

In-medium parton showers with overlapping emissions

Peter Arnold
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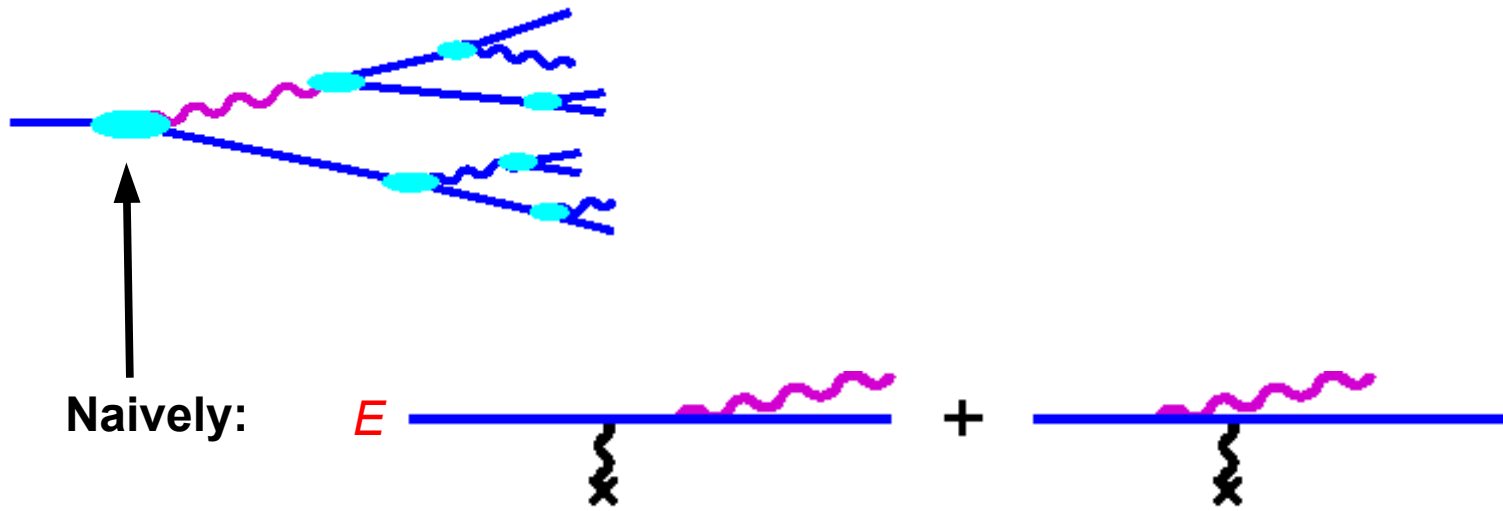
Reporting (eventually) on recent work with



Shahin Iqbal Omar Elgedawy

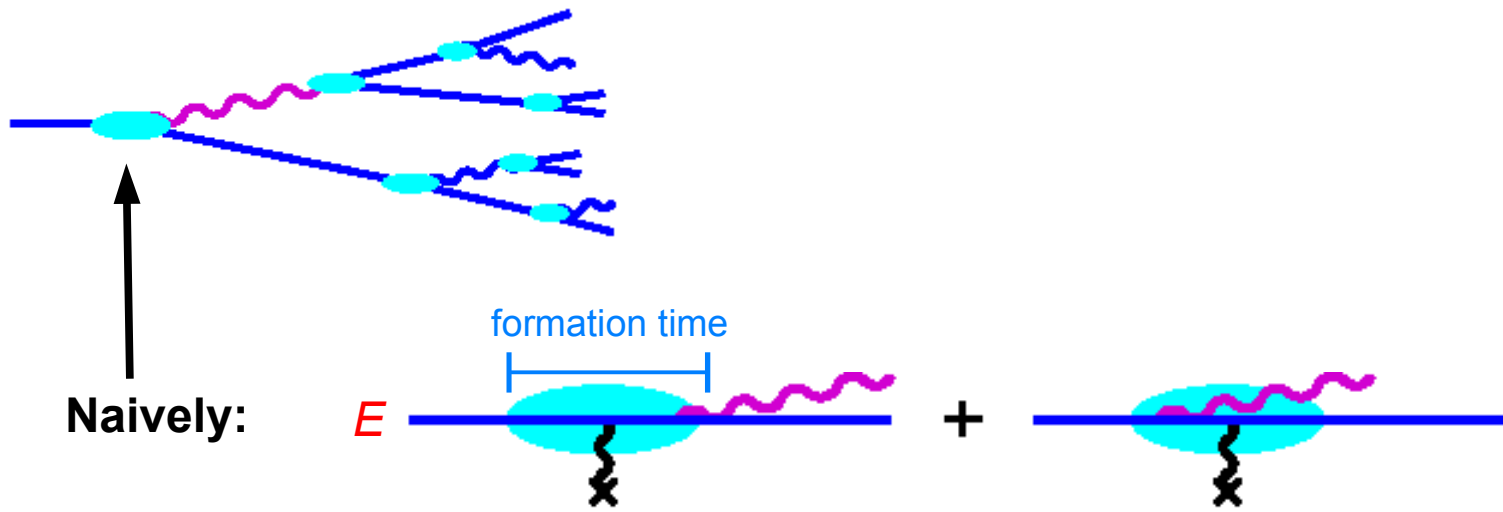
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Medium-induced showering



Prob. of brem $\sim \alpha$ per collision with medium
(up to logs)

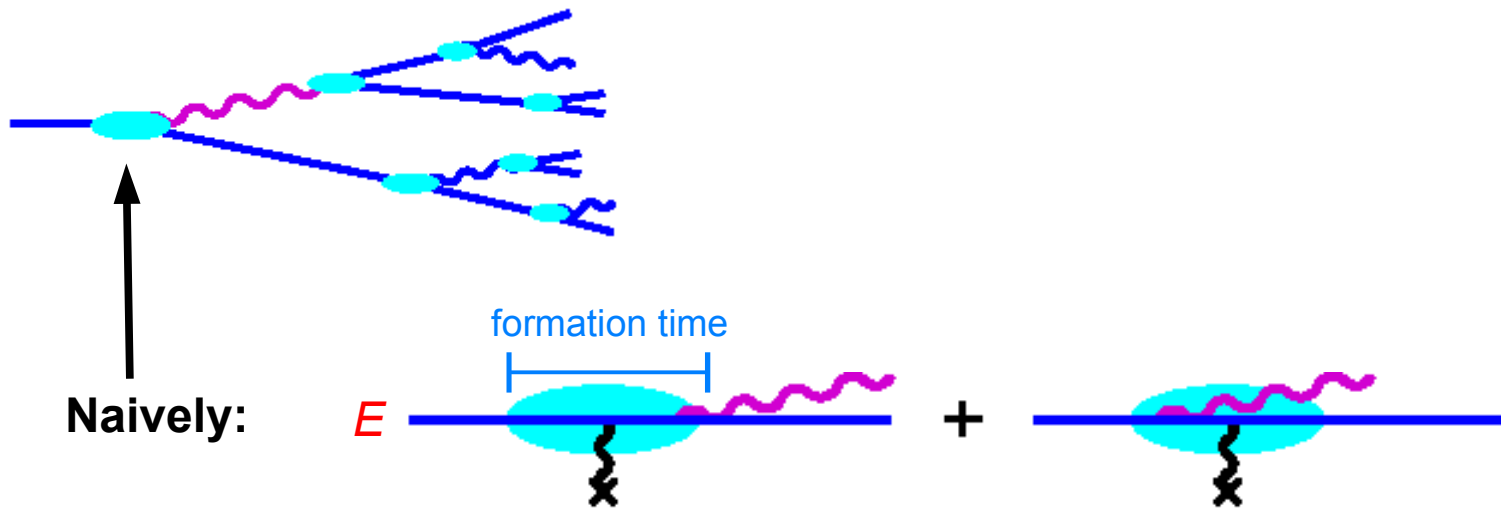
Medium-induced showering



Formation time means quantum duration of splitting process.

Formation time **grows** with energy E .

Medium-induced showering

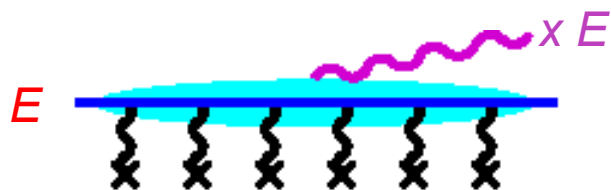


Formation time means quantum duration of splitting process.

Formation time **grows** with energy E .

LPM Effect:

What happens when formation time \gg mean free time between collisions w/ medium?



Prob. of brem $\sim \alpha$ per formation time

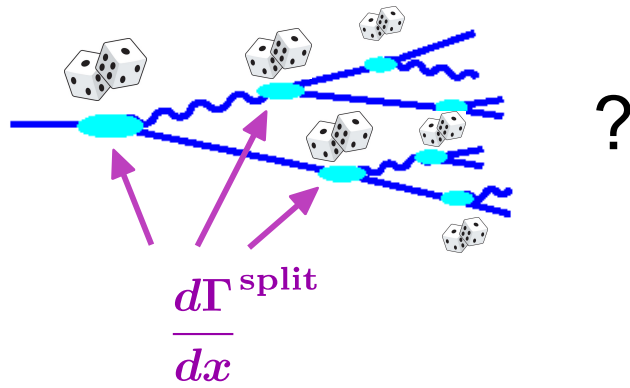
QED (1950s): LPM [Landau-Pomeranchuk & Migdal]

QCD (1990s): BDMPS-Z + many later variations



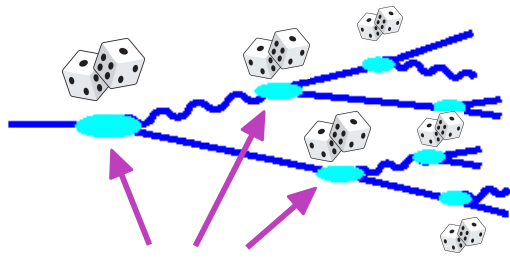
calculation of splitting rates $\frac{d\Gamma^{\text{split}}}{dx}$

Can we then describe in-medium shower development by



(LPM splitting rates)

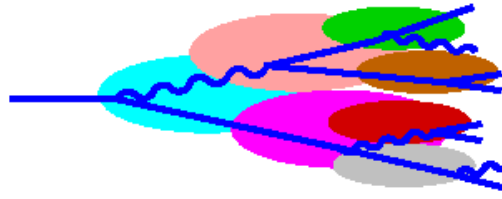
Or can splittings overlap?



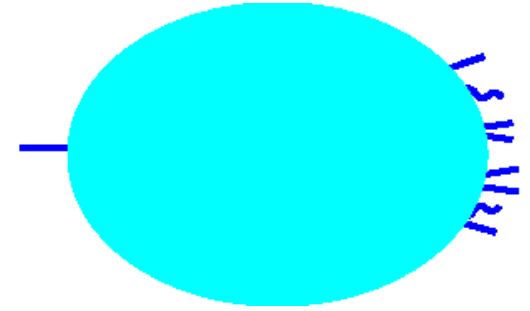
$$\frac{d\Gamma^{\text{split}}}{dx}$$

(LPM splitting rates)

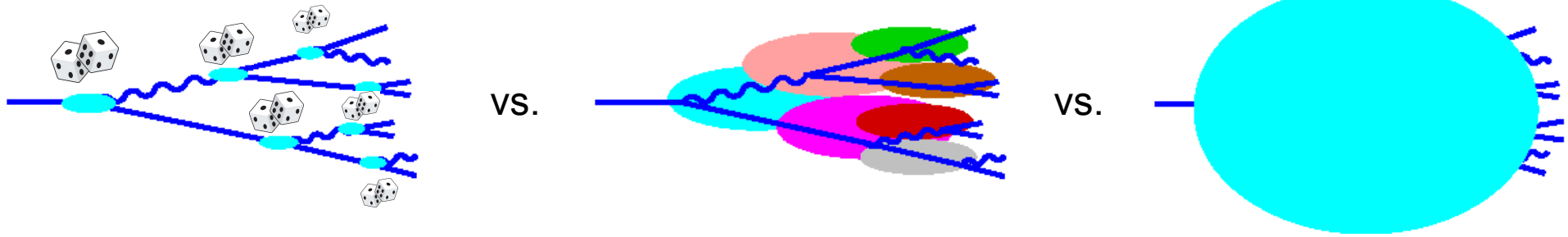
vs.



vs.



Or can splittings overlap?



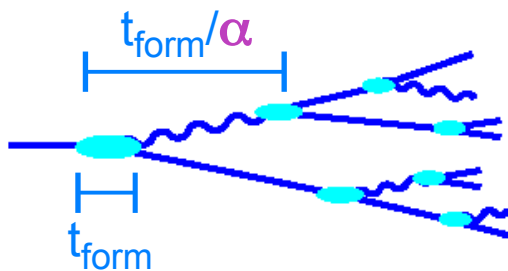
Prob. of brem $\sim \alpha(\mu)$ per formation time

→ Prob. two consecutive splittings overlap $\sim \alpha(\mu)$

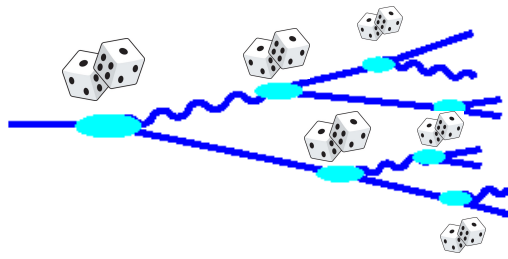
↑
a scale depending on energy
Not the $\alpha(T)$ of the QGP

All depends on how big $\alpha(\mu)$ is!

For small α , there is a hierarchy of scales that (typically) separates the splittings:



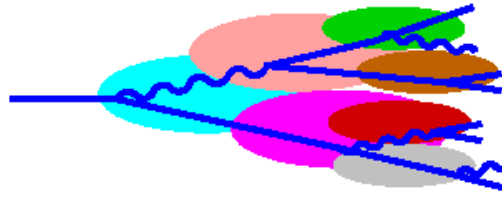
Summary so far



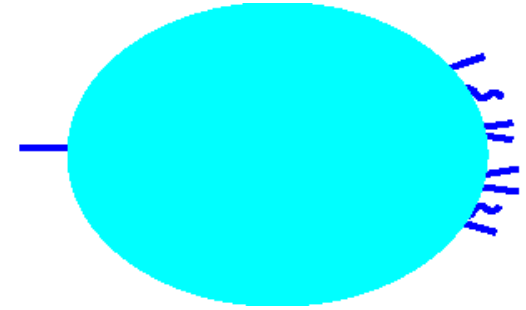
$\alpha_s(\mu)$ small

a “standard” picture
of a shower

vs.



vs.

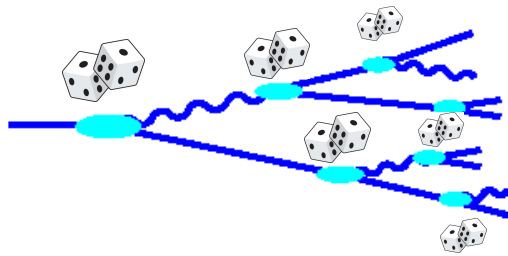


$\alpha_s(\mu)$ big

HELP!

Turn to AdS/CFT for
qualitative insight

How do we tell if



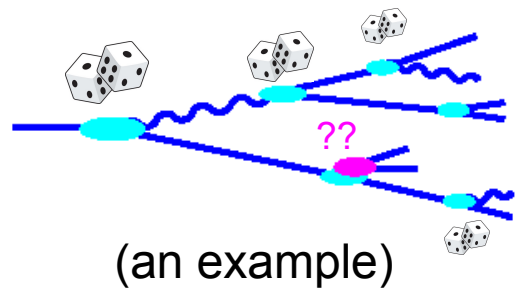
is a good or bad picture for reasonable values of $\alpha_s(\mu)$?

Two approaches

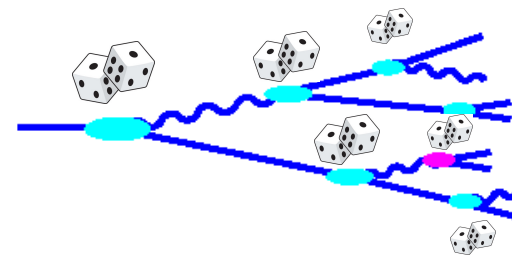
- (1) EXTERNAL VALIDATION: Confront w/ experiment.
But.... many confounding factors.
- (2) INTERNAL CONSISTENCY: Test with theory!

Question:

Are the first corrections



to



small for reasonable values of $\alpha_s(\mu)$?

Perks for theorists:

- May avoid confounding factors by testing in simplified situations.
- Can test on simple shower characteristics not accessible to experiment.

So...

A theorist thought experiment

Simplifying assumptions

- Treat elastic scattering w/ medium in the \hat{q} approximation:

$$\langle (\text{change in } p_{\perp})^2 \rangle = \hat{q} \cdot (\text{distance traveled})$$

- A static, homogeneous, “infinite”-size QGP

I can now reveal that scale for $\alpha_s(\mu)$ is
and formation times are

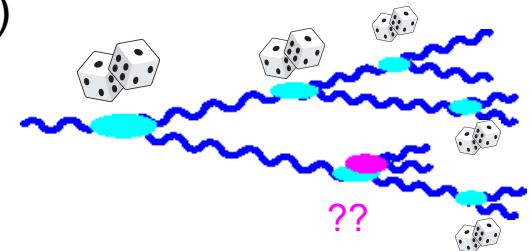
$$\mu \sim (\hat{q}E)^{1/4}$$

$$t_{\text{form}} \sim \sqrt{E/\hat{q}}$$

- Start with a parton that is (approx.) on-shell.
- Study gluon-initiated showers in large- N_c limit (w/ N_f fixed)



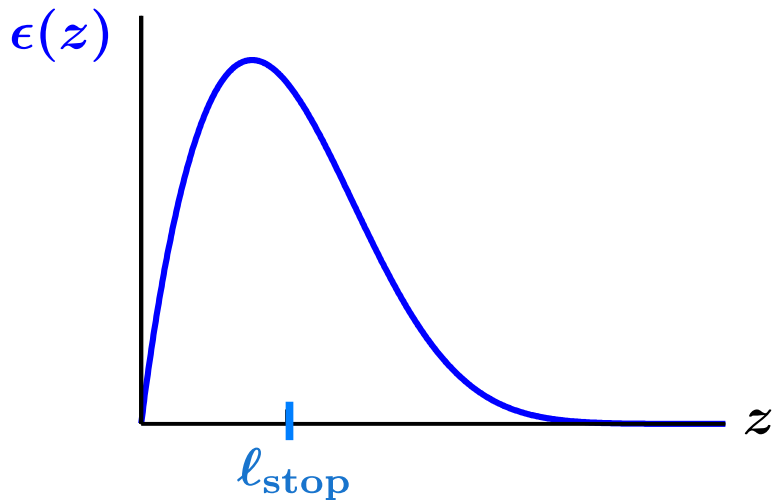
Only $g \rightarrow gg$ splittings consider (so far!)



A theorist thought experiment

Something theorists could “observe”:

(statistically averaged) distribution of energy deposited by shower as a function of distance z



$l_{\text{stop}} \equiv \langle z \rangle$ (1st moment of energy deposition distribution)

$$l_{\text{stop}} \sim \frac{t_{\text{form}}}{\alpha} \sim \frac{1}{\alpha} \sqrt{\frac{E}{\hat{q}}}$$

Note: l_{stop} depends on \hat{q}

How big are the overlap corrections to $\epsilon(z)$?

Answer: **BIG!** ... which has been known since

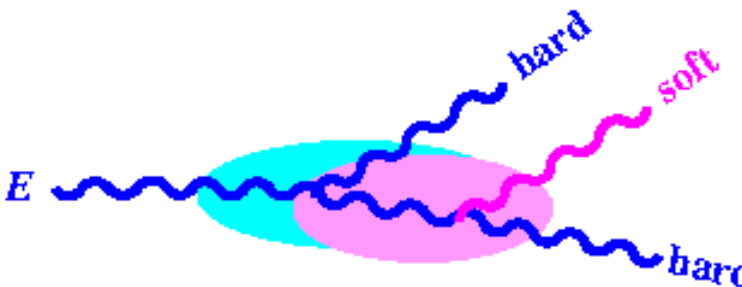
Iancu (2014)

Blaizot and Mehtar-Tani (2014)

Wu (2014)

[building on radiative corrections to \hat{q} found by Liou, Mueller, Wu (2013)]

(1) BIG because there is a double-log enhancement coming from **SOFT** radiation:

Prob. of overlap  suppressed by (in my application)

$\alpha_s(\mu) \ln^2\left(\frac{E}{T}\right) \Rightarrow$ BIG result for large E

(2) But these BIG soft-radiation effects can be absorbed into an effective value of \hat{q} :

$$\hat{q} \longrightarrow \hat{q}_{\text{eff}}(E) = \hat{q} \left[1 + \# \alpha_s \ln^2\left(\frac{E}{T}\right) \right]$$

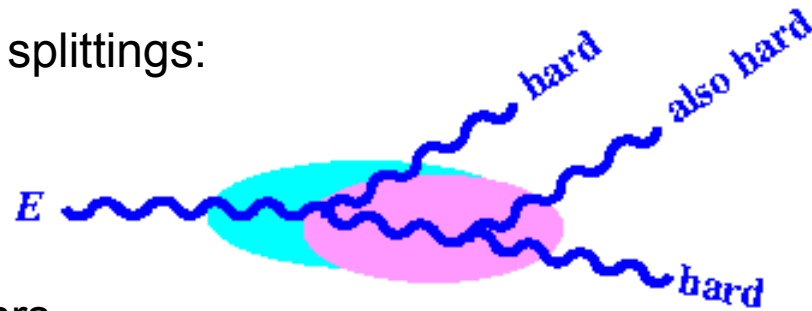
Can even be re-summed at leading log to all orders in α_s

A REFINED QUESTION

How big are overlap effects that cannot be absorbed in \hat{q} ?

(1) Need to calculate overlap of two hard splittings:

Extremely difficult calculation.



After lots of QFT and many (!!) years ...

Completed (for gluons) in 2022 with S. Iqbal and



Tyler Gorda

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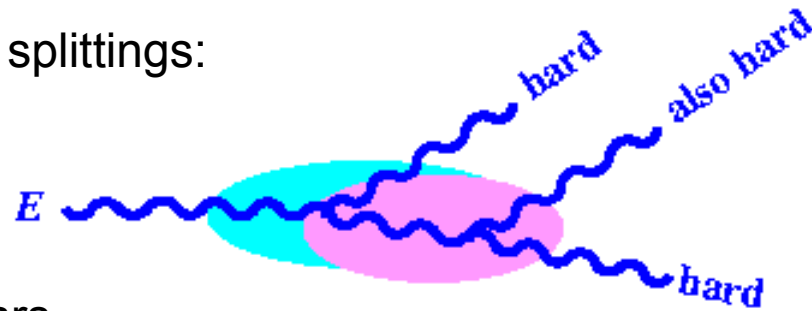
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Technical note

The drawing above is short-hand for what we call

 $\Delta \frac{d\Gamma}{dx dy} \equiv$ the overlap correction to two independent splittings

$$= \left[\left\langle \left| \int_0^\infty d(\Delta t) \left[\text{full calculation of double splitting rate} \right] + \dots \right|^2 \right\rangle_{\text{medium avg}} \right] - \left[\frac{d\Gamma^{\text{split}}}{dx} \text{ and } \frac{d\Gamma^{\text{split}}}{dy} \right]$$

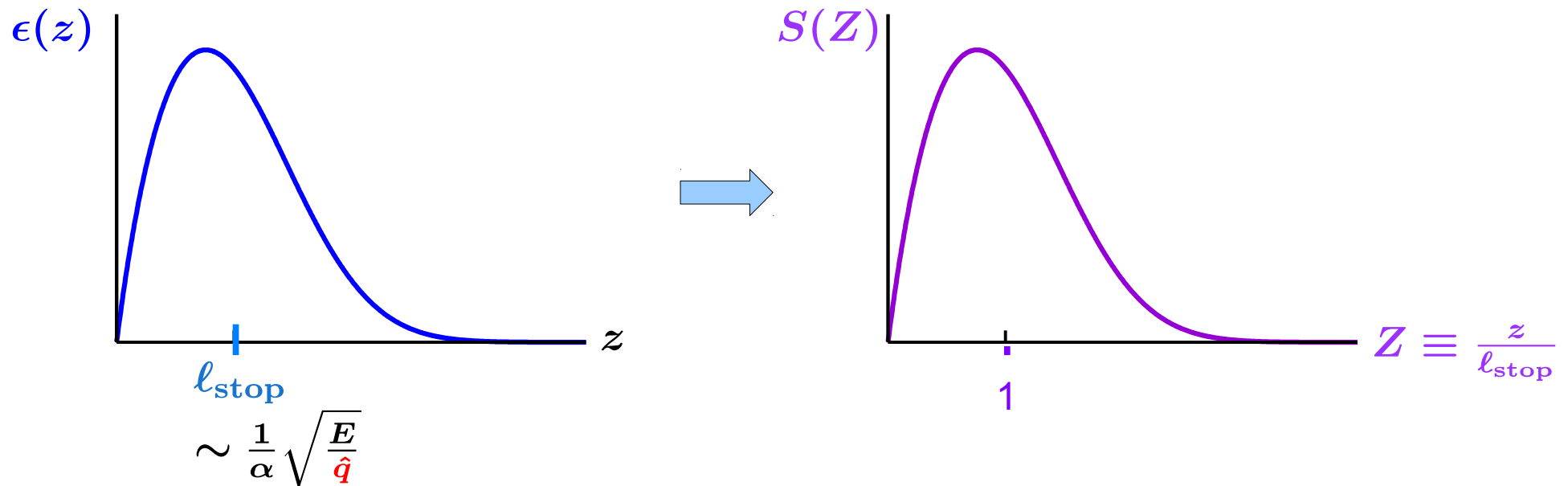
which cancels except for contributions from splittings separated by $\Delta t \lesssim t_{\text{form}}$

A REFINED QUESTION

How big are overlap effects that cannot be absorbed in \hat{q} ?

(2) Choose a theorist observable that is insensitive to \hat{q} :

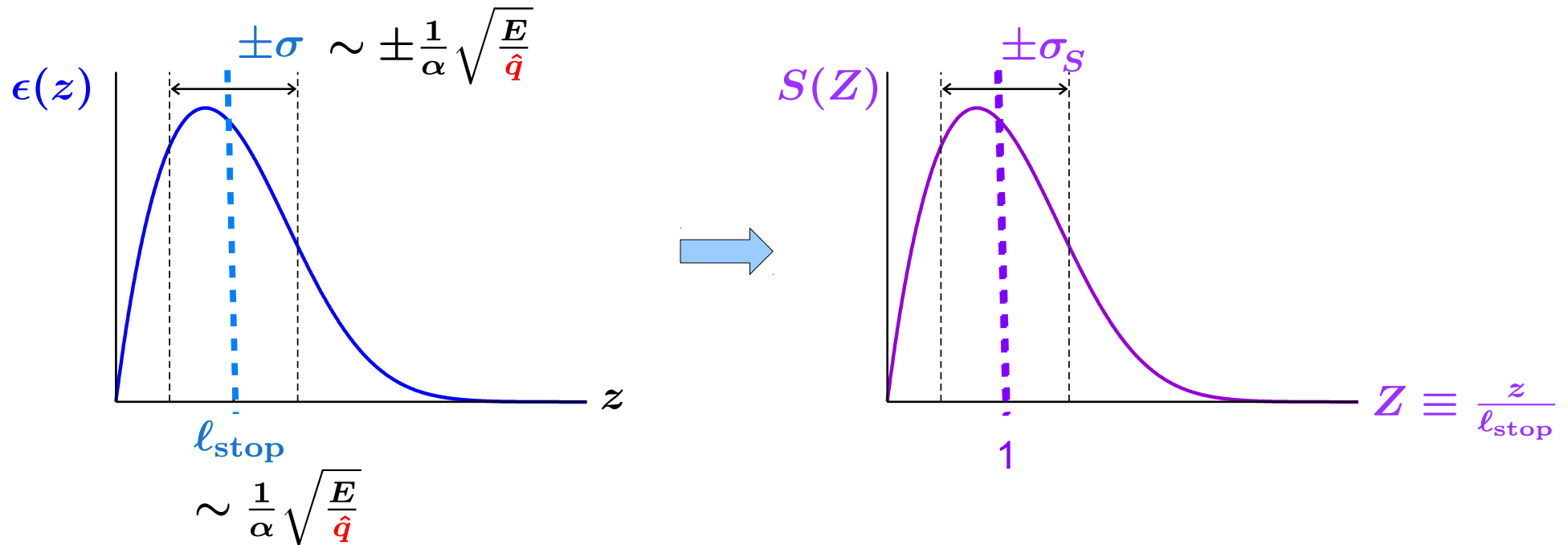
consider the shape $S(Z)$ of the energy deposition distribution:



A REFINED QUESTION

How big are overlap effects that cannot be absorbed in \hat{q} ?

Example

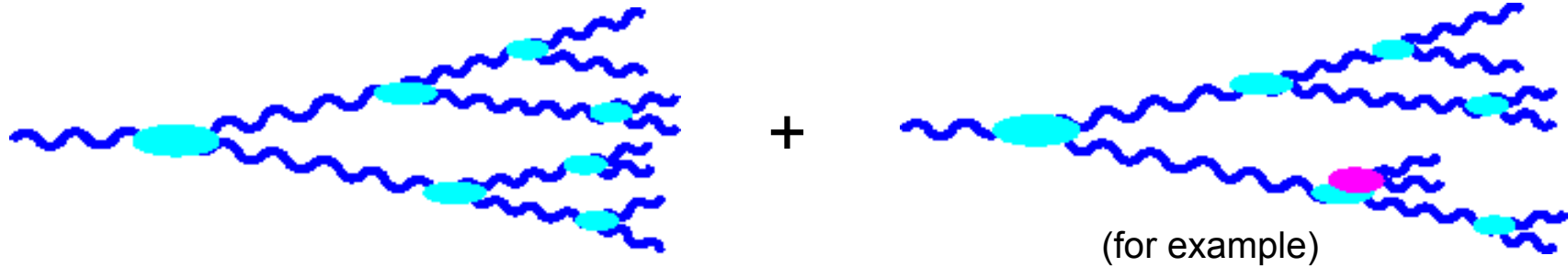


$$\sigma_S = \frac{\sigma}{l_{\text{stop}}} \text{ is independent of } \hat{q} \quad *$$

* Important, interesting, and resolvable caveats that I may not have time to explain.

How to account for overlaps in showers

Think of



as “standard” shower development with independent splittings but two types of localized, independent vertices:



1→2 (normal LPM)

$$\frac{d\Gamma^{\text{split}}}{dx}$$



1→3 (overlap correction)

$$\Delta \frac{d\Gamma}{dx dy}$$

Then treat these “splitting” probabilities as purely classical.

RESULTS

To start: the width of the shape $S(Z)$ of energy deposition

Large- N_f QED [2018 w/ S. Iqbal]:

charge deposition

$$\sigma_S = \frac{\sigma}{\ell_{\text{stop}}} = \left(\frac{\sigma}{\ell_{\text{stop}}} \right)_{\text{LO}} [1 - 0.87 N_f \alpha(\mu)]$$

“LO” means “ignoring overlaps”

Large- N_c QCD (gluons only) [2022 w/ S. Iqbal and O. Elgedawy]:

energy deposition

$$\sigma_S = \frac{\sigma}{\ell_{\text{stop}}} = \left(\frac{\sigma}{\ell_{\text{stop}}} \right)_{\text{LO}} [1 + \text{???} N_c \alpha_s(\mu)]$$

DRUM ROLL
PLEASE

RESULTS

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$$\sigma_S = \frac{\sigma}{\ell_{\text{stop}}} = \left(\frac{\sigma}{\ell_{\text{stop}}} \right)_{\text{LO}} [1 - 0.02 N_c \alpha_s(\mu)]$$

RESULTS

To start: the width of the shape S(Z) of energy deposition

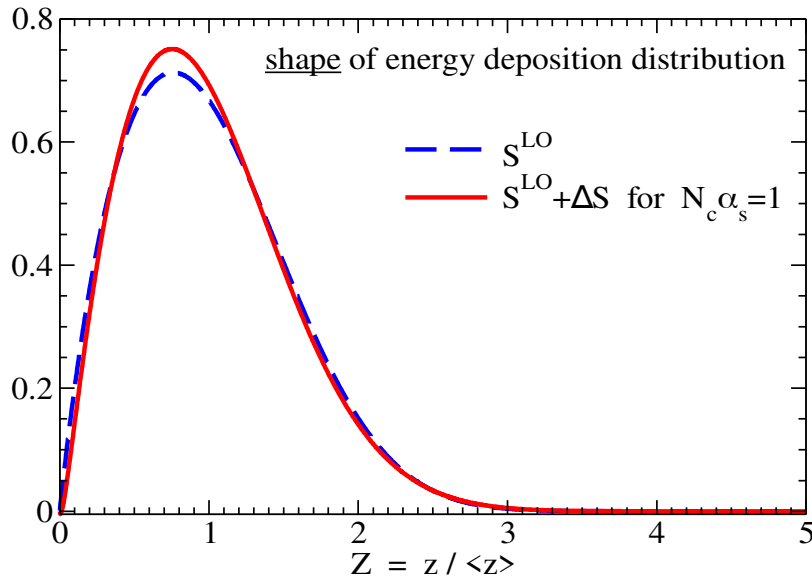
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Conclusion for this test

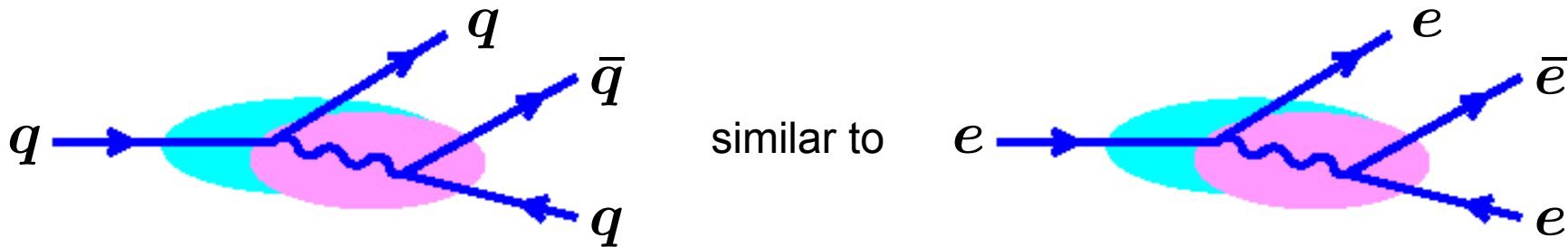
Overlap corrections that cannot be absorbed into \hat{q} are **negligible**.

The QED and gluon results are very different: Discuss!

Large- N_f QED $\sigma_S = \frac{\sigma}{l_{\text{stop}}} = \left(\frac{\sigma}{l_{\text{stop}}} \right)_{\text{LO}} [1 - 0.87 N_f \alpha(\mu)]$

Large- N_c gluons $\sigma_S = \frac{\sigma}{l_{\text{stop}}} = \left(\frac{\sigma}{l_{\text{stop}}} \right)_{\text{LO}} [1 - 0.02 N_c \alpha_s(\mu)]$

A concern: QCD with quarks has some overlap diagrams that look similar to QED



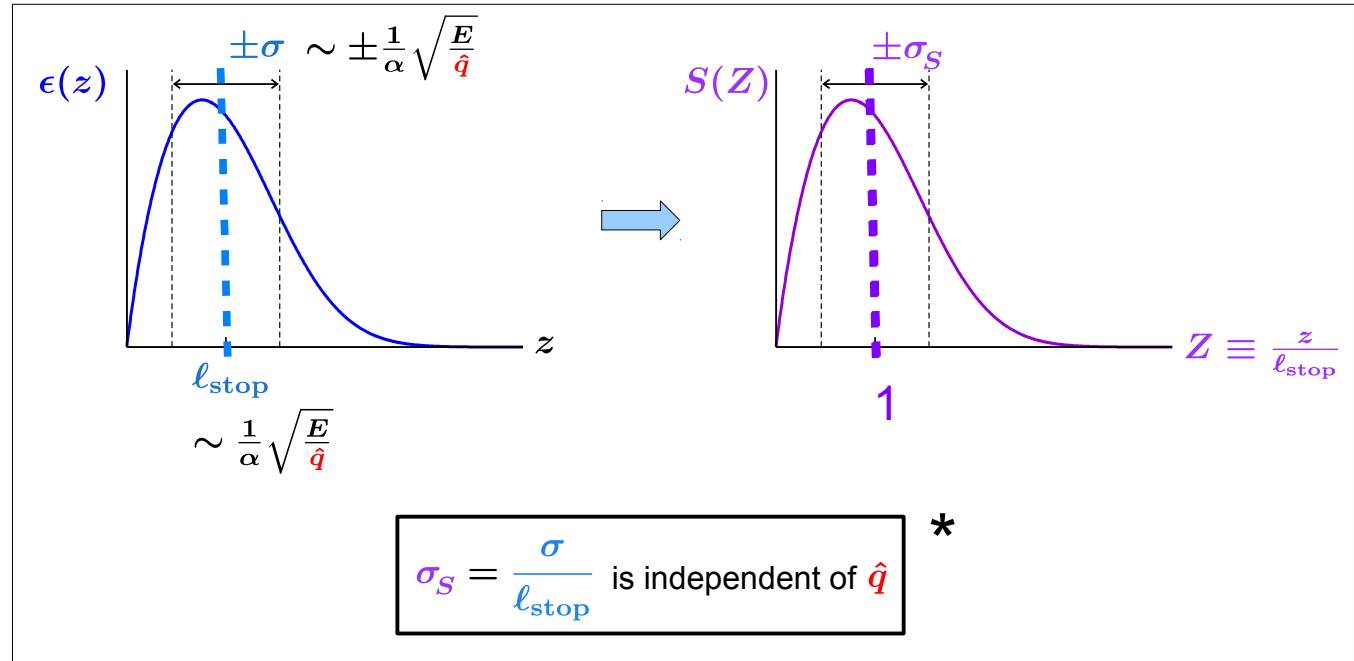
Will adding quarks to the analysis qualitatively change the conclusion for QCD?

Answer: Work in progress.

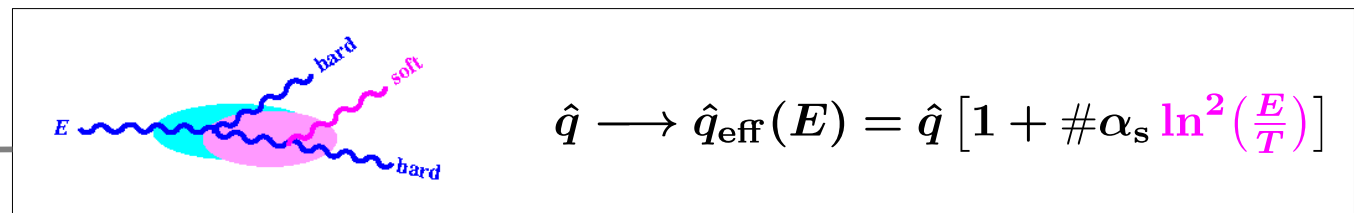
Shrouded from view in this presentation ...

I half-lied about something

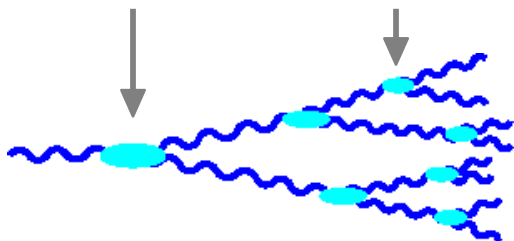
Remember



and why we did that:



But then $\hat{q}_{\text{eff}}(E)$ is different here and there.



Those difference don't quite cancel in $\sigma_S = \sigma / \ell_{\text{stop}}$ and $S(Z)$. They cancel at leading log but leave behind BIG single-log corrections to σ_S and $S(Z)$:

overlap corrections $\sim \alpha_s(\mu) \ln\left(\frac{E}{T}\right)$

Factorization

Remember that soft radiation can be absorbed into \hat{q} .

When factorizing away some IR or UV physics in QFT, we must introduce a factorization scale to do NLO calculations.

Examples

UV divergences absorbed into couplings:

renormalization scale μ

Collinear divergences absorbed into PDFs:

factorization scale M_{fac}

Such factorization scales appear explicitly inside logarithms in NLO results.

- Set them to the appropriate physics scale for the process.
- Check sensitivity to the precise choice of scale.

Our problem

To factorize *all* the soft radiation effects into \hat{q}_{eff} , we introduce an energy factorization scale

$$\Lambda_{\text{fac}} = \# \left(\text{min energy of daughters of } \begin{array}{c} \text{hard} \\ \text{hard} \end{array} \right)$$

where $\# =$
any reasonable $O(1)$ number.

The overlap result shown earlier was the result for $\# = 1$.

Now showing dependence on the normalization # of the factorization scale:

$$\sigma_S = \frac{\sigma}{l_{\text{stop}}} = \left(\frac{\sigma}{l_{\text{stop}}} \right)_{\text{LO}} \left[1 - (0.02 + 0.001 \ln \#) N_c \alpha_s(\mu) \right]$$

Extremely weak dependence on factorization scale.

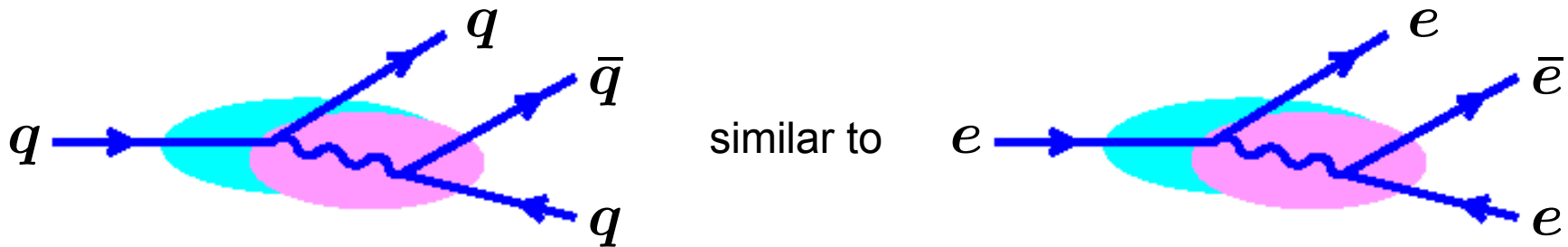


Return to Conclusions

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