

# **Searching for gravitational wave bursts from cosmic string cusps with the Parkes Pulsar Timing Array**

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in collaboration with  
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and Parkes Pulsar Timing Array team

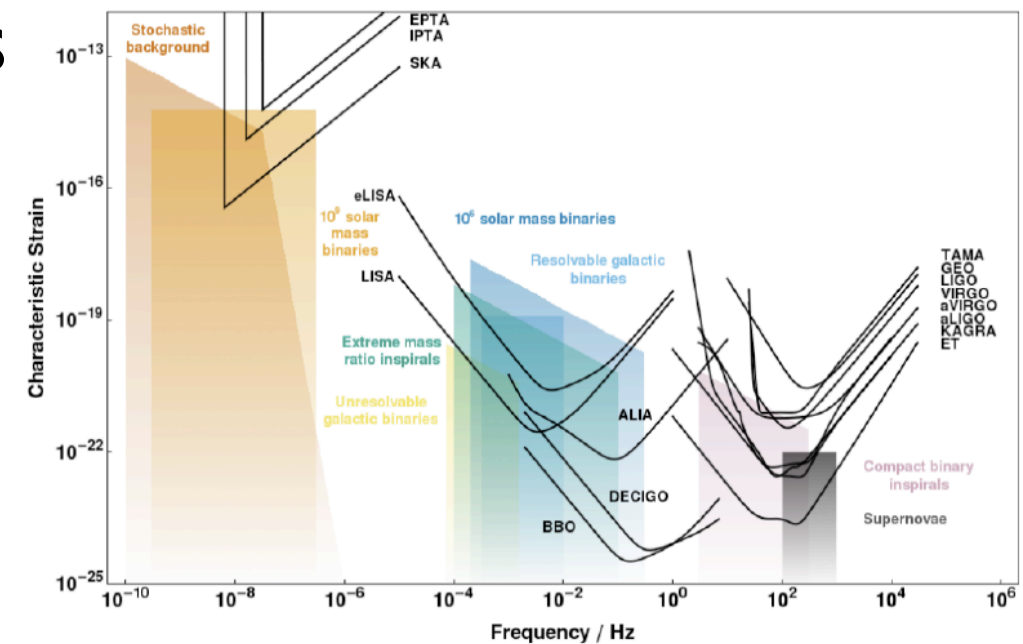
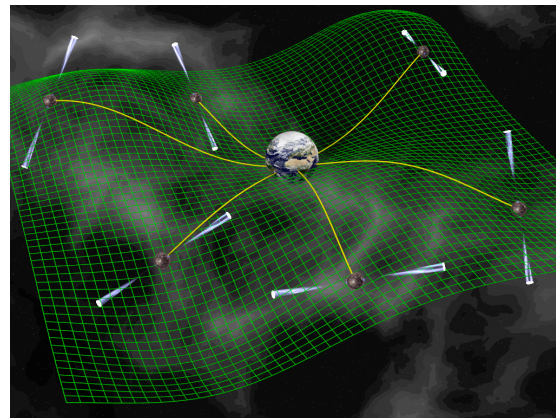
based on *Yonemaru et al. MNRAS 501, 701–712 (2021)*

**12 July 2021**

Gravitational Wave Probes of Physics Beyond Standard Model

# Introduction

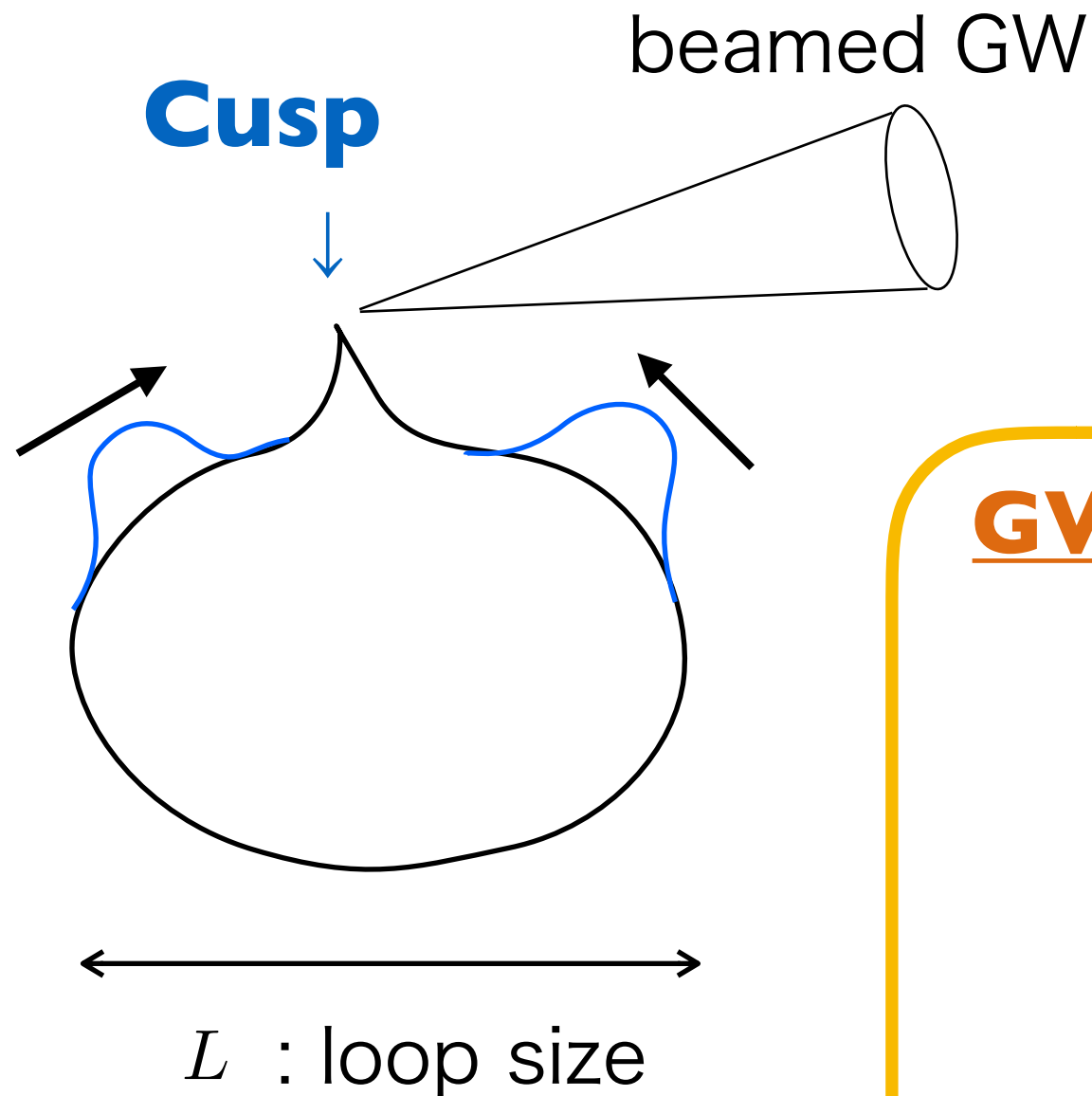
- Cosmic strings may have been generated in the early Universe  
→ **unique way to access the early universe physics**
- **Gravitational Waves (GWs)** is a powerful observational tool to probe cosmic strings
- **Pulsar timing** probes nanoHz GWs



- GWs coming from different directions overlap one another and form a **stochastic GW background**
- Nearby events are observed as a **single burst event**

# GW bursts from cosmic strings

Strong GW bursts are emitted from singular points called **cusps**



Damour & Vilenkin, PRD 64, 064008 (2001)

**GW amplitude** (in Fourier space)

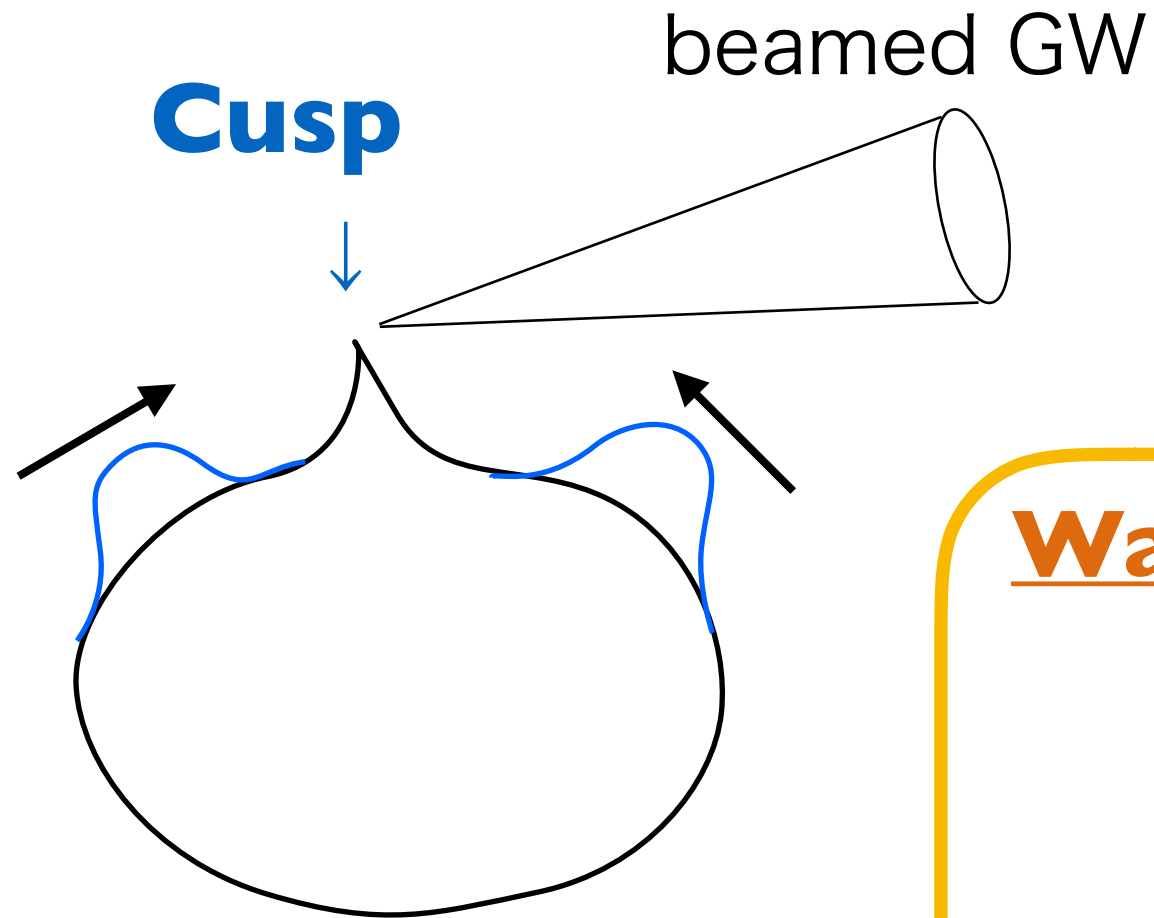
$$\tilde{h}(f) = \frac{G\mu L}{[(1+z)fL]^{1/3}r(z)f}$$

$G\mu$  : tension = line density

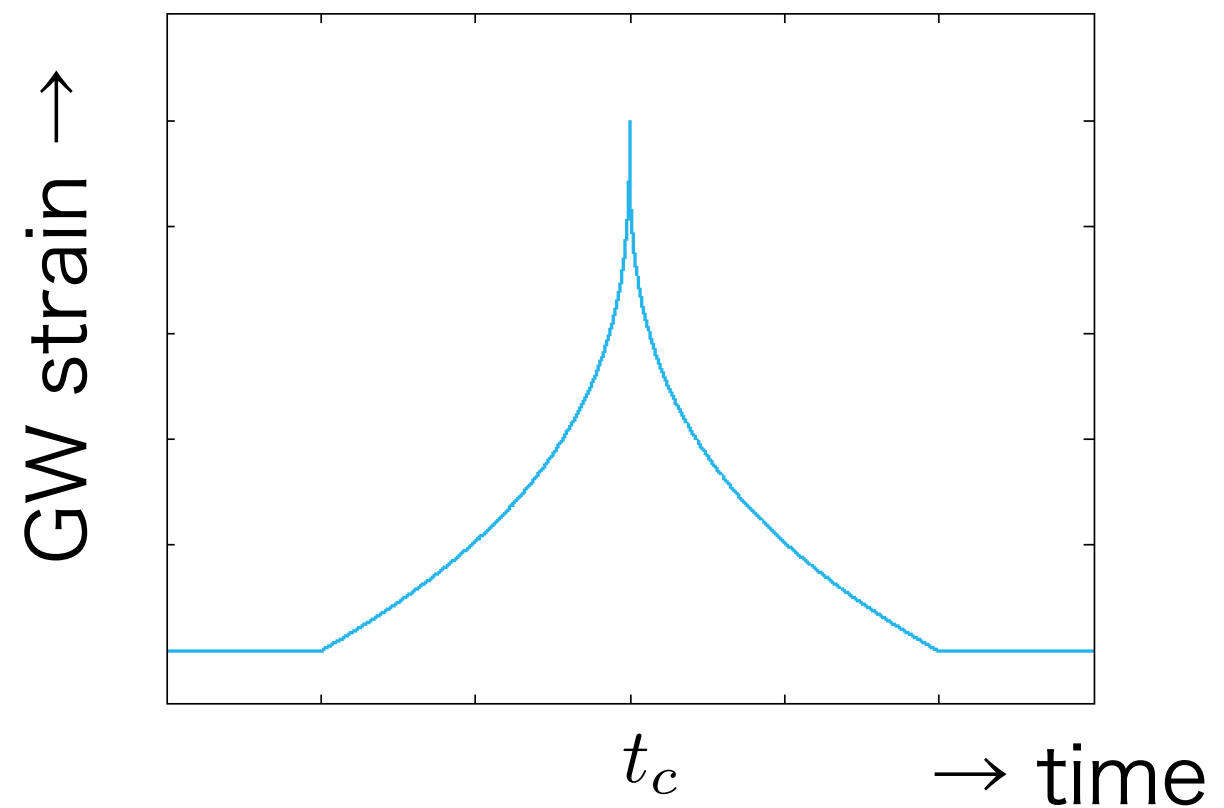
$r(z)$  : distance to the source

# GW bursts from cosmic strings

Strong GW bursts are emitted from singular points called **cusps**

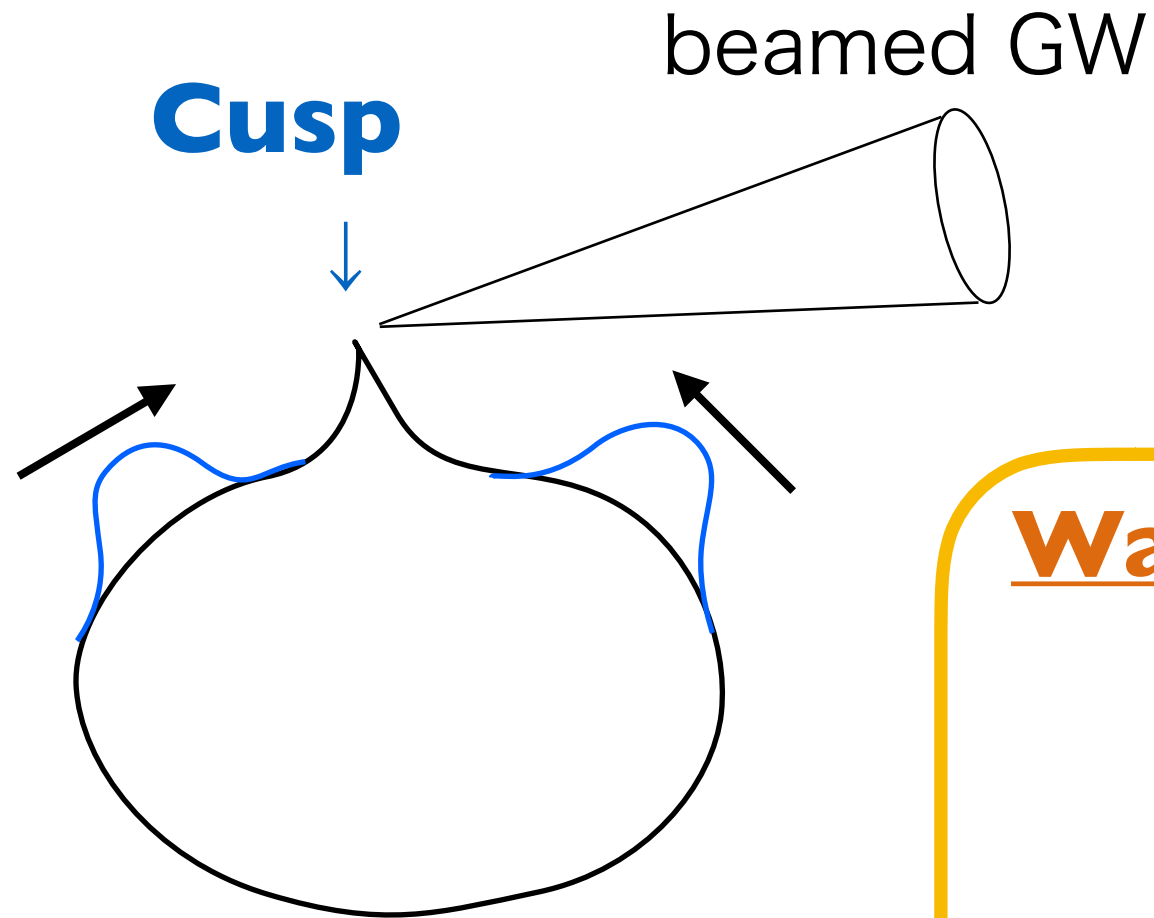


## Waveform (in time domain)

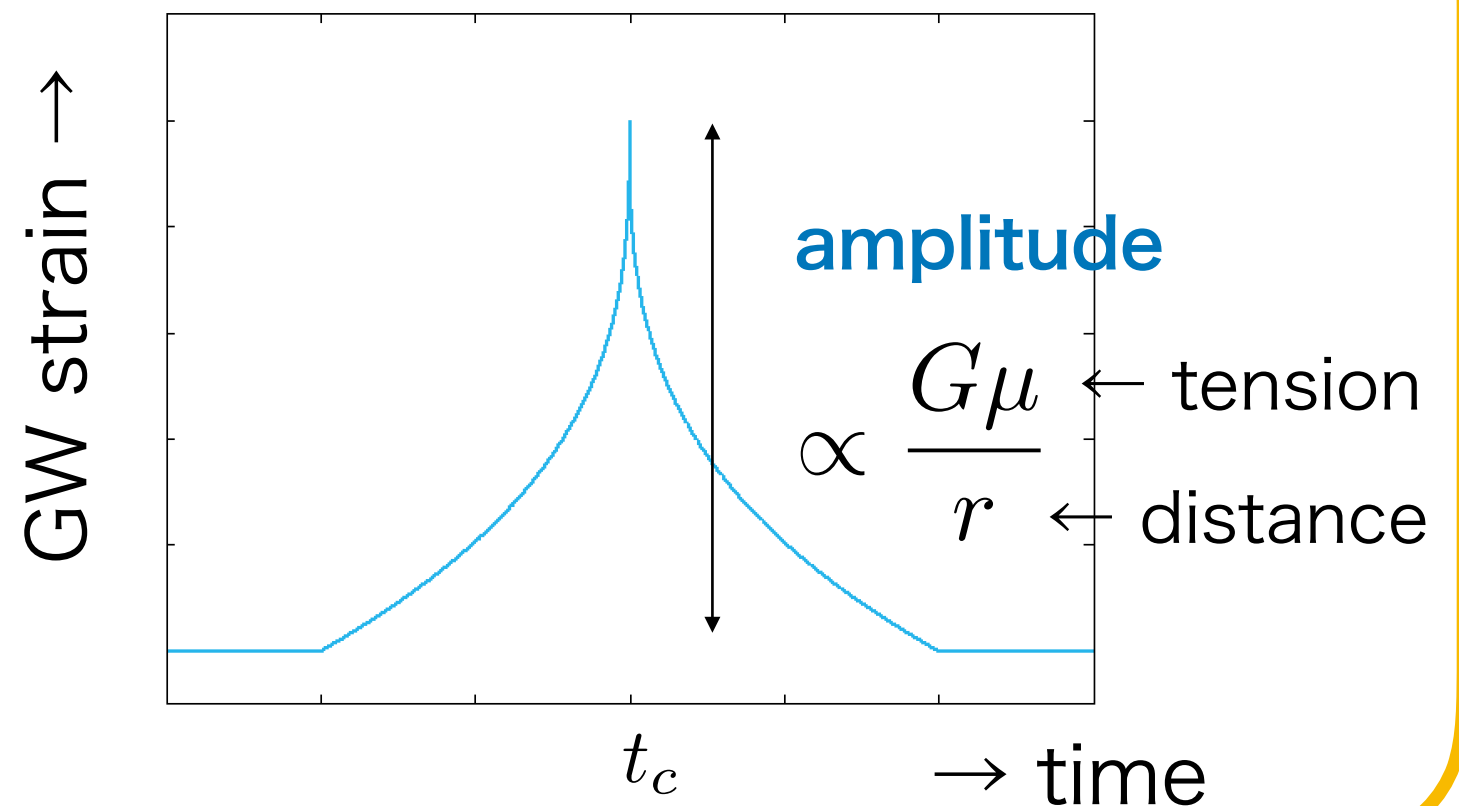


# GW bursts from cosmic strings

Strong GW bursts are emitted from singular points called **cusps**

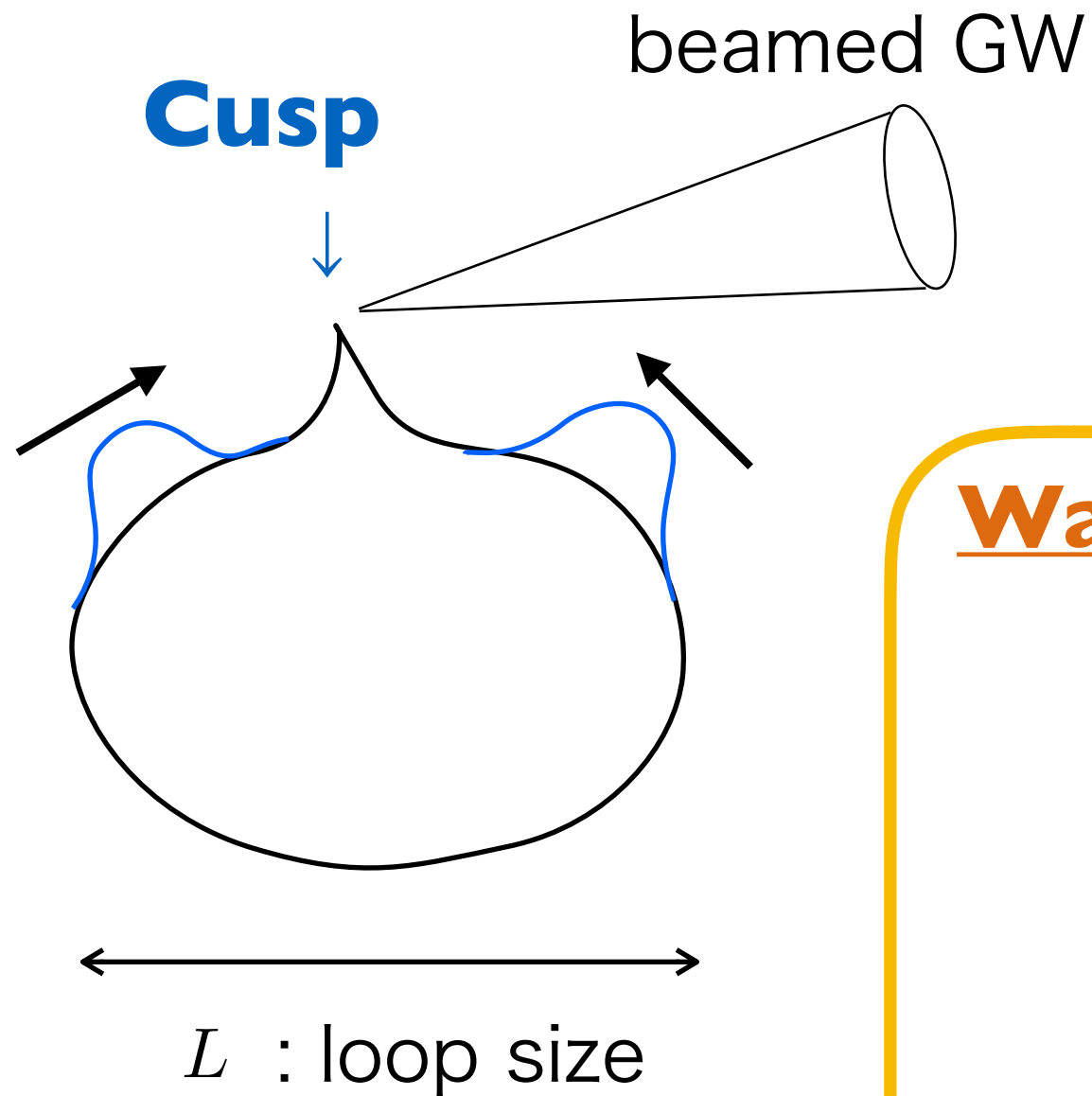


## Waveform

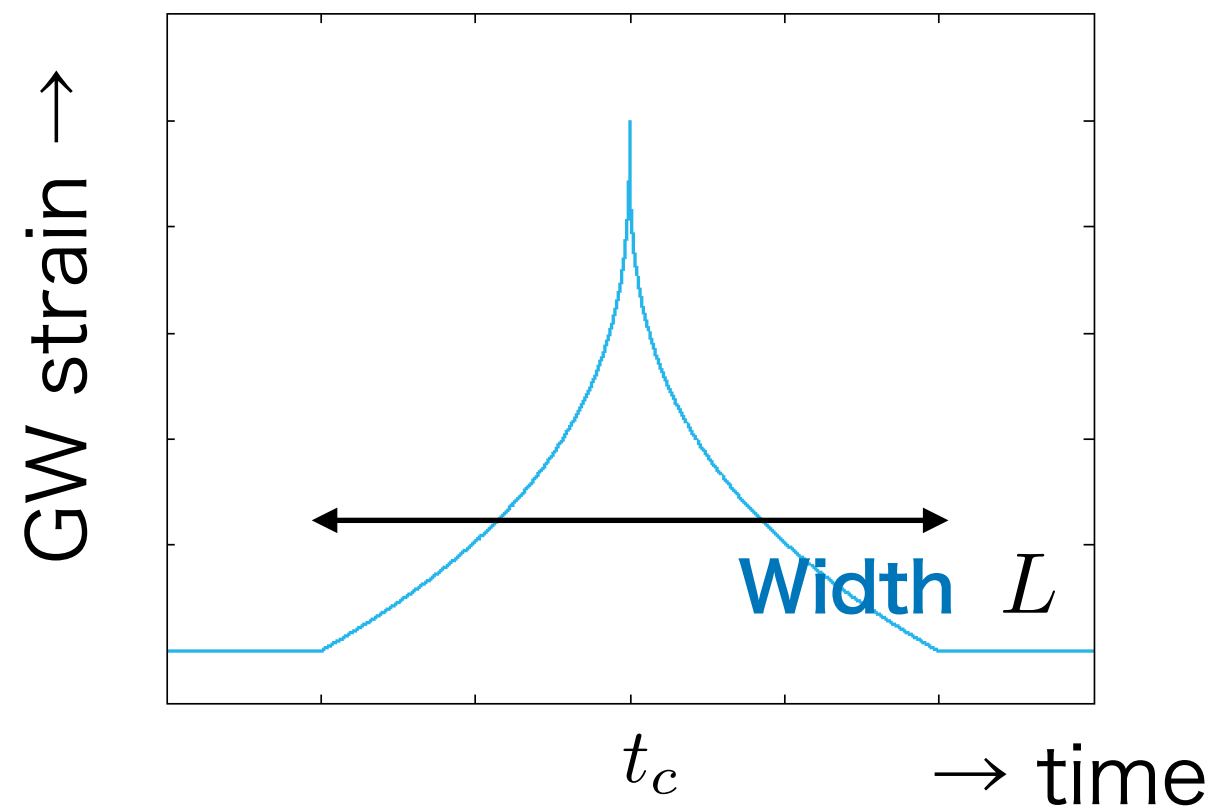


# GW bursts from cosmic strings

Strong GW bursts are emitted from singular points called **cusps**

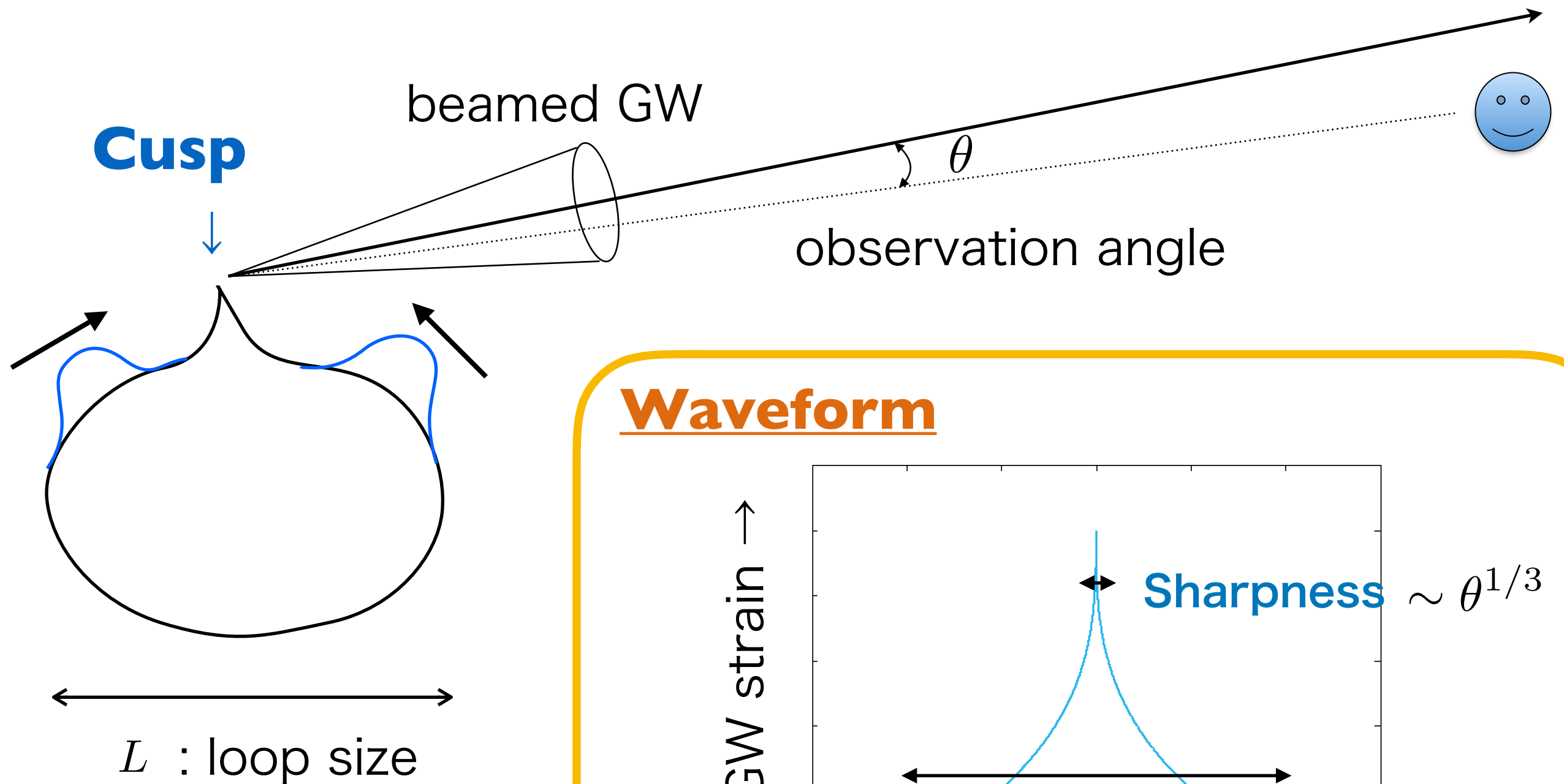


## Waveform

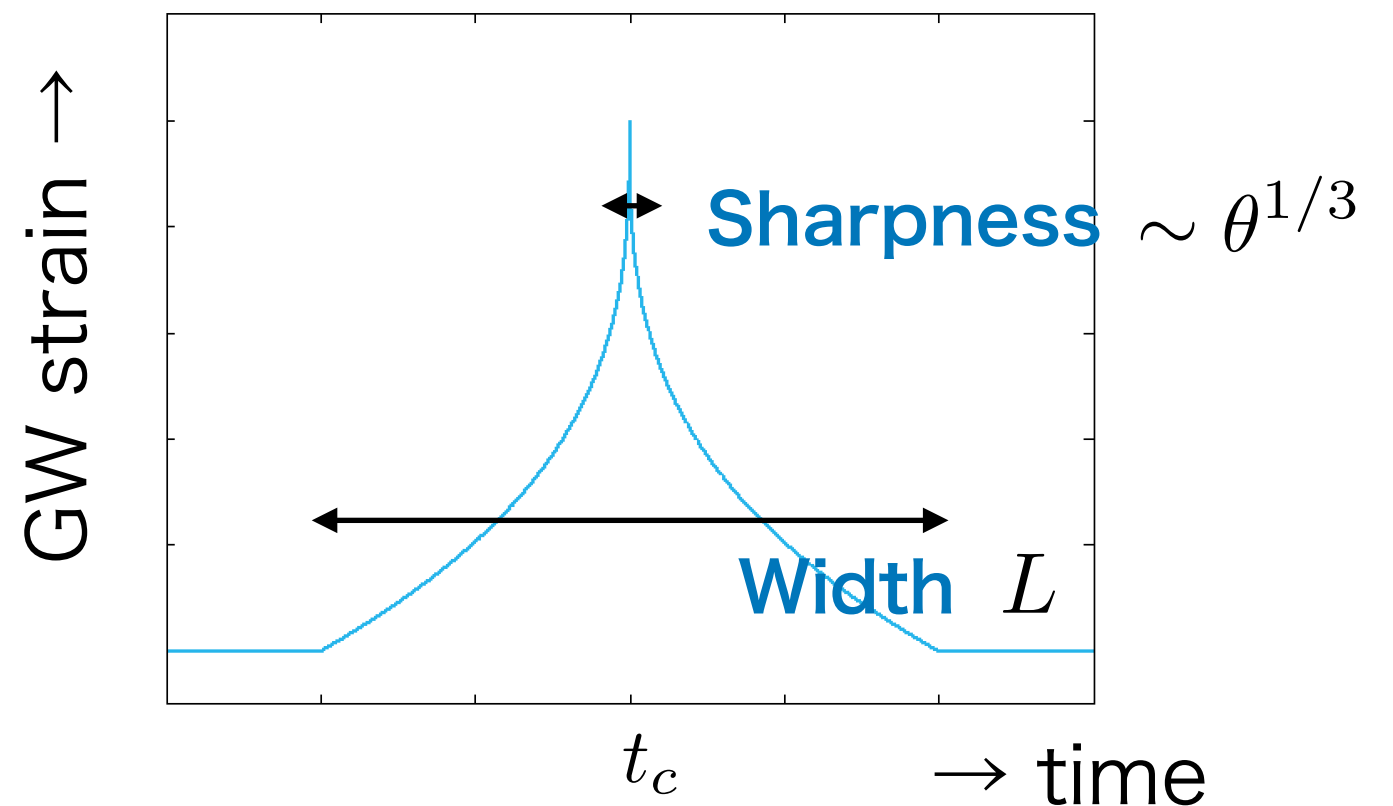


# GW bursts from cosmic strings

Strong GW bursts are emitted from singular points called **cusps**



## Waveform



# Observation

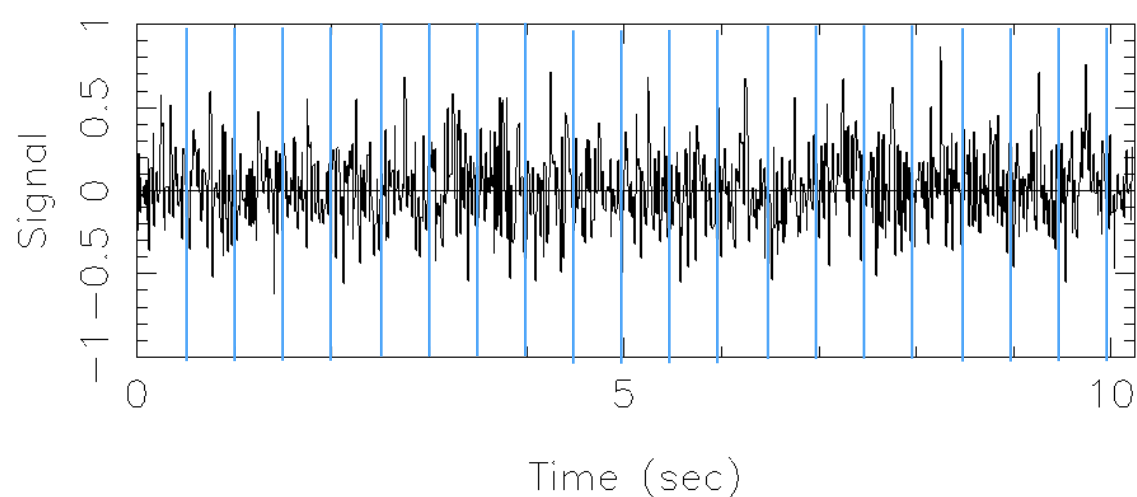


## Parkes Pulsar Timing Array

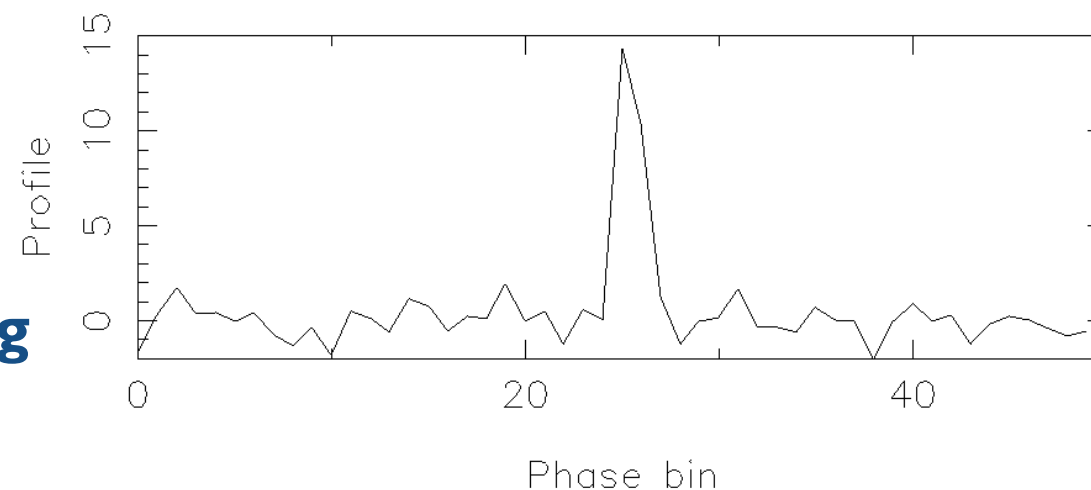
~ 20 pulsars

1 hour for each every 2 weeks

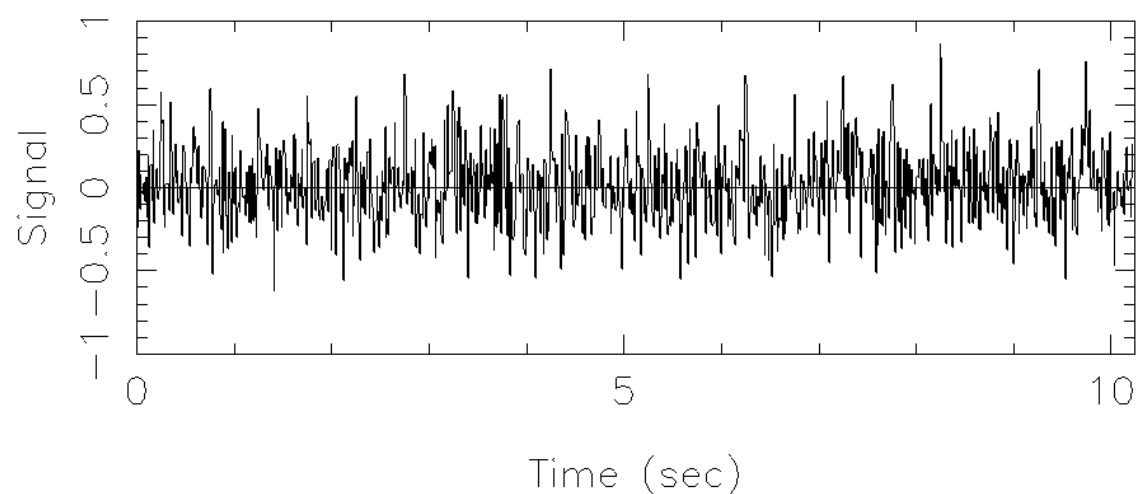
sum up all the bins



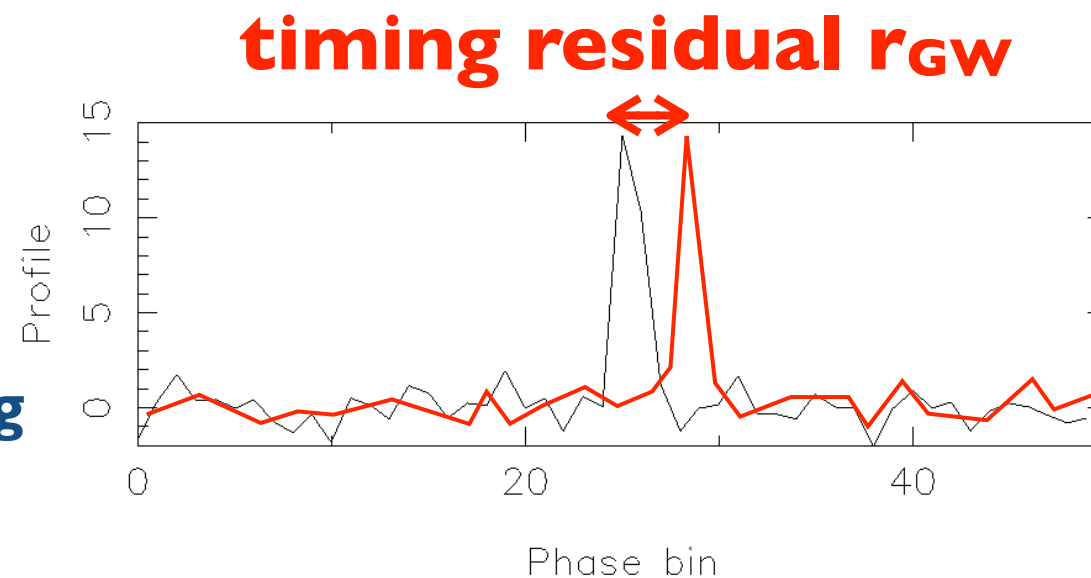
➔  
**folding**



2 weeks later



➔  
**folding**





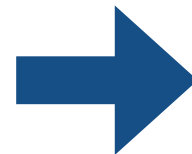
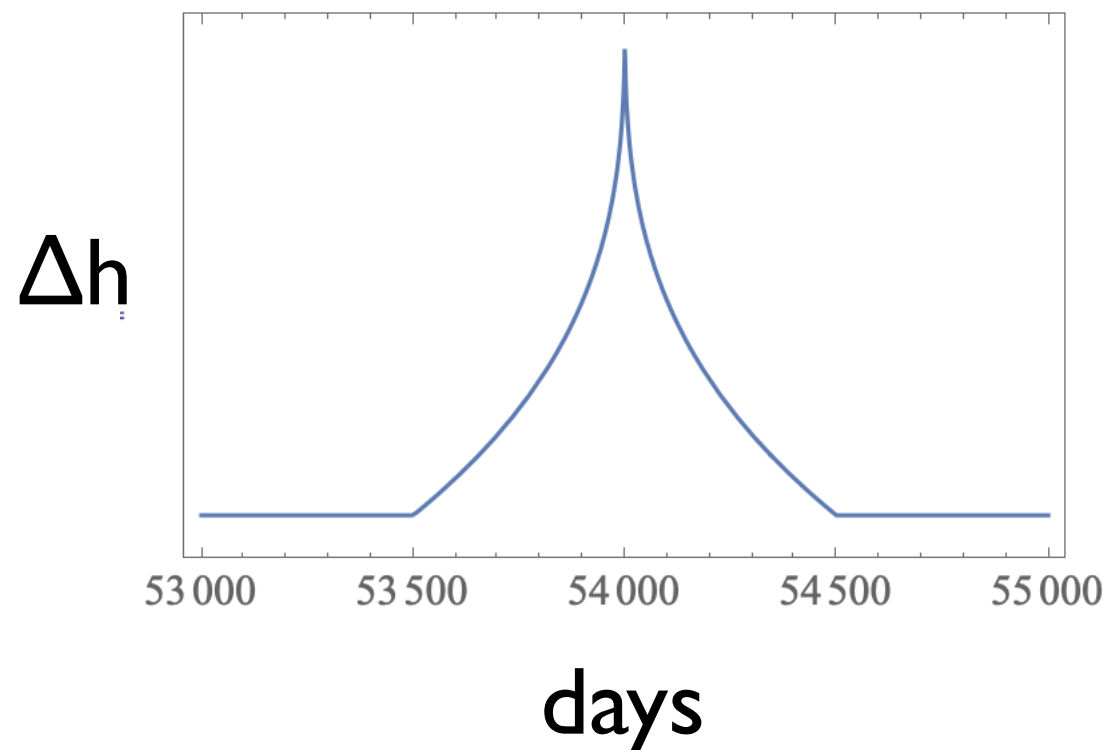
# Timing residual induced by GWs

Timing residual: 
$$r_{\text{GW}}(t) = \sum_{A=+, \times} \underbrace{F^A(\hat{\Omega}, \hat{p})}_{\text{antenna pattern (response function)}} \int \underbrace{\Delta h_A(t', \hat{\Omega})}_{\text{GW amplitude}} dt'$$

antenna pattern  
(response function)

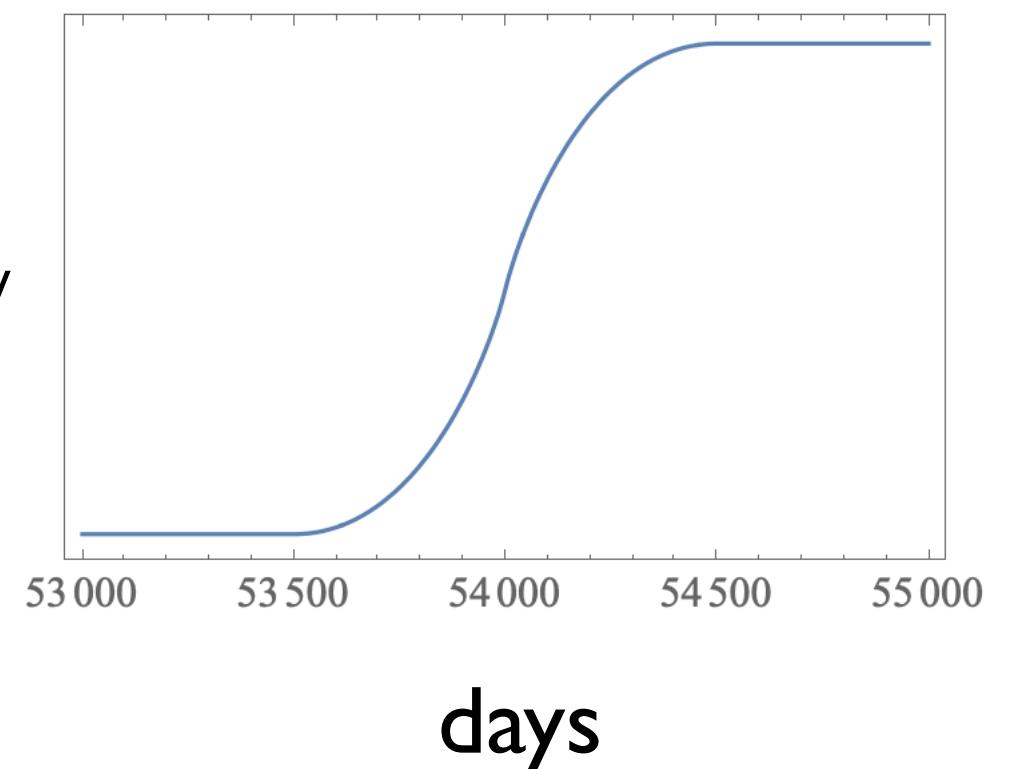
GW amplitude

**GW amplitude**



$r_{\text{GW}}$

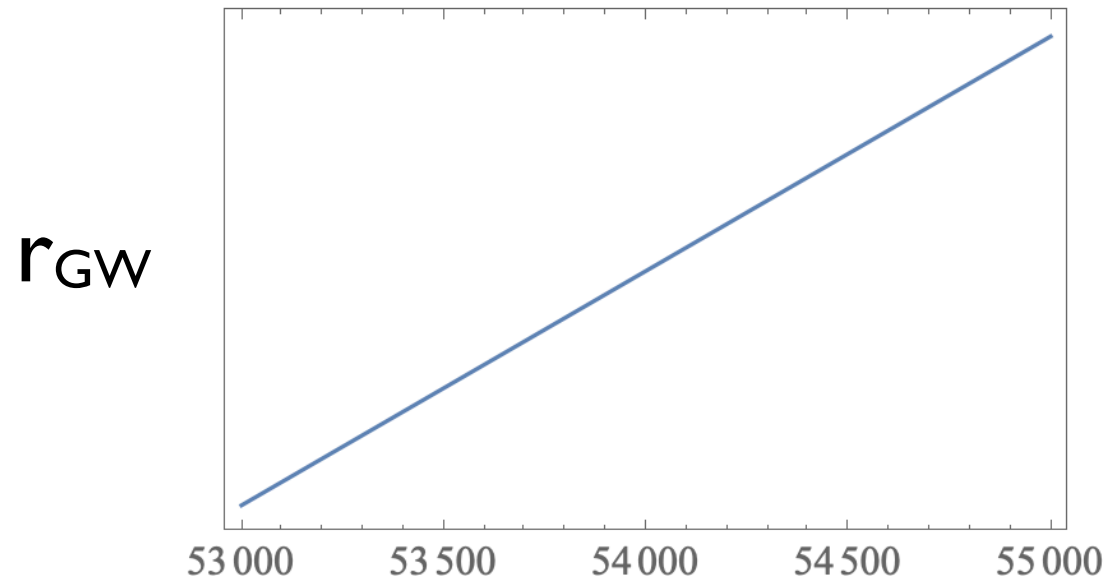
**Timing residual**



# Post-fit effect

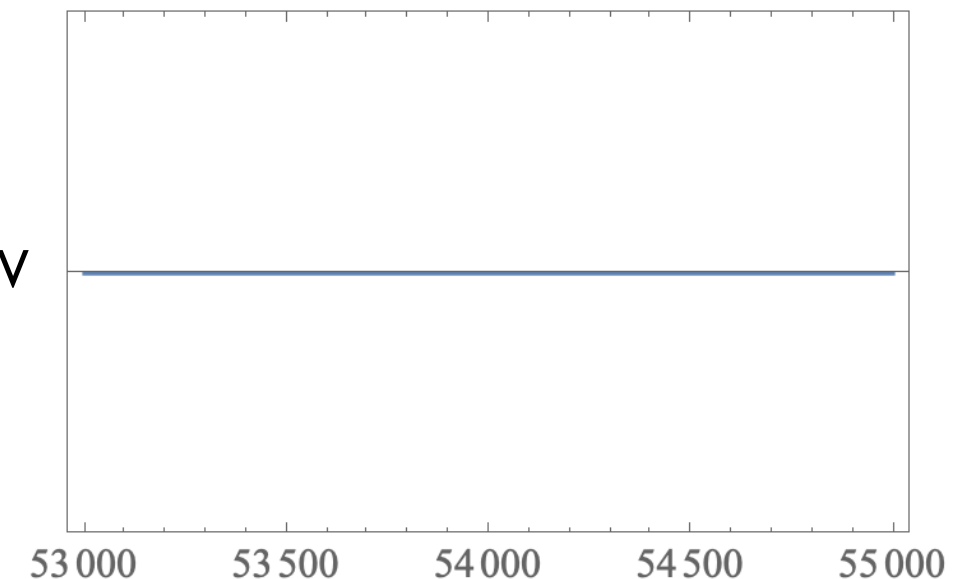
We do not know the true pulsar period

Even without GW signal, we get  $r_{\text{GW}}$  by assuming wrong pulse period

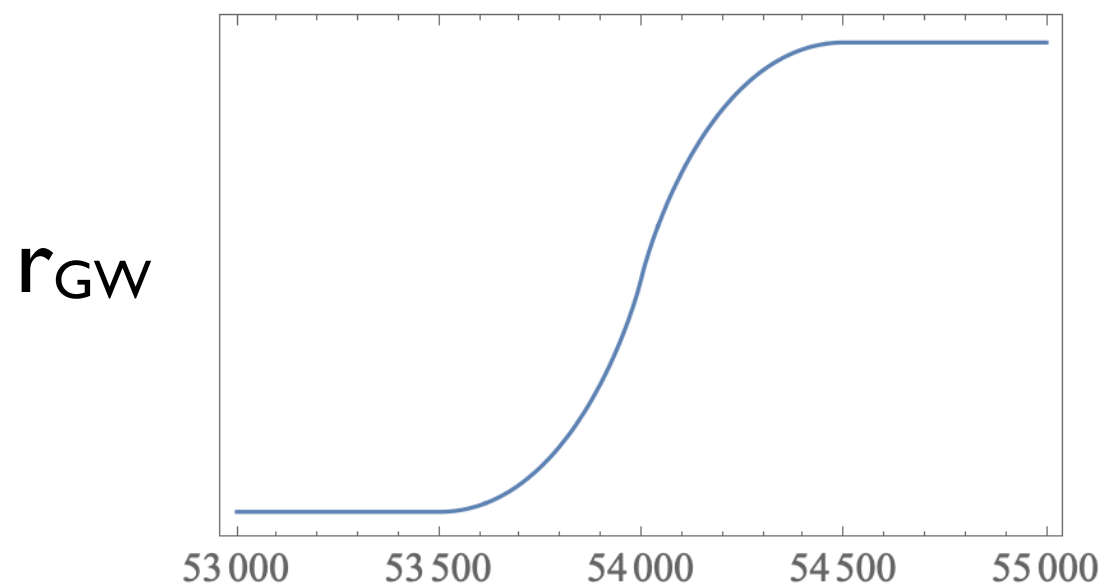


➔  
**Post-fit**

Correct the period and remove the linear trend

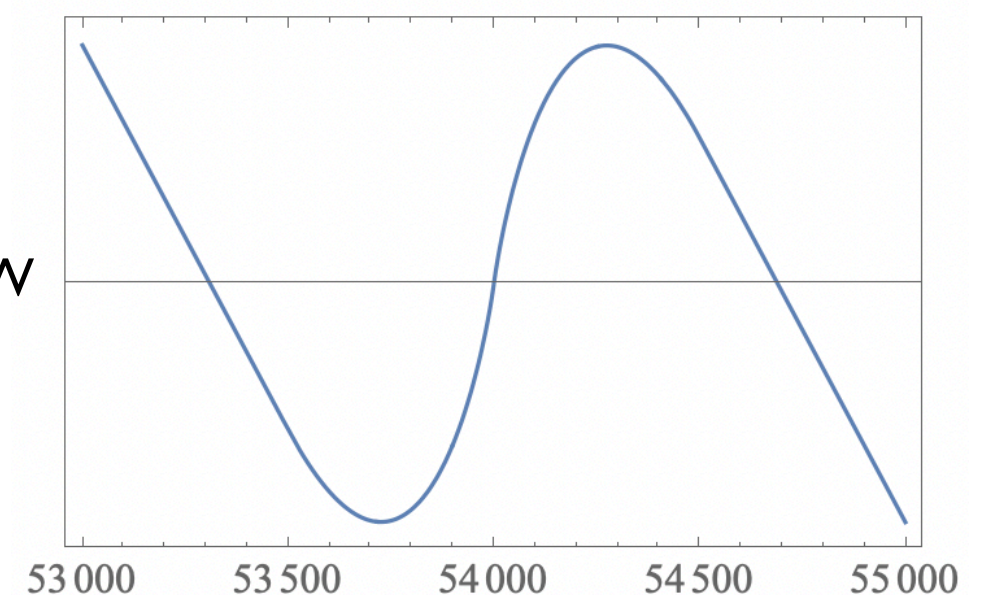


**Timing residual by cosmic string GWs**

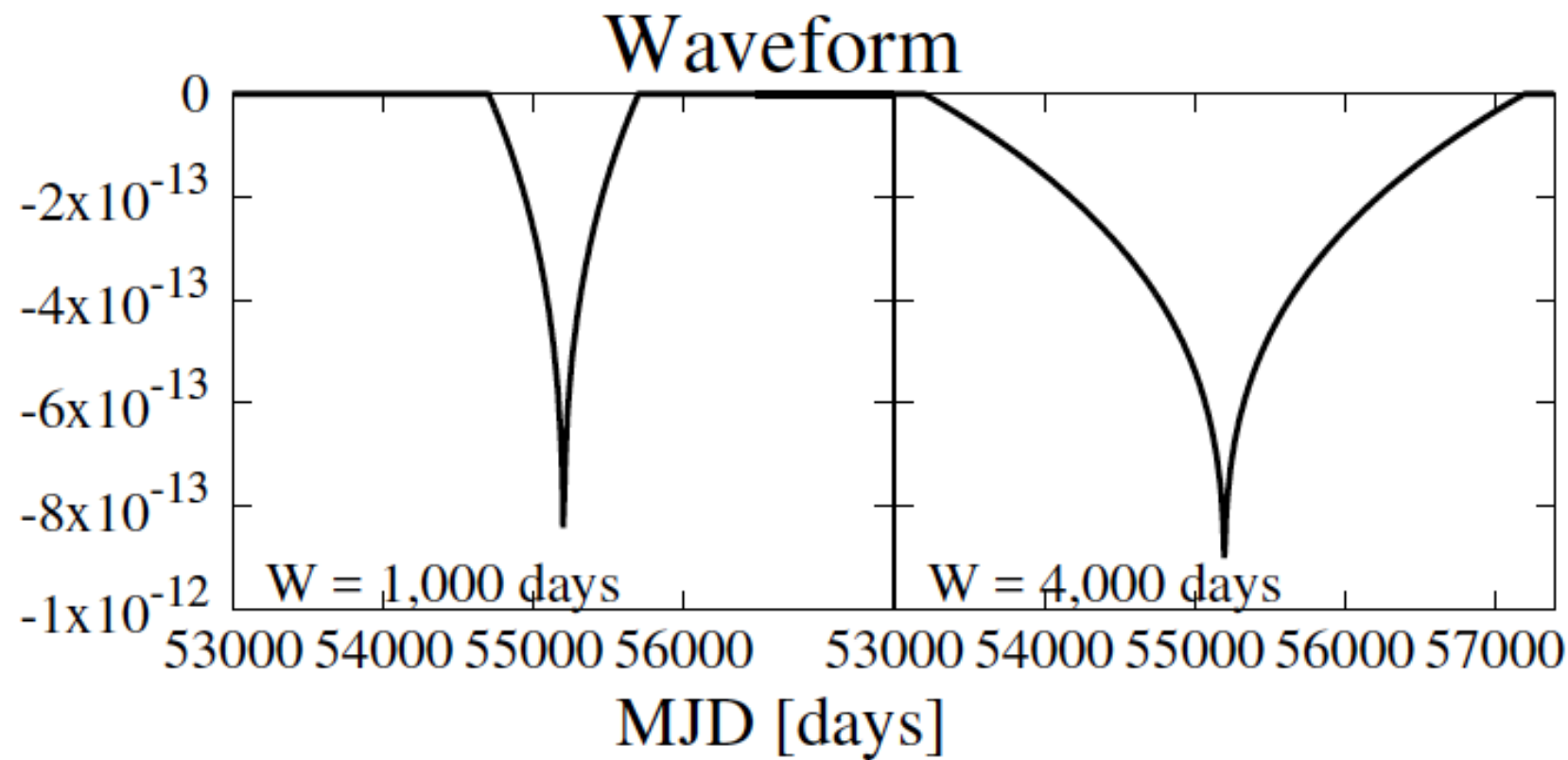


➔  
**Post-fit**

**Post-fit residual**



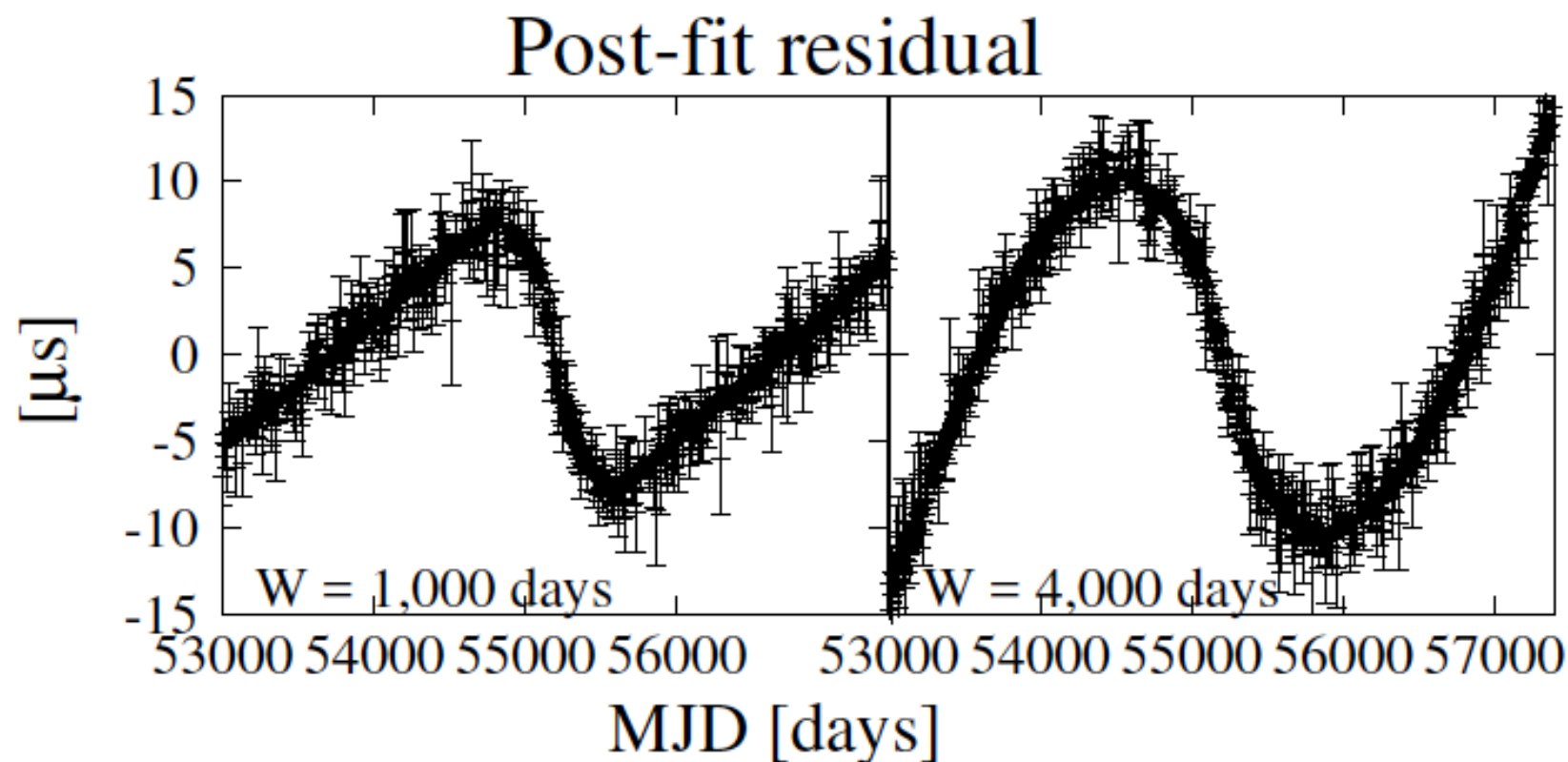
# Simulated timing residual



parameters

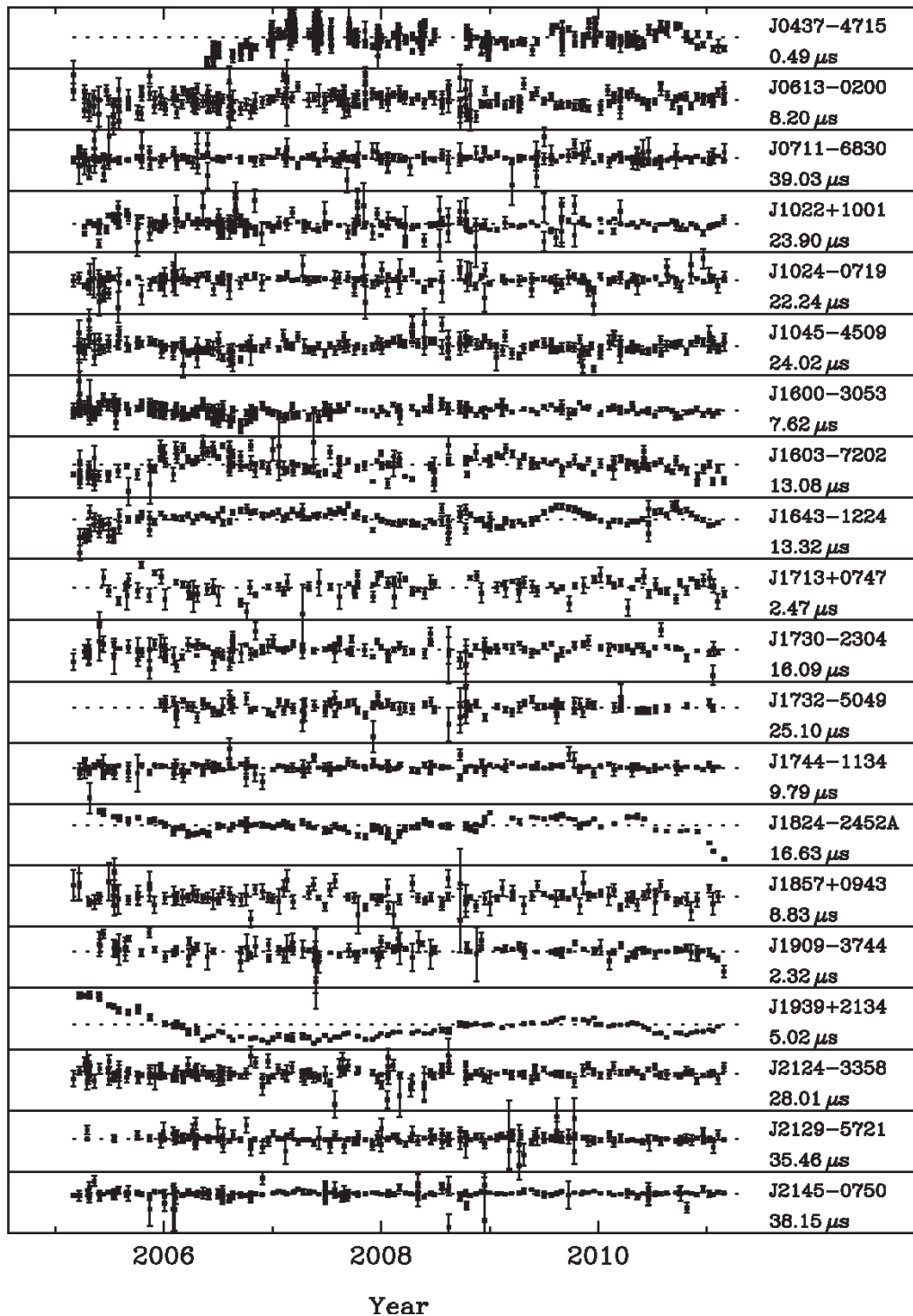
W: width

$t_0$ : epoch

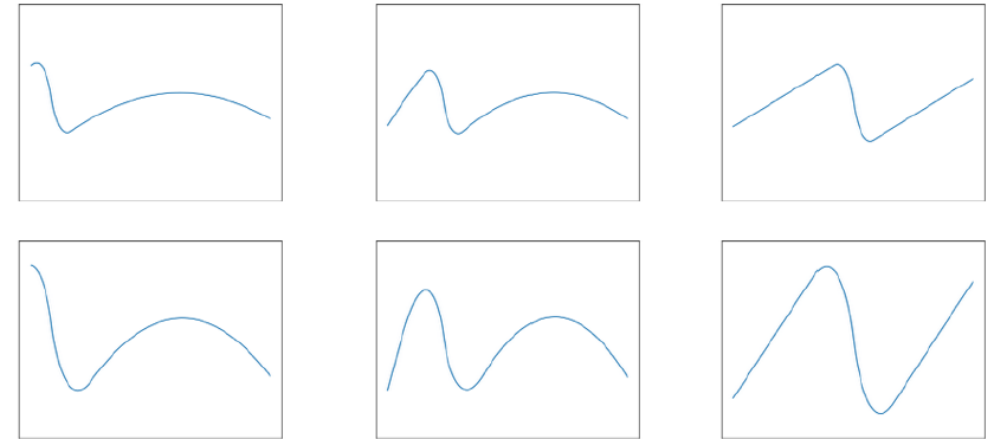


# Template search

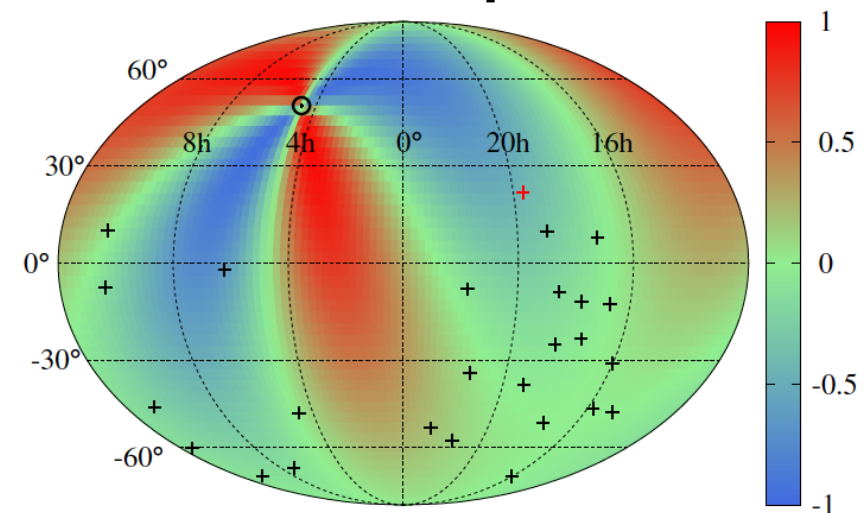
## Data (Parkes PTA DR2)



## Timing residual templates



## Search the quadratic moment pattern



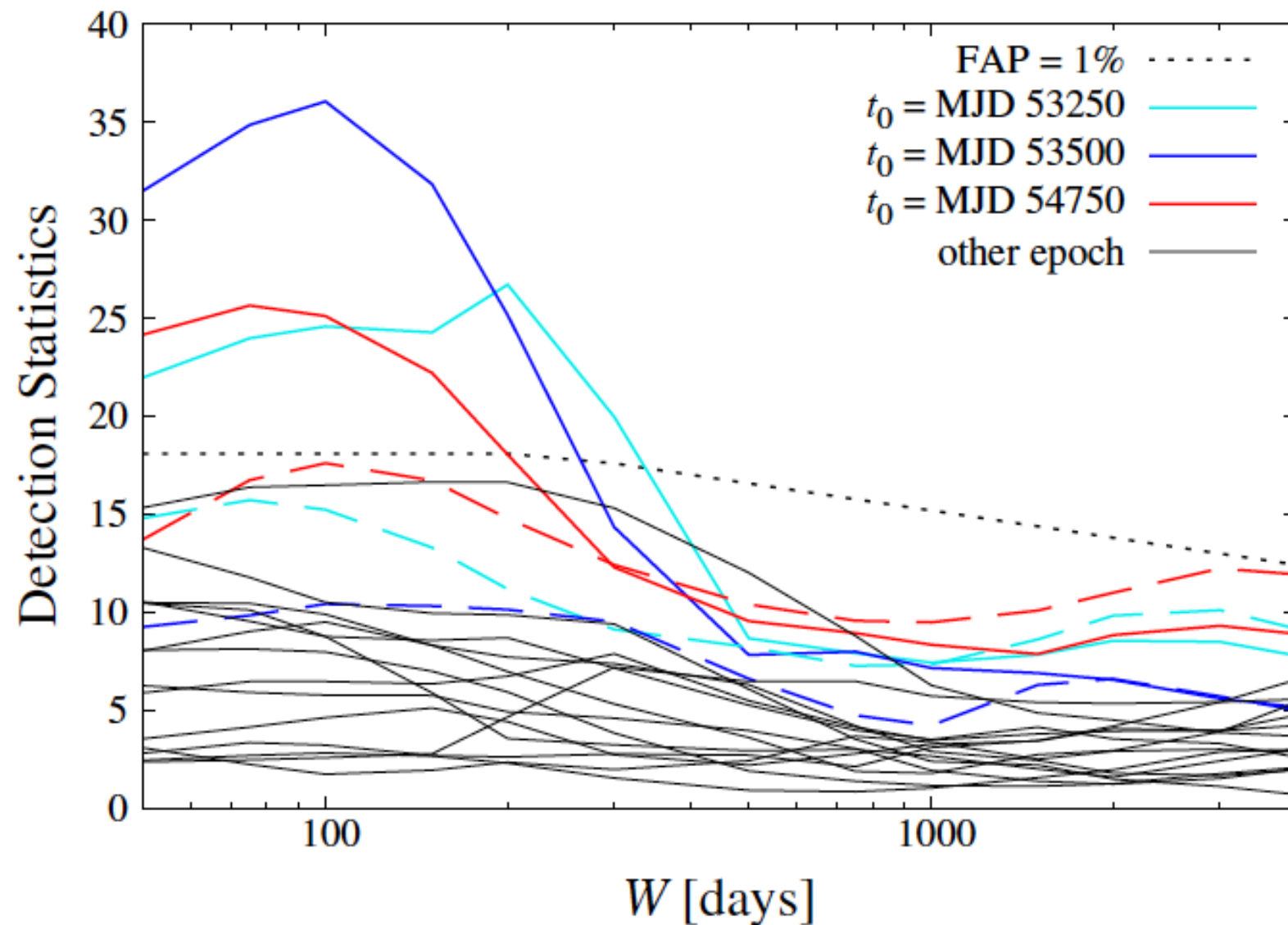
Detection statistics:  $D = \mathbf{A}^t \mathbf{C}_0^{-1} \mathbf{A}$

amplitude  $\mathbf{A} = [A_1; A_2]$

covariance matrix

# Search result

$D_{\max}$  (maximum value over the sky)



Duration of the burst event

parameters

$W$ : width

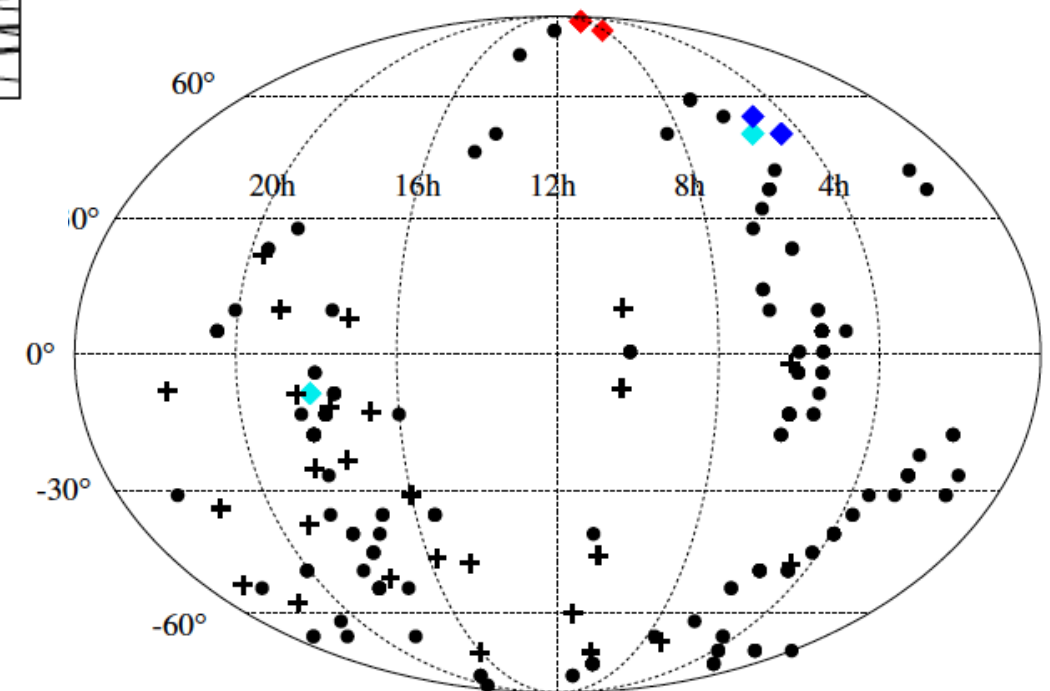
$t_0$ : epoch

- Location of  $D_{\max}$
- + Location of pulsars

◆  $t_0 = \text{MJD53250}$

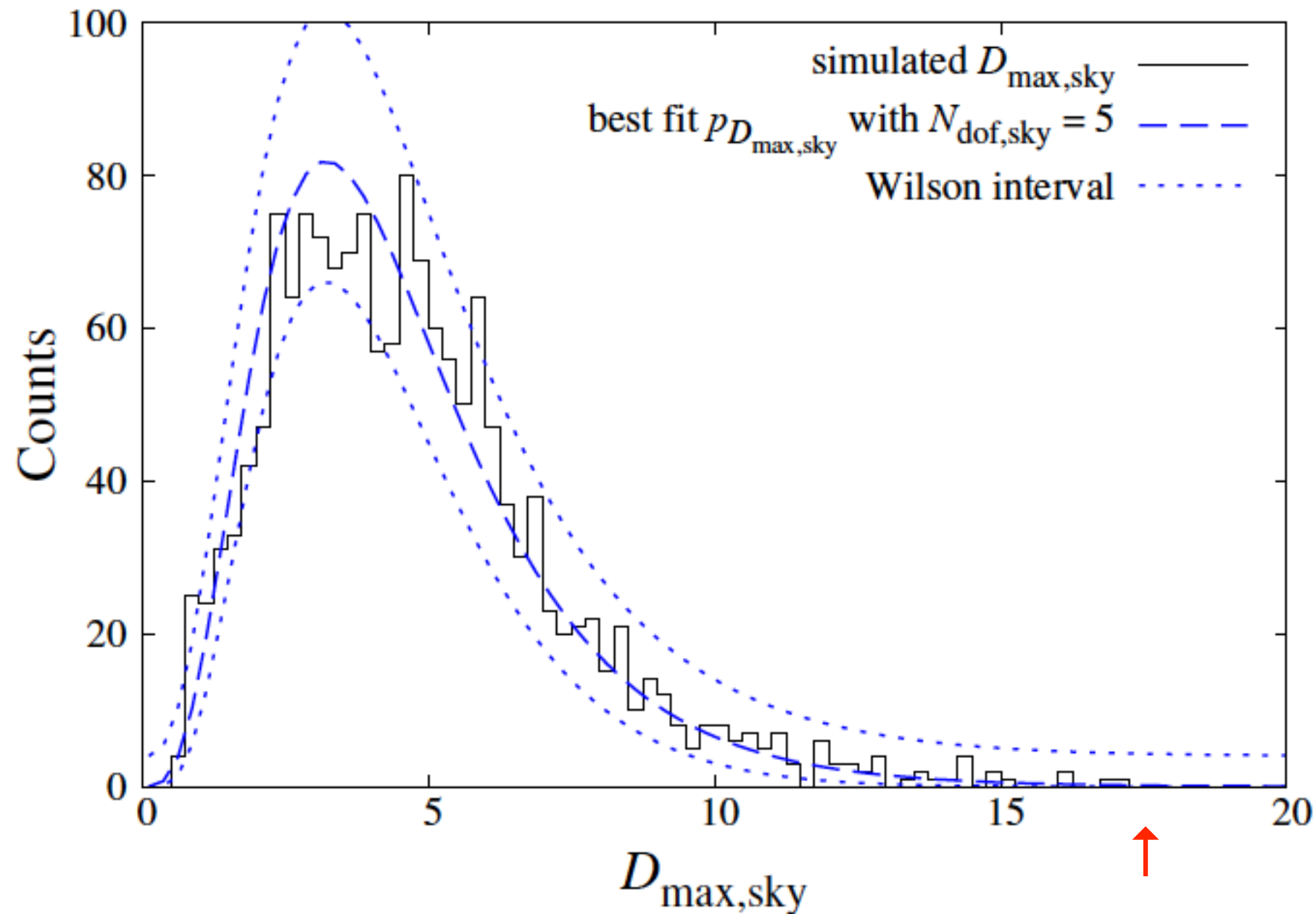
◆  $t_0 = \text{MJD53500}$

◆  $t_0 = \text{MJD54750}$



# Simulation

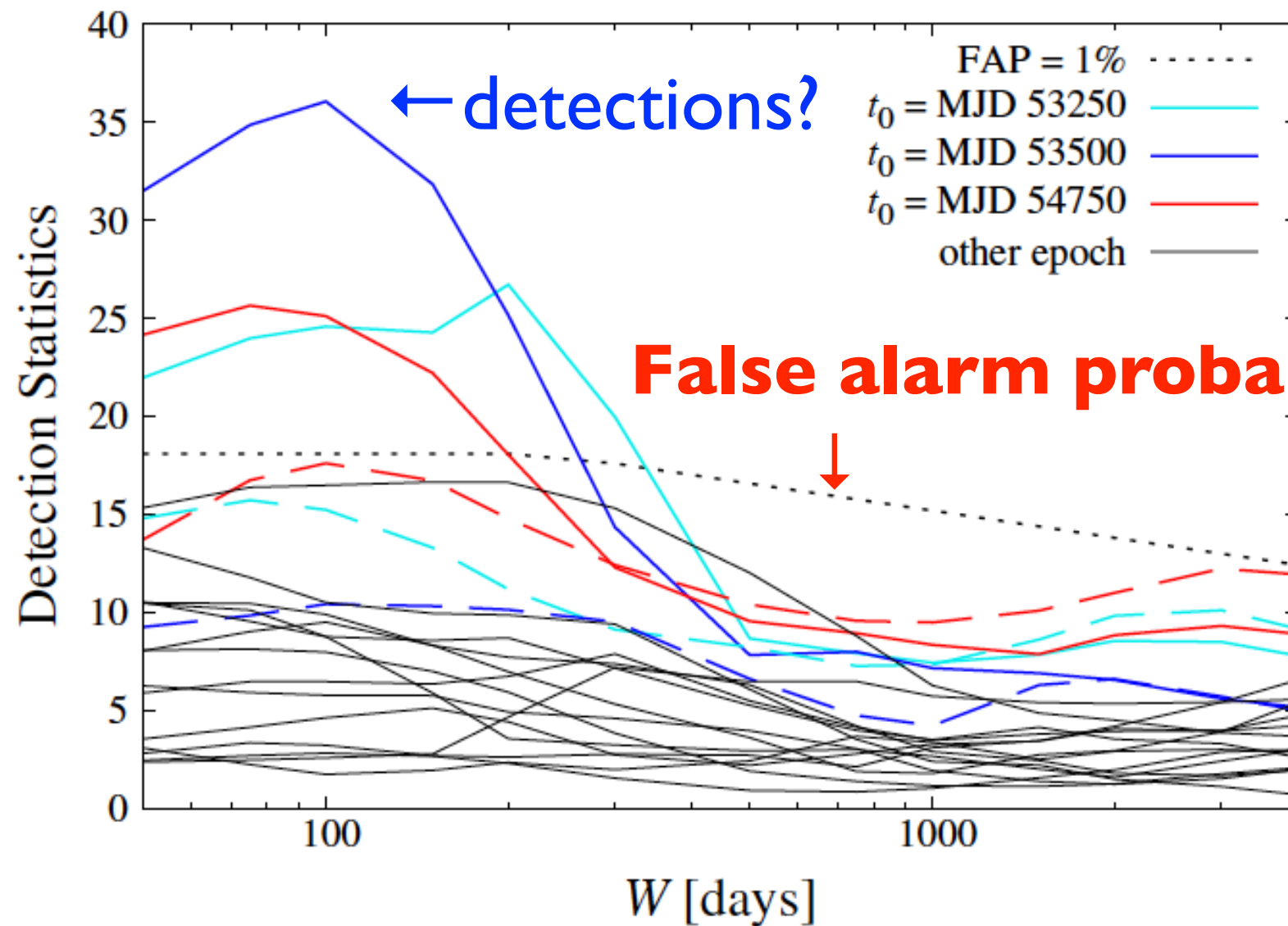
maximum values of  $D$  in simulated PPTA dataset (100 realizations)  
without GW injection



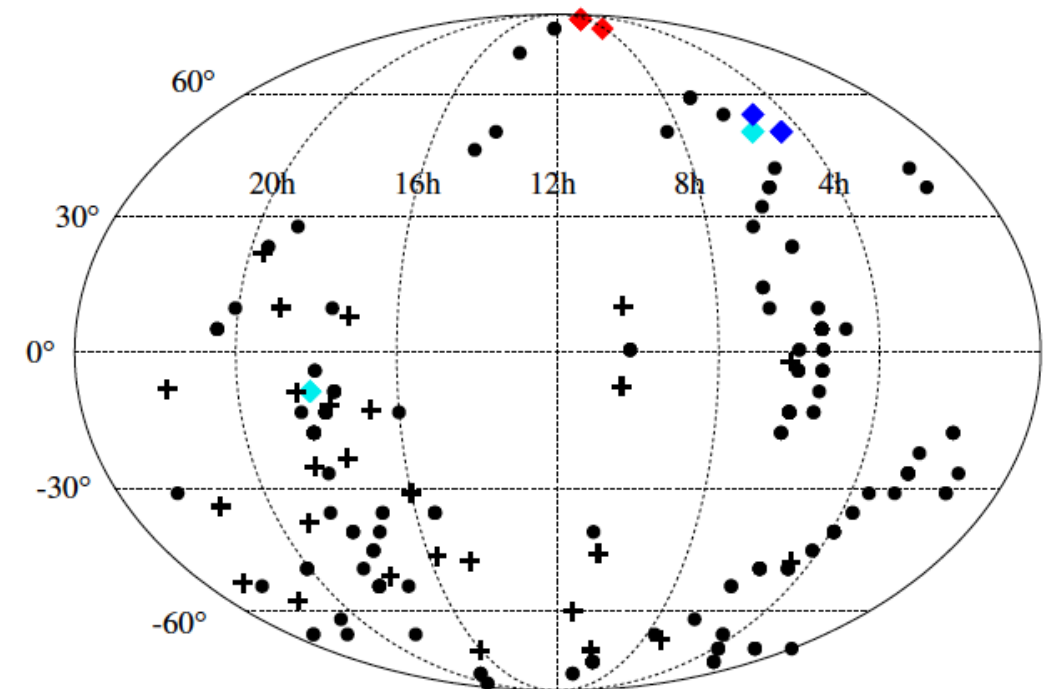
↑  
 $D \sim 17-18$  is possible



# Search result



↑ Detection statistics for different epochs

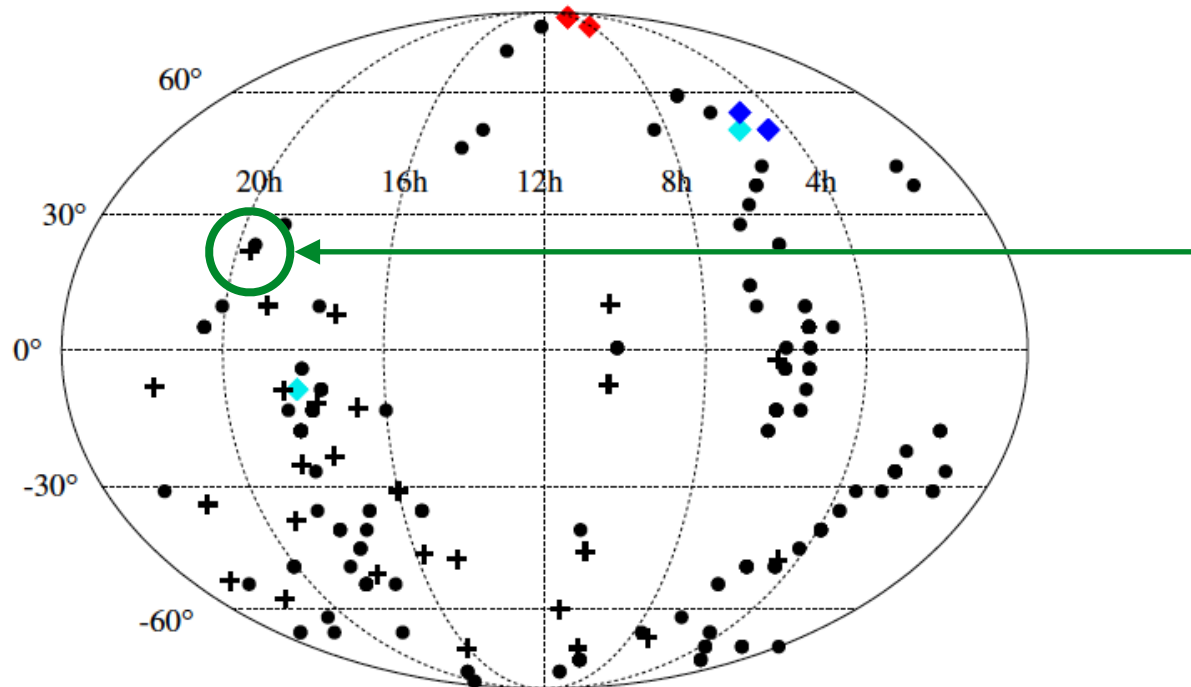


# What is causing high $D_{\max}$ ?

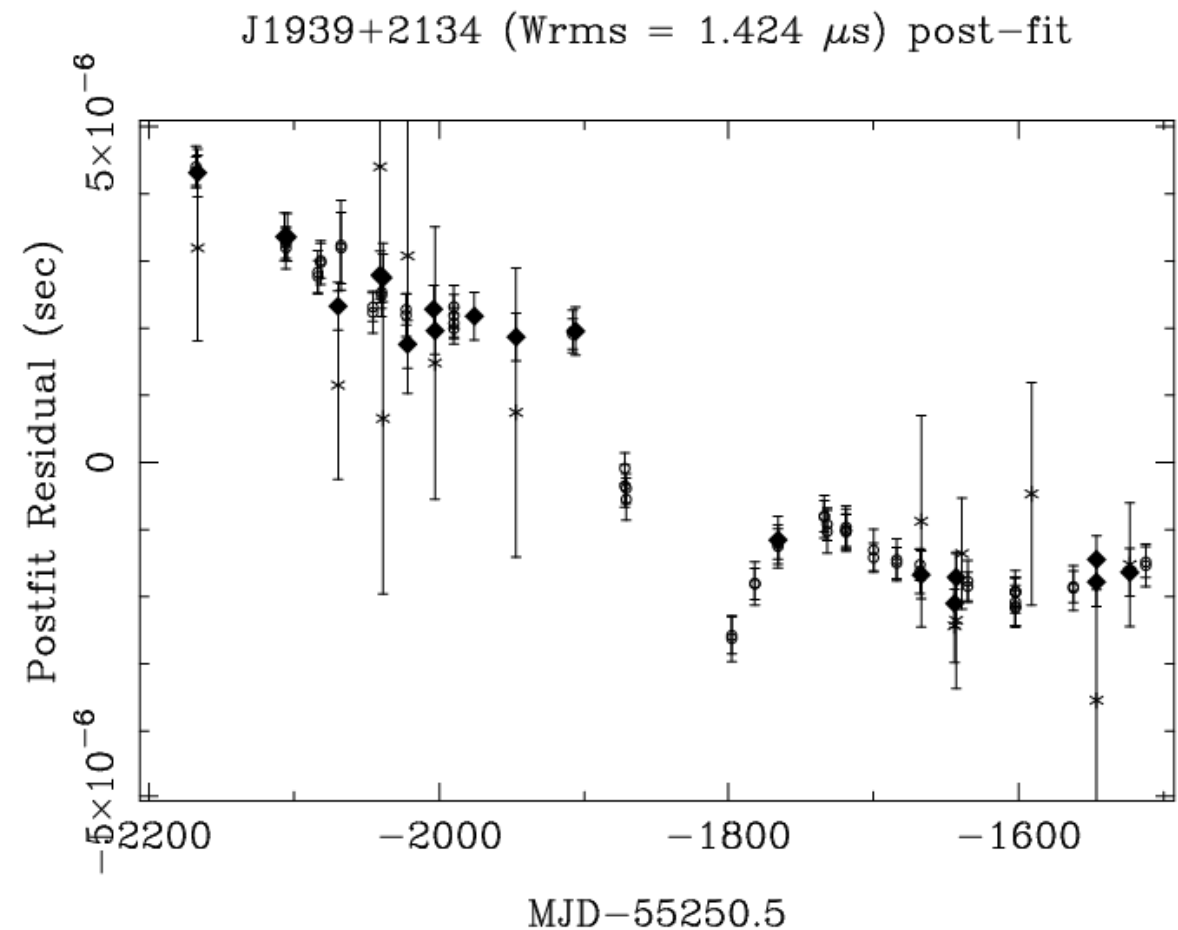
- Location of  $D_{\max}$
- + Location of pulsars

◆  $t_0 = \text{MJD}53250$

◆  $t_0 = \text{MJD}53500$



## Post-fit residual of PSR J1939+2134



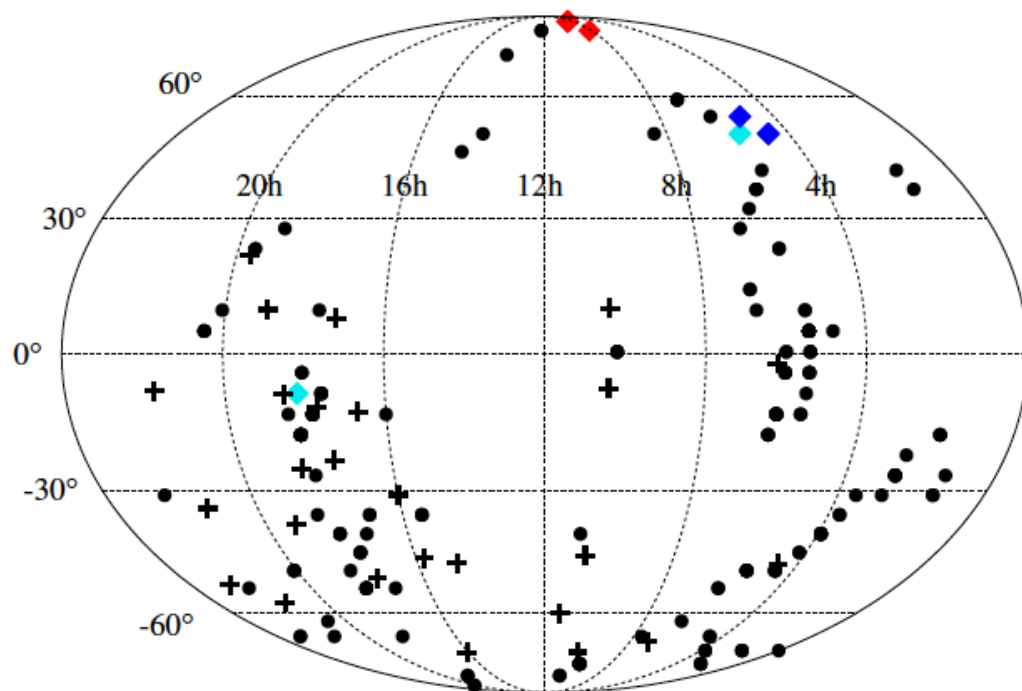
- A poor red noise modeling
  - No other pulsars to confirm
- (International PTA can help in future)



# What is causing high $D_{\max}$ ?

- Location of  $D_{\max}$
- + Location of pulsars

◆  $t_0 = \text{MJD}54750$



Removal of **PSR J1939+2134** does not eliminate the detection at MJD54750

**But the detection is not convincing because...**

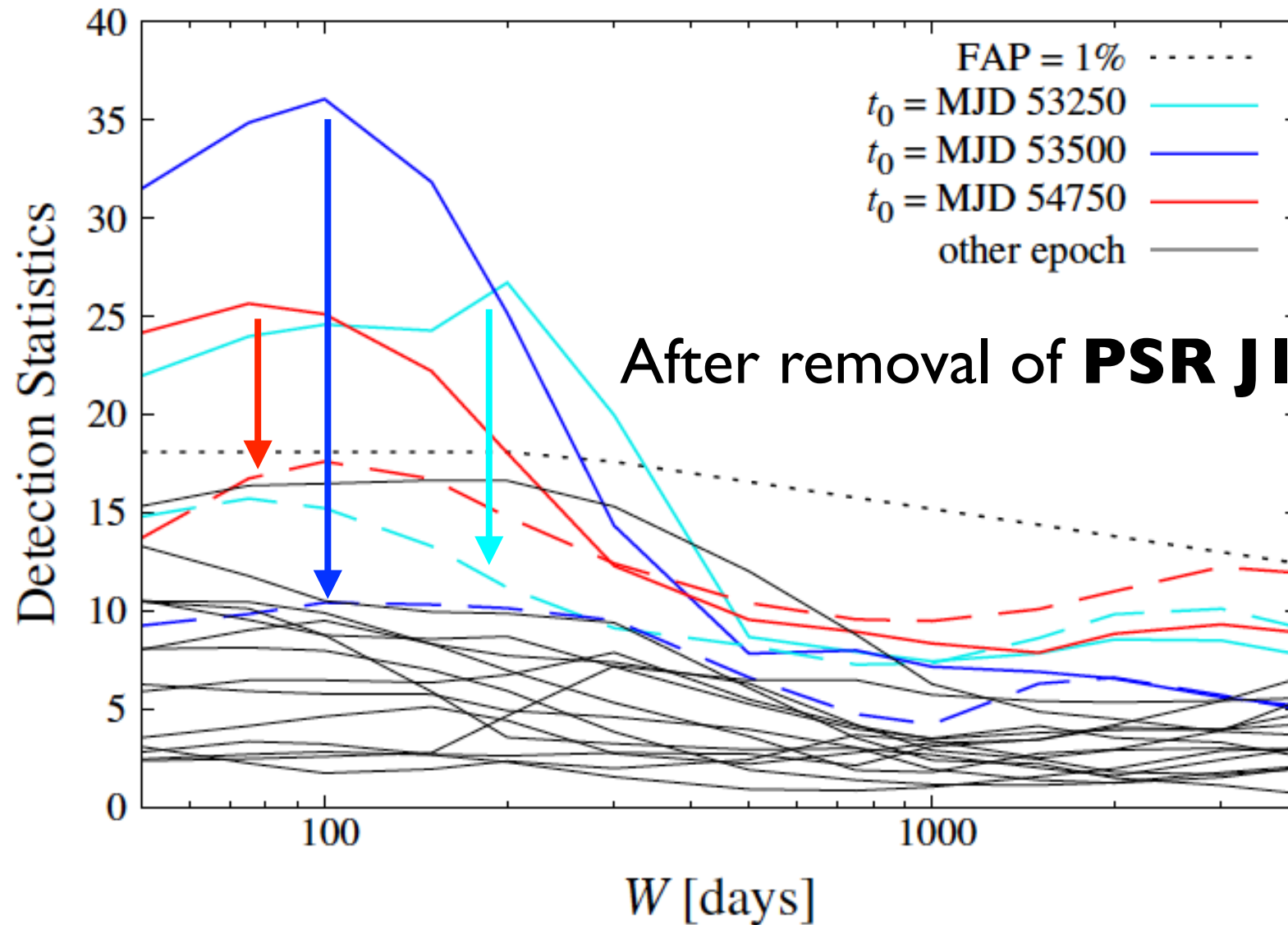
1. We had a **receiver configuration change** at MJD 54751 (11th Oct. 2008)

2. Removal of **PSR J0437-4715** reduces  $D_{\max}$  at MJD 54750

low white noise and high red noise

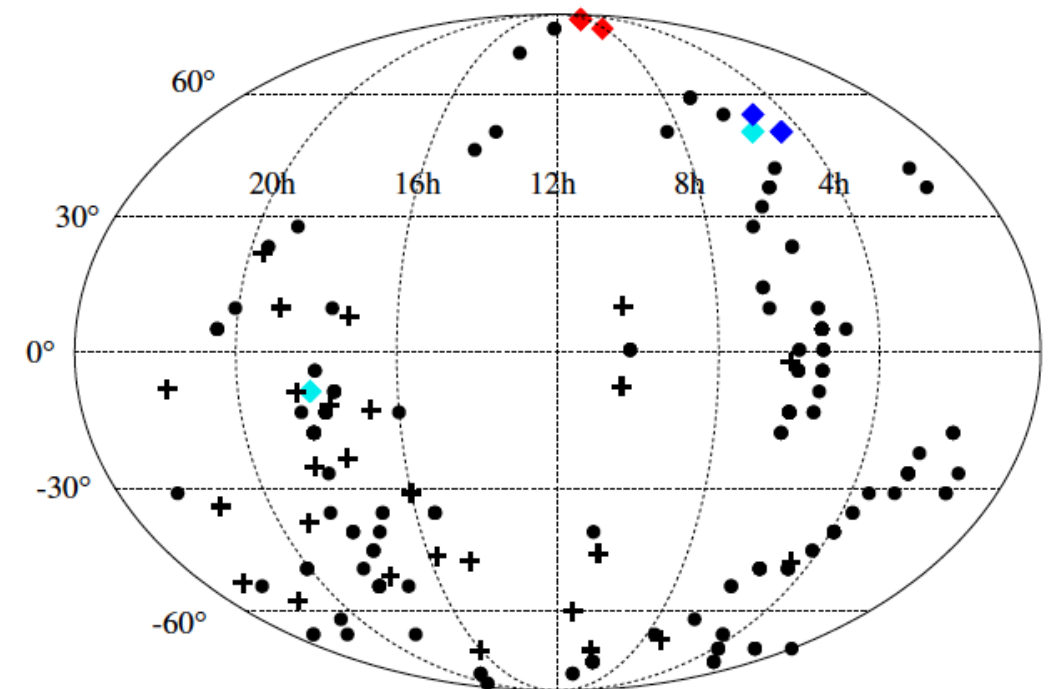
→ among the pulsars most difficult to model

# Search result

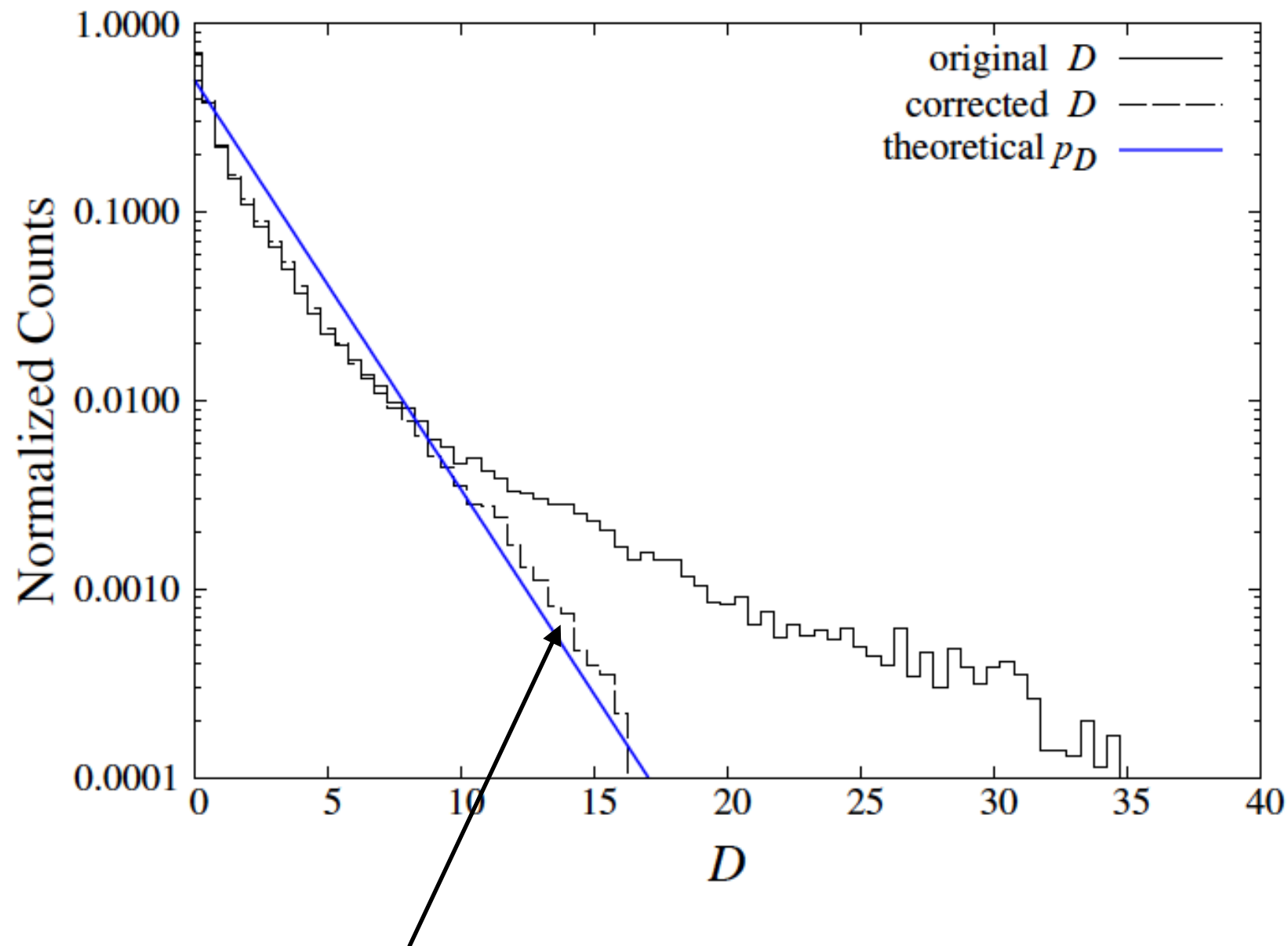


After removal of **PSR J1939+2134** and **J0437-4715**

↑ Detection statistics for different epochs



# Histogram of $D$

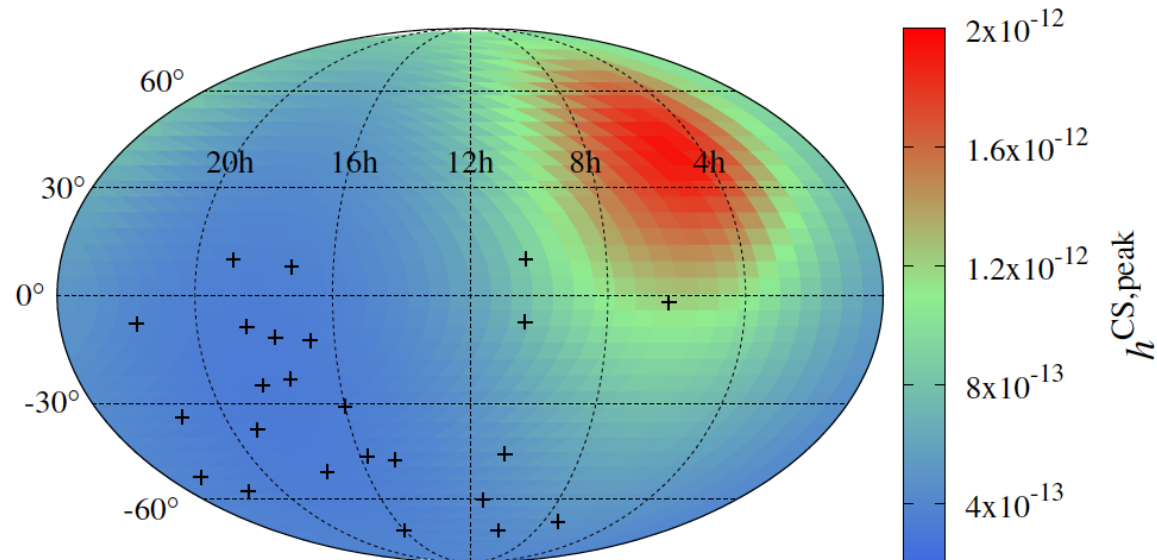


After removal of **PSR J1939+2134** and **J0437-4715**

**Consistent with no detection of GWs**

# Upper bound on GW amplitude

Sensitivity depends on the direction of GWs



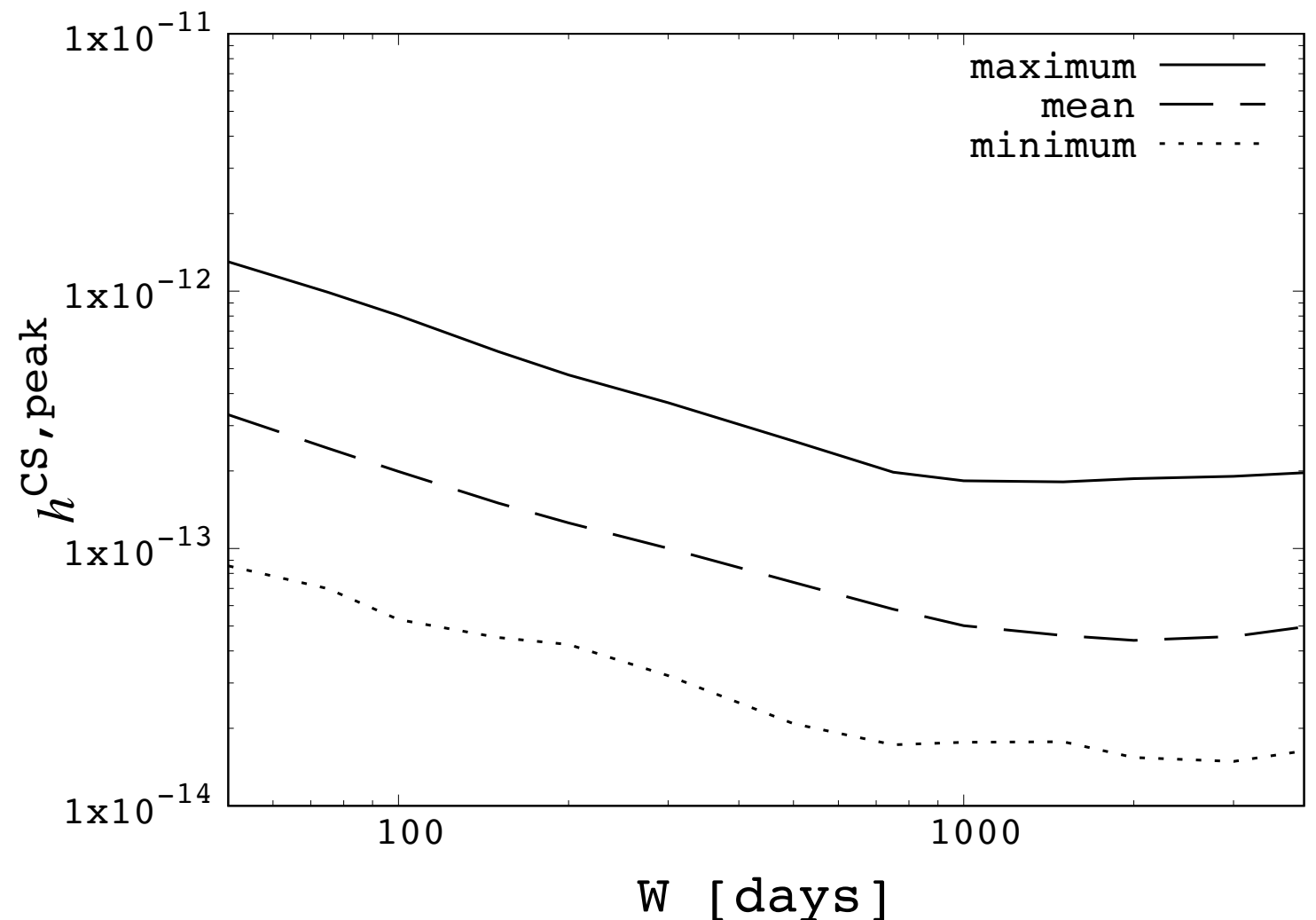
## Sensitivity map

$t_0 = \text{MJD } 54750$

$W = 100$  days

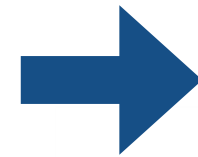
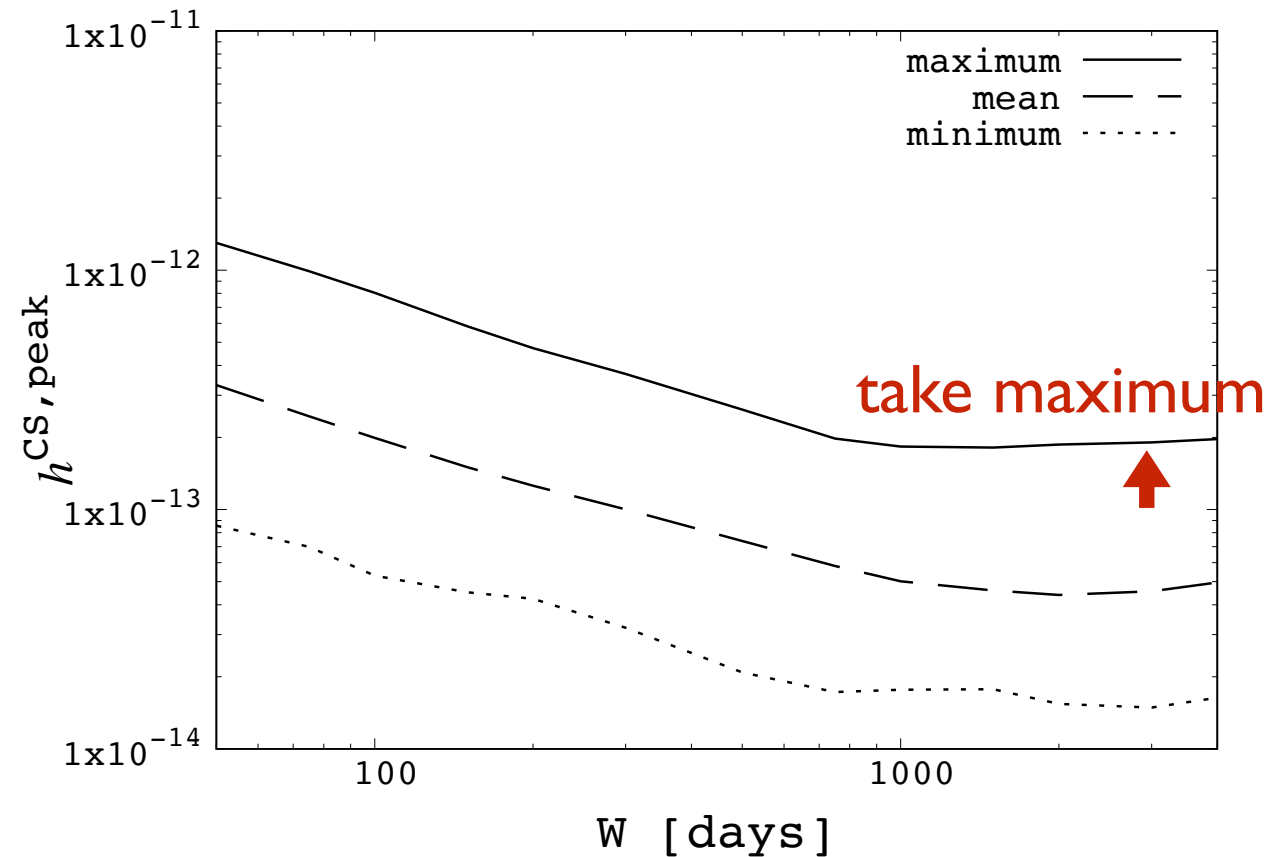
is assumed

Dependence of sensitivity on  $W$   
(averaged over the epoch)

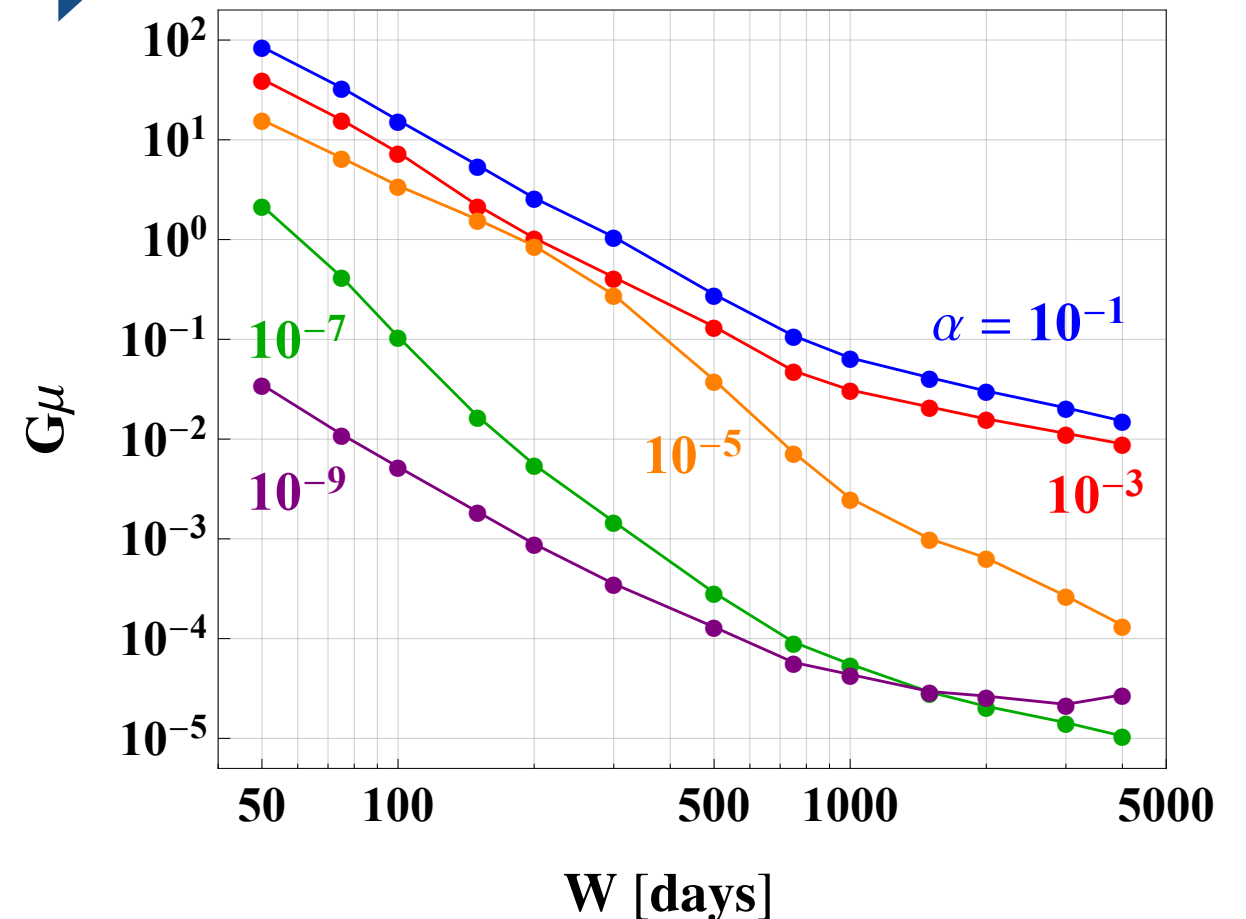


# Constraints on cosmic strings

## Upper bound on GW amplitude



## Constraint on $G\mu$



### GW amplitude (in Fourier space)

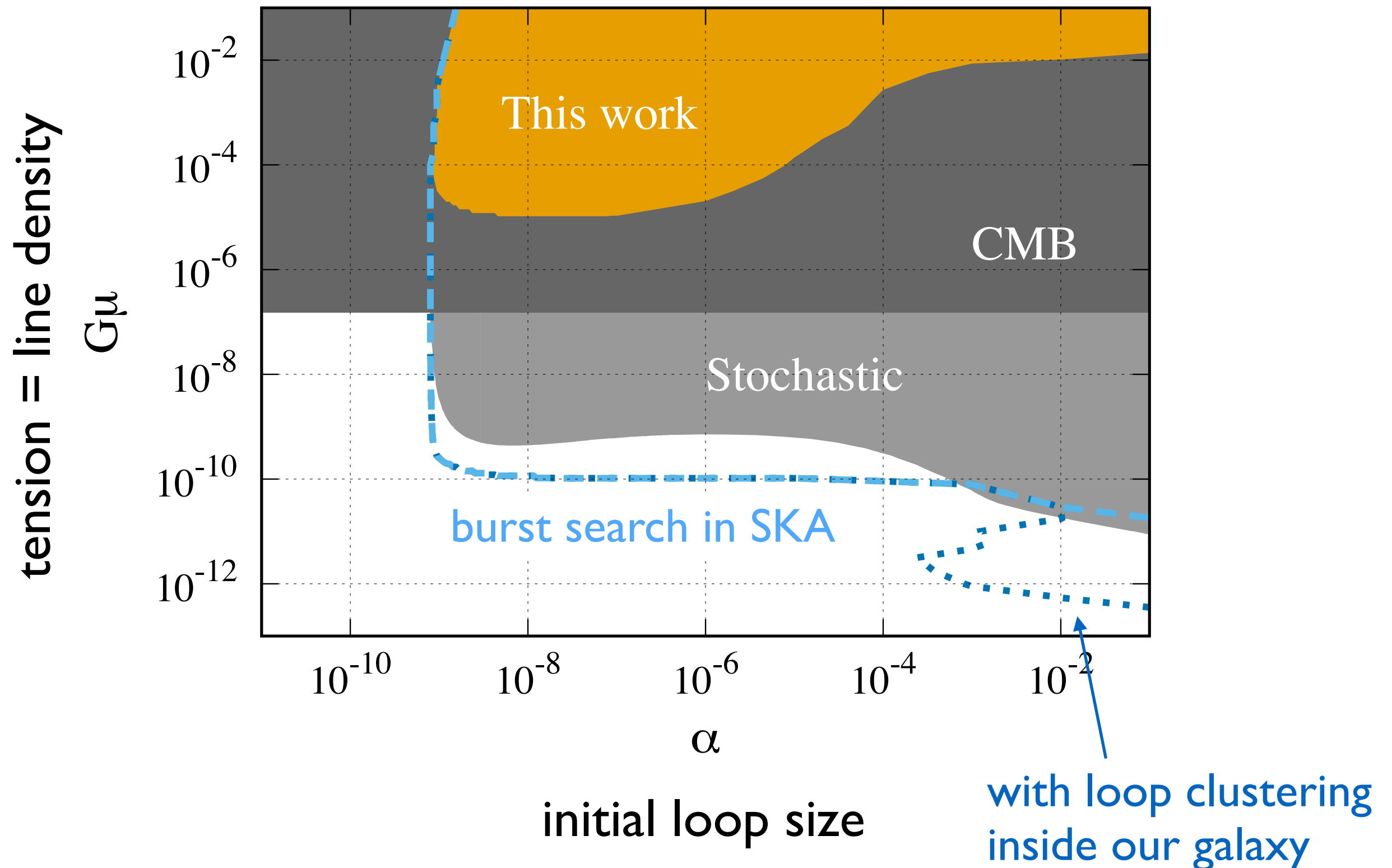
$$\tilde{h}(f) = \frac{G\mu L}{[(1+z)fL]^{1/3} r(z)f}$$

$G\mu$  : tension = line density  
 $r(z)$  : distance to the source

$\alpha$ : initial loop size  
 (changes the loop distribution)

can be obtained by modeling  
 distribution of loops

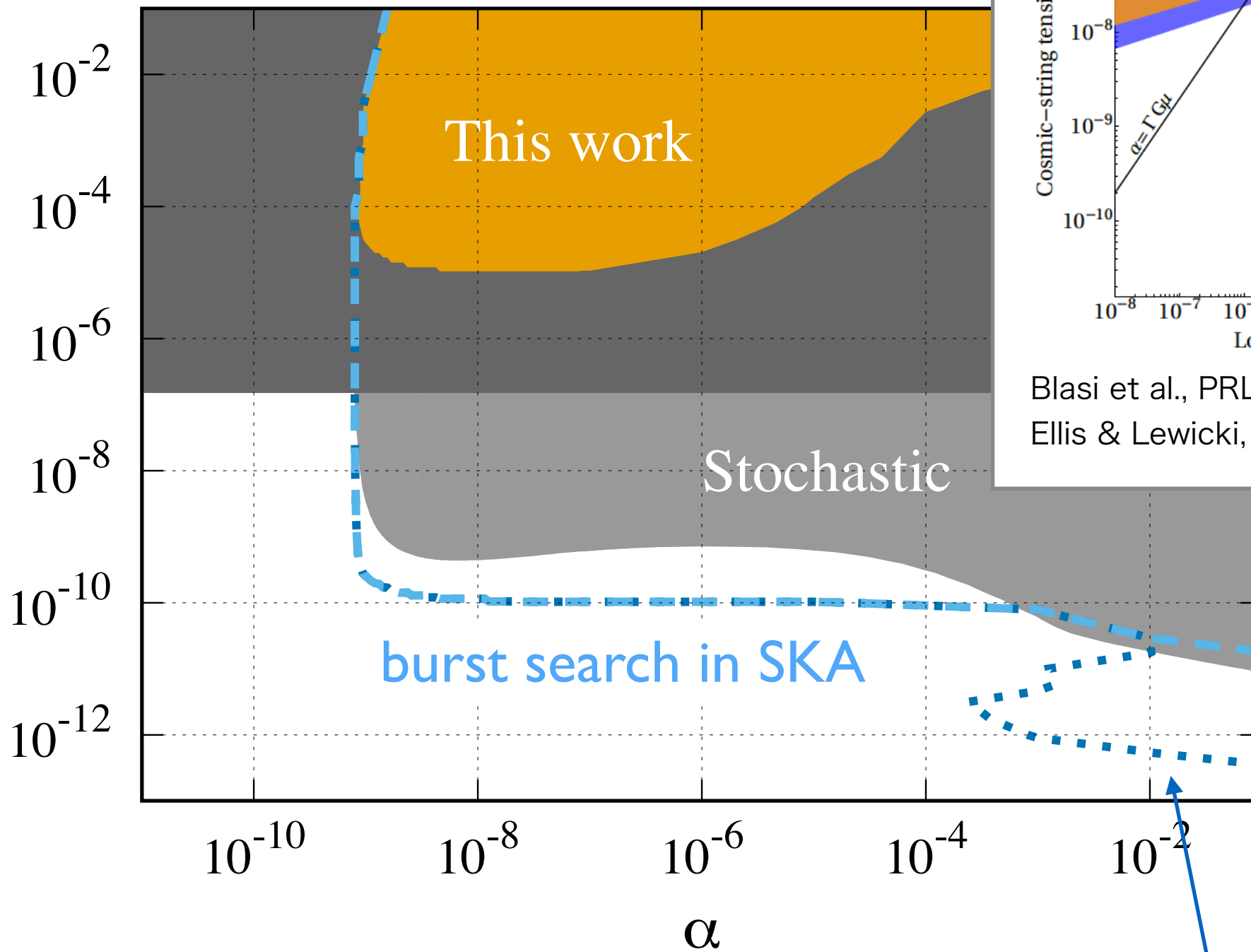
# Constraint on cosmic string parameters



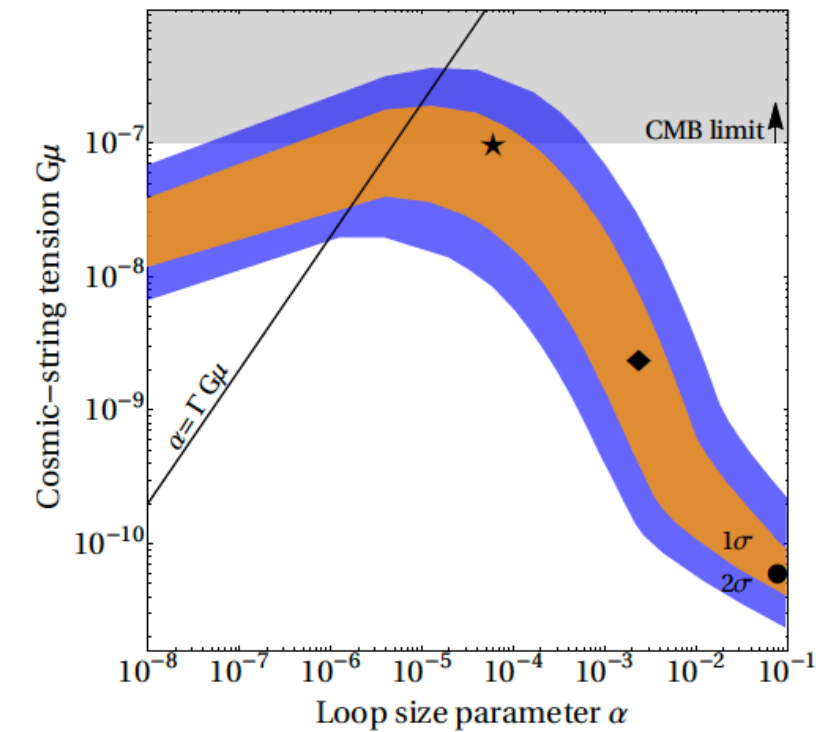
# Constraint on cosmic strin

tension = line density

$G\mu$



Note: NANOGrav 12.5 yr



Blasi et al., PRL 126, 041305 (2021)

Ellis & Lewicki, PRL 126, 041304 (2021)

with loop clustering  
inside our galaxy

# Summary

- We searched for a signal of cosmic string GW burst in **Parkes Pulsar Timing Array** data.
- **No detection of GWs** provided constraints on cosmic string parameters.
- It turned out to be weaker than the other types of observations, but it's independent test of cosmic strings.
- In future, **SKA** will improve the sensitivity.