



Beam-cavity interaction studies for the FCC-ee – RF frequency considerations

Ivan Karpov, Rama Calaga, Heiko Damerau, Linhao Zhang, Elena Shaposhnikova, Philippe Baudrenghien, Franck Peauger, Olivier Brunner, Igor Syratchev, Sosoho-Abasi Udongwo, Shahnám Gorgi Zadeh

Acknowledgements:

Tor Raubenheimer, Katsunobu Oide, and many other members of FCC-ee optics and FCC SRF

Interaction of beam with cavity impedance

Main effects that need to be considered:

- Higher-order-mode power losses
- Beam loading (steady-state and transient)
- Coupled-bunch instabilities (longitudinal and transverse)

→ Operation at Z energy is the most challenging

Most of them were addressed for the CDR parameters, but

- Beam and accelerator parameters keep changing
- Alternative scenarios emerge

→ Re-evaluation of beam-cavity interaction aspects is needed

Update of parameters

Layout & placement optimization results in a smaller FCC circumference (PA31-1.0)

→ A decrease of the beam current by about 8% for all energies

Optimization of luminosity for 4 IPs

→ Higher bunch charge

→ Higher RF voltage

Parameter	FCC week 2022	CDR
Total current, J_A	1.28 A	1.39 A
Bunch intensity, N_p	2.43×10^{11}	1.7×10^{11}
Number of bunches, M_b	10000	16640
Bunch length (BS), σ	14.5 mm	12.1 mm
Total RF voltage, V_{tot}	120 MV	100 MV

These parameters will change after taking into account the precise choice based on FCC-hh RF synchronization aspects (see slides of L. Zhang)

→ What is the impact on operation at Z energy?

HOM power losses

Simulated cavity
impedance

Normalized Fourier harmonics
of beam current

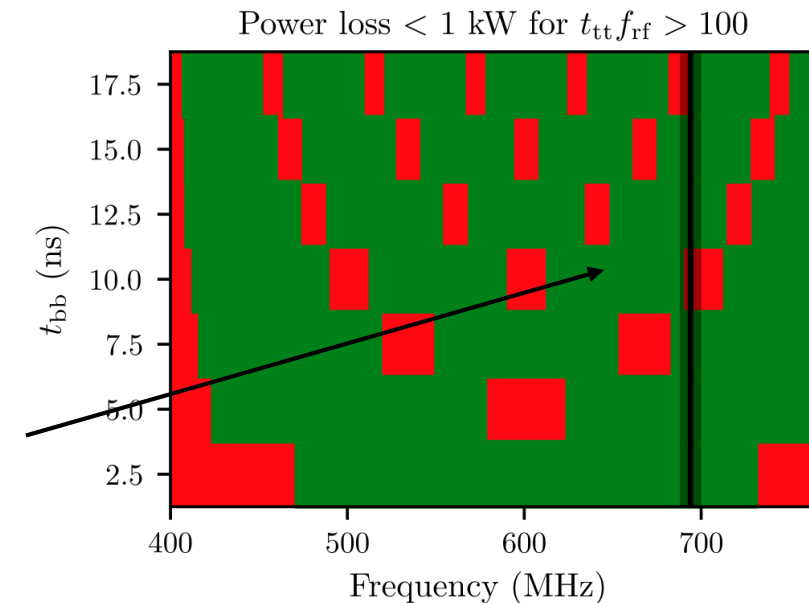
$$P = I_{b,DC}^2 \sum_{k=-\infty}^{\infty} \text{Re}[Z_{||}(k f_{\text{rev}})] |I_k|^2$$

$I_{b,DC}$ – average beam current

f_{rev} – revolution frequency

k – revolution harmonic number

Detailed analysis was performed for single-cell cavity design of 2015* with HOM below cut-off frequency of the beam pipe
→ Acceptable filling schemes were defined for operation
→ For all recent cavity designs, the first HOM is above cut-off frequency



*I. Karpov, R. Calaga, E. Shaposhnikova, PRAB 21, 071001 (2018)

Impact of higher bunch charge

HOM power loss for broadband impedance can be approximated

$$P_{\text{HOM}} = I_{b,\text{DC}}^2 \sum_{k=-\infty}^{\infty} \text{Re}[Z_{\parallel}(k f_{\text{rev}})] |I_k|^2$$

Gaussian bunches

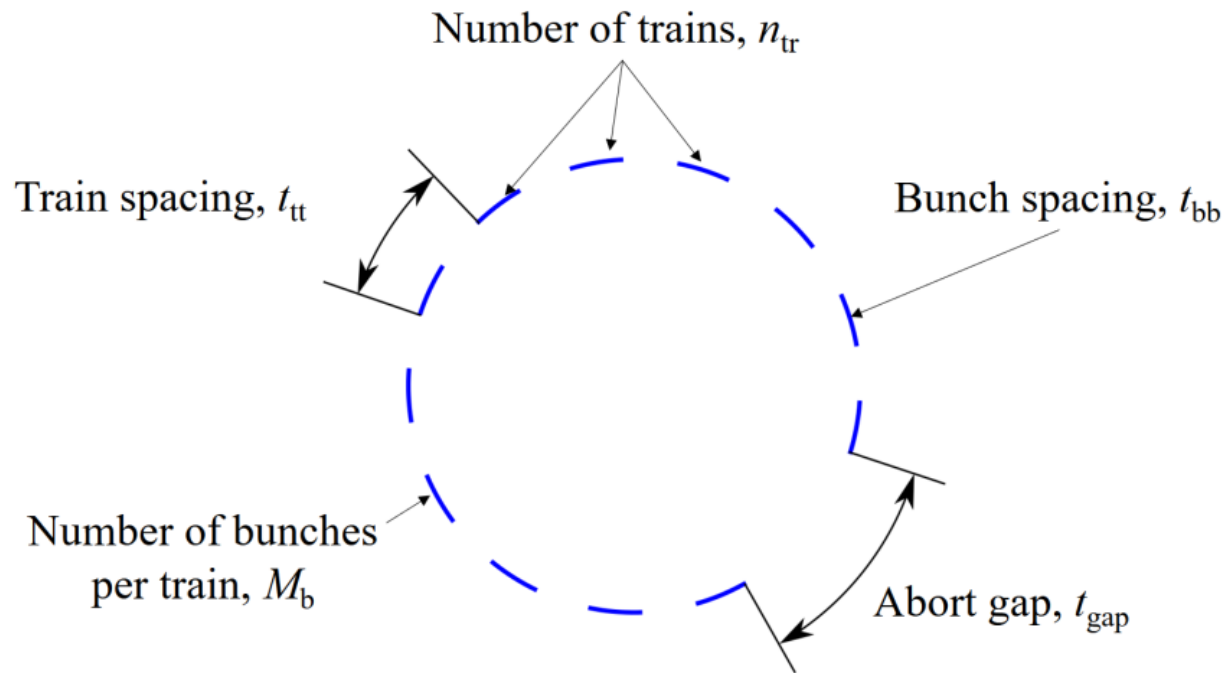
$$\approx e f_0 I_{b,\text{DC}} N_p \sum_{k=-\infty}^{\infty} \text{Re}[Z_{\parallel}(k f_{\text{rev}})] e^{-(2\pi k f_0 \sigma / c)^2}$$

Since bunch length scales* as $\sigma \propto \sqrt{N_p}$

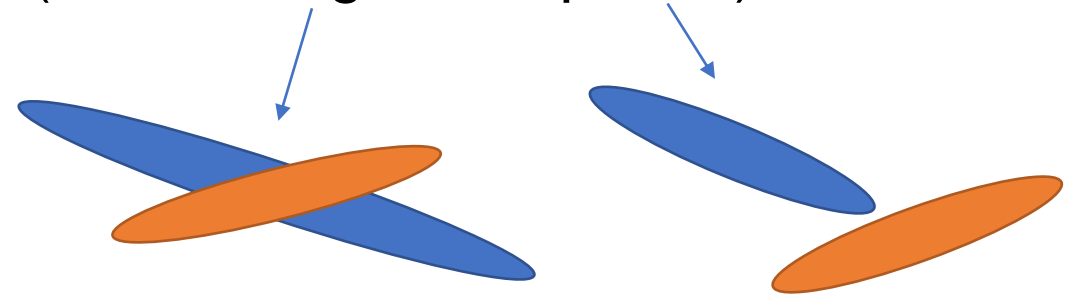
→ For the worst-case scenario, $\text{Re}Z_{\parallel}(k f_{\text{rev}}) = \text{const.}$, $P_{\text{HOM}} \propto \sqrt{N_p}$ and thus weakly depends on parameter variations

*D. Shatilov, ICFA Beam Dyn. Newslett. 72, 30 (2017)

Transient beam loading



Gaps in machine filling will result in modulation beam parameters (bunch length and phase)



→ Might have impact on luminosity

Conventional approaches:

- Small-signal model in frequency domain*, which assumes small modulations (but we have 100% modulation of beam current!)
 - Particle tracking simulations (difficult for 10000 bunches in FCC-ee Z)
- We use steady-state time domain method**

* *F. Pedersen, RF Cavity feedback, CERN/PS 92-59 (1992)*

** *J. Tückmantel, CERN Report No. CERN-ATS-Note-2011- 002 TECH, 2011*

Example for single-cell cavity in FCC Z

→ There is a strong modulation due to the abort gap and a fine structure due to the gaps between trains

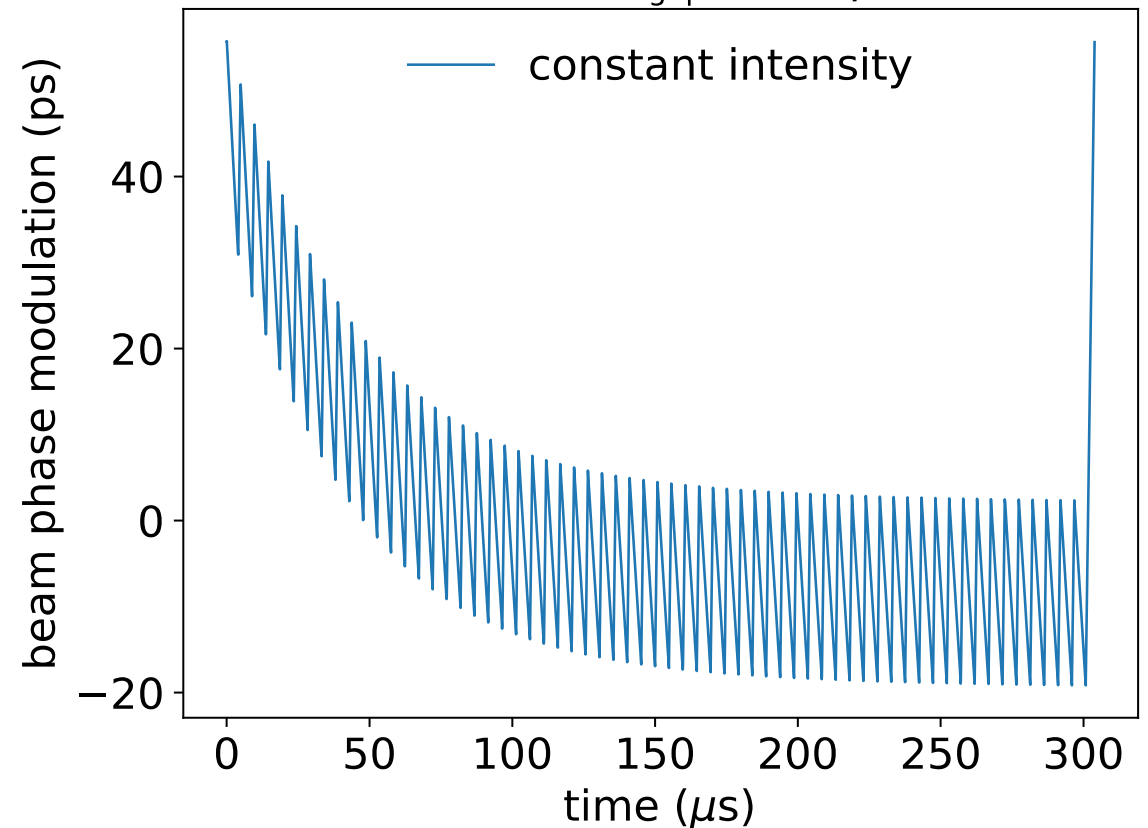
→ For identical rings, transients can be compensated by matching abort gaps (e.g., in PEPII, LHC,...); one gap is sufficient for 4 symmetric IPs

Imbalance of charge results in different detuning for electron and positron beams

→ Slightly different transients

→ The collision point shift is negligible for $\pm 5\%$ random spread of bunch charge

$$t_{tt}f_{rf} = 1950, t_{bb} = 25.0 \text{ ns}, M_b = 161, \\ n_{tr} = 62, t_{gap} = 3.1 \mu\text{s}$$



Example for single-cell cavity in FCC Z

→ There is a strong modulation due to the abort gap and a fine structure due to the gaps between trains

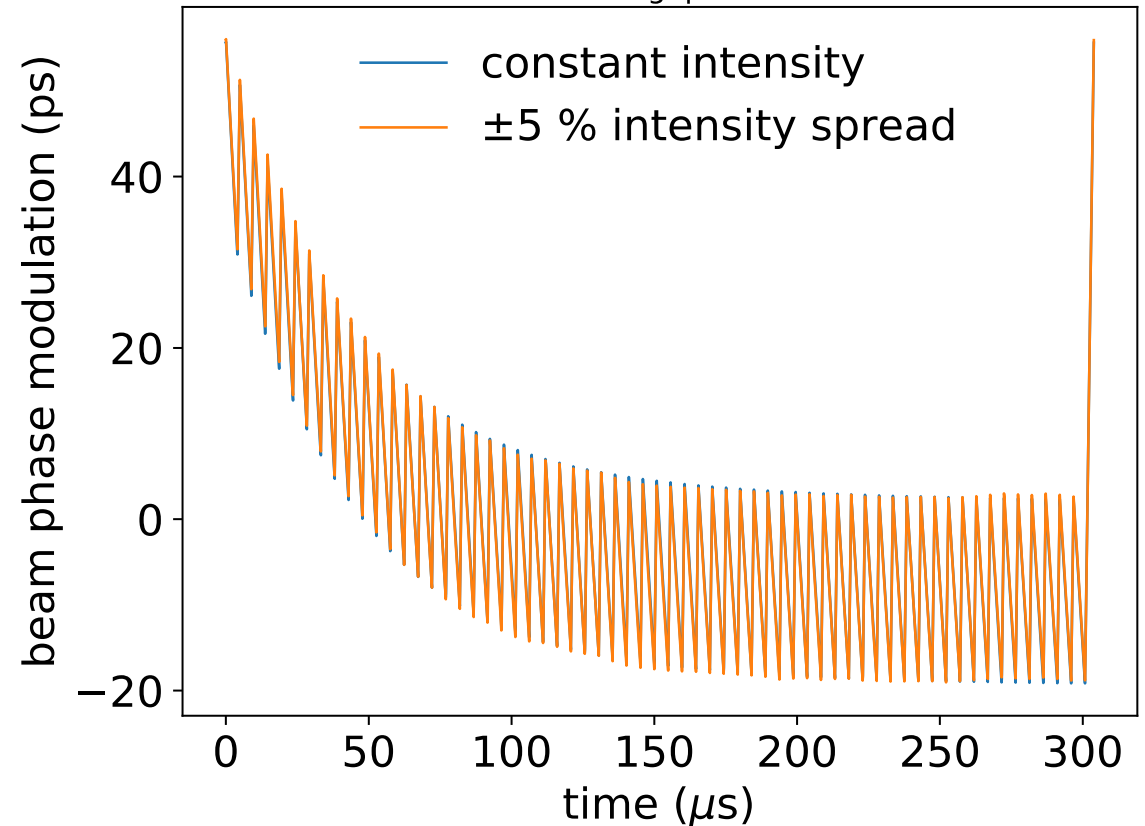
→ For identical rings, transients can be compensated by matching abort gaps (e.g., in PEPII, LHC,...); one gap is sufficient for 4 symmetric IPs

Imbalance of charge results in different detuning for electron and positron beams

→ Slightly different transients

→ The collision point shift is negligible for $\pm 5\%$ random spread

$$t_{\text{tt}}f_{\text{rf}} = 1950, t_{\text{bb}} = 25.0 \text{ ns}, M_{\text{b}} = 161, \\ n_{\text{tr}} = 62, t_{\text{gap}} = 3.1 \mu\text{s}$$



Example for single-cell cavity in FCC Z

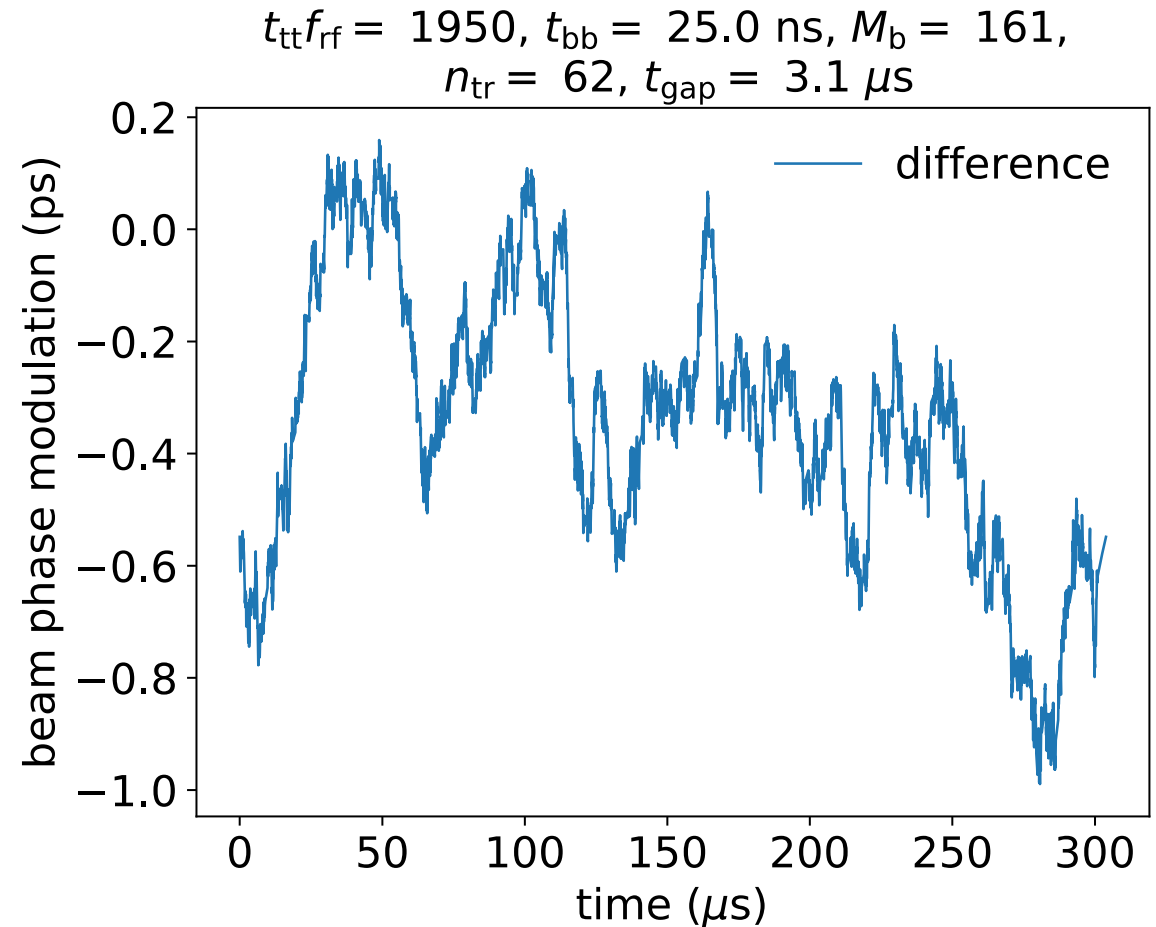
→ There is a strong modulation due to the abort gap and a fine structure due to the gaps between trains

→ For identical rings, transients can be compensated by matching abort gaps (e.g., in PEPII, LHC,...); one gap is sufficient for 4 symmetric IPs

Imbalance of charge results in different detuning for electron and positron beams

→ Slightly different transients

→ The collision point shift is negligible for $\pm 5\%$ random spread



Longitudinal CBI due to fundamental mode

For short Gaussian bunches the growth rate of the mode m is*

$$\frac{1}{\tau_m} \approx \frac{e\eta I_{b,DC} N_{cav}}{4\pi E_b Q_s} \omega_{RF} \{ \text{Re}[Z_{cl}(\omega_+)] - \text{Re}[Z_{cl}(\omega_-)] \},$$

with $\omega_{\pm} = \omega_{RF} \pm (m + Q_s)\omega_{rev}$

Closed loop impedance with direct RF feedback**

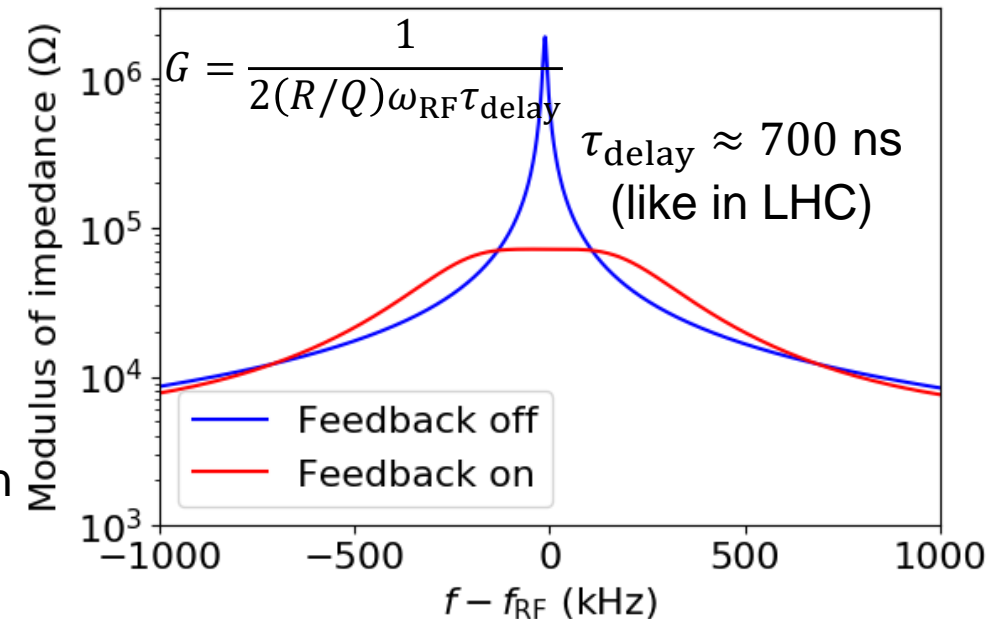
$$Z_{cl}(\omega) = \frac{Z(\omega)}{1 + e^{-i\tau_{delay}\omega} G e^{i\phi} Z(\omega)}$$

τ_{delay} – overall delay

$G = \frac{1}{2(R/Q)\omega_{RF}\tau_{delay}}$ - FB gain

ϕ – phase adjustment

Cavity impedance at fundamental, $Z(\omega)$



Passive damping sufficient if $\tau_m > \tau_{SR}$ - synchrotron radiation damping time (1170 turns at Z pole)
 → Since, $\eta, I_{b,DC}, N_{cav}, Q_s$ have changed slightly, we don't expect any significant impact on beam stability

* J. L. Laclare, CAS, (1985)

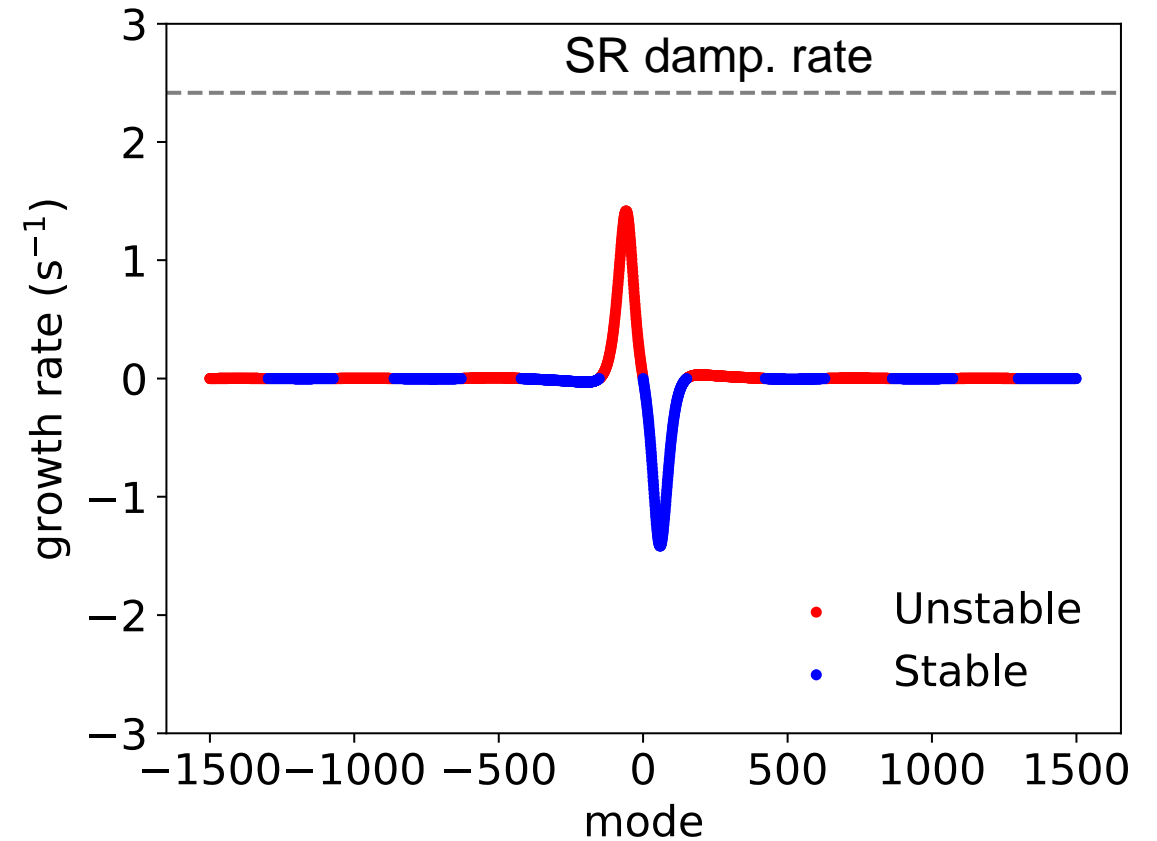
** F. Pedersen, RF Cavity feedback, CERN/PS 92-59 (1992)

Growth rates vs bunch mode number

→ Direct RF feedback reduces CBI growth rates below the threshold.

→ An additional feedback (1-turn delay, multi-harmonic?) could be implemented to provide additional margin.

Case of direct RF feedback only

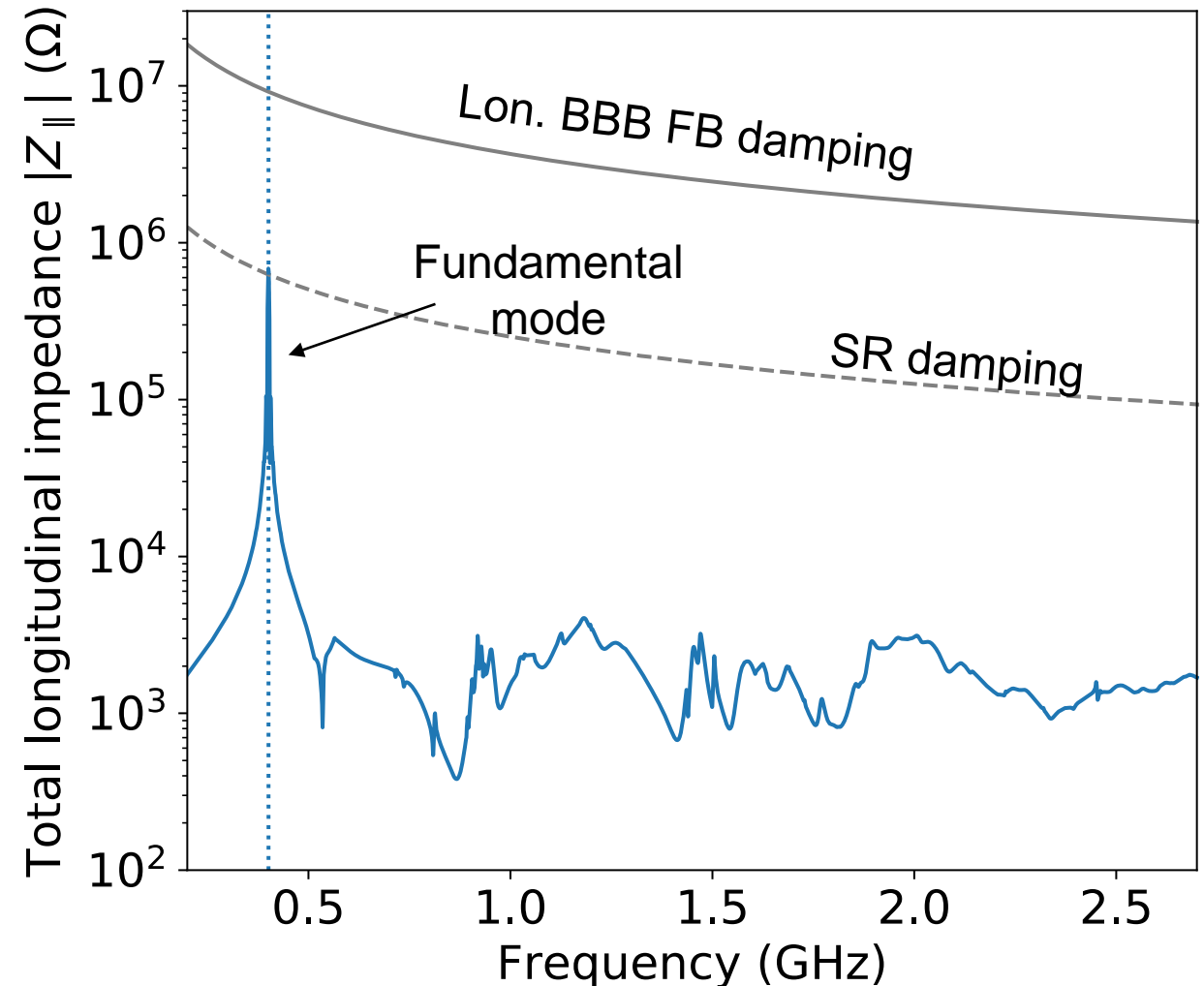


Longitudinal CBI due to HOMs at Z pole

A standard formula for threshold
(only one sideband contributes)

$$Z_{\parallel}^{\text{th}}(f) = \frac{2E_b Q_s}{e I_{b,DC} \eta \tau_{SR}} \frac{1}{f}$$

→ CBI instabilities due to HOMs
will be suppressed by
synchrotron radiation

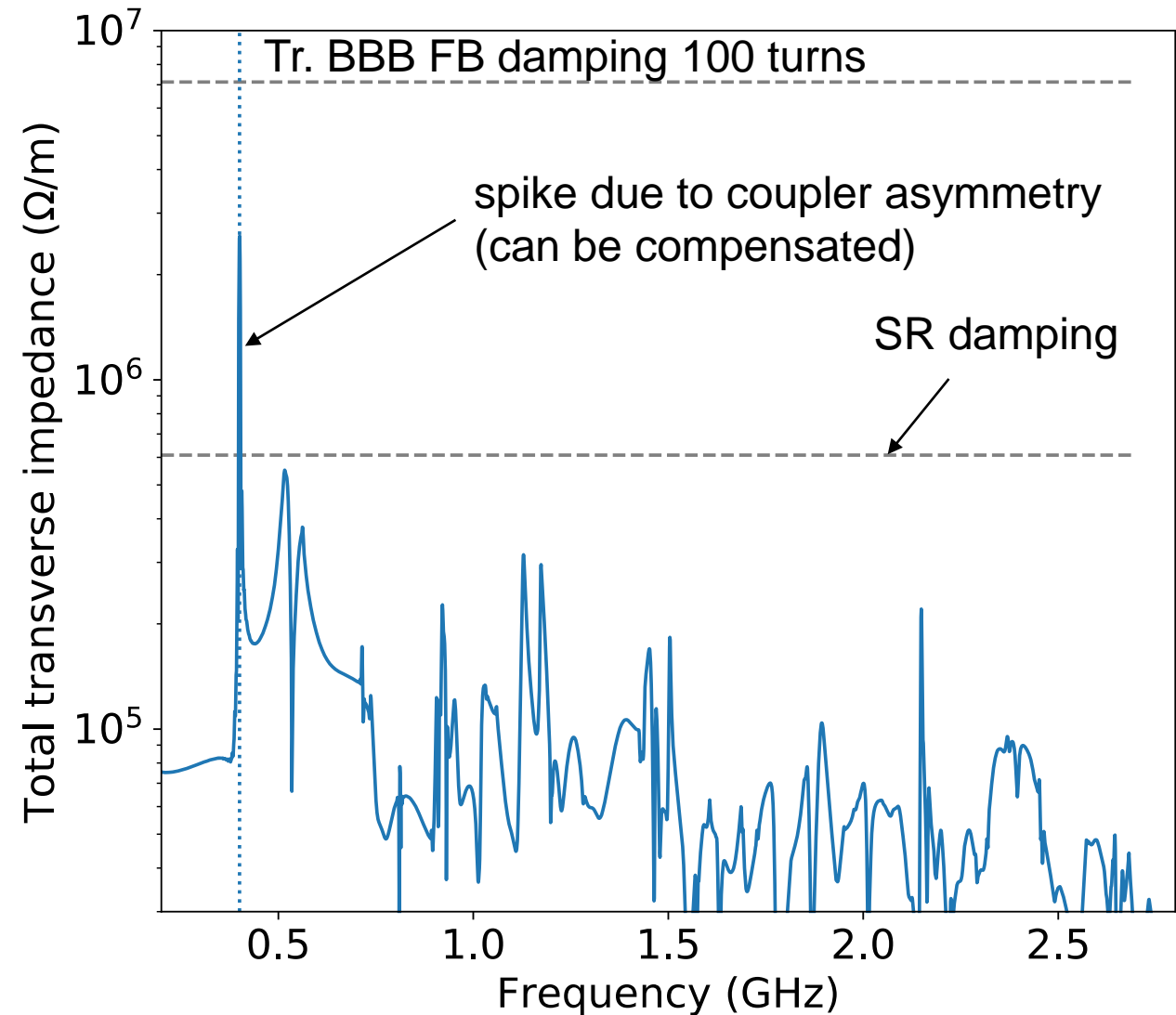


Transverse CBI due to HOMs

Similar expression for threshold

$$Z_{\perp}^{\text{th}} = \frac{E_b}{e f_{\text{rev}} I_{b,\text{DC}} \beta_{xy} \tau_{\text{SR}}}$$

- HOM below the cut-off frequency is close to the CBI threshold
- Bunch-by-bunch feedback system damping time of about 100 turns should be sufficient to suppress instabilities due to HOMs



Alternative scenario

In CDR 400 MHz 4-cell cavities are used from W + for ttbar 800 MHz, 5-cell cavities will be installed in addition

In 2018*, a “Hybrid scheme” was proposed to split RF systems for low energy (Z, W) and high energy (H,ttbar) operation points (additionally to replace 4-cell with 2-cell cavities to be below the threshold of transverse CBI)

→ The present scenario also assumes the same RF system for e- and e+ beams for the H pole (similar to ttbar)**

→ RF voltage for ttbar is dominated by 800 MHz (its contribution is optimized as suggested by T. Raubenheimer and K. Oide***)

Parameter	Unit	H (ZH)	ttbar
Total current, J_A	mA	26.7	5.0
Total RF voltage, V_{tot}	GV	2.1	2.1/8.2
SR energy loss/turn, U_0	GeV	1.9	10.0
RF frequency, f_{RF}	MHz	400	400/800

* S. Gorgi Zadeh et al, IPAC, 2018

** F. Peauger, V. Parma, O. Brunner, 151st FCC-ee Optics Design Meeting & 22nd FCCIS WP2.2 Meeting, Mar. 17, 2022

*** K.Oide, FCC-ee parameter meeting, Nov. 16, 2021

preliminary

Summary

Beam-cavity interaction for FCC-ee operating at the Z pole is the most challenging (high beam current, a large number of bunches, etc).

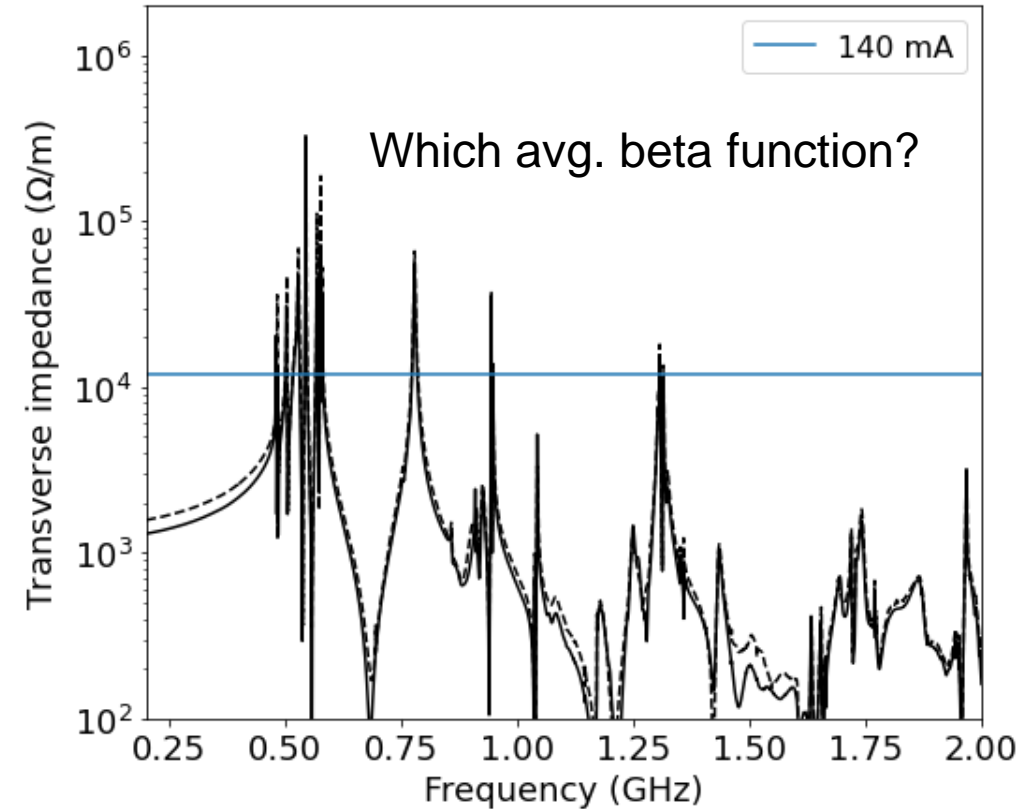
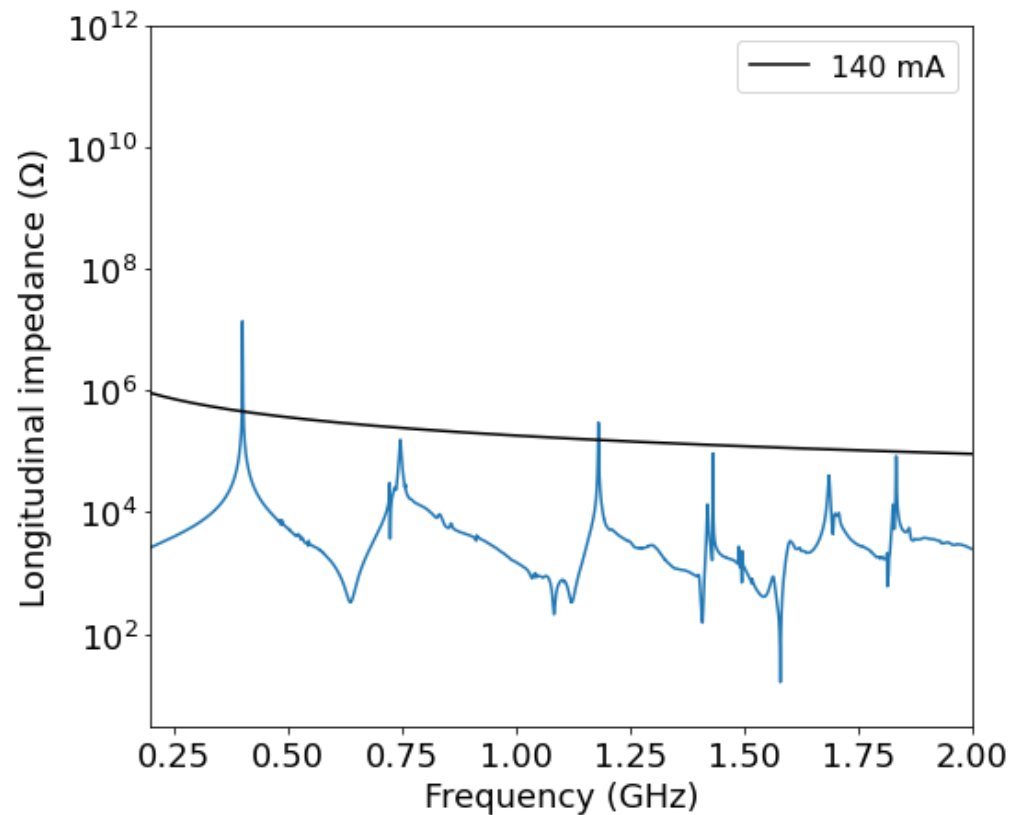
They are re-evaluated for the latest parameter set:

- HOM power losses weakly depend on single-bunch intensity.
- The impact of transient beam loading on the displacement of collision points is negligible.
- Longitudinal coupled-bunch instabilities are under control.
- To suppress transverse CBI due to HOMs bunch-by-bunch transverse feedback system is required.

Next steps:

- CBI and beam loading for W, H, ttbar calibration at Z energy
- HOM power and stability in the booster (24s at 20 GeV with the beam current of 140 mA)

Coupled-bunch instabilities in Booster



Longitudinal is very close to the threshold.

Thank you for your attention!

Steady-state beam loading

RF power per cavity in presence of beam loading $P = V_{\text{cav}} I_{b,\text{DC}} \cos \phi_s = 50 \text{ MW}/N_{\text{cav}}$ is minimized by using*

Optimal quality factor $Q_L = \frac{V_{\text{cav}}}{2(R/Q) I_{b,\text{DC}} \cos \phi_s}$ Red – fixed parameters

Optimal detuning $\Delta\omega = \omega_0 - \omega_{\text{RF}} = -\frac{\omega_{\text{RF}}(R/Q) I_{b,\text{DC}} \sin \phi_s}{V_{\text{cav}}}$

Lower voltage requires less RF power but results in **larger detuning**.

→ Transient beam loading can potentially affect luminosity

→ Longitudinal coupled bunch instability (CBI) due to fundamental mode can be an issue

* *D. Boussard, Control of cavities with high beam loading, IEEE Trans. Nucl. Sci. 32, 1852 (1985)*

- Next steps:
- HOM power and stability in Full energy booster (should be easy to get longitudinal CBI threshold curve at 20 GeV)
- CBI and beam loading for H calibration at Z energy