

**IS THE UNDERLYING PHYSICS OF COLLECTIVE FLOW IN SMALL & DILUTE  
“IN ESSENCE” THE SAME AS IN LARGE & DENSE SYSTEMS?**

**“YES” (MAYBE, DON’T KNOW)**

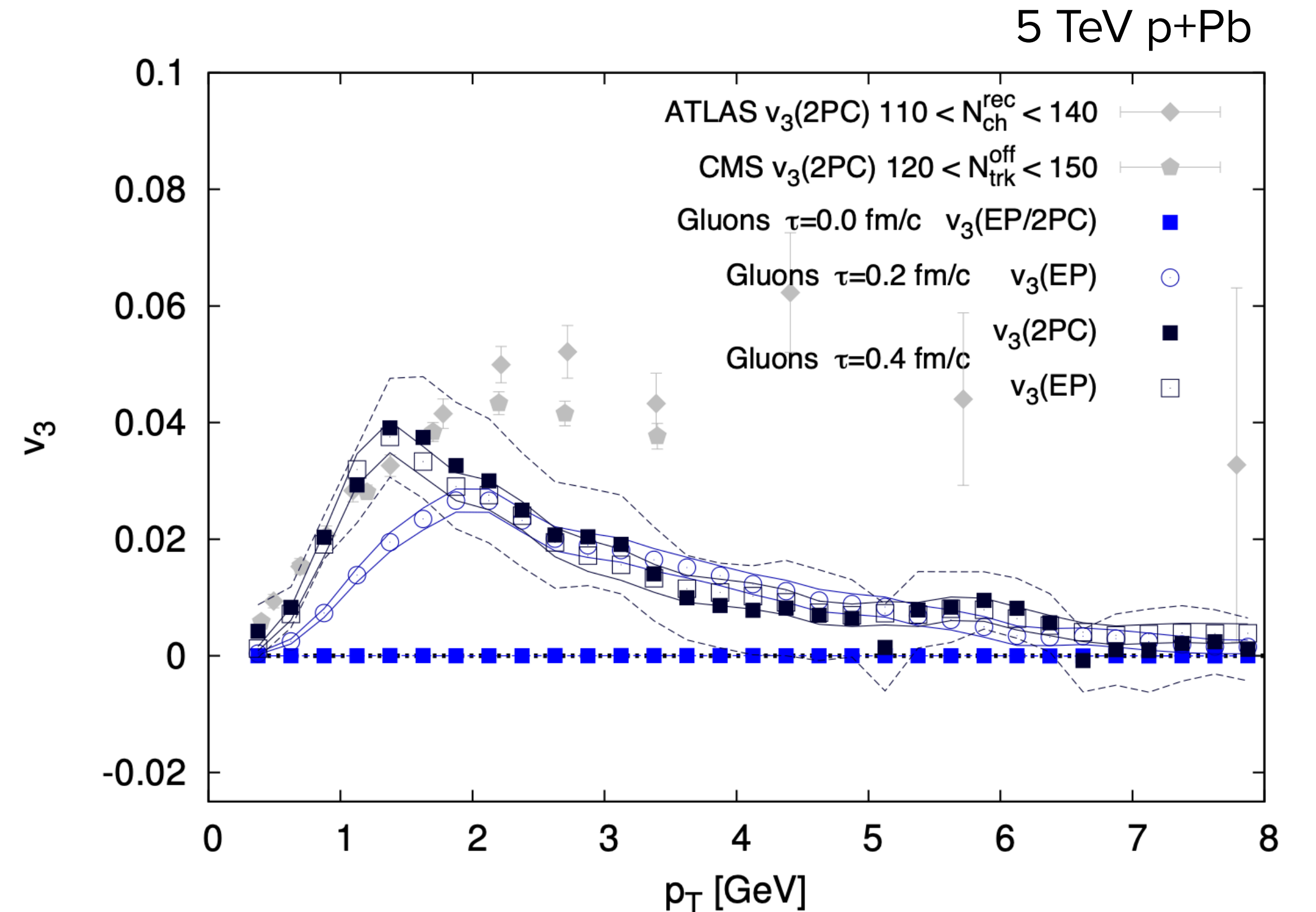
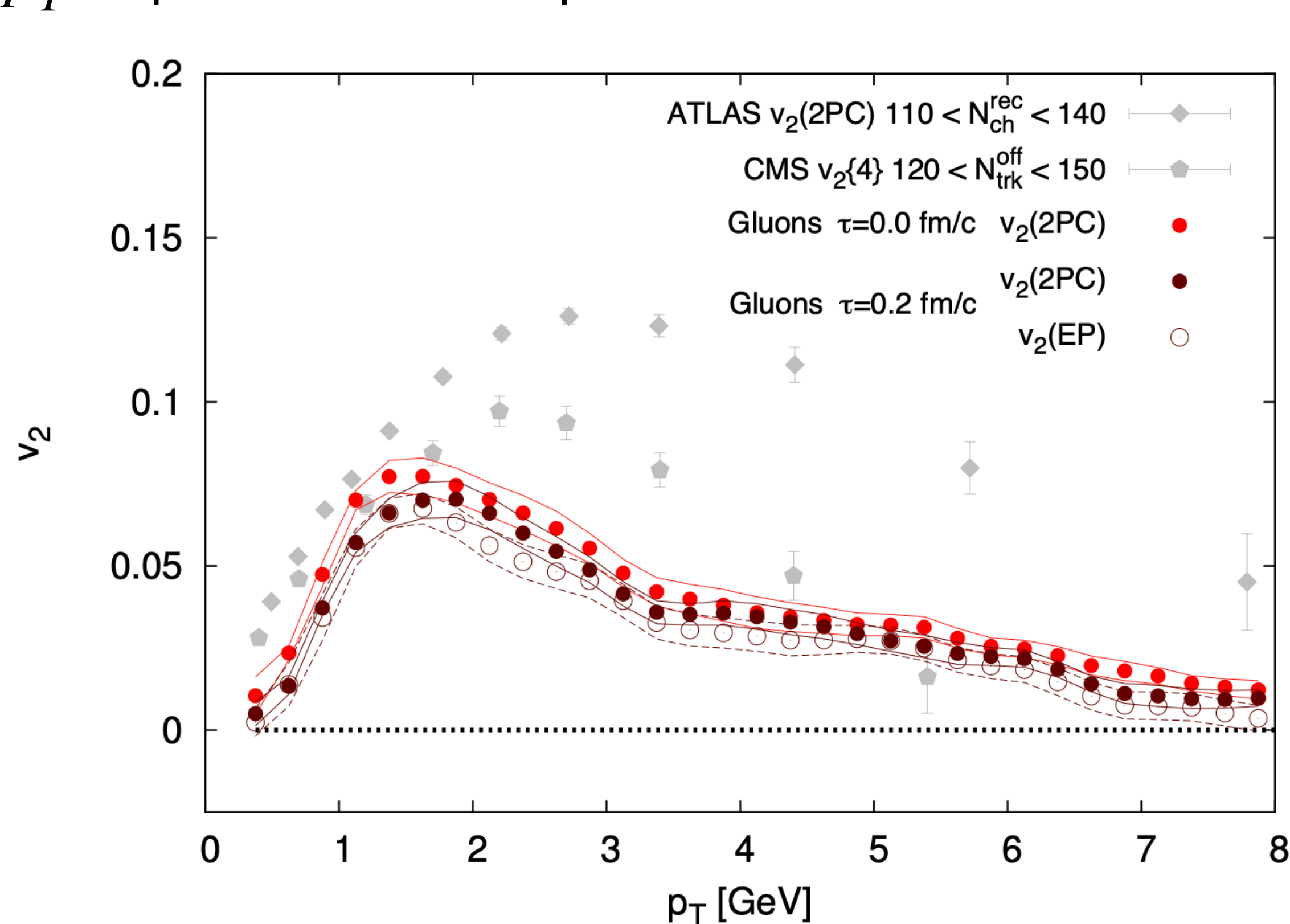
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Initial Stages 2023  
Copenhagen, Denmark  
06/22/2023

# 2015: Boost-invariant dense-dense calculation of initial gluon $v_n$

B. Schenke, S. Schlichting, R. Venugopalan, Phys.Lett.B 747 (2015) 76-82, 1502.01331

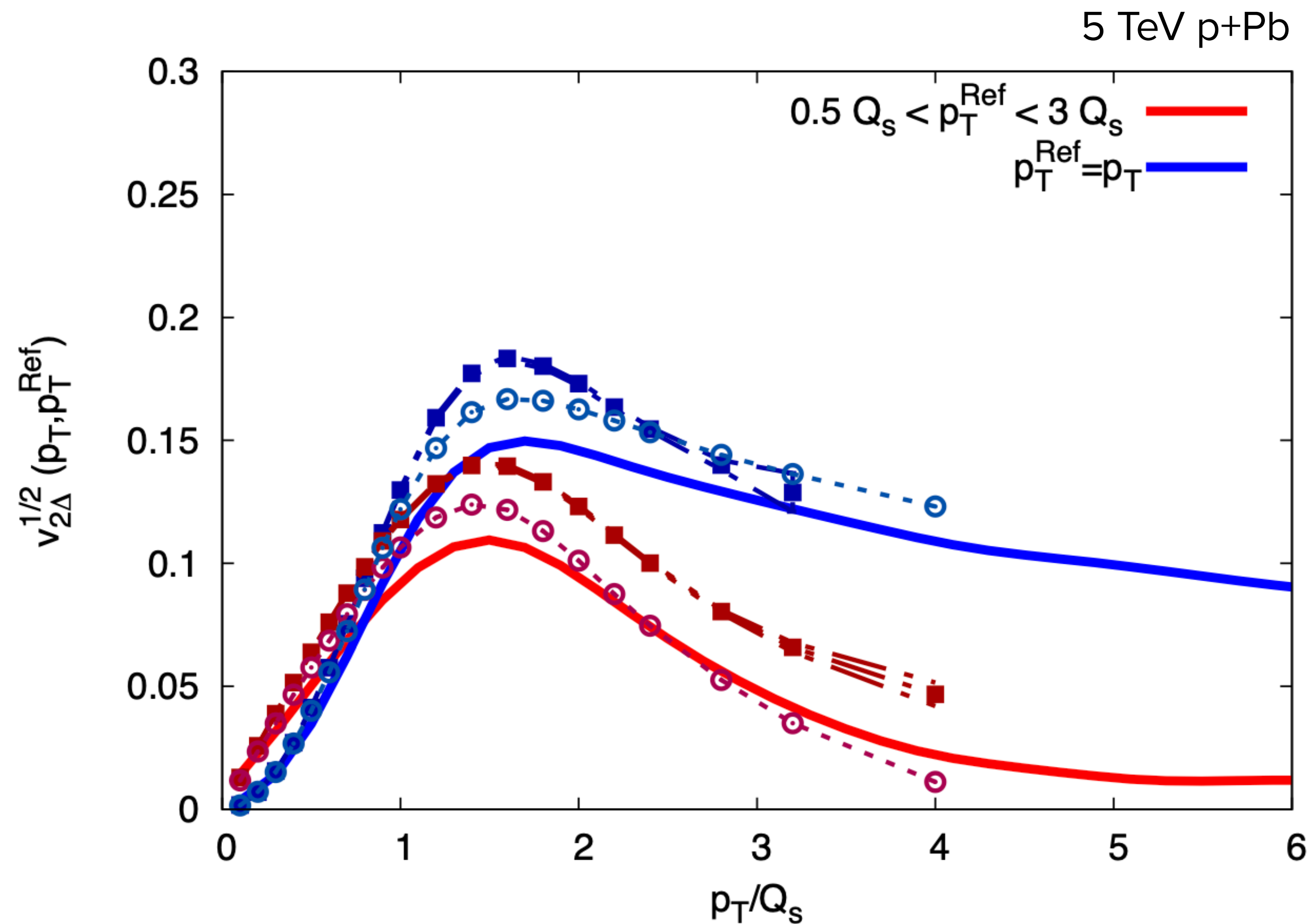
- Compute single and double inclusive gluon distributions in classical Yang-Mills simulations (dense-dense)
- Gluons have a large  $v_2$  in the initial state, while odd harmonics vanish identically at the initial time  $\tau = 0+$ .
- Yang-Mills evolution generates  $v_3$
- $p_T$  dependence not quite like the data



# 2015: Strong dependence on $p_T$ difference in two-particle correlations

T. Lappi, B. Schenke, S. Schlichting, R. Venugopalan., JHEP 01 (2016) 061, 1509.03499

- Decorrelation in  $p_T$  stronger than in the data



- Solid lines: lattice simulation without additional approximation
- Dash-dotted lines (with squares): non-linear Gaussian approximation
- Dotted line (with circles): Glasma graph approximation



# 2022: Rapidity decorrelation of the initial momentum anisotropy

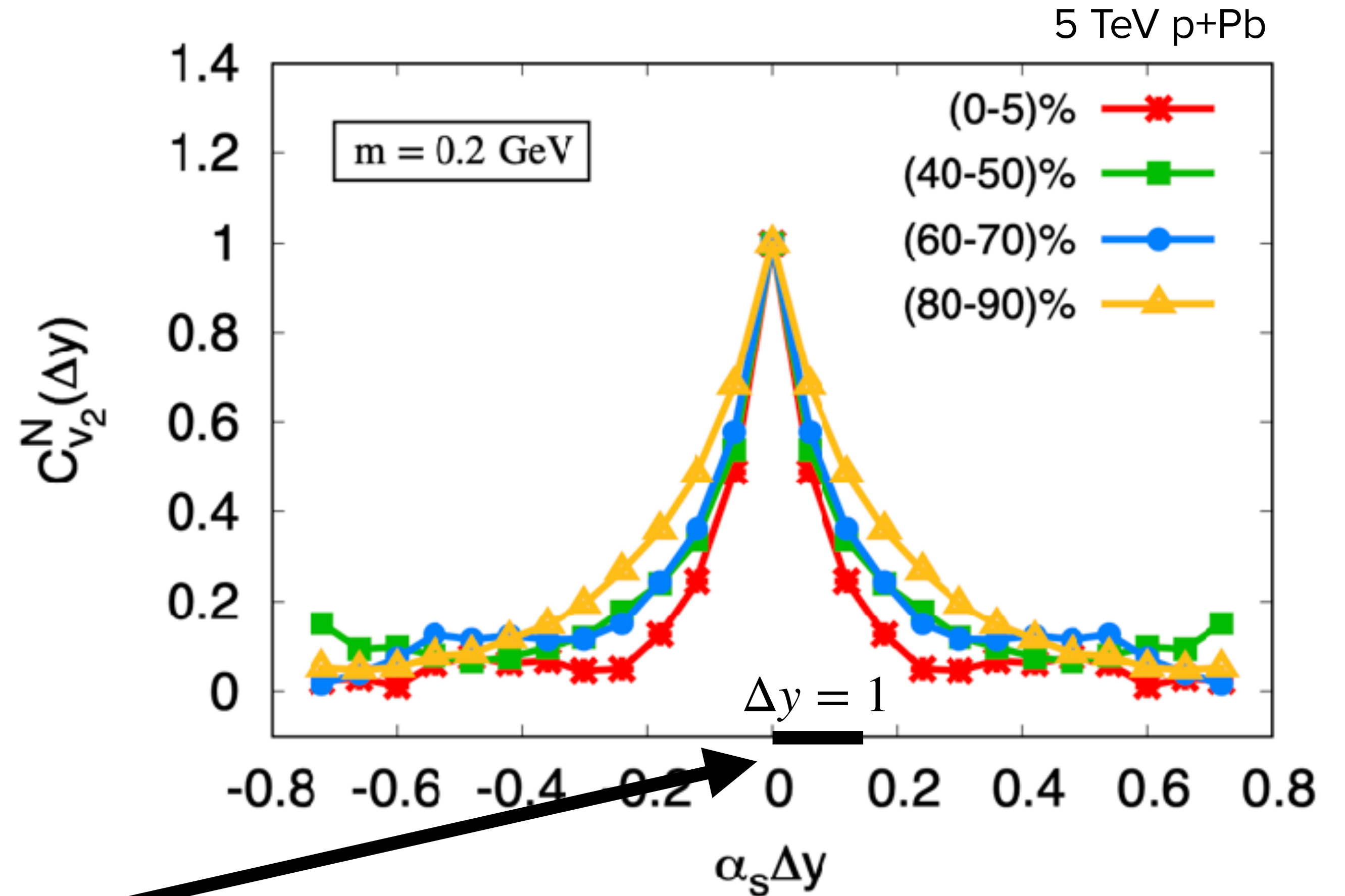
B.Schenke, S. Schlichting, and Pragma Singh, Phys.Rev.D 105 (2022) 9, 094023, e-Print: 2201.08864 [nucl-th]

$$v_2^g(y) = \frac{\int d^2\mathbf{k}_\perp |\mathbf{k}_\perp| \frac{dN_g}{dy d^2\mathbf{k}_\perp}(y) e^{2i\phi_{\mathbf{k}_\perp}}}{\int d^2\mathbf{k}_\perp |\mathbf{k}_\perp| \frac{dN}{dy d^2\mathbf{k}_\perp}(y)}$$

The initial anisotropy of the gluons decorrelates quickly (like typical non-flow)

$$C_{\mathcal{O}}^N(\eta_1, \eta_2) = \frac{\langle \text{Re}(\mathcal{O}(\eta_1) \mathcal{O}^*(\eta_2)) \rangle}{\sqrt{\langle |\mathcal{O}(\eta_1)|^2 \rangle \langle |\mathcal{O}(\eta_2)|^2 \rangle}}$$

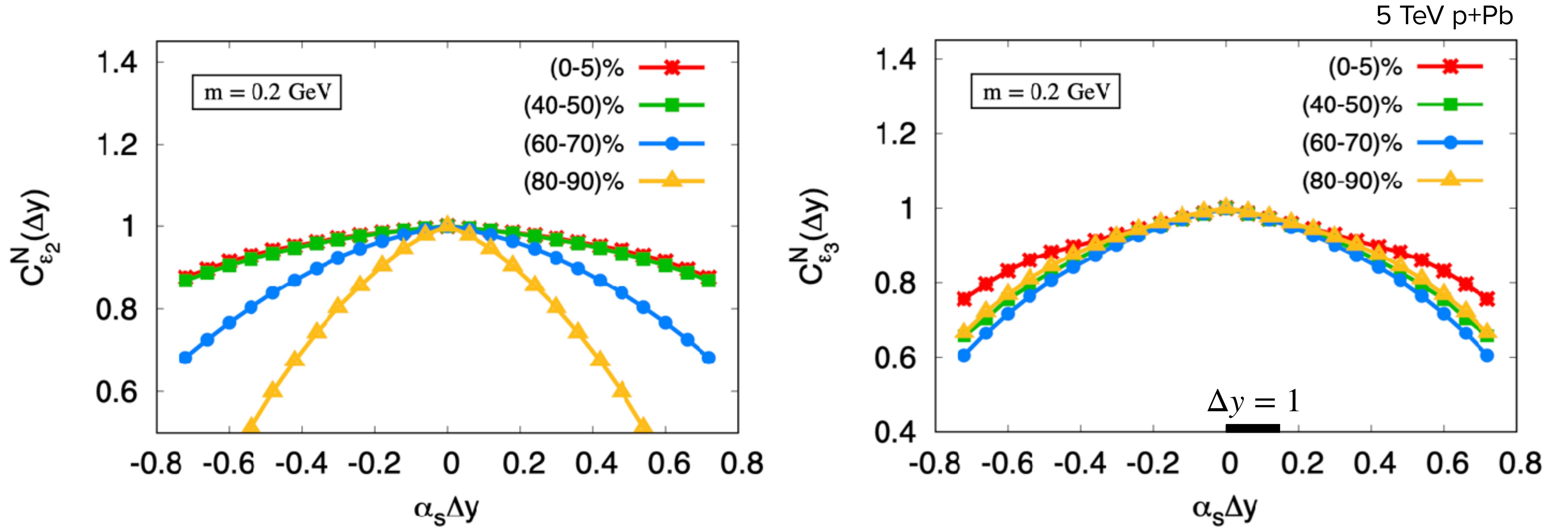
Plug in  $\alpha_s = 0.15$ , which works well for the rapidity dependence of the charged hadron distribution



If experiments employ large rapidity gaps, this effect could remove the initial momentum anisotropy like any other non-flow

# 2022: Rapidity decorrelation of the geometry

B.Schenke, S. Schlichting, P. Singh, Phys.Rev.D 105 (2022) 9, 094023, e-Print: 2201.08864 [nucl-th]



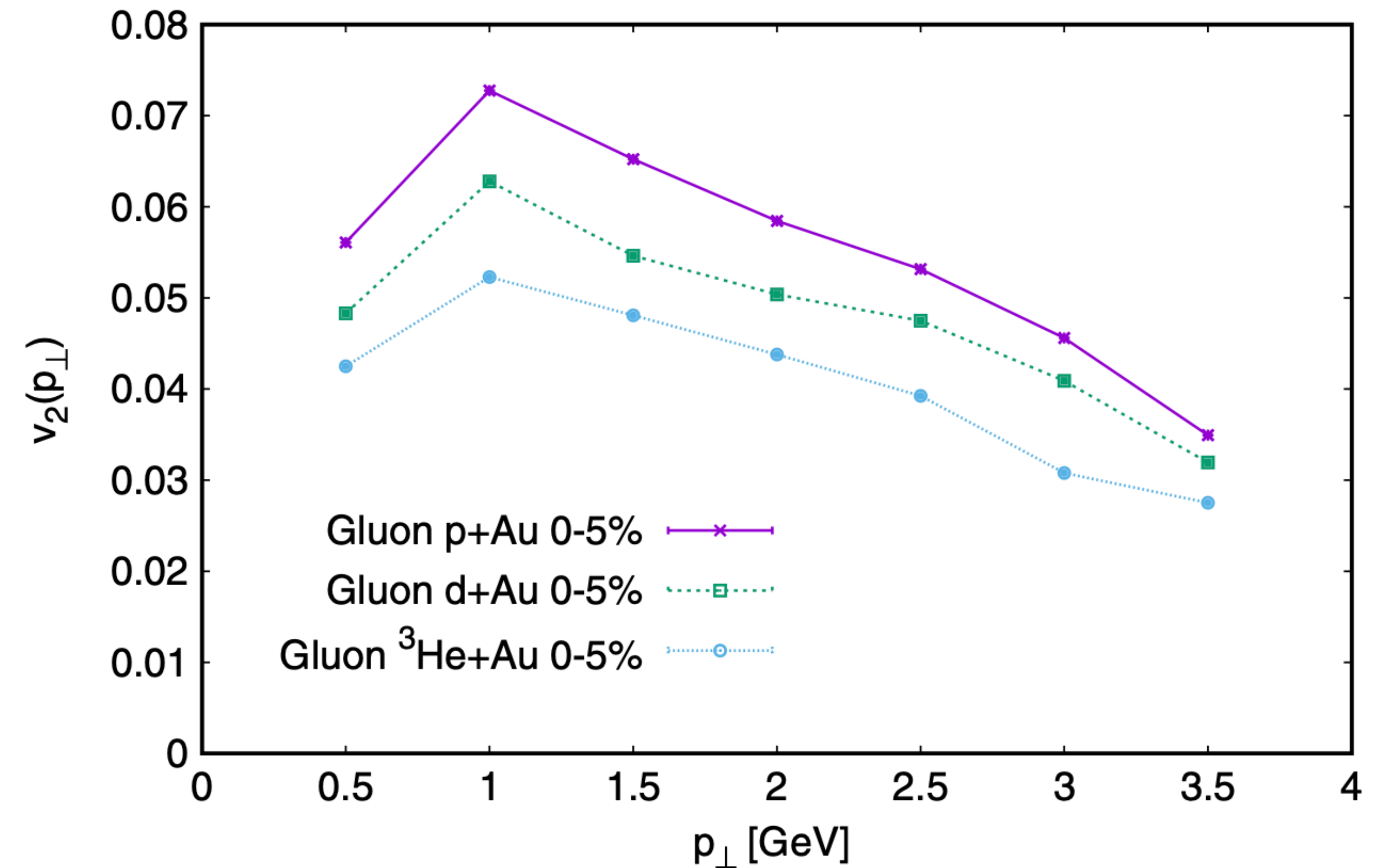
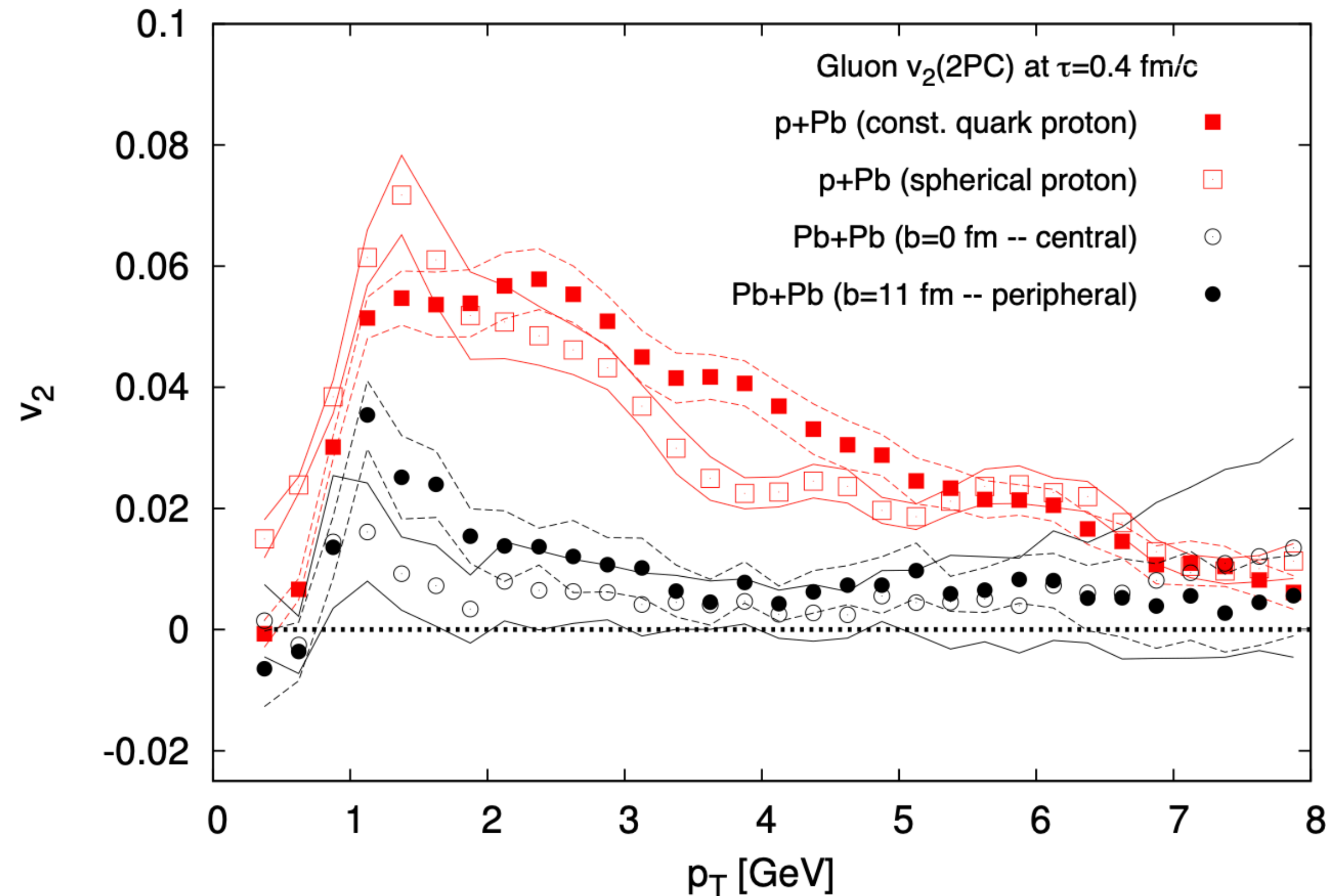
The geometry, quantified here with  $\epsilon_2$  and  $\epsilon_3$ , decorrelates slowly, at least for not too peripheral events

$$C_{\mathcal{O}}^N(\eta_1, \eta_2) = \frac{\langle \text{Re}(\mathcal{O}(\eta_1) \mathcal{O}^*(\eta_2)) \rangle}{\sqrt{\langle |\mathcal{O}(\eta_1)|^2 \rangle \langle |\mathcal{O}(\eta_2)|^2 \rangle}}$$

# Systematics with system size

B. Schenke, S. Schlichting, R. Venugopalan, Phys.Lett.B 747 (2015) 76-82, 1502.01331

M. Mace, V. V. Skokov, P. Tribedy, R. Venugopalan, Phys. Rev. Lett. 121, 052301 (2018), PRL123, 039901(E) (2019)



- Trend with system size is opposite to the data (when using large rapidity gaps)



# Discussion

- CGC initial state momentum distributions:  
**Systematics** with system size and  $p_T$  **do not agree with data**
- Full 3+1D Yang Mills will evolve towards isotropization at early times  
This would reduce the initial anisotropy even more
- At small  $p_T \lesssim Q_s$  and large rapidity separation, **we will not observe a momentum anisotropy that is not driven by geometry and strong final state effects**
- **Short range rapidity correlations** may be there but mixed with other non-flow
- Note: Remember that the CGC does also predicts how energy and momentum is deposited. Along with sampled nucleon positions, it provides the initial geometry that, in combination with hydrodynamics, successfully describes a lot of data
- Hydrodynamics has issues in small systems - very large viscous corrections/problems with causality

# BACKUP



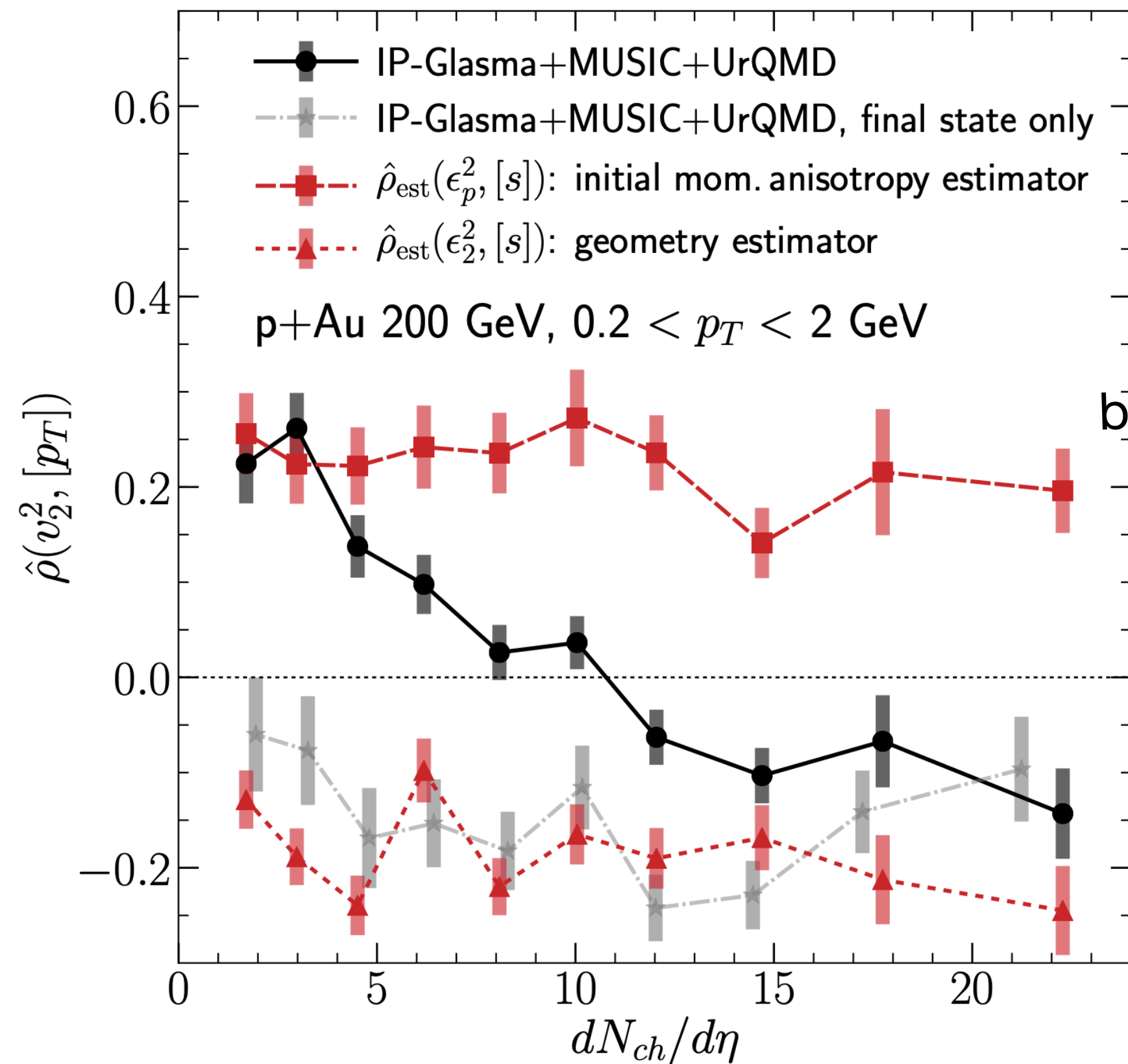
# Initial pressure anisotropy

B.Schenke, S. Schlichting, and Pragma Singh, Phys.Rev.D 105 (2022) 9, 094023, e-Print: 2201.08864 [nucl-th]

$$\varepsilon_p(y) = \frac{\int d^2\mathbf{r}_\perp T^{xx}(y, \mathbf{r}_\perp) - T^{yy}(y, \mathbf{r}_\perp) + 2iT^{xy}(y, \mathbf{r}_\perp)}{\int d^2\mathbf{r}_\perp T^{xx}(y, \mathbf{r}_\perp) + T^{yy}(y, \mathbf{r}_\perp)}$$

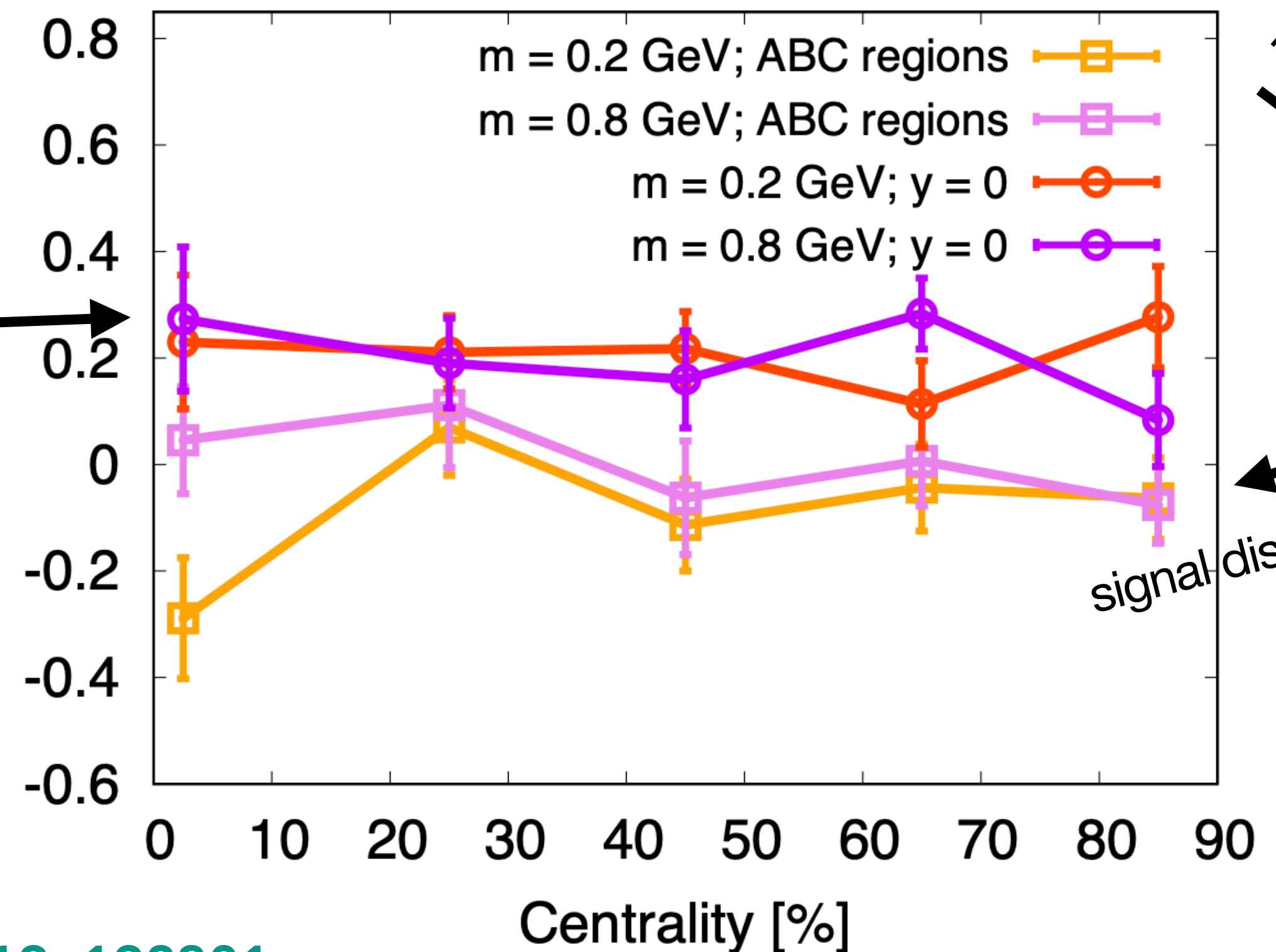
in the quasi-particle picture this is equivalent to the  $v_2^g$  definition

This anisotropy is present in the initial condition for hydrodynamics  $T^{\mu\nu}$ :



But:

boost invariant  $\hat{\rho}_{\text{est}}(\epsilon_p^2, [s])$



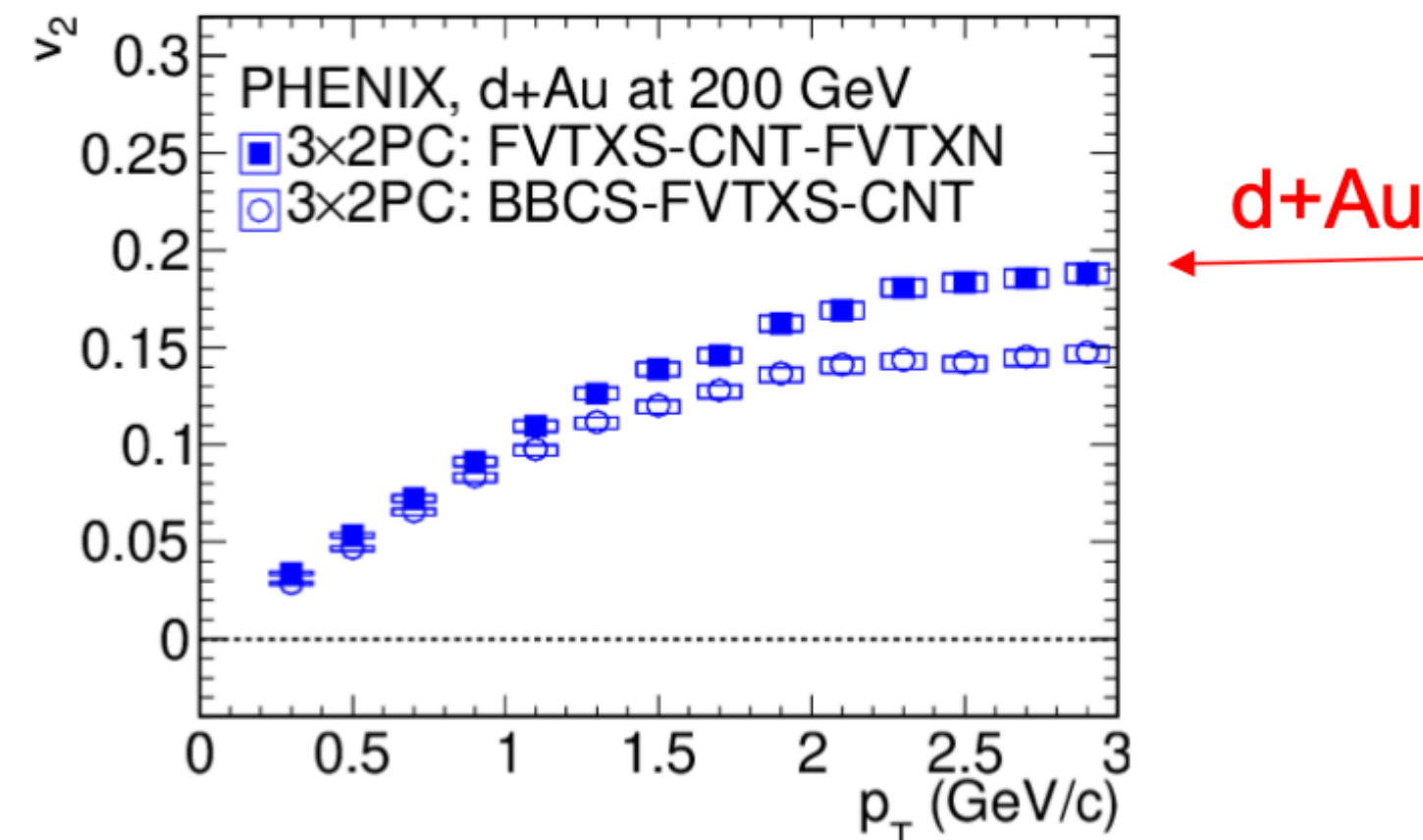
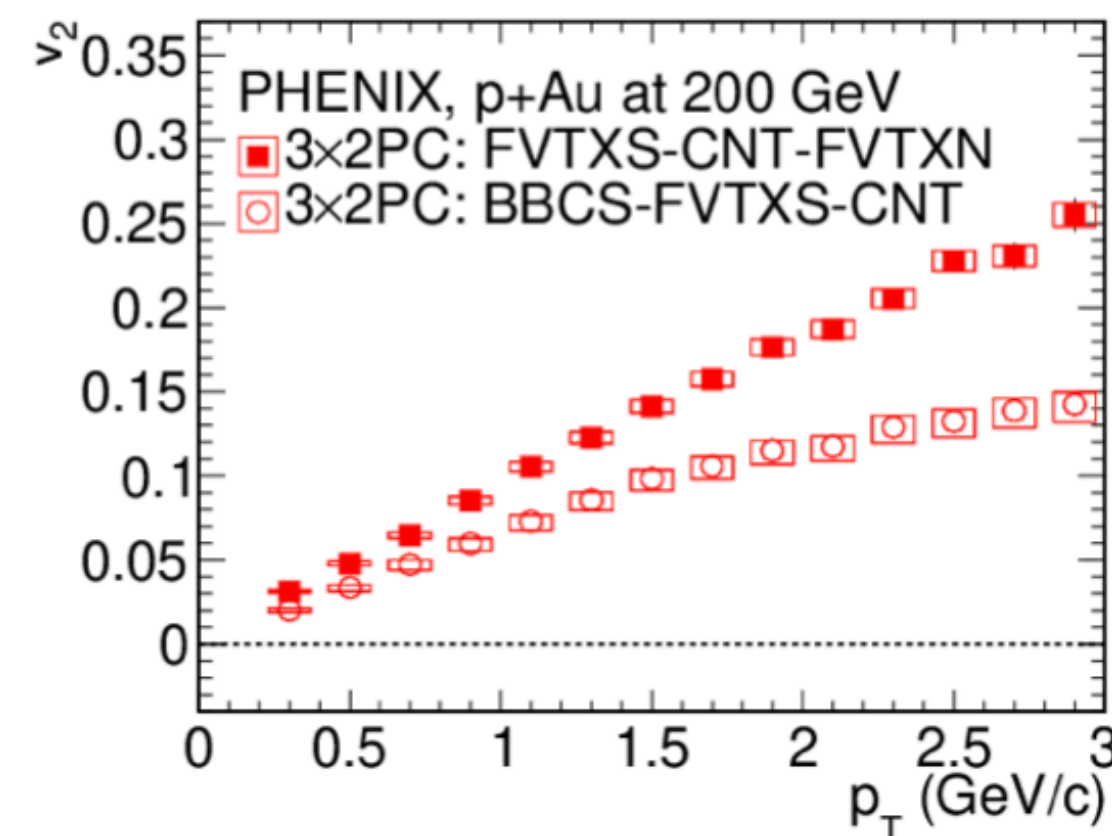
experimental rapidity regions

signal disappears

# Non-flow $v_n$ at forward rapidities

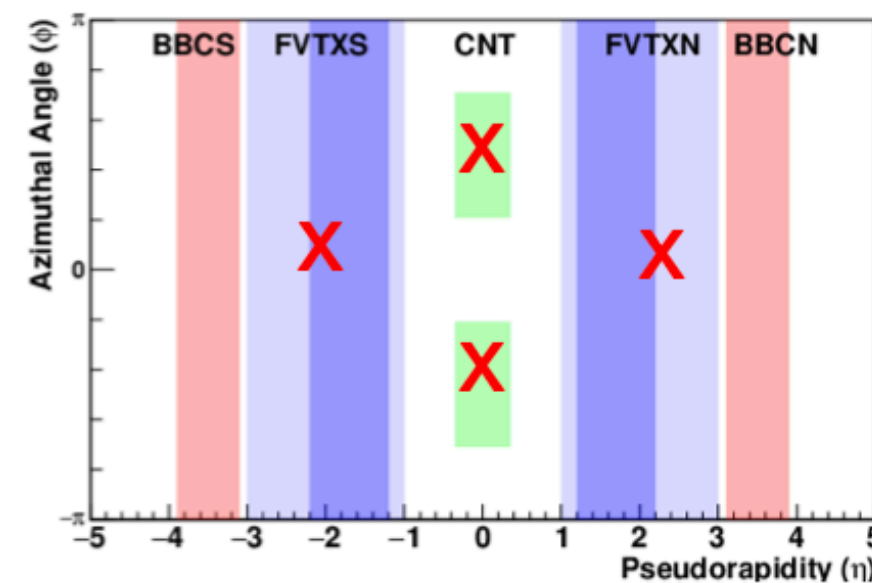
Reverse ordering for  $v_2$

$$v_2(p+Au) > v_2(d+Au)$$



not consistent with hydro/final-state effect

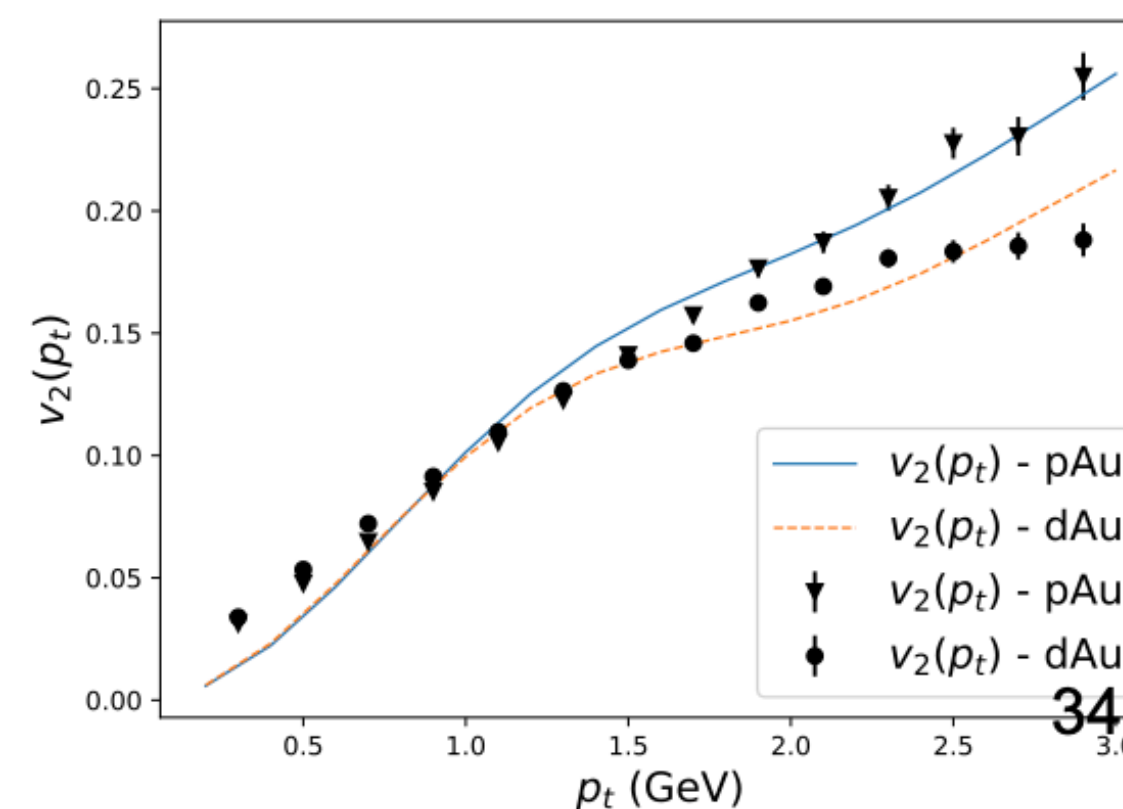
PHENIX collaboration (2022)



PHENIX sub-detectors used

can be reproduced  
with simple model

Kohara and CM  
to appear



But this could be other  
non-flow as well  
should check PYTHIA