



## Gravitational waves from metastable cosmic strings

Based on work in collaboration with W. Buchmüller, V. Domcke, and H. Murayama [1912.03695, 2009.10649, 2107.04578 (today)]

Kai Schmitz
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Gravitational Wave Probes of Physics Beyond Standard Model
Online workshop | Session on topological defects | 12 July 2021

## Gravitational waves from stable cosmic strings

## Cosmic strings:

[See also talk by Hitoshi + all other talks today]


- Topological defects after $U(1)$ breaking in the early Universe
[Ringeval: 1005.4842]


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Gravitational waves (GWs):

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Assumption: Energy loss via particle emission off closed loops is negligible
[Matsunami, Pogosian, Saurabh, Vachaspati: 1903.05102] [Hindmarsh, Lizarraga, Urio, Urrestilla: 2103.16248]

## Stable cosmic strings and NANOGrav

[Blasi, Brdar, KS: 2009.06607]
[See also Ellis, Lewicki: 2009.06555]

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(2) GUT scale $\Lambda \sim 10^{15 \cdots 16} \mathrm{GeV}$ points to $G \mu \sim 10^{-(7 \cdots 8)}$ (smaller $\alpha$ ?)
(2) Signal at higher frequencies too small for LIGO, Virgo, KAGRA

## Cosmic strings and grand unification

[Dror, Hiramatsu, Kohri, Murayama, White: 1908.03227]
[See also King, Pascoli, Turner, Zhou: 2005.13549, 2106.15634]


UV embedding of the seesaw mechanism in GUT models:
Neutrino mass, leptogenesis, cosmic strings, GWs, proton decay

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Assumption: Inflation dilutes monopoles; otherwise string-monopole gas

## Monopole pair production



## Decay rate per string length: <br> [Vilenkin: Nucl. Phys. B 196 (1982) 240] <br> [Preskill, Vilenkin: hep-ph/9209210] <br> [Monin, Voloshin: 0808.1693] <br> $$
\begin{equation*} \Gamma_{d}=\frac{d \#}{d t d \ell}=\frac{\mu}{2 \pi} e^{-\pi \kappa}, \quad \kappa=\frac{m^{2}}{\mu} \tag{1} \end{equation*}
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Monopoles with and without unconfined magnetic flux:

- Unconfined flux: $M \bar{M}$ annihilation, emission of massless gauge bosons
- No unconfined flux: energy loss only via emission of gravitational waves


## Possible scenarios

$$
\begin{equation*}
W_{B-L}=\lambda T\left(S \bar{S}-\frac{1}{2} v_{B-L}^{2}\right)+\frac{h_{i}}{M_{*}} S^{2} N_{i}^{2} \tag{2}
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$B-L$ phase transition after
supersymmetric hybrid inflation:

- T: inflaton, S, $\bar{S}$ : Higgs / waterfall fields, $N_{i}$ : right-handed neutrinos


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[Buchmüller, Domcke, Murayama, KS: 1912.03695]

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$v_{B-L} \sim(3 \cdots 6) \times 10^{15} \mathrm{GeV}$
[Buchmüller, KS, Vertongen: 1008.2355, 1104.2750]
[Buchmüller, Domcke, KS: 1111.3872, 1202.6679, 1203.0285]
[Buchmüller, Domcke, Kamada, KS: 1305.3392, 1309.7788]


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$\underset{\text { MBuchnuiler: 2102.09923] }}{\text { Minimal alternative: } S U(2) \times U(1) \xrightarrow{\text { triplet }} U(1) \times U(1) \xrightarrow{\text { doublets }} U(1), ~(1) ~}$


## Strategy

End of scaling when long string segments begin to enter the horizon:
[Leblond, Shlaer, Siemens: 0903.4686]

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\begin{equation*}
\Gamma_{d} \ell t_{s} \sim \Gamma_{d} H^{-1} t_{s} \sim \Gamma_{d} t_{s}^{2} \sim 1 \quad \Rightarrow \quad t_{s} \sim \frac{1}{\sqrt{\Gamma_{d}}} \tag{3}
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Scaling regime, $t<t_{s}$

- Loops: emit GWs, decay into segments negligible
- Long strings: decay into segments on superhorizon scales, chop off closed loops, GW emission negligible


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## Decay regime, $t>t_{s}$

- Loops: emit GWs and decay into segments
- Segments from loops and long strings: emit GWs and decay into segments; no production of new loops


## Formal description

Kinetic equation for the number densities of loops and segments, $n$ and $\tilde{n}$ :

$$
\begin{equation*}
\partial_{t} n(\ell, t)=S(\ell, t)-\partial_{\ell}[u(\ell, t) n(\ell, t)]-\left[3 H(t)+\Gamma_{d} \ell\right] n(\ell, t) \tag{4}
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Time derivative of the string length $u=\dot{\ell}$ :

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Challenge: Solve set of partial integro-differential equations in both the scaling and decay regimes, match solutions at $t=t_{s}$. (Plus, RD / MD.)

## Number densities

Loop number density during the decay regime in the radiation era:
[Cf. Blanco-Pillado, Olum, Shlaer: 1309.6637] [Cf. Blanco-Pillado, Olum: 1709.02693] [See also talk by Jose Juan today]

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- Exponential suppression at $\ell t>1 / \Gamma_{d}=t_{s}^{2}$ or $t^{2}>2 /\left(\Gamma_{d} \Gamma G \mu\right)=t_{e}^{2}$ because of new exponential suppression factor:

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 $\tilde{n}_{>}^{(s) \mathrm{mm}}, \tilde{n}_{>}^{(1) \mathrm{rr}}, \tilde{n}_{>}^{(1) \mathrm{rm}}, \tilde{n}_{>}^{(1) \mathrm{mm}}$. The integro-differential equation for $\tilde{n}_{>}^{(1)}$ is solved by an infinite series that needs be evaluated order by order.


## Spectrum

Compute GW spectrum following the standard procedure:

$$
\begin{equation*}
\Omega_{\mathrm{gw}}(f)=\frac{G \mu^{2}}{\rho_{\text {crit }}} \sum_{k} P_{k} \frac{2 k}{f} \int_{t_{\mathrm{ini}}}^{t_{0}} d t\left[\frac{a(t)}{a\left(t_{0}\right)}\right]^{5} n\left(\frac{a(t)}{a\left(t_{0}\right)} \frac{2 k}{f}, t\right) \tag{7}
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- Loop contributions scales like $f^{2}$ at low $f$ [cf. Buchuiler, Domcce, Murayama, Ks: 1912.03695]


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\begin{equation*}
\Omega_{\mathrm{gw}}(f)=\frac{G \mu^{2}}{\rho_{\text {crit }}} \sum_{k} P_{k} \frac{2 k}{f} \int_{t_{\mathrm{ini}}}^{t_{0}} d t\left[\frac{a(t)}{a\left(t_{0}\right)}\right]^{5} n\left(\frac{a(t)}{a\left(t_{0}\right)} \frac{2 k}{f}, t\right) \tag{7}
\end{equation*}
$$




- Loop contribution almost always dominant [cf. Leelond, Shaer, Siemens: 0003.4686]
- Loop contributions scales like $f^{2}$ at low $f$ [cf. Buchmiller, Domcce, Murayama, ks: 1912.03695]
- Suppress spectrum in nHz range, explain NANOGrav for larger $G \mu$


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## Thank you very much for your attention!

