

The Primitive Domain for Analytic Off-shell Correlators in SFT

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Based on the work

“Analyticity of off-shell Green’s functions in superstring field theory”

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- In complex external momenta variables, Feynman loop diagrams those with no massless internal propagator in closed SFT are analytic on a domain – work of *de Lacroix, Erbin, and Sen (LES)* form 2018. We prove an analytic extension of the LES domain to a larger domain
- Explicit applications of our formula to 2-, 3-, 4-point functions show: the extension equals the well-known **primitive domain**. In case of 5-point functions, we do the analysis for limited subcases

will start by posing a question. then, to answer this, will discuss our proposed extension and its applications

Green's functions in Local QFTs: some old results

Assumptions –

1. Commutator of local operators vanishes outside lightcone,
2. Existence of a complete system of (physical) states with +ve energy

– put constraints on position space correlators

⇒ primitive analyticity of off-shell p -space Green's functions in LQFTs without massless states [BOGOLYUBOV ET AL. book]

Here, off-shell p -space correlators are defined as **Fourier Transforms** of position space correlators

Primitive analyticity

Analyticity of $G(p_1, \dots, p_n)$ on primitive domain in complex external D -momenta

Define, $P_I = \sum_{a \in I} p_a$, $I \subsetneq \{1, \dots, n\} \setminus \emptyset$

A collection of points (p_1, \dots, p_n) satisfying 2 conditions –

C1. external momentum conservation: $p_1 + \dots + p_n = 0$

C2. For each I ,

if $\text{Im } P_I \neq 0$, then $\text{Im } P_I$ must be timelike,

if $\text{Im } P_I = 0$, then $-P_I^2 < M_I^2$, where M_I is the threshold mass for producing any (multi-particle) state in the collision of particles carrying total momentum P_I

– Primitive domain (PD) $\subsetneq \mathbb{C}^{(n-1)D}$

Consequences

PD $\not\subset$ on-shell external momenta

Holomorphic extension of primitive domain includes on-shell external momenta \implies can compute physical amplitudes

Properties of a subregion of PD \implies crossing symmetry of amplitudes for $2 \rightarrow 2$ scattering [BROS, EPSTEIN, GLASER (1965)]

Analytic extension of the same subregion leads to JLD domain in one P_i variable keeping other independent variables real at forward lightcone [BROS, MESSIAH, STORA (1961)]

These results are due to property of PD, irrespective of functions that are analytic on it

Green's functions in closed SFT: recent results

Action directly written in momentum space. Has non-local vertices. No position space construction known

Work with perturbative expansion of $G(\overbrace{p_1, \dots, p_n}^p)$. Consider any Feynman diagram $F(p)$ with n amputated external legs and which $\not\supseteq$ any massless internal propagator

$F(p)$ is analytic on a domain defined as

C1, C2, and C3: each $\text{Im } p_a$ lies on a 2d Lorentzian plane $p^0 - p^1$ if $\text{Im } p_a \neq 0$

– DE LACROIX, ERBIN, SEN (2018)

PD \supseteq LES domain \implies crossing symmetry, JLD domain

Green's functions in closed SFT: recent results

Complex matrix $\tilde{\Lambda}$, s.t. $\tilde{\Lambda}^T \eta \tilde{\Lambda} = \eta = \text{Minkowski metric in } \mathbb{R}^D$

Define its action on p as $\tilde{\Lambda}p \equiv (\tilde{\Lambda}p_1, \dots, \tilde{\Lambda}p_n)$ – complex Lorentz transformation

2 corollaries from the LES paper:

Cor1. $F(p)$ remains analytic at $\tilde{\Lambda}p$, if $p \in \text{LES domain}$

Cor2. at each $p \in \text{LES domain}$, we can allow a small open neighbourhood where $F(p)$ remains analytic

Q. Are $F(p)$ analytic on the full primitive domain?

Outline of our strategy

$F(p)$ is analytic on extended LES domain (LESD+Cor1+Cor2). No further input from SFT

Known that, PD \supset union of a family of tube domains \mathcal{T}_λ , λ : some index

Primitive tube: $\mathcal{T}_\lambda = \mathbb{R}^{(n-1)D} + i\mathcal{C}_\lambda$, \mathcal{C}_λ : cone in $\mathbb{R}^{(n-1)D}$

In \mathcal{T}_λ , identify an open tube (call: LES tube) from (LESD+Cor1+Cor2)

We prove that LES tube admits non-trivial holomorphic extension in \mathcal{T}_λ for $n \geq 3$, due to *Bochner*

possibly obtaining full of \mathcal{T}_λ ? explicitly analyse when $n = 3, 4, 5$

Primitive tube, \mathcal{T}_λ

A specific way to divide space \mathbb{R}^{n-1} into finite no. of cells

– each of which defines a sign-valued map: $\lambda(I)$, $I \subsetneq \{1, \dots, n\} \setminus \emptyset$

$$\begin{aligned} \mathcal{T}_\lambda &= \left\{ p \equiv (p_1, \dots, p_n) : \sum_{a=1}^n p_a = 0, \lambda(I) \text{Im } P_I \in V^+ \quad \forall I \right\} \\ &= \mathbb{R}^{(n-1)D} + i \underbrace{\left\{ \text{Im } p \in \mathbb{R}^{(n-1)D} : \lambda(I) \text{Im } P_I \in V^+ \quad \forall I \right\}}_{\equiv \mathcal{C}_\lambda} \end{aligned} \quad (1)$$

V^+ : open forward lightcone in \mathbb{R}^D ; \mathcal{C}_λ : called the base of \mathcal{T}_λ [ARAKI, BURGOYNE '60; BROS, EPSTEIN, GLASER '64; LASSALLE '74; BROS, LASSALLE '75]

For $n = 3$, # of cells = 6; for $n = 4$ it is 32; for $n = 5$ it is 370, etc

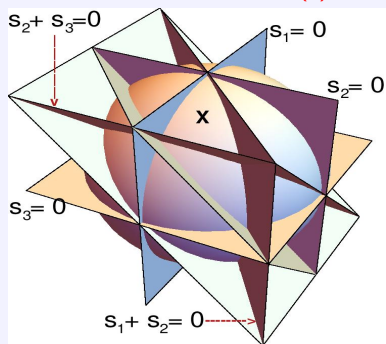
$\lambda(I)$: more details

$$\mathbb{R}^{n-1} = \left\{ s \equiv (s_1, \dots, s_n) : \sum_{a=1}^n s_a = 0 \right\};$$

$$S_I \equiv \sum_{a \in I} s_a, \quad I \subsetneq \{1, \dots, n\} \setminus \emptyset$$

Family of planes $\{S_I = 0\}$ divides \mathbb{R}^{n-1} into open convex cones with common apex at origin. $S_I = 0$, $S_{X \setminus I} = 0$ are same plane. Each such cone is called a cell, γ_λ . Inside a cell each S_I is of definite sign $\lambda(I)$

E.g., for $n = 4$, \mathbb{R}^3 is divided into 32 cells. One such cell has been marked in the figure



LES tube, $\cup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$

$\vec{\theta}$: set of $D - 2$ angles specifying a 2d Lorentzian plane $p^0 - p^{\vec{\theta}}$, $\vec{\theta} = 0$ specifying $p^0 - p^1$ plane

$$\begin{aligned} \mathcal{T}_{\lambda}^{\vec{\theta}} &= \left\{ p \equiv (p_1, \dots, p_n) : \forall a \operatorname{Im} p_a \in p^0 - p^{\vec{\theta}} \text{ plane}; \sum_{a=1}^n p_a = 0, \right. \\ &\quad \left. \text{and } \lambda(l) \operatorname{Im} P_l \in V^+ \quad \forall l \right\} \\ &= \mathbb{R}^{(n-1)D} + i\mathcal{C}_{\lambda}^{\vec{\theta}}, \\ \mathcal{C}_{\lambda}^{\vec{\theta}} &= \left\{ (\operatorname{Im} p_1, \dots, \operatorname{Im} p_n) \text{ on manifold } \sum_{a=1}^n \operatorname{Im} p_a = 0 \text{ such that} \right. \\ &\quad \left. \forall a \operatorname{Im} p_a \in p^0 - p^{\vec{\theta}} \text{ plane and } \lambda(l) \operatorname{Im} P_l \in V^+ \quad \forall l \right\} \end{aligned} \quad (2)$$

$\mathcal{C}_{\lambda}^{\vec{\theta}} \subset \mathbb{R}^{(n-1)D}$. $\mathcal{T}_{\lambda}^{\vec{\theta}}$ can be obtained by acting a real rotation on $\mathcal{T}_{\lambda}^{\vec{\theta}=0}$.
Analyticity of $F(p)$ on $\cup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$ – guaranteed by Cor1

Extension of LES tube

$\bigcup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$ is given by $\mathbb{R}^{(n-1)D} + i(\bigcup_{\vec{\theta}} \mathcal{C}_{\lambda}^{\vec{\theta}}) \subset \mathcal{T}_{\lambda}$

Although each $\mathcal{T}_{\lambda}^{\vec{\theta}}$ is convex, we prove: $\bigcup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$ is non-convex when $n > 2$

Furthermore, we prove: $\bigcup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$ is path-connected

LES tube can be thickened in order to make it open. To holomorphically extend $\bigcup_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}$, we apply Bochner's tube theorem

BOCHNER (1937)

Any open connected tube $\mathbb{R}^m + iA$ has a holomorphic extension to the domain $\mathbb{R}^m + i\text{Ch}(A)$

$\text{Ch}(A)$: smallest convex set containing A , called the convex hull of A

Extension of LES tube

Extension of $\bigcup_{\vec{\theta}} \mathcal{T}_\lambda^{\vec{\theta}}$ to $\mathbb{R}^{(n-1)D} + i\text{Ch}(\bigcup_{\vec{\theta}} \mathcal{C}_\lambda^{\vec{\theta}})$ is non-trivial when $n > 2$

$\mathcal{C}_\lambda \supset \text{Ch}(\bigcup_{\vec{\theta}} \mathcal{C}_\lambda^{\vec{\theta}})$, since \mathcal{C}_λ is a convex cone containing $\bigcup_{\vec{\theta}} \mathcal{C}_\lambda^{\vec{\theta}}$

For 3-, 4-point functions, for each \mathcal{C}_λ above extension yields the full of \mathcal{C}_λ , i.e. $\text{Ch}(\bigcup_{\vec{\theta}} \mathcal{C}_\lambda^{\vec{\theta}}) = \mathcal{C}_\lambda$, and for 5-point functions we obtained subcases in which we are able to prove this equality

For 2-point functions, $\mathcal{T}_\lambda = \bigcup_{\vec{\theta}} \mathcal{T}_\lambda^{\vec{\theta}}$. That is, whenever $p \in \mathcal{T}_\lambda$, $\text{Im } p_1$ ($= -\text{Im } p_2$) lies on some 2d Lorentzian plane

Four-point functions

Primitive tubes, 32 in numbers

$$\mathcal{T}_a^\pm = \left\{ p \in \mathbb{C}^{3D} : \text{Im } p \in \mathcal{C}_a^\pm \right\}, \quad \mathcal{T}_{ab}^\pm = \left\{ p \in \mathbb{C}^{3D} : \text{Im } p \in \mathcal{C}_{ab}^\pm \right\}, \quad (3)$$

$$p = (p_1, \dots, p_4) \text{ linked by } p_1 + \dots + p_4 = 0, \quad 1 \leq a, b \leq 4, \quad a \neq b$$

Conical bases

$$\mathcal{C}_a^+ = -\mathcal{C}_a^- = \left\{ \text{Im } p : \text{Im } p_b, \text{Im } p_c, \text{Im } p_d \in V^+ \right\}$$

$$\mathcal{C}_{ab}^+ = -\mathcal{C}_{ab}^- = \left\{ \text{Im } p : -\text{Im } p_b, \text{Im } (p_b + p_c), \text{Im } (p_b + p_d) \in V^+ \right\} \quad (4)$$

$(abcd)$ = a permutation of (1234)

For each of 32 bases, the no. of $\text{Im } P_i$ that are required to be in specific lightcones = 3. This in turn, determines all other $\text{Im } P_i$ be in specific lightcones for a base

Four-point functions

$$(\mathbf{p}_b, \mathbf{p}_c, \mathbf{p}_d) \longleftrightarrow (-\mathbf{p}_b, \mathbf{p}_b + \mathbf{p}_c, \mathbf{p}_b + \mathbf{p}_d)$$

↖ INVERTIBLE CHANGE OF BASIS

$\text{Im } p_1 + \dots + \text{Im } p_4 = 0 \implies$ each of 32 bases $\subset \mathbb{R}^{3D}$. They can be brought to a common form: $\mathcal{C} = \{\vec{Q} : P_\alpha, P_\beta, P_\gamma \in V^+\}$

$\mathcal{C} \ni \vec{Q} = (P_\alpha, P_\beta, P_\gamma)$ can be written as a $D \times 3$ matrix:

$$\vec{Q} = \begin{pmatrix} P_\alpha^0 & P_\beta^0 & P_\gamma^0 \\ P_\alpha^1 & P_\beta^1 & P_\gamma^1 \\ \vdots & \vdots & \vdots \\ P_\alpha^{D-1} & P_\beta^{D-1} & P_\gamma^{D-1} \end{pmatrix}, \quad (5)$$

where $P_r^0 > +\sqrt{\sum_{i=1}^{D-1} (P_r^i)^2} \quad \forall r = \alpha, \beta, \gamma$

$\text{Ch}(U_{\vec{\theta}} \mathcal{T}_{\lambda}^{\vec{\theta}}) = \mathcal{T}_{\lambda}$ holds

Consider $\vec{Q}_1 \in \mathcal{C}^{\vec{\theta}_1}$, $\vec{Q}_2 \in \mathcal{C}^{\vec{\theta}_2}$, $\vec{Q}_3 \in \mathcal{C}^{\vec{\theta}_3}$ as follows

$$\vec{Q}_1 = 3 \begin{pmatrix} P_{\alpha}^0 - \epsilon & \epsilon/2 & \epsilon/2 \\ P_{\alpha}^1 & 0 & 0 \\ \vdots & \vdots & \vdots \\ P_{\alpha}^{D-1} & 0 & 0 \end{pmatrix}, \quad \vec{Q}_2 = 3 \begin{pmatrix} \epsilon/2 & P_{\beta}^0 - \epsilon & \epsilon/2 \\ 0 & P_{\beta}^1 & 0 \\ \vdots & \vdots & \vdots \\ 0 & P_{\beta}^{D-1} & 0 \end{pmatrix}, \quad (6)$$

$$\vec{Q}_3 = 3 \begin{pmatrix} \epsilon/2 & \epsilon/2 & P_{\gamma}^0 - \epsilon \\ 0 & 0 & P_{\gamma}^1 \\ \vdots & \vdots & \vdots \\ 0 & 0 & P_{\gamma}^{D-1} \end{pmatrix}$$

Choose: $0 < \epsilon < \min \{P_r^0 - \sqrt{\sum_i (P_r^i)^2}, r = \alpha, \beta, \gamma\}$. Implies for each r
 $P_r^0 - \epsilon > +\sqrt{\sum_{i=1}^{D-1} (P_r^i)^2}$. And, we have

$$\frac{\vec{Q}_1}{3} + \frac{\vec{Q}_2}{3} + \frac{\vec{Q}_3}{3} = \vec{Q} \quad (7)$$

Five-point functions, no. of primitive tubes = 370

Bases for 20 of them are: $C_{ab}^{(5)+}$, $C_{ab}^{(5)-} = -C_{ab}^{(5)+}$

$$C_{ab}^{(5)+} = \left\{ \text{Im } p : \text{Im } (p_a + p_c), \text{Im } (p_a + p_d), \text{Im } (p_a + p_e), \right. \\ \left. \text{Im } (p_b + p_c), \text{Im } (p_b + p_d), \text{Im } (p_b + p_e) \in V^+ \right\} \quad (8)$$

$p = (p_1, \dots, p_5)$ linked by $p_1 + \dots + p_5 = 0$

$(abcde) =$ a permutation of (12345)

$C_{ab}^{(5)+} \ni \vec{Q}$ will have 4 independent D -momenta. The no. of $\text{Im } P_i$ that are required to be in specific lightcones = 6

Formula to decompose a general \vec{Q} into terms, each from $\bigcup_{\vec{\theta}} C_{ab}^{(5)+, \vec{\theta}}$ is still missing *ongoing work in progress*

Five-point functions: remaining tubes *work in progress*

– can be brought to the form: $\mathcal{C}^{(5)+} =$

$$\left\{ \vec{Q} = (P_1, P_2, P_3, P_4) : P_1, P_2, P_3, P_4, \underbrace{P_2 + P_4 - P_1}_{P_5}, \underbrace{P_3 + P_4 - P_1}_{P_6} \in V^+ \right\}$$

Here, $P_i - P_j$, $i \neq j$ are timelike/null/spacelike

\vec{Q} : $P_1 - P_4$ is timelike or null. Then \vec{Q} can be decomposed into terms, each of which belongs to the LES tube

\vec{Q} : $P_1 - P_4$ is spacelike. Then go to a Lorentz frame s.t. $P_1^0 - P_4^0 = 0$

Now, $s_r = P_r^0 - \|\mathbf{P}_r\| > 0$, $r = 1, \dots, 6$; $l_1 = \|\mathbf{P}_1 - \mathbf{P}_4\| \geq 0$; and various inequalities among them – allowed primitive space in (s_1, \dots, s_6, l_1)

Various decompositions of \vec{Q} into points from LES tube are found covering various parts of allowed region. Searching left-over points, if any ...

Summary and outlook

Extended the LES domain to a larger subset of the primitive domain, holomorphically. For 2-, 3- and 4-point functions in closed SFT, the full primitive domain is recovered. For 5-point functions, analysed limited subcases. More complete analysis – *work in progress*

Rather dealing case by case for all n , we need to automate the analysis in a computer to carry out the proof for arbitrary n -point functions

Superstring amplitudes might have potential singularities on primitive domain arising from non-perturbative effects. Local QFTs are free from those. Will be interesting to find any differences in context of analyticity

Holomorphic envelope, i.e. largest analytic extension of PD – still remains a challenging open problem since BROS, MESSIAH, STORA (1961)

THANK YOU!