Institute for Research in Fundamental Sciences

FCC-ee PHYSICS MEETING

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Searching For Lepton Flavor Violating Interactions At Future Electron-positron Colliders

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30 Aug. 2021

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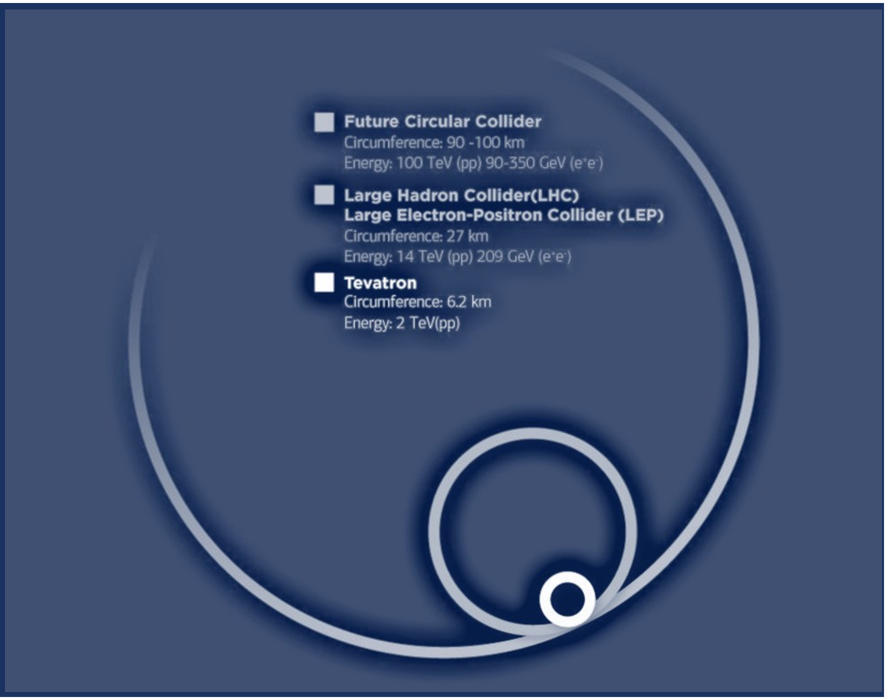
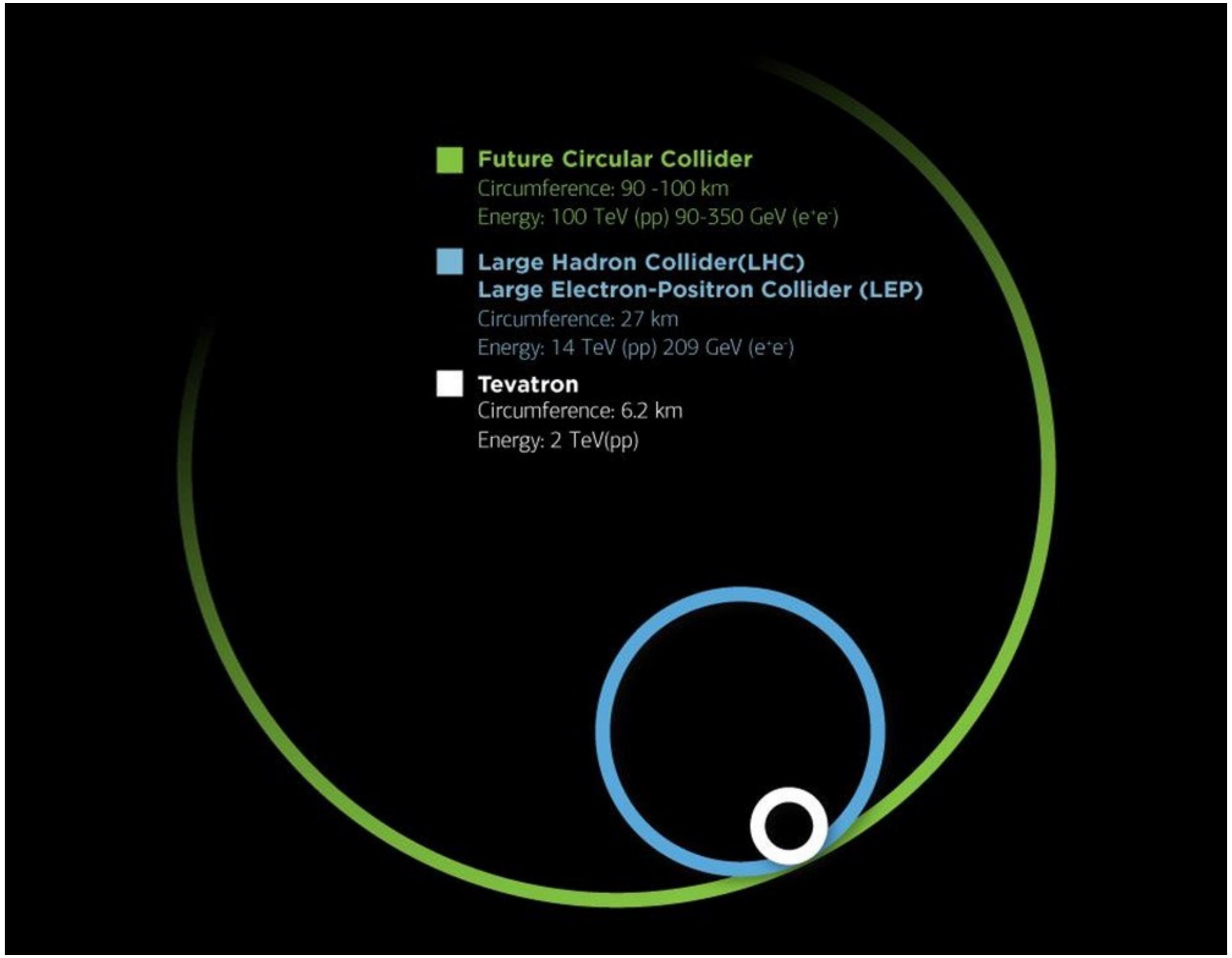


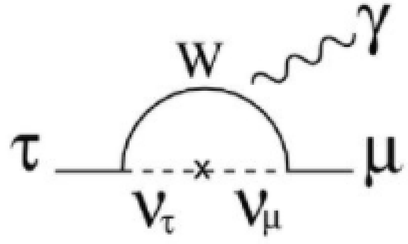
OUTLINE

- Motivation & introduction
- Theoretical framework
- Data Simulation
- Analysis strategy
- Results & discussion
- Conclusion



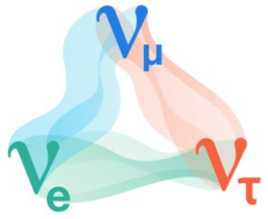
MOTIVATION & INTRODUCTION

- **In the Standard Model (SM) neutrinos are massless.**
- **Lepton Flavor Violating (LFV) interactions are forbidden in the SM.**
- **Neutrino oscillations have been observed → Neutrinos are massive.**
- **This leads to LFV. But …**



$$\boxed{BR(\tau^- \to \ell^+\ell^-\ell'^-) \lesssim 10^{-54}}$$

[arXiv:1912.09862]



- **An increase of several orders of magnitude is predicted in some SM extensions.** [arXiv:0406039]

Any detection of LFV signal → Clear evidence for BSM

MOTIVATION & INTRODUCTION

- **So far, no cLFV has been observed and there are several strong constraints from various experiments.**

$$\mathcal{B}(\tau^- \to e^- e^+ e^-) \leq 2.9 \times 10^{-8} \text{ (BaBar)}$$

$$\mathcal{B}(\tau^- \to e^- e^+ e^-) \leq 2.7 \times 10^{-8} \text{ (Belle)}$$

- **The Belle II prospect at 90% CL with 50 ab^{-1} :**

$$\mathcal{B}(\tau^- \to e^- e^+ e^-) \lesssim 10^{-10}$$

arXiv:1808.10567

Future lepton colliders are expected to provide an extraordinary place in flavor physics.

- **What have been studied?**
  - **The LFV through Z and H**
  - **H and scalar LFV decays**

If the new degrees of freedom contributing to LFV are heavy enough, the LFV couplings could be reasonably parameterized via the effective contact interactions.

MOTIVATION & INTRODUCTION

- **LFV $eee\tau$ contact interactions previous studies:**

The LFV contact operators probed: [arXiv1410.1485]

- **Via $e^-e^+ \rightarrow e^{\pm}\tau^{\mp}$ process at $\sqrt{s}$ = 250, 500, 1000, 3000 GeV.**
- **Considering two main background sources, $\tau\tau$ and $e\tau\upsilon\upsilon$.**

Similar study at $\sqrt{s}$ = 250, 500, 1000 GeV: [arXiv1803.10475]

- **The effects of polarization beams.**
- **Detector response**
- **The main source of backgrounds of $e\tau\upsilon\upsilon$.**

- **In this study:**

Four FCC-ee benchmarks	ISR effect
other main backgrounds	statistical data combination

THEORETICAL FRAMEWORK

Four Fermi contact interactions ($ee e\tau$)

To characterize the new physics effects

In a model-independent framework

- Six chirality conserving four-Fermi operators (ΔL=1):
  - Vector type
  - Scalar type
- LFV operators containing dipole structures: Tightly constrained by radiative LFV decays

Not considered in this work.

THEORETICAL FRAMEWORK

- **The effective Lagrangian and the relevant operators:** [arXiv:9909265]

$$\mathcal{L}_{\text{eff}} \supset \sum_{\alpha,\beta} \sum_{ij} \frac{c^{ij}_{\alpha\beta}}{\Lambda^2} \mathcal{O}^{ij}_{\alpha\beta} \,,$$

$$\begin{aligned}
\mathcal{O}^{S,ij}_{RL} &= (\bar{\ell}_{jL}\ell_{iR})\,(\bar{\ell}_{jL}\ell_{jR})\,, & \mathcal{O}^{S,ij}_{LR} &= (\bar{\ell}_{iR}\ell_{jL})\,(\bar{\ell}_{jR}\ell_{jL})\,, \\
\mathcal{O}^{V,ij}_{RR} &= (\bar{\ell}_{iR}\gamma^{\mu}\ell_{jR})\,(\bar{\ell}_{jR}\gamma_{\mu}\ell_{jR})\,, & \mathcal{O}^{V,ij}_{LL} &= (\bar{\ell}_{iL}\gamma^{\mu}\ell_{jL})\,(\bar{\ell}_{V,jL}\gamma_{\mu}\ell_{jL})\,, \\
\mathcal{O}^{V,ij}_{LR} &= (\bar{\ell}_{iL}\gamma^{\mu}\ell_{jL})\,,(\bar{\ell}_{jR}\gamma_{\mu}\ell_{jR})\,, & \mathcal{O}^{V,ij}_{RL} &= (\bar{\ell}_{iR}\gamma^{\mu}\ell_{jR})\,(\bar{\ell}_{iL}\gamma_{\mu}\ell_{iL})\,,
\end{aligned}$$

- **Λ : the energy scale of new physics**
- **$c_{i,j}$:Wilson couplings between leptons of flavor i and j**
- **$O_{i,j}$: four fermion leptonic operators, invariant under the SM gauge symmetry**

THEORETICAL FRAMEWORK

- **LFV among 1st and 2nd generations: tightly constrained by experimental constraints arising from:**

Nucl. Phys. B 299 (1988)	arXiv:1605.05081	Eur. Phys. J. C 47, 337-346 (2006)
$\mu \to 3e$ at SINDRUM experiment	**Muon transition to $e\gamma$**	**$e-\mu$ conversion**

However,

- **Constraints on LFVs between e and τ, and μ and τ are much looser $\Rightarrow$ eeτ couplings**
- **In addition to the $ee e\tau$ four-Fermi contact interactions**

Other favorite interactions

arXiv: 1602.01698 , 2101.05286 , …

- **Leptons and quarks** $eeqq'$
- **Electrons and Higgs-Z** $eeHZ$

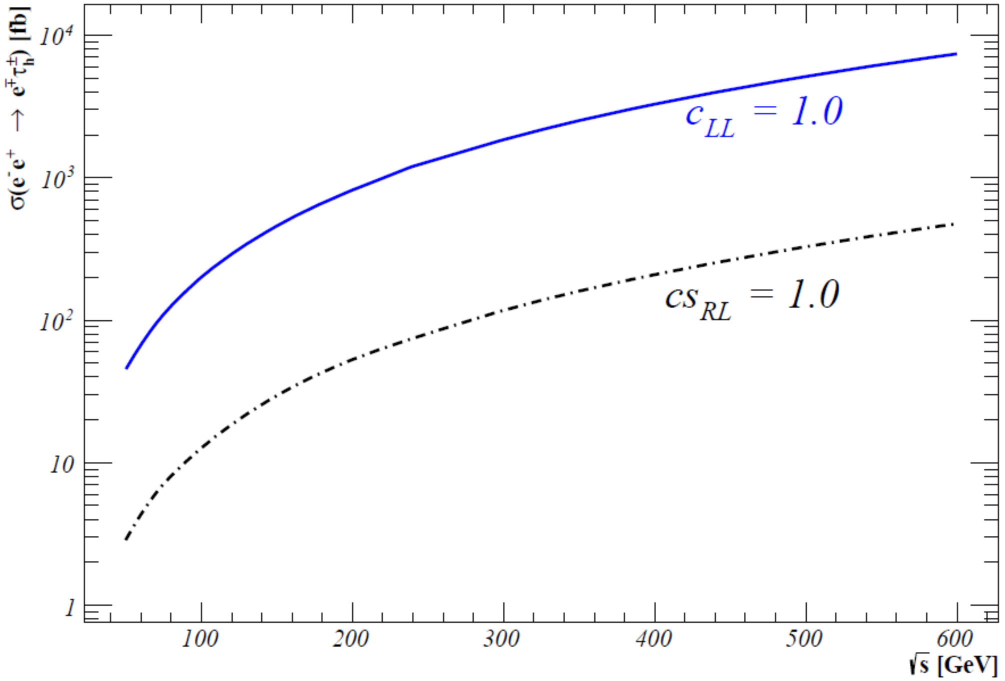
THEORETICAL FRAMEWORK

- **The theoretical cross section of $e^-e^+ \rightarrow e^\pm\tau^\mp$:**

[arXiv: 0611222]

$$\sigma(s) = \frac{s}{96\pi\Lambda^4}\left\{(|c^S_{LR}|^2 + |c^S_{RL}|^2) + 16(|c^V_{LL}|^2 + |c^V_{RR}|^2 + |c^V_{LR}|^2 + |c^V_{RL}|^2)\right\}$$

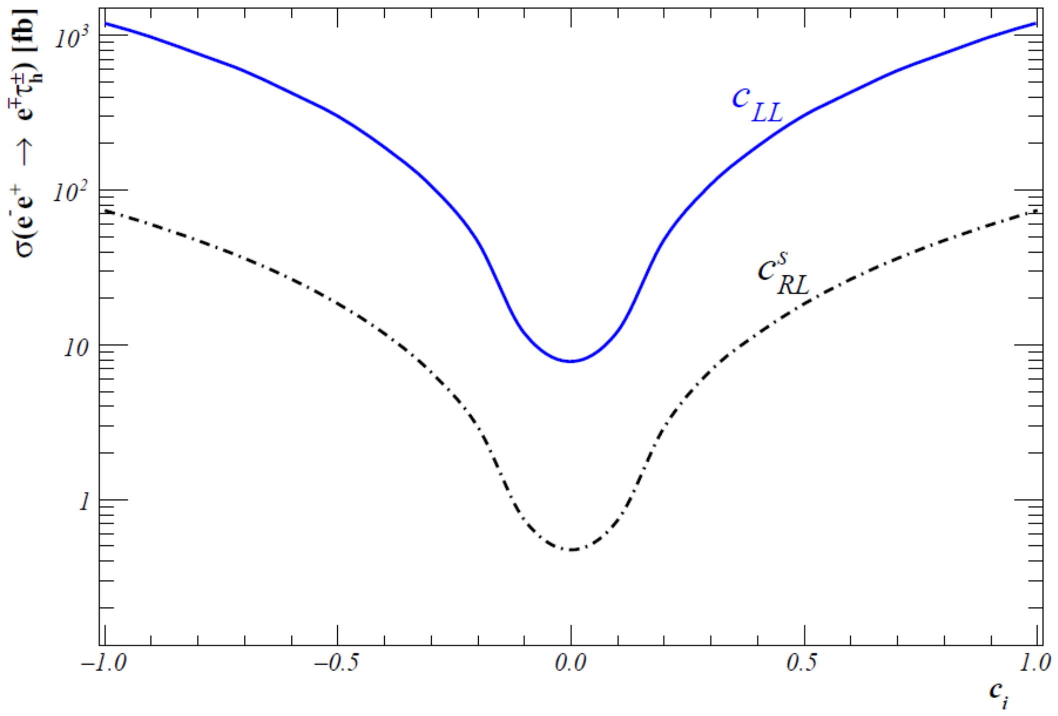
$$\boxed{\sigma(e^-e^+ \rightarrow e\tau) \propto s}$$



- **Vector type operators are larger than the scalar type ones by a factor of 16 .**

DATA SIMULATION

Effective Lagrangian → FeynRule program → Universal FeynRules Output (UFO) → MADGRAPH5_ aMC@NLO → PYTHIA 8 (showering) → DELPHES 3 (ILD-like)



- **Six different signal samples ≡ Six operators**
- **Tau-lepton hadronic decay.**
- **Wilson coefficients $c_{i,j}$ =0.1, with i, j = L,R, and Λ = 1 TeV.**

DATA SIMULATION

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- **Four FCC-ee benchmarks**

C.M. Energy (GeV)	365	240	162.5	157.5
Integrated luminosity (ab^{-1})	1.5	5	5	5

- **Background processes:**

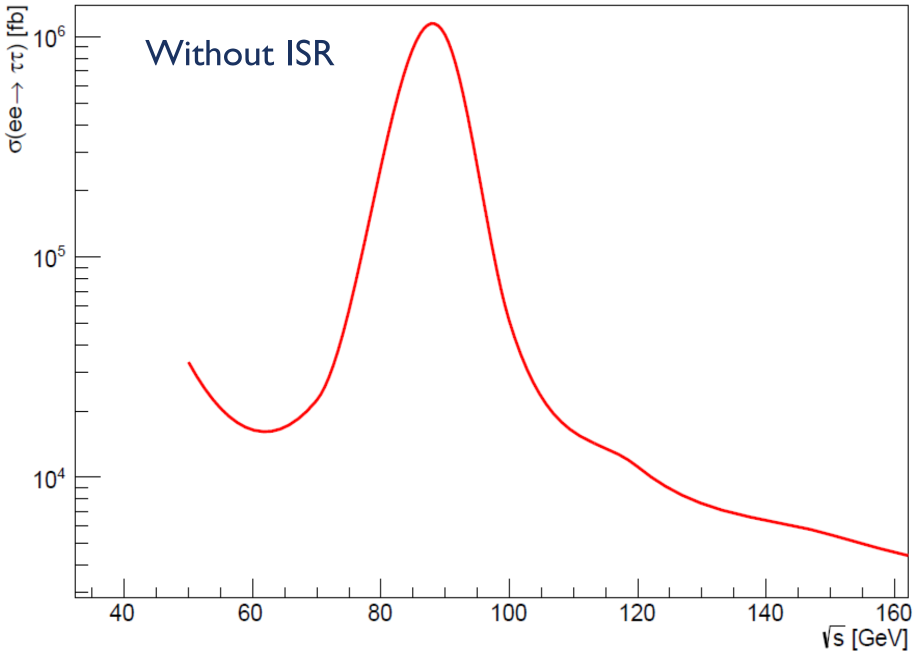
$$
\begin{array}{ll}
\text{(I)} & e^-e^+ \to e^\pm \tau^\mp \nu\bar{\nu}, \\
\text{(II)} & e^-e^+ \to \tau^+\tau^-, \\
\text{(III)} & e^-e^+ \to \ell^\pm\ell^\mp\ell'^\pm\ell'^\mp \; (\ell, \ell' = e, \mu, \tau), \\
\text{(IV)} & e^-e^+ \to \ell^\pm\ell^\mp jj \; (\ell = e, \mu, \tau), \\
\text{(V)} & e^-e^+ \to \ell^\pm \nu jj (\ell = e, \mu, \tau), \\
\text{(VI)} & e^-e^+ \to jj.
\end{array}
$$

DATA SIMULATION

- **ISR effects** are considered using the MGISR plugin (MadGraph5 version: 2.6.6) [arXiv:1705.04486] [arXiv:1804.00125]

$\sqrt{s}$ [GeV]	$\tau\bar{\tau}$(without ISR)	$\tau\bar{\tau}$(with ISR)
157.5	4869.4	11076.5
162.5	4514.9	10275.8
240	1910.5	4196.8
365	804.15	1803.6

ee $\to \tau\tau$ cross section [fb]



ANALYSIS STRATEGY

Event selection:

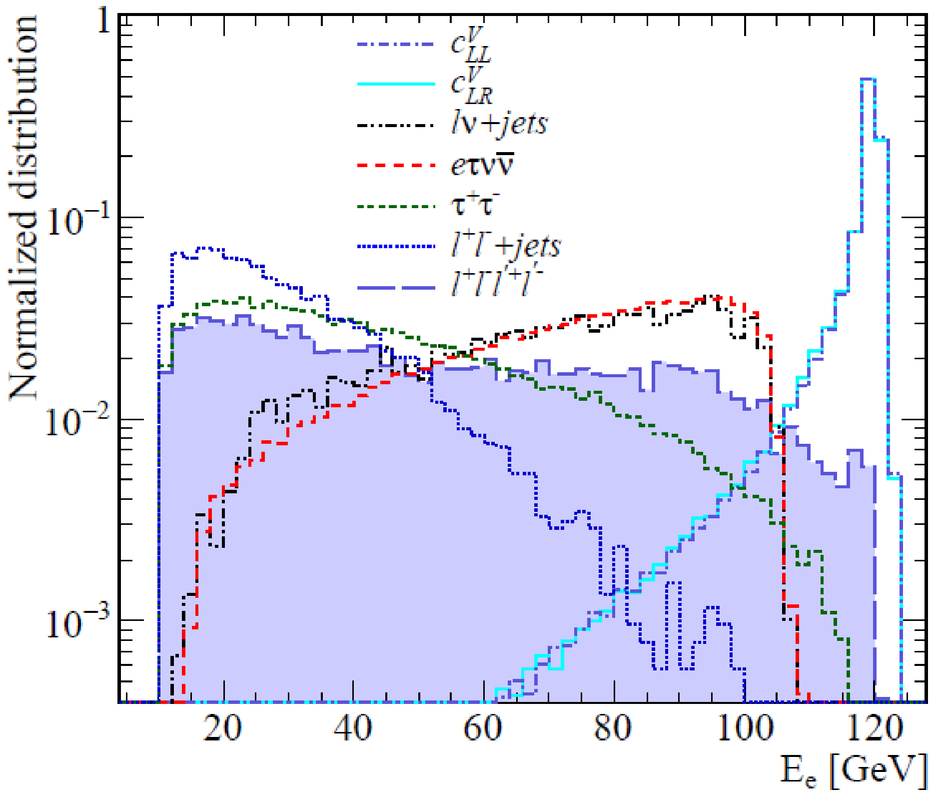
- Exactly 1 tau-tagged jet (Hadronic decay)
- Exactly 1 electron (positron)
- Opposite sign leptons

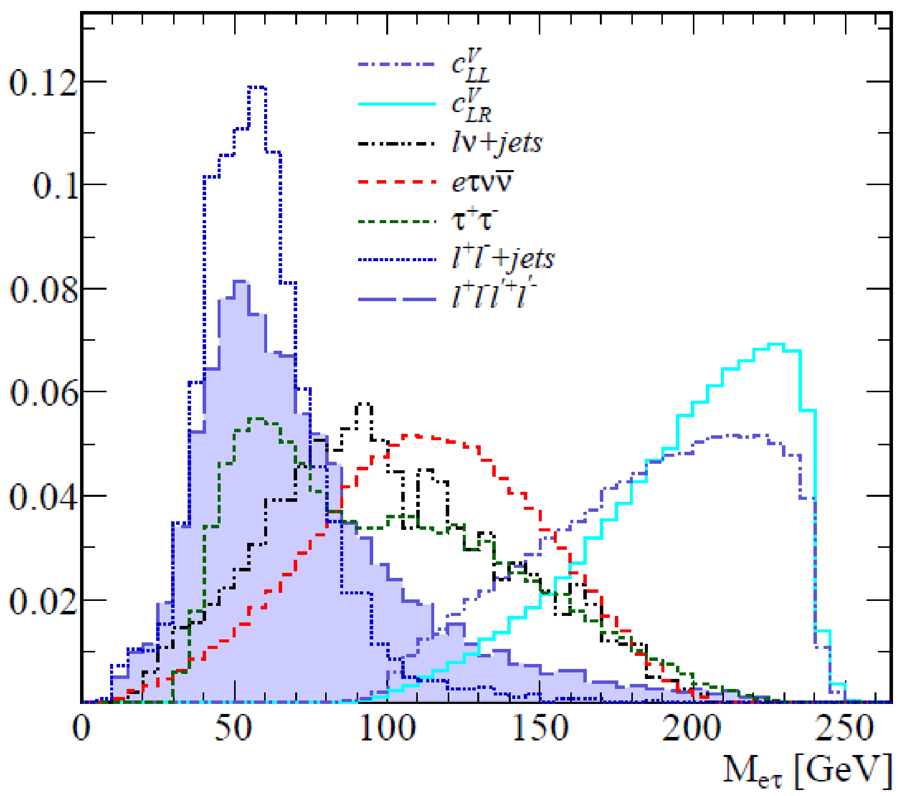
- $P_T > 20$ GeV for tau
- $P_T > 10$ GeV for electron (positron)
- $|\eta| \leq 2.5$ for all objects
- $\Delta R > 0.5$ GeV for all objects

- RelIso < 0.15; The ratio of the sum of P_T of charged particle tracks inside a cone of size 0.5 around the electron track to P_T of the electron.

ANALYSIS STRATEGY

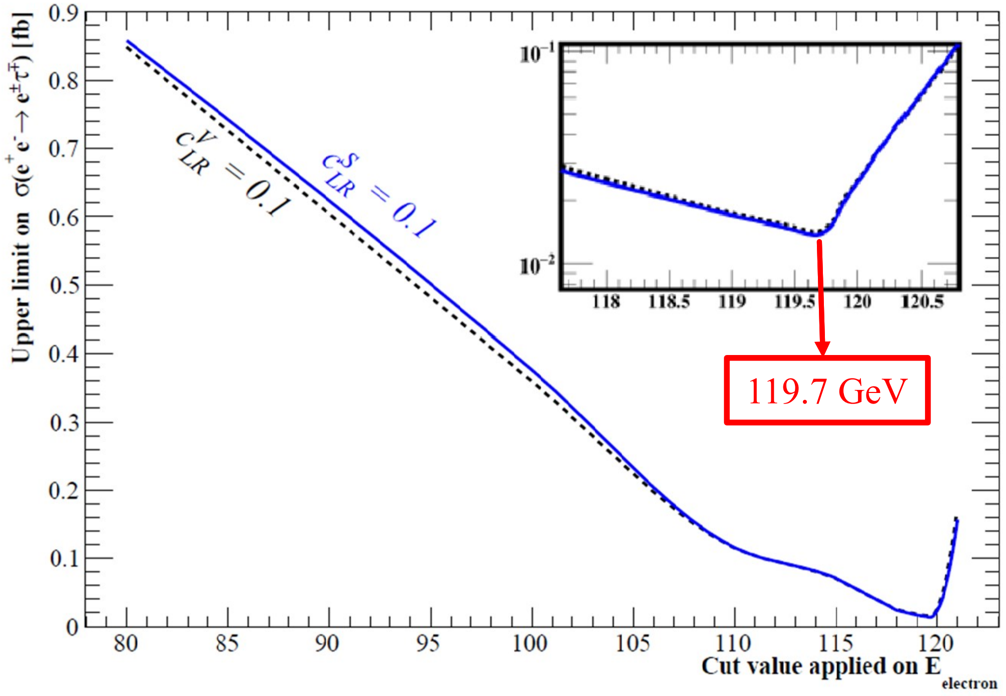
- **To enhance the sensitivity, we apply additional cuts on (for $\sqrt{s}$ =240 GeV):**

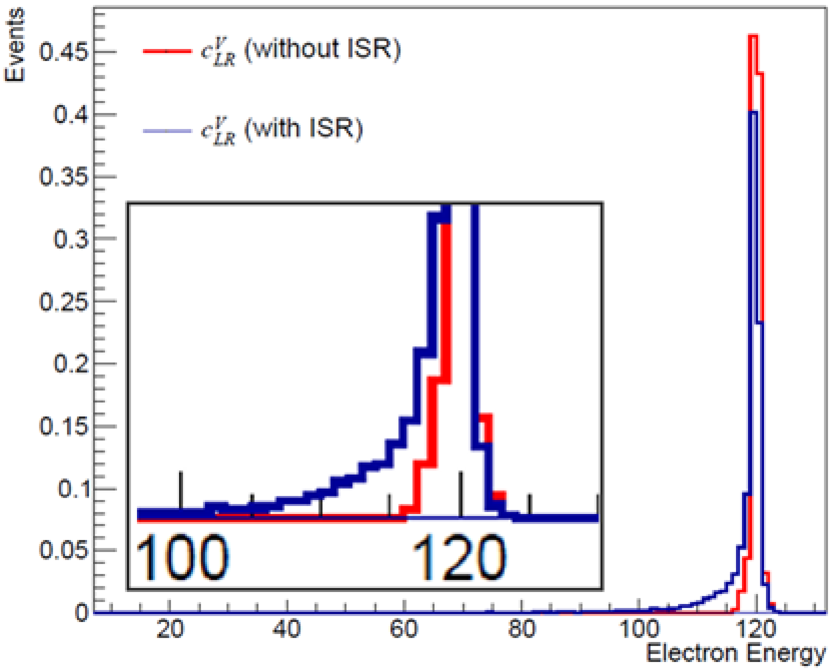




ANALYSIS STRATEGY

- **The optimized lower cuts on the energy of electron are obtained to be 78.6, 81.0, 119.7 and 182.0 GeV for center-of-mass energy of 157.5, 162.5, 240 and 365 GeV, respectively.**





ANALYSIS STRATEGY

About the $ee \rightarrow jj$ background:

- **This process contributes to the backgrounds. But …**
- **The jet fake probability is expected to be 0.1%.**
- **The rate of this background is assessed to be less than 5% of the total background contributions after event selection criteria.**

ANALYSIS STRATEGY

	Signal		SM Backgrounds				
$\sqrt{s} = 157.5$ GeV	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.1746	0.1698	0.099	0.045	4.9×10^{-3}	1.4×10^{-3}	3.3×10^{-4}
(II): $M_{e\tau} > 65$ GeV	0.1741	0.1697	0.038	0.019	2.2×10^{-3}	1.8×10^{-4}	7.5×10^{-5}
(III): $E_e > 78.6$ GeV	0.0984	0.0831	2.8×10^{-8}	1.5×10^{-7}	6.02×10^{-6}	1.7×10^{-7}	0.0
$\sqrt{s} = 162.5$ GeV	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.1727	0.1711	0.106	0.048	4.9×10^{-3}	1.6×10^{-3}	4.5×10^{-4}
(II): $M_{e\tau} > 65$ GeV	0.1727	0.1710	0.041	0.025	2.4×10^{-3}	2.1×10^{-4}	1.0×10^{-4}
(III): $E_e > 81$ GeV	0.1122	0.0949	6×10^{-8}	2.0×10^{-7}	3.61×10^{-6}	2.1×10^{-7}	0.0
$\sqrt{s} = 240$ GeV	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.2156	0.2137	0.131	0.037	8.8×10^{-3}	6.2×10^{-3}	4.9×10^{-4}
(II): $M_{e\tau} > 100$ GeV	0.2150	0.2134	0.084	0.017	1.6×10^{-3}	2.4×10^{-4}	2.0×10^{-4}
(III): $E_e > 119.7$ GeV	0.1072	0.0989	2.1×10^{-8}	1.5×10^{-7}	1.2×10^{-5}	2.4×10^{-7}	0.0
$\sqrt{s} = 365$ GeV	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.2093	0.2097	0.133	0.066	0.012	6.0×10^{-3}	5.0×10^{-4}
(II): $M_{e\tau} > 150$ GeV	0.2053	0.2051	0.093	0.041	2.0×10^{-3}	1.5×10^{-4}	2.4×10^{-4}
(III): $E_e > 182$ GeV	0.0993	0.0986	2.6×10^{-8}	3.2×10^{-7}	2.6×10^{-5}	1.4×10^{-7}	0.0

ANALYSIS STRATEGY

$\sqrt{s} = 157.5$ GeV	Signal		SM Backgrounds				
	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.1746	0.1698	0.099	0.045	4.9×10^{-3}	1.4×10^{-3}	3.3×10^{-4}
(II): $M_{e\tau} > 65$ GeV	0.1741	0.1697	0.038	0.019	2.2×10^{-3}	1.8×10^{-4}	7.5×10^{-5}
(III): $E_e > 78.6$ GeV	0.0984	0.0831	2.8×10^{-8}	1.5×10^{-7}	6.02×10^{-6}	1.7×10^{-7}	0.0

$\sqrt{s} = 162.5$ GeV	Signal		SM Backgrounds				
	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.1727	0.1711	0.106	0.048	4.9×10^{-3}	1.6×10^{-3}	4.5×10^{-4}
(II): $M_{e\tau} > 65$ GeV	0.1727	0.1710	0.041	0.025	2.4×10^{-3}	2.1×10^{-4}	1.0×10^{-4}
(III): $E_e > 81$ GeV	0.1122	0.0949	6×10^{-8}	2.0×10^{-7}	3.61×10^{-6}	2.1×10^{-7}	0.0

$\sqrt{s} = 240$ GeV	Signal		SM Backgrounds				
	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.2156	0.2137	0.131	0.037	8.8×10^{-3}	6.2×10^{-3}	4.9×10^{-4}
(II): $M_{e\tau} > 100$ GeV	0.2150	0.2134	0.084	0.017	1.6×10^{-3}	2.4×10^{-4}	2.0×10^{-4}
(III): $E_e > 119.7$ GeV	0.1072	0.0989	2.1×10^{-8}	1.5×10^{-7}	1.2×10^{-5}	2.4×10^{-7}	0.0

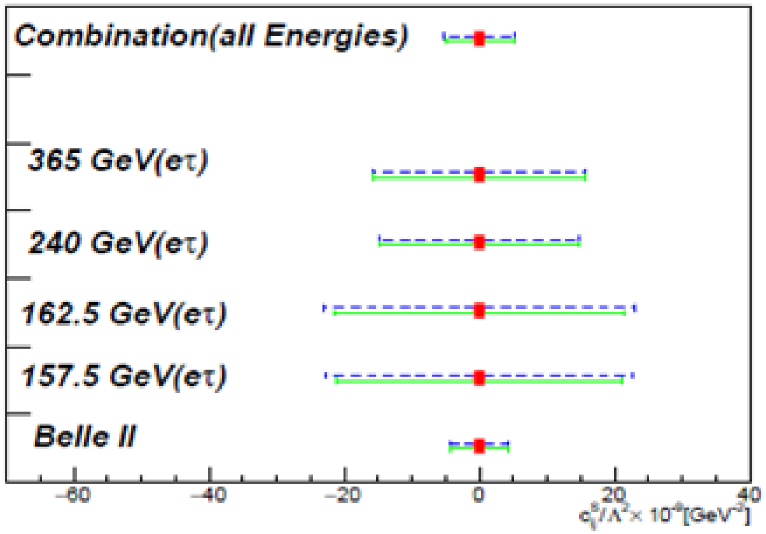
$\sqrt{s} = 365$ GeV							
	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\ell\ell'\ell'$	$\ell\ell jj$	$\ell\nu jj$
(I): Pre-selection cuts	0.2093	0.2097	0.133	0.066	0.012	6.0×10^{-3}	5.0×10^{-4}
(II): $M_{e\tau} > 150$ GeV	0.2053	0.2051	0.093	0.041	2.0×10^{-3}	1.5×10^{-4}	2.4×10^{-4}
(III): $E_e > 182$ GeV	0.0993	0.0986	2.6×10^{-8}	3.2×10^{-7}	2.6×10^{-5}	1.4×10^{-7}	0.0

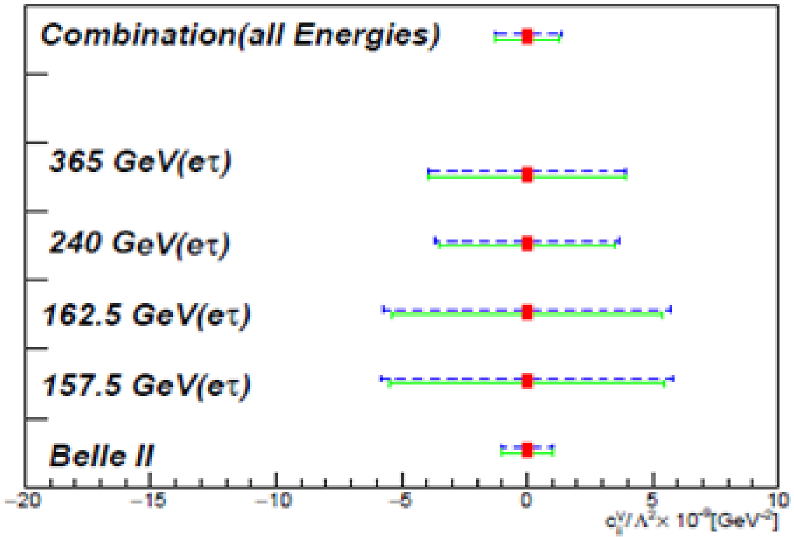
RESULTS & DISCUSSION

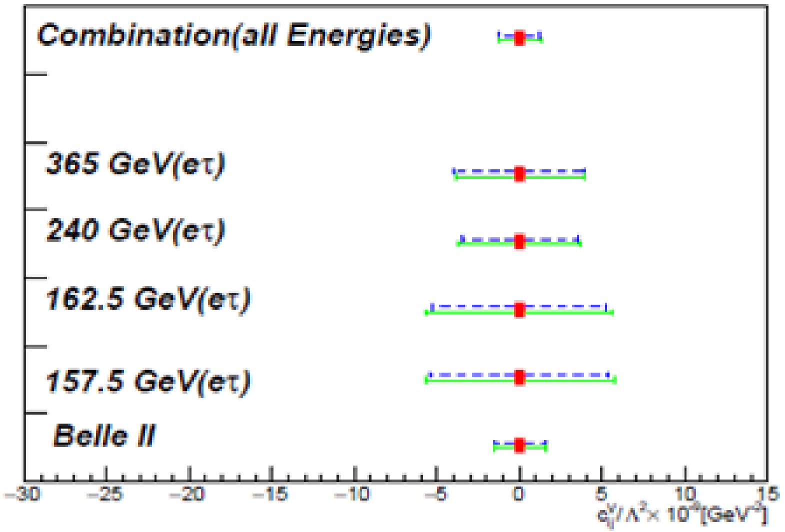
- **In order to achieve better sensitivity, the results from four energy benchmarks are combined.**
- **Comparison to the Belle-II experiment with $50\ ab^{-1}$ data** [arXiv:1808.10567]
- **Comparison to a study at $\sqrt{s} = 1$ TeV with beam polarization: $P(e^-) = 0.8, P(e^+) = -0.3$** [arXiv:1803.10475]

$\sqrt{s}$ (GeV) , $\mathcal{L}$ (ab^{-1})	$\frac{c^V_{LL}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})	$\frac{c^V_{RR}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})	$\frac{c^V_{RL}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})	$\frac{c^V_{LR}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})	$\frac{c^S_{RL}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})	$\frac{c^S_{LR}}{\Lambda^2}[\times 10^{-9}]$(GeV^{-2})
157.5 , 5	5.82	5.46	5.74	5.36	21.18	22.61
162.5 , 5	5.71	5.36	5.62	5.29	21.42	23.12
240 , 5	3.69	3.50	3.73	3.53	14.81	14.74
365 , 1.5	3.93	3.94	3.92	3.93	15.80	15.80
Combination	1.32	1.25	1.32	1.25	5.1	5.3
Belle II	1.06	1.06	1.55	1.55	4.29	4.29
$\sqrt{s} = 1$ TeV, pol. beam	4.3	1.1	1.6	1.8	13	5.9

RESULTS & DISCUSSION







RESULTS & DISCUSSION

Impact of systematic uncertainties on the results

Based on a search for LFV at LEP2 with the OPAL detector: [arXiv:0109011]

- **The systematic uncertainty on the signal efficiency is 3.5%**
- **On the number of expected background events is 5%.**

- **5% uncertainty on both signal efficiency and on background expectation at 365 GeV:**

c^V_{LL}/Λ^2	c^V_{RR}/Λ^2	c^V_{RL}/Λ^2	c^V_{LR}/Λ^2	c^S_{RL}/Λ^2	c^S_{LR}/Λ^2
4.14	4.15	4.12	4.13	16.57	16.65

$[\times 10^{-9}\ (\mathrm{GeV}^{-2})]$

CONCLUSION

- LFV processes are absent in the SM but appear in some extensions of the SM.
- The sensitivity of the FCC-ee, to the LFV couplings is examined using $e\tau$ production.
- Effective Lagrangian: four Fermi contact interactions with vector and scalar types
- The events are generated using MadGraph5 considering ISR effect and passed through PYTHIA8 and Delphes for using the ILD detector card.
- The hadronic tau decay channel and the main sources of background are considered.
- Cuts on E_e and $M_{e\tau}$ are applied to suppress the background contributions.
- Limits at 95% CL on the LFV have been obtained for the four center-of-mass energies. Finally, a statistical combination of results is performed.
- We show that the statistical combination increases the sensitivity to the LFV couplings significantly with respect to the individual energies.

THANKS FOR YOUR ATTENTION!

BACKUP

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TABLE I. The cross sections of signal $e^-e^+ \to e^\pm\tau^\mp$ and main background processes with their corresponding uncertainties are presented. The cross section of two signal scenarios are given assuming $c^V_{LR} = 0.1$, $c^S_{LR} = 0.1$, and $\Lambda = 1$ TeV. The cross sections are in the unit of fb and include the ISR effects.

$\sqrt{s}$ [GeV]	$c^V_{LR} = 0.1$	$c^S_{LR} = 0.1$	$e\tau\nu\bar{\nu}$	$\tau\bar{\tau}$	$\ell\bar{\ell}\ell'\bar{\ell}'$	$\ell\bar{\ell}jj$	$\ell\nu jj$	jj
157.5	4.72 ± 0.007	0.29 ± 0.0004	22.33 ± 0.07	11076.5 ± 3.4	39.86 ± 0.08	80.95 ± 0.2	272.9 ± 0.4	32032 ± 8.1
162.5	5.02 ± 0.007	0.31 ± 0.0004	102.12 ± 0.3	10275.8 ± 2.9	42.23 ± 0.08	83.06 ± 0.3	1198.05 ± 0.8	29133 ± 6.2
240	10.98 ± 0.04	0.69 ± 0.0008	415.63 ± 0.6	4196.8 ± 1.2	86.24 ± 0.2	217.8 ± 0.5	4552.7 ± 1.3	10481 ± 3.5
365	25.26 ± 0.07	1.57 ± 0.002	327.59 ± 0.5	1803.6 ± 0.6	85.05 ± 0.1	195.13 ± 0.3	3247.02 ± 1.1	4306 ± 1.2

BACKUP

SM relevant input values:

$$
\begin{aligned}
M_Z &= 91.188 \ \text{GeV} \\
m_\tau &= 1.777 \ \text{GeV} \\
G_F &= 1.166 \times 10^{-5} \ \text{GeV}^{-2} \\
\alpha_e &= 1/127.9 \\
\alpha_s &= 0.118
\end{aligned}
$$

BACKUP

```
#############
# tau-tagging
#############


module TauTagging TauTagging {
  set ParticleInputArray Delphes/allParticles
  set PartonInputArray Delphes/partons
  set JetInputArray JetEnergyScale/jets

  set DeltaR 0.5

  set TauPTMin 1.0

  set TauEtaMax 4.0

  # add EfficiencyFormula {abs(PDG code)} {efficiency formula as a function of eta and pt}

  # default efficiency formula (misidentification rate)
  add EfficiencyFormula {0} {0.001}
  # efficiency formula for tau-jets
  add EfficiencyFormula {15} {0.4}
}
```

BACKUP

```
######################
# Electron efficiency
######################

module Efficiency ElectronEfficiency {
  set InputArray ElectronFilter/electrons
  set OutputArray electrons

  # set EfficiencyFormula {efficiency formula as a function of eta and pt}

  # efficiency formula for electrons
  set EfficiencyFormula {                                     (pt <= 10.0) * (0.00) +
                                           (abs(eta) <= 1.5) * (pt > 10.0)  * (0.95) +
                         (abs(eta) > 1.5 && abs(eta) <= 2.5) * (pt > 10.0)  * (0.95) +
                         (abs(eta) > 2.5)                                   * (0.00)}
}
```

```
######################
# Electron isolation
######################

module Isolation ElectronIsolation {
  set CandidateInputArray ElectronEfficiency/electrons
  set IsolationInputArray EFlowFilter/eflow

  set OutputArray electrons

  set DeltaRMax 0.5

  set PTMin 0.5

  set PTRatioMax 0.12
}
```

BACKUP

- For electrons with $P_T > 10$ GeV and $|\eta| \leq 2.5$, the identification efficiency in the ILD card is 95%.

- The efficiency in the ILD simulation card is 40% and the tau misidentification rate is assumed to be equal 0.1%.
- considering the τ tagging efficiency, a jet is considered potentially as a τ candidate if a generated τ exists within a bellow distance from the jet axis.

$$\Delta R = \sqrt{(\eta_{\text{jet}} - \eta_{\tau})^2 + (\phi_{\text{jet}} - \phi_{\tau})^2} = 0.3$$

BACKUP

- **It is notable that the Meτ distribution has different behaviours for cV LL and cV LR which arises from the fact that for LL coupling** $d\sigma/d\cos\theta \propto (1+\cos\theta)^2$

  for RL coupling $d\sigma/d\cos\theta \propto (1-\cos\theta)^2$

BACKUP

Limit setting method

- The CLs technique is exploited to find upper limits on the signal cross section
- The **RooStats** package is used to perform the numerical evaluation of the CLs.
- **CLs technique**: we define log-likelihood functions L_{Bkg} and $L_{Signal+Bkg}$ for the background hypothesis, and for the signal+background hypothesis as the multiplication of Poissonian likelihood functions.
- The p-value for hypothesis of signal+background and for the background hypothesis are determined using the log-likelihood ratio:

$$\mathrm{Q} = -2\ln(\mathrm{L}_{\mathrm{Signal+Bkg}}/\mathrm{L}_{\mathrm{Bkg}})$$

- The signal cross section is constrained using

$$\mathrm{CL_s} = \mathrm{P}_{\mathrm{Signal+Bkg}}(\mathrm{Q} > \mathrm{Q}_0)/(1 - \mathrm{P}_{\mathrm{Bkg}}(\mathrm{Q} < \mathrm{Q}_0)) \leq 0.05$$

which is corresponding to 95% CL and Q_0 is the expected value of test statistics Q.

BACKUP

- **Other info:**
- There are a variety of theories that give rise to LFV. For instance, additional fermions present in the type IIIseesaw model or in the low-scale seesaw models give rise to large LFV effects
- production rate of the four-fermion interactions grows linearly with the squared center-of-mass energy s, and diverge when $s \to \infty$. However, one should note that we are working in a non-renormalizable formalism and these operators provide an acceptable description of physics at high energy up to an energy scale Λ.