

The University of Manchester

Searching for light long-lived neutralinos at Super-Kamiokande

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Based on

P. Candia, G. Cottin, A. Méndez, V. Muñoz, Phys.Rev.D 104 (2021) 5, 055024

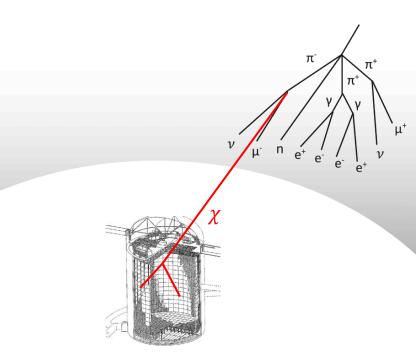
Cosmic-ray air showers

- The earth is constantly being bombarded with cosmic rays
- These cosmic rays collide with air nuclei producing a big amount of unstable hadrons
- These hadrons decay into other particles that form part of the atmospheric shower
- Atmospheric neutrino experiments look at neutrinos that are produced in these showers

			/	
	/	π'	π^+	τ⁺ \
	1	v A	Λv	
ν μ΄	n e+	e-		ν
			e e⁺	

BSM particles in air showers

- Hypothetical long-lived particles could be produced in cosmic ray air showers
- The particles could travel and decay inside the detection volume of atmospheric neutrino detectors
- This would result in an excess of events
 - Kusenko, Pascoli, Semikoz, JHEP 11 (2005) 028 (Sterile neutrinos at SK)
 - Asaka, Watanabe, JHEP 07 (2012) 112 (Sterile neutrinos at SK)
 - Argüelles, Coloma, Hernandez, Muñoz, JHEP 02 (2020) 190 (HNL and dark photons at SK and IC)
- Advantages:
- Long lifetimes
- > There is data already available!



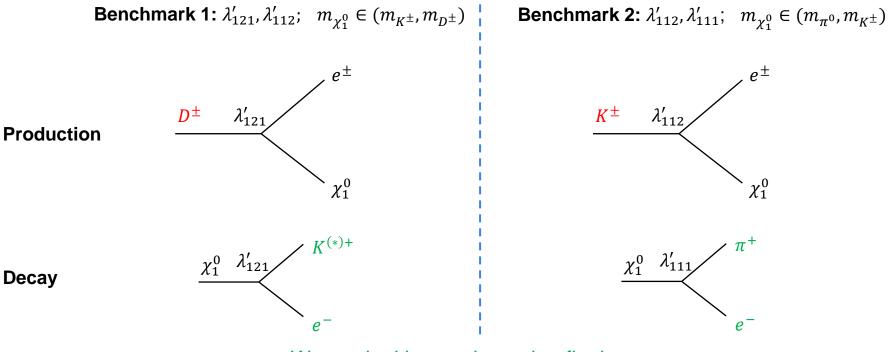
Case study: Long-lived light neutralino

- We apply this strategy to search for long-lived light neutralinos in Super-Kamiokande
- Light neutralinos are allowed in the RPV MSSM. We consider the term

$$W = \lambda_{ijk}' \hat{L}_i \hat{Q}_j \hat{D}_k$$

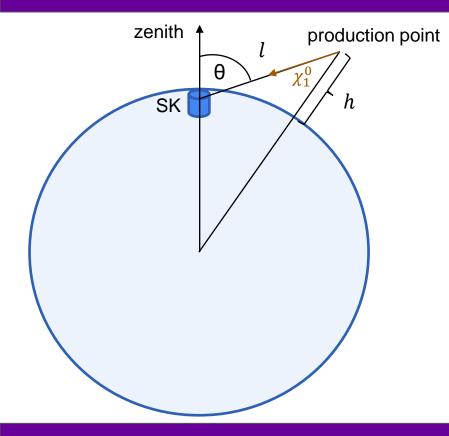
- The widths are parametrised by λ'_{ijk} (if they are smaller, it lives longer)
- Some studies:
 - SHiP, ATLAS: de Vries, HD, Schmeier, PRD 94 (2016) 035006; Gorbunov, Timiryasov, PRD 92 (2015) 7, 075015
 - > CODEX-b, FASER, MATHUSLA: Dercks, de Vries, Dreiner, Wang; PRD 99 (2019) 055039
- We focus on production from meson decays and degenerate sfermion masses. Specifically, we choose λ'_{121} , λ'_{112} and λ'_{111}

Neutralino production and decay



We are looking at showering final states

Neutralino trajectories



The following coordinates are relevant for the production of the neutralino:

- *l*: Distance between the production point and the detector
- θ: zenith angle of the trajectory of the neutralino
- *h*: height at which the neutralino is produced

Column depth:
$$X = \int \rho(h, l) dl$$

Cascade equations and event rate

Cascade equation

$$\frac{d\Phi_{\chi_1^0}}{dE_{\chi_1^0}d\Omega dX} = \sum_M \int dE_M \frac{1}{\rho\lambda_M} \frac{d\Phi_M}{dE_{\chi_1^0}d\Omega} \frac{dn}{dE_{\chi_1^0}}$$

Differential neutralino production rate

$$\int dE_M \frac{1}{\rho \lambda_M} \frac{d\Phi_M}{dE_{\chi_1^0} d\Omega} \frac{dn}{dE_{\chi_1^0}}$$
from MCEq

https://github.com/afedynitch/MCEg

 ρ : atmospheric density at depth X

 λ_M : decay length of the meson

 $\frac{d\Phi_M}{dE_{\chi_1^0}d\Omega}$: production rate of the parent meson

 $\frac{dn}{dE_{\chi_1^0}}$: distribution of neutralinos produced in meson decays

Differential event rate

Probability of decay in the detector volume

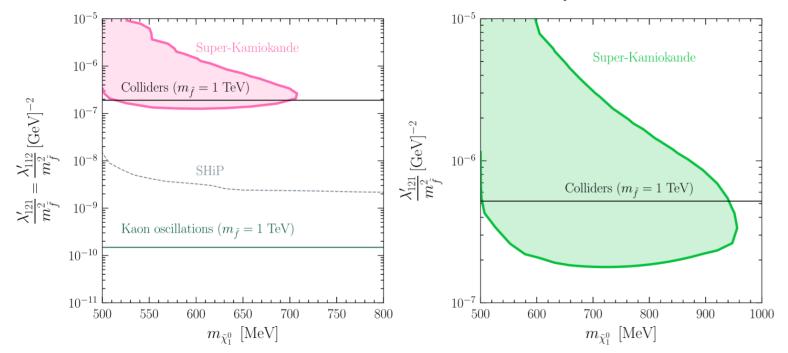
$$\frac{dN}{dE_{\chi_1^0}d\Omega} = \int dS_{\perp} \int dX \frac{d\Phi_{\chi_1^0}}{dE_{\chi_1^0}d\Omega dX} e^{-l/\lambda_{\chi_1^0}} (1 - e^{-\Delta l_{det}/\lambda_{\chi_1^0}})$$

Depends on neutralino lifetime $\lambda_{\chi_1^0}$

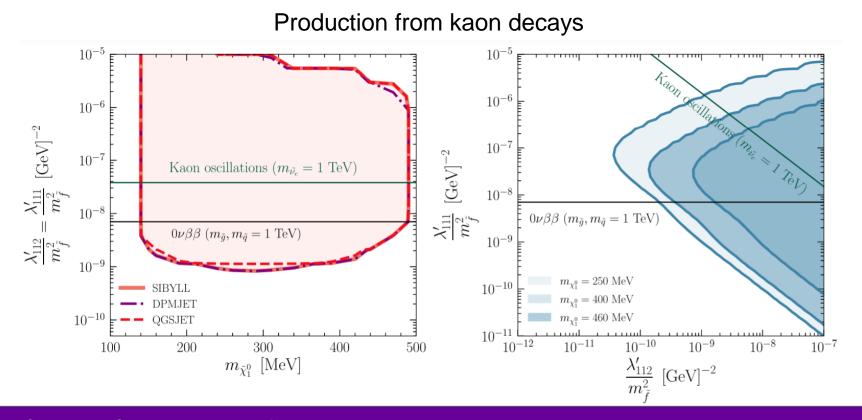
The differential event rate can be multiplied by the exposure time and integrated to get the expected number of events

Results for benchmark 1

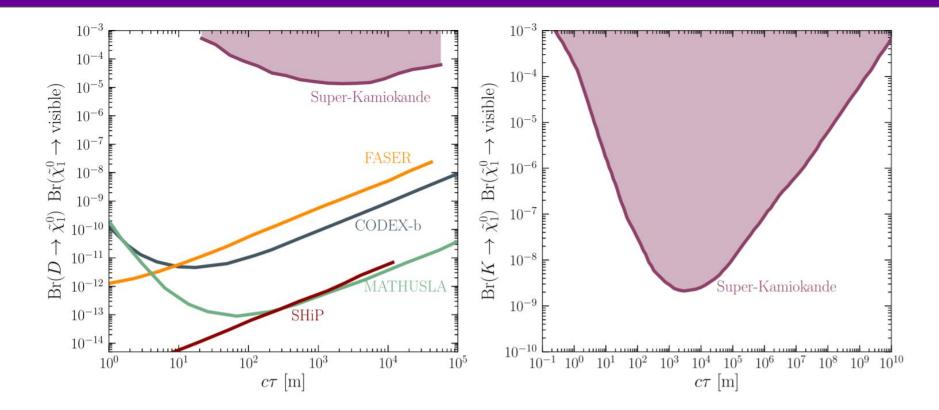
Production from *D* mesons decays



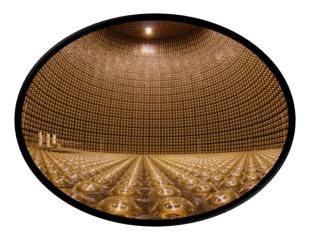
Results for benchmark 2

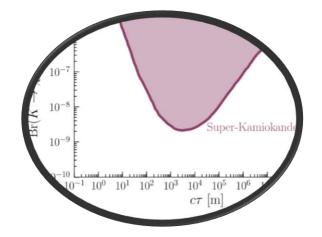


Results in the Br vs. $c\tau$ plane



To summarise



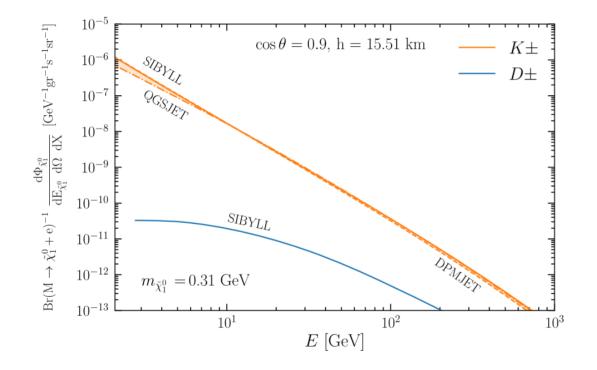


Water Cherenkov detectors can be good at searching for atmospheric long lived particles from meson decays This search can probe a wide range of lifetimes (peaking at order 1 km)

Many possible directions to follow from here!

BACK-UP SLIDES

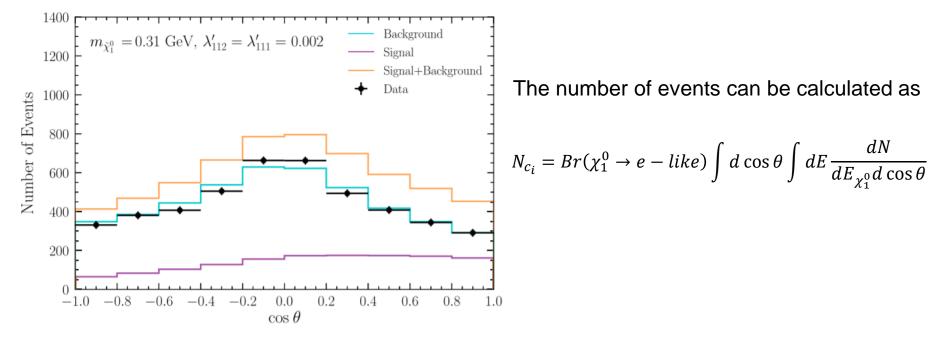
Neutralino production rate



Super-K e-like events dataset

Super-K runs I-IV (5326 days)

*Data obtained from [1710.09126]



Production and decay channels

	RPV coupling	Production	Decay mode
			$\tilde{\chi}^0_1 \xrightarrow{\lambda'_{121}} K^0_S + \nu_e$
B1	$\lambda'_{121}, \lambda'_{112}$		${\widetilde \chi}^0_1 \xrightarrow{\lambda'_{121}} K^{*0} + \nu_e$
		$D^{\pm} \stackrel{\lambda_{121}'}{\longrightarrow} e^{\pm} + ilde{\chi}_1^0$	${ ilde \chi}^0_1 \stackrel{\lambda'_{112}}{\longrightarrow} K^{(*)+} + e^-$
			$\tilde{\chi}^0_1 \xrightarrow{\lambda'_{112}} K^0_S + \nu_e$
			${\widetilde \chi}^0_1 \xrightarrow{\lambda'_{112}} K^{*0} + \nu_e$
B2	$\lambda'_{112}, \lambda'_{111}$	<i>λ</i> 'μ2	$ ilde{\chi}^0_1 \stackrel{\lambda'_{111}}{\longrightarrow} \pi^+ + e^-$
		$K^{\pm} \xrightarrow{\lambda'_{112}} e^{\pm} + ilde{\chi}^0_1$	$\tilde{\chi}_1^0 \xrightarrow{\lambda'_{111}} \pi^0 + \nu_e$

Bounds on RPV parameters

Kaon oscillations
$$\begin{cases} |\lambda'_{112}\lambda'_{121}| \le 2.2 \times 10^{-8} \left(\frac{m_{\tilde{\nu}_e}}{1 \text{ TeV}}\right)^2 \\ |\lambda'_{111}\lambda'_{112}| \le 1.5 \times 10^{-3} \left(\frac{m_{\tilde{\nu}_e}}{1 \text{ TeV}}\right)^2 \\ \lambda'_{112} \le 0.16 \frac{m_{\tilde{s}_R}}{1 \text{ TeV}} + 0.030 \\ \lambda'_{121} \le 0.34 \frac{m_{\tilde{q}}}{1 \text{ TeV}} + 0.18 \\ 0\nu\beta\beta \qquad \lambda'_{111} \le 2.2 \times 10^{-3} \left(\frac{m_{\tilde{q}}}{1 \text{ TeV}}\right)^2 \left(\frac{m_{\tilde{g}}}{1 \text{ TeV}}\right)^{1/2} \end{cases}$$