# Cosmic Inflation after Planck

## Jerome Martin

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VOI

23-(3-2-) (101) 2

 $V\left(\phi\right) = M^{4} \left[1 - 2e^{-2\phi/(\sqrt{6}M_{\rm P})}\right]$ 

°⁄0

V@ Mecontr

L(0)=11-1

 $V(\phi) \equiv M^4 \left| 1 \right|$ 

 $V(\phi) = M^{*}$ 

V(0) = M

V @ M

ar

Gravitational Waves Probes of Physics beyond Standard Model 12-16 July 2021

+ 05

 $\left(\phi | M_{\rm Pl} \right)^2$ 

 $\alpha + (\phi | M_{\rm Pl})^2$ 

V (\$)

 $= M^4 \ln^2 \left( \frac{\phi}{\phi_0} \right)$ 

NP1

Øo

2 m (\$

 $\left[\left(3-\frac{2}{3}\right)^{2}\left(\frac{2}{3}\right)^{2}\left(\frac{2}{3}\right)^{2}\left(\frac{2}{3}\right)^{2}V(\phi) = M^{4}$ 

Ø



## <u>Outline</u>

## □ Observational status of inflation after Planck 2013 & 2015 & 2018

□ Assessing the theoretical status of inflation after Planck

Going beyond the present state of the art ... GW & PBHs from "vanilla" inflationary models

□ Conclusions.



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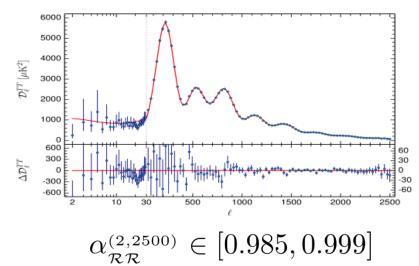
□ Conclusions.

#### Observational status

- Universe spatially flat
- Phase coherence

- Adiabatic perturbations
- Gaussian perturbations
- Almost scale invariant power spectrum
- Background of quantum gravitational waves

$$\Omega_{\kappa} = -0.011^{+0.013}_{-0.012}$$



$$f_{_{\rm NL}}^{\rm loc} = 0.8 \pm 5$$

 $n_{\rm S} = 0.9649 \pm 0.0042$ 

r < 0.056



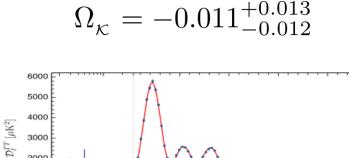
#### Observational status

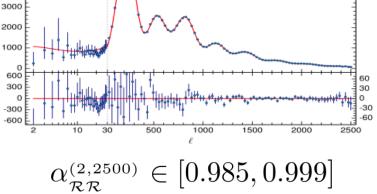
- Universe spatially flat
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Single field slow-roll models, with minimal kinetic terms, are preferred

 $\Delta D_{\ell}^{TT}$ 





$$f_{\rm \scriptscriptstyle NL}^{\rm loc}=0.8\pm5$$

 $n_{\rm S} = 0.9649 \pm 0.0042$ 



#### Which models?



Displayed Models: 193/193 VHI. VHI -1 -21  $\ln(\mathcal{E}/\mathcal{E}_{\mathrm{HI}})$ -40 No unconstrained parameter -50 Model performance (Bayesian <sup>•</sup>ste Evidence) Number of unconstrained parameters STE<sub>1</sub>  $STE_{1-2-3}$ STE1\_ STE, STE<sub>2-3</sub> STE<sub>3</sub> Displayed Models: 151/193 KKLTL. KKLT -1.1 -2.5 GMSSMI  $\ln({\cal E}/{\cal E}_{\rm HI})$ GMSSMI IMI1 Number of unconstrained parameters

 $STE_{1-2}$ 

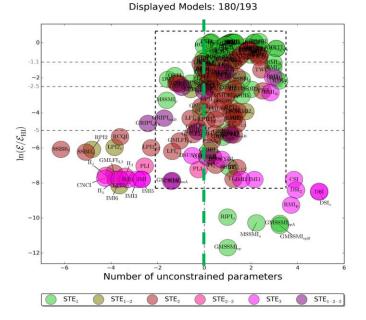
STE,

STE<sub>2</sub>

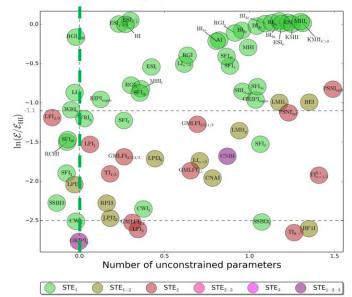
STE<sub>2-3</sub>

STE<sub>3</sub>

 $\mathsf{STE}_{1-2-3}$ 

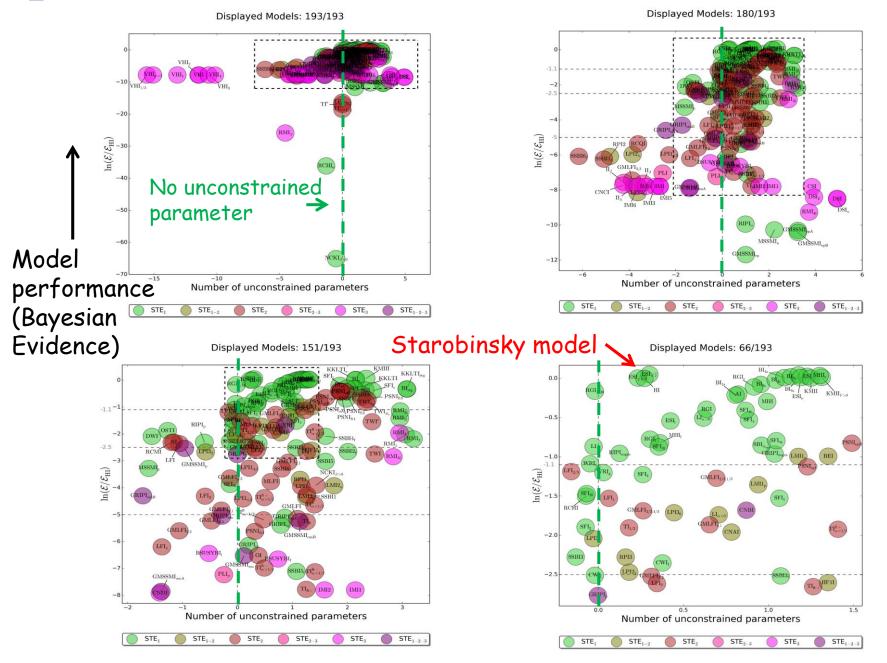


Displayed Models: 66/193



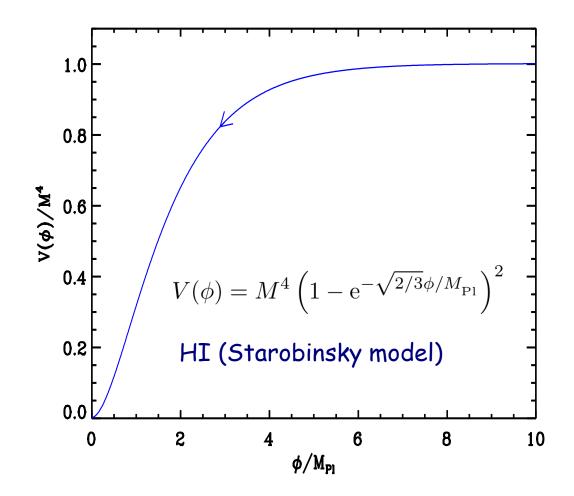
#### Which models?





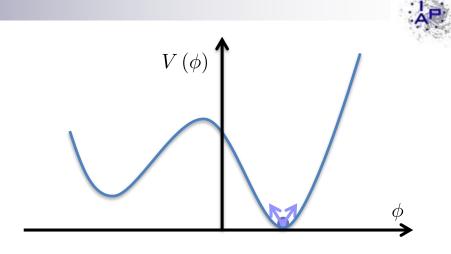


Plateau inflationary models are the winners! ... examplified by the Starobinsky model



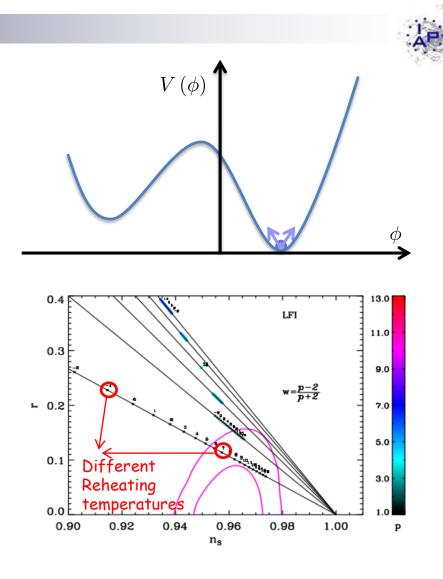
J. Martin, C. Ringeval and V. Vennin, Phys. Dark Univ. 5-6 (2014) 75, arXiv:1303.3787 J. Martin, C. Ringeval, R. Trotta and V. Vennin, JCAP 1403 (2014) 039, arXiv1312.3529

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- Changing reheating duration moves the observational window along the inflaton potential

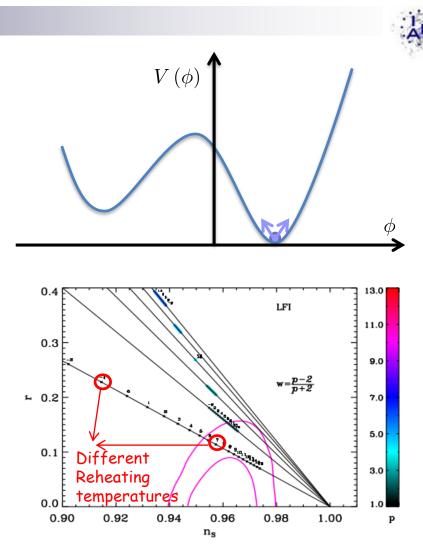


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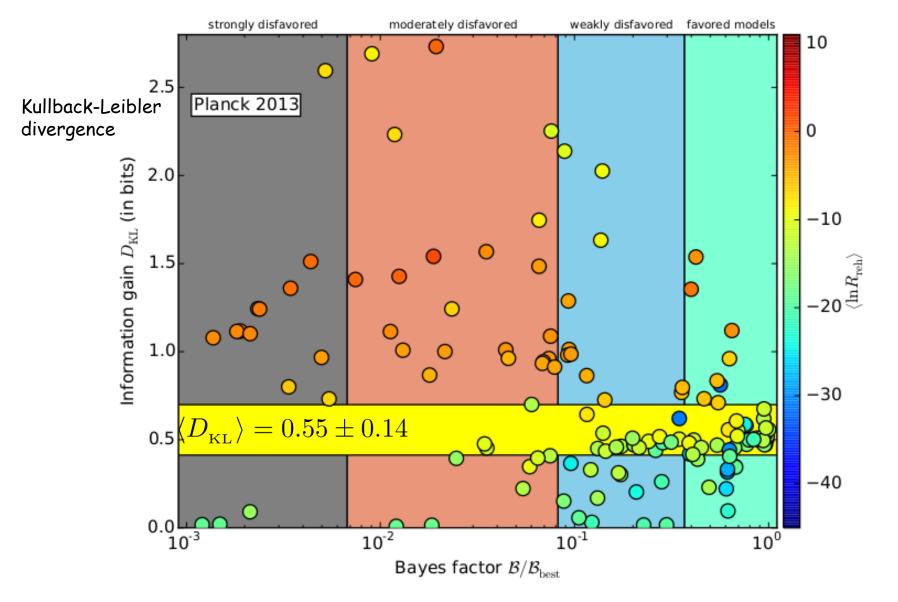
- Reheating can be very complicated (as the process of re-ionization) but, as long as the CMB is concerned, only the reheating parameter is important

$$\ln R_{\rm rad} = \frac{1 - 3\overline{w}_{\rm reh}}{12 + 12\overline{w}_{\rm reh}} \ln \left(\frac{\rho_{\rm reh}}{\rho_{\rm end}}\right)$$



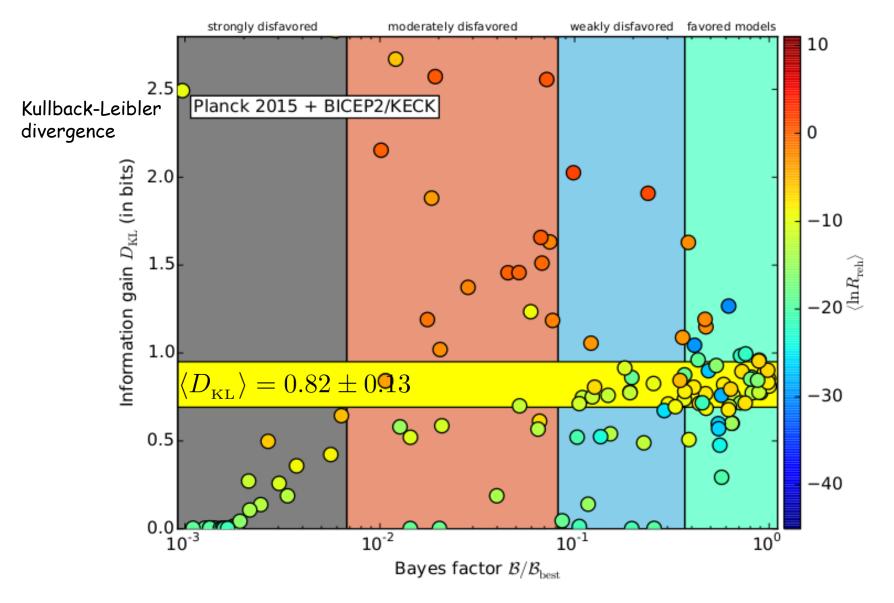
J. Martin and C. Ringeval, Phys. Rev. D82 (2010) 023511, arXiv:1004.5525





J. Martin, C. Ringeval and V. Vennin, Phys. Rev. Lett. 114 (2015) 8, 081303, arXiv:1410.7958





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- More precisely, it rests on the fundamental principle that, in GR, every form of energy weighs including pressure

$$\frac{\ddot{a}}{a} = -\frac{1}{6M_{\rm Pl}^2} \left(\rho + 3p\right)$$

Also at play in: - dark energy

- Neutron stars
- BBN



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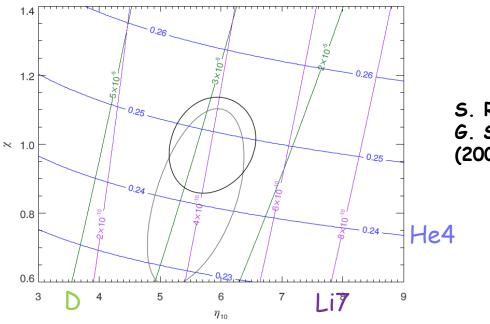
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S. Rappaport, J. Schwab, S. Burles, and G. Steigman, Phys. Rev. D77, 023515 (2008), arXiv:0710.5300

## This is similar to the Schwinger effect: interaction of a quantum field with a classical source

J. Martin, Lect. Notes Phys. 738 (2008), 195 arXiv:0704.3540

#### Schwinger effect

- Electron and positron fields
- Classical electric field
- Amplitude of the effect controlled by E

Inflationary cosmological perturbations

- Inhomogeneous gravity field
- Background gravitational field: scale factor
- Amplitude controlled by the Hubble parameter H

See also dynamical Schwinger effect, dynamical Casimir effect etc ...



## Theoretical open issues

- Model building & physical nature of the inflaton field:
  - Has inflation some connections with the Higgs?
  - Extensions to the standard model contains mainy degrees of freedom ... how a single field description emerges from that?
  - Can we observe the effect of these extra degrees of freedom (multifield inflation, gauge fields, inflationary magnetogenesis, PBHs etc ...)
- How does inflation end?
  - Reheating/preheating/thermalization, PBHs ...
- Inflationary mechanism for structure formation
  - Quantum-to-classical transition of inflationary perturbations: inflation is the only known system that uses GR and QM and where high-accuracy data are available



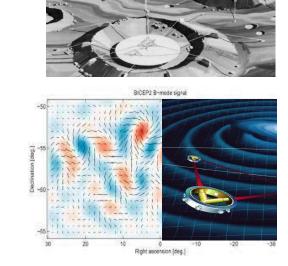
## **Observational tools**

- Non-Gaussianity

- Gravitational Waves

<u>Remark:</u> these probes can be used to study inflation beyond simple models but they can also be useful (although challenging) to further test the vanilla scenarios.











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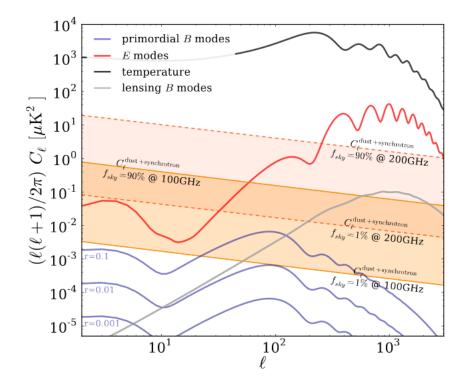


## <u>1- the presence of a "quantum" GW background from inflation</u>



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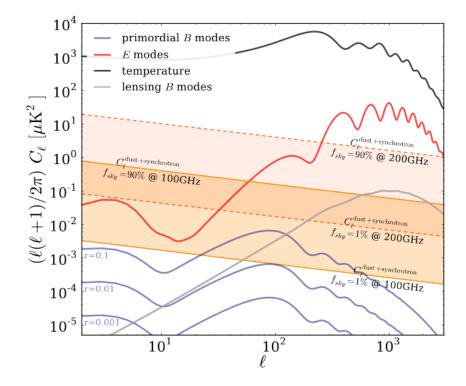
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- Next generation of CMB mission with a target: r ~  $10^{-4}$  [Starobinsky model, r ~ (2-4) x10<sup>-3</sup>, Planckian excursion r~  $10^{-3}$ ]



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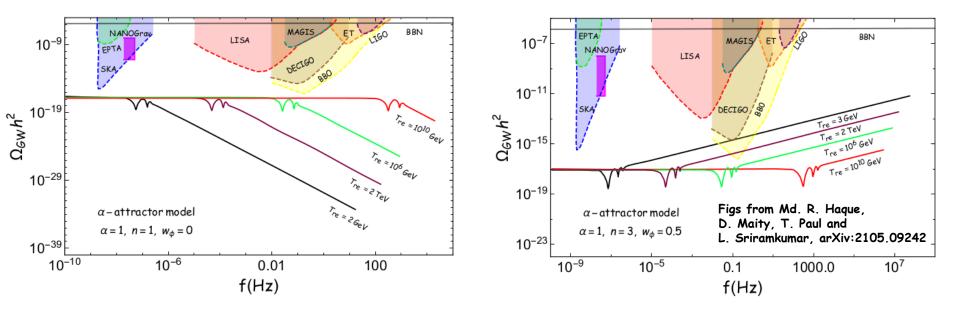
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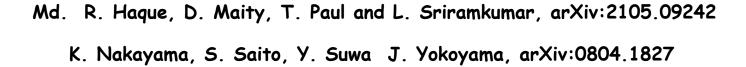
- What can be learned from a detection?
- Energy scale of inflation
- $f \square$  Final proof of vanilla inflation: consistency check (but needs n\_T)  $n_{
  m T}=-r/8$
- Measurement of the first derivative of the potential
- □ Field excursion
- Greatly improve model selection
- Greatly improve constraints on reheating



## 1- the presence of a "guantum" GW background from inflation

- On small scales can be revealed (maybe?) by direct detection

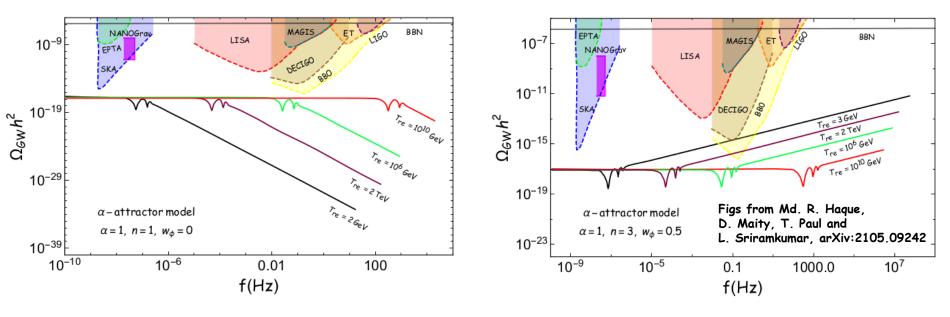






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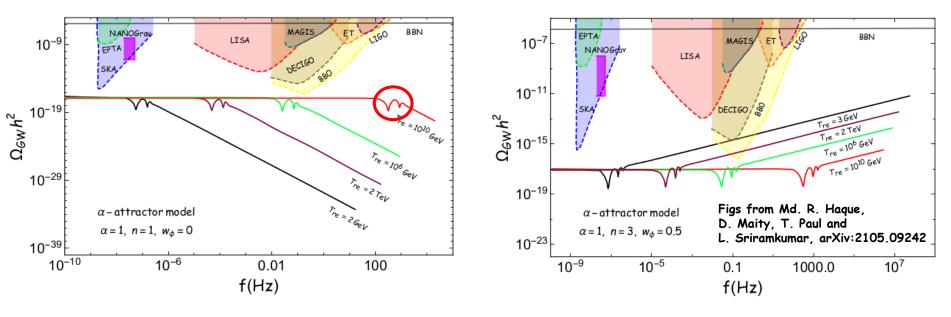


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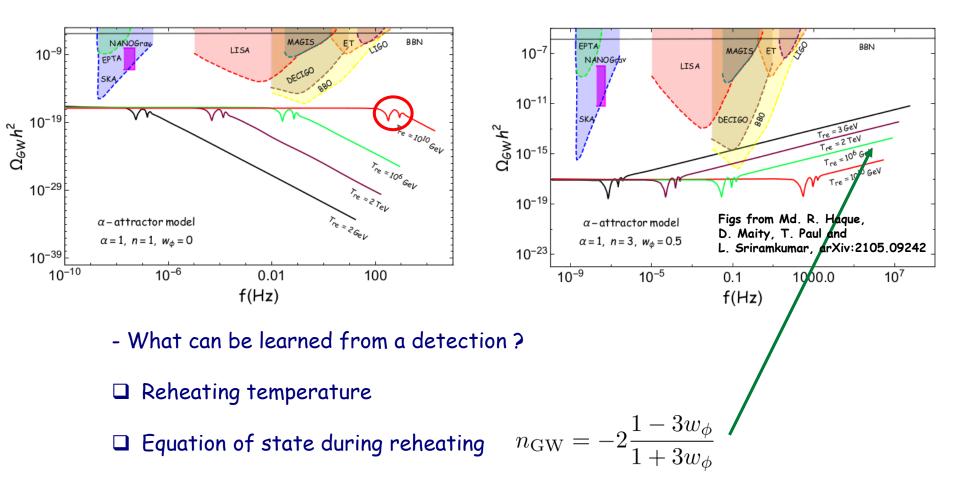
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Reheating temperature



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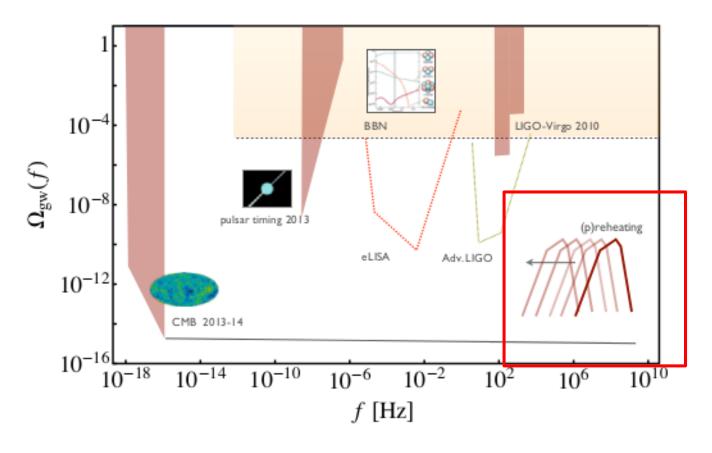


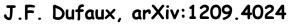


## 2- the presence of a GW background from background preheating



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D. Figueroa & F. Torrenti, arXiv:1707.04533

3- GW and ultra-lights PBHs from metric preheating

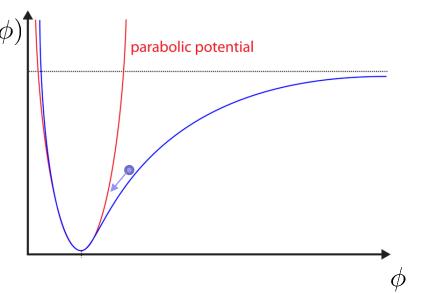


Other key predictions of vanilla inflation

3- GW and ultra-lights PBHs from metric preheating

 $\hfill\square$  In the vicinity of the minimum, any  $V(\phi)$  potential can be written

$$V(\phi) = \frac{1}{2}m^2\phi^2 + \cdots$$





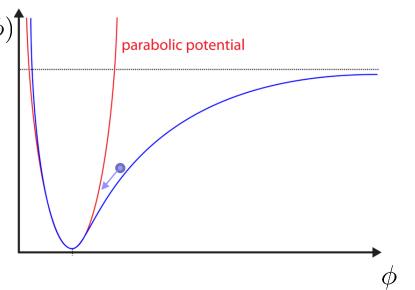
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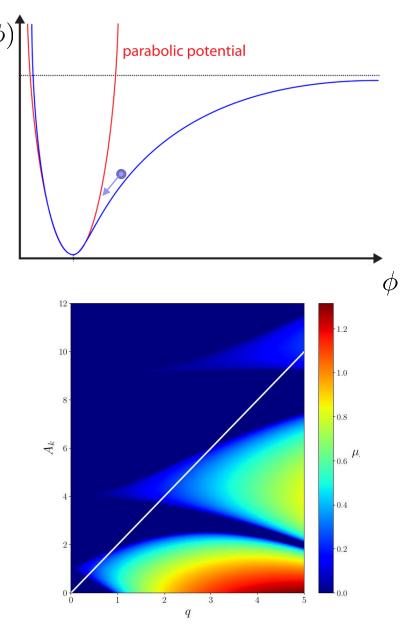
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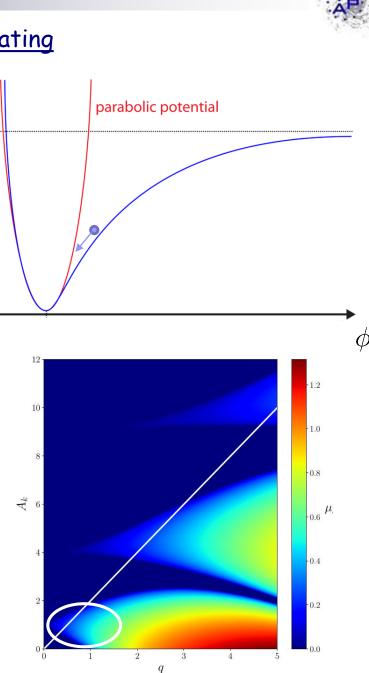
□ When the field oscillates, the universe is effectively matter-dominated

During this phase, perturbations obey a Mathieu equation. Modes in the first instability band, ie

$$0 < \frac{k}{a} < \sqrt{3Hm}$$

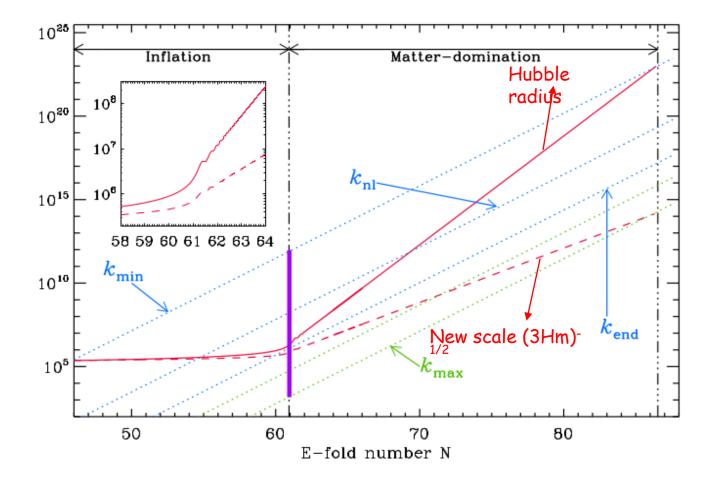
undergo parametric reasonance and the density contrast grows like the scale factor, like a matter dominated Universe

K. Jedamzik, M. Lemoine, J. Martin, arXiv:1002.3039



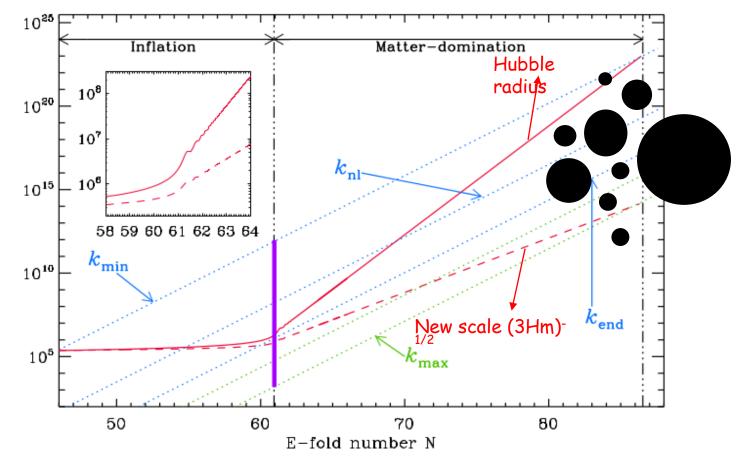


## 3- GW and ultra-lights PBHs from metric preheating





## 3- GW and ultra-lights PBHs from metric preheating



Lead to the formation of ultra-light PBHs at the end of inflation

J. Martin, T. Papanikolaou, V. Vennin, arXiv:1907.04236

Remark: no need to "distord" the potential, USR, multifield etc ...



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# Å.

### Recap

□ Single field slow-roll (with canonical kinetic terms) models are solid, robust and fasifiable scenarios of the early Universe

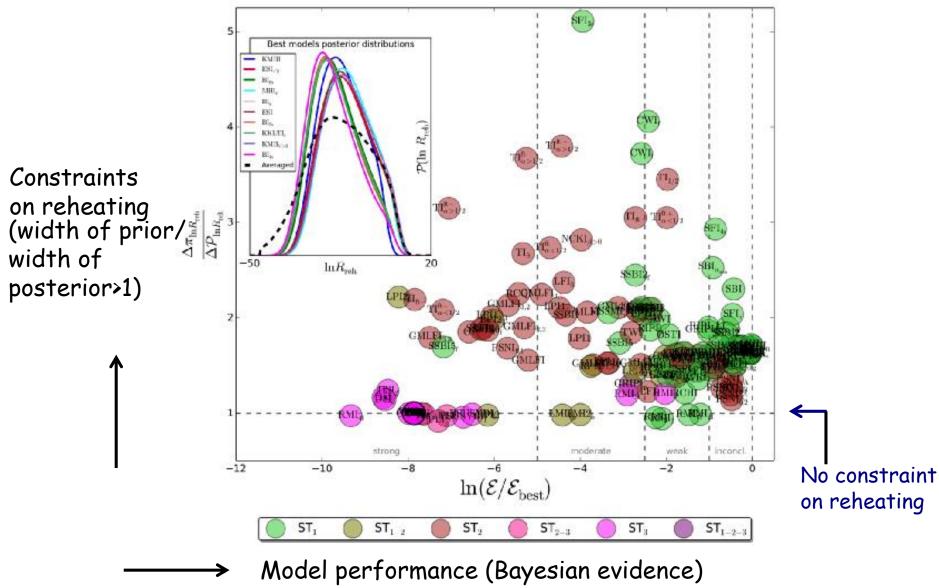
□ They have made predictions (ie not postdictions) that fit well the data

There are still open issues, eg how to understand (single field) inflation from a high energy point of view, reheating ...

□ There are still predictions that have not been verified, GW and PBHs

- Recently, many new interesting mechanisms for non-minimal inflation have been studied (multi-field inflation, inflation and gauge fields etc ...)
- To see to which realizations of inflation we deal with, GW, PBHs and NG are keys

#### Planck 2013 constraints on reheating

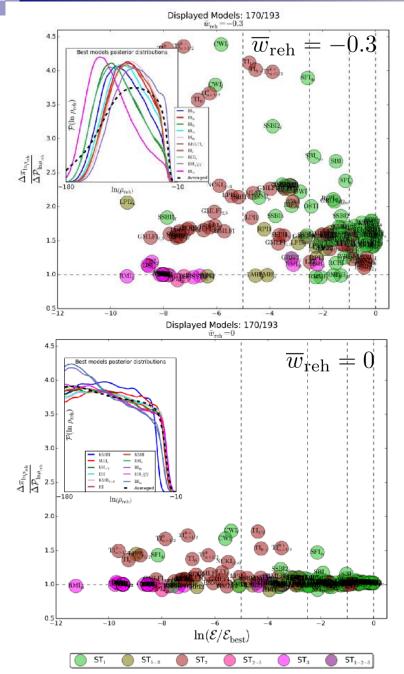


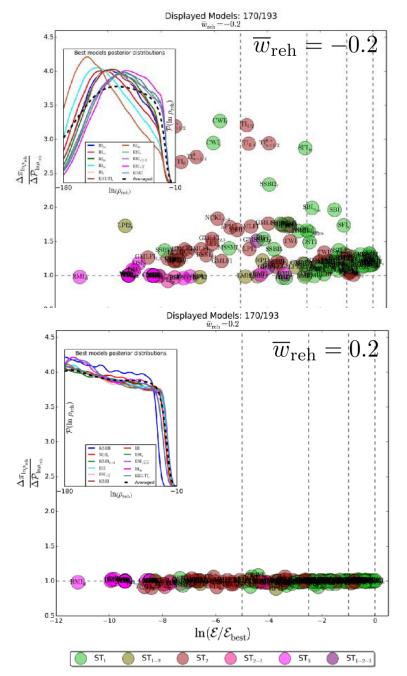
Displayed Models: 170/193

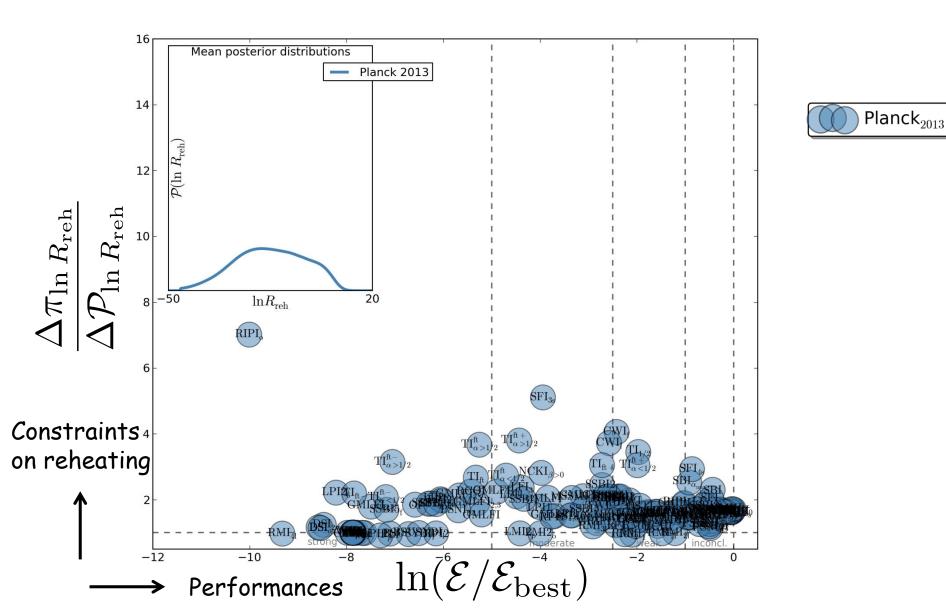
J. Martin, C. Ringeval and V. Vennin, Phys. Rev. Lett. 114 (2015) 8, 081303, arXiv:1410.7958

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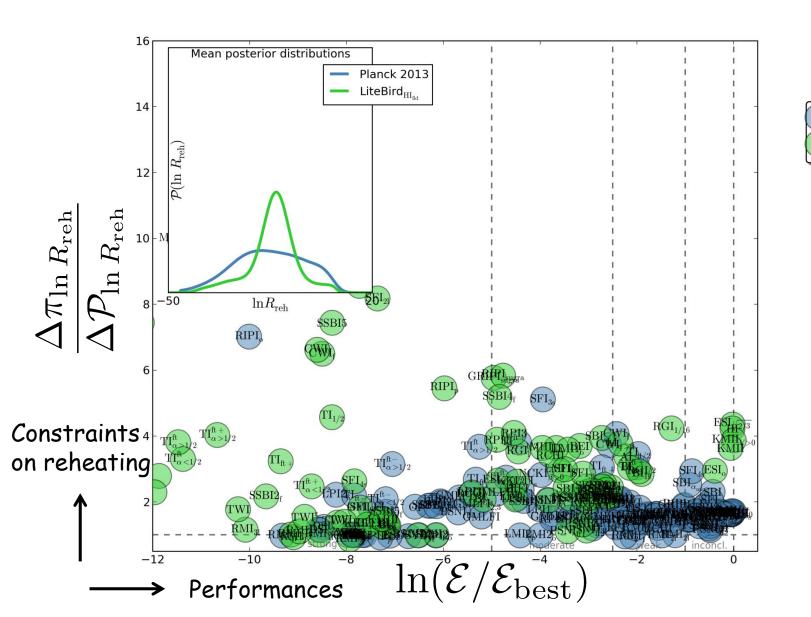


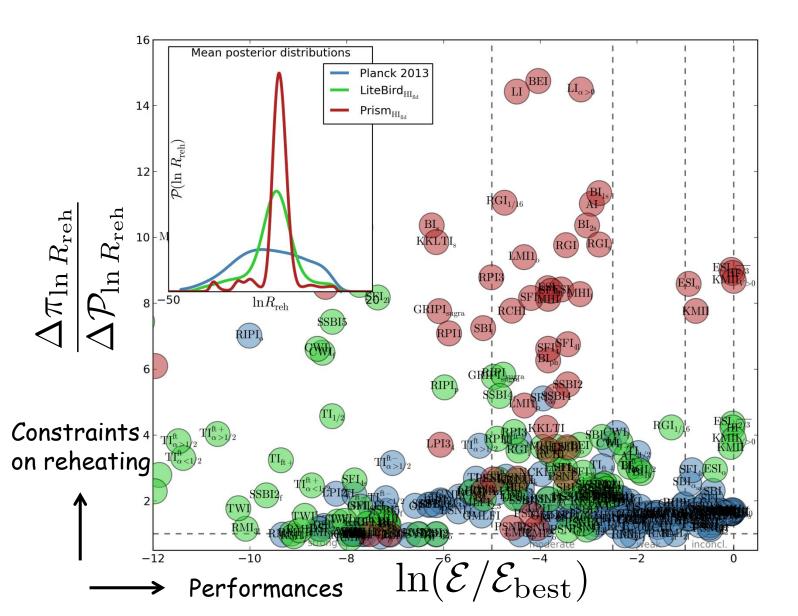


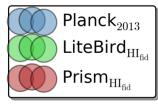


Planck<sub>2013</sub>

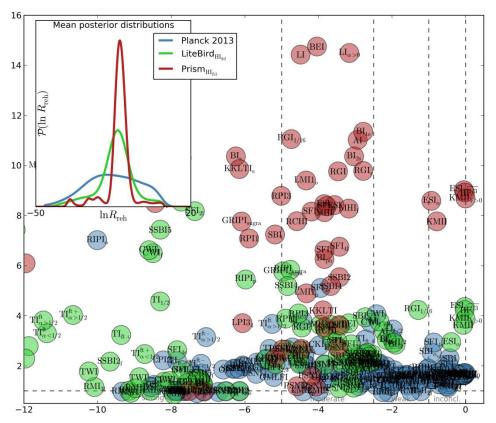
 $\mathsf{LiteBird}_{\mathrm{HI}_{\mathrm{fid}}}$ 











J. Martin, C. Ringeval and V. Vennin, Phys. Rev. Lett. 114 (2015) 8, 081303, arXiv:1410.7958

J. Martin, C. Ringeval and V. Vennin, JCAP 1410 (2014) 10, 038, arXiv:1407.4034

#### Planck 2013

$$\left\langle \frac{\Delta \pi_{\ln R_{\rm reh}}}{\Delta \mathcal{P}_{\ln R_{\rm reh}}} \right\rangle \simeq 40\%$$

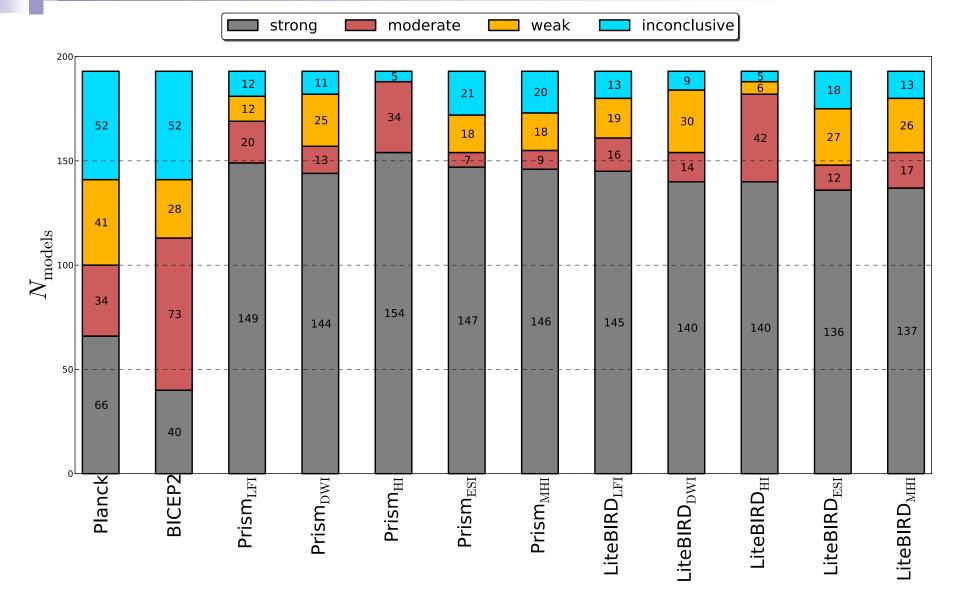
#### LiteBIRD HI

$$\left\langle \frac{\Delta \pi_{\ln R_{\rm reh}}}{\Delta \mathcal{P}_{\ln R_{\rm reh}}} \right\rangle \simeq 73\%$$

#### Prism HI

$$\left\langle \frac{\Delta \pi_{\ln R_{\rm reh}}}{\Delta \mathcal{P}_{\ln R_{\rm reh}}} \right\rangle \simeq 88\%$$

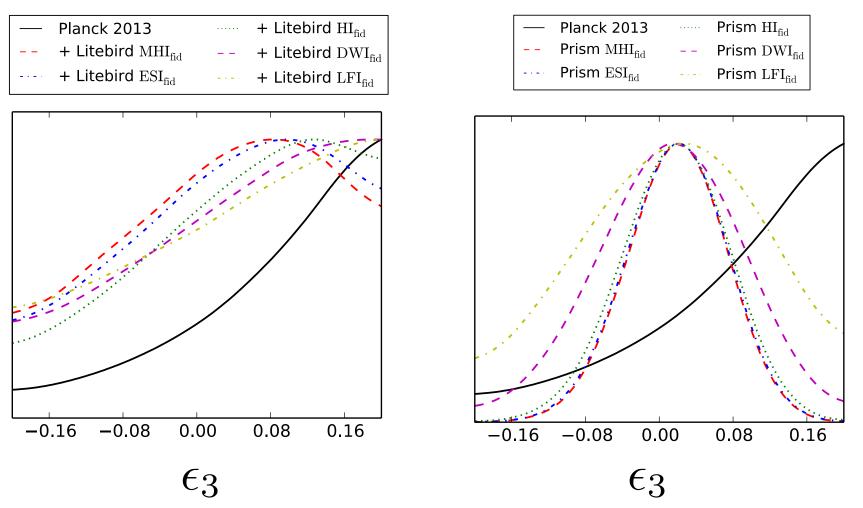
#### Consequences of a B-modes detection



Planck: 1/3 of the models excluded; PRISM & LiteBIRD > 4/5

J. Martin, C. Ringeval and V. Vennin, JCAP 1410 (2014) 10, 038, arXiv:1407.4034

#### Message 5: Prism can detect the slow-roll running ...



J. Martin, C. Ringeval and V. Vennin, JCAP 1410 (2014) 10, 038, arXiv:1407.4034



- In order to evaluate the performance of a model, one can compare its predictions to the data in the parameter space  $(n_{\rm S},r)$ 

- But how can we compare the performance of models with each other?

- The performance of a model can be described by its <u>Bayesian evidence</u> which is the integral of the likelihood over prior space

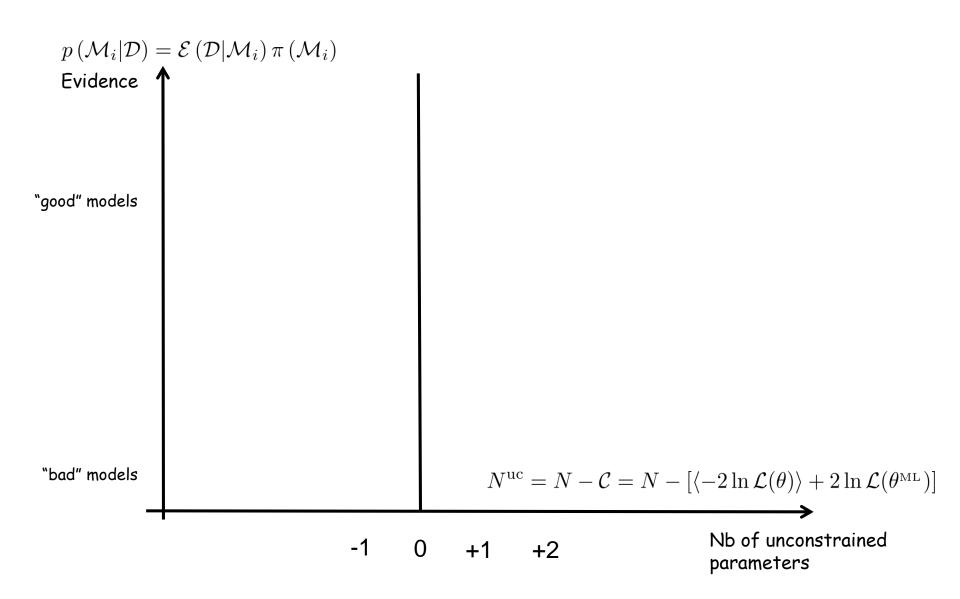
$$p\left(\mathcal{M}_{i}|\mathcal{D}\right) = \mathcal{E}\left(\mathcal{D}|\mathcal{M}_{i}\right)\pi\left(\mathcal{M}_{i}\right)$$

- Another number is needed in order to describe the performance of a model: the <u>effective number of unconstrained parameters</u> or <u>Bayesian complexity</u>

$$N^{\rm uc} = N - \mathcal{C} = N - \left[ \langle -2\ln \mathcal{L}(\theta) \rangle + 2\ln \mathcal{L}(\theta^{\rm ML}) \right]$$

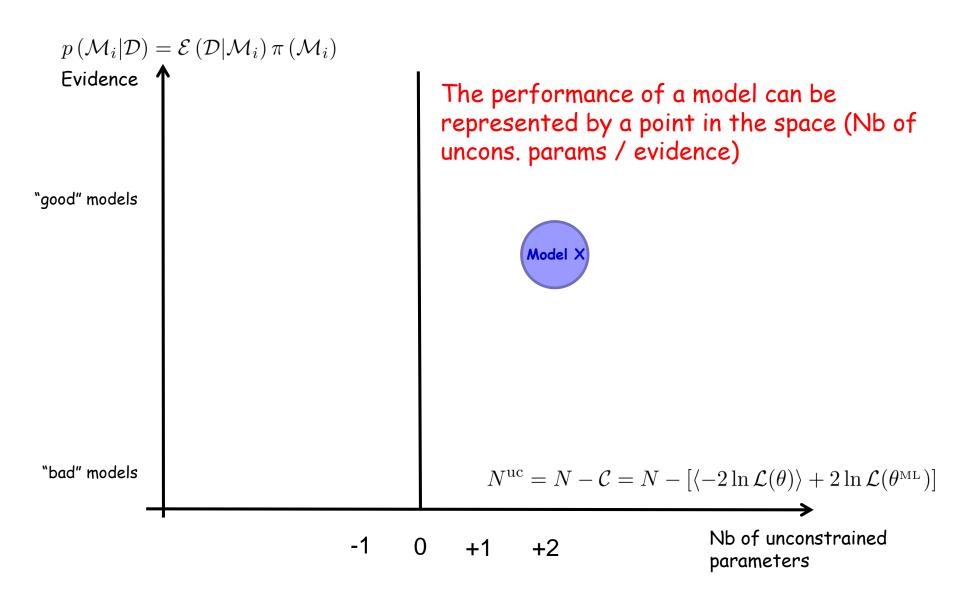


The performance of a model can be represented in the space (complexity, evidence)



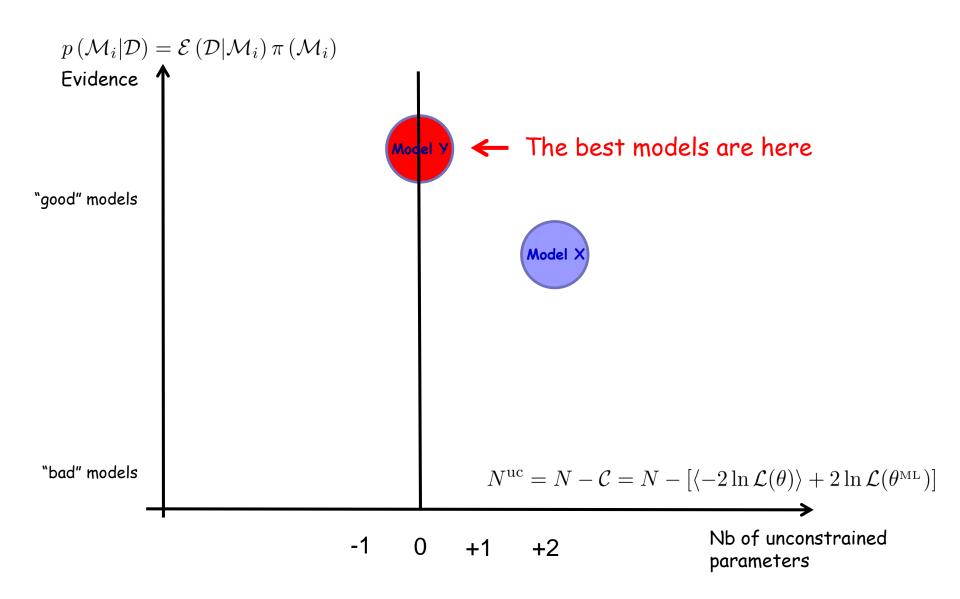


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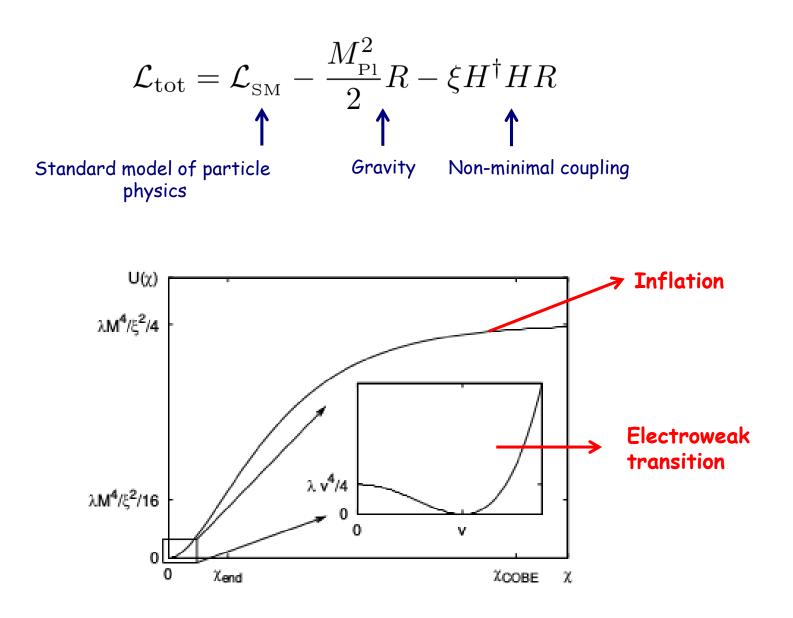


The performance of a model can be represented in the space (complexity, evidence)





Can the inflaton field be the Higgs boson?





Small field inflation

$$\rho_{\rm reh}^{1/4} > 400 \,\mathrm{TeV} \qquad \overline{w}_{\rm reh} = -0.3$$

$$\rho_{\rm reh}^{1/4} > 90 \,\mathrm{TeV} \qquad \overline{w}_{\rm reh} = -0.2$$

Loop inflation

$$\begin{split} \rho_{\rm reh}^{1/4} &< 1.8 \times 10^7 \, {\rm TeV} & \overline{w}_{\rm reh} = -0.3 \\ \rho_{\rm reh}^{1/4} &< 6.5 \times 10^7 \, {\rm TeV} & \overline{w}_{\rm reh} = -0.2 \\ \rho_{\rm reh}^{1/4} &< 4 \times 10^{10} \, {\rm TeV} & \overline{w}_{\rm reh} = 0 \end{split}$$