

SMASH hadronic transport

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Physical idea: separation of hard and soft physics

- Do you know what **hard** and **soft** \equiv **bulk** physics means in context of heavy ion collisions?
Press yes/no. Take 60 seconds to write it, after 60 seconds post in the chat.

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Imagine a fountain of water (jet) shooting through a dense fog (bulk)
or a jet in jacuzzi

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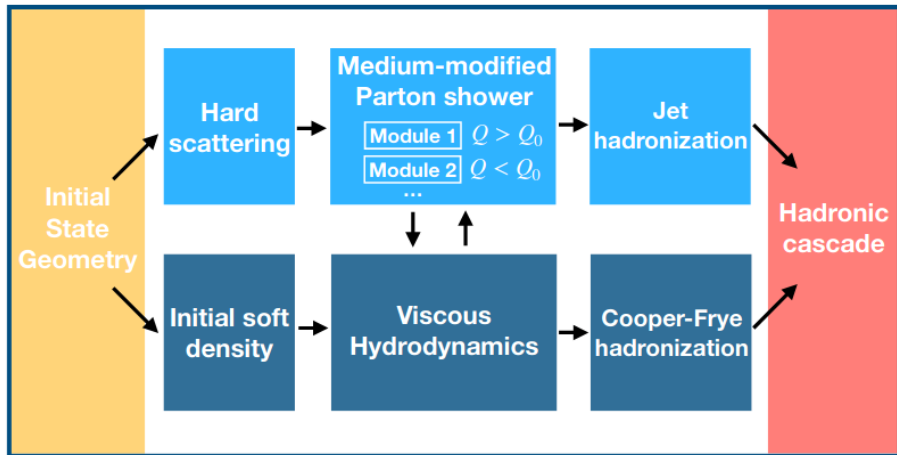
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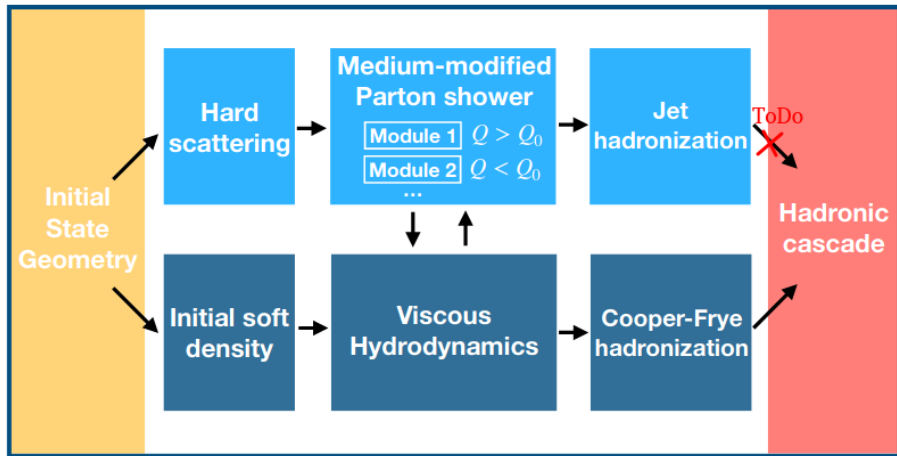
JETSCAPE

Hard and soft / bulk physics separated explicitly



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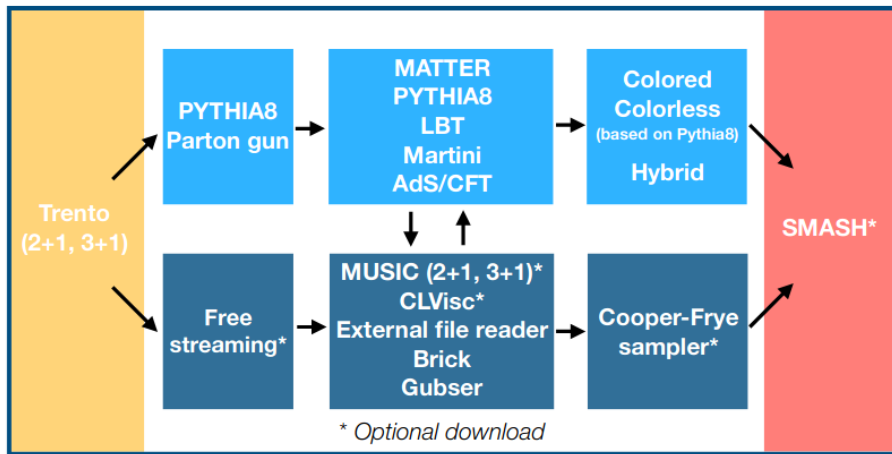
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Kudos to James Mulligan for nice illustrations

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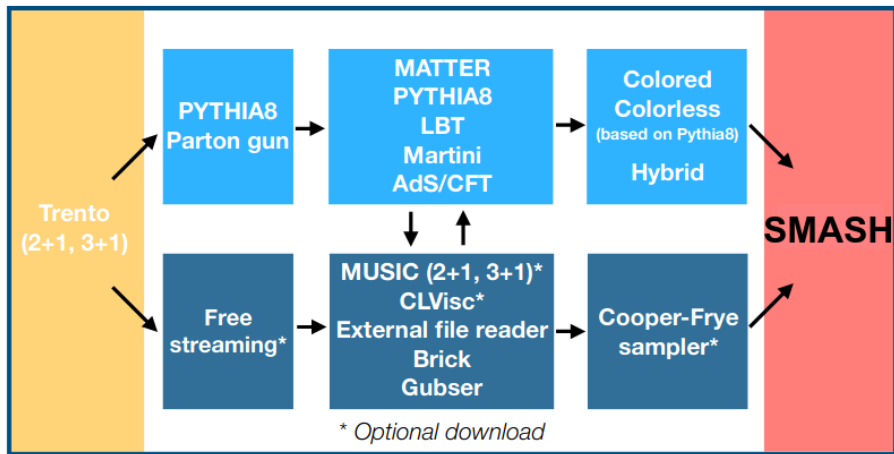
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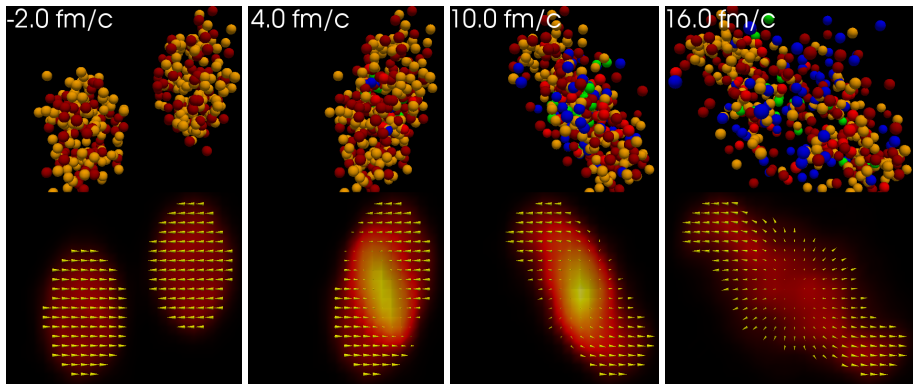
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Transport is simple: particles propagate, collide, decay

Au+Au, $\sqrt{s_{NN}} = 3$ GeV, $b = 5$ fm



... but the devil is in the details

Applications of hadronic transport

- Full simulations of ion collisions at lower and intermediate energies ($\sqrt{s} \lesssim 20$ GeV, e.g. SMASH, UrQMD, HSD, IQMD), at higher energy they tend to undershoot v_2
- With some partonic part: full simulations at any energies (e.g. PHSD, AMPT)
- Higher energies ion collisions ($\sqrt{s} \gtrsim 20$ GeV) as a **hadronic afterburner** after hydrodynamics
- $e + A$, $\nu + A$, e.g. GiBUU
- Participate in simulations of air-shower from cosmic rays, e.g. UrQMD

Some theoretical foundations

- Conceptually transport codes rely on Vlasov and Boltzmann equations
- Have you heard about Vlasov and Boltzmann equations before? Press yes/no. If yes, take 60 seconds to write 1-2 random facts you know about them. After 60 seconds post it in the slack chat.

Vlasov equation (non-relativistic version)

Motion of particles in self-generated mean field

$$\frac{\partial}{\partial t} f(t, \vec{r}, \vec{p}) + \frac{\vec{p}}{m} \nabla_{\vec{r}} f(t, \vec{r}, \vec{p}) - \nabla_{\vec{r}} U(\vec{r}) \nabla_{\vec{p}} f(t, \vec{r}, \vec{p}) = 0$$

$$U(\vec{r}) = \int d^3 r' d^3 p V(\vec{r} - \vec{r}') f(t, \vec{r}', \vec{p})$$

Classical single-particle equations of motion:

$$\frac{d\vec{r}}{dt} = \frac{\vec{p}}{m}$$

$$\frac{d\vec{p}}{dt} = -\nabla_{\vec{r}} U(\vec{r})$$

$$\frac{df(t, \vec{r}, \vec{p})}{dt} = \left(\frac{\partial}{\partial t} + \dot{\vec{r}} \nabla_{\vec{r}} + \dot{\vec{p}} \nabla_{\vec{p}} \right) f = 0$$

Easy way to think: f is number of particles per $d^3x d^3p$. Conserving number of particles and phase space volume (Liouville theorem).

Boltzmann equation (non-relativistic)

Neglect quantum effects like interference, assume 2-body correlations are local in space-time (space and time span of collisions \ll mean free path). Same left side as for Vlasov equation, but there is right side responsible for collisions.

$$\frac{df(t, \vec{r}, \vec{p})}{dt} = \frac{\partial}{\partial t} f(t, \vec{r}, \vec{p}) + \frac{\vec{p}}{m} \nabla_{\vec{r}} f(t, \vec{r}, \vec{p}) - \nabla_{\vec{r}} U(\vec{r}) \nabla_{\vec{p}} f(t, \vec{r}, \vec{p}) = \left(\frac{\partial f}{\partial t} \right)_{coll}$$

$$I_{coll} \equiv \left(\frac{\partial f}{\partial t} \right)_{coll} = \left(\frac{\partial f}{\partial t} \right)_{gain} - \left(\frac{\partial f}{\partial t} \right)_{loss}$$

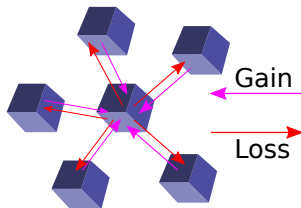
Boltzmann equation: gain and loss terms

Number of particles $dN(t, r, p)$ in the phase-space cell $d^3\vec{r}d^3\vec{p}$:

$$\frac{d}{dt}N(t, r, p) = dN_{coll}(p', \dots \rightarrow p, \dots) - dN_{coll}(p, \dots \rightarrow p', \dots)$$

$$dN(t, r, p) = f(t, r, p)d^3\vec{r}d^3\vec{p}$$

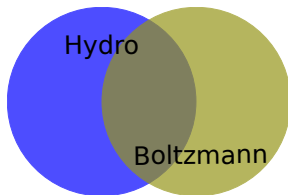
- Assumptions to calculate dN_{coll} :
 - only $2 \rightarrow 2$ scattering, neglect many-particle scatterings as rare
 - incoming particles uncorrelated
 - separation between long- and short-range interactions



$$I_{coll} = \int \frac{d^3p_2}{E_2} \frac{d^3p'_1}{E_1} \frac{d^3p'_2}{E'_2} \times W(p, p_2 \rightarrow p'_1, p'_2) \times (f(p'_1)f(p'_2) - f(p)f(p_2))$$

Ideal hydro follows from equilibrated Boltzmann

Regions of applicability for hydro and Boltzmann



Regardless of cross sections if one waits long enough then entropy reaches maximum (H-theorem). With corresponding equilibrium distribution

$$f_0(r, p) = \exp((-p^\mu u_\mu + \mu(r))/T(r))$$

right hand side of Boltzmann equation vanishes.

From $p^\mu \frac{\partial f_0}{\partial x^\mu} = 0$ follows $\partial_\mu T^{\mu\nu} = 0$ and $\partial_\mu j^\mu = 0$, where

$$T^{\mu\nu} = \int \frac{d^3 p}{(2\pi\hbar)^3} \frac{p^\mu p^\nu}{p^0} f(p) \quad \text{and} \quad j^\mu = \int \frac{d^3 p}{(2\pi\hbar)^3} \frac{p^\mu}{p^0} f(p).$$

Solving coupled Boltzmann equations in practice

- All (or almost all) hadron species known: $\pi, \rho, K, a_2, f_1, \phi, \dots, N, \Delta(1232), N(1440), \dots$: more than 100 species without accounting for charges
- Solve coupled equations, $D \equiv \frac{d}{dt}$:

$$Df_\pi = I_{coll}(f_\pi, f_N, f_\Delta, \dots)$$

$$Df_N = I_{coll}(f_\pi, f_N, f_\Delta, \dots)$$

$$Df_\Delta = I_{coll}(f_\pi, f_N, f_\Delta, \dots)$$

...

- Left hand side: testparticle ansatz

$$f \sim \frac{1}{N_{test}} \sum_i^{N_{test}} \delta^4(x - x_i) \delta^4(p - p_i) \delta(p^\mu p_\mu = m^2)$$

- Collision integrals: Monte-Carlo approach

Treatment of potentials in transport codes

Have you heard terms QMD and BUU before? Press yes/no.

Treatment of potentials in transport codes

- Boltzmann - Ühling - Uhlenbeck (BUU) approach
 - Mean-field potentials depend on densities: $U = U(\rho(\{\vec{r}_1, \vec{r}_2, \dots\}))$
 - Utilize testparticle ansatz: $N \rightarrow N \cdot N_{test}, \sigma \rightarrow \sigma/N_{test}$
 - Precise energy and momentum conservation only in the limit $N_{test} \rightarrow \infty$
 - Solve Boltzmann equations in the limit of $N_{test} \rightarrow \infty$
 - No correlations in the limit $N_{test} \rightarrow \infty$
- Quantum Molecular Dynamics (QMD) approach
 - Pairwise potentials depend on coordinates $U(r_{12})$
 - Energy and momentum conserved exactly event-by-event
 - Does not solve any particular equation for distribution function

Treatment of collisions in transport codes

- A) Geometrical criterion: $d_{ij} \leq \sqrt{\sigma/\pi}$ (d_{ij} in the CM frame of colliding particles)
 - Only allows $2 \rightarrow 2$, not $3 \rightarrow 2$ or $3 \rightarrow 1$: detailed balance violation
 - Collision time: time of closest approach
 - Sort by collision time, perform the earliest
 - Time sorting depends on frame, problems with Lorentz-invariance
 - Kodama criterion: smaller Lorentz-invariance troubles
- B) Stochastic rates: choose two random particles in cell and collide with some probability

Cassing, NPA 700, 618 (2002); Xu and Greiner, PRC 71, 064901 (2005)

- No problems with Lorentz-invariance
- Allows $3 \rightarrow 2$ or $3 \rightarrow 1$ collisions
- Not applicable with QMD
- Needs care concerning cell size and timestep
- So far only in BAMPS and (P)HSD, recently implemented in SMASH

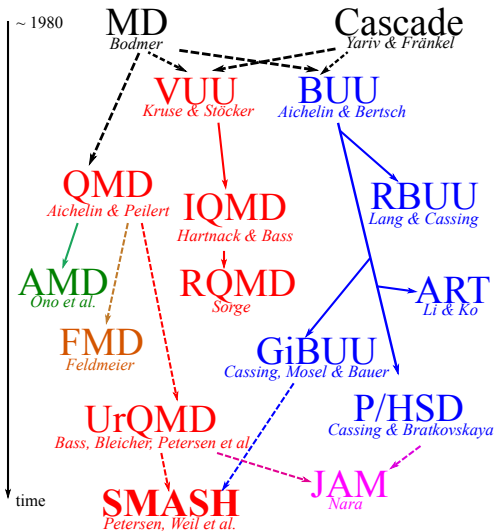
String models: $\sqrt{s} > 3 - 4 \text{ GeV}$

- Pick up $q\bar{q}$ or quark-diquark from colliding hadrons
- They form a string, which undergoes a sequence of decays: “string fragmentation”
- Different models of string fragmentation: FRITIOF (Lund), PYTHIA (Lund), HERWIG
- In Lund string: tunnelling through QCD potential

$$\mathcal{P} \sim \exp\left(\frac{-\pi p_{\perp}^2}{\kappa}\right) \exp\left(-\frac{\pi m_q^2}{\kappa}\right)$$

- Flavor composition, longitudinal momenta $\rightarrow \simeq 10$ parameters

SMASH and its ancestors



- SMASH :
Simulating
Many
Accelerated
Strongly-interacting
Hadrons
- first C++ code in this historical chain
- written from scratch
- started coding in 2013
- under *git* version control from the very beginning

thanks to Steffen Bass

SMASH: general properties

J. Weil *et al.*, Phys.Rev. C94 (2016) no.5, 054905

- Monte-Carlo solver of relativistic Boltzmann equations

BUU type approach, testparticles ansatz: $N \rightarrow N \cdot N_{test}, \sigma \rightarrow \sigma/N_{test}$

- Degrees of freedom

- most of established hadrons from PDG up to mass 2.5 GeV
- strings: do not propagate, only form and decay to hadrons
- leptons and photons production, decoupled from hadronic evolution

- Propagate from action to action (timesteps only for potentials)
action \equiv collision, decay, wall crossing

- Geometrical collision criterion: $d_{ij} \leq \sqrt{\sigma/\pi}$

- Interactions: $2 \leftrightarrow 2$ and $2 \rightarrow 1$ collisions, decays, potentials, string formation (soft - SMASH, hard - Pythia 8) and fragmentation via Pythia 8

- C++ code, git version control, public on github

SMASH: initialization

- “collider” - elementary or AA reactions, $E_{beam} \gtrsim 0.5 A \text{ GeV}$
- “box” - infinite matter simulations
 - detailed balance tests, computing transport coefficients, thermodynamics of hadron gas
[Rose et al., PRC 97 \(2018\) no.5, 055204](#)
- “sphere” - expanding system
 - testing collision term via comparison to analytical solution of Boltzmann equation,
[Tindall et al., Phys.Lett. B770 \(2017\) 532-538](#)
- “list” - hadronic afterburner after hydrodynamics

SMASH: degrees of freedom

N	Δ	Λ	Σ	Ξ	Ω	Unflavored				Strange
N ₉₃₈	Δ_{1232}	Λ_{1116}	Σ_{1189}	Ξ_{1321}	Ω_{1672}	π_{138}	f_{0980}	f_{21275}	π_{21670}	K_{494}
N ₁₄₄₀	Δ_{1620}	Λ_{1405}	Σ_{1385}	Ξ_{1530}	Ω_{2250}	π_{1300}	f_{01370}	$f_{2'1525}$		K^*_{892}
N ₁₅₂₀	Δ_{1700}	Λ_{1520}	Σ_{1660}	Ξ_{1690}		π_{1800}	f_{01500}	f_{21950}	ρ_{31690}	K_{11270}
N ₁₅₃₅	Δ_{1905}	Λ_{1600}	Σ_{1670}	Ξ_{1820}			f_{01710}	f_{22010}		K_{11400}
N ₁₆₅₀	Δ_{1910}	Λ_{1670}	Σ_{1750}	Ξ_{1950}		η_{548}		f_{22300}	ϕ_{31850}	K^*_{1410}
N ₁₆₇₅	Δ_{1920}	Λ_{1690}	Σ_{1775}	Ξ_{2030}		η'_{958}	a_{0980}	f_{22340}		K_{0^*1430}
N ₁₆₈₀	Δ_{1930}	Λ_{1800}	Σ_{1915}			η_{1295}	a_{01450}		a_{42040}	K_{2^*1430}
N ₁₇₀₀	Δ_{1950}	Λ_{1810}	Σ_{1940}			η_{1405}		f_{11285}		K^*_{1680}
N ₁₇₁₀		Λ_{1820}	Σ_{2030}			η_{1475}	ϕ_{1019}	f_{11420}	f_{42050}	K_{21770}
N ₁₇₂₀		Λ_{1830}	Σ_{2250}				ϕ_{1680}			K_{3^*1780}
N ₁₈₇₅		Λ_{1890}				σ_{800}		a_{21320}		K_{21820}
N ₁₉₀₀		Λ_{2100}					h_{11170}			K_{4^*2045}
N ₁₉₉₀		Λ_{2110}				ρ_{776}		π_{11400}		
N ₂₀₈₀		Λ_{2350}				ρ_{1450}	b_{11235}	π_{11600}		
N ₂₁₉₀						ρ_{1700}				
N ₂₂₂₀							a_{11260}	η_{21645}		
N ₂₂₅₀						ω_{783}				
						ω_{1420}		ω_{31670}		
						ω_{1650}				

- Isospin symmetry
- Perturbative treatment of non-hadronic particles (photons, dileptons)

Hadrons and decay modes configurable via human-readable files

Interactions in SMASH

- Resonance formation and decay

Ex. $\pi\pi \rightarrow \rho \rightarrow \pi\pi$, quasi-inelastic scattering
 $\pi\pi \rightarrow f_2 \rightarrow \rho\rho \rightarrow \pi\pi\pi\pi$

- (In)elastic $2 \rightarrow 2$ scattering

parametrized cross-sections $\sigma(\sqrt{s}, t)$ or
 isospin-dependent matrix elements $|M|^2(\sqrt{s}, I)$

- String formation/fragmentation

$2 \rightarrow n$ processes

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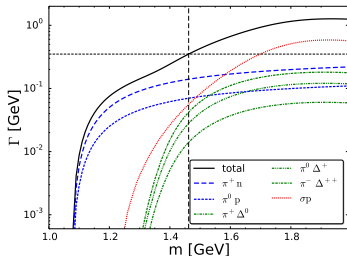
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$2 \rightarrow n$ processes

$N^*(1440)^+$



For every resonance:

- Breit-Wigner spectral function

$$\mathcal{A}(m) = \frac{2\mathcal{N}}{\pi} \frac{m^2 \Gamma(m)}{(m^2 - M_0^2)^2 + m^2 \Gamma(m)^2}$$

- Mass dependent partial widths $\Gamma_i(m)$

Manley formalism for off-shell width [Manley and Saleski, Phys. Rev. D 45, 4002 \(1992\)](#)

$$\text{Total width } \Gamma(m) = \sum_i \Gamma_i(m)$$

Interactions in SMASH

- Resonance formation and decay

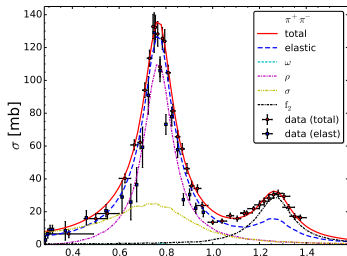
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- $2 \rightarrow 1$ cross-sections from detailed balance relations

Interactions in SMASH

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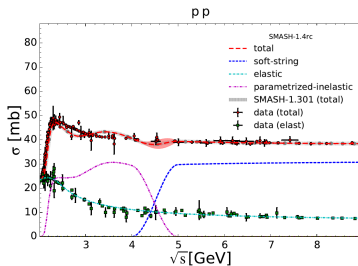
- String formation/fragmentation

$2 \rightarrow n$ processes

- $NN \rightarrow NN^*$, $NN \rightarrow N\Delta^*$, $NN \rightarrow \Delta\Delta$, $NN \rightarrow \Delta N^*$,
 $NN \rightarrow \Delta\Delta^*$

angular dependencies of $NN \rightarrow XX$ cross-sections implemented

- Strangeness exchange $KN \rightarrow K\Delta$, $KN \rightarrow \Lambda\pi$, $KN \rightarrow \Sigma\pi$



Interactions in SMASH

- Resonance formation and decay

Ex. $\pi\pi \rightarrow \rho \rightarrow \pi\pi$, quasi-inelastic scattering
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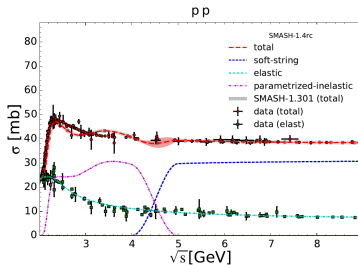
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Interactions in SMASH

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- **String formation/fragmentation**

$2 \rightarrow n$ processes

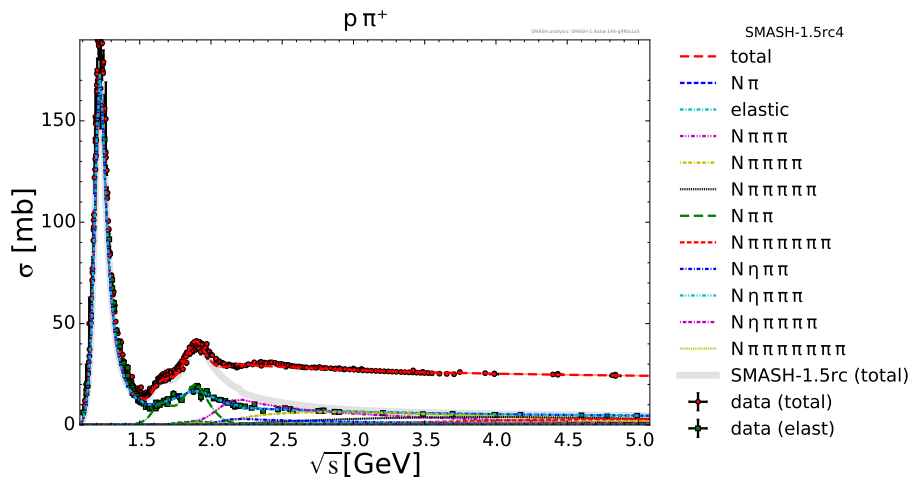
string model parameters
 tuned to NA61 pp

Mohs et al., J.Phys.G 47 (2020) 6, 065101

- String (soft or hard) fragmentation: always via Pythia 8
- Hard scattering and string formation: Pythia
- Soft string formation: SMASH
 - single/double diffractive
 - $B\bar{B}$ annihilation
 - non-diffractive

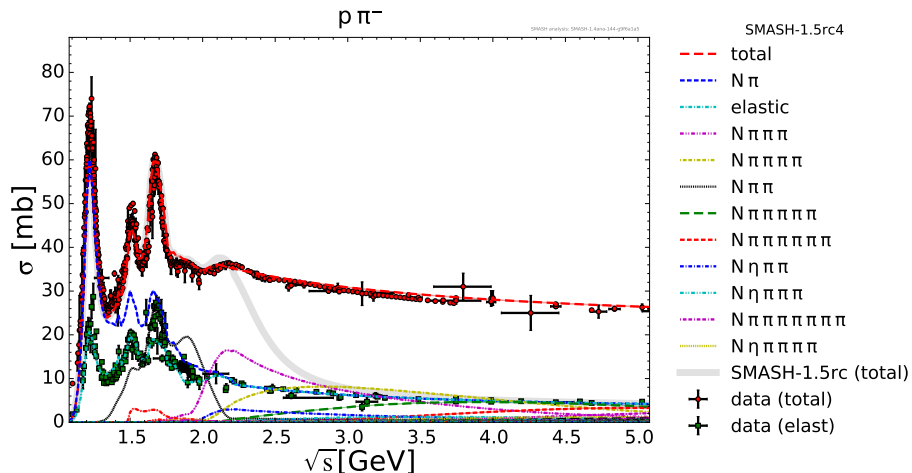
Cross-sections

Resonances at lower energies, string models at $\sqrt{s} \gtrsim 3 - 4$ GeV



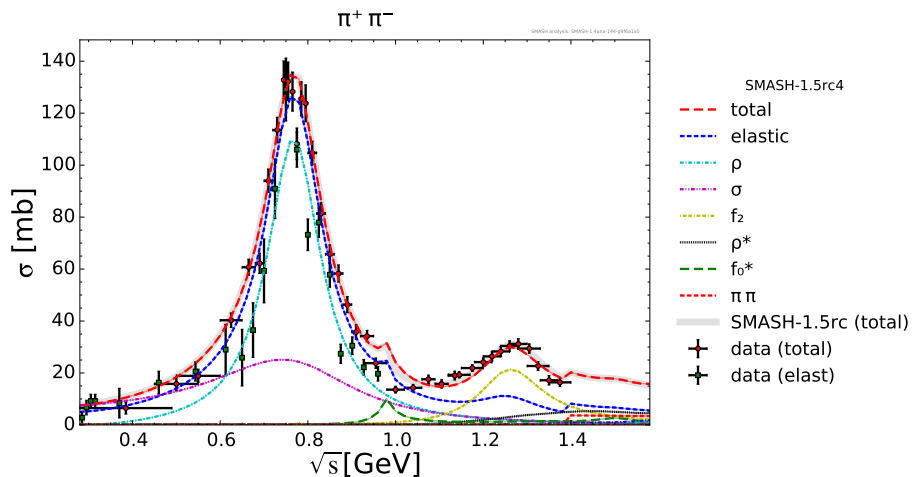
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Cross-sections

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SMASH: analysis suite

SMASH analysis suite

Quick test

Did you learn anything from this lecture? Press yes/no.
If yes then write 1-2 random things you learned in the chat.