

# The Hubble Tension and the Magnetic Universe

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with Karsten Jedamzik (LUPM, Montpellier),  
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arXiv:2004.09487, Phys. Rev. Lett.  
arXiv:2109.03816, submitted to PRD

# This Talk

The Hubble Tension: how fast is the universe expanding today?

Primordial Magnetic Fields

How magnetic fields help to relieve the  $H_0$  tension

Further tests



# The Hubble Constant

## *A RELATION BETWEEN DISTANCE AND RADIAL VELOCITY AMONG EXTRA-GALACTIC NEBULAE*

BY EDWIN HUBBLE

MOUNT WILSON OBSERVATORY, CARNEGIE INSTITUTION OF WASHINGTON

Communicated January 17, 1929

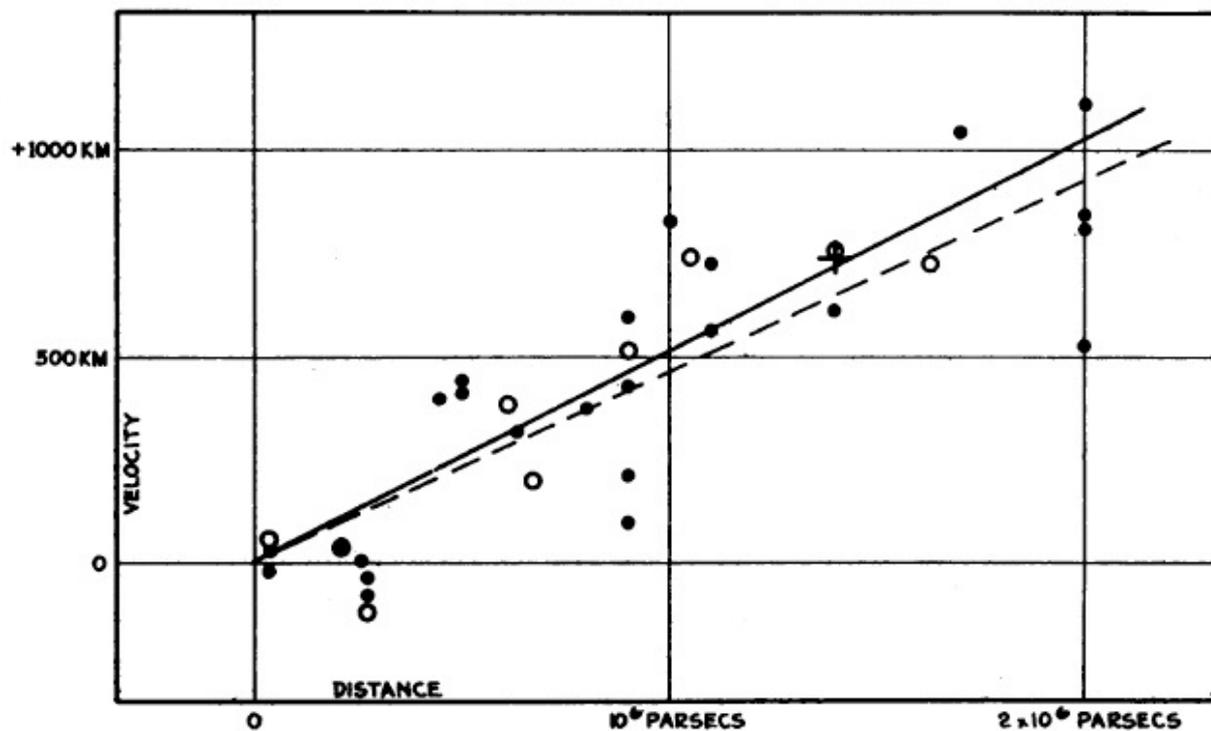
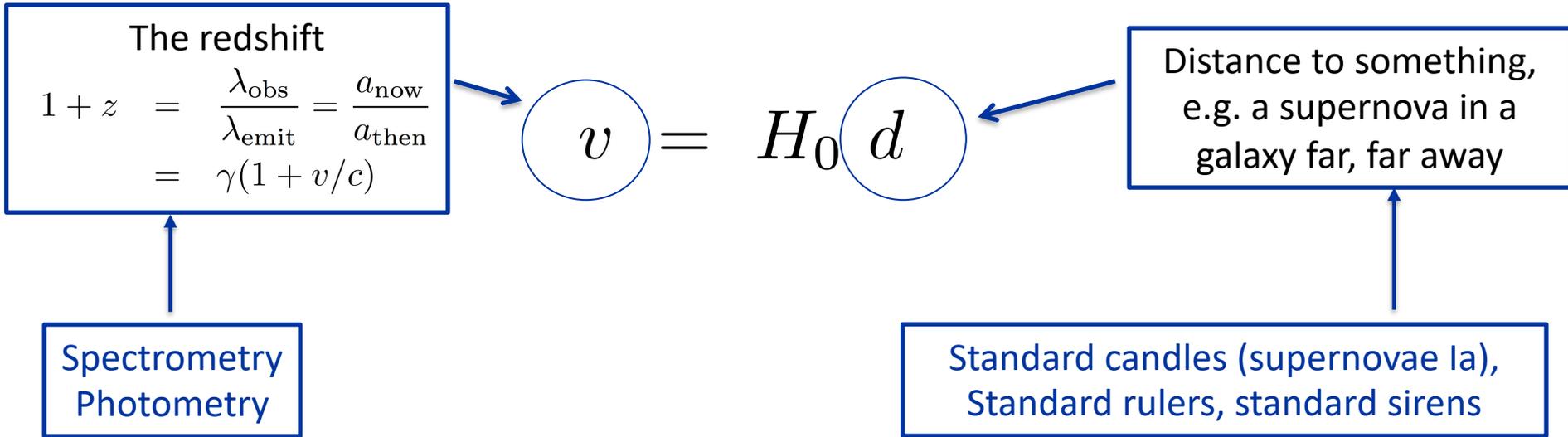


FIGURE 1

Velocity-Distance Relation among Extra-Galactic Nebulae.

# Measuring the expansion rate *a la* Hubble



$$\text{Observed Luminosity} = \frac{\text{Intrinsic Luminosity}}{4\pi \text{ Distance}^2} = \frac{L}{4\pi d_L^2}$$

$$d_A = \frac{\text{Known Physical Length}}{\text{Observed Angular Size}} = \frac{S}{\theta}$$

$$[H_0] = [\text{velocity/distance}] = [1/\text{time}] = \text{km/s/Mpc}$$

## It's a bit more complicated...

H is not a constant

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2(z) = H_0^2 [\Omega_r(1+z)^4 + \Omega_M(1+z)^3 + \Omega_k(1+z)^2 + \Omega_{DE}]$$

$$\Omega_r + \Omega_M + \Omega_k + \Omega_{DE} = 1$$

Distances depend on H(z)

$$d_L(z) = (1+z) \int_0^z \frac{c dz'}{H(z')}$$

$$d_A(z) = \frac{1}{(1+z)} \int_0^z \frac{c dz'}{H(z')}$$

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Distances depend on H(z)

$$d_L(z) = (1+z) \int_0^z \frac{c dz'}{H(z')} \xrightarrow{z \ll 1} cH_0^{-1}z \approx H_0^{-1}v$$

$$d_A(z) = \frac{1}{(1+z)} \int_0^z \frac{c dz'}{H(z')} \xrightarrow{z \ll 1} cH_0^{-1}z \approx H_0^{-1}v$$



## OPEN ACCESS

# Large Magellanic Cloud Cepheid Standards Provide a 1% Foundation for the Determination of the Hubble Constant and Stronger Evidence for Physics beyond $\Lambda$ CDM

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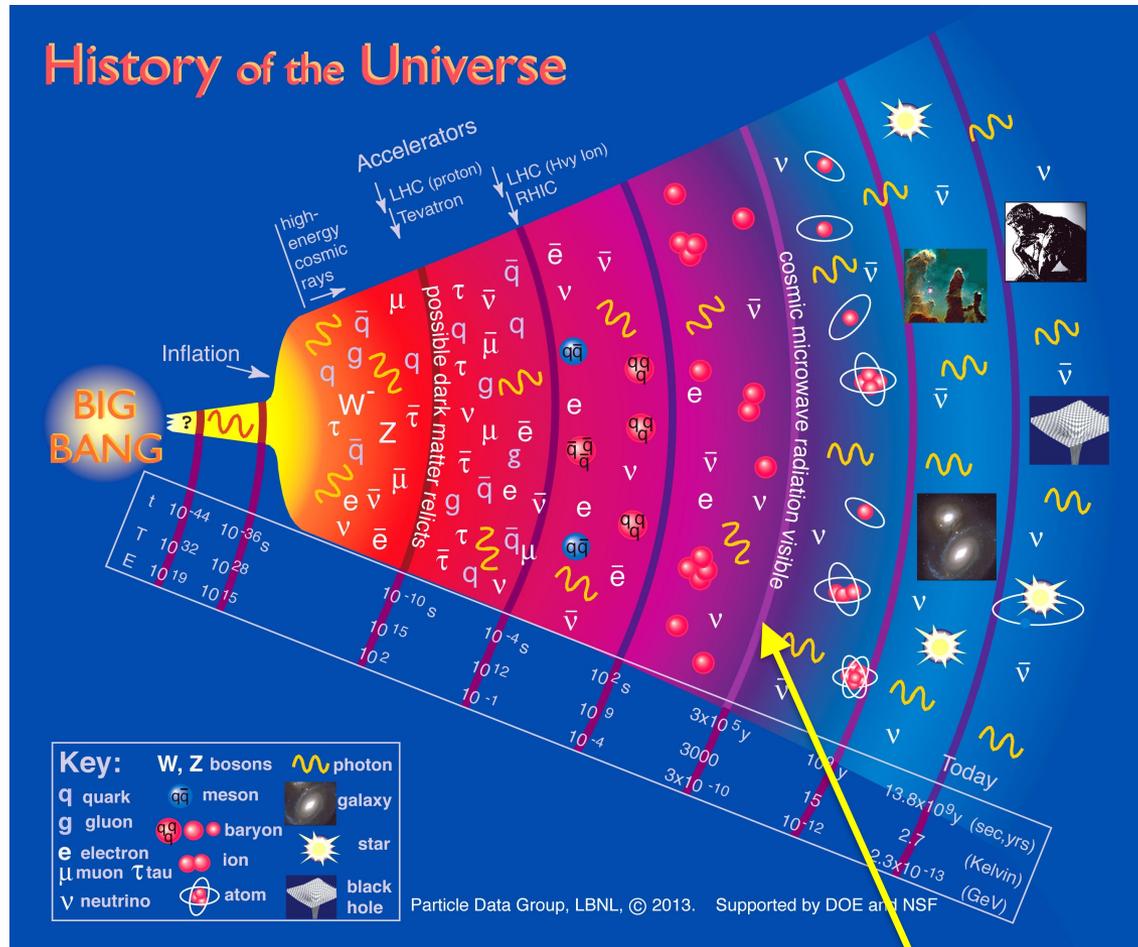
Received 2019 March 18; revised 2019 March 25; accepted 2019 March 26; published 2019 May 7

## Abstract

We present an improved determination of the Hubble constant from *Hubble Space Telescope* (*HST*) observations of 70 long-period Cepheids in the Large Magellanic Cloud (LMC). These were obtained with the same WFC3 photometric system used to measure extragalactic Cepheids in the hosts of SNe Ia. Gyroscopic control of *HST* was employed to reduce overheads while collecting a large sample of widely separated Cepheids. The Cepheid period–luminosity relation provides a zero-point-independent link with 0.4% precision between the new 1.2% geometric distance to the LMC from detached eclipsing binaries (DEBs) measured by Pietrzyński et al. and the luminosity of SNe Ia. Measurements and analysis of the LMC Cepheids were completed prior to knowledge of the new DEB LMC distance. Combined with a refined calibration of the count-rate linearity of WFC3-IR with 0.1% precision, these three improved elements together reduce the overall uncertainty in the geometric calibration of the Cepheid distance ladder based on the LMC from 2.5% to 1.3%. Using only the LMC DEBs to calibrate the ladder, we find  $H_0 = 74.22 \pm 1.82 \text{ km s}^{-1} \text{ Mpc}^{-1}$  including systematic uncertainties, 3% higher than before for this particular anchor. Combining the LMC DEBs, masers in NGC 4258, and Milky Way parallaxes yields our best estimate:  $H_0 = 74.03 \pm 1.42 \text{ km s}^{-1} \text{ Mpc}^{-1}$ , including systematics, an uncertainty of 1.91%–15% lower than our best previous result. Removing any one of these anchors changes  $H_0$  by less than 0.7%. The difference between  $H_0$  measured locally and the value inferred from Planck CMB and  $\Lambda$ CDM is  $6.6 \pm 1.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$  or  $4.4\sigma$  ( $P = 99.999\%$  for Gaussian errors) in significance, raising the discrepancy beyond a plausible level of chance. We summarize independent tests showing that this discrepancy is not attributable to an error in any one source or measurement, increasing the odds that it results from a cosmological feature beyond  $\Lambda$ CDM.

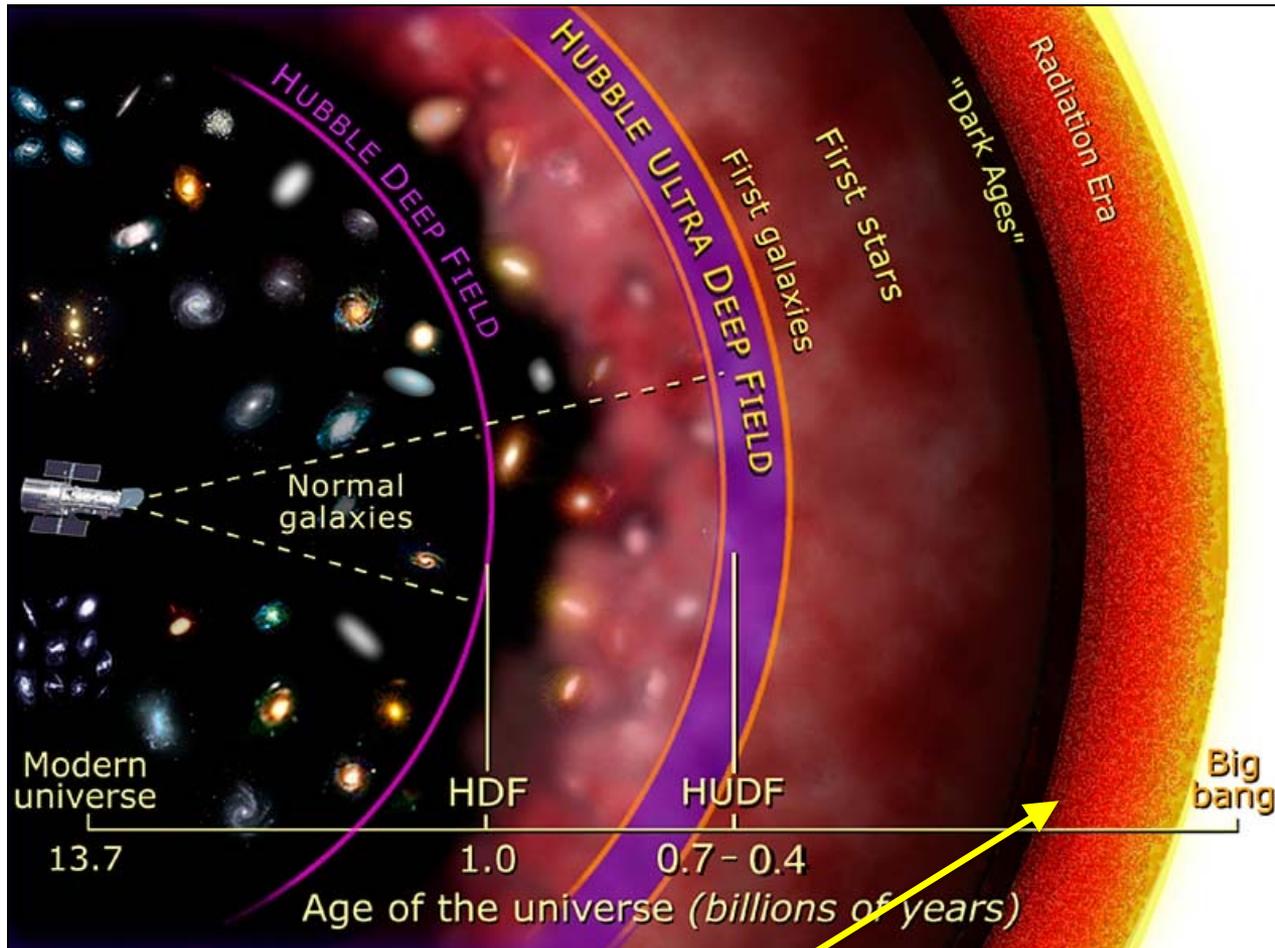
How does Cosmic Microwave Background (CMB)  
determine  $H_0$ ?

# A poster in a typical physics department



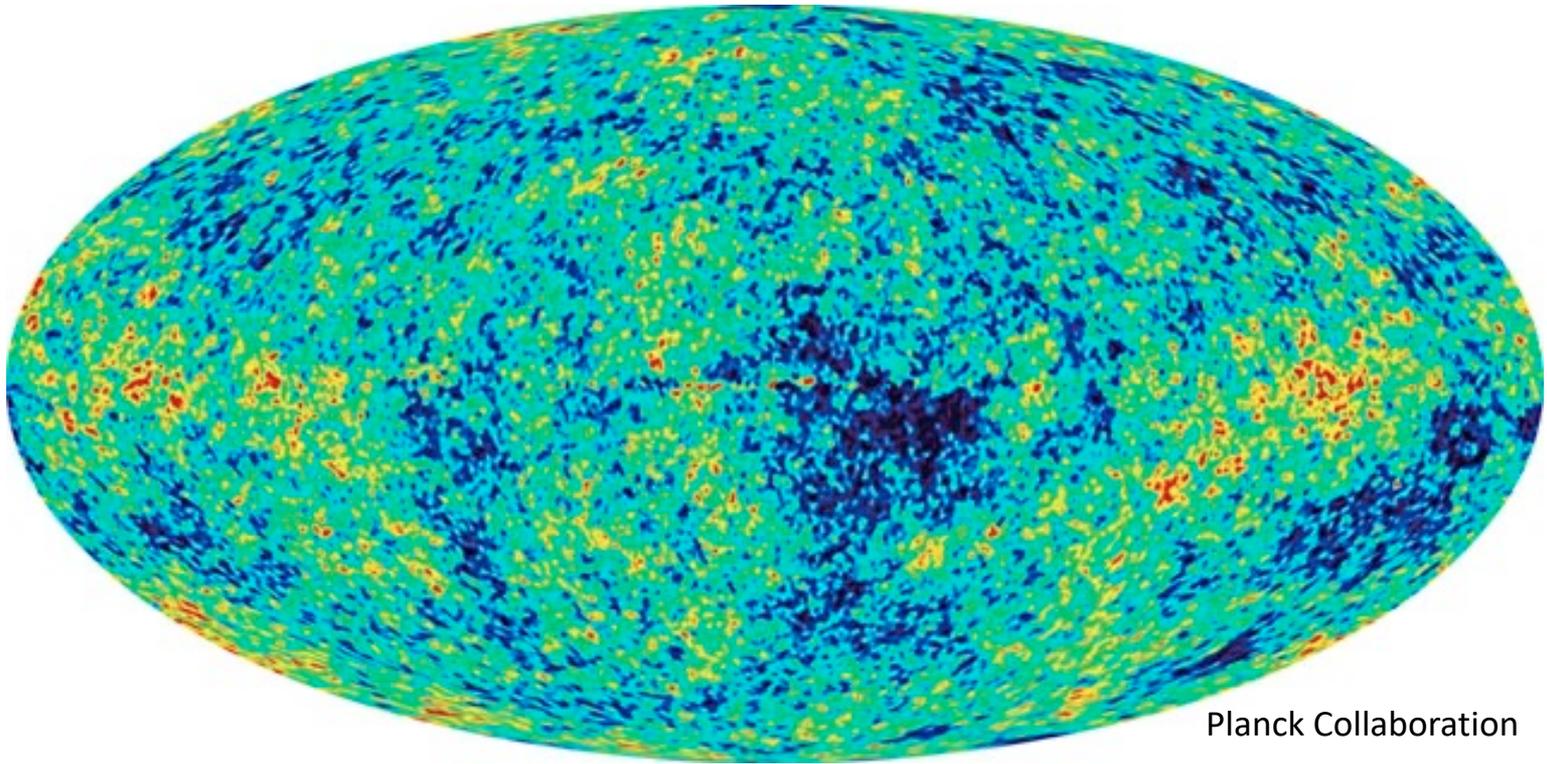
The epoch of “recombination”

# A poster in a typical astronomy department



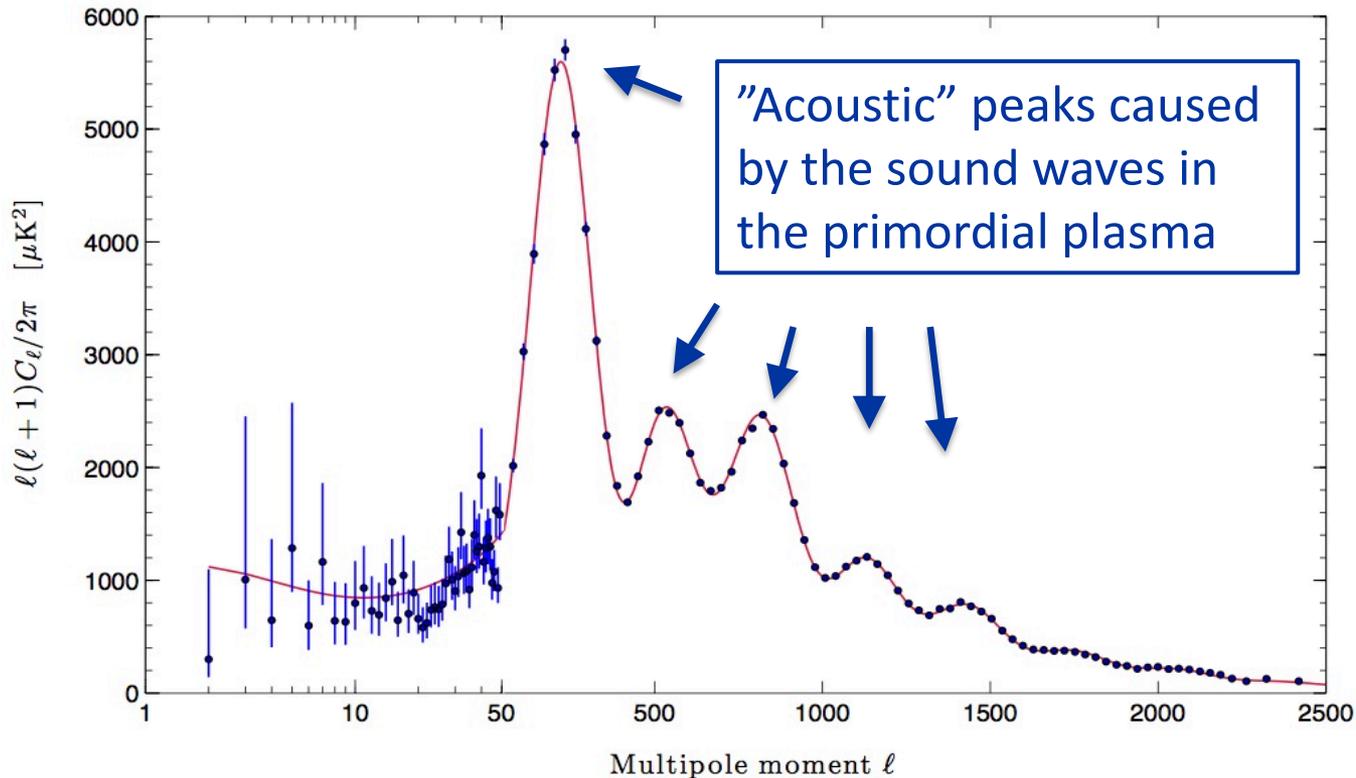
The transition from transparent to opaque  
We observe it as Cosmic Microwave Background (CMB)

# CMB temperature fluctuations



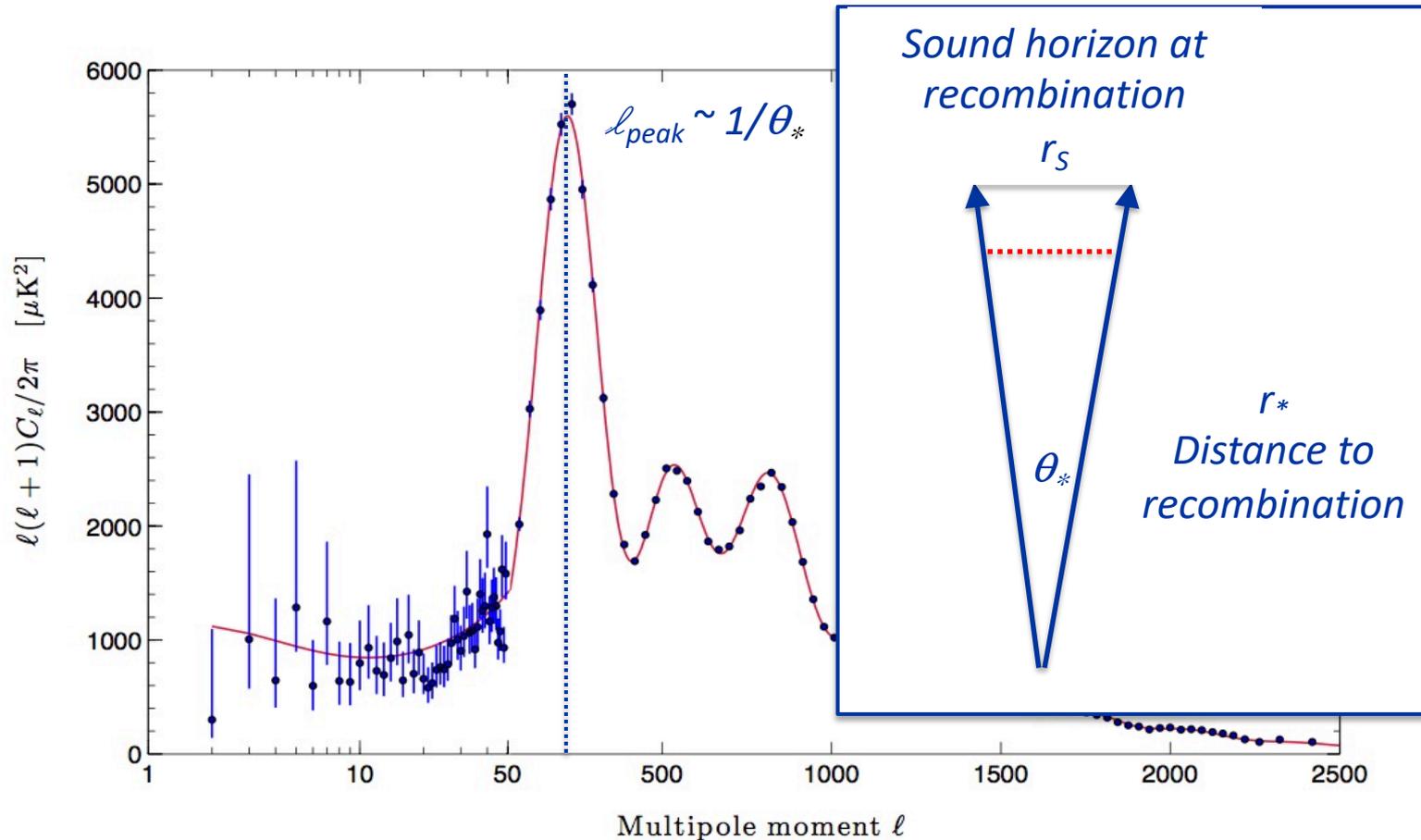
- Nearly perfect black body radiation with average  $T=2.7\text{K}$
- Temperature fluctuations of order  $1/100,000$

# The power spectrum of CMB temperature fluctuations



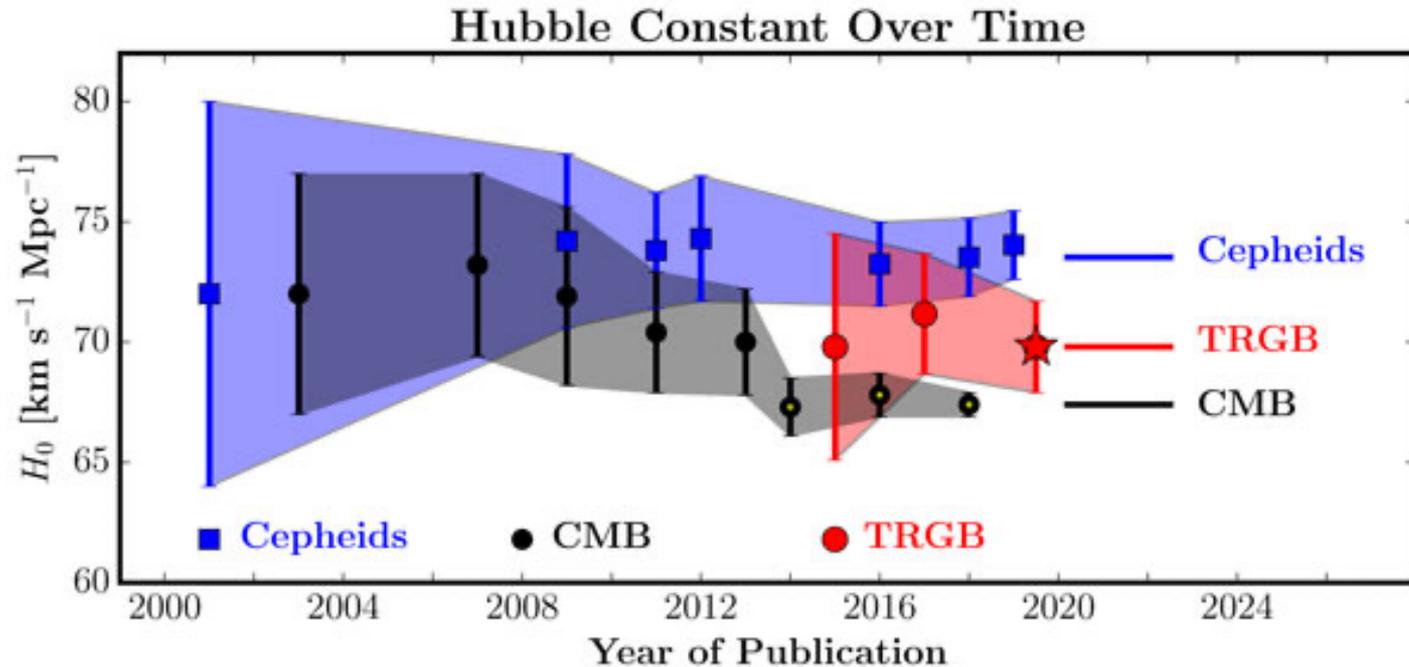
By fitting to the CMB spectrum, one constrains all parameters of the LCDM model, including  $H_0$ .

# How does CMB constrain $H_0$ ?



- Positions of the acoustic peaks tell us the angular size of the sound horizon at recombination,  $\theta_* = r_s/r_*$
- A smaller sound horizon  $r_s$  would imply a shorter distance to the redshift of recombination  $r_*$ , implying a larger  $H_0$

# The Road to the Hubble Tension



*Freedman et al. / Astrophysical Journal*

CMB (Planck):  $H_0 = 67.36 \pm 0.54$  km/s/Mpc  
Cepheid calibrated SNIa (SHOES):  $H_0 = 73.2 \pm 1.3$  km/s/Mpc

## 'Serious gap' in cosmic expansion rate hints at new physics

By Paul Rincon  
Science editor, BBC News website, Washington DC

🕒 11 January 2018



The New York Times



# *Have Dark Forces Been Messing With the Cosmos?*

Axions? Phantom energy? Astrophysicists scramble to patch a hole in the universe, rewriting cosmic history in the process.

OUT THERE

## *Cosmos Controversy: The Universe Is Expanding, but How Fast?*

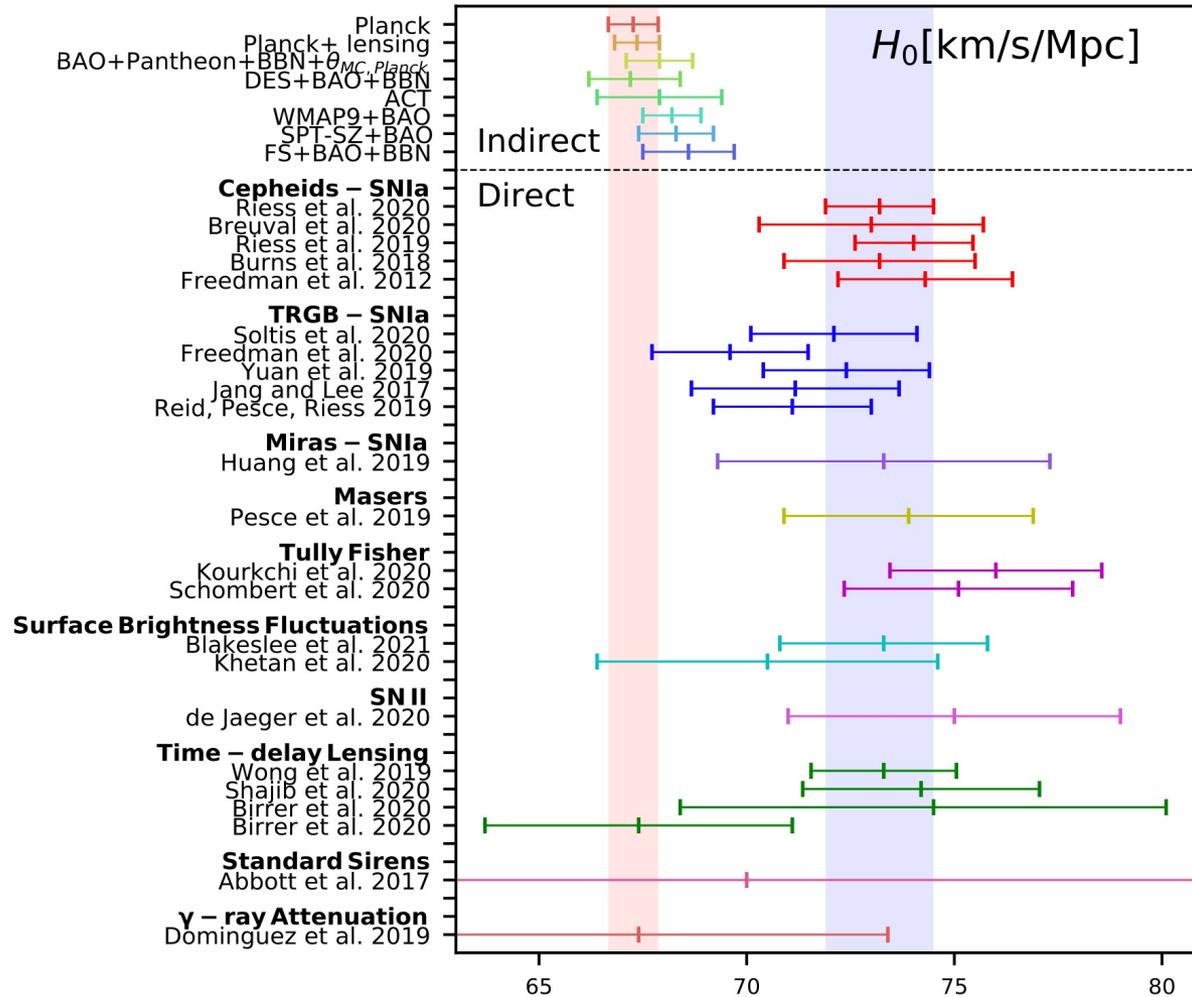
A small discrepancy in the value of a long-sought number has fostered a debate about just how well we know the cosmos.

Artwork: The expansion of the Universe

A mathematical discrepancy is "pretty serious", and could be said by a Nobel laureate.

# The Hubble tension

from E. Di Valentino,  
arXiv:2011.00246



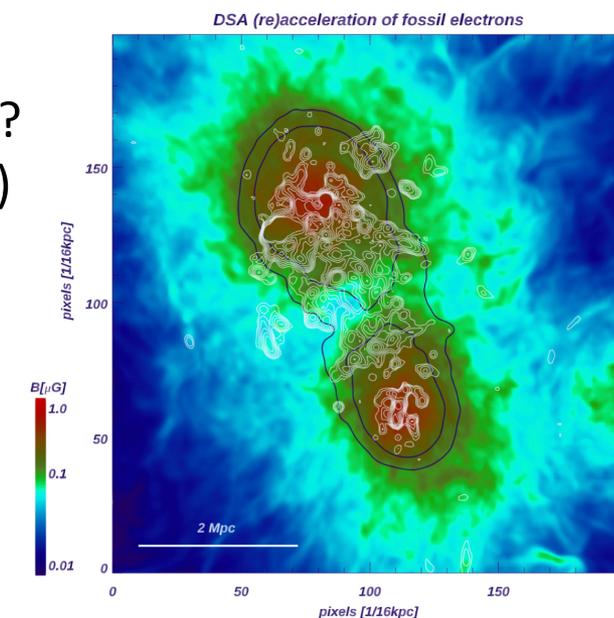
The tension is between measurements that rely on a model to determine *the sound horizon at recombination* and those that do not

# What kind of new physics can help reduce the sound horizon at recombination?

- Many models proposed with the aim of solving the Hubble tension
- Primordial Magnetic Fields

# Cosmic Magnetic Fields

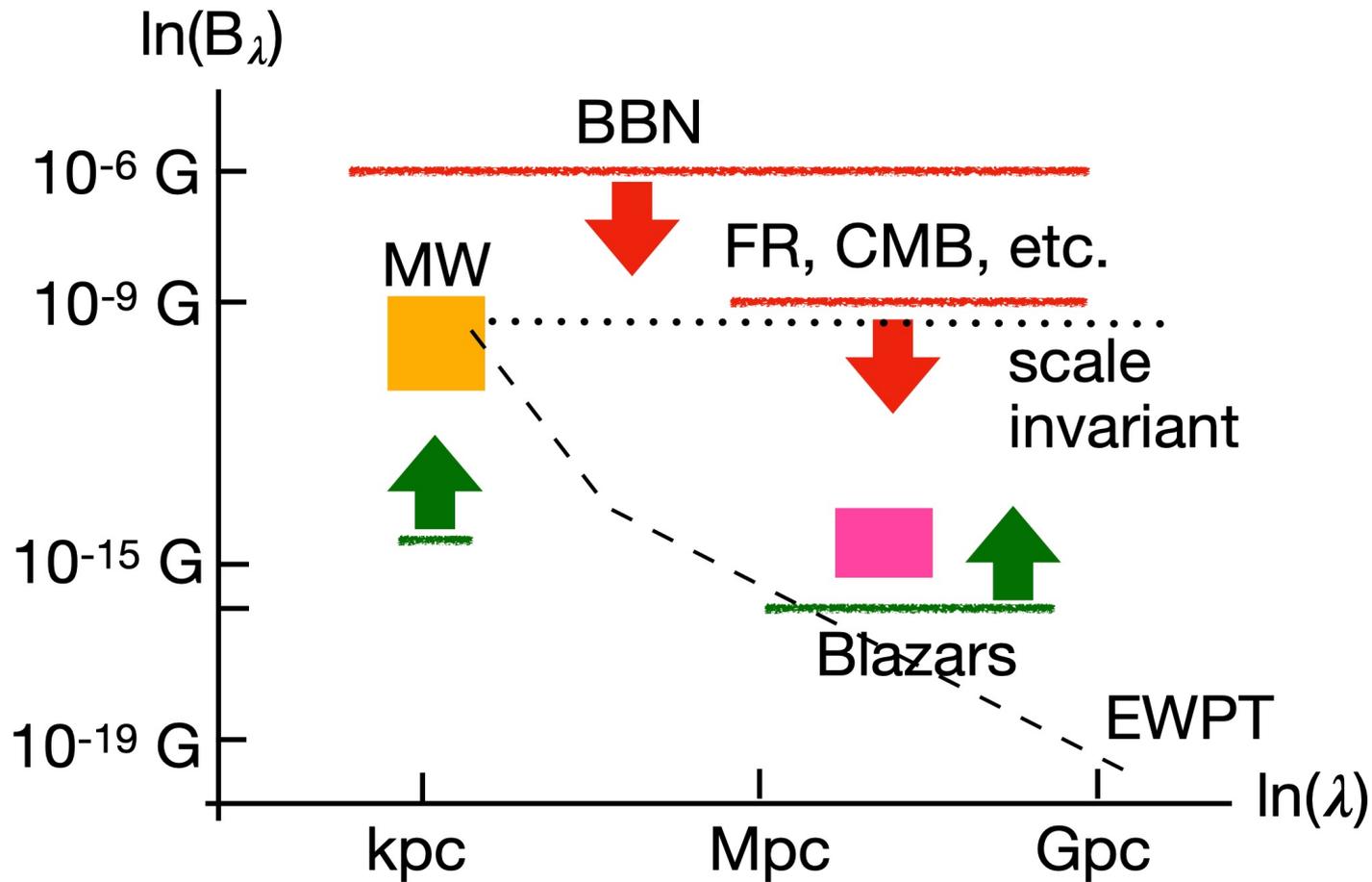
- Micro-Gauss ( $\mu\text{G}$ ) fields in galaxies and clusters
  - produced during galaxy formation via dynamo?
  - primordial origin? (need 0.01-0.1 nano-Gauss)
  - $\mu\text{G}$  fields seen in proto-galaxies that haven't turned enough times for the dynamo to work!
- Evidence of magnetic fields in voids
  - missing GeV  $\gamma$ -ray halos around TeV blazars  
*A. Neronov and I. Vovk, arXiv:1006.3504, Science (2010)*
- Magnetic fields in filaments
  - LOFAR observation of a  $\sim 3$ -10 Mpc radio emission ridge connecting two merging galaxy clusters suggests  $\sim 0.1$ -0.3  $\mu\text{G}$  fields in the filament  
*F. Govoni et al, arXiv:1906.07584, Science (2019)*
- Generated in the early universe – not “if”, but “how much”
  - phase transitions
  - inflationary mechanisms
  - a window into the early universe



# Stochastic Primordial Magnetic Field (PMF)

- Generated in the early universe, e.g. during the electroweak phase transitions or inflation
- Frozen in the plasma on large scales, the amplitude decreases with the expansion as  $B(a)=B_0/a^2$
- PMF generated in a phase transition would have most of its power on small scales
- The simplest Inflation based models predict a scale-invariant PMF

# Cosmological Magnetic Fields



# How do magnetic fields help to relieve the Hubble tension?

## In two sentences:

- Magnetic fields present in the plasma prior to recombination induce baryon inhomogeneities (clumping) on very small ( $\sim 1$  kpc) scales, speeding up the recombination  
*Jedamzik & Abel, arXiv:1108.2517, JCAP (2013); Jedamzik & Saveliev, arXiv:1804.06115, PRL (2019)*
- An earlier completion of recombination results in a smaller sound horizon at decoupling, helping to relieve the  $H_0$  tension  
*Jedamzik & LP, arXiv:2004.09487, PRL (2020)*

# Magneto-Hydro-Dynamics

Navier-Stokes:

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \underbrace{c_s^2 \frac{\nabla \rho}{\rho}}_{\text{Pressure} = c_s^2 \rho} + \underbrace{\nabla \Phi}_{\text{Gravity}} = \underbrace{\nu \nabla^2 \mathbf{v}}_{\text{Viscosity}} - \underbrace{\frac{1}{4\pi\rho} \mathbf{B} \times (\nabla \times \mathbf{B})}_{\text{Lorentz force}}$$

Continuity:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Induction:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}$$

# MHD at Recombination

Navier-Stokes: 
$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + c_s^2 \frac{\nabla \rho}{\rho} + \cancel{\nabla \Phi} = \overset{-\alpha v}{\nu \nabla^2 \mathbf{v}} - \frac{1}{4\pi\rho} \mathbf{B} \times (\nabla \times \mathbf{B})$$

Continuity: 
$$\frac{\partial \mathbf{v}}{\partial t} + \nabla(\rho \mathbf{v}) = 0$$

Induction: 
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \cancel{\eta \nabla^2 \mathbf{B}}$$

# Magnetic field induces density inhomogeneities on scales below the photon mean free path

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + c_s^2 \frac{\nabla \rho}{\rho} = -\alpha \mathbf{v} - \frac{1}{4\pi\rho} \mathbf{B} \times (\nabla \times \mathbf{B})$$

$\alpha \sim 1/l_\gamma$

$\frac{1}{2} \nabla B^2 - (\mathbf{B} \cdot \nabla) \mathbf{B}$

$c_s^2 = 1/3$  for  $L > l_\gamma$   
 $c_s^2 \ll 1$  for  $L < l_\gamma$

Drag due to photons with mean free path  $l_\gamma$

Pushes baryons towards regions of low magnetic energy density

$L > l_\gamma$       tightly coupled incompressible baryon-photon fluid

$L < l_\gamma$       viscous compressible baryon gas

Plasma develops density fluctuations on small scales  
(below the photon mean free path)

# Inhomogeneities enhance the recombination rate

$$\frac{dn_e}{dt} + 3Hn_e = -C \left( \alpha_e n_e^2 - \beta_e n_{H^0} e^{-h\nu_\alpha/T} \right)$$

The probability of a proton and an electron combining to form H is proportional to  $n_p n_e = n_e^2$

# Inhomogeneities enhance the recombination rate

$$\left\langle \frac{dn_e}{dt} + 3Hn_e = -C \left( \alpha_e n_e^2 - \beta_e n_{H^0} e^{-h\nu_\alpha/T} \right) \right\rangle$$

# Inhomogeneities enhance the recombination rate

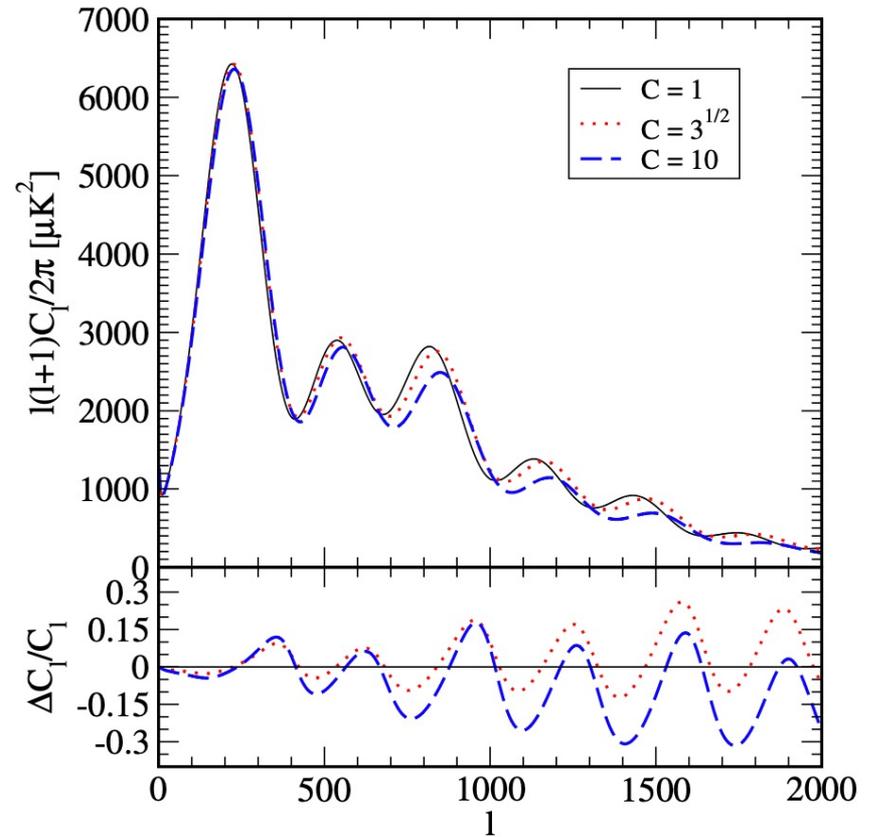
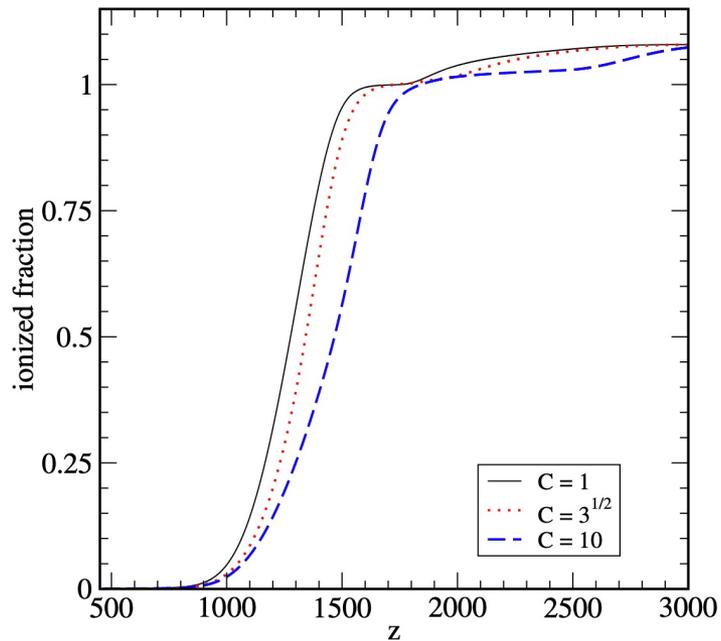
$$\left\langle \frac{dn_e}{dt} + 3Hn_e = -C \left( \alpha_e n_e^2 - \beta_e n_{H^0} e^{-h\nu_\alpha/T} \right) \right\rangle$$

$$n_e = \langle n_e \rangle + \delta n_e \quad \rightarrow \quad \langle n_e^2 \rangle > \langle n_e \rangle^2$$

# Inhomogeneities enhance the recombination rate

$$\left\langle \frac{dn_e}{dt} + 3Hn_e = -C \left( \alpha_e n_e^2 - \beta_e n_{H^0} e^{-h\nu_\alpha/T} \right) \right\rangle$$

$$\langle n_e^2 \rangle > \langle n_e \rangle^2$$



# Implementation

An additional baryon clumping parameter:

$$b = (\langle n_b^2 \rangle - \langle n_b \rangle^2) / \langle n_b \rangle^2$$

Considered two different models of baryon density distribution

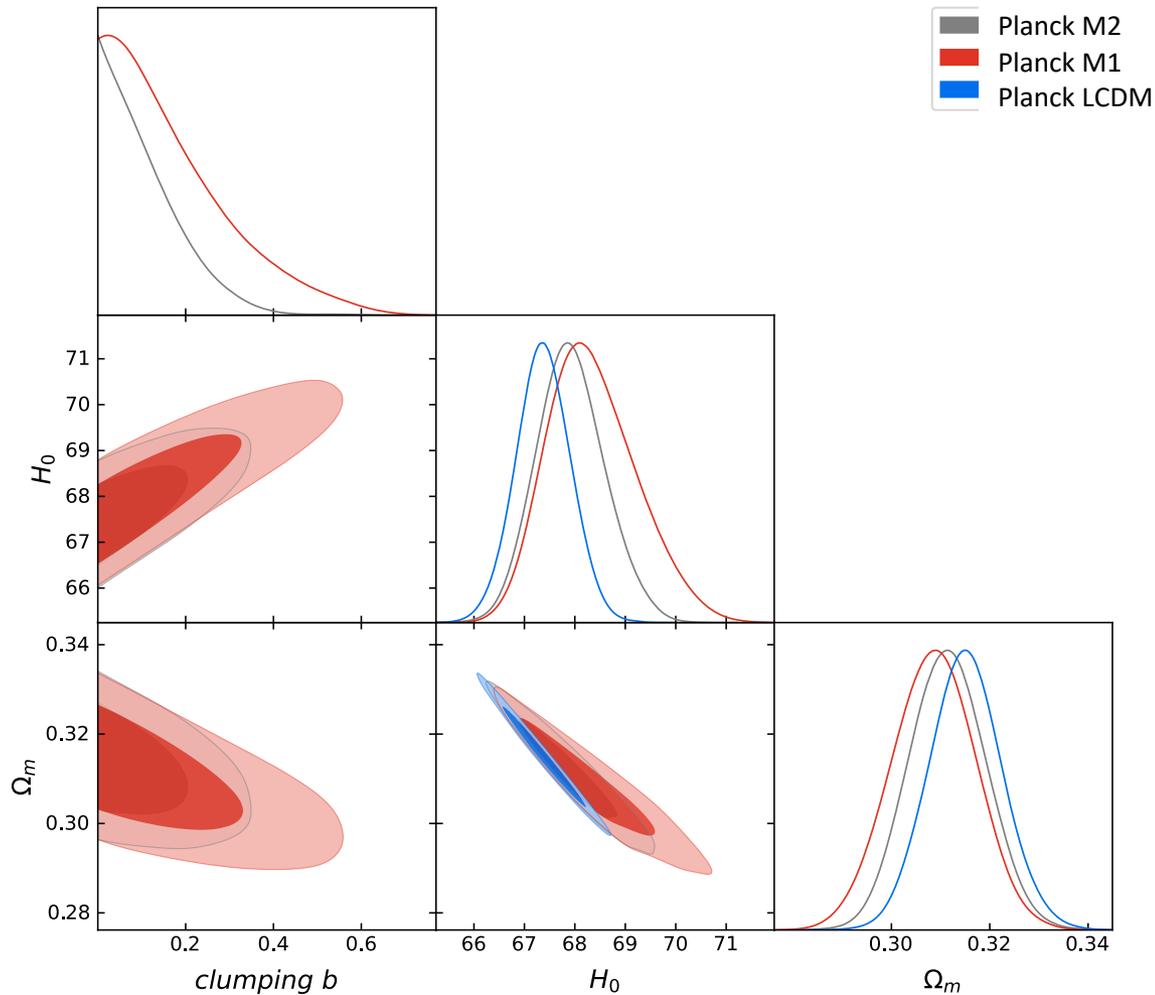
- M1, with the same baryon density PDF as in Jedamzik and Abel (2013)
- M2, using a different PDF

(The exact PDF determination from MHD simulations is in progress)

Datasets considered in this work:

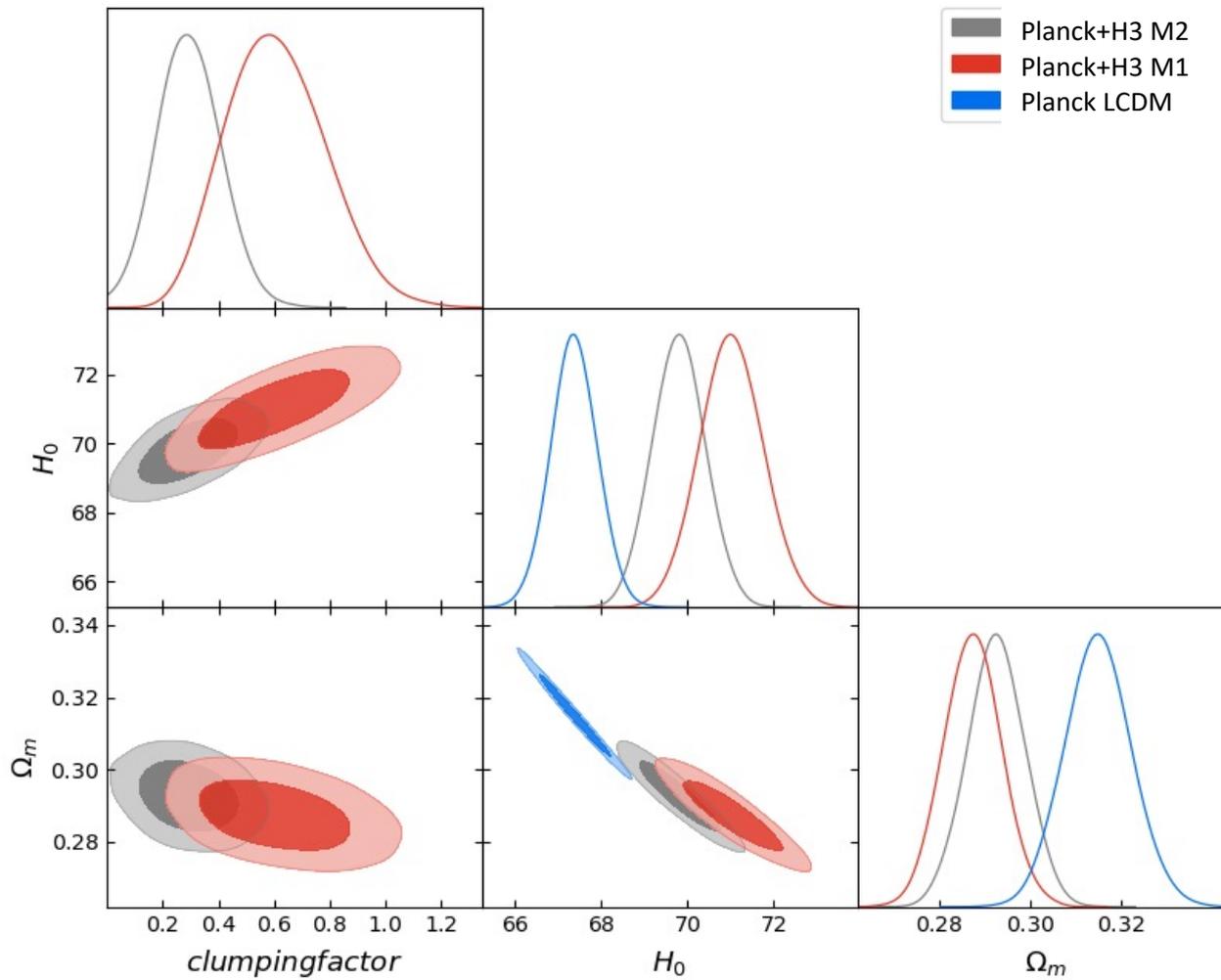
- CMB temperature and polarization from Planck 2018
- SH0ES, H0LiCOW and MCP determinations of  $H_0$  (H3)
- Baryon Acoustic Oscillations (BAO) and supernovae (SN) data

# Fitting to Planck CMB only



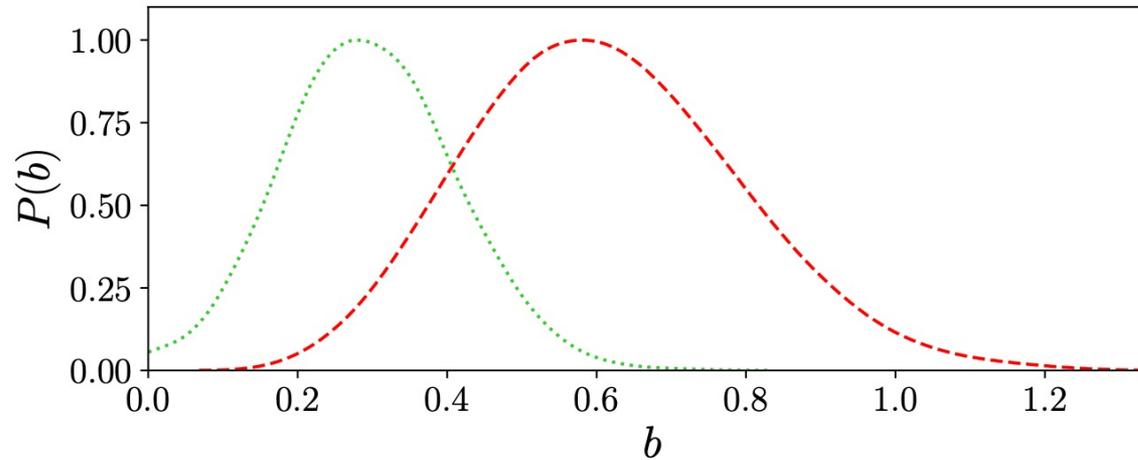
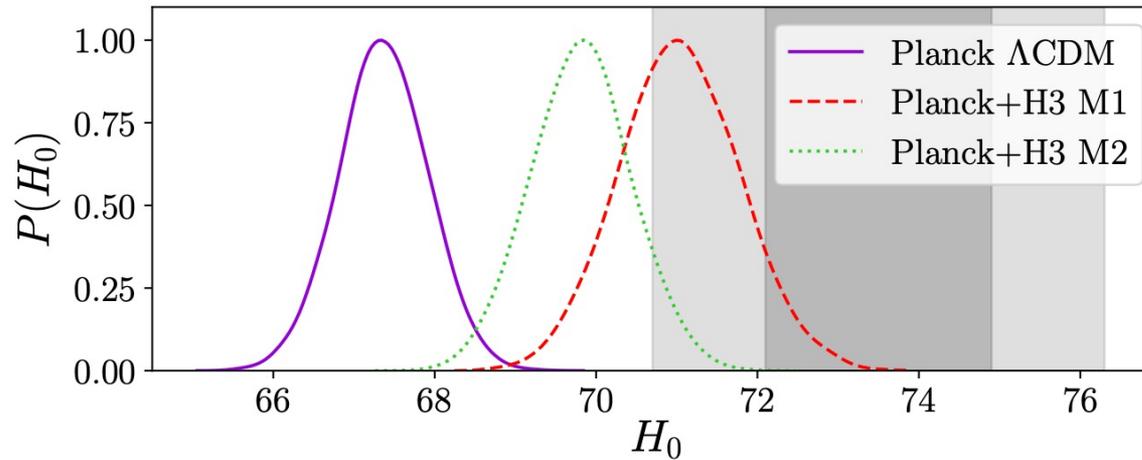
- Strong degeneracy between the clumping parameter  $b$  and  $H_0$
- No preference for a non-zero value of  $b$

# Fitting to Planck + H3

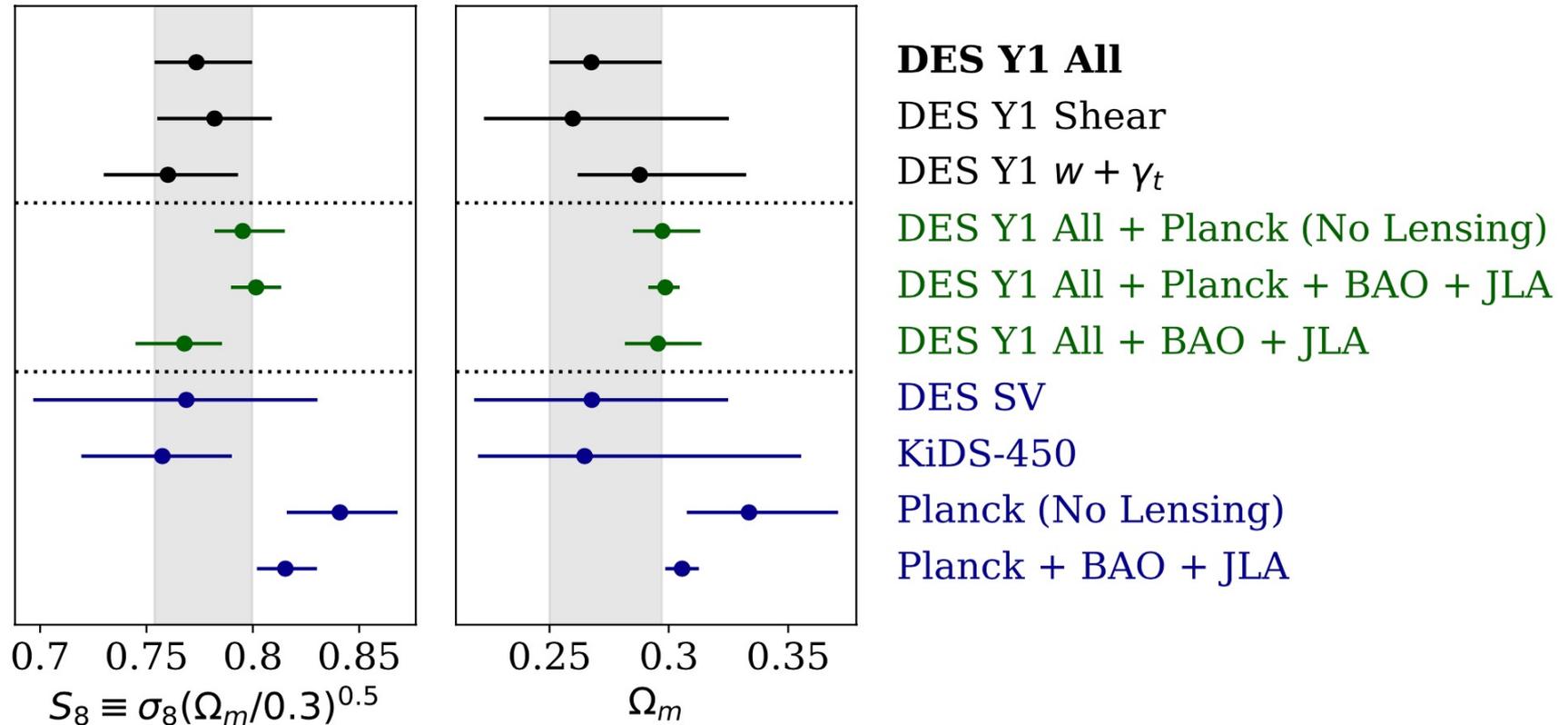


a clear detection of clumping

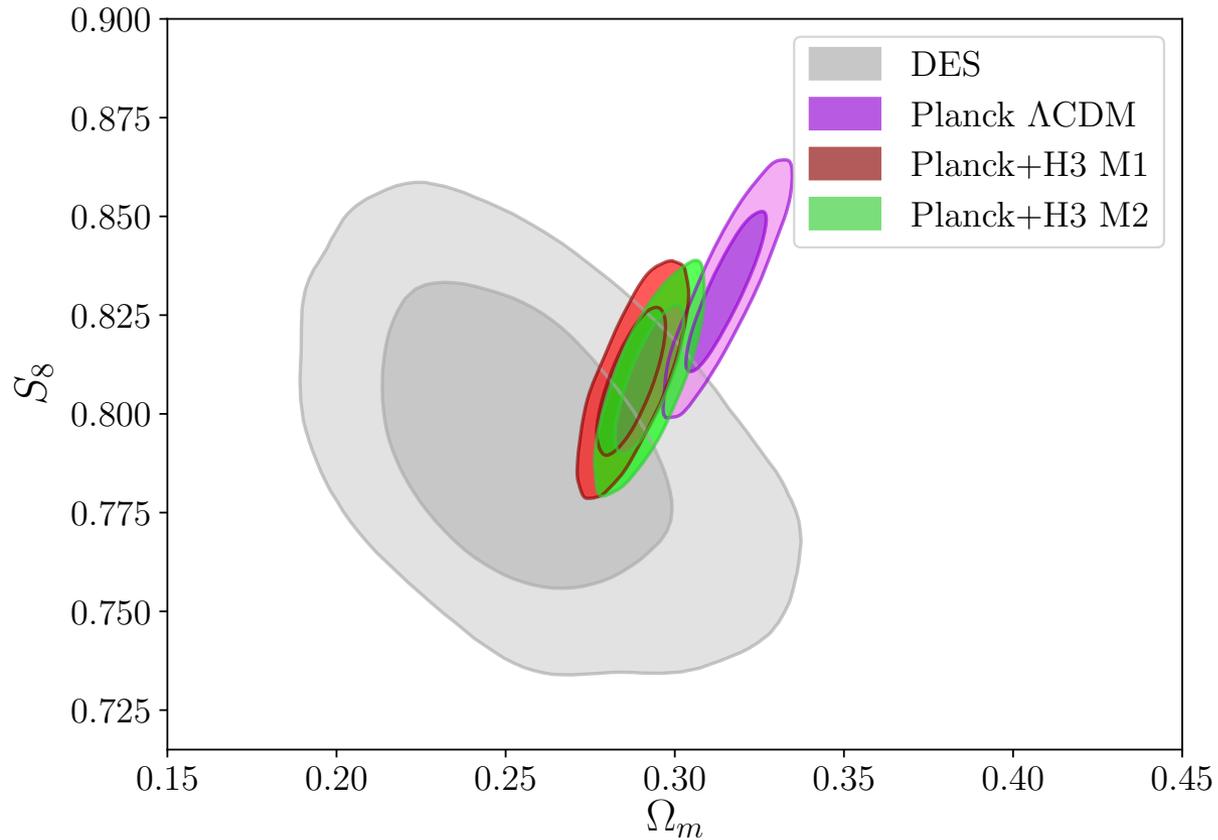
# Relieving the Hubble tension



# Another (milder) tension, cosmic matter is less clustered than predicted

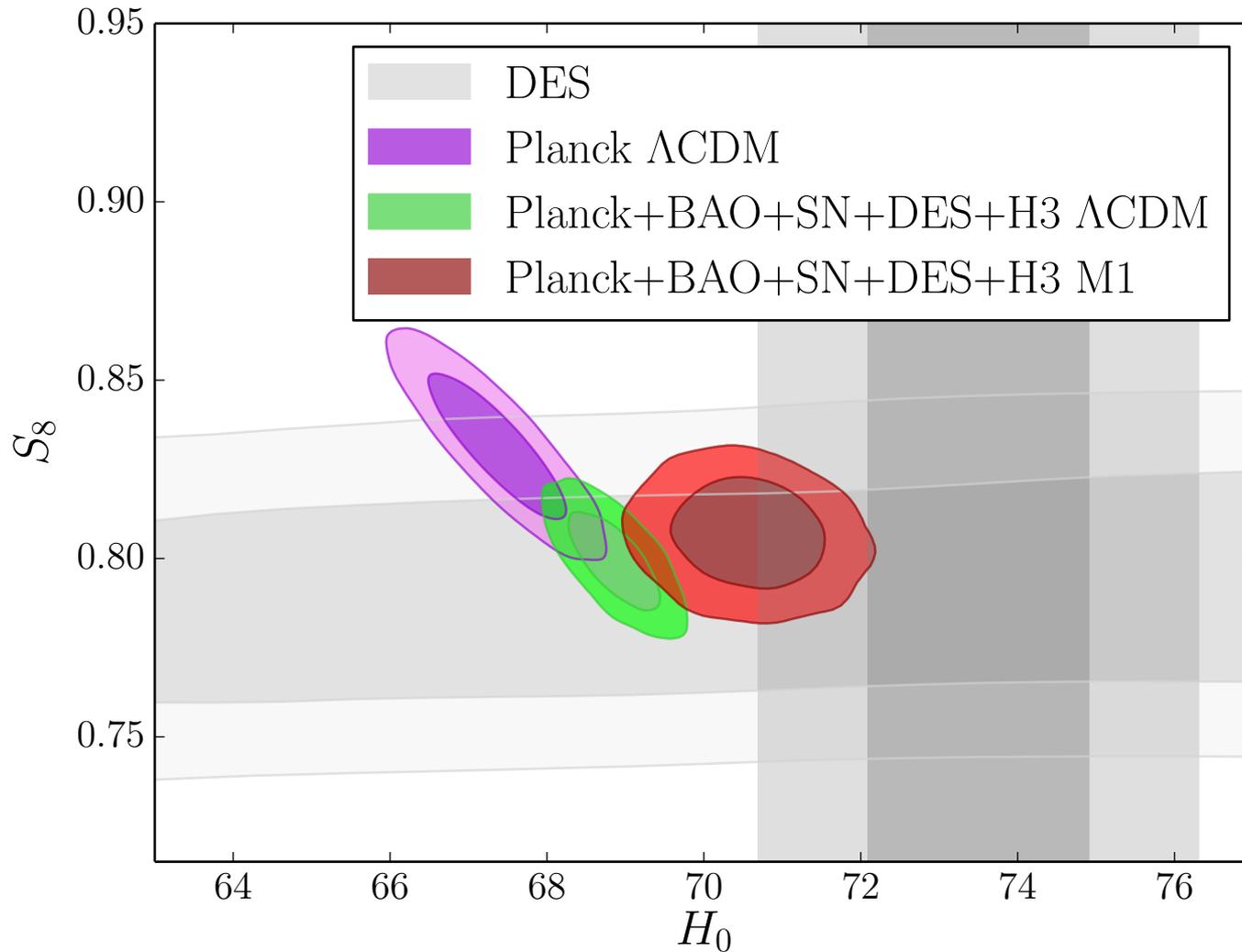


# Relieving the $S_8$ - $\Omega_m$ tension



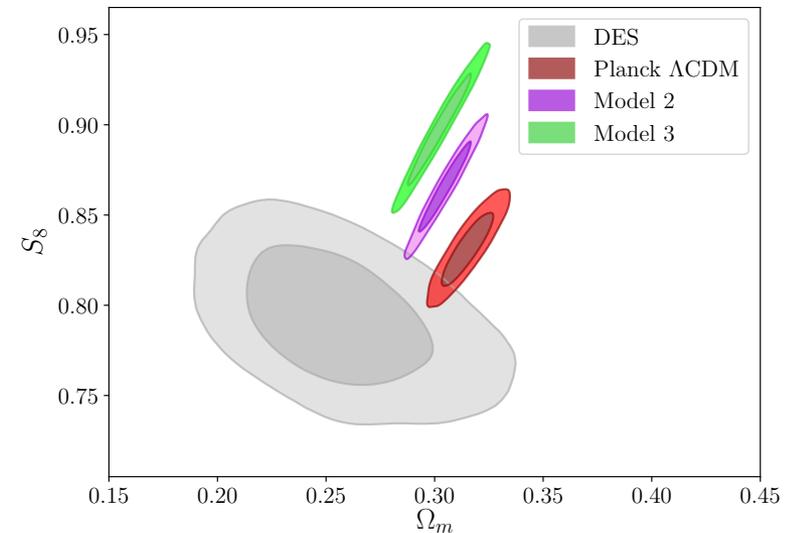
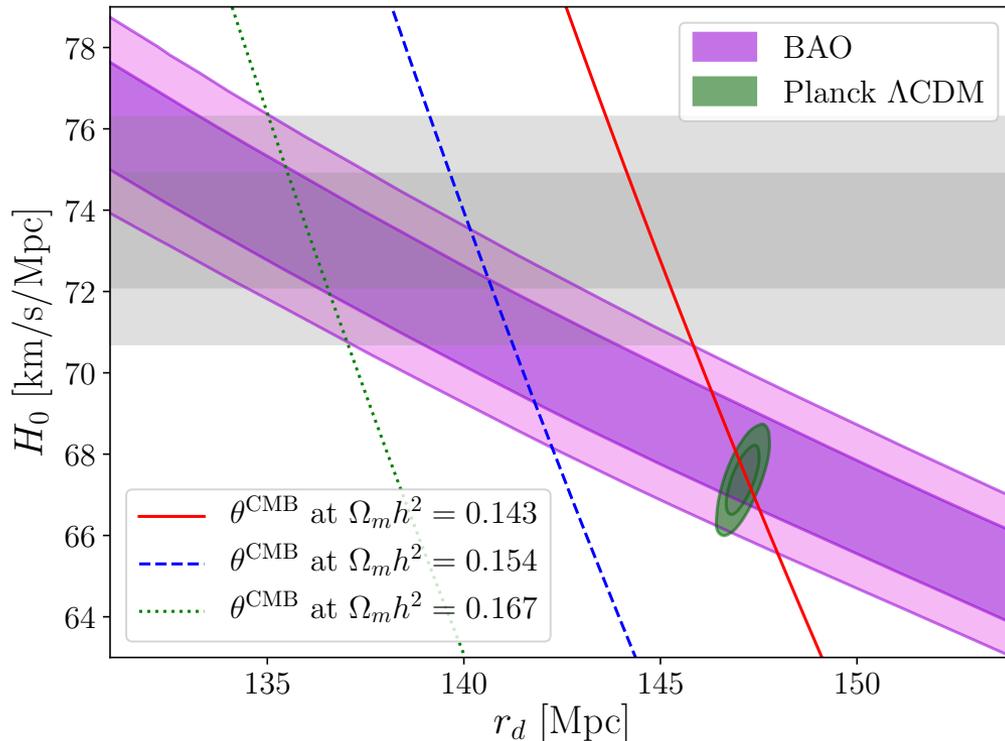
As a byproduct, clumping models help to relieve the  $S_8$ - $\Omega_m$  tension

# Fitting to all data



# Why reducing the sound horizon cannot (by itself) fully relieve the Hubble tension

- The relevant parameters are the sound horizon at decoupling,  $r_d$ ,  $h$  and  $\Omega_m h^2$
- To simultaneously fit the CMB, Baryon Acoustic Oscillations (BAO) and SHOES data one needs to **increase  $\Omega_m h^2$**
- A larger  $\Omega_m h^2$  creates **tension with weak lensing** data, e.g. DES and KiDS

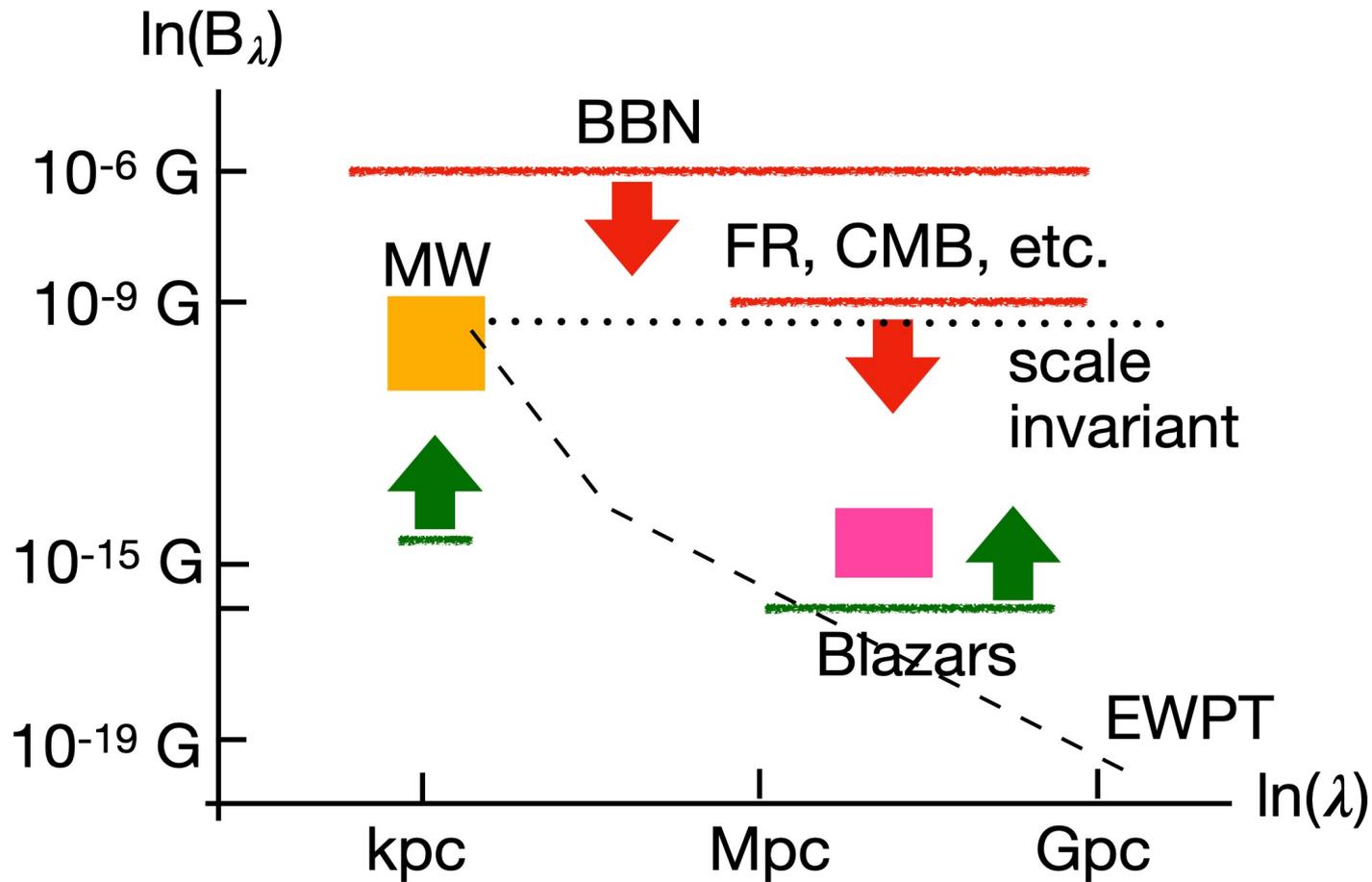


*K. Jedamzik, LP, G.-B. Zhao,  
Physics Communications (Nature) arXiv:2010.04158*

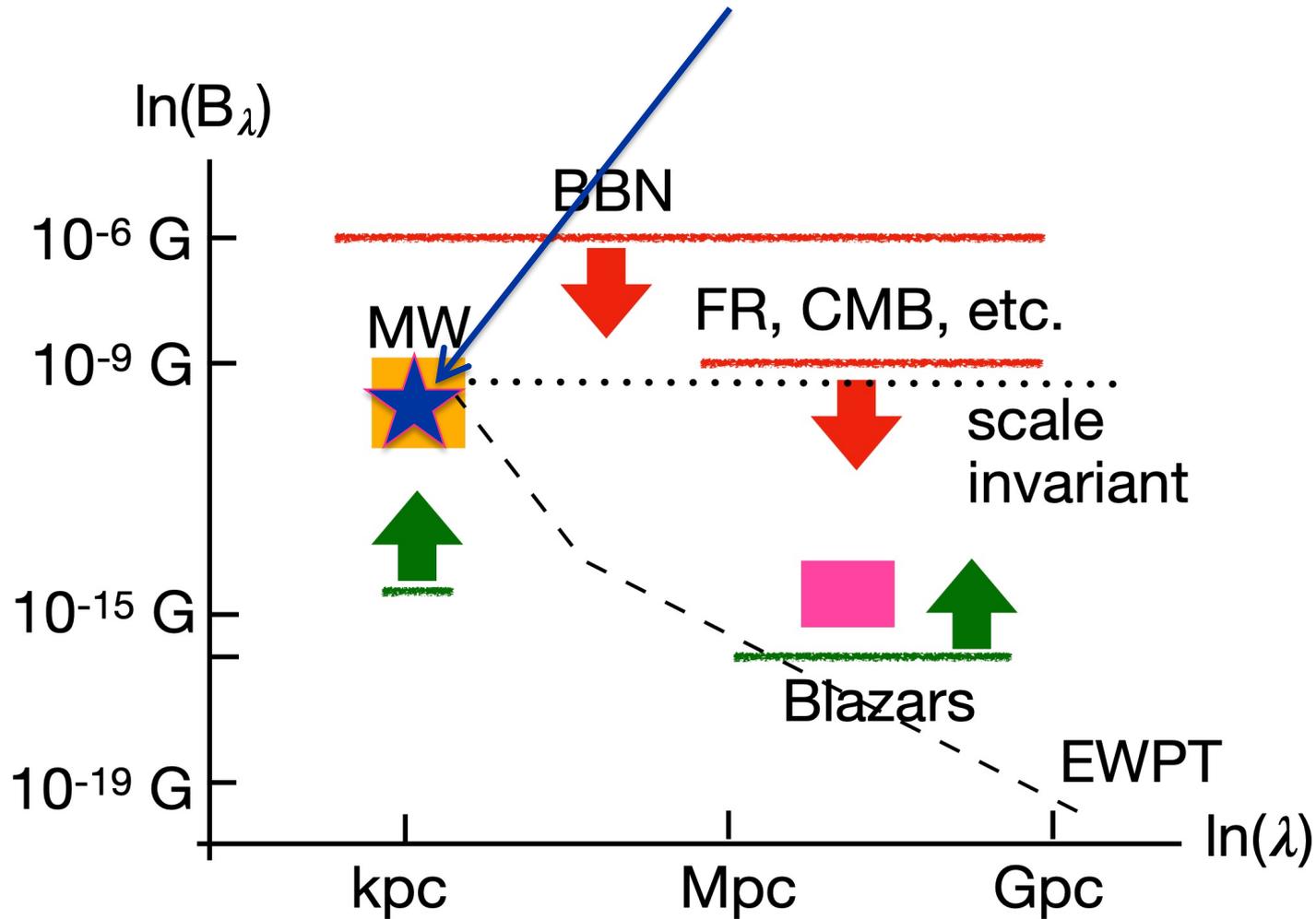
# Implications

- Magnetic fields can raise the CMB+BAO inferred  $H_0$  to  $\sim 70$  km/s/Mpc
- The amount of clumping needed for this corresponds to  $\sim 0.05-0.1$  nano-Gauss pre-recombination magnetic field

# Cosmological Magnetic Fields



# Clumping required to relieve the $H_0$ tension

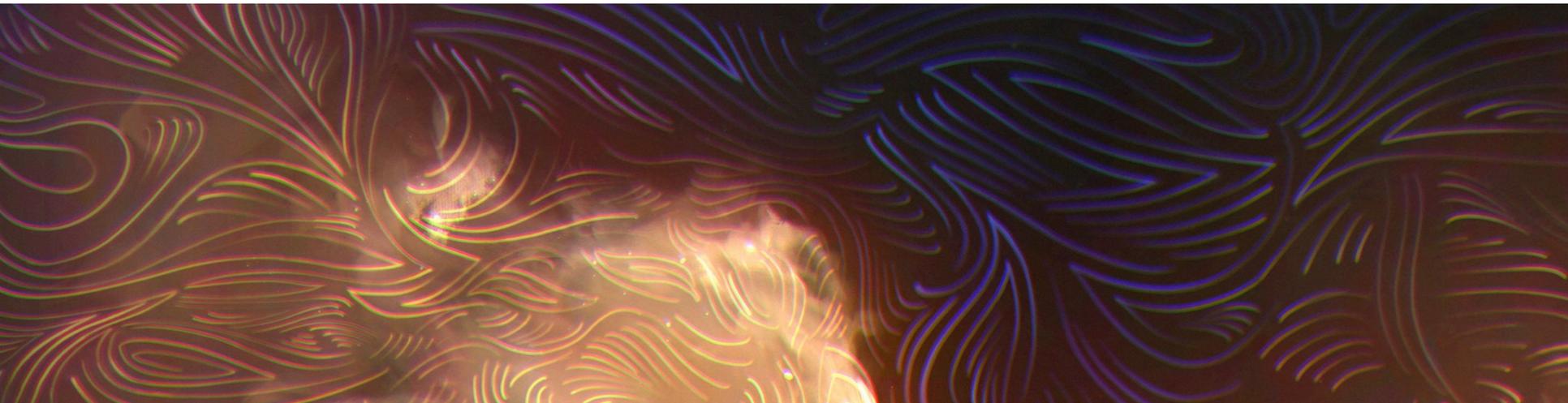


COSMOLOGY

# The Hidden Magnetic Universe Begins to Come Into View

 39 | 

*Astronomers are discovering that magnetic fields permeate much of the cosmos. If these fields date back to the Big Bang, they could solve a major cosmological mystery.*

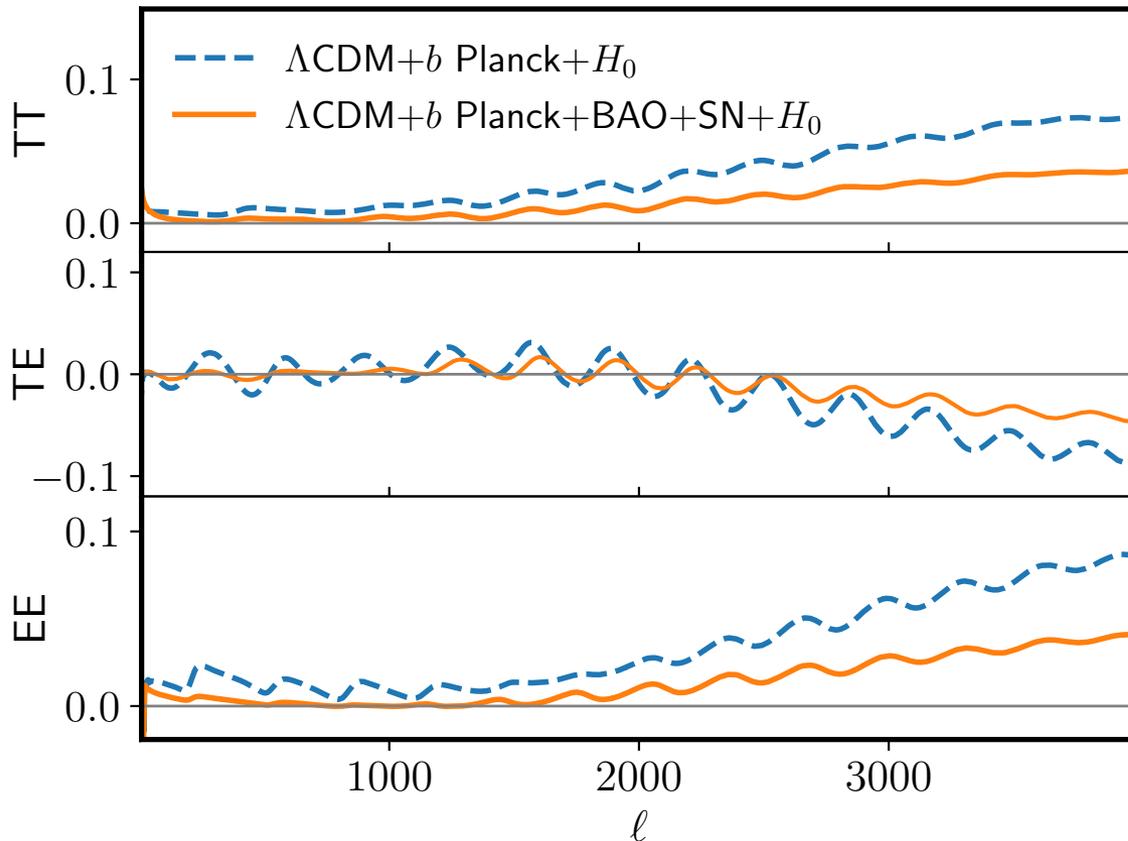


# Implications

- Magnetic fields can raise the CMB+BAO inferred  $H_0$  to  $\sim 70$  km/s/Mpc
- The amount of clumping needed for this corresponds to  $\sim 0.05$ - $0.1$  nano-Gauss pre-recombination magnetic field, which is what one would need to explain the observed galactic, cluster and intergalactic fields
- This is [a highly falsifiable proposal](#) -- future observations will rule it out or land further support
- Clumping affects the amount of Silk damping that determines the anisotropy power at the high- $l$  end of CMB spectra
- How about the recent high resolution CMB data from ACT and SPT-3G?  
*(see also Thiele et al, arXiv:2105.03003, for ACT DR4 constraints on clumping)*

# The Silk Damping Tail

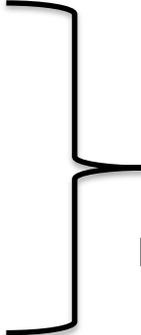
$$(C_\ell - C_\ell^{\Lambda\text{CDM}}) / C_\ell^{\Lambda\text{CDM}}$$



$\Lambda\text{CDM}$  and  $\Lambda\text{CDM}+b$  make comparable predictions for CMB Temperature (T) and polarization (E) spectra for  $\ell < 2000$ , but the differences becomes large at higher  $\ell$

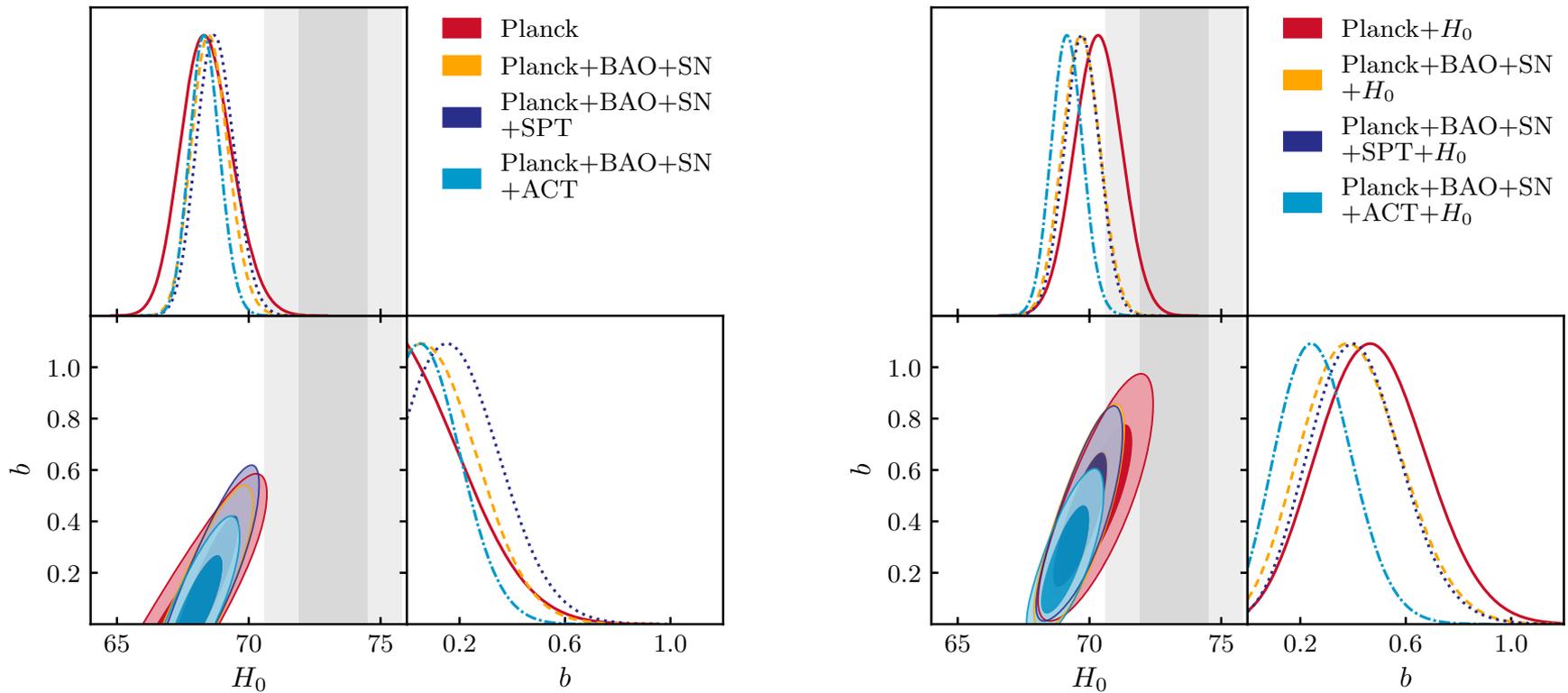
# The new data (since Spring 2020)

- New BAO from eBOSS DR16  
*Alam et al, arXiv:2007.08991*
- New SHOES,  $H_0=73.2 \pm 1.3$  km/s/Mpc  
*Riess et al, arXiv:2012.08534*
- ACT DR4 TT ( $600 < l < 4000$ ), TE and EE ( $350 < l < 4000$ )  
*Choi et al, arXiv:2007.07289*
- SPT-3G Year 1, TE and EE ( $300 < l < 3000$ )  
*Dutcher et al, arXiv:2101.01684*



Do not notably change  
clumping constraints  
based on  
DR12 BAO and 2019 SHOES

# New constraints on clumping



without SHOES

with SHOES

Planck+BAO+SN

$b < 0.47$  (95%CL),  $H_0 = 68.57 \pm 0.68$

$b = 0.42 \pm 0.18$ ,  $H_0 = 69.68 \pm 0.66$

with SPT

$b < 0.50$  (95%CL),  $H_0 = 68.73 \pm 0.64$

$b = 0.43 \pm 0.17$ ,  $H_0 = 69.74 \pm 0.61$

with ACT

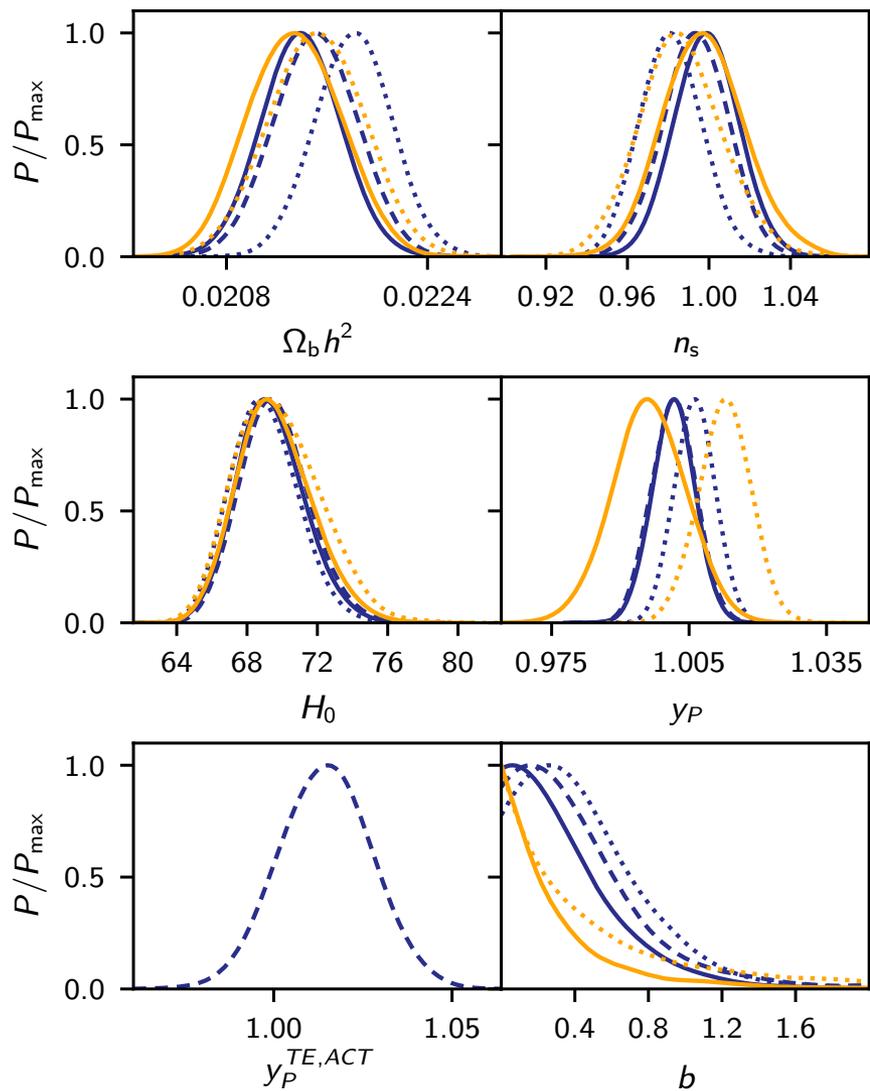
$b < 0.34$  (95%CL),  $H_0 = 68.30 \pm 0.55$

$b = 0.28 \pm 0.14$ ,  $H_0 = 69.14 \pm 0.56$

## Why is ACT DR4 so much more constraining compared to SPT-3G Y1?

- Not the 3000<l<4000 band powers
- Not the TT: ACT constraints on  $b$  get stronger when TT is removed
- LCDM based mock simulations show that ACT and SPT-3G TE+EE spectra should yield comparable constraints on  $b$ , while adding ACT TT should make them tighter
- Anomalously strong constraints coming from ACTDR4 TE+EE
- A  $2.7\sigma$  tension between Planck and ACT DR4 in LCDM can be partially resolved by a 5% re-calibration of TE ( $Y_p^{\text{TE}}=1.05$ , *Aiola et al, arXiv:2007.07288*)
- While there is no apparent physical reason for recalibrating TE, doing so significantly relaxes the ACTDR4 constraints on clumping

- ACT DR4 TE,EE
- ⋯ ACT DR4 TE,EE,  $y_p^{TE}=1.05$
- ACT DR4
- ⋯ ACT DR4,  $y_p^{TE}=1.05$
- - - ACT DR4, free  $y_p^{TE}$



# Conclusions

- The Hubble tension hints at a missing ingredient in the physics of recombination. That missing ingredient could be a primordial magnetic field of strength that happens to be of the right order to also explain the observed galactic, cluster and intergalactic fields
- This can only raise the value of  $H_0$  up to 70 km/s/Mpc (not 73.2, but it could be all we need!)
- Primordial magnetic fields were not invented to solve the Hubble tension. A detection of clumping is important by itself, as a solution of a much older puzzle and a tantalizing evidence of new physics in the early universe
- Future high resolution CMB temperature and polarization anisotropy data (Simons Observatory, CMB-S4), along with comprehensive MHD simulations, will provide a conclusive test of this scenario