

Early Matter Domination from LLPs in the Visible Sector

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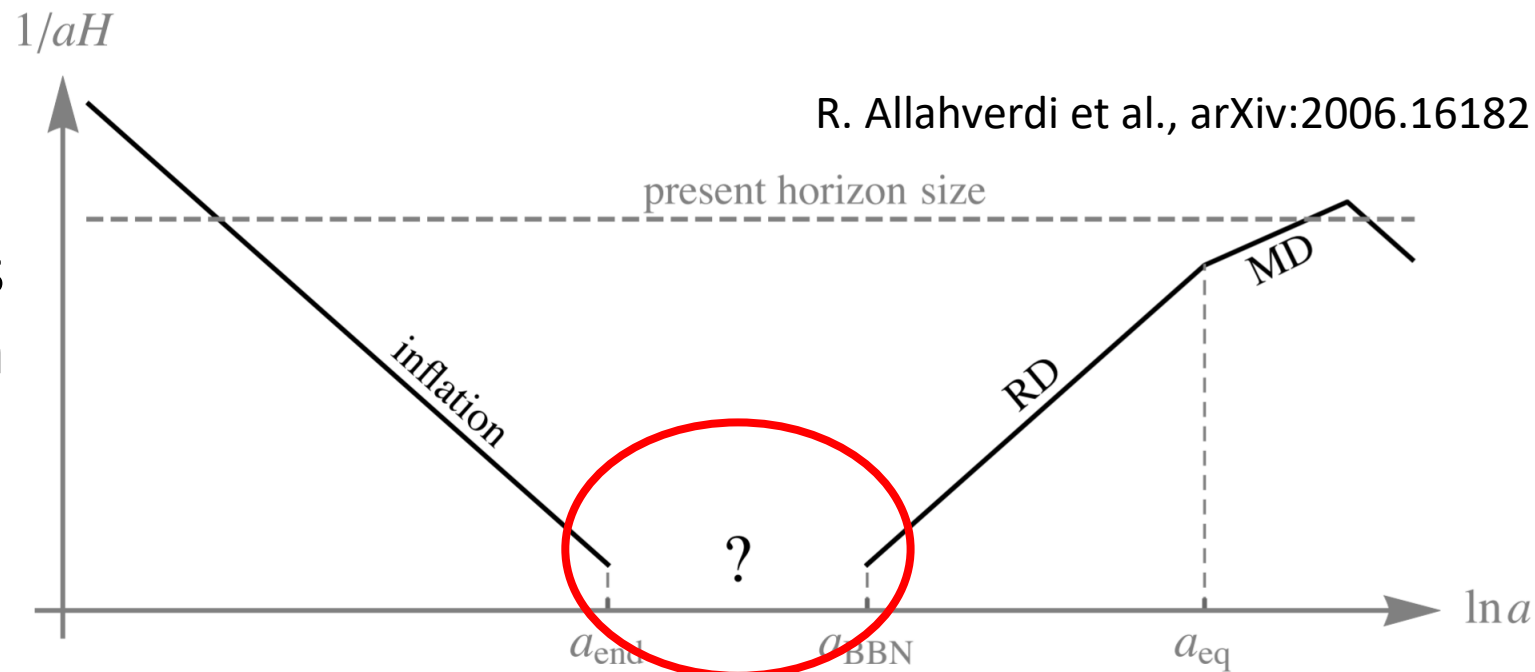
ASTROCENT



Early matter domination (EMD)

- History between inflation and BBN unknown
- Standard assumption is radiation domination (RD)
- EMD well-motivated class of nonstandard expansion (inflationary reheating, decoupled heavy particles, moduli, primordial black holes, ...)
- Constrained by BBN (must recover RD before $T \sim \text{MeV}$)

- Can be probed with cosmological observations and DM indirect detection
- What about colliders?



The scenario

$$\mathcal{L}_{\text{new}} \supset hXN\psi + h'X^\dagger\psi\psi + \text{h.c.}$$

- X : scalar with Standard Model (SM) charges
- N : SM singlet (Majorana fermion) – the LLP
- ψ : SM fermions

$$\begin{aligned} m_X &\gtrsim 1 \text{ TeV} \\ m_N &\sim \text{weak} \\ m_N &\ll m_X \end{aligned}$$

- Effective four-fermion interaction at low energies ($E \ll m_X$):

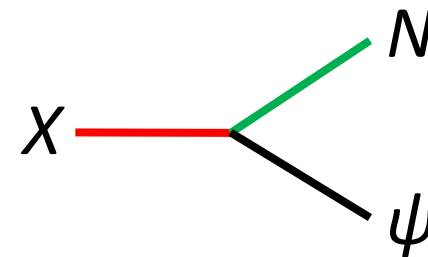
$$hh'^\dagger N\psi\psi\psi / m_X^2$$

- X starts in thermal equilibrium, decays to N which acquires thermal abundance, N freezes-out, dominates, and decays to SM with long lifetime \rightarrow 5 conditions needed
- Can also accommodate observed dark matter relic abundance

Stages of the scenario's history

(1) $H \gtrsim H(T = m_X)$

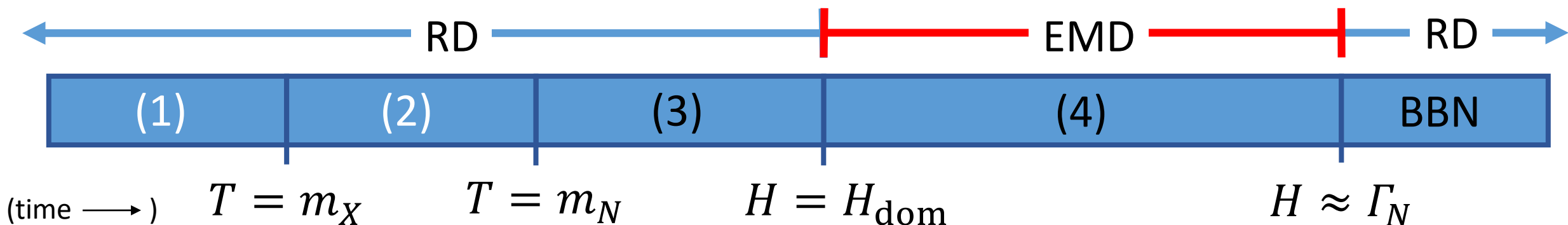
- Start RD at some high temperature
- X in thermal equilibrium due to interactions with SM
- N acquires thermal abundance from X decay, provided that:



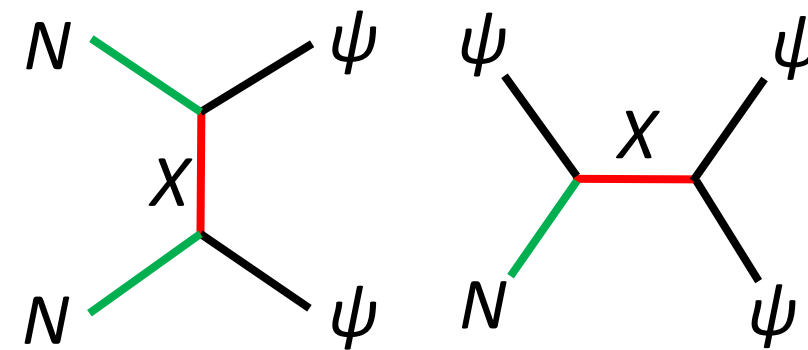
$$\Gamma_{X \rightarrow N} \gtrsim H(T = m_X)$$

(2) $H(T = m_X) > H \gtrsim H(T = m_N)$

- N is relativistic and in equilibrium



Stages of the scenario's history



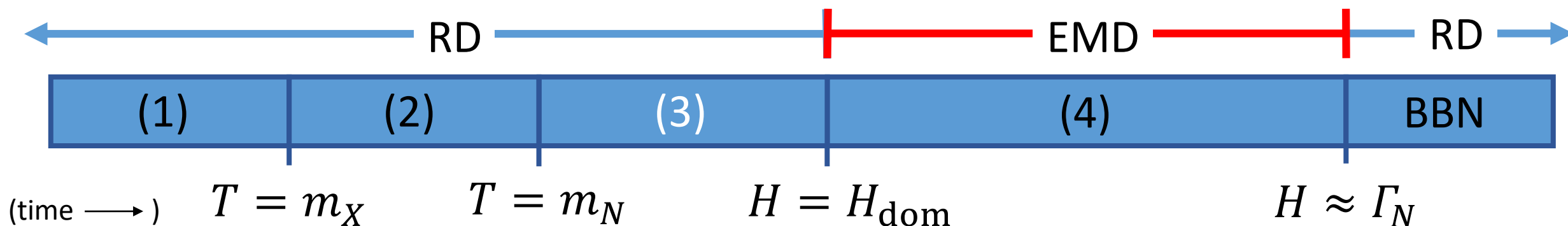
$$(3) \quad H(T = m_N) > H \gtrsim H_{\text{dom}}$$

- N becomes nonrelativistic and freezes-out provided that scatterings/annihilations are inefficient:
- N dominates over radiation, provided that lifetime is long enough:

$$\Gamma_{NN \rightarrow \psi\psi^*} < H(T = m_N)$$

$$\Gamma_{N\psi \rightarrow \psi^*\psi^*} < H(T = m_N)$$

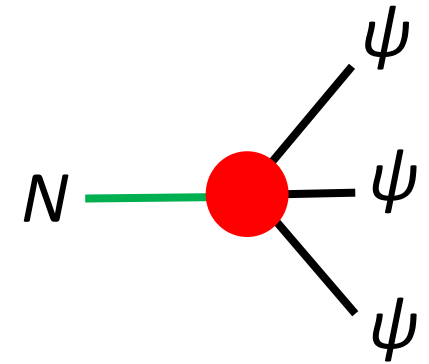
$$\Gamma_N < H_{\text{dom}}$$



Stages of the scenario's history

$$(4) \quad H_{\text{dom}} > H \gtrsim \Gamma_N$$

- EMD period driven by N , diluting prior abundances
- Must end before BBN, limiting N lifetime:

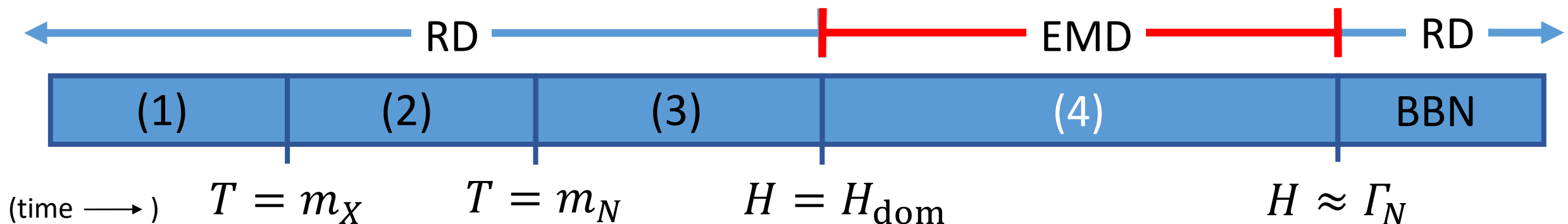


$$T_{\text{BBN}} \simeq 4 \text{ MeV}$$

T. Hasegawa et al., arXiv:1908.10189

P.F. de Salas et al., arXiv:1511.00672

$$\Gamma_N \gtrsim H_{\text{BBN}} \sim \mathcal{O}(10) \text{ s}^{-1}$$



Conditions for successful scenario

- N reaches equilibrium early on:

$$\Gamma_{X \rightarrow N} \gtrsim H(T = m_X)$$

- N doesn't maintain equilibrium for too long:

$$\Gamma_{NN \rightarrow \psi\psi^*} < H(T = m_N) \quad \Gamma_{N\psi \rightarrow \psi^*\psi^*} < H(T = m_N)$$

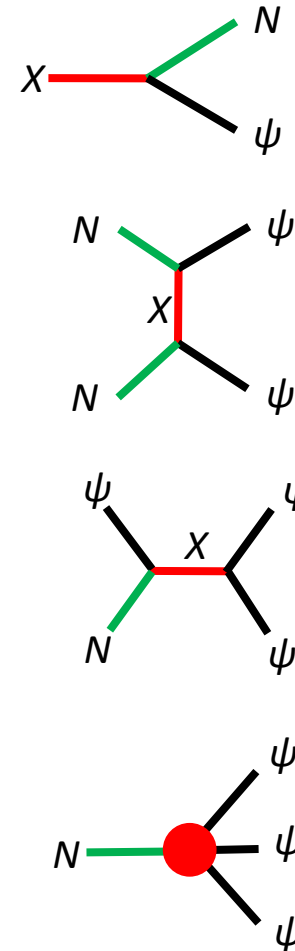
- Long enough lifetime to dominate energy density:

$$\Gamma_N < H_{\text{dom}}$$

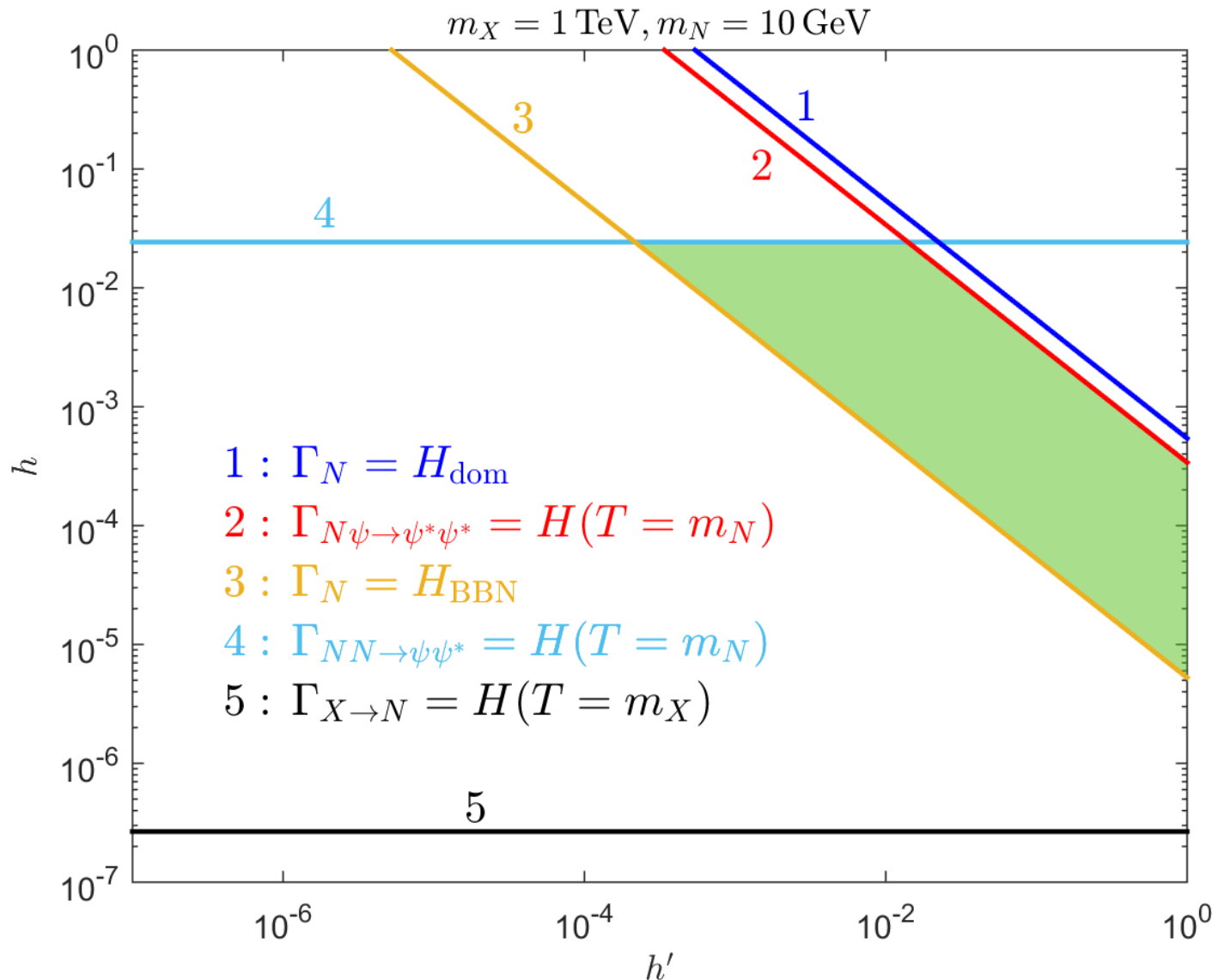
- But not too long... (BBN):

$$\Gamma_N \gtrsim H_{\text{BBN}} \sim \mathcal{O}(10) \text{ s}^{-1}$$

- They can all be met simultaneously!



Allowed parameter space – couplings



$$\Gamma_{X \rightarrow N} \sim |h|^2 m_X$$

$$\Gamma_{NN \rightarrow \psi\psi^*} \sim |h|^4 \frac{m_N^5}{m_X^4}$$

$$\Gamma_{N\psi \rightarrow \psi^*\psi^*} \sim |h|^2 |h'|^2 \frac{m_N^5}{m_X^4}$$

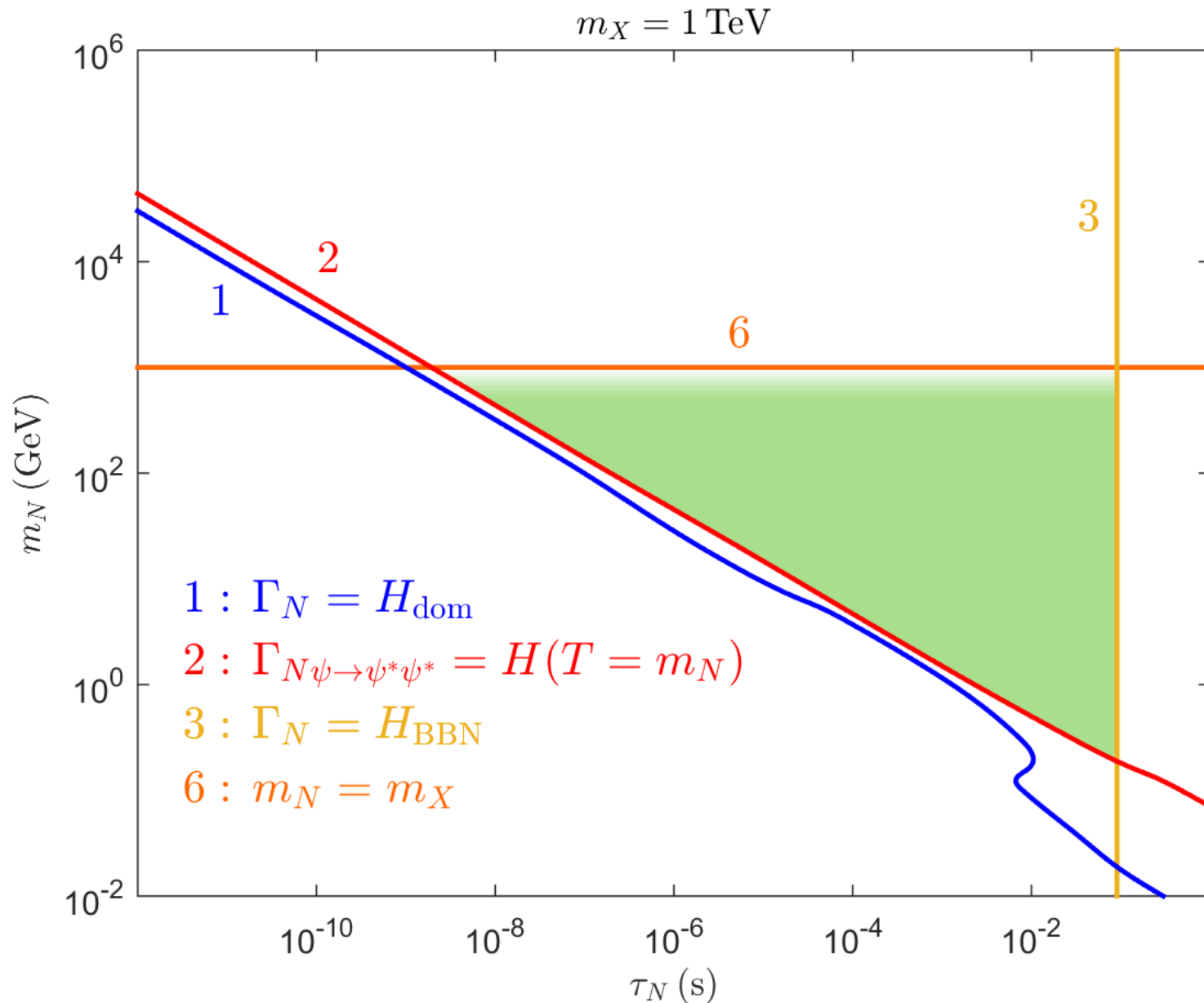
$$\Gamma_N \sim |h|^2 |h'|^2 \frac{m_N^5}{m_X^4}$$

Similar allowed region for

$$m_X = 1 - 10 \text{ TeV}$$

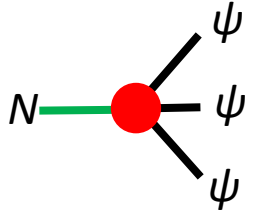
$$m_N = 10 \text{ GeV} - 1 \text{ TeV}$$

Main parameter space of interest – mass vs lifetime ⁹



$$\tau_N \equiv \Gamma_N^{-1}$$

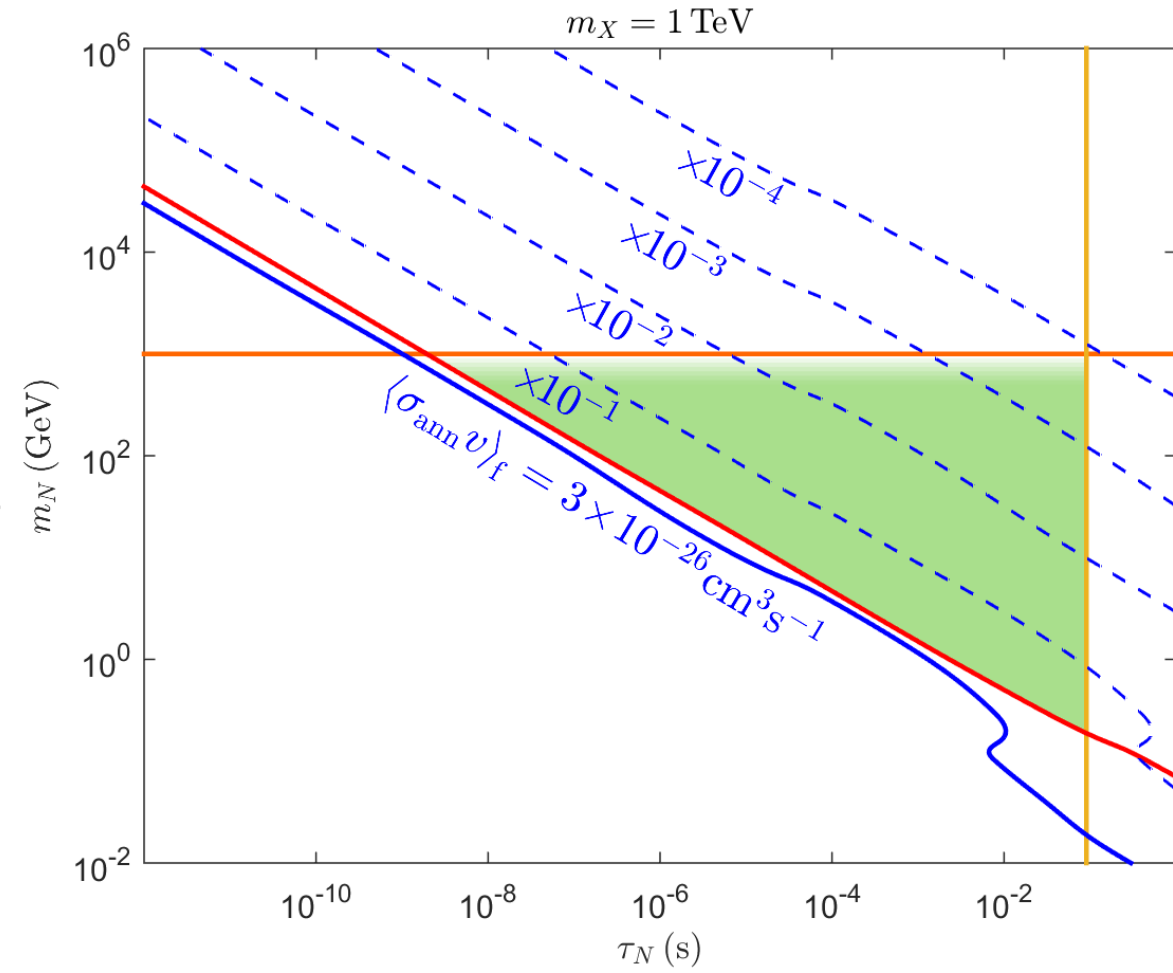
(three-body decay of $N \rightarrow \text{SM fermions}$)



Only dependence on m_X is in $m_N = m_X$ line

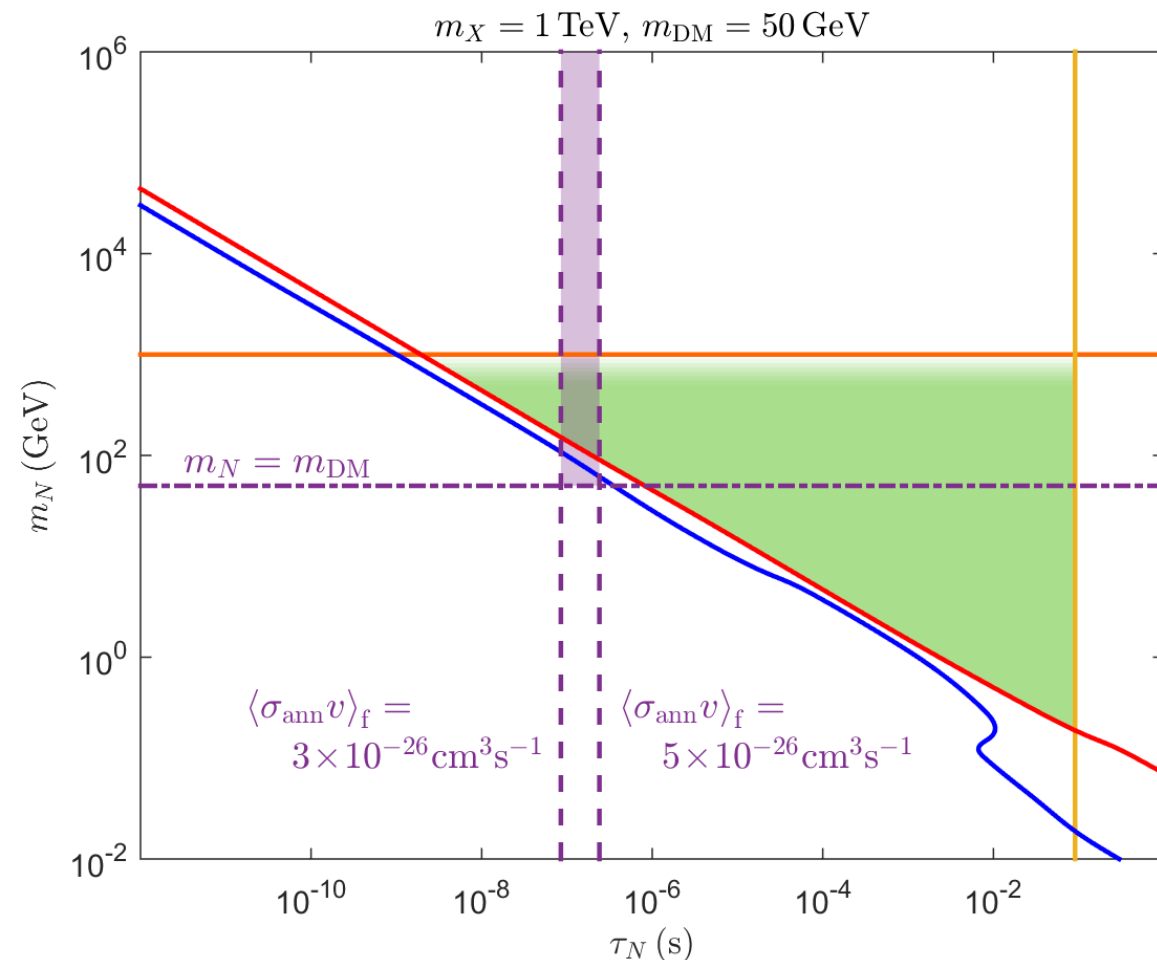
Implications for dark matter (DM)

- Scenario can accommodate DM abundance for both large and small annihilation rates (relative to thermal WIMPs)
- Small: $\langle \sigma_{\text{ann}} v \rangle_f < 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$
- Thermal overproduction in standard RD history
- If $m_N < m_{\text{DM}}$, DM freezes-out prior to EMD and gets diluted



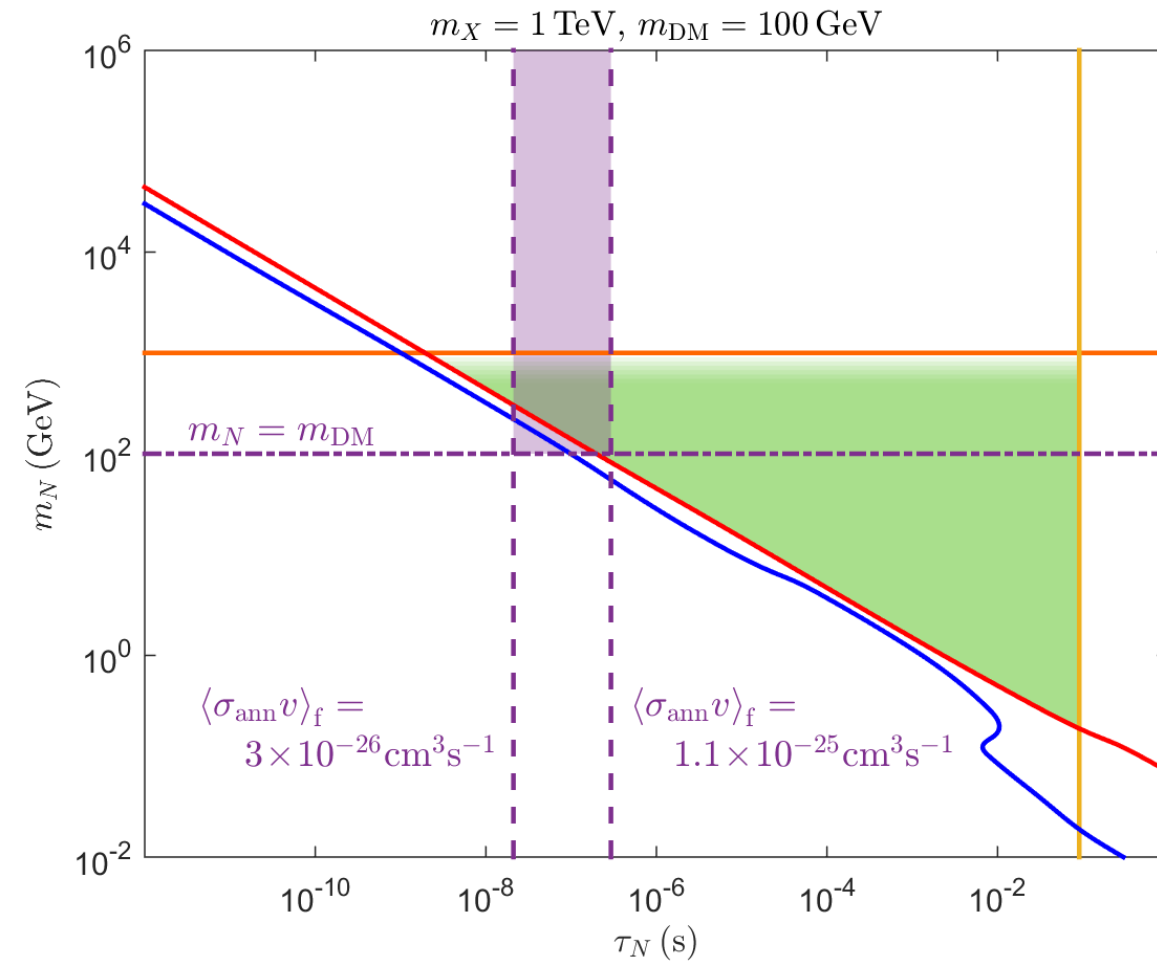
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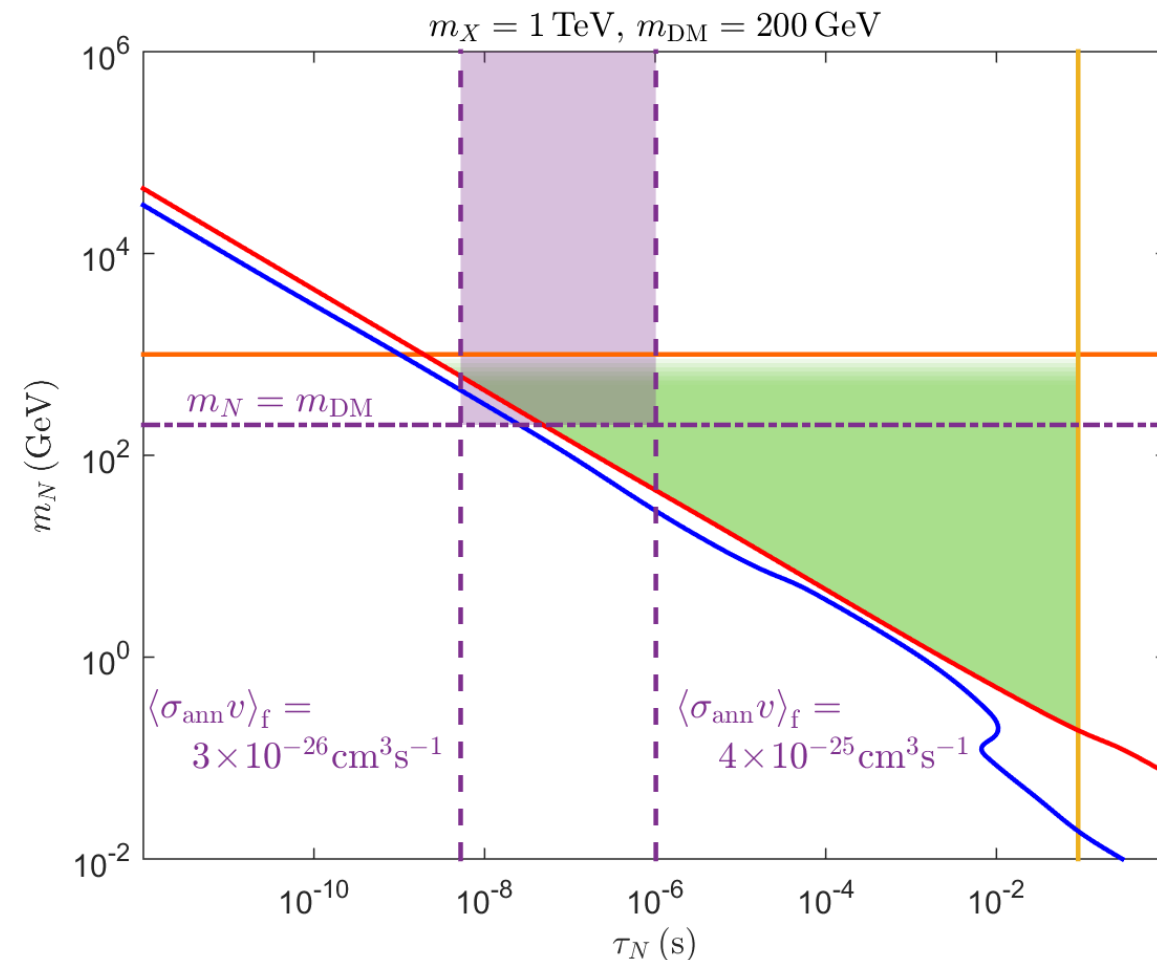
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Collider prospects

- N (LLP) production from X decay
- N decay to 3 SM fermions
- N mass and lifetime give start and end of EMD phase

$$m_N \sim \sqrt{M_P H_{\text{dom}}} \sim T_{\text{dom}}$$



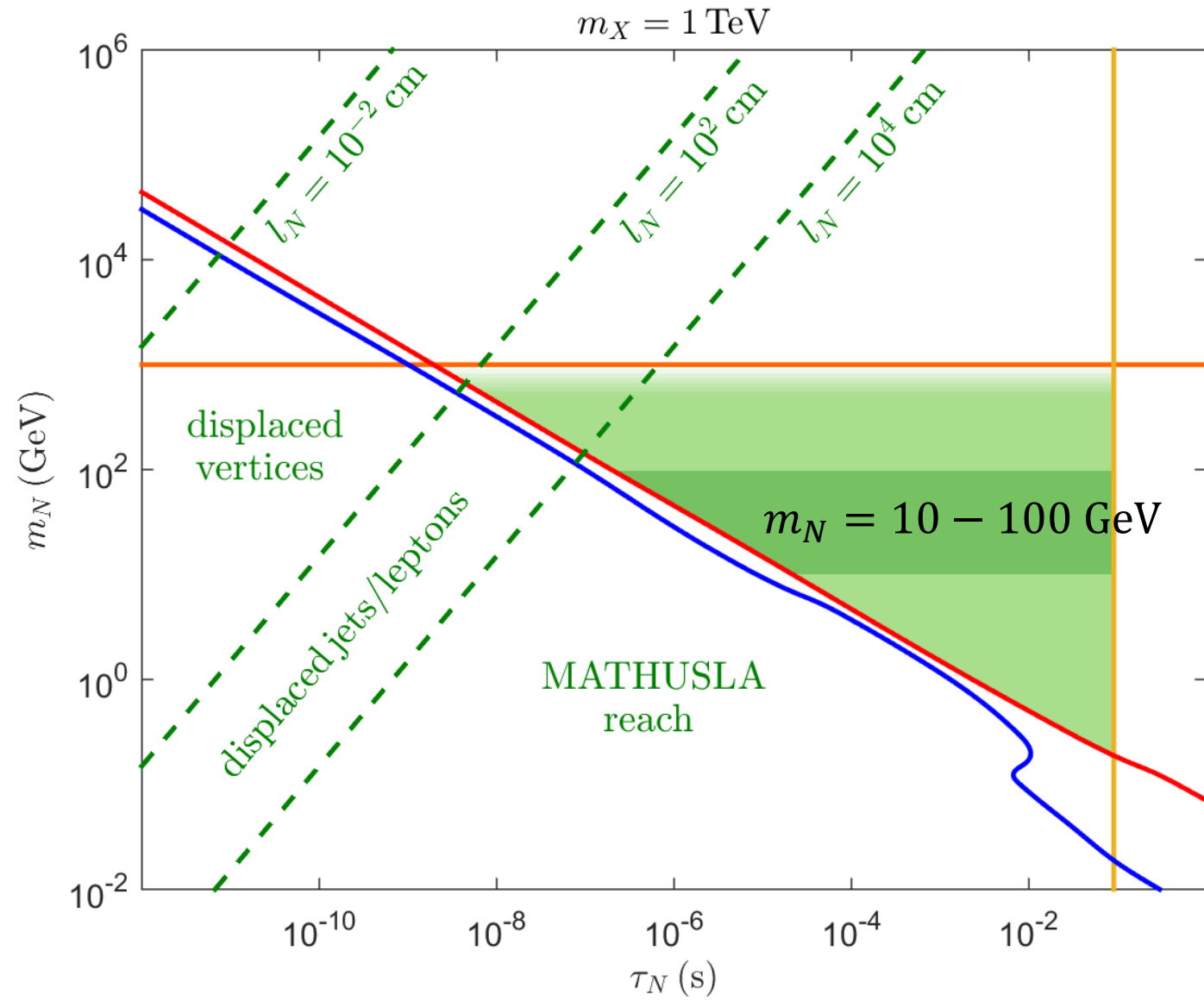
Start of EMD

$$\tau_N \equiv \Gamma_N^{-1} \sim H_{\text{dec}}^{-1} \sim M_P / T_{\text{dec}}^2$$



End of EMD,
start of RD for BBN

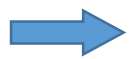
$$l_N = \bar{b} c \tau_N \quad \bar{b} \sim m_X / 2m_N$$



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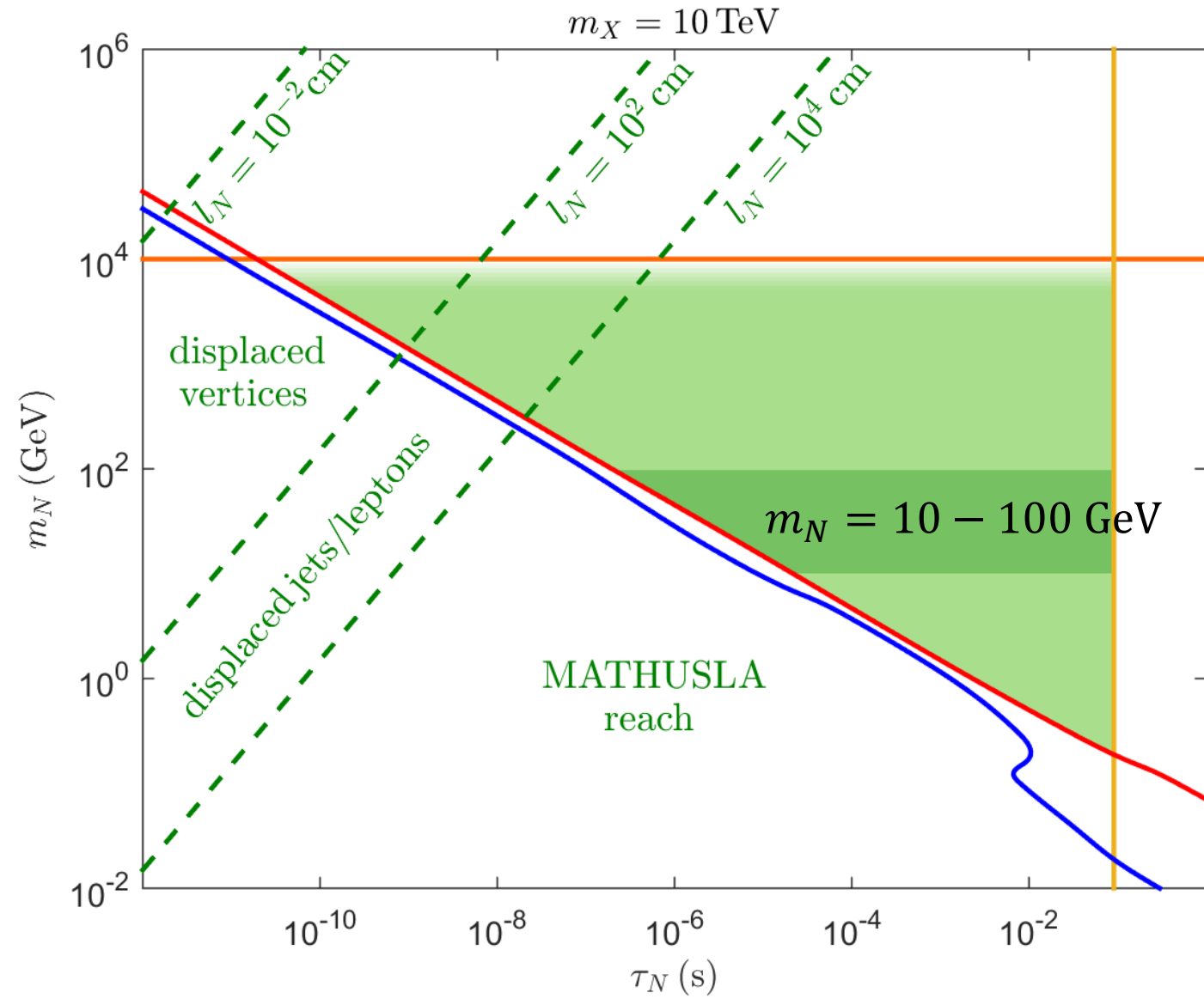
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End of EMD,
start of RD for BBN

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Summary

- Weak-scale LLP can drive EMD period rather naturally (thermal bath)
- Cosmological history conditions are not difficult to satisfy
- LLP lifetime can be quite long (up to BBN), mass relatively high
- Connection between collider observables and pre-BBN era

Thank you!

Backup

Some probes of EMD

- Indirect detection signals of DM annihilation in microhalos boosted due to EMD:
 - M. Sten Delos et al., arXiv:1910.08553
 - C. Blanco et al., arXiv:1906.00010
 - A.L. Erickcek et al., arXiv:1510.04291
 - A.L. Erickcek, arXiv:1504.03335
- DM free-streaming from relativistic velocities at production makes it harder to form MW satellite galaxies and can alter Lyman- α forest:
 - C. Miller et al., arXiv:1908.10369

Rates in more detail

$$\Gamma_{X \rightarrow N} \simeq \frac{|h|^2}{16\pi} m_X$$

$$\Gamma_{NN \rightarrow \psi\psi^*} \simeq C_1 \frac{|h|^4}{16\pi} \frac{E^2}{m_X^4} n_N$$

$$\Gamma_N = 2C_2 \frac{|h|^2 |h'|^2}{128 \cdot 192\pi^3} \frac{m_N^5}{m_X^4}$$

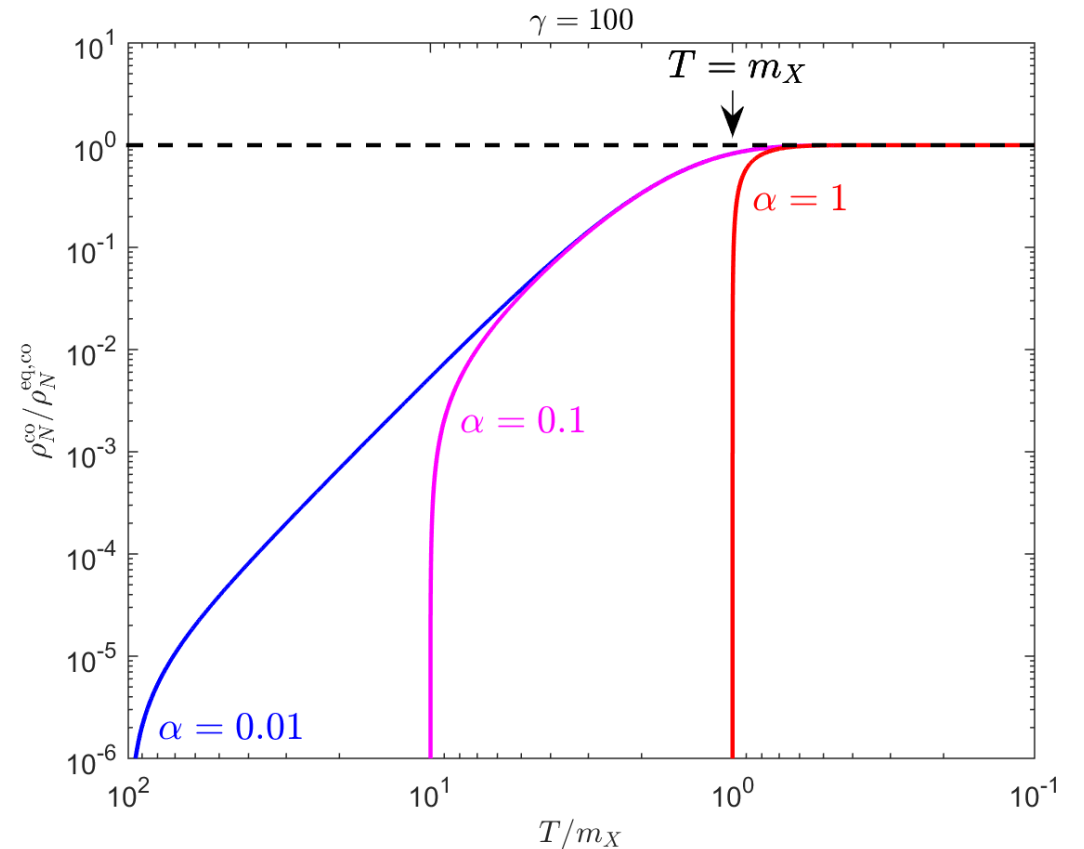
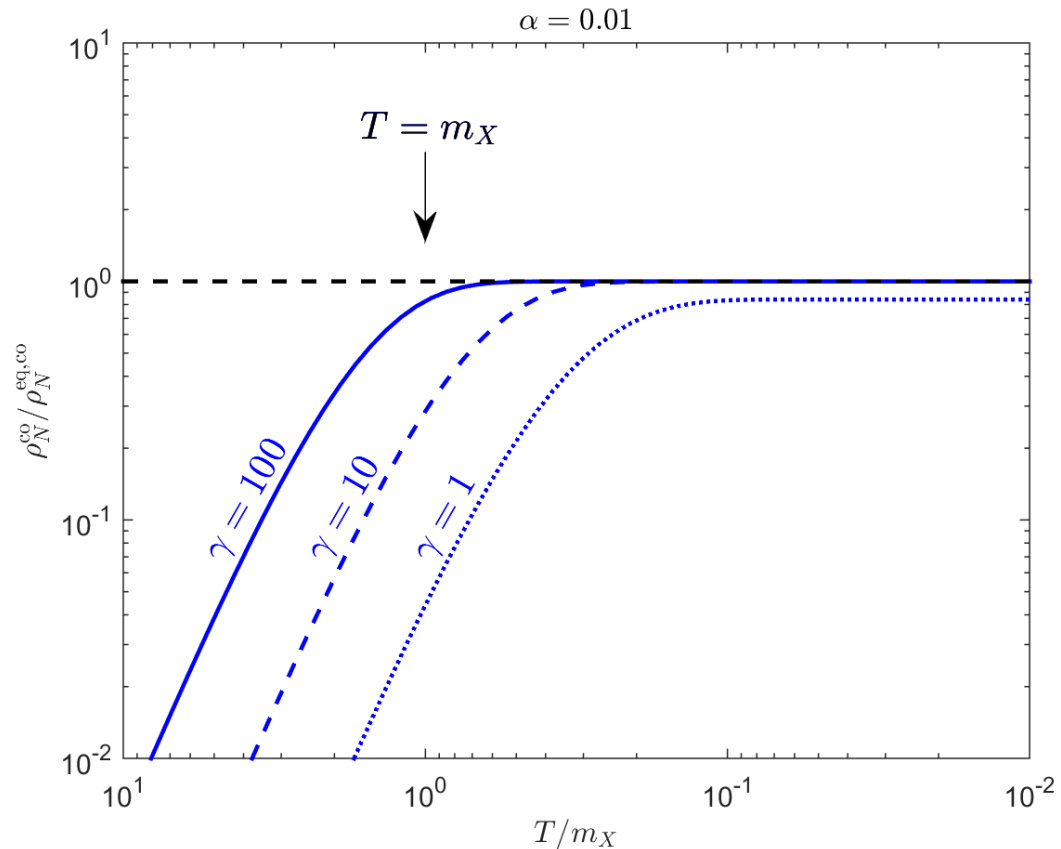
$$\Gamma_{N\psi \rightarrow \psi^*\psi^*} \simeq 3C_2 \frac{|h|^2 |h'|^2}{16\pi} \frac{E^2}{m_X^4} n_\psi$$

- ψ SM fermions of a specific chirality
- Consider X coupling to given flavor combinations in NN and $N\psi$ for simplicity (need to sum over all flavor combinations for full)
- C 's are $O(1)$ model-dependent multiplicity factors
- n 's are fermionic relativistic equilibrium number densities
- $E \sim m_N$ (evaluated at $T = m_N$ because of decoupling condition)

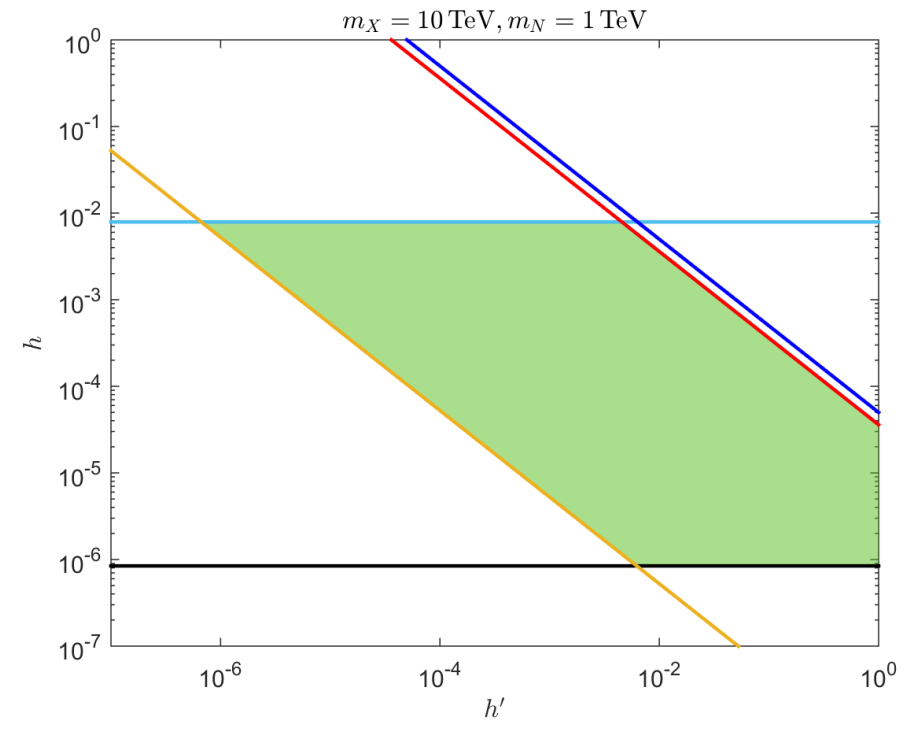
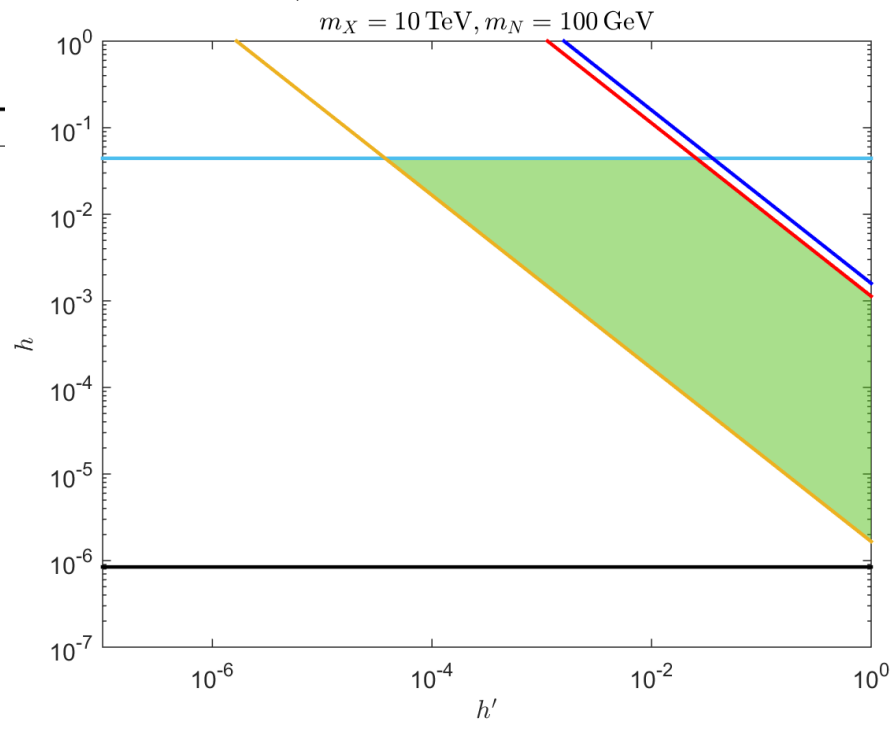
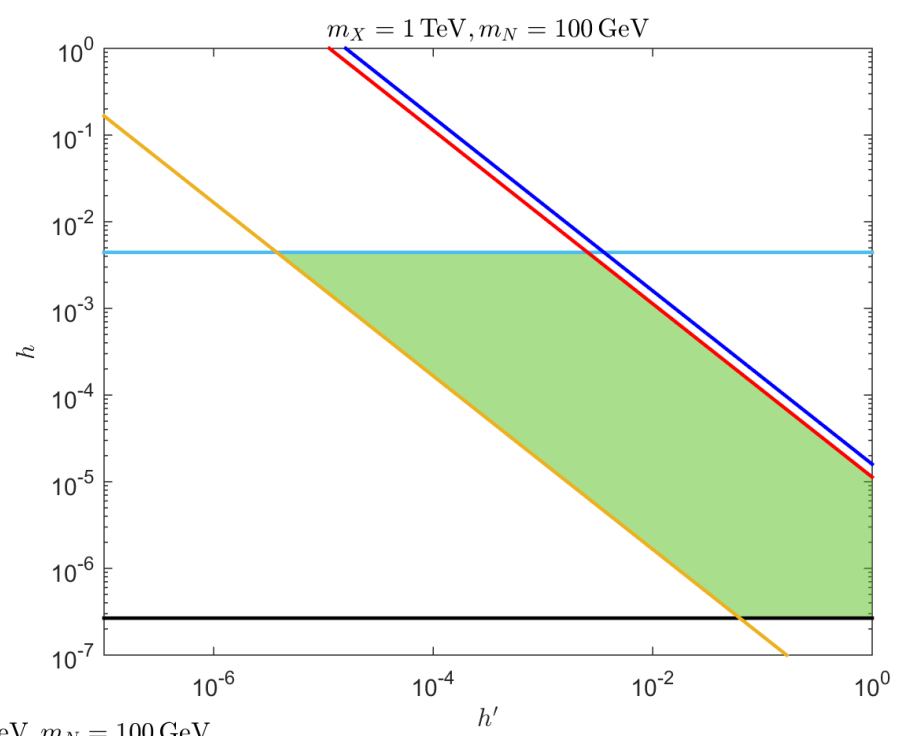
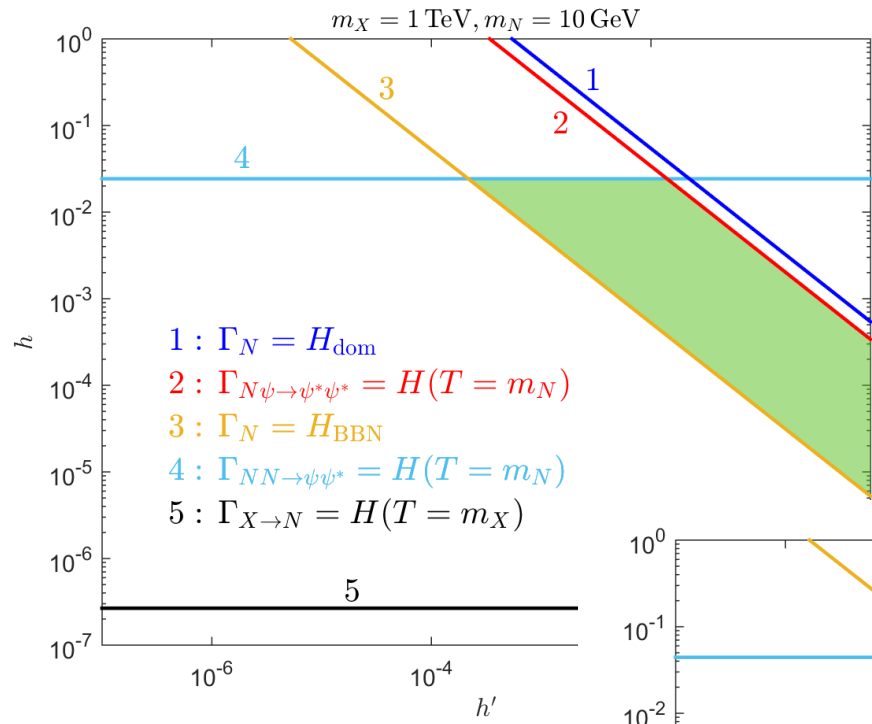
Reaching equilibrium in early Universe

- Thermal abundance of N maximizes duration of EMD
- Boltzmann suppression if X (or N) becomes nonrelativistic while maintaining equilibrium
- Comoving energy densities normalized to equilibrium shown below

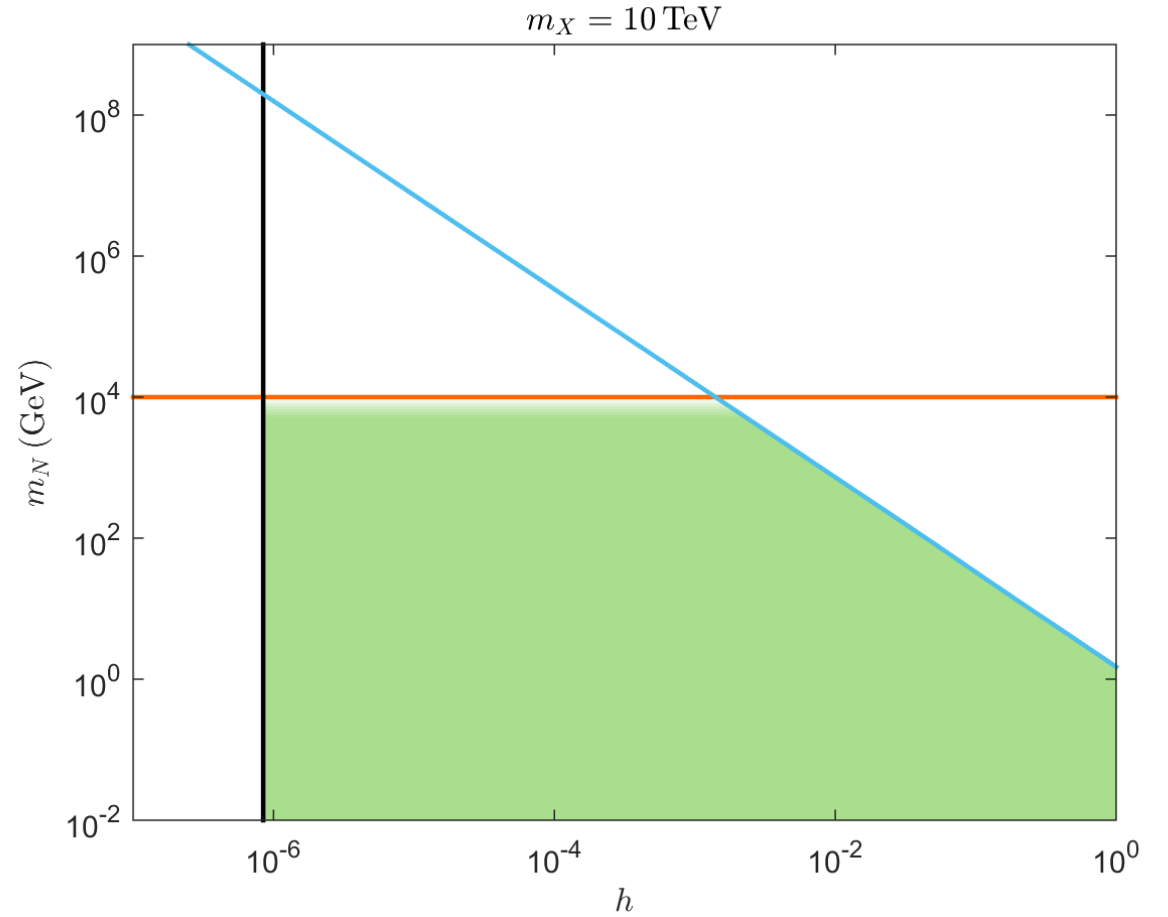
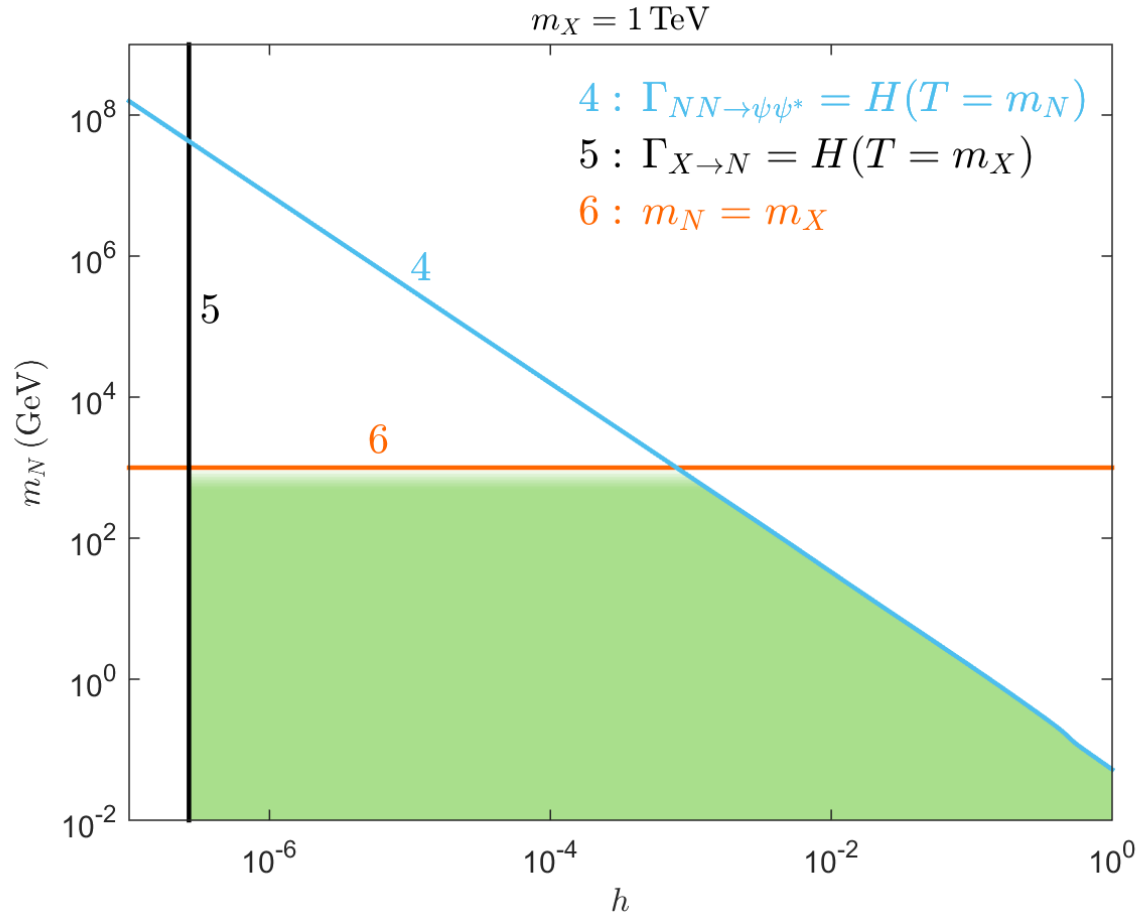
$$\gamma \equiv \Gamma_{X \rightarrow N} / H(T = m_X) \quad \alpha \equiv m_X / T_i$$



More $h - h'$ plane



Relating the two parameter spaces



An explicit model

$$\mathcal{L} \supset (h_i X N u_i^c + h'_{ij} X^* d_i^c d_j^c + \frac{1}{2} m_N N N + \text{h.c.}) + m_X^2 |X|^2$$

- i, j, k flavor indices
- u/d left-handed up/down-type antiquarks
- X iso-singlet color-triplet scalar with hypercharge $+4/3$
- Effective four-fermion interaction at low energies: $N u_i^c d_j^c d_k^c$
- N decay can source baryon asymmetry
- Also has DM candidates
- Testable at colliders

Model:

R. Allahverdi and B. Dutta, arXiv:1304.0711

K.S. Babu et al., arXiv:hep-ph:0612357

R. Allahverdi et al., arXiv:1305.0287

Testability:

B. Dutta et al., arXiv:1401.1825

R. Allahverdi et al., arXiv:1507.02271

R. Allahverdi et al., arXiv:1712.02713

P.S. Bhupal Dev et al., arXiv:1504.07196