

Subtractions for SCET

in collaboration with:

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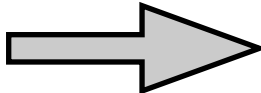
U. Washington

Particle, Field, and String Theory Group

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Why subtractions?

- future in SCET: jet algorithms
- jet functions and soft functions with algs are hard
- only handful of (relatively simple!) examples:
 - N-jets in e^+e^- with the shape of N measured Ellis, Hornig, Lee, Walsh, Vermilion
 - N-jet threshold with $N > 0$ (or 1) Bauer, Dunn, Hornig
 - now, N-jettiness (see Teppo's talk) Tackmann, Stewart, Waalewijn, Jouttenus
- alg dependence on measured jets fncs power suppressed
 -  focus on soft functions

Original Subtractions

- Add & subtract to get finite Real and Virtual:

Ellis, Kunzst, Soper
Catani, Seymour

$$\sigma^{NLO} = \int_{m+1} \left[\left(d\sigma^R \right)_{\epsilon=0} - \left(d\sigma^A \right)_{\epsilon=0} \right] + \int_m \left[d\sigma^V + \int_1 d\sigma^A \right]_{\epsilon=0}$$

- subtractions are universal and analytically calculable, same singularities as Virtual and Real (point-by-point!)
- difference here Soft/Jet Functions: depend on observable/algorithm
- however, will see that all obs/algs belong in one of 2 universality classes

Outline

- UV study: all alg/obs combos 2 types of subtractions
- Angularities as basis for all UV limits for φ -symm obs
- Examples:
 - cross-check theta'-phi from theta-phi jets
 - jet shapes for eta-phi jets @ pp
 - 1-jettiness
 - others...
- Compare/contrast with related work (see Teppo's talk on Tues)

Soft Functions:

$$\mathcal{S}^{(1)}(\mathcal{A}, \mathcal{M}; \{\sigma_m\}) = \sum_{\langle i, j \rangle} \int \frac{d^d k}{(2\pi)^d} N_{ij}(k) \sum_{k=0}^N \Theta_{\mathcal{A}}^k \Delta_{\mathcal{M}}^k$$

- integrand:

$$N_{ij} = -g^2 \mu^{2\epsilon} \mathbf{T}_i \cdot \mathbf{T}_j \frac{n_i \cdot n_j}{n_i \cdot k n_j \cdot k} 2\pi \delta(k^2) \theta(k^0)$$

- measurement:

$$\Delta_{\mathcal{M}}^k(k) = \delta(\sigma_k - \sigma_k(k)) \prod_{l \neq k} \delta(\sigma_l)$$

- algorithm: $\Theta_{\mathcal{A}}^k$

- e.g., $\Theta_{\text{cone}}^k(k) = \theta\left(\frac{n_k \cdot k}{\bar{n}_k \cdot k} < \tan^2 \frac{R}{2}\right)$

$$\Theta_{\text{cone}}^0(k) = \left(1 - \sum_{k=1}^N \Theta_{\text{cone}}^k(k)\right) \theta(k^0 < \Lambda)$$

Divergences in Soft Functions

$$N_{ij} = -g^2 \mu^{2\epsilon} \mathbf{T}_i \cdot \mathbf{T}_j \frac{n_i \cdot n_j}{n_i \cdot k n_j \cdot k} 2\pi \delta(k^2) \theta(k^0)$$

- IR: $k^0 \rightarrow 0$
 - Collinear: $n_i \cdot k \rightarrow 0$
 - UV: $k^0 \rightarrow \infty$
- } $1/\epsilon^2$
- } $1/\epsilon^2$

Divergences in *Differences of Soft Functions*

- take moments for easier comparison:

$$\langle \sigma_k^n \rangle = \prod_l \int_0^1 d\sigma_l \sigma_k^n \mathcal{S}^{(1)}(\mathcal{A}, \mathcal{M}; \{\sigma_m\}) \equiv \int \frac{d^d k}{(2\pi)^d} I_k^n(k)$$

- IR & Coll. trivial due to IR safety:

- $n > 0$: $\lim_{k^0 \rightarrow 0} \sigma_i(k) = 0$ \Rightarrow $\lim_{k^0 \rightarrow 0} (I_k^n(k) - \tilde{I}_k^n(k)) = 0$

- $n > 0$: $\lim_{n_i \cdot k \rightarrow 0} \sigma_i(k) = 0$ \Rightarrow $\lim_{n_i \cdot k \rightarrow 0} (I_k^n(k) - \tilde{I}_k^n(k)) = 0$

- $n = 0$: $\lim_{k^0 \rightarrow 0} (I^0(k) - \tilde{I}^0(k)) \propto \sum_{k=0}^N [\Theta_{\mathcal{A}}^k(k) - \Theta_{\tilde{\mathcal{A}}}^p(k)] = 0$

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Divergences in *Differences of Soft Functions*

- take moments for easier comparison:

$$\langle \sigma_k^n \rangle = \prod_l \int_0^1 d\sigma_l \sigma_k^n \mathcal{S}^{(1)}(\mathcal{A}, \mathcal{M}; \{\sigma_m\}) \equiv \int \frac{d^d k}{(2\pi)^d} I_k^n(k)$$

- UV safety:

- $n > 0$:

$$\lim_{k^0 \rightarrow \infty} (I^0(k) - \tilde{I}^0(k)) \propto \sum_{k \in \text{meas}} \left[\Theta_{\mathcal{A}}^k(k) \theta(\sigma_k(k) < 1) - \Theta_{\tilde{\mathcal{A}}}^k(k) \theta(\tilde{\sigma}_k(k) < 1) \right] + \sum_{k \notin \text{meas}} \left[\Theta_{\mathcal{A}}^k(k) - \Theta_{\tilde{\mathcal{A}}}^k(k) \right]$$

what matters depends on whether alg. or obs. “wins”

only alg. matters

- $n = 0$:

$$\lim_{k^0 \rightarrow \infty} (I_k^n(k) - \tilde{I}_k^n(k)) \propto \sigma_k^n(k) \Theta_{\mathcal{A}}^l(k) \theta(\sigma_k(k) < 1) - \tilde{\sigma}_k^n(k) \Theta_{\tilde{\mathcal{A}}}^k(k) \theta(\tilde{\sigma}_k(k) < 1)$$

⇒ 2 types of soft functions

Type I vs. Type II

- Type I: $\lim_{k^0 \rightarrow \infty} \Theta_{\mathcal{A}}^l(k) \theta(\sigma_l(k) < 1) = \lim_{k^0 \rightarrow \infty} \theta(\sigma_l(k) < 1)$ (obs. “wins”)
 - incl. kt/cone with any angularity $a < 2$ (or obs. with same UV)
 - excl. kt with $1 < a < 2$
 - N-jettiness
 - N-jet threshold at e+e- and pp colliders
- Type II: $\lim_{k^0 \rightarrow \infty} \Theta_{\mathcal{A}}^l(k) \theta(\sigma_l(k) < 1) = \lim_{k^0 \rightarrow \infty} \Theta_{\mathcal{A}}^l(k)$ (alg. “wins”)
 - “unmeas.” jets (no shape is probed)
 - excl. kt with $a < 1$

Angularities as Basis of Obs. in UV

- def'n (Q arbitrary normalization):

Berger, Kucs, Sterman

$$\tau_a = \frac{1}{Q} \sum_i E_i (\sin \theta_i)^a (1 - |\cos \theta_i|)^{1-a} = \frac{1}{Q} \sum_i |\mathbf{p}_i^T| e^{-|\eta_i|(1-a)}$$

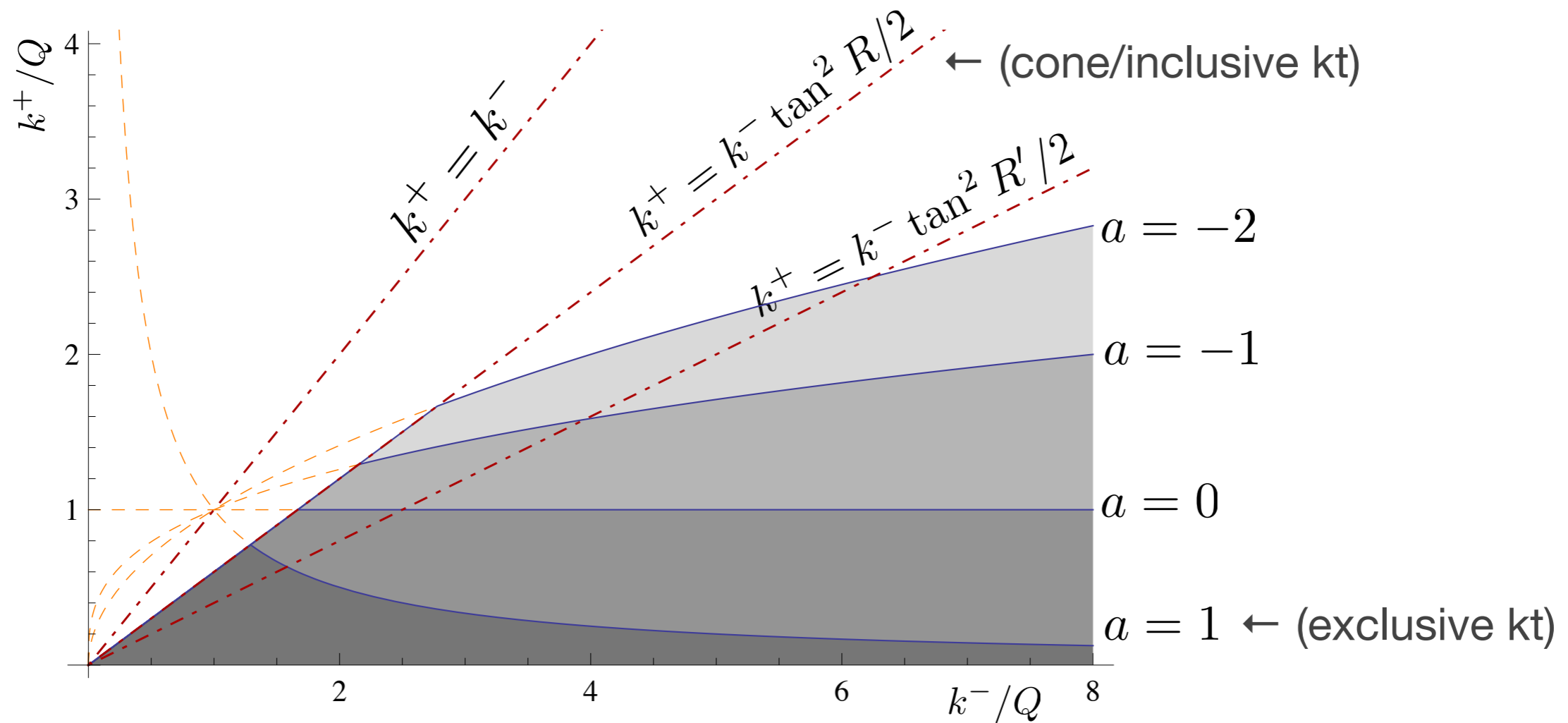
- for fixed τ , as $E \rightarrow \infty$ (so $\theta \rightarrow 0$):

$$\tau_a \rightarrow E_i \theta^{2-a}$$

- this is the most general, IR safe (for $a < 2$) result in $\theta \rightarrow 0$ limit
- results exist for jet shapes for all $a < 1$ (will use as “subtractions”)...

Ellis, Hornig, Lee, Walsh, Vermilion

Demonstration of Type I vs. II in k^+k^- -plane



- Note: any measurement (w/ $a < 1$) does not change divergences of Type II

Final Result

$$\mathcal{S}^{(1)}(\mathcal{A}, \mathcal{M}; \{\sigma_m\}) = \tilde{\mathcal{S}}^{(1)}(\tilde{\mathcal{A}}, \tilde{\mathcal{M}}; \{\tilde{\sigma}_m \rightarrow \sigma_m\}) + \mathcal{D}^0 \prod_i \delta(\sigma_i) + \sum_i \mathcal{D}_i^1 \prod_{j \neq i} \delta(\sigma_j) \left(\frac{1}{\sigma_i}\right)_+$$

$$\mathcal{D}^0 = \langle \sigma^0 \rangle - \langle \tilde{\sigma}^0 \rangle \quad \mathcal{D}_k^n = \langle \sigma_k^n \rangle - \langle \tilde{\sigma}_k^n \rangle$$

(0th moment w.r.t. all obs) (nth moment w.r.t. kth obs)

- we use jet shape calc. as subtraction (variation needed for threshold)
- moments only involve $\int d\Omega$:

$$\mathcal{D}^0 = -\frac{\alpha_s}{\pi} \sum_{\langle i,j \rangle} \mathbf{T}_i \cdot \mathbf{T}_j \int \frac{d\Omega}{4\pi} \frac{n_i \cdot n_j}{(1 - \hat{k} \cdot \hat{n}_i)(1 - \hat{k} \cdot \hat{n}_j)}$$

$$\times \left[\underbrace{\sum_{k=1}^N (\Theta_{\mathcal{A}}^k - \Theta_{\tilde{\mathcal{A}}}^k) \log \frac{\tilde{k}_{\max}^k(\Omega)}{\Lambda}}_{\text{diff. from alg.}} + \underbrace{\left(1 - \sum_{k=1}^N \Theta_{\mathcal{A}}^k\right) \log \frac{E_{\text{cut}}(\Omega)}{\Lambda}}_{\text{diff. from veto}} + \underbrace{\sum_{k=1}^N \Theta_{\mathcal{A}}^k \log \frac{k_{\max}^k(\Omega)}{\tilde{k}_{\max}^k(\Omega)}}_{\text{diff. from obs.}} \right]$$

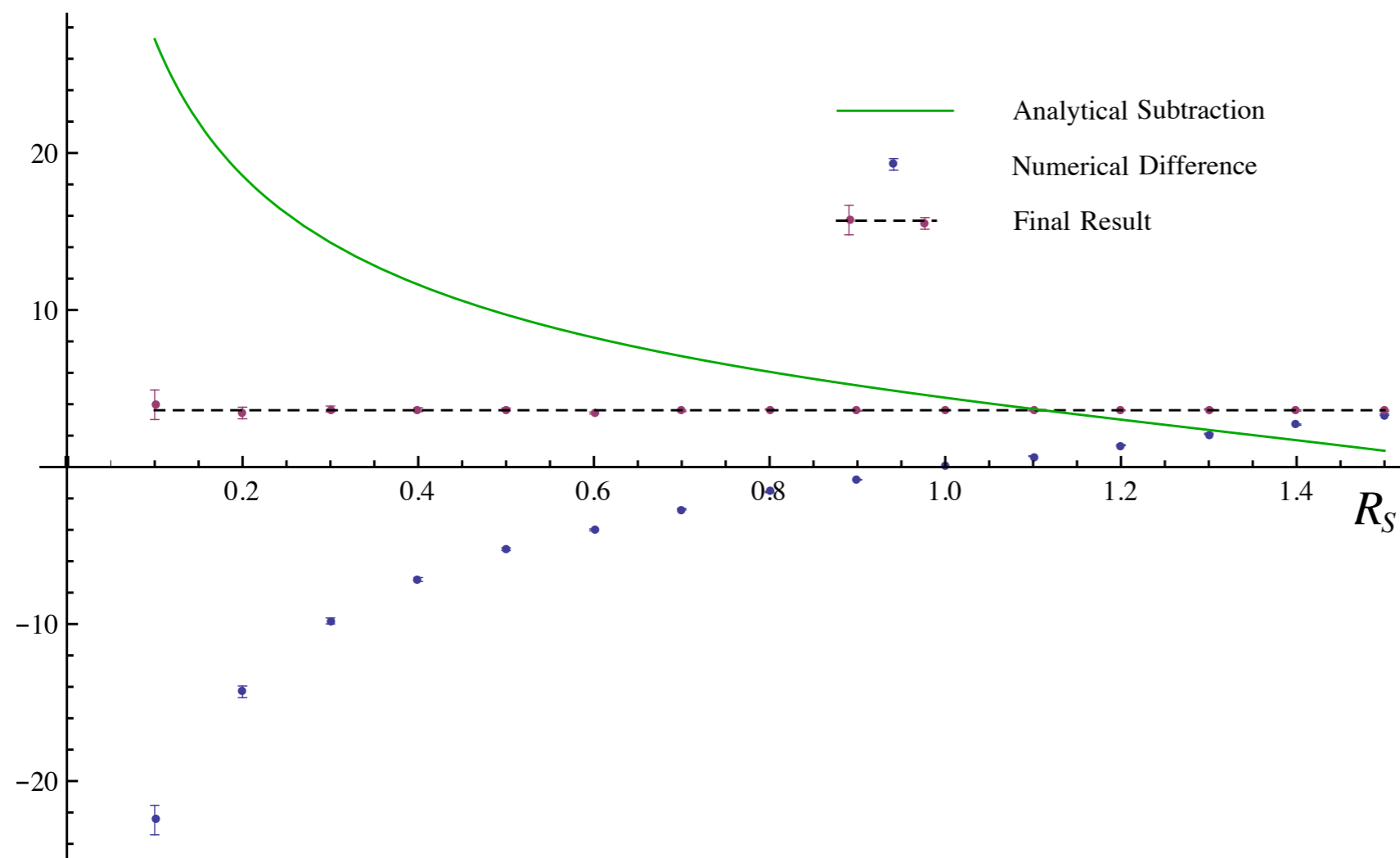
$$k_{\max}^k(\Omega) = k^0 / \sigma_k(k)$$

$$\tilde{k}_{\max}^k(\Omega) = Q_i (1 + \hat{k} \cdot n_k)^{-a/2} (1 - \hat{k} \cdot n_k)^{-1+a/2}$$

$$E_{\text{cut}}(\Omega) = \begin{cases} E_{\text{cut}} \\ p_{\text{cut}}^T / \sin \theta \\ \vdots \end{cases}$$

$\theta\varphi$ jets of $R = R_t$ from $R = R_s$ jets (cross-check)

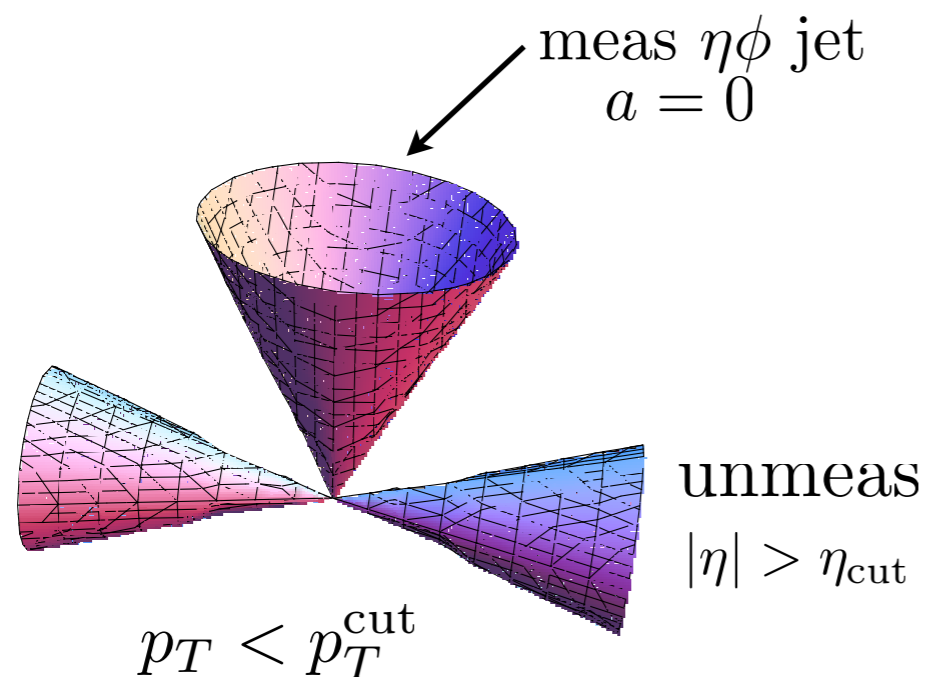
- Changing subtraction R_s in mercedes benz (target R_t fixed):



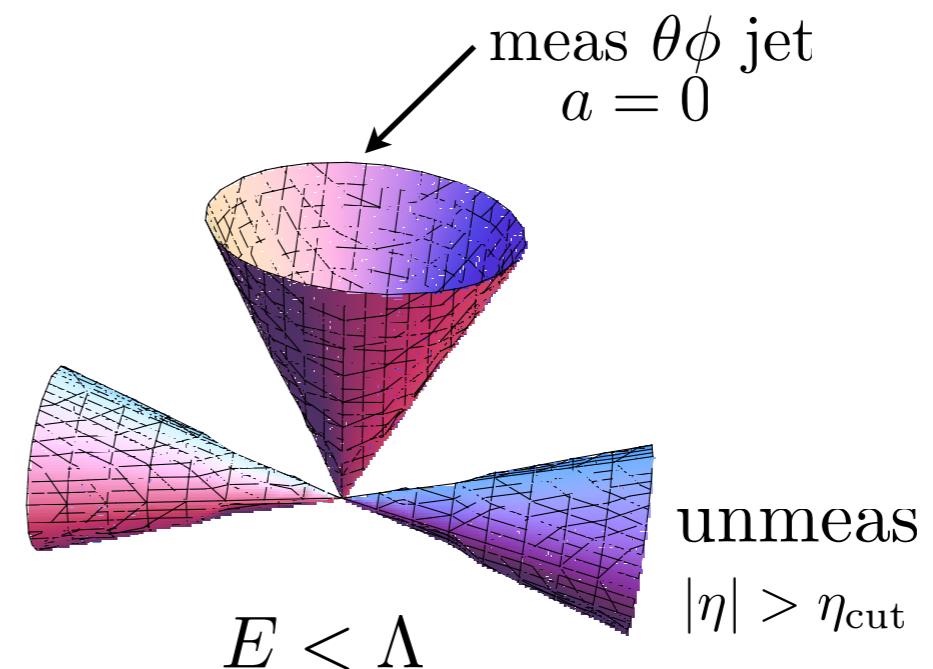
jet shapes of $\eta\phi$ jets (w/ $p_{T\text{cut}}$ & η_{cut})

- one jet w/ angularity measured, two beams with $|\eta| > \eta_{\text{cut}}$ cut-out
 - needs Type I for meas. jet, Type II for beams

Target:



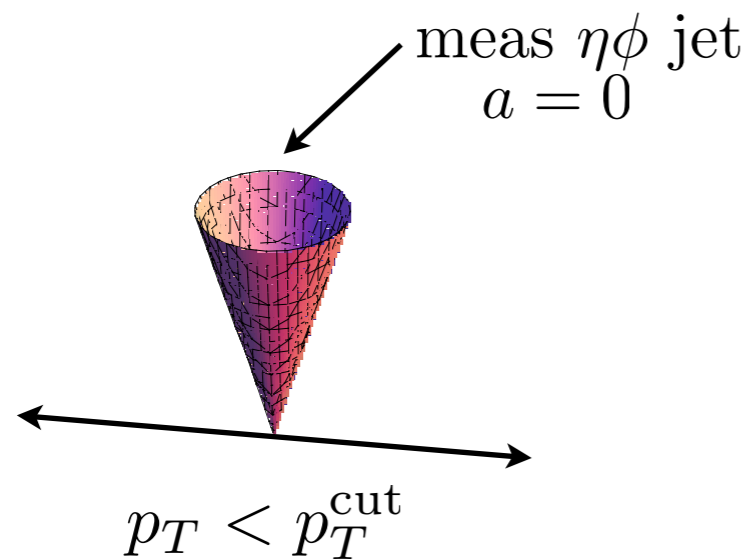
Subtraction:



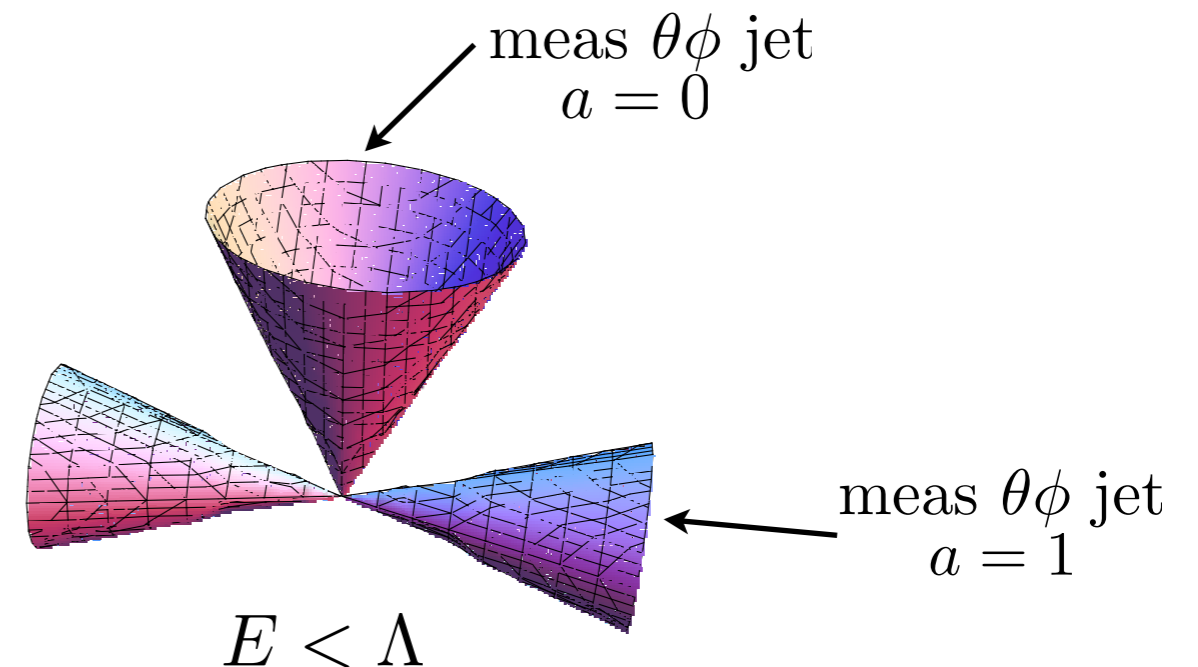
jet shapes of $\eta\phi$ jets (w/ $p_{T\text{cut}}$ & η_{cut})

- note: could also do IF we had $a=1$ (no logs of η_{cut})

Target:



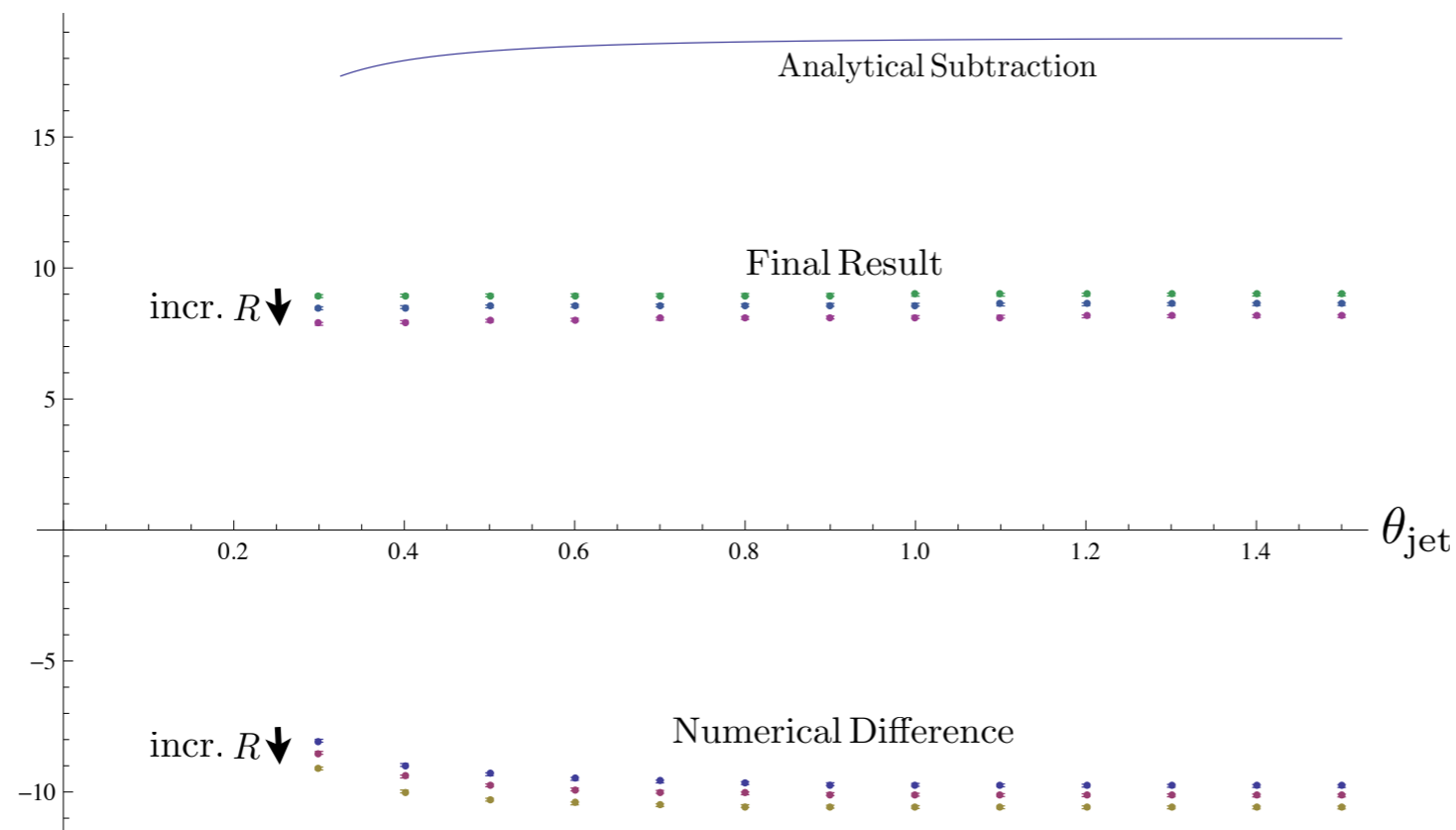
Subtraction:



- question: ok that $a=1$ weights near beam more than $a=0$ (θ^1 vs θ^2)

jet shapes of $\eta\phi$ jets (w/ pT_{cut} & η_{cut})

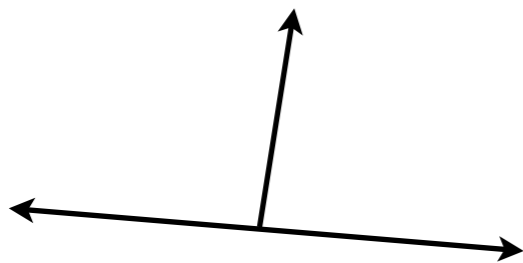
- one jet w/ angularity measured, two beams with $|\eta| > \eta_{\text{cut}}$ cut-out
 - needs Type I for meas. jet, Type II for beams
- beam-beam color-connection contribution as a function of R_{jet} :



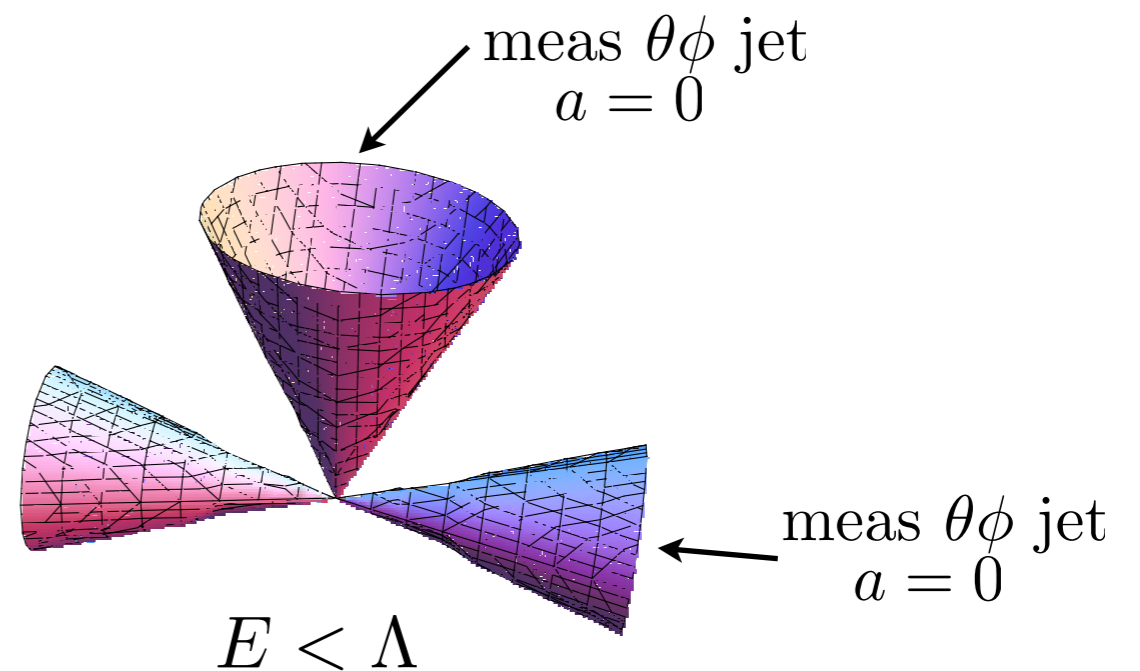
1-Jettiness from $\theta\phi$ jets w/ E_{cut}

- 1-jettiness has measurement everywhere:

Target:

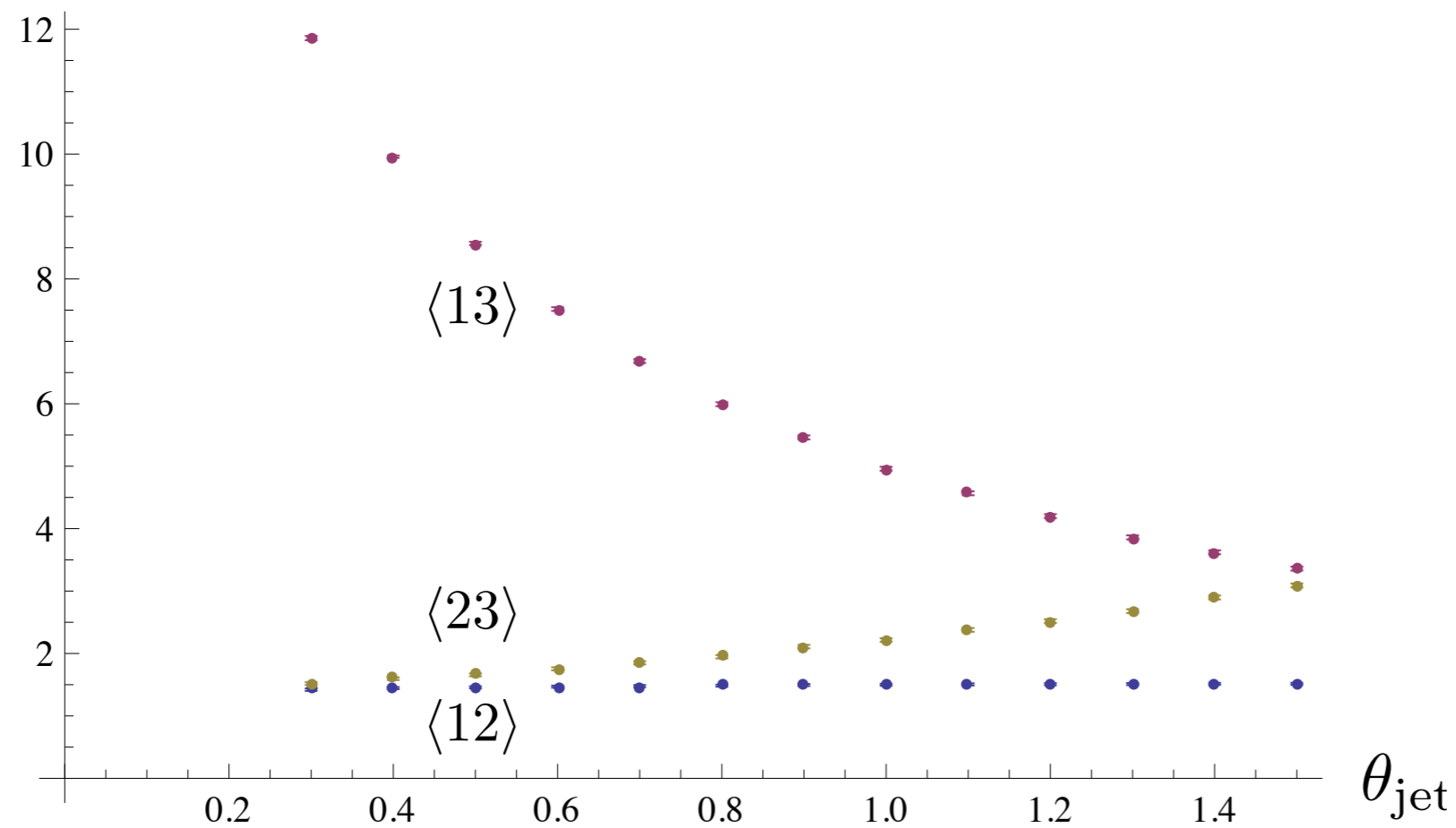


Subtraction:



1-Jettiness from $\theta\varphi$ jets w/ E_{cut}

- 1,2 = beams; 3 = jet



More examples

- jet shapes that only asymptote to an angularity
- N-jet threshold at pp (subtraction = N-jet threshold at e+e-)
- pT_{cut} everywhere (no large logs of η_{cut})
- azimuthally asymmetric observables (e.g., planarity)

Comparison to the “Hemisphere Decomposition” (MIT) (see Teppo’s talk)

$$F(\{k_i\}, p) = F_{ij,\text{hemi}}(\{k_i\}, p) + F_{ji,\text{hemi}}(\{k_i\}, p) + \sum_{m \neq i} F_{ij,m}(\{k_i\}, p) + \sum_{m \neq j} F_{ji,m}(\{k_i\}, p)$$

Analytical, Divergent:

$$F_{ij,\text{hemi}}(\{k_i\}, p) = \theta(p^j - p^i) \delta[k_i - f_i(p)] \prod_{l \neq i} \delta(k_l)$$

No R-Dependence



Numerical, Finite:

$$F_{ij,m}(\{k_i\}, p) = \theta(p^j - p^i) \Theta_m(p) \prod_{l \neq i,m} \delta(k_l) \left\{ \delta(k_i) \delta[k_m - f_m(p)] - \delta[k_i - f_i(p)] \delta(k_m) \right\}$$

- $F_{ij,\text{hemi}}$ easier than $S_{ij}^{\text{meas}}(\tau^k)$, but not known for a $\neq 0$
- Scope: only works for Type I (UPDATE: maybe Type II ??)

Comparison to the “Hemisphere Decomposition” (MIT) (see Teppo’s talk)

- Numerics: unbounded y integral
 \Rightarrow analytical evaluation needed (just for N-jettiness)

$$I_0(\alpha, \beta, \{\alpha_l, \beta_l, \phi_l\}) = \frac{1}{\pi} \int_{-\pi}^{\pi} d\phi \int \frac{dy}{y} \theta(y - \sqrt{\beta/\alpha}) \theta(1/\alpha - 1 - y^2 + 2y \cos \phi) \\
\times \prod_l \theta[\alpha_l - 1 + (\beta_l - 1)y^2 - 2y[\sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l) - \cos \phi]],$$

$$I_1(\alpha, \beta, \{\alpha_l, \beta_l, \phi_l\}) = \frac{1}{\pi} \int_{-\pi}^{\pi} d\phi \int \frac{dy}{y} \ln(1 + y^2 - 2y \cos \phi) \theta(y - \sqrt{\beta/\alpha}) \theta(1/\alpha - 1 - y^2 + 2y \cos \phi) \\
\times \prod_l \theta[\alpha_l - 1 + (\beta_l - 1)y^2 - 2y[\sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l) - \cos \phi]].$$

Unbounded dy/y

Quadratic Eq. in y:

$$y_{\pm}(\phi, \alpha_l, \beta_l, \phi_l) = \frac{1}{1 - \beta_l} \left\{ \cos \phi - \sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l) \pm \sqrt{[\cos \phi - \sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l)]^2 - (1 - \alpha_l)(1 - \beta_l)} \right\}$$

1. $\beta_l < 1$:

(a) $\alpha_l \geq 1$: $y \leq y_+(\phi, \alpha_l, \beta_l, \phi_l)$

(b) $\alpha_l < 1$: $y_-(\phi, \alpha_l, \beta_l, \phi_l) \leq y \leq y_+(\phi, \alpha_l, \beta_l, \phi_l)$
 $\cos \phi - \sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l) \geq \sqrt{(1 - \alpha_l)(1 - \beta_l)}$

2. $\beta_l > 1$:

(a) $\alpha_l \leq 1$: $y \geq y_+(\phi, \alpha_l, \beta_l, \phi_l)$

(b) $\alpha_l > 1$: $y \leq y_-(\phi, \alpha_l, \beta_l, \phi_l)$ or $y \geq y_+(\phi, \alpha_l, \beta_l, \phi_l)$
 $\cos \phi - \sqrt{\alpha_l \beta_l} \cos(\phi + \phi_l) \geq \sqrt{(1 - \alpha_l)(1 - \beta_l)}$

3. $\beta_l = 1$:

(a) $\alpha_l \leq 1$: $y \geq \frac{1 - \alpha_l}{2 \cos \phi - 2\sqrt{\alpha_l} \cos(\phi + \phi_l)}$

$\cos \phi \geq \sqrt{\alpha_l} \cos(\phi + \phi_l)$

(b) $\alpha_l > 1$: $y \leq \frac{\alpha_l - 1}{2\sqrt{\alpha_l} \cos(\phi + \phi_l) - 2 \cos \phi}$

$\cos \phi \leq \sqrt{\alpha_l} \cos(\phi + \phi_l)$

Conclusions

- subtractions can easily be used to get LHC-ready soft functions
- ingredients for all azimuthally symmetric obs ($a < 1$) with Type I algorithms
- need for future:
 - $a \geq 1$
 - $S_{ij}(\tau^k)$, S_{ij}^k for azimuthally-asymmetric classes
 - 2-loop anom dims of jet/soft fncs for $a \neq 0$ for NNLL
 - understanding resummation for Type II