

The EFT measurement problem

Gauthier Durieux
(CERN)

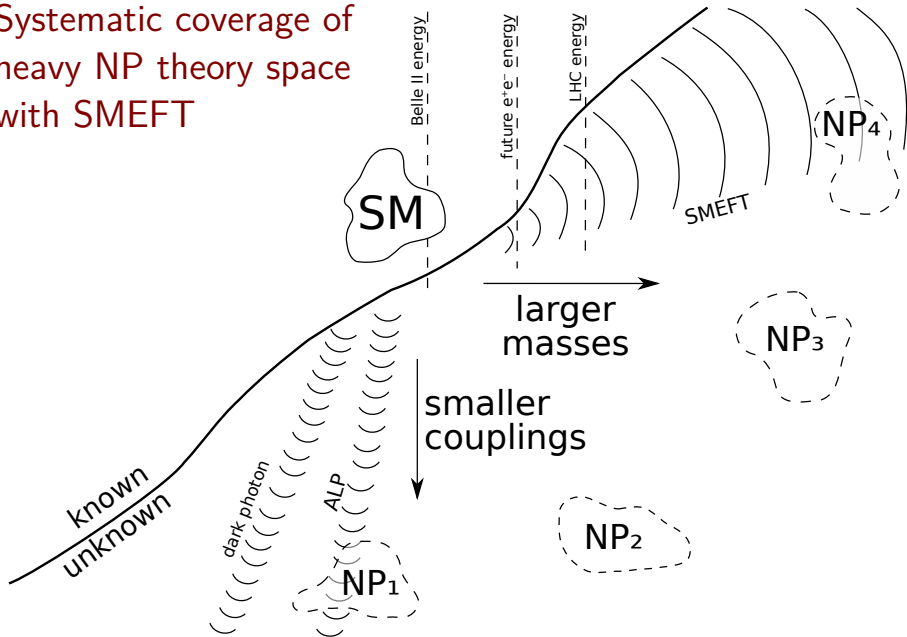
LHC EFT WG [webpage](#), [mailing list](#), [agenda](#)

see [Jan 11 meeting](#): *Experimental measurements and observables*

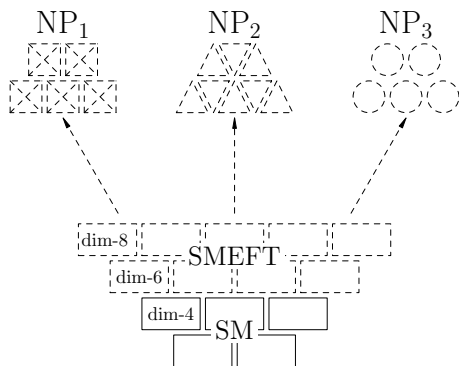
IML meeting
13 July 2021



Systematic coverage of heavy NP theory space with SMEFT



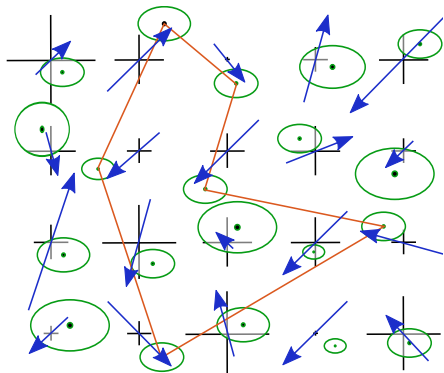
Taking the SM to higher dimensions



- using established bricks (fields and symmetries)
- deformations organised in layers (by dimension)
- including all deformations (theory space coverage)

→ a global approach

Identifying the pattern of new physics



design sensitive observables

- precise SM predictions
 - precise SMEFT predictions
 - precise measurements
- leverage correlations

EFT measurement challenges

- multidimensional parameter space
expensive optimization at every point

- suppressed linear sensitivities

$$d\sigma \sim |A_{\text{SM}}|^2 + 2C_i \text{Re}\{A_{\text{SM}}^* A_i\} + C_i C_j A_i^* A_j + \dots$$

with intricate kinematics & impacting validity

- re-interpretability
more operators, or more precise EFT predictions

- global combinations

→ compromise btw sensitivity, practicality and usefulness

Statistical optimality

Matrix element method

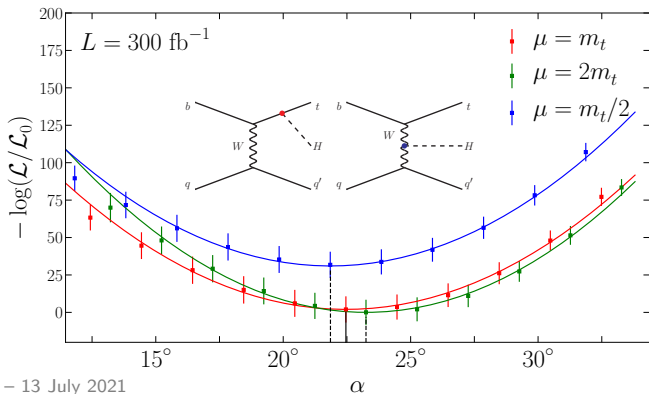
$$L(\theta|\{x_n\}) = \prod_{n=1}^N p(x_n|\theta) \quad \text{with} \quad p(x|\theta) = \frac{1}{\sigma(\theta)} \int dz \frac{d\sigma(z|\theta)}{dz} W(z, x)$$

possibly adding total rate information

Mostly 1D parameter space

Demanding when $W(z, x) \neq \delta(z, x)$: partial or imperfect measurements

Make the most of limited data, without training samples



- fNLO = Born + Virtual + Real
with virtual and real occupying different phase-spaces (n vs. $n + 1$)
and IR divergences cancelling between them

- ‘jet events’ instead of parton events
intrinsic inclusiveness guaranteeing IR safety
for a given jet clustering algorithm
integrating out soft and collinear radiations

$$\frac{d^r \sigma^{\text{NLO}}}{dx_1 \dots dx_r} = \frac{d^r \sigma^{\text{BV}}}{dx_1 \dots dx_r} + \sum_i \sum_{\substack{j \\ j \neq i}} \frac{d^r \sigma_{\mathcal{R}_{ij}}^{\text{R}}}{dx_1 \dots dx_r} + \sum_i \frac{d^r \sigma_{\mathcal{R}_i}^{\text{R}}}{dx_1 \dots dx_r} + \sum_i \frac{d^r \sigma_{\bar{\mathcal{R}}_i}^{\text{R}}}{dx_1 \dots dx_r}$$

- phase space parametrization with variables
already non-trivial at LO
not allowing reconstruction of
 - jet invariant masses
 - overall transverse momentum

FNLO QCD corrections to event weights

e.g. in $pp \rightarrow t j$ use $x = \{\eta_t, E_j, \eta_j, \phi_j\}$

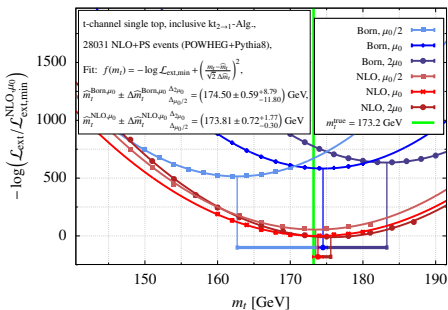
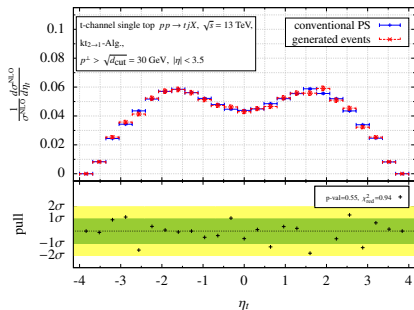
$$J_t = (E_t, -J^\perp \cos \phi_j, -J^\perp \sin \phi_j, J^\perp \sinh \eta_t)$$

$$J_j = (E_j, J^\perp \cos \phi_j, J^\perp \sin \phi_j, J^\perp \sinh \eta_j)$$

with

$$J^\perp = J_t^\perp = J_j^\perp = \frac{\sqrt{E_j^2 - J_j^2}}{\cosh \eta_j}$$

$$E_t = \sqrt{J^{\perp 2} \cosh^2 \eta_t + J_t^2}$$



Addressing multidimensionality

Statistically optimal observables

[Atwood,Soni '92]
[Diehl,Nachtmann '94]

- discrete set (one number per $\{x_n\}$ for each C_i)
- amenable to usual studies
 - higher-order effects
 - systematic uncertaintiesmost easily on LO parton-level definition (i.e. optimization)
- optimized at one point (e.g. $\{C_i = 0 \forall i\}$)
- for linear sensitivity
 - used for $ee \rightarrow WW$ at LEP2 [Opal, L3, ALEPH, DELPHI]
 - similar to MELA discriminant averages

Statistically optimal observables

minimize the one-sigma ellipsoid in target parameter space

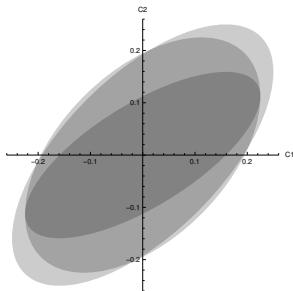
(joint efficient set of estimators, saturating the Cramér-Rao bound: $V^{-1} = I$)

For small C_i , with a phase-space distribution $\sigma(x) = \sigma_{SM}(x) + \sum_i C_i \sigma_i(x)$,

The stat. optimal observables: $OO_i \equiv \left\langle n \sigma_i(x) / \sigma_{SM}(x) \right\rangle_{\{x_n\}}$

The associated covariance:

$$\text{cov}(C_i, C_j)^{-1} = \epsilon \mathcal{L} \int dx \frac{\sigma_i(x) \sigma_j(x)}{\sigma_{SM}(x)} + \mathcal{O}(C_k).$$



e.g. $\sigma(\phi) = 1 + \cos(\phi) + C_1 \sin(\phi) + C_2 \sin(2\phi)$

1. asymmetries $\sim \langle \text{sign}\{\sin(i\phi)\} \rangle$

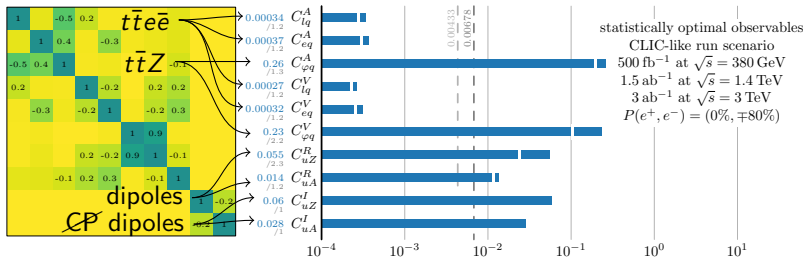
2. moments $\sim \langle \sin(i\phi) \rangle$

3. statistically optimal $\sim \left\langle \frac{\sin(i\phi)}{1 + \cos \phi} \right\rangle$

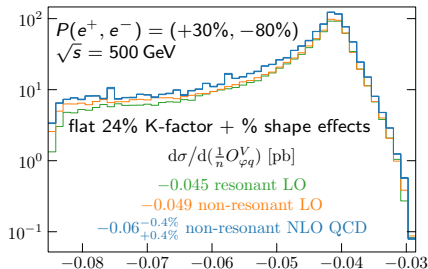
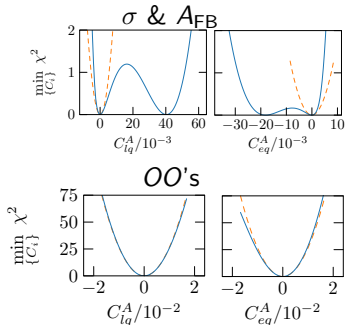
\Rightarrow area ratios 1.9 : 1.7 : 1

$e^+e^- \rightarrow t\bar{t}$ application

[Grzadkowski, Hioki '00], [Janot '15]
 [GD, Perelló, Vos, Zhang '18]
 [CLICdp '18]



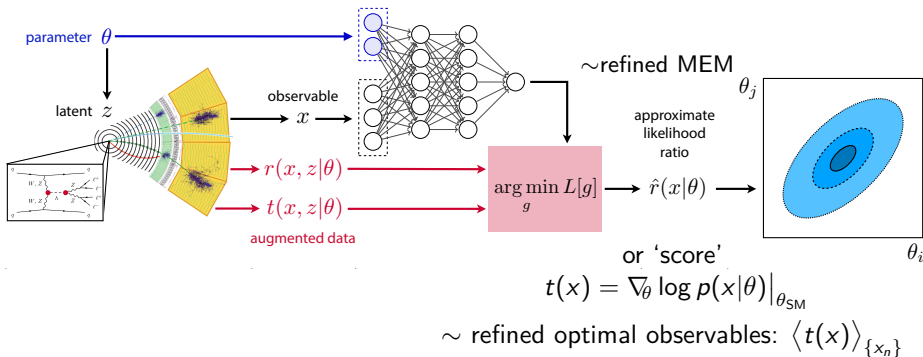
7.8% efficiency vs. $A_{FB} : 4.7\%$ ($\sigma : 12.9\%$)
 semileptonic $t\bar{t}$, $\sqrt{s} = 380$ GeV, $P(e^-) = -80\%$



Integration, shower and detector

Integration, shower and detector

- extrapolate from truth to reco, from training sample
- avoid integration over unmeasured phase-space and transfer fct.
- include shower, hadronization and detector effects



Re-interpreting

see also [LHC Re-interpretation Forum](#)

Reweighting

Full truth distribution $p_0(z)$

Reconstructed distributions $p_0(x)$

Abstract *folding* operator $W(x|z)$

$$p_0(x) = \int dz p_0(z) W(x|z) \quad \text{through MC sampling}$$

Reweighting from theory hypothesis 0 to 1

given a set of pre-folded events (x_k, z_k) distributed following $p_0(z)$

$$p_1(x) = \int dz p_0(z) \frac{p_1(z)}{p_0(z)} W(x|z)$$

Recipe:

1. take events $\{x_k, z_k, p_0(z_k)\}$ (published?!)
2. reweight by $p_1(z_k)/p_0(z_k)$ (no new sample generation)
3. get arbitrarily fine-grained and multidimensional reco. $p_1(x)$
4. change observable too, if x is rich enough

The EFT measurement problem

Compromise btw sensitivity, practicality, usefulness?

Cover multidimensional spaces?

Re-interpretability?

Global combinations?

Use of reweighting and fNLO weights?

Enhanced EFT validity?