

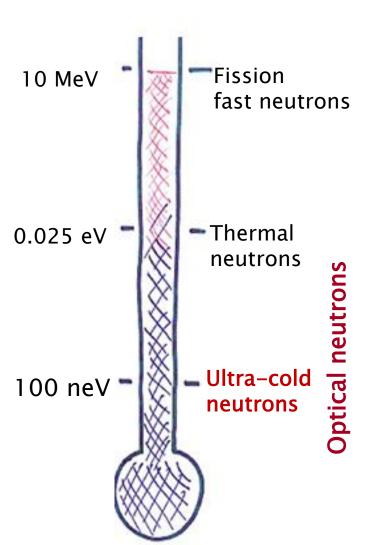


# Spatial resolution of CCD based UCN detectors



Benoit Clément - VCI - 02/2022

# **Ultra-cold neutrons (UCN)**



When the neutron wavelength is greater than the interatomic distance

**⇒** Coherent scattering on surfaces

**Characterised by an effective Fermi potential:** 

$$V = \frac{2\pi\hbar^2}{m_n} < b^{coh} >$$

For neutron energy < Fermi potential: neutrons are reflected on a surface at every incidence angle.

These are Ultra-cold neutrons:

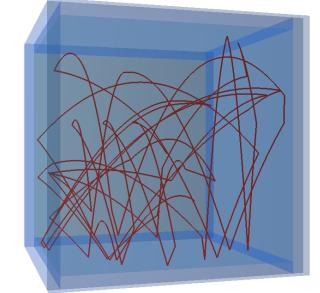
$$E < 250 \text{ neV } v < 7\text{m/s } \lambda > 0.2\text{nm}$$

4m/s neutron in a 50cm cubic box

UCNs can be stored in boxes, guided by material tube,...

**UCNS** are sensitive to gravity:

$$m_n g x (1m) = 100 neV$$



... but UCNs bounce of (most) detectors

### **UCNs** detection

Neutron detection required a conversion material to produce charges particles (<sup>3</sup>He, <sup>10</sup>B, <sup>6</sup>Li, <sup>235</sup>U, <sup>238</sup>Pu) or gamma rays (Gd,...)

For UCN, the Fermi potential must be taken into account

#### **Gaseous detector (3He), limiting factors:**

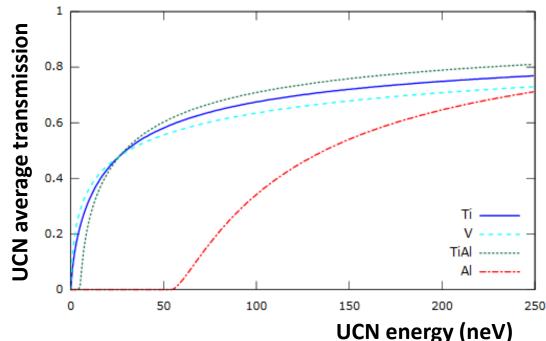
- Fermi potential of entrance window:
- Absobtion/scattering losses within the window: thin foils

commonly used: Aluminum, titanium

#### **Solid detector**

- Fermi potential of the conversion layer
- Converted particle must escape the layer (Boron on silicon) or the conversion layer must be an active medium (LiF scintillators)

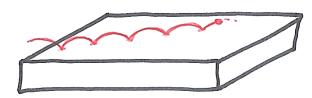
 $^{10}B: V_f = -3 \text{ neV (20\% of nat. boron)}$ 



n + 
$$^{10}\text{B} \rightarrow ^{7}\text{Li} + \alpha$$
 , BR = 6.3%  $\text{E}_{\text{Li}}\text{= }1014 \text{ keV}$   $\text{E}_{\alpha}$  = 1775 keV

n + 
$$^{10}$$
B  $\rightarrow$   $^{7}$ Li\* +  $\alpha$   $\rightarrow$   $^{7}$ Li +  $\alpha$  +  $\gamma$  , BR = 93.7%   
E<sub>Li</sub>= 841 keV E $_{\alpha}$  = 1471 keV

### **Quantum freefall and position sensitive detectors**



We study the quantum free fall of a neutron : Bound state in a vertical potential mgz

Schrödinger equation : 
$$\left(\frac{\hbar^2}{2m_i}\frac{\partial^2}{\partial z^2} + m_g g z\right) \psi = E \psi \ \Rightarrow \left(\frac{\partial^2}{\partial Z^2} + Z\right) \psi = 0$$

With : 
$$z_0=\left(\frac{\hbar^2}{2m_im_ag}\right)^{\frac{1}{3}}\approx 5.87\mu m$$
  $E_0=m_gg~z_0\approx 0.602~{
m peV}$   $\epsilon=\frac{E}{E_0}$   $Z=\frac{z}{z_0}-\epsilon$ 

$$E_0=m_g g \ z_0pprox 0.602 \ \mathrm{peV}$$

Airy equation

$$\epsilon = \frac{E}{E_0}$$

$$Z = \frac{z}{z_0} - \epsilon$$

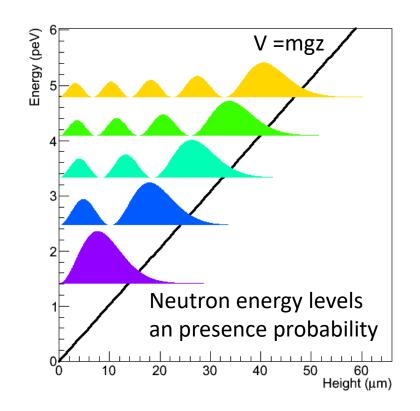
Boundary conditions  $\psi(0) = \psi(+\infty) = 0$  quantify the energies  $E_k$ :

### Two kind of observables

Quantified energies: transition between levels

Wavefunction: shape, nodes positions.

**Need position sensitive detector with** micrometric resolution



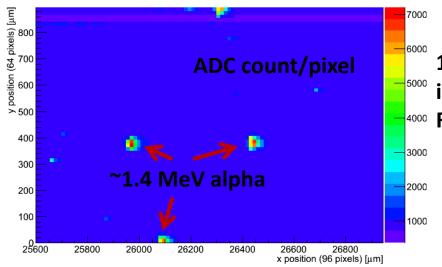
### The UCNBox detector

To get a spatial information we use CCD sensors as charge particle detectors

- window less back-thinned CCD (Hamamatsu)
- 2048x64 pixels, 14x14 μm

#### **UCN Boron piXels detector:**

- line of 8 sensors (30 cm) to adapt to the GRANIT installation at ILL to observe neutron wavefunctions
  - mechanical support for alignment
  - dedicated electronics O. Bourrion et al. NIM. A880 (2018)



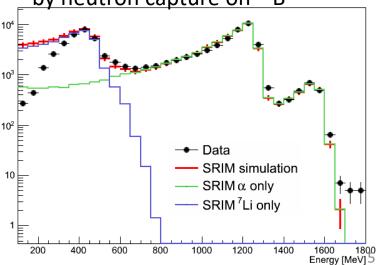
1 charge particle (1.4 MeV  $\alpha$ ) deposits energy in several pixels

#### 4000 Reconstruction strategy:

- build 11x11 pixels clusters around «seed» pixels
- the cluster barycenter gives the particle position
- the cluster mean ADC count gives the energy

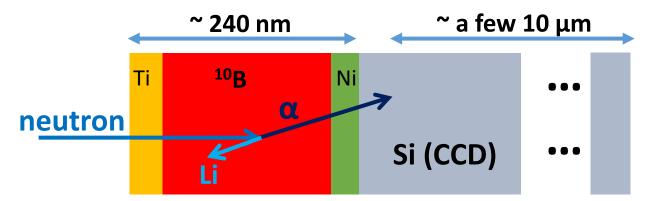


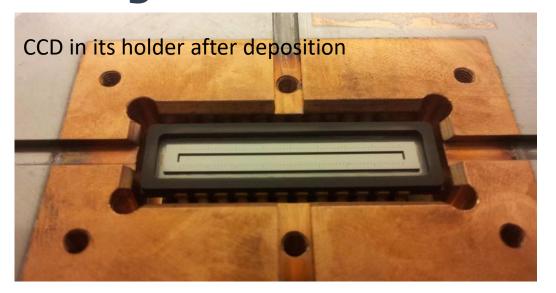
Measured energy of  $\alpha$ /Li produced by neutron capture on  $^{10}$ B



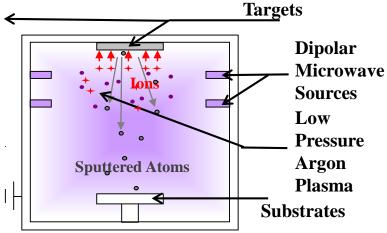
# Realisation of the multilayer

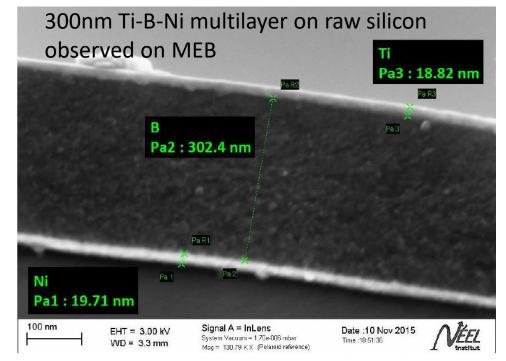
Multilayer Ti-B-Ni deposed using microwave plasma-assisted cosputtering method



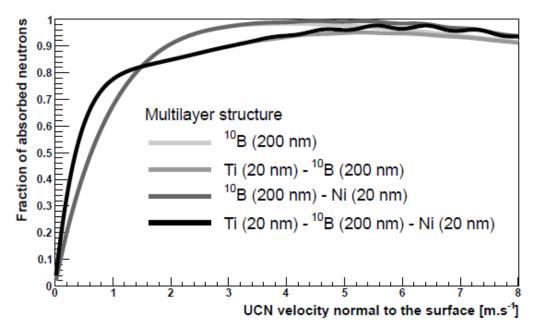


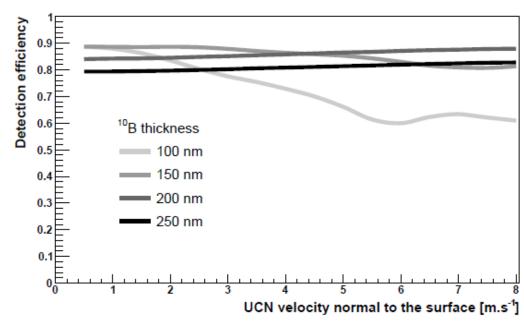






### Ti-B-Ni conversion multilayer and efficiency





B.C. et al, *JINST* 14 (2019) 09, P09003

#### The conversion layer devised is:

20 nm entrance Ti layer

- → reduce reflection at low velocities
- → protects boron from oxidation

200 nm conversion <sup>10</sup>B layer

→ optimum to convert UCNs and let charged particles out of the layer

20 nm back Ni layer

- → reflect hich velocities UCN
- → improve boron adhesion on silicon substrate

# Measuring the spacial resolution

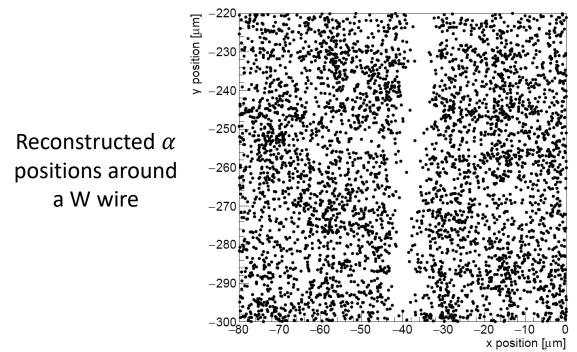
Place a piece of wire on the sufrace of the CCD  $\,$  « Light » with neutrons or  $\alpha$  particles and detect the wire shadow

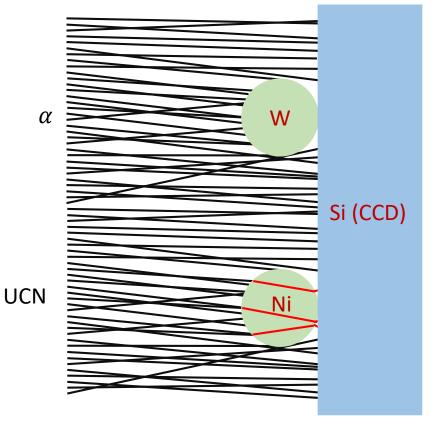
#### Two datasets:

- 1.5MeV  $\alpha$  particles over 6 $\mu$ m tungsten wire (no boron layer)
- UCN over 10μm Ni wires -> at ILL PF2 Test line

Approximtively 2.4mm of usable Ni wire and 3.5mm of W.

A few days of data, 0.5 to 0.7 particle per μm<sup>2</sup>





#### **Problems**

- Reconstruction the position and direction of the wire
- Projection along this direction and combining several wires
- Exctraction of the resolution : need MC simulation
- Quantifying the systematics

### **Estimation of density of points**

#### **Delaunay triangulation:**

Given a set of points in 2D, find the triangle meshing that maximise the smallest angle of each triangle (tringle less flats as possible)

Standard tool for meshing in finite elements calculation: many optimized algorithms and

freely vailable codes!

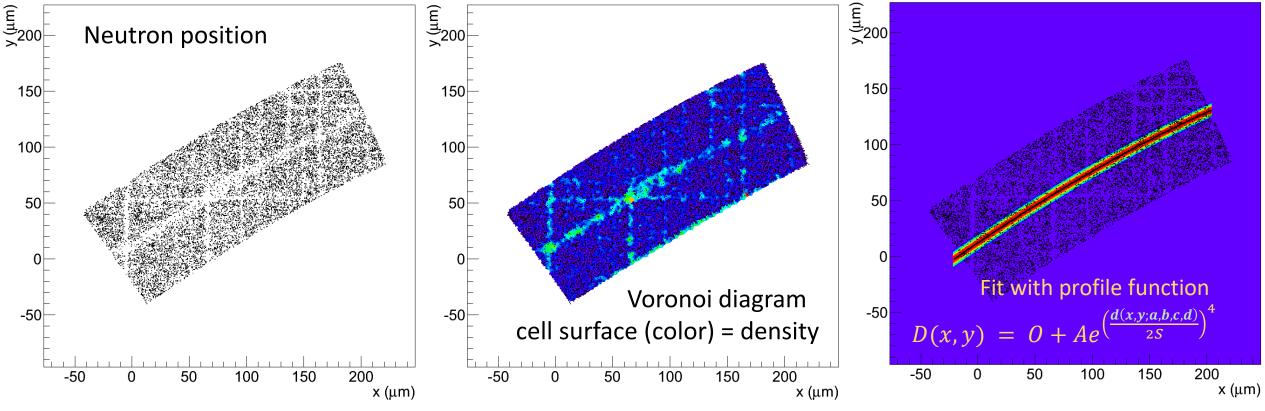
### Voronoi diagram:

compute the median of each segment in the Delaunay trangulation and compute their first instersections.

each original point is contained within one Voronoi cell, the surface can be used as a (inverse) density estimator: used in cellular biology, geography,...

### Wire direction

Cut each wire in ~200µm segments, and manually select a band around the wire

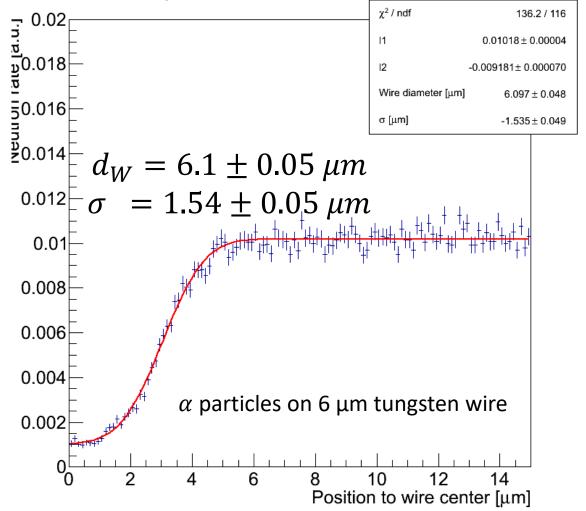


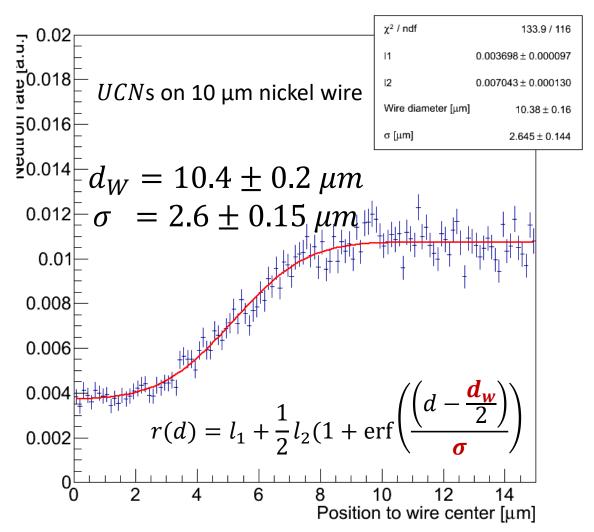
d(x, y; a, b, c, d): distance of point (x, y) to the polynomial curve  $y = ax^3 + b^{x^2} + cx + d$ 

Histograms of d gives the projected wire profile

### **Projected wires**

### The wire profile is fitted with an Erf function





The fitted  $\sigma$  is not the resolution yet has it includes shadowing effects due to the wire shape and diffusion of fastets UCNS through the wire.

### **Conversion to resolution**

Use Monte-Carlo simulation to compute the conversion curve from fitted resolution to actual resolution.

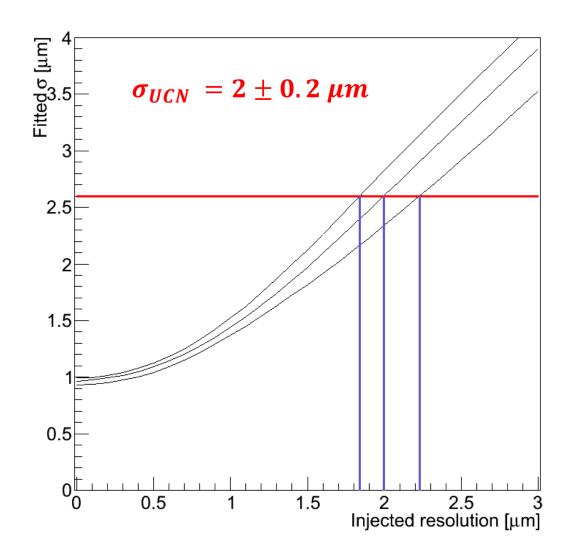
Some uncontrolled parameters

- diffusivity of the wire surface
- velocity spectrum of UCNs

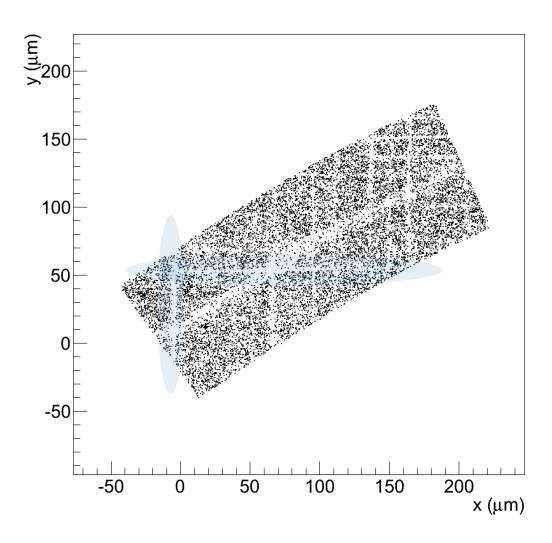
Similar exercise for alpha particles leads to

$$\sigma_{\alpha} = 1 \pm 0.1 \,\mu m$$

The systematic effect due to the reconstruction and fitting procedure is estimated by building fake wire images using the fitted parameters :  $\sim 0.3 \mu m$ 



# Sensor aging and hot pixels



Looking at the data, one noticies depleted vertical and horizontal lines

- → 11x11 squares around « hot » pixels (pixels with sometimes large noise that tend to displace the reconstructed hit toward them)
- → didn't appear in older measurement : degradation of sensor over time (multilayers are >3 years old, sensors kept in contact with air in a clean room)
- → Impossible to correct, but can be simulated : this effect alone leads to :

$$\sigma_{fit} = 1.9 \pm 0.2 \ \mu m \Rightarrow \sigma_{UCN} = 1.5 \ \mu m$$

Quadratically substracting to the previous result, one could extrapolated an expected resolution for a freshly coated sensor to

$$\sigma_{UCN} = \sim 1.3 \, \mu m$$

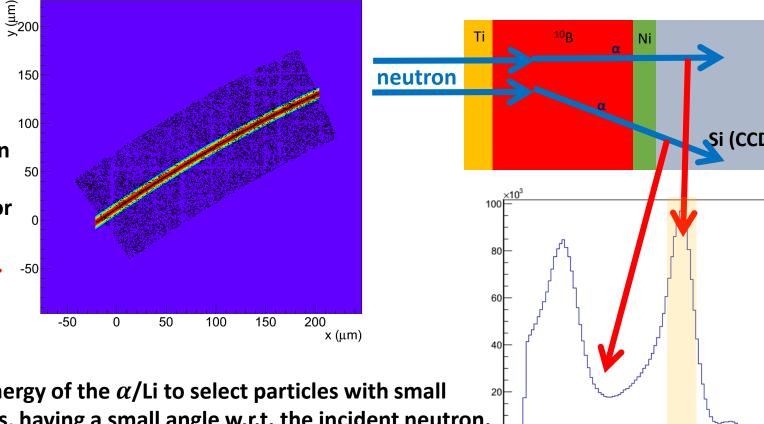
# **Summary and conclusion**



A complete procedure to determine experimentaly the spatial resolution has been developed using the shadow of a thin wire in contact with the conversion layer The measured resolution on an «old» sensor is:

 $\sigma_{UCN}=2\pm0,2(stat)\pm0,3~(syst)$  µm. -50 and the resolution of a fresh sensor could lower to 1,3µm.

UCNBox is a UCN detectors using 8 CCD sensors coated with a Ti-<sup>10</sup>B-Ni multilayer with an efficiency of 84% designed to measure wavefunctions of free falling neutrons bouncing on a mirror.



One possible improvement: Use the energy of the  $\alpha$ /Li to select particles with small energy loss, having a small angle w.r.t. the incident neutron. This should improve the resolution at the cost of statistics.

ADC count ∝ energv