Search for new light particles using MeV photon beam at the ILC

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in collaboration with:
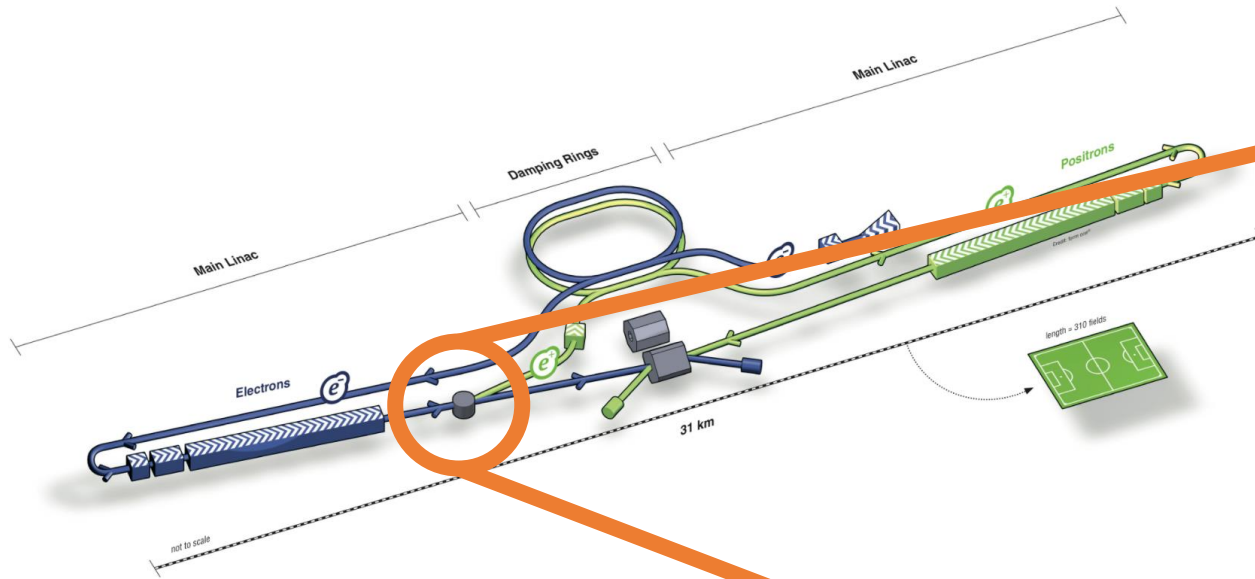
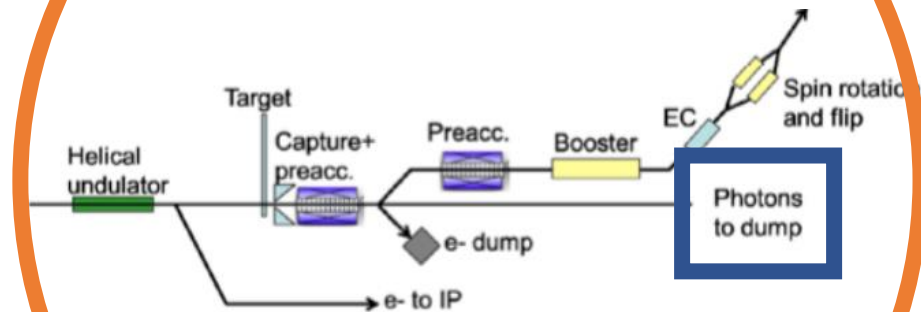Hidetoshi Otono (Kyushu U.) Satoshi Shirai (Kavli IPMU)

Positron source at the ILC

- It is a challenge to generate enough e^+ at the ILC
- Two methods are proposed to create e^+e^- pair:
  - Use a 3 GeV e^- beam created at an independent LINAC
  - **Use a γ beam created from the 125 GeV e^- beam at a undulator**

Let's focus on this design

Positron from the 125 GeV electron

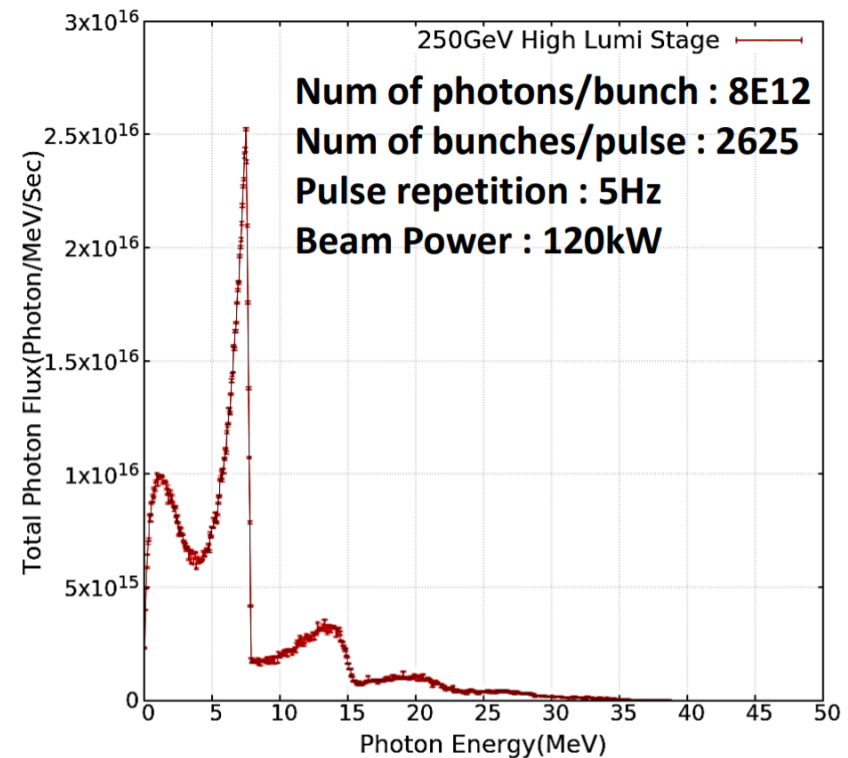From ILC collaboration

Riemann et al. 2002.10919

- An oscillating magnetic field is imposed at the undulator and a photon beam is created
- The photon beam then create e+e- pair at a target

Photon beam

- As a biproduct, an energetic and intense photon beam is obtained

- Energy: $\mathcal{O}(1-10)$ MeV

- Intensity: $\mathcal{O}(10^{16})$ γ/sec

more intense than any other MeV

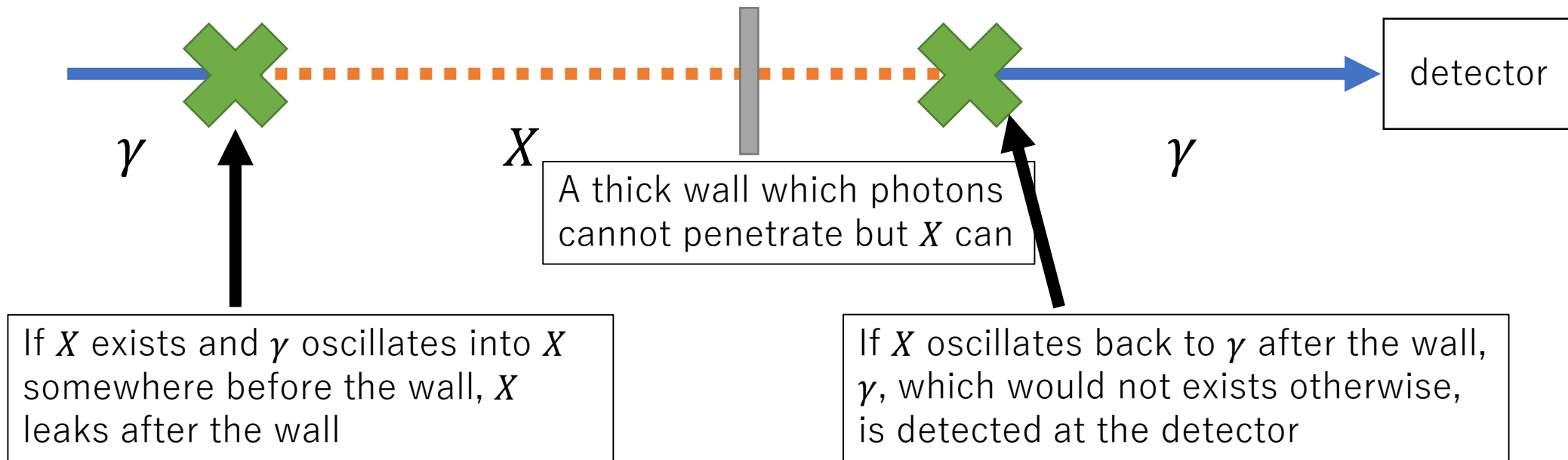γ source up to the present

→Our motivation

Can we do anything
with this photon beam?

250GeV High Lumi Stage

Num of photons/bunch : 8E12
Num of bunches/pulse : 2625
Pulse repetition : 5Hz
Beam Power : 120kW

From Morikawa-san's
slide at POSIPOL-2018

Search for new light particles

- A new light particle X may mix with the photon

γ

X

A thick wall which photons cannot penetrate but X can

γ

detector

If X exists and γ oscillates into X, somewhere before the wall, X leaks after the wall

If X oscillates back to γ after the wall, γ, which would not exists otherwise, is detected at the detector

Called the light-shining-through-the-wall (LSW) experiment

Light particle models

- QCD axion
- Dark photon

- Today, we focus on the QCD axion

QCD Axion

- A hypothetical particle to solve the strong CP problem
- Lagrangian

$$\mathcal{L} = \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{N_{DW} \alpha_s}{8\pi f_a} a G^a_{\mu\nu} \tilde{G}^{a\mu\nu} - \frac{c\alpha}{4\pi f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$
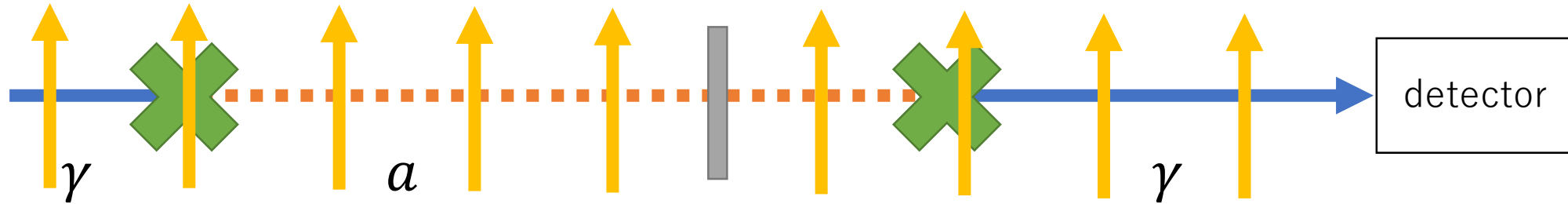
- Very light
  - $m_a \simeq 10^{-2} \left(\frac{10^9 \text{ GeV}}{f_a} \right) \text{ eV}$
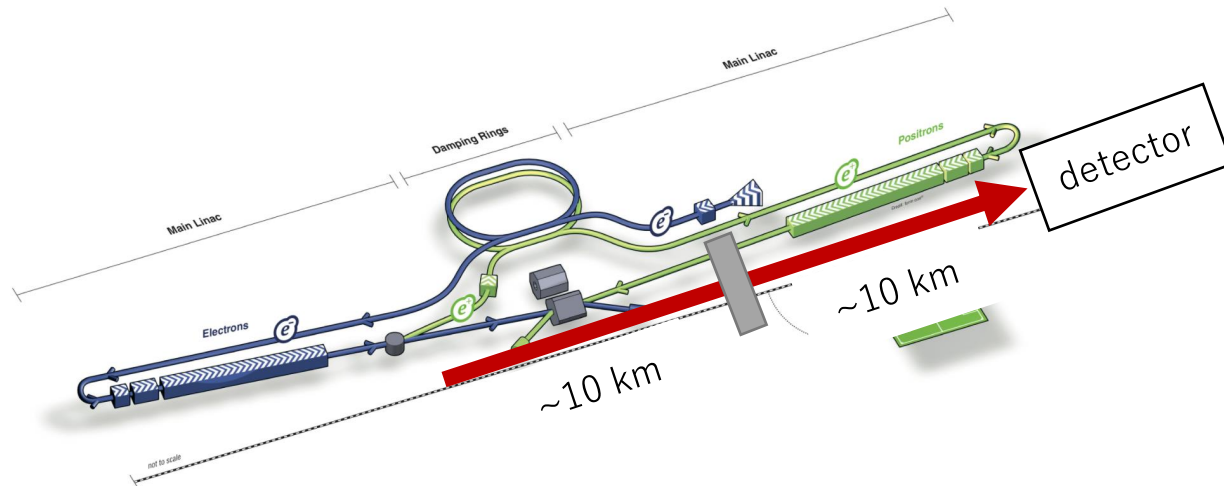- Mix with photon in the presence of a magnetic field
  - $a F_{\mu\nu} \tilde{F}^{\mu\nu} \rightarrow B_0 (a \, \partial A)$

Experimental setup

- Magnetic fields are imposed on the entire path

γ a γ detector

- We assume the path is of the similar length as the ILC itself

The photon beam is very collimated; the experimental facility may be not under but on the ground.

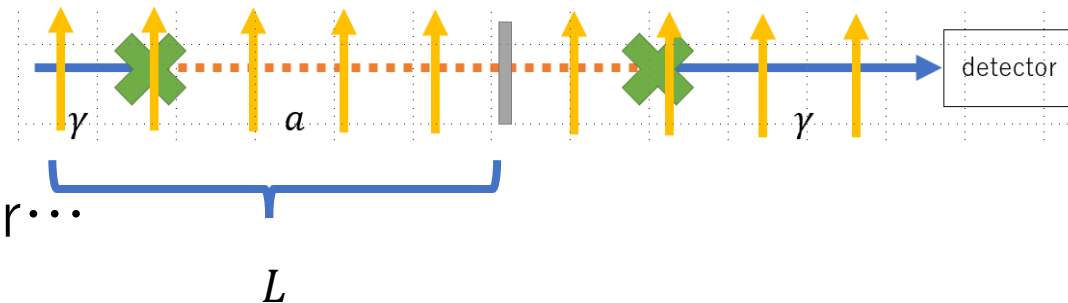# Advantage to use energetic photons

- Conversion probability

$$P(\gamma \to a) = g_{a\gamma\gamma}^2 \left| \int^L dz \; e^{iqz} \, B(z) \right|^2$$

where ⸱   ld configuration

For the longer L to maximize P,
we need larger E!

$$q = E - \sqrt{E^2 - m_a^2} \simeq \frac{m_a}{2E} \sim (10 \text{ km})^{-1} \left(\frac{m_a}{10^{-2} \text{ eV}} \right) \left(\frac{2.5 \text{ MeV}}{E} \right)$$
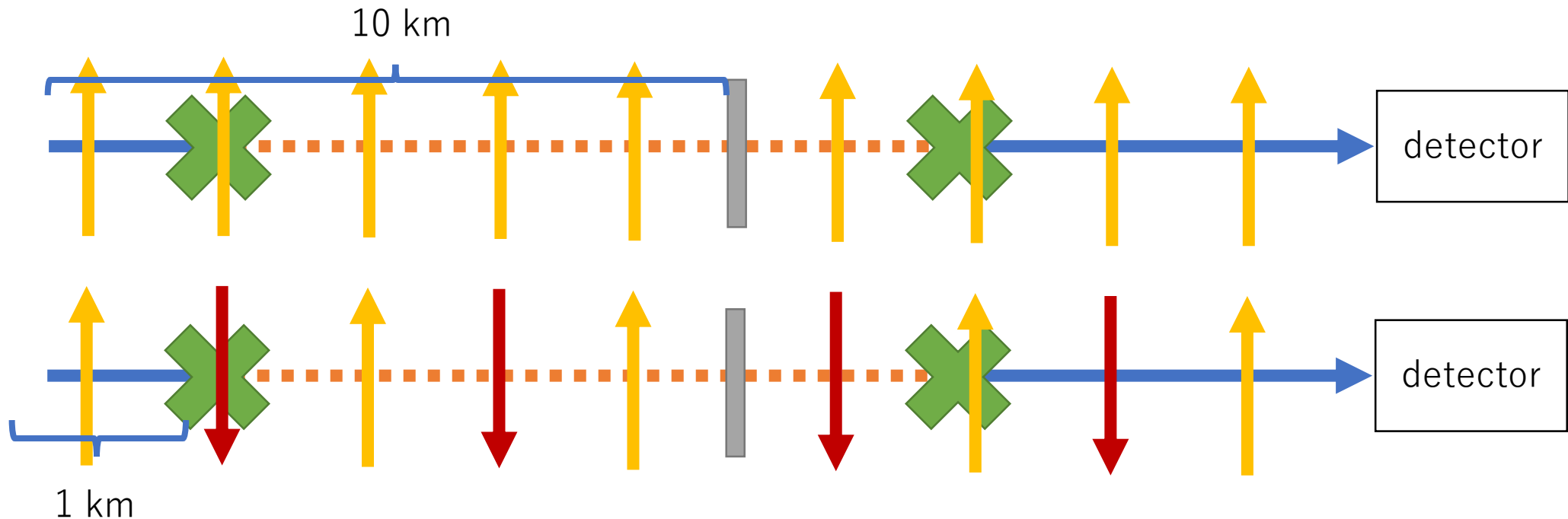
- To maximize P,
  - Longer BL is better
  - But then, a momentum $\sim L^{-1}$ is smaller⋯
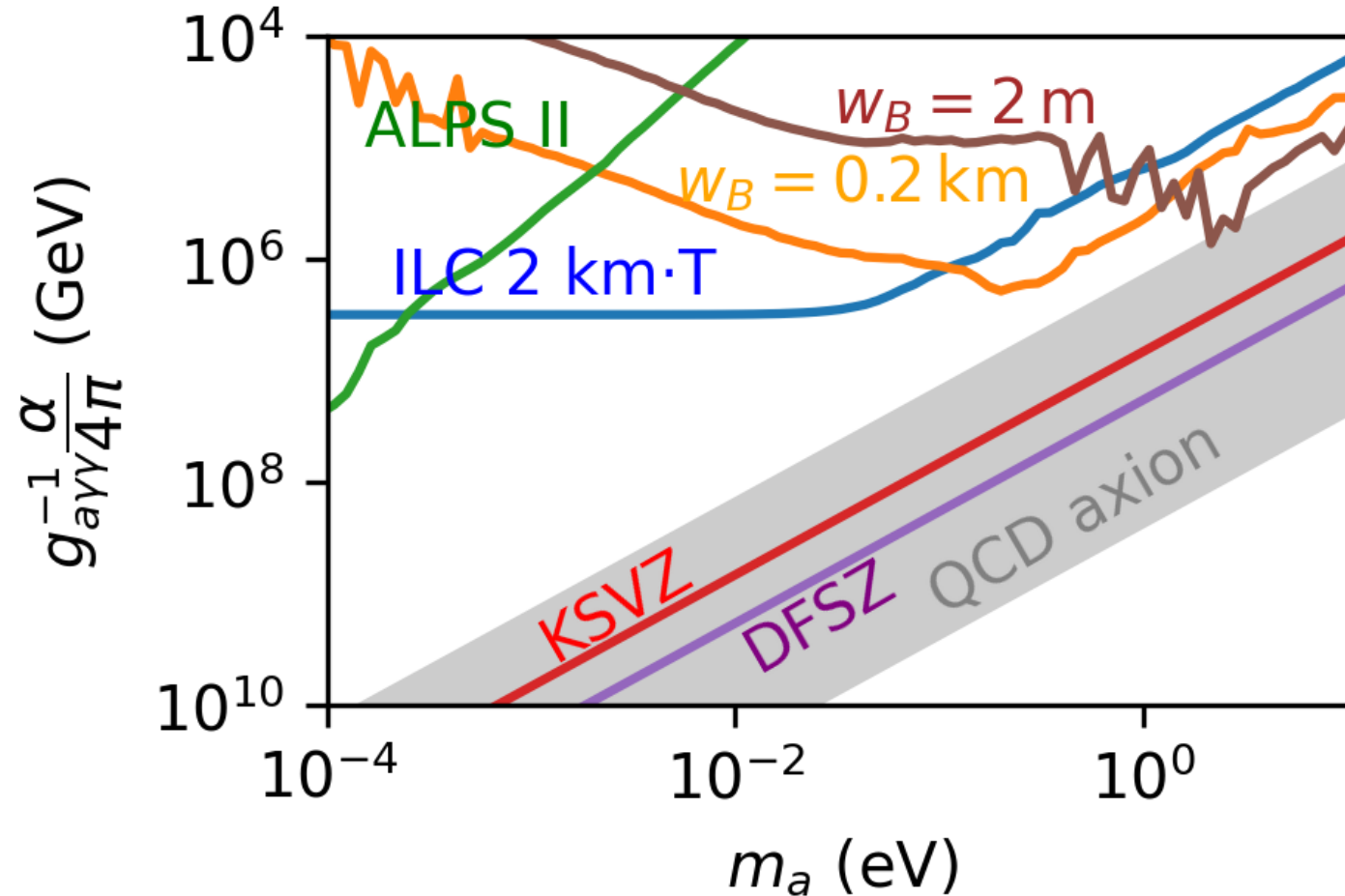
detector

γ a γ

L

Magnetic field configuration

- If we flip magnetic fields over shorter scale w, the "momentum" can be larger
  - But for too short w, the "momentum" is peaked too sharply and the sensitivity is limited
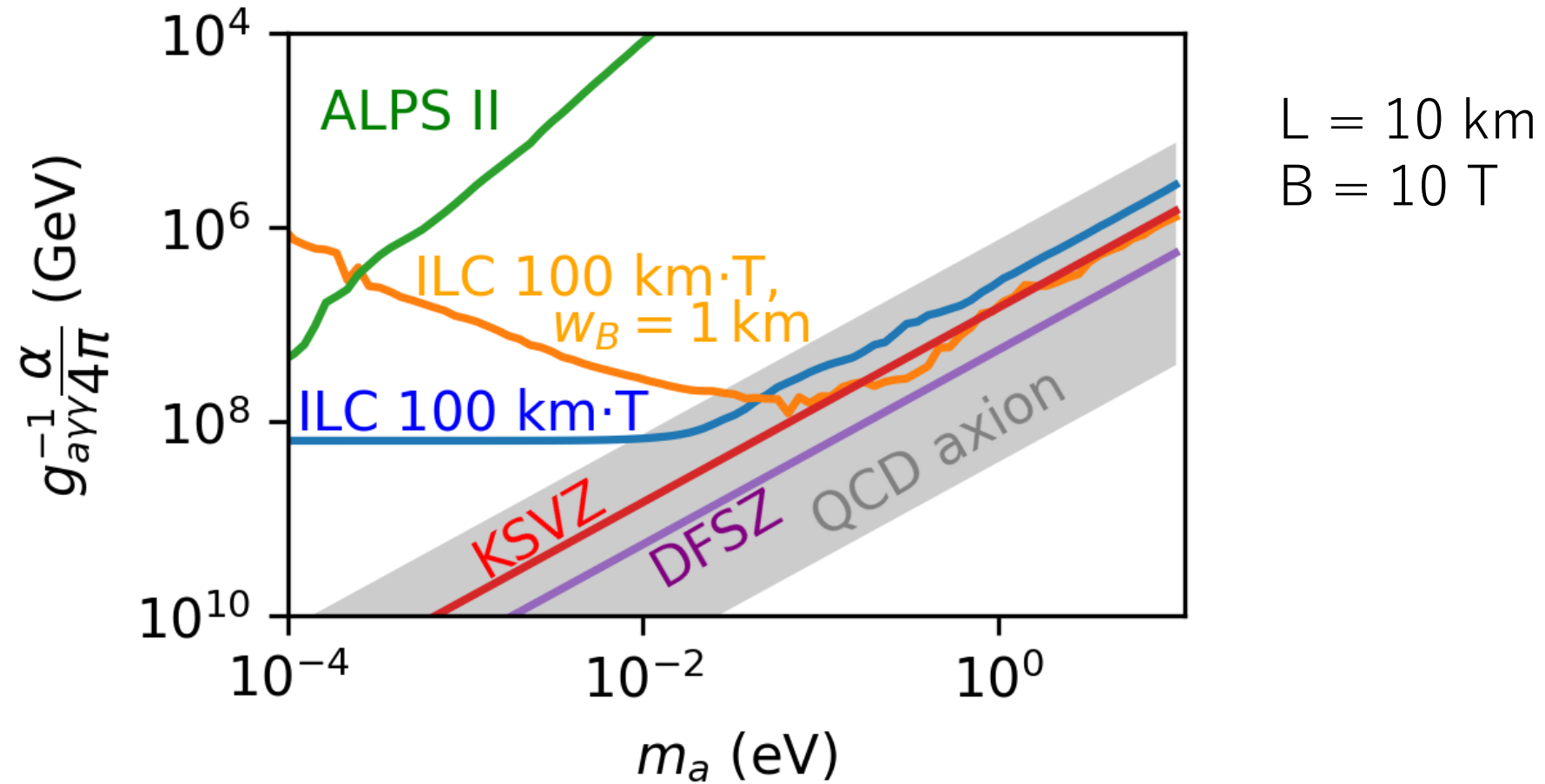
Result (1-year), conservative

L = 2 km
- The distance between γ source and dump in the current design

B = 1 T

Result (1-year)

L = 10 km
B = 10 T

Summary

- It is proposed to generate MeV photon beam as the positron source at the ILC

- The high intensity and large energy are advantageous for light particle search by the LSW experiment

- For the QCD axion, this may constrain the KSVZ model, where no other ground experiments are expected to do.