
  - Steel → LLPs decay **products can shower**
  - Shower → **signature tracks**  $E_{\text{LLP}}$  rather than  $m_{\text{LLP}}$
  - Steel → exceptional **shielding from SM backgrounds**
- we use public data to recast the analysis for
  - **Higgs decays into scalar/dark photon** (closest to the model considered in the original CMS analysis, decoupled production and decay channels)
  - **Inelastic Dark Matter** (decouples  $E_{\text{LLP}}$  from MET)
  - **Axion-like particles** (production and decay channels controlled by the same parameter, enhanced production at high  $E_{\text{LLP}}$ )
  - **Dark showers** (studies the impact of isolation cuts)

# RECAST STRATEGY

## Event Generation

**MadGraph + Pythia8** generate the hard processes, shower, and hadronize them

**loose generator level cuts** (MET cut, isolation veto, ecc..) are applied to increase statistics

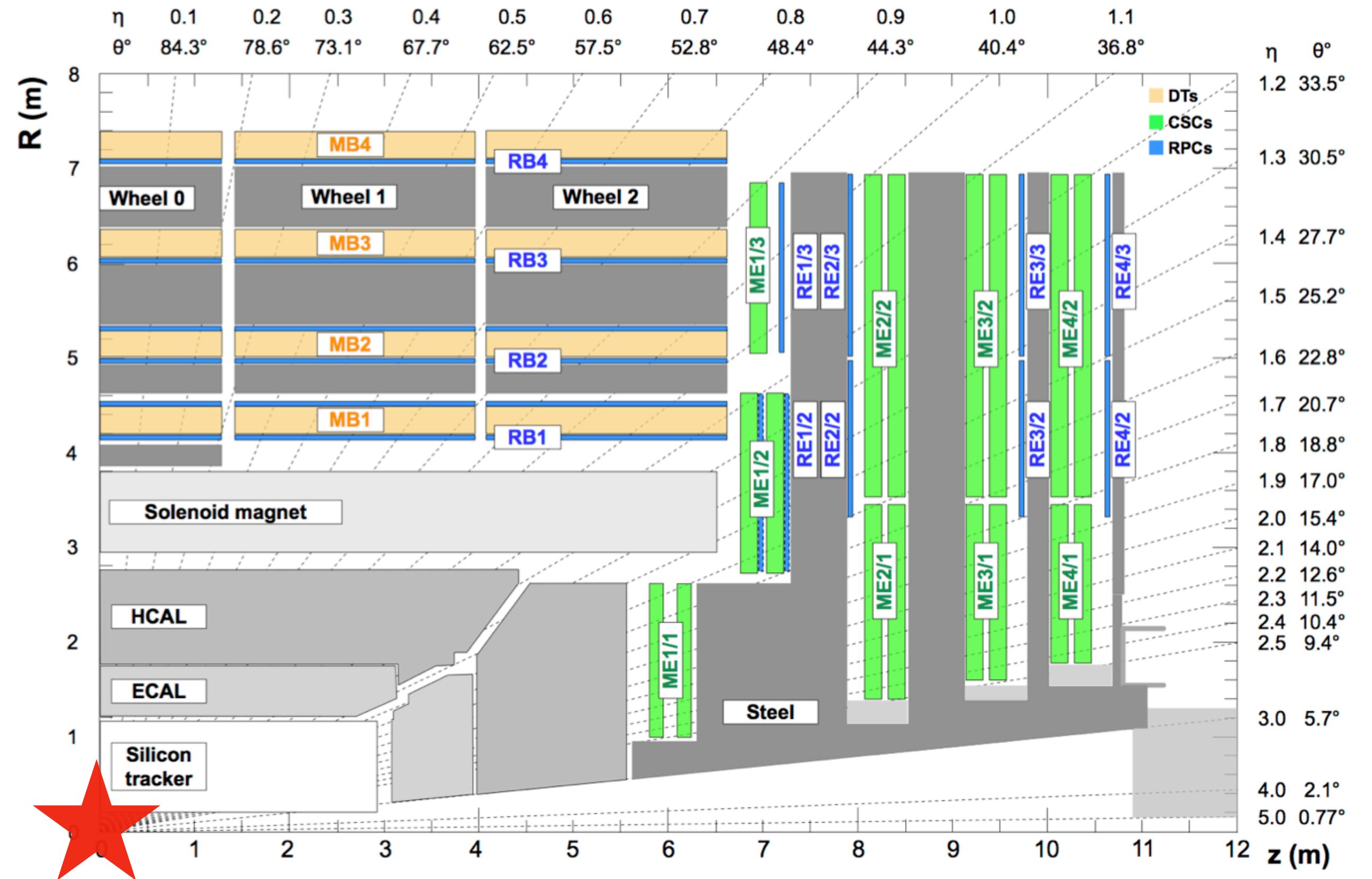
## Custom Decayer

given LLP lifetime and energy derive the probability to decay in each detector region

by using Pythia8 **generates a decay event in each detector region** and weights it with the probability to decay in that region

## Detector simulation

uses Delphes, with the publicly available module to take to account **cluster** and **cut-based ID** efficiencies for LLPs decays



# RECAST STRATEGY

## Event Generation

**MadGraph + Pythia8** generate the hard processes, shower, and hadronize them

**loose generator level cuts** (MET cut, isolation veto, ecc..) are applied to increase statistics

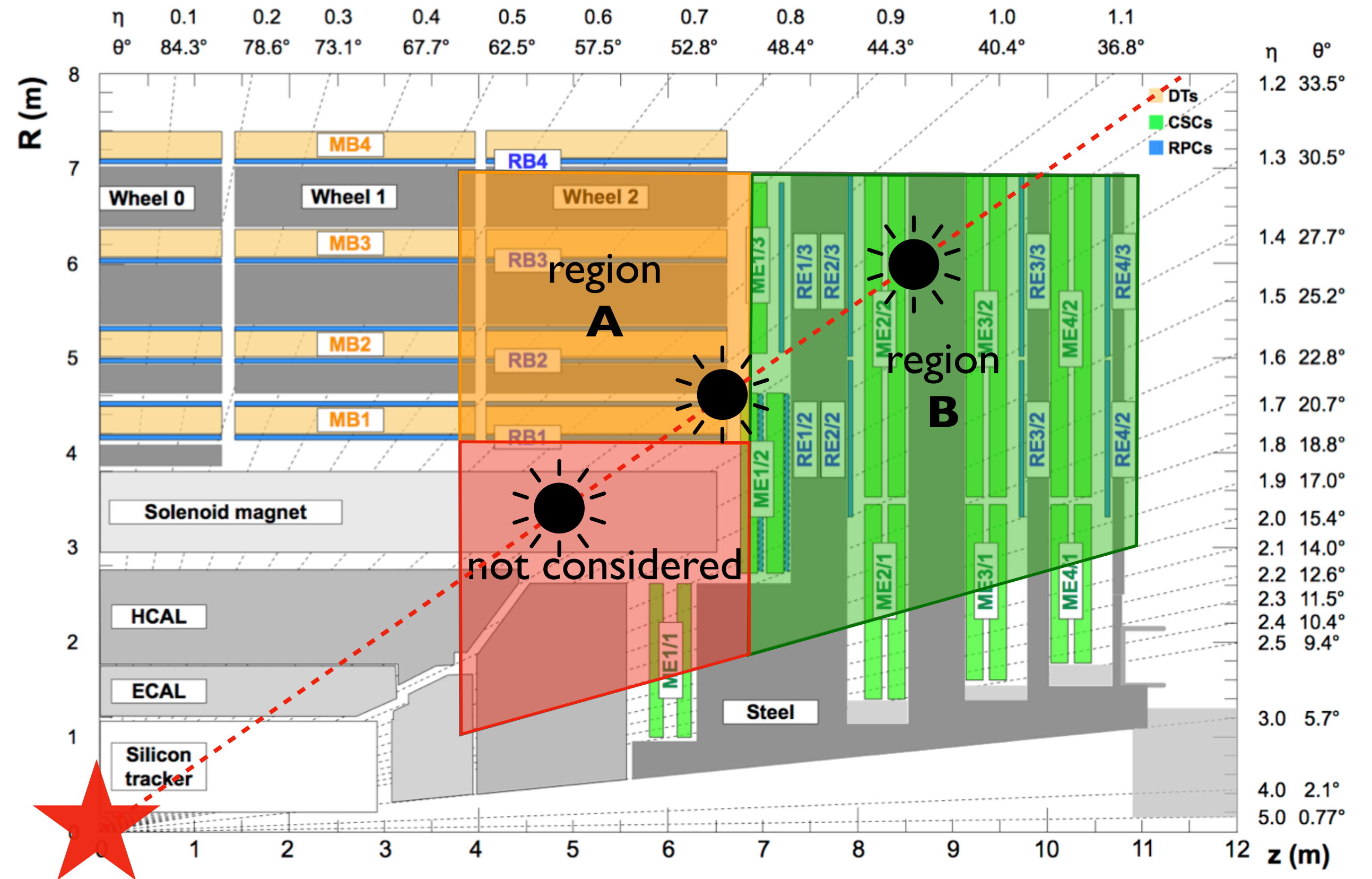
## Custom Decayer

given LLP lifetime and energy derive the probability to decay in each detector region

by using Pythia8 **generates a decay event in each detector region** and weights it with the probability to decay in that region

## Detector simulation

uses Delphes, with the publicly available module to take to account **cluster** and **cut-based ID** efficiencies for LLPs decays





# RECAST STRATEGY

## Event Generation

**MadGraph + Pythia8** generate the hard processes, shower, and hadronize them

**loose generator level cuts** (MET cut, isolation veto, ecc..) are applied to increase statistics

## Custom Decayer

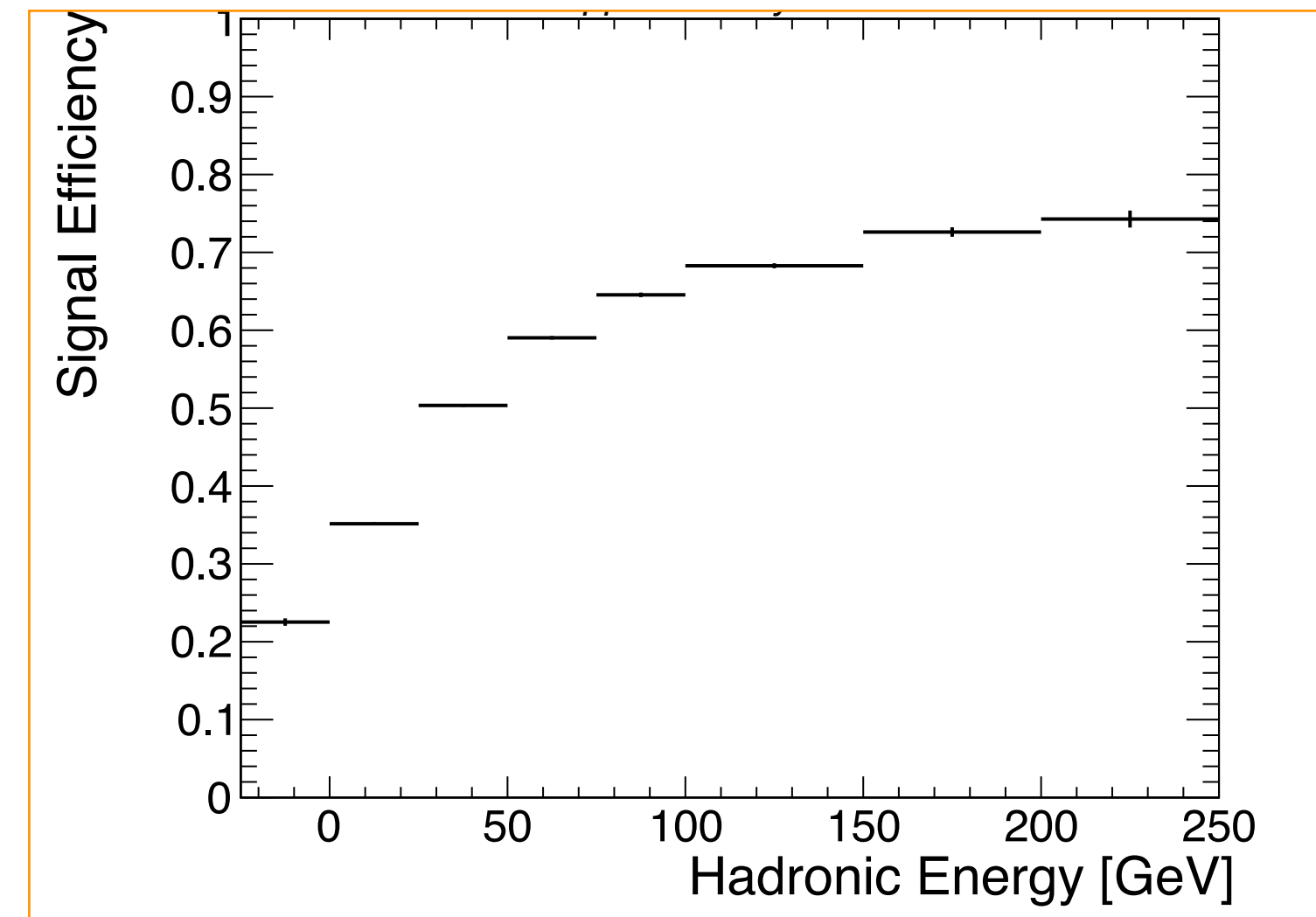
given LLP lifetime and energy derive the probability to decay in each detector region

by using Pythia8 **generates a decay event in each detector region** and weights it with the probability to decay in that region

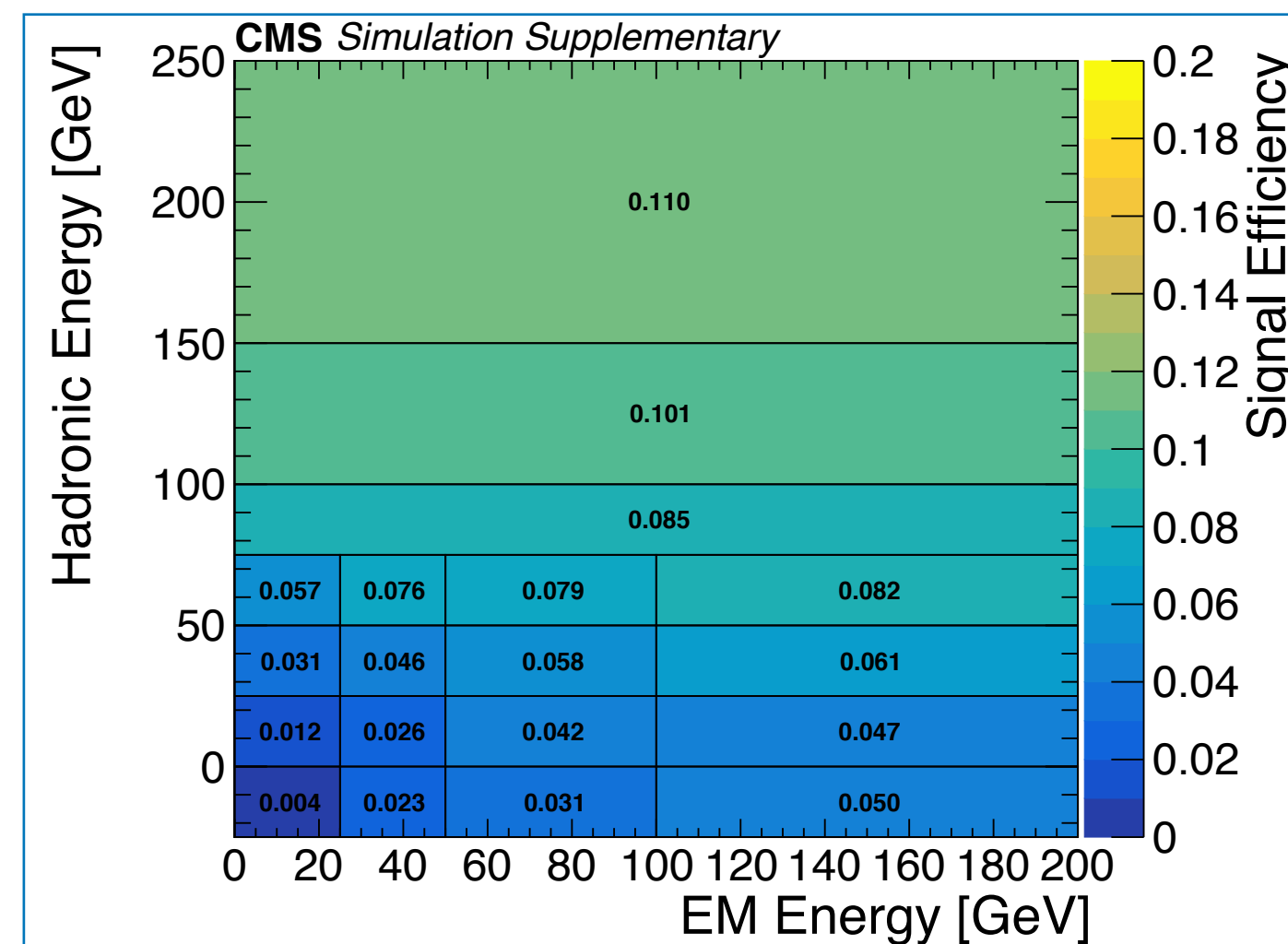
## Detector simulation

uses Delphes, with the publicly available module to take into account **cluster** and **cut-based ID** efficiencies for LLPs decays

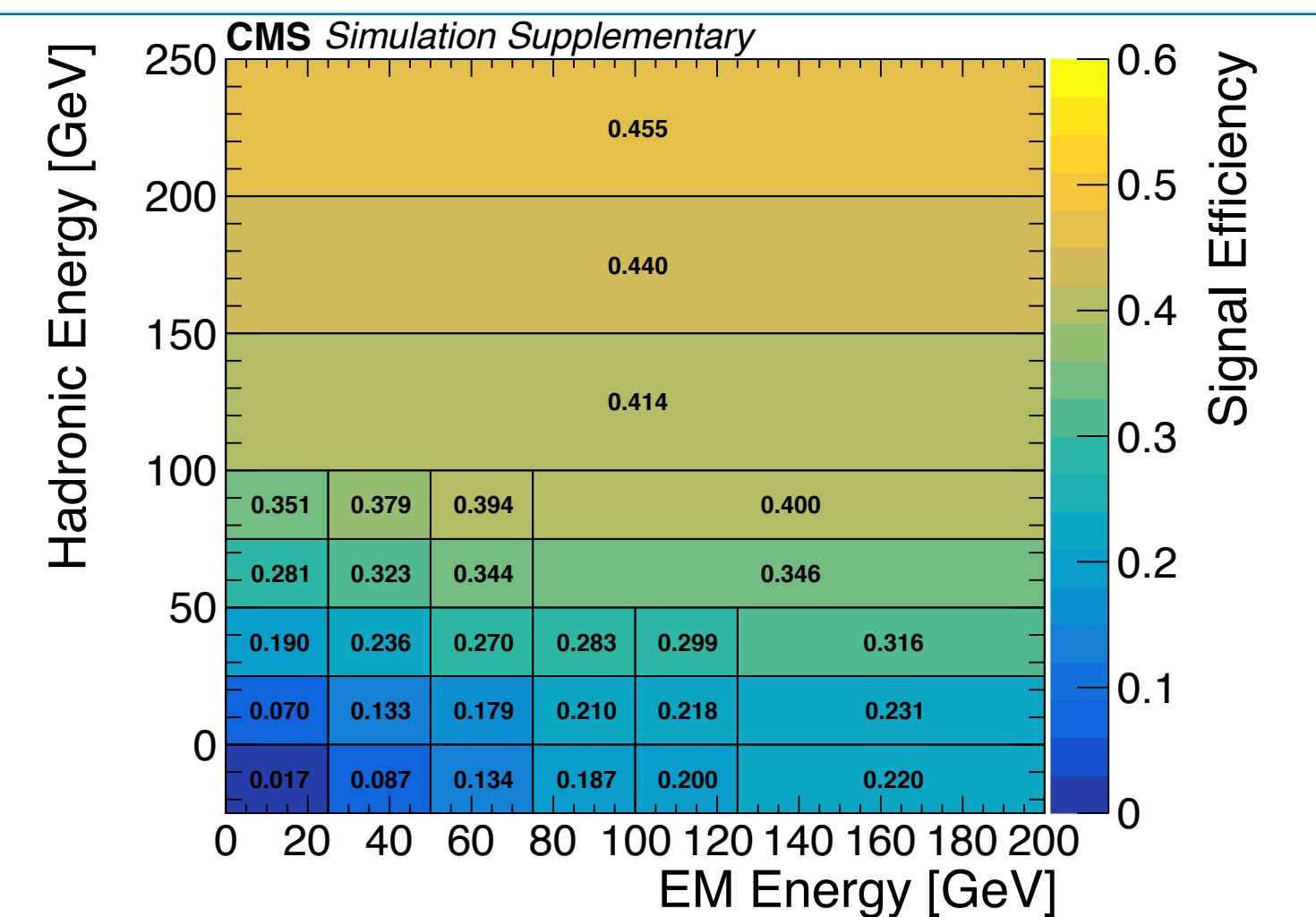
## cut-based ID for region B



## cluster efficiency region A



## cluster efficiency region B





# RESULTS FOR BENCHMARK MODELS



# LIGHT SCALAR

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \overbrace{\frac{1}{2} \partial_\mu \hat{S} \partial^\mu \hat{S} - \frac{\mu_S^2}{2} \hat{S}^2}^{\mathcal{L}_{DS}} - \overbrace{\left( A_{HS} \hat{S} + \lambda_{HS} \hat{S}^2 \right) \hat{H}^\dagger \hat{H}}^{\text{Higgs portal}}$$

- potential for S chosen such that S does not develop a vev
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \ll m_H$



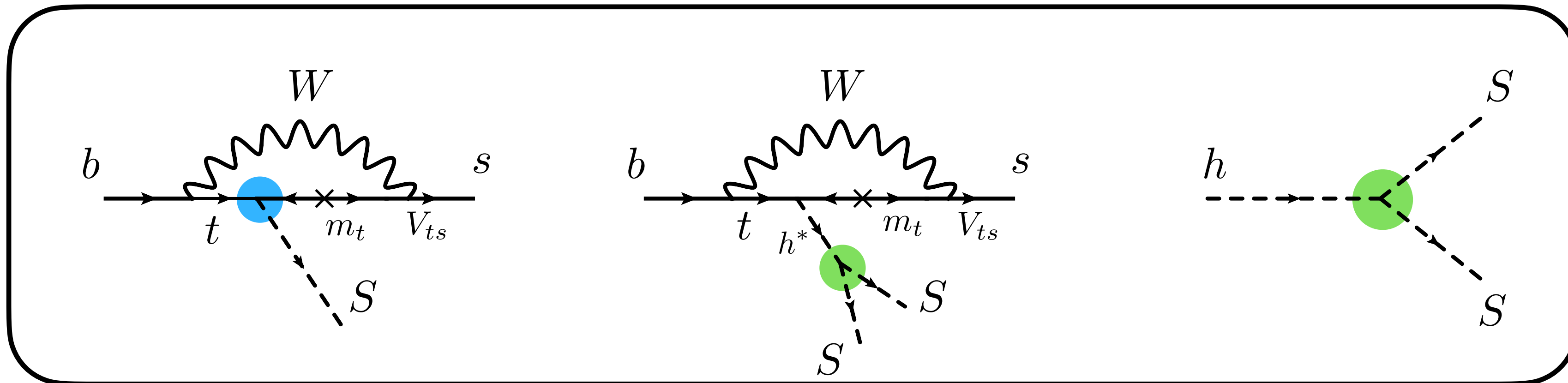
# LIGHT SCALAR

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \underbrace{\frac{1}{2} \partial_\mu \hat{S} \partial^\mu \hat{S} - \frac{\mu_S^2}{2} \hat{S}^2}_{\mathcal{L}_{DS}} - \underbrace{\left( A_{HS} \hat{S} + \lambda_{HS} \hat{S}^2 \right) \hat{H}^\dagger \hat{H}}_{\text{Higgs portal}}$$

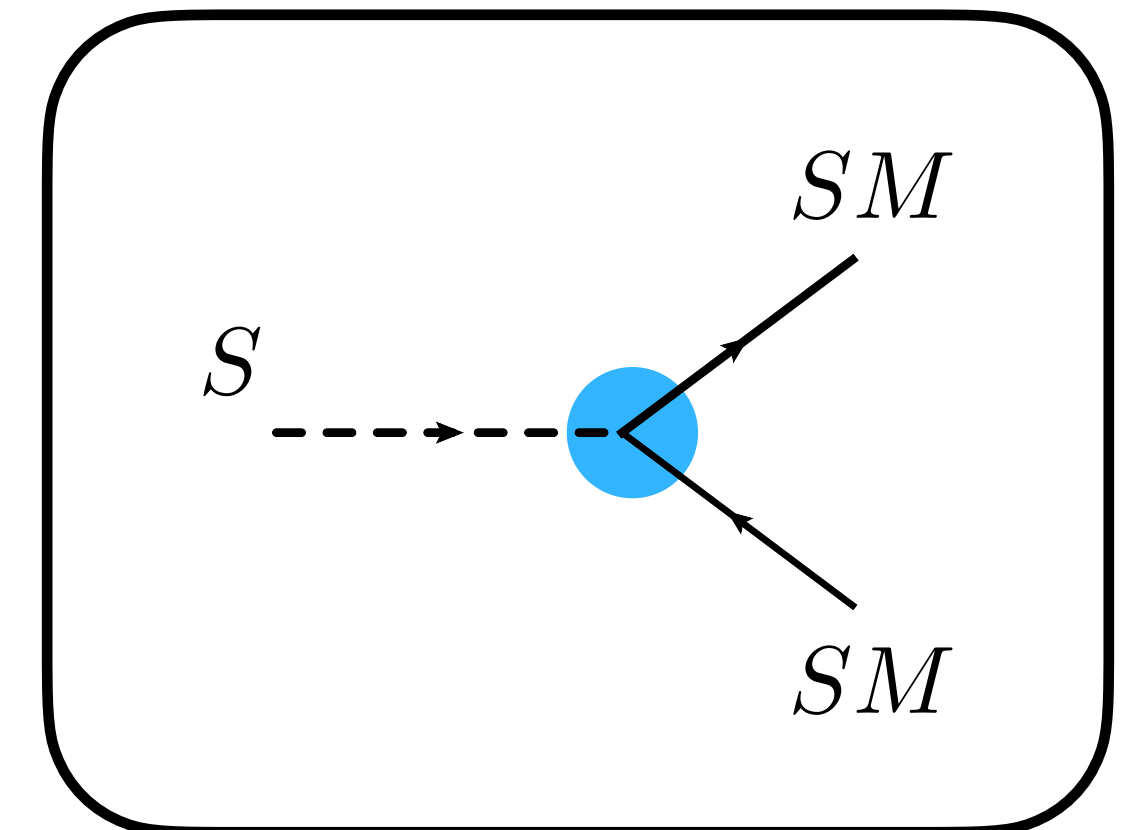
controls the  $\hat{H} - \hat{S}$  mixing     controls  $Br(H \rightarrow SS)$

- potential for S chosen such that S does not develop a vev
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \ll m_H$

## production



## decay





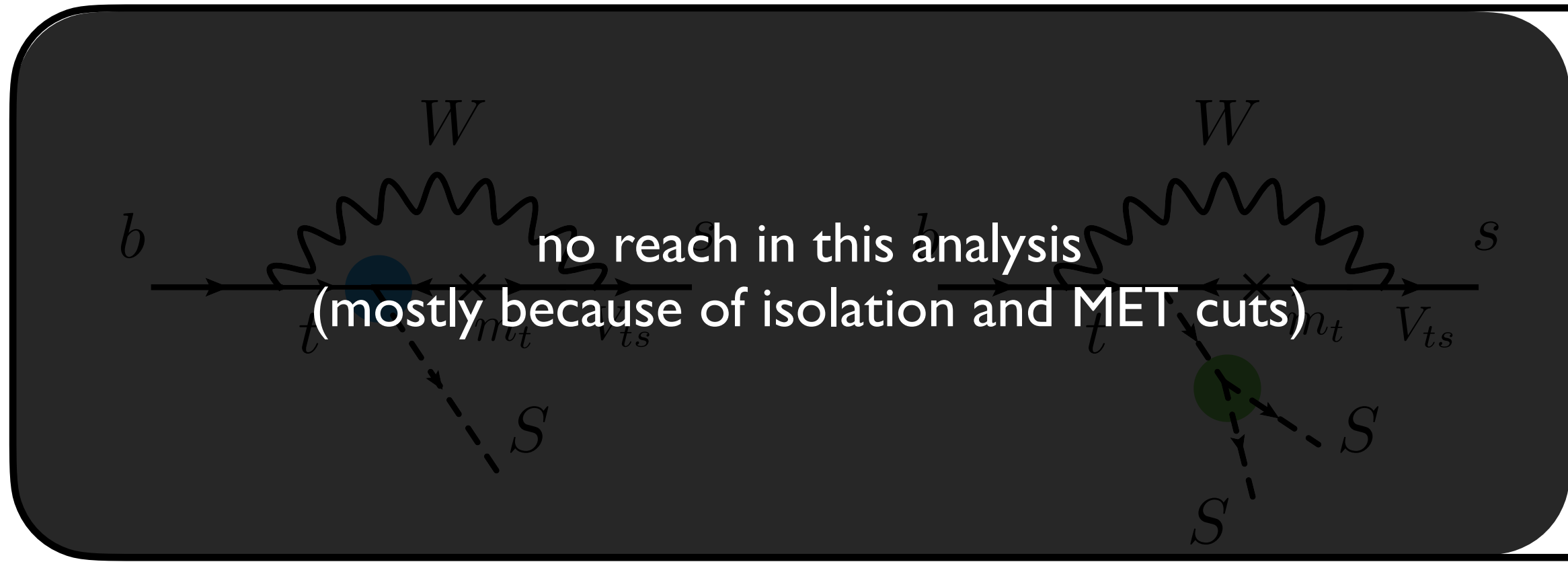
# LIGHT SCALAR

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \underbrace{\frac{1}{2} \partial_\mu \hat{S} \partial^\mu \hat{S} - \frac{\mu_S^2}{2} \hat{S}^2}_{\mathcal{L}_{DS}} - \underbrace{\left( A_{HS} \hat{S} + \lambda_{HS} \hat{S}^2 \right) \hat{H}^\dagger \hat{H}}_{\text{Higgs portal}}$$

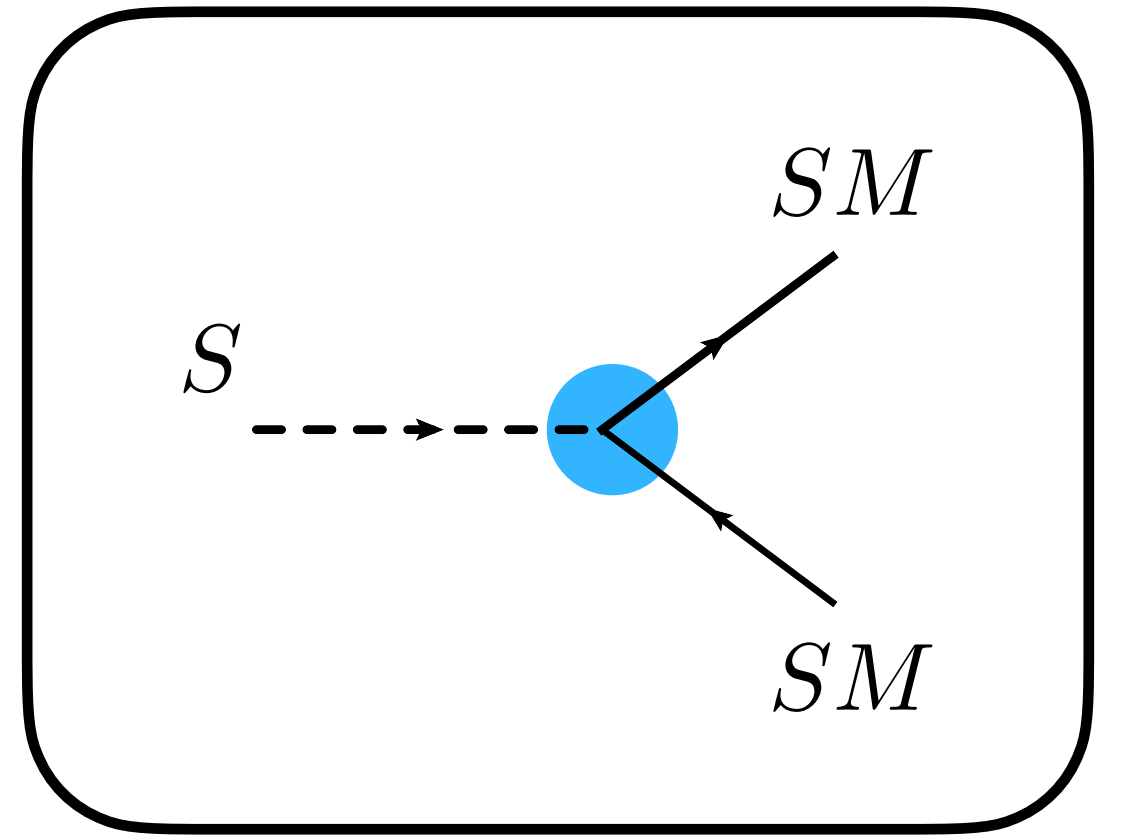
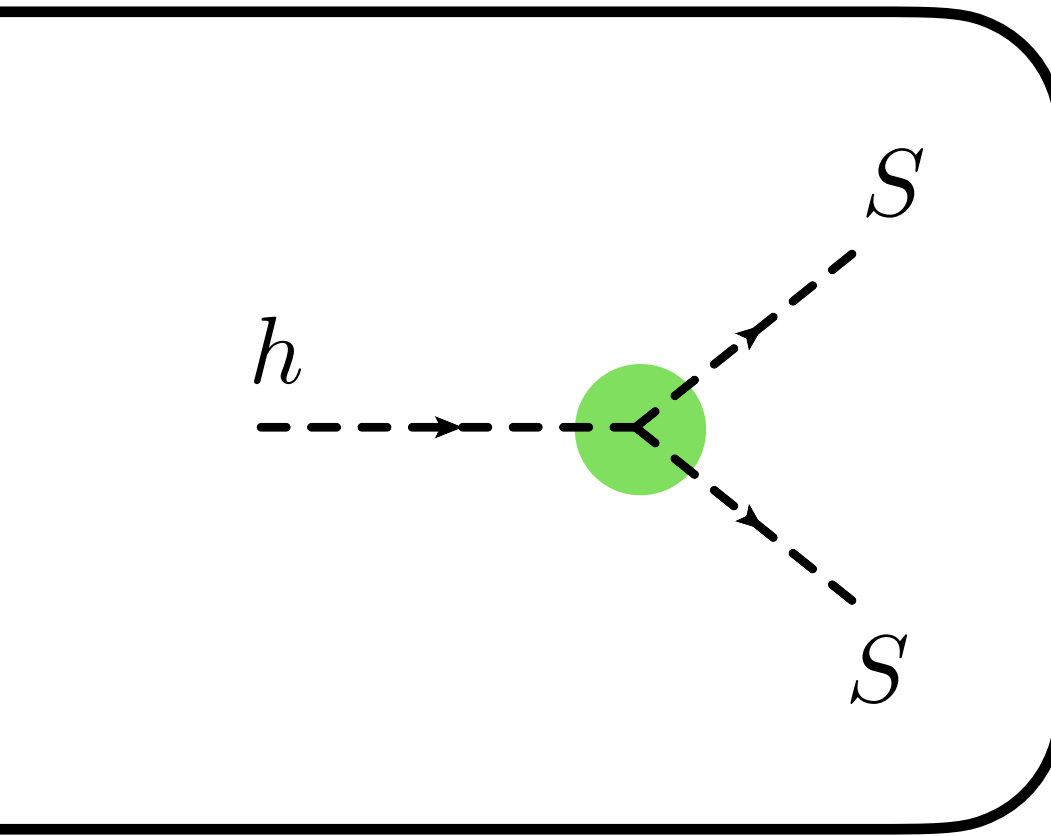
controls the  $\hat{H} - \hat{S}$  mixing     controls  $Br(H \rightarrow SS)$

- potential for S chosen such that S does not develop a vev
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \ll m_H$

## production



## decay

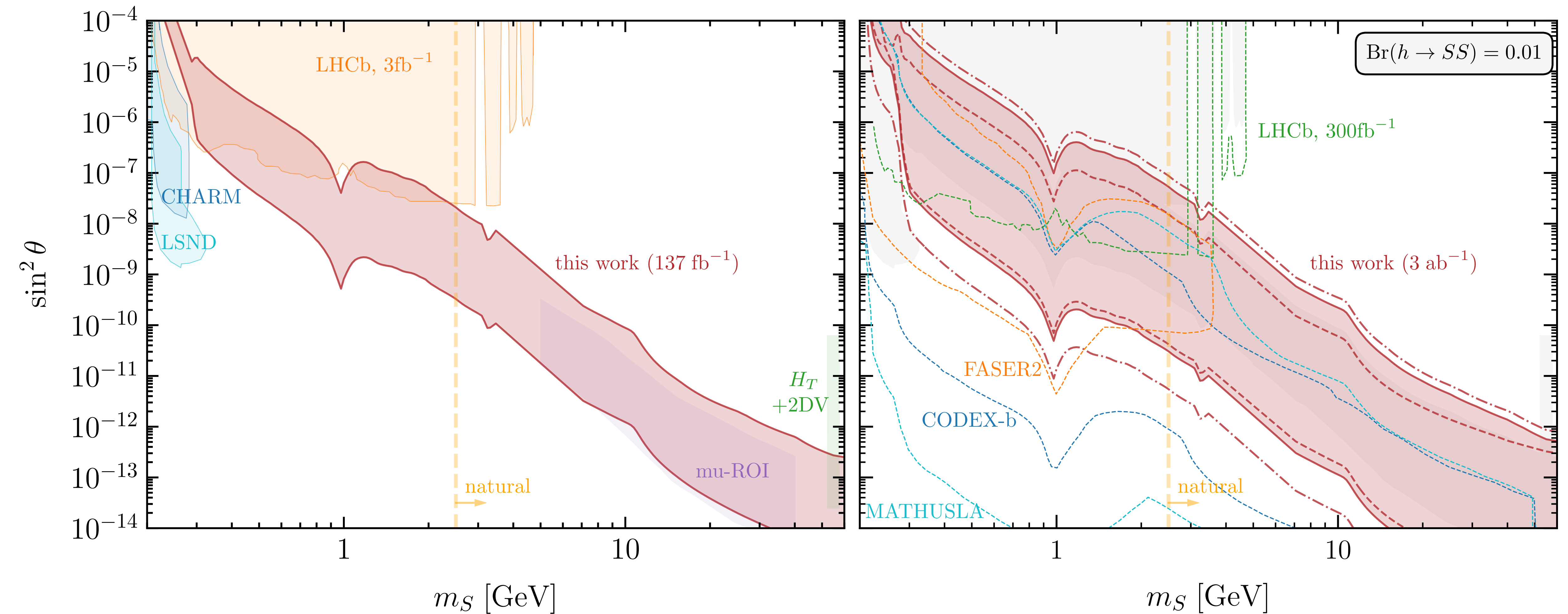


**production and decay** channels are **decoupled**



# LIGHT SCALAR

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \frac{1}{2} \partial_\mu \hat{S} \partial^\mu \hat{S} - \frac{\mu_S^2}{2} \hat{S}^2 - \left( A_{HS} \hat{S} + \lambda_{HS} \hat{S}^2 \right) \hat{H}^\dagger \hat{H}$$





# INELASTIC DM

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + ie_D \hat{X}^\mu \bar{\chi}_1 \gamma^\mu \chi_2 - \frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}$$

the mass hierarchy of the model  
is

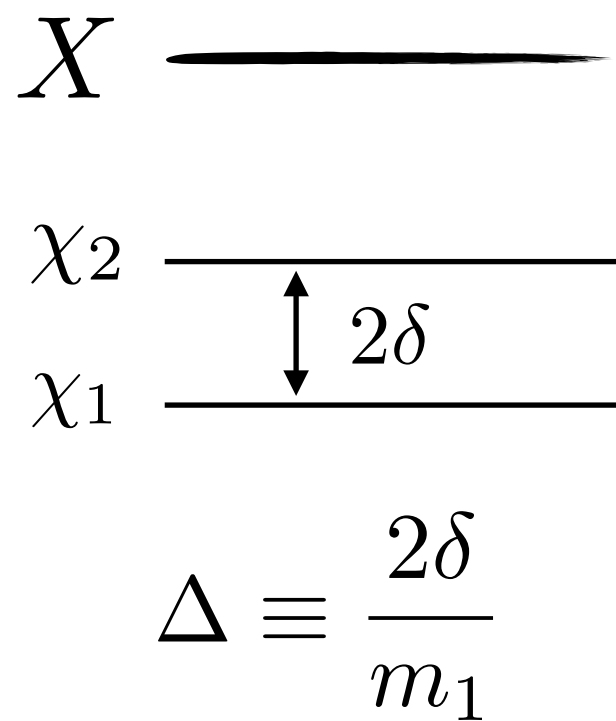
$$\begin{array}{c} X \text{ —————} \\ \chi_2 \text{ —————} \\ \quad \updownarrow 2\delta \\ \chi_1 \text{ —————} \\ \Delta \equiv \frac{2\delta}{m_1} \end{array}$$



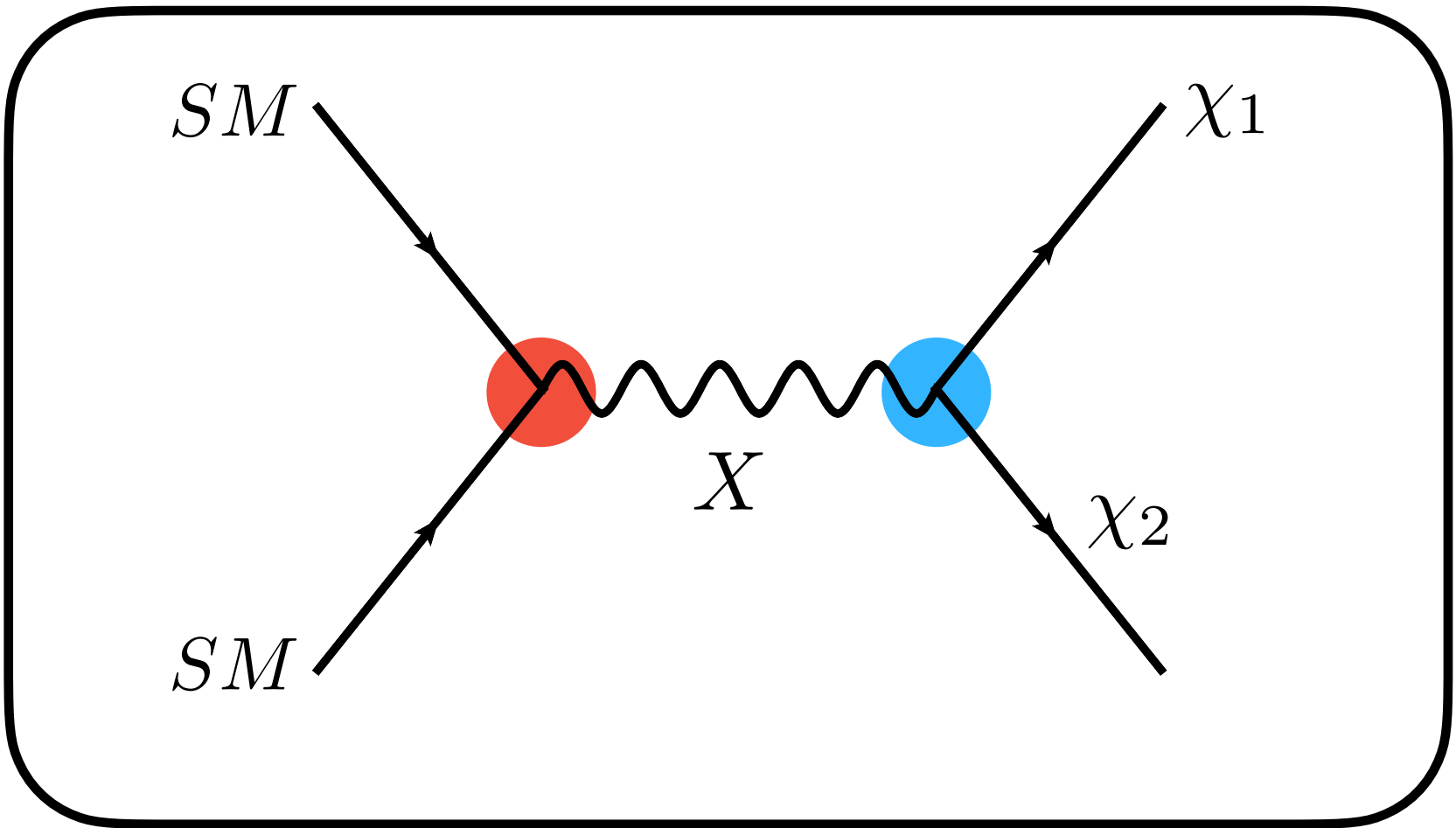
# INELASTIC DM

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + ie_D \hat{X}^\mu \bar{\chi}_1 \gamma^\mu \chi_2 - \frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}$$

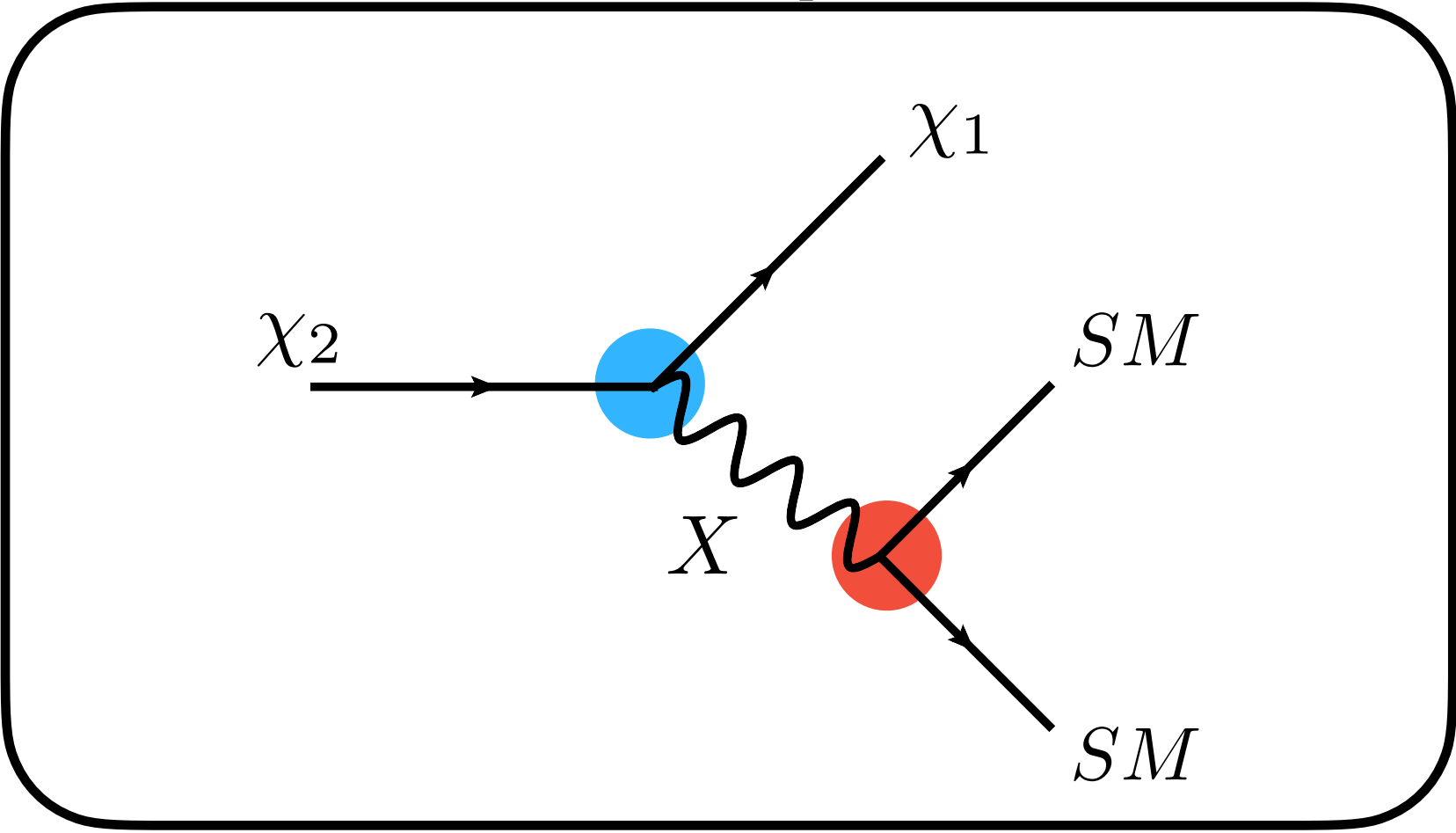
the mass hierarchy of the model is



**production**



**decay**

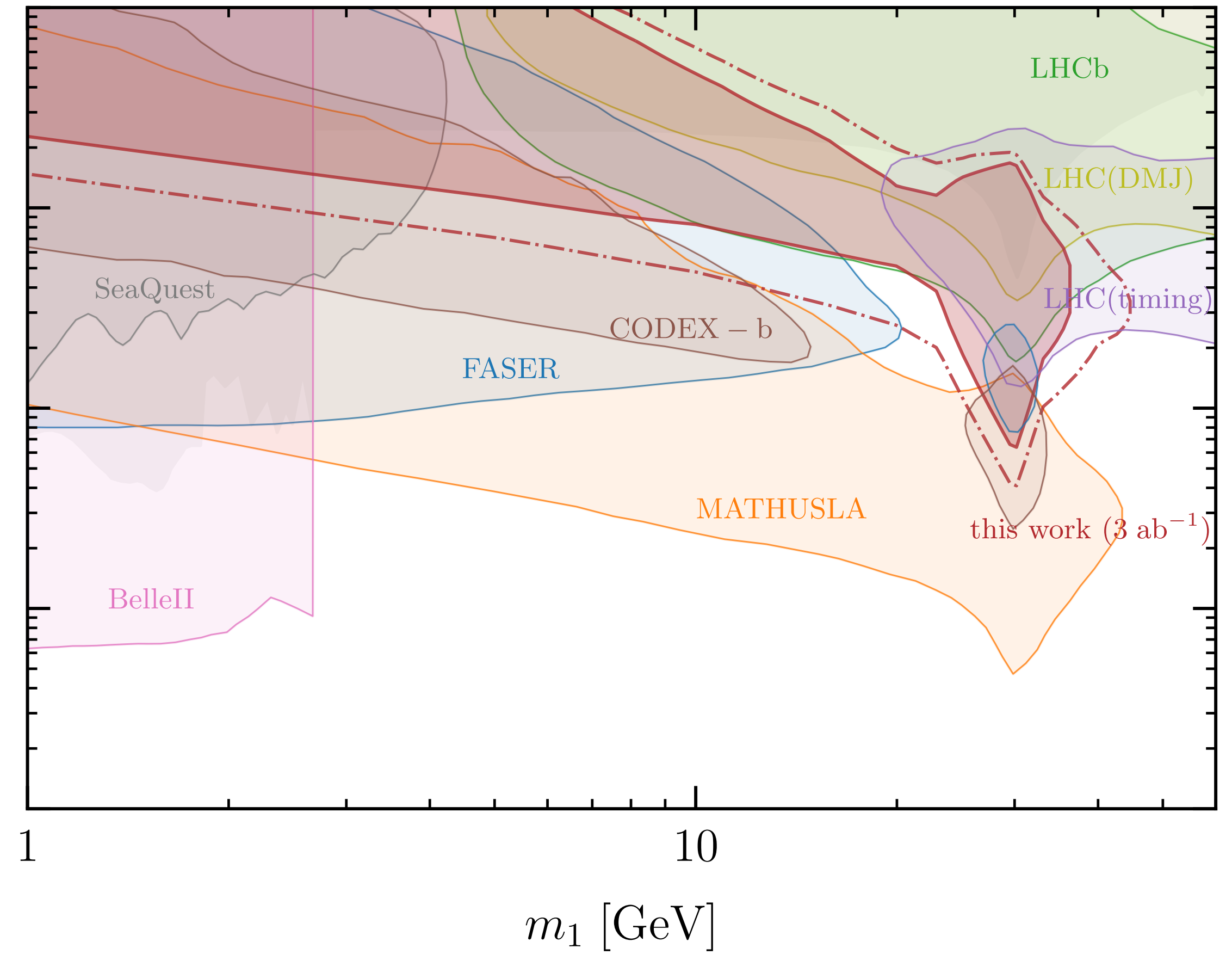
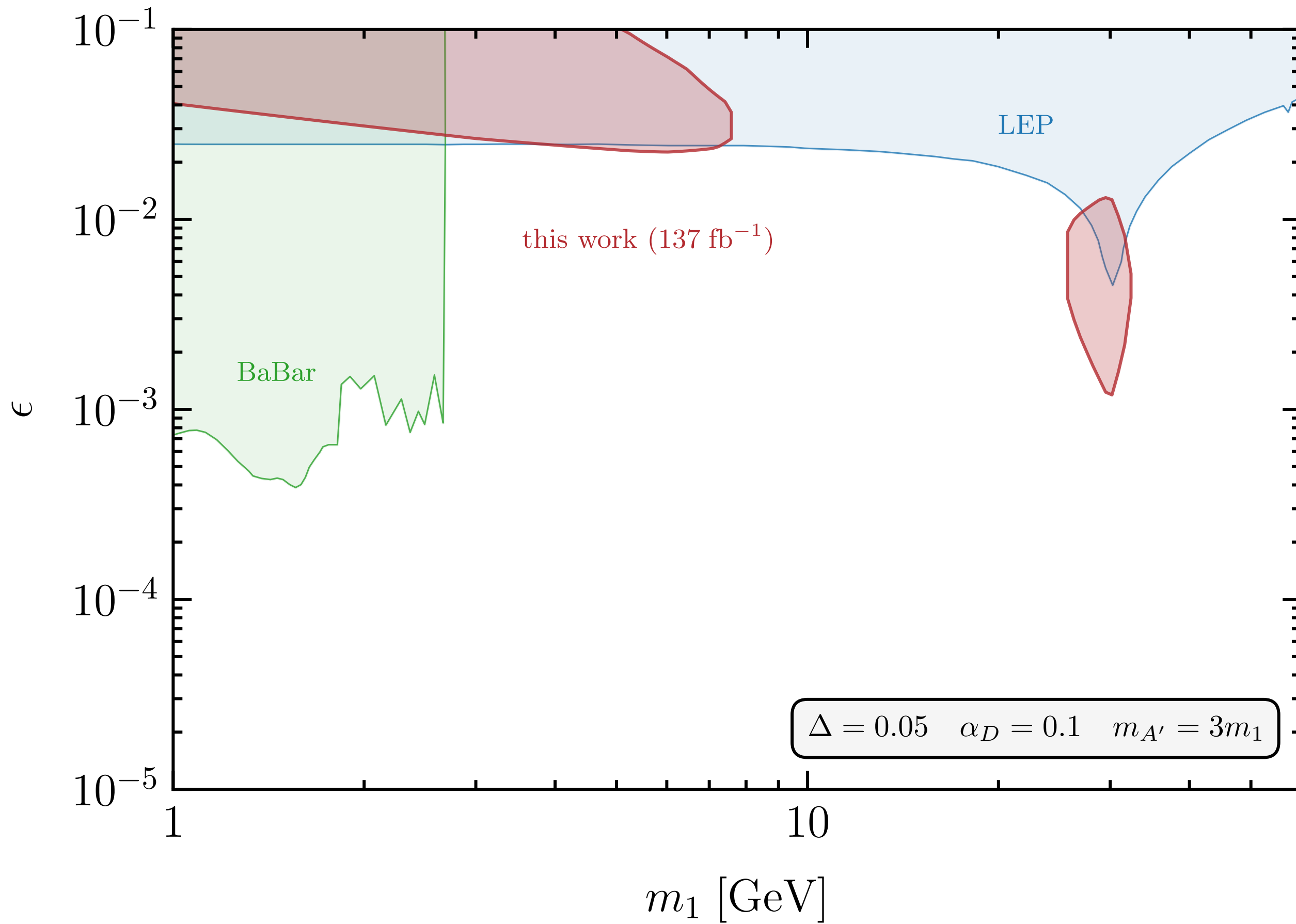


$E_{\text{LLP}}$  and MET are decoupled



# INELASTIC DM

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + ie_D \hat{X}^\mu \bar{\chi}_1 \gamma^\mu \chi_2 - \frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}$$

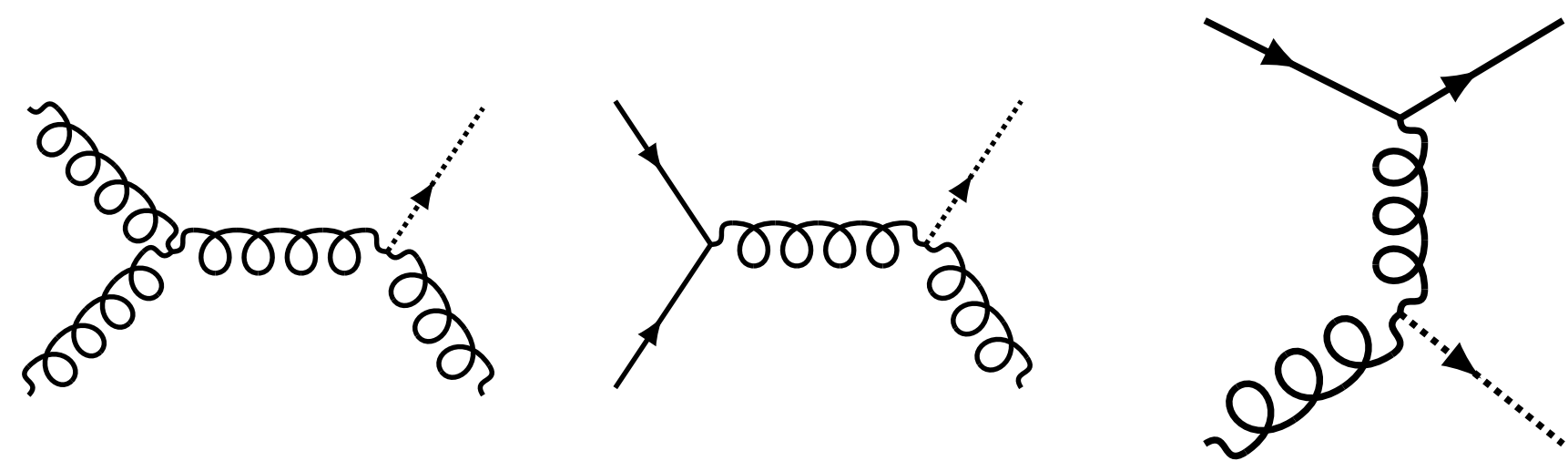




# ALP

$$\mathcal{L} = \frac{1}{2} (\partial_\mu a)^2 + \frac{a}{4\pi f_a} \left( \alpha_s c_{GG} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} + \alpha_2 c_{WW} W_{\mu\nu}^a \tilde{W}^{a,\mu\nu} + \alpha_1 c_{BB} B_{\mu\nu} \tilde{B}^{\mu\nu} \right) + \dots$$

gluon-coupled ALP  
 $(c_{GG} \neq 0, c_{BB} = c_{WW} = 0)$



$$a \rightarrow \gamma\gamma$$

$$a \rightarrow \eta\pi\pi$$

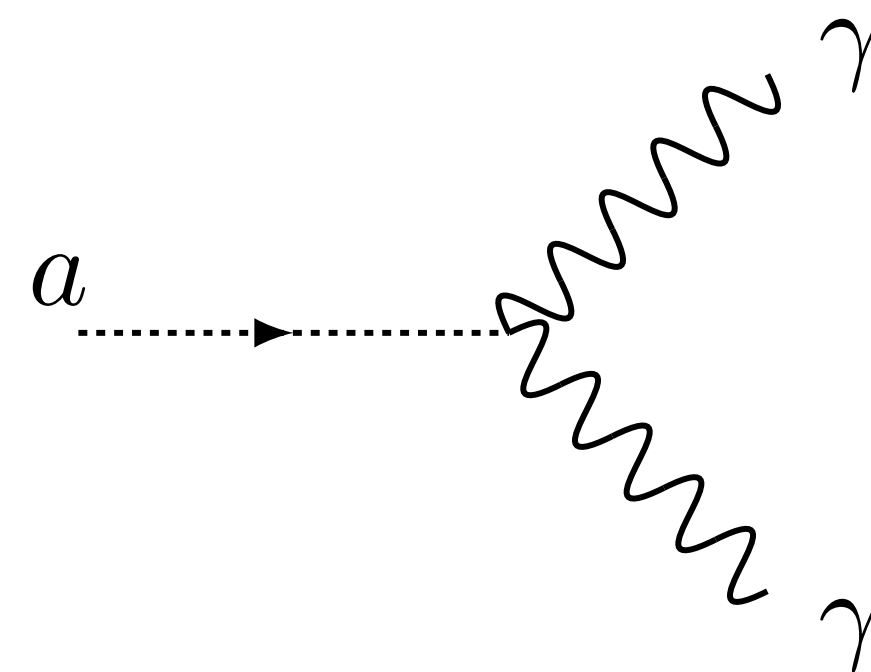
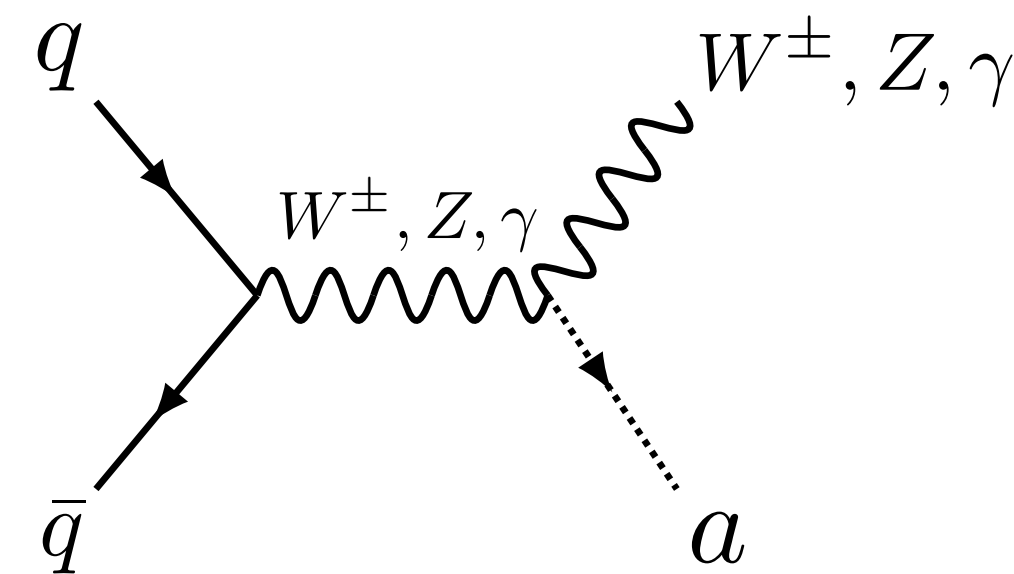
$$a \rightarrow 3\pi$$

$$a \rightarrow \pi\pi\gamma$$

main production channels

main decay channels  
for current reach

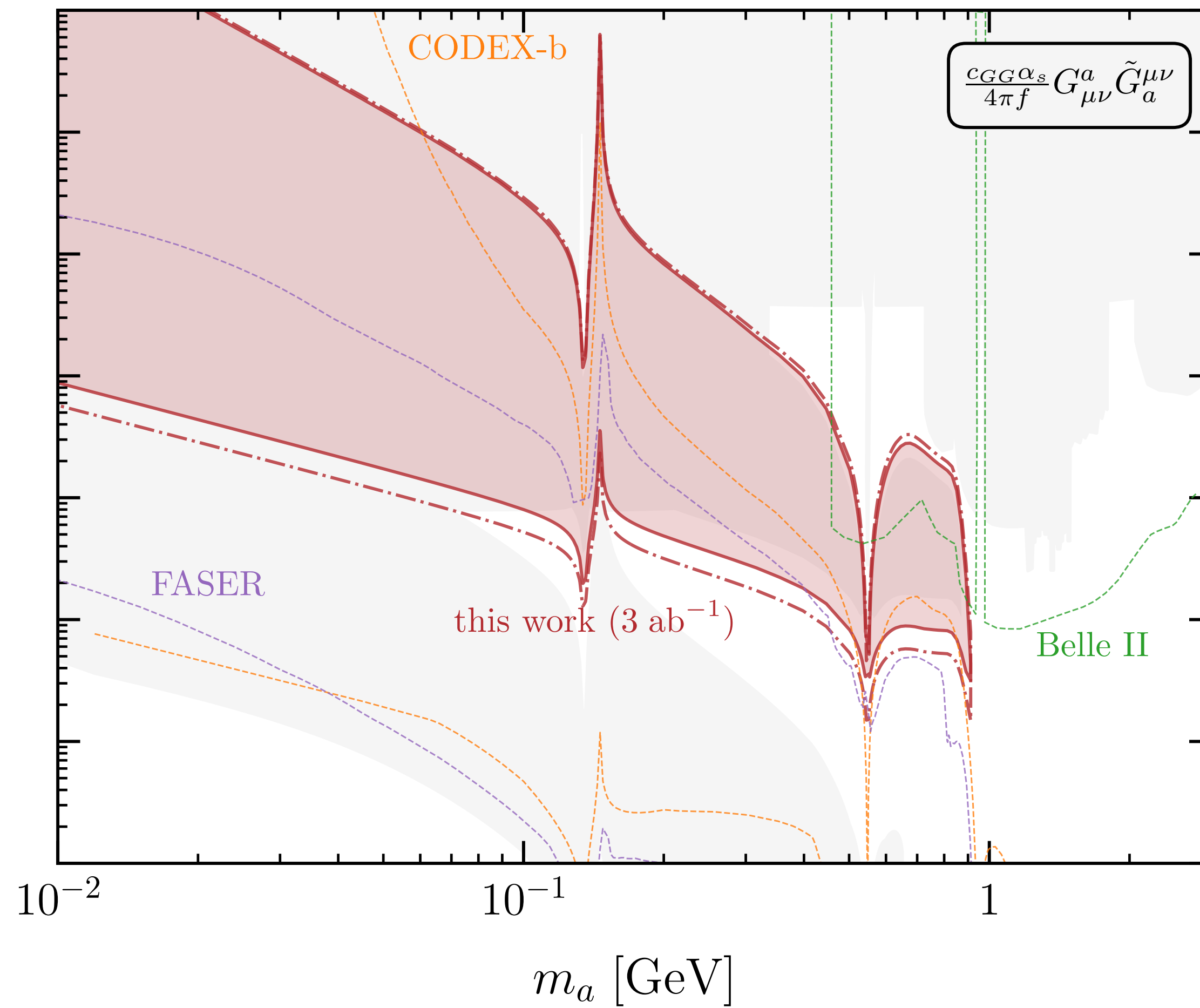
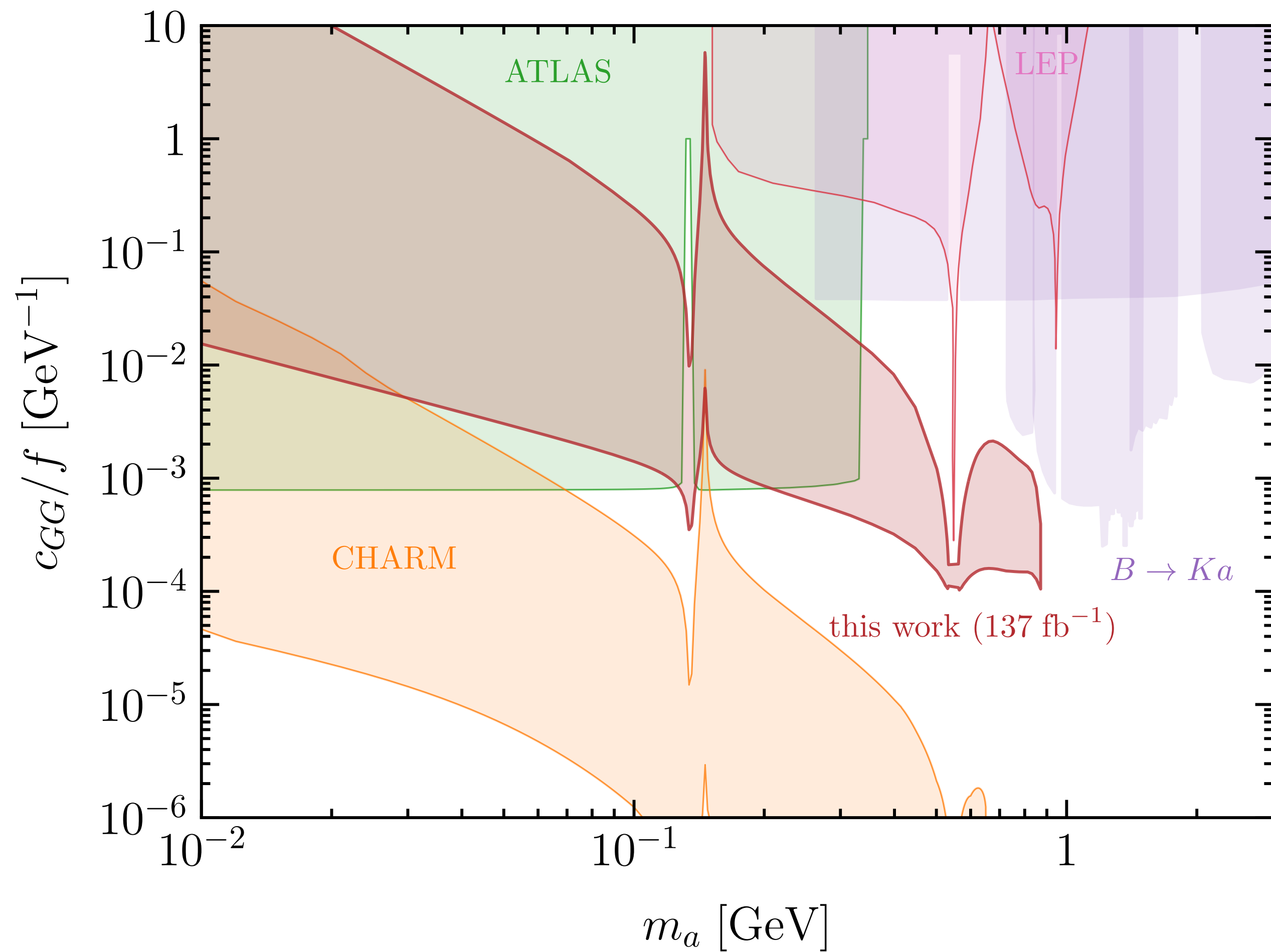
photon-coupled ALP  
 $(c_{GG} \neq 0, c_{BB} = c_{WW} = 0)$





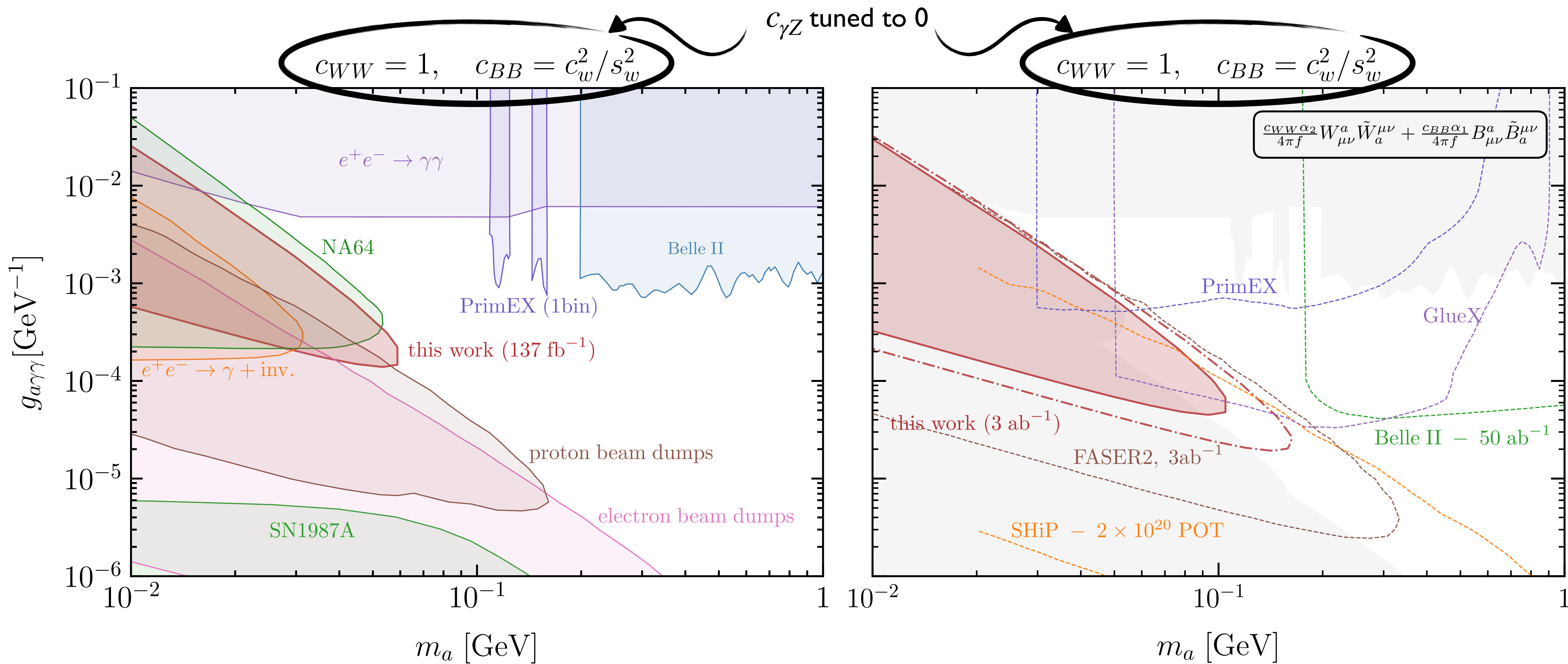
# ALP

$$\mathcal{L} = \frac{1}{2} (\partial_\mu a)^2 + \frac{a}{4\pi f_a} \left( \alpha_s c_{GG} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} + \alpha_2 c_{WW} W_{\mu\nu}^a \tilde{W}^{a,\mu\nu} + \alpha_1 c_{BB} B_{\mu\nu} \tilde{B}^{\mu\nu} \right) + \dots$$





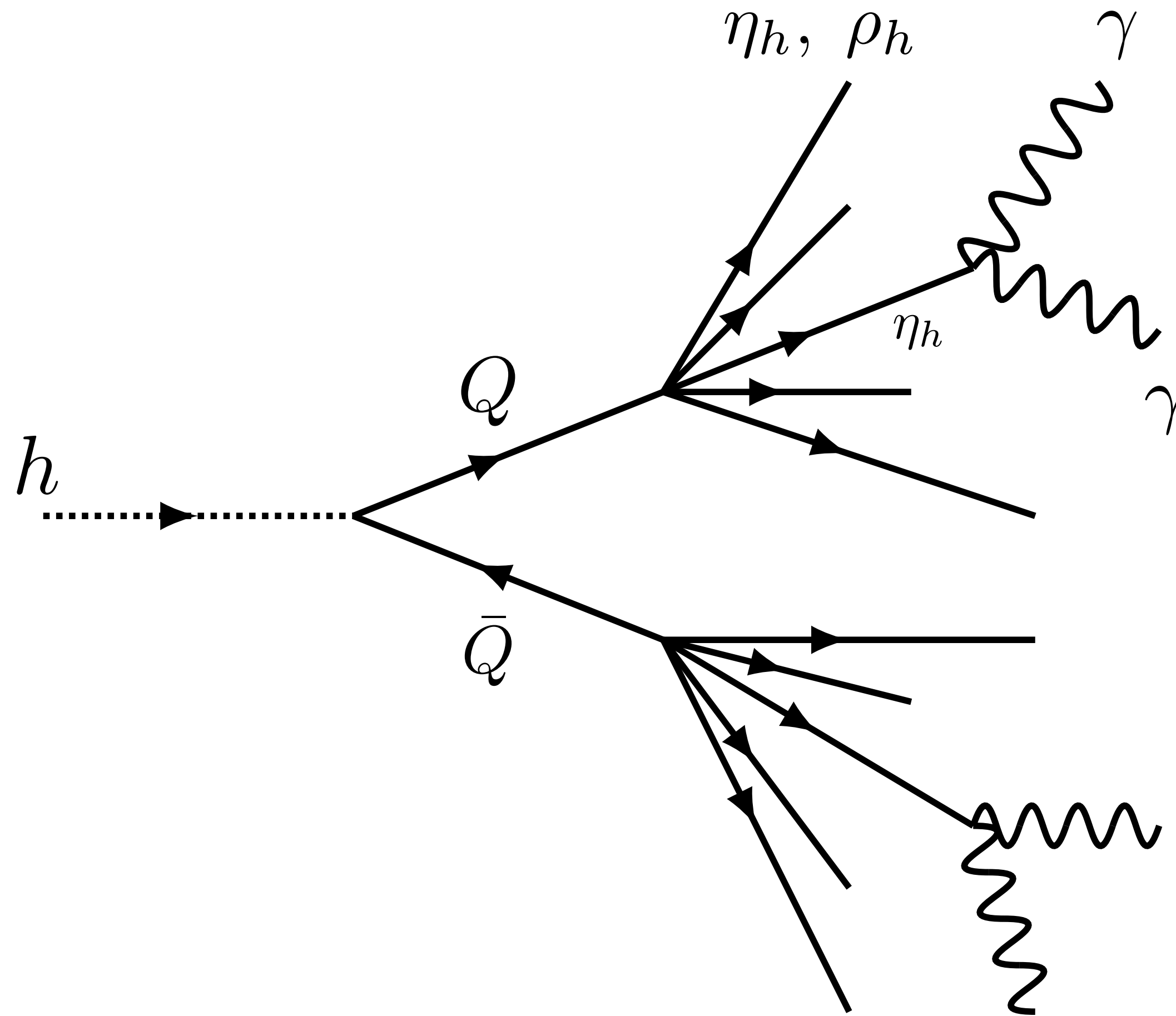
$$\mathcal{L} = \frac{1}{2} (\partial_\mu a)^2 + \frac{a}{4\pi f_a} \left( \alpha_s c_{GG} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} + \alpha_2 c_{WW} W_{\mu\nu}^a \tilde{W}^{a,\mu\nu} + \alpha_1 c_{BB} B_{\mu\nu} \tilde{B}^{\mu\nu} \right) + \dots$$





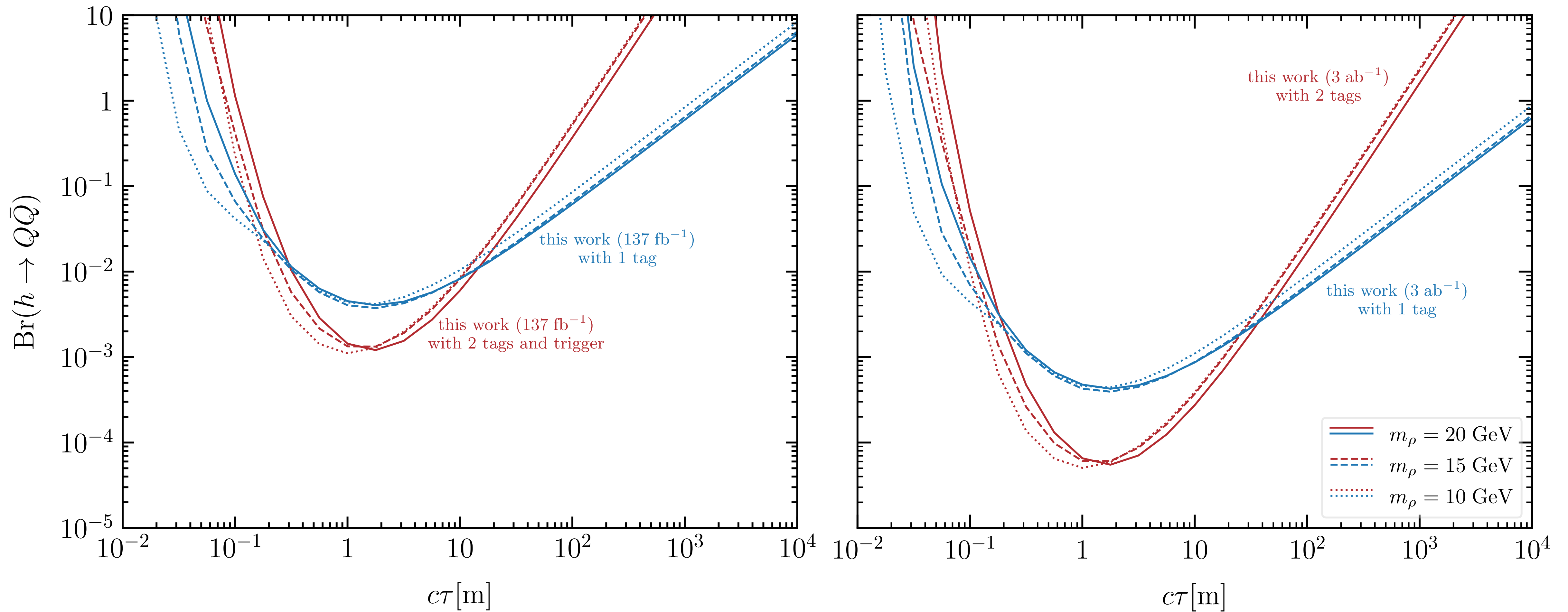
# HIDDEN VALLEY

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \lambda h \bar{Q} Q$$



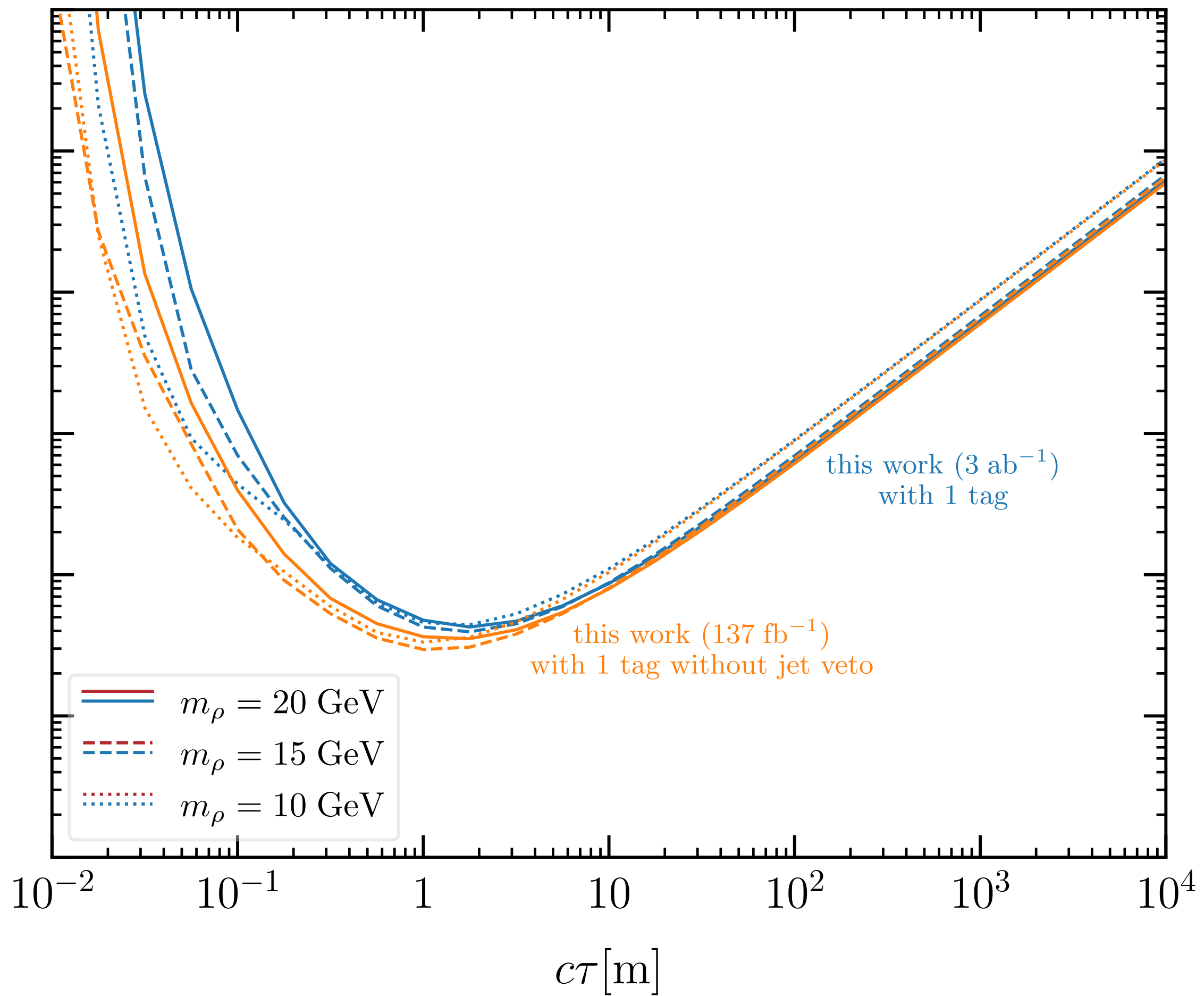
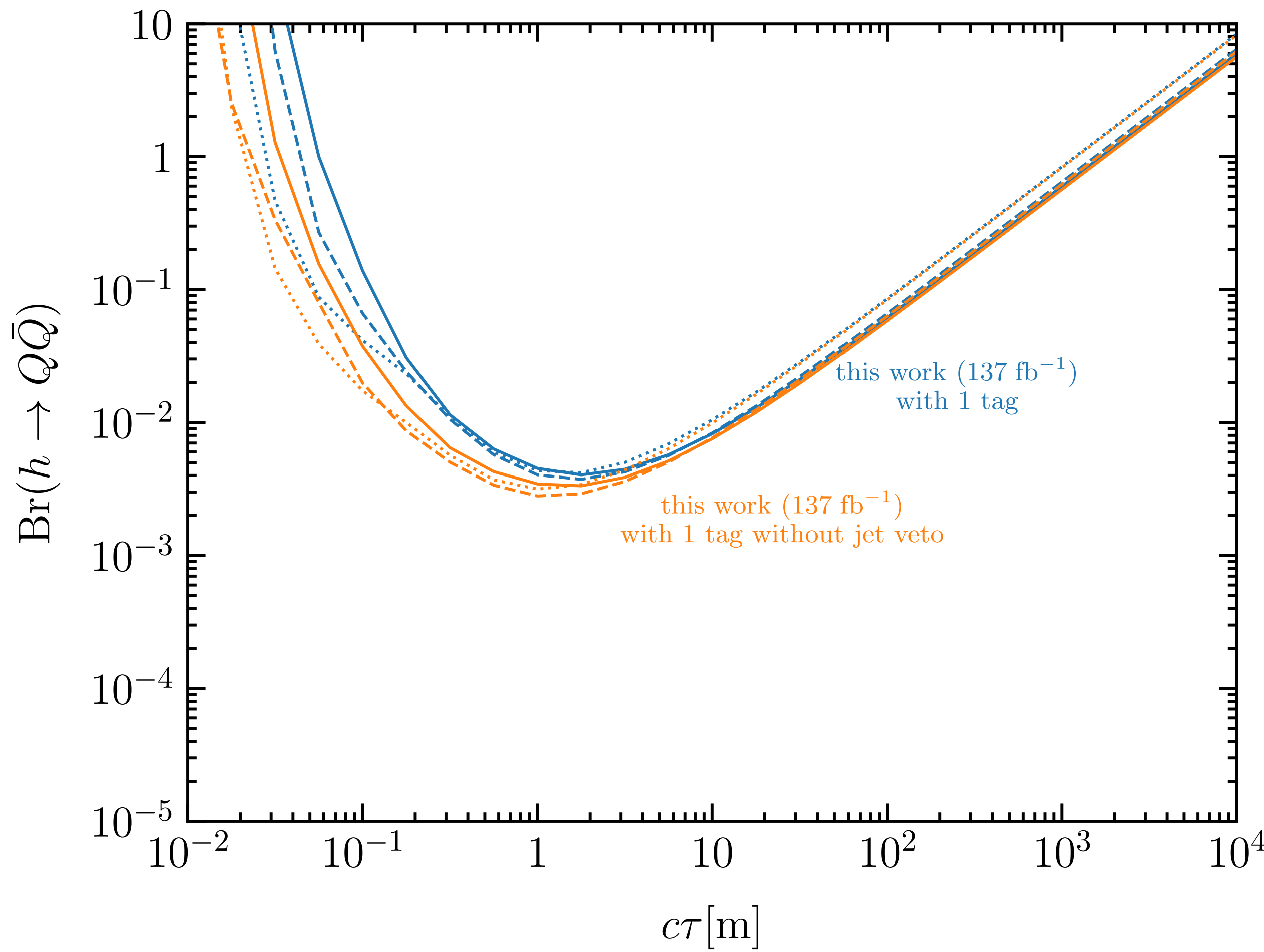
- we treat the  $\eta_h$  lifetime as a free parameter
- we take  $\text{Br}(\eta_h \rightarrow \gamma\gamma) = 1$  (hard to probe with other searches)
- we assume  $\rho_h$  to decay into  $\eta_h\eta_h$  to maximize self-veto effects
- we implement the dark shower with Pythia treating  $\eta_h$  as a pion

# HIDDEN VALLEY





# HIDDEN VALLEY



# TAKEAWAYS

- for most of the benchmark models, *the analysis covers previously unconstrained regions of the parameter space*
- ***MET cut boon and bane:***
  - allows to probe lower lifetimes
  - reduces cross-section
  - sensitivity could be improved by using model-specific techniques to suppress the backgrounds
- Isolation cuts reduce sensitivity for models where multiple LLPs are produced in the same direction (e.g. hidden valley)



# BACKUP SLIDES

# DARK PHOTON

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \underbrace{\mathcal{L}_{DS}}_{\text{Higgs portal}} - \lambda_{HS} \hat{S}^2 \hat{H}^\dagger \hat{H} - \underbrace{\frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}}_{\text{vector portal}}$$

$$\mathcal{L}_{DS} = -\frac{1}{4} \hat{X}_{\mu\nu} \hat{X}^{\mu\nu} + \frac{1}{2} m_{X,0}^2 \hat{X}_\mu \hat{X}^\mu + \mu_S^2 \hat{S}^2 - \lambda_S \hat{S}^4 + |(\partial_\mu + ig_D \hat{X}_\mu) \hat{S}|^2$$

- dark scalar breaks U(1) and gives a mass to the dark photon
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \gg m_H$



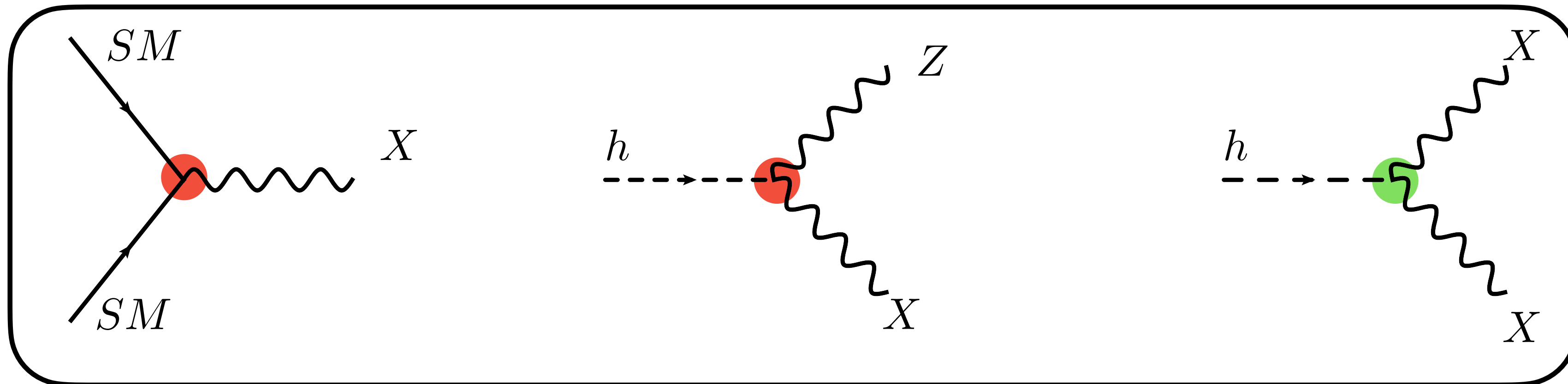
# DARK PHOTON

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \mathcal{L}_{DS} - \underbrace{\lambda_{HS} \hat{S}^2 \hat{H}^\dagger \hat{H}}_{\text{Higgs portal}} - \underbrace{\frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}}_{\text{vector portal}}$$

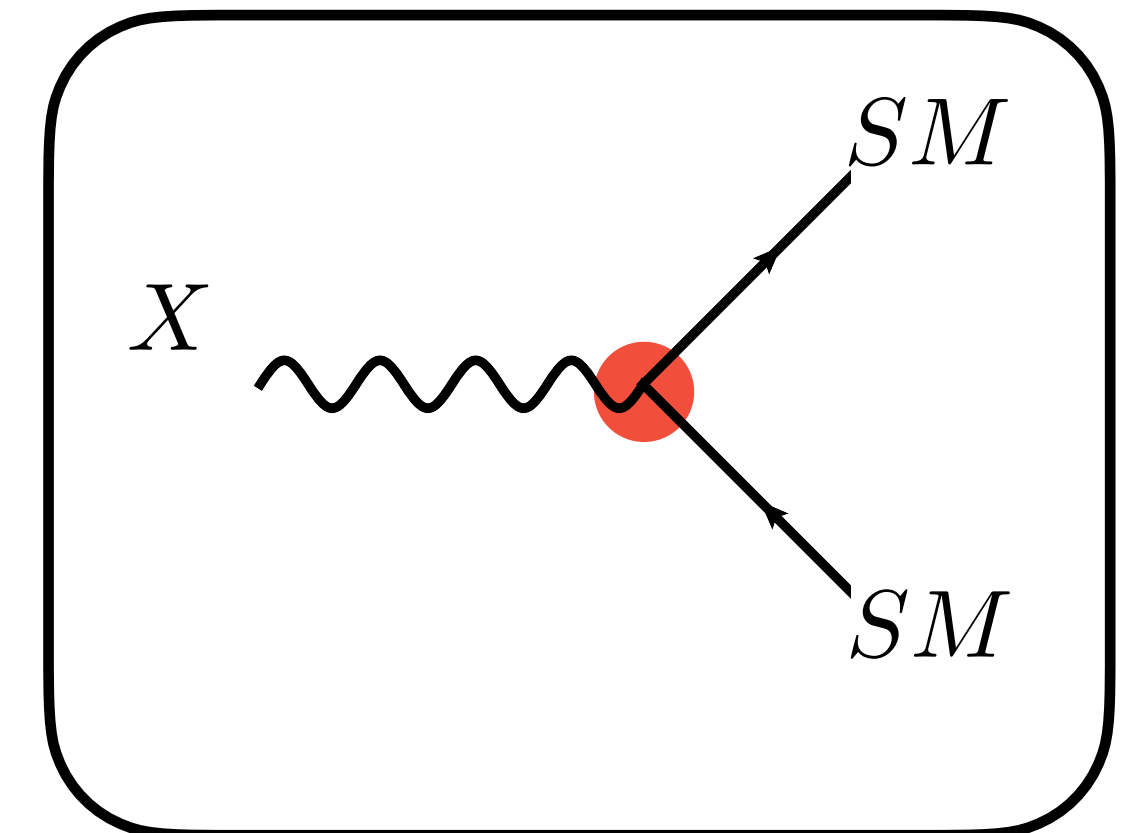
$$\mathcal{L}_{DS} = -\frac{1}{4} \hat{X}_{\mu\nu} \hat{X}^{\mu\nu} + \frac{1}{2} m_{X,0}^2 \hat{X}_\mu \hat{X}^\mu + \mu_S^2 \hat{S}^2 - \lambda_S \hat{S}^4 + |(\partial_\mu + ig_D \hat{X}_\mu) \hat{S}|^2$$

- dark scalar breaks U(1) and gives a mass to the dark photon
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \gg m_H$

## production



## decay



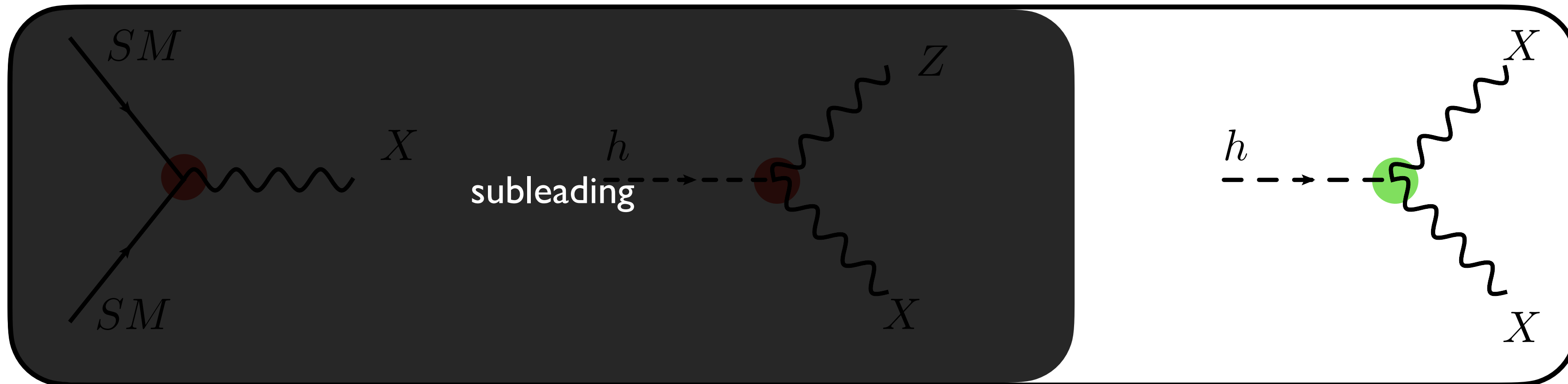
# DARK PHOTON

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \mathcal{L}_{DS} - \underbrace{\lambda_{HS} \hat{S}^2 \hat{H}^\dagger \hat{H}}_{\text{Higgs portal}} - \underbrace{\frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}}_{\text{vector portal}}$$

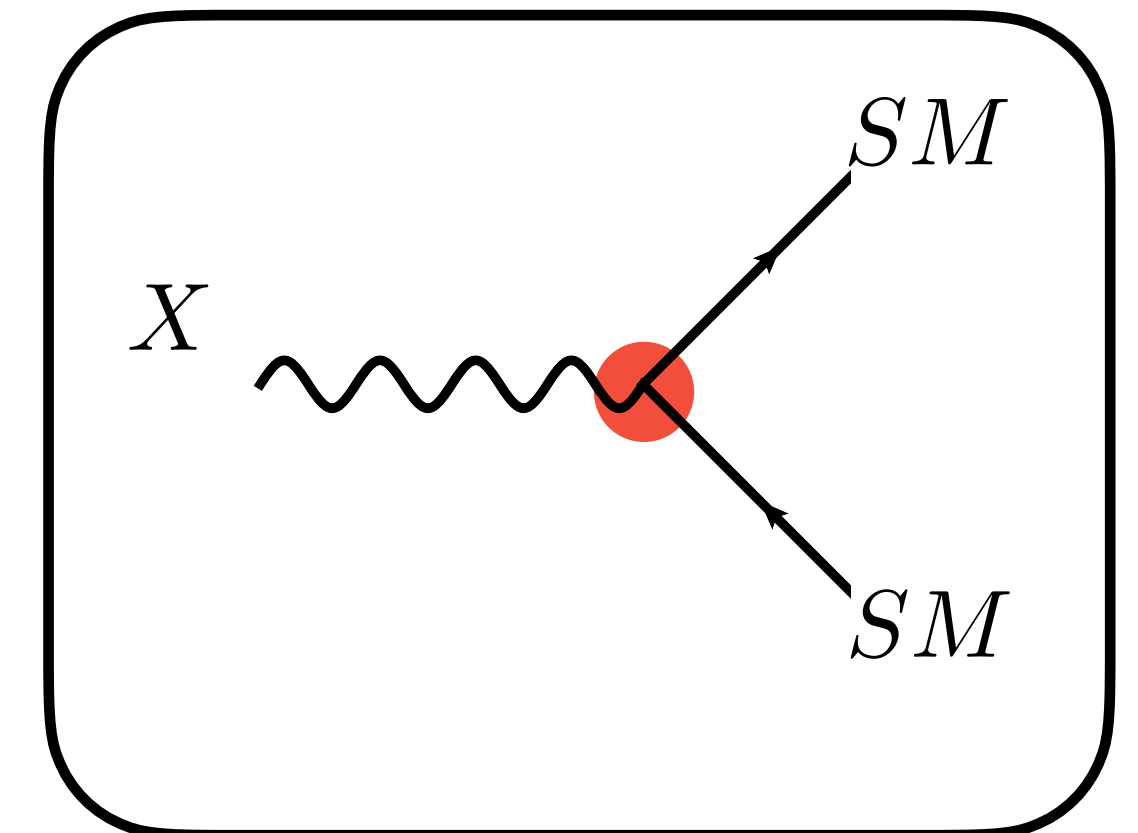
$$\mathcal{L}_{DS} = -\frac{1}{4} \hat{X}_{\mu\nu} \hat{X}^{\mu\nu} + \frac{1}{2} m_{X,0}^2 \hat{X}_\mu \hat{X}^\mu + \mu_S^2 \hat{S}^2 - \lambda_S \hat{S}^4 + |(\partial_\mu + ig_D \hat{X}_\mu) \hat{S}|^2$$

- dark scalar breaks U(1) and gives a mass to the dark photon
- $m_S \sim \mu_S^2 + \lambda_{HS} v^2 \gg m_H$

**production**



**decay**

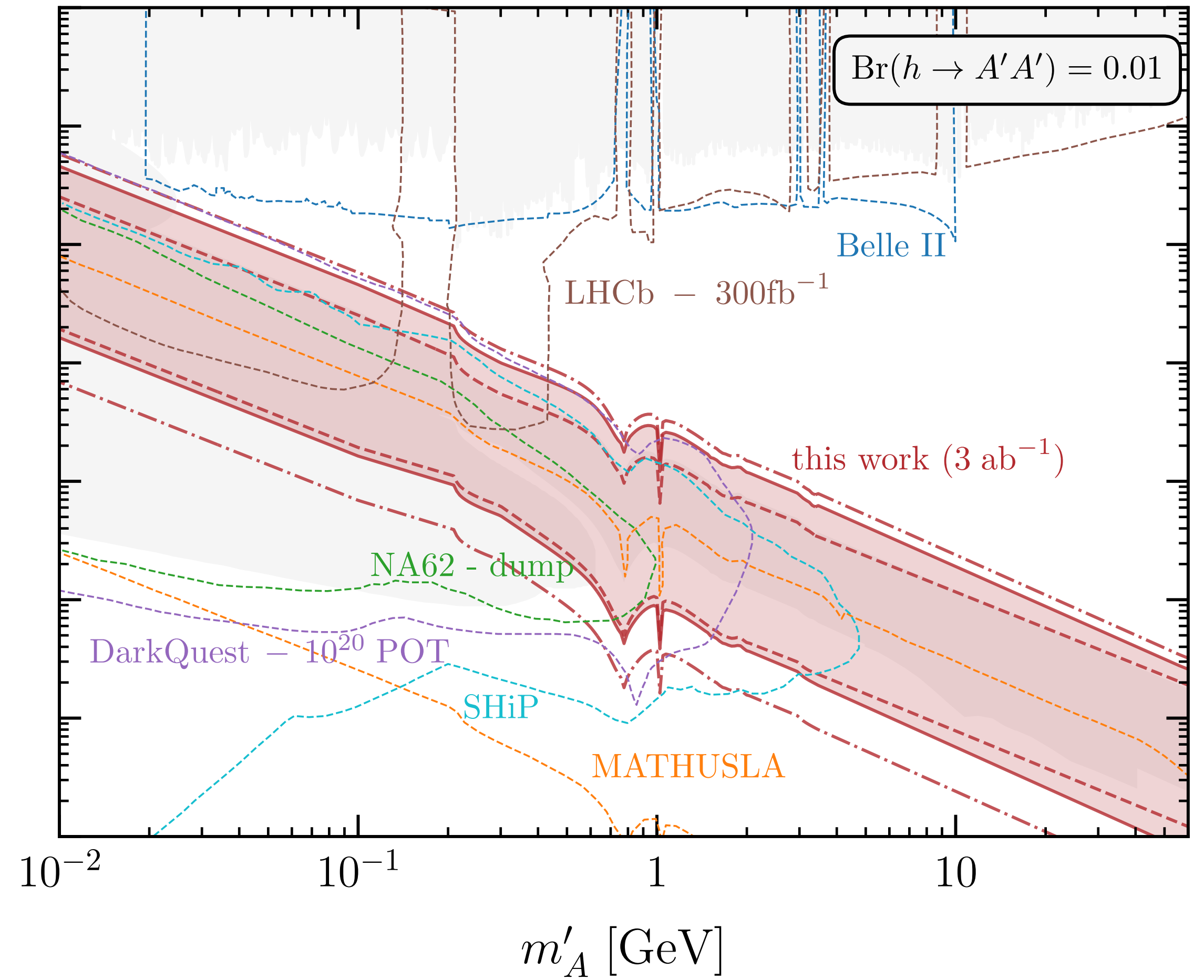
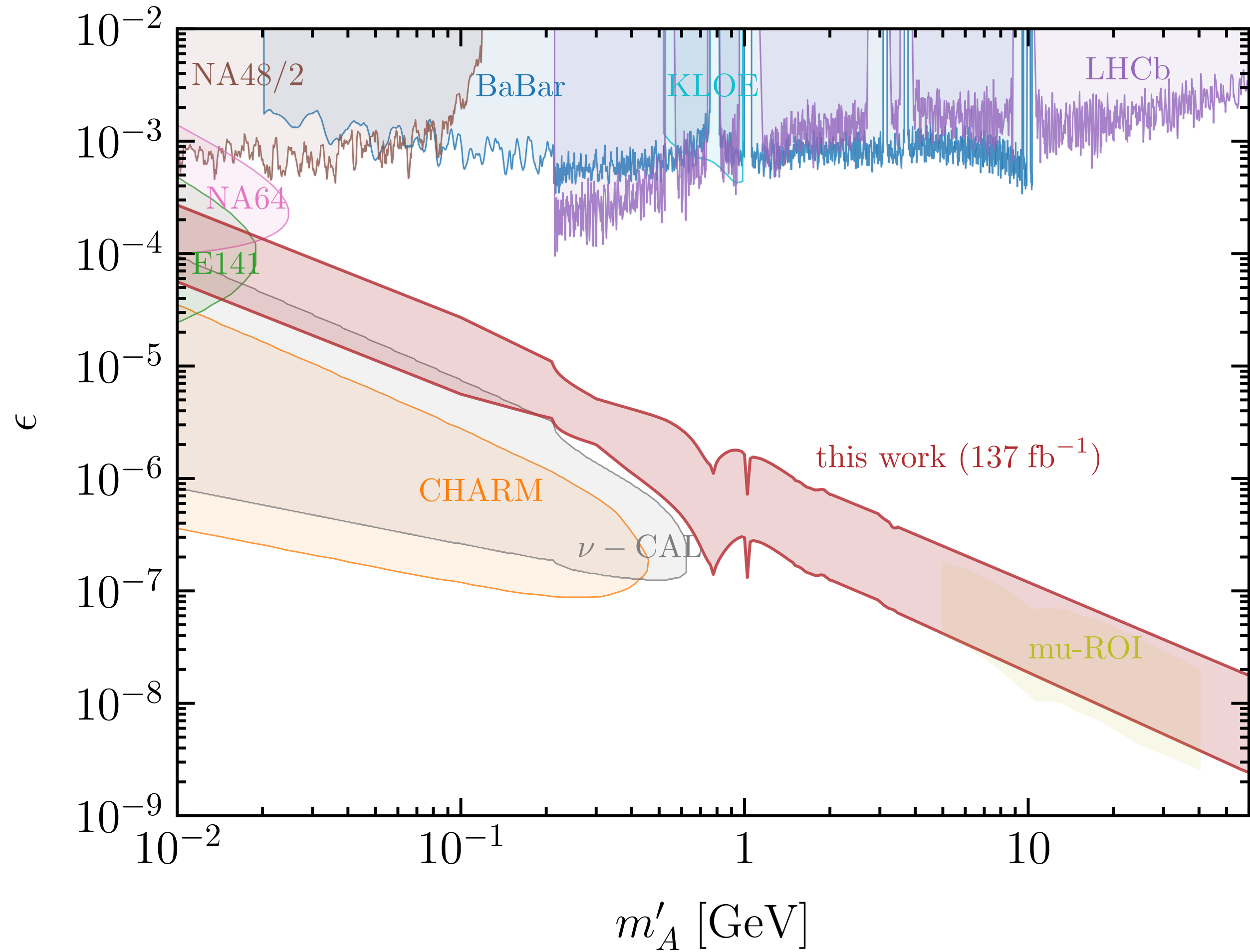


**production and decay** channels are **decoupled**



# DARK PHOTON

$$\mathcal{L}_{SH} = \mathcal{L}_{SM} + \mathcal{L}_{DS} - \lambda_{HS} \hat{S}^2 \hat{H}^\dagger \hat{H} - \frac{\epsilon}{2 \cos \theta_W} \hat{X}_{\mu\nu} \hat{B}^{\mu\nu}$$



# ALP

$$\mathcal{L} = \frac{1}{2} (\partial_\mu a)^2 + \frac{a}{4\pi f_a} \left( \alpha_s c_{GG} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} + \alpha_2 c_{WW} W_{\mu\nu}^a \tilde{W}^{a,\mu\nu} + \alpha_1 c_{BB} B_{\mu\nu} \tilde{B}^{\mu\nu} \right) + \dots$$

