

Based on work in collaboration with W. Buchmüller, V. Domcke, and H. Murayama [1912.03695, 2009.10649, 2107.04578 (today)]

Kai Schmitz MSCA Fellow in the CERN Theory Group Gravitational Wave Probes of Physics Beyond Standard Model Online workshop | Session on topological defects | 12 July 2021



[See also talk by Hitoshi + all other talks today]

Cosmic strings:

• Topological defects after *U*(1) breaking in the early Universe

[Ringeval: 1005.4842]



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Gravitational waves (GWs):

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Gravitational waves (GWs):

• Loop oscillations + GW bursts from cusps and kinks on loops

Assumption: Energy loss via particle emission off closed loops is negligible [Matsunami, Pogosian, Saurabh, Vachaspati: 1903.05102] [Hindmarsh, Lizarraga, Urio, Urrestilla: 2103.16248]

Stable cosmic strings and NANOGrav

[Blasi, Brdar, KS: 2009.06607] [See also Ellis, Lewicki: 2009.06555]



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- \odot GUT scale $\Lambda \sim 10^{15 \cdots 16} \text{ GeV}$ points to $G\mu \sim 10^{-(7 \cdots 8)}$ (smaller α ?)
- Signal at higher frequencies too small for LIGO, Virgo, KAGRA

Cosmic strings and grand unification

[Dror, Hiramatsu, Kohri, Murayama, White: 1908.03227] [See also King, Pascoli, Turner, Zhou: 2005.13549, 2106.15634]



UV embedding of the seesaw mechanism in GUT models: Neutrino mass, leptogenesis, cosmic strings, GWs, proton decay

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Assumption: Inflation dilutes monopoles; otherwise string-monopole gas



Decay rate per string length:

[Vilenkin: Nucl. Phys. B 196 (1982) 240] [Preskill, Vilenkin: hep-ph/9209210] [Monin, Voloshin: 0808.1693]

$$\Gamma_d = \frac{d\#}{dtd\ell} = \frac{\mu}{2\pi} e^{-\pi\kappa}, \qquad \kappa = \frac{m^2}{\mu} \qquad (1)$$

String tension µ, monopole mass m



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Monopoles with and without unconfined magnetic flux:

- Unconfined flux: $M\bar{M}$ annihilation, emission of massless gauge bosons
- No unconfined flux: energy loss only via emission of gravitational waves

Possible scenarios

$$W_{B-L} = \lambda T \left(S\bar{S} - \frac{1}{2} v_{B-L}^2 \right) + \frac{h_i}{M_*} S^2 N_i^2$$
(2)

B-L phase transition after supersymmetric hybrid inflation:

• T: inflaton, S, \overline{S} : Higgs / waterfall fields, N_i : right-handed neutrinos

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[Buchmüller, KS, Vertongen: 1008.2355, 1104.2750] [Buchmüller, Domcke, KS: 1111.3872, 1202.6679, 1203.0285] [Buchmüller, Domcke, Kamada, KS: 1305.3392, 1309.7788]

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 $\underset{\text{[Buchmüller: 2102.08923]}}{\text{Minimal alternative: }} SU(2) \times U(1) \stackrel{\text{triplet}}{\longrightarrow} U(1) \times U(1) \stackrel{\text{doublets}}{\longrightarrow} U(1)$

Strategy

End of scaling when long string segments begin to enter the horizon: [Leblond, Shlaer, Siemens: 0903.4686]

$$\Gamma_d \,\ell \, t_s \sim \Gamma_d H^{-1} \, t_s \sim \Gamma_d t_s^2 \sim 1 \quad \Rightarrow \quad t_s \sim \frac{1}{\sqrt{\Gamma_d}} \tag{3}$$



Scaling regime, $t < t_s$

- Loops: emit GWs, decay into segments negligible
- Long strings: decay into segments on superhorizon scales, chop off closed loops, GW emission negligible

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Decay regime, $t > t_s$

- Loops: emit GWs and decay into segments
- Segments from loops and long strings: emit GWs and decay into segments; no production of new loops

Kinetic equation for the number densities of loops and segments, $\stackrel{\circ}{n}$ and $\tilde{n}{:}$

$$\partial_t n(\ell, t) = S(\ell, t) - \partial_\ell \left[u(\ell, t) n(\ell, t) \right] - \left[3H(t) + \Gamma_d \ell \right] n(\ell, t) \quad (4)$$

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Source term *S*:

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Time derivative of the string length $u = \dot{\ell}$:

• Long strings during scaling: $u = 3H(t) \ell - 2\ell/t$

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Time derivative of the string length $u = \dot{\ell}$:

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Challenge: Solve set of partial integro-differential equations in both the scaling and decay regimes, match solutions at $t = t_s$. (Plus, RD / MD.)

Loop number density during the decay regime in the radiation era:

[Cf. Blanco-Pillado, Olum, Shlaer: 1309.6637] [Cf. Blanco-Pillado, Olum: 1709.02693] [See also talk by Jose Juan today]

$$\stackrel{\circ}{n_{>}}^{\rm rr}(\ell,t) = \frac{B \, e^{-\Gamma_{d} \left[\ell(t-t_{s})+\frac{1}{2}\Gamma G \mu(t-t_{s})^{2}\right]}}{t^{3/2} \left(\ell+\Gamma G \mu t\right)^{5/2}} \,\Theta\left(\alpha t_{s}-\bar{\ell}\left(t_{s}\right)\right) \,\Theta\left(t_{\rm eq}-t\right)$$
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• Exponential suppression at $\ell t > 1/\Gamma_d = t_s^2$ or $t^2 > 2/(\Gamma_d \Gamma G \mu) = t_e^2$ because of new exponential suppression factor:

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Similar results for $\overset{\circ}{n_{<}}^{rr}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{>}}^{rm}$, $\overset{\circ}{n_{<}}^{rm}$, $\overset{\circ}{n_{<}}^{rr}$, $\overset{\circ}{n_{<}}^{(s) rr}$, $\overset{\circ}{n_{>}}^{(s) rr}$, $\overset{\circ}{n_{>}}^{(s)$

Compute GW spectrum following the standard procedure:

$$\Omega_{\rm gw}(f) = \frac{G\mu^2}{\rho_{\rm crit}} \sum_k P_k \frac{2k}{f} \int_{t_{\rm ini}}^{t_0} dt \left[\frac{a(t)}{a(t_0)}\right]^5 n\left(\frac{a(t)}{a(t_0)}\frac{2k}{f}, t\right)$$
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- Suppress spectrum in nHz range, explain NANOGrav for larger $G\mu$



Extrapolate spectrum to large *f* and compare with LIGO, Virgo, KAGRA:



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- Tilt at PTA frequencies correlated with amplitude at LVK frequencies
- LISA will probe the entire parameter space consistent with NANOGrav



Metastable cosmic strings:

• Prediction in many GUT models when combined with inflation to solve the monopole problem



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Thank you very much for your attention!