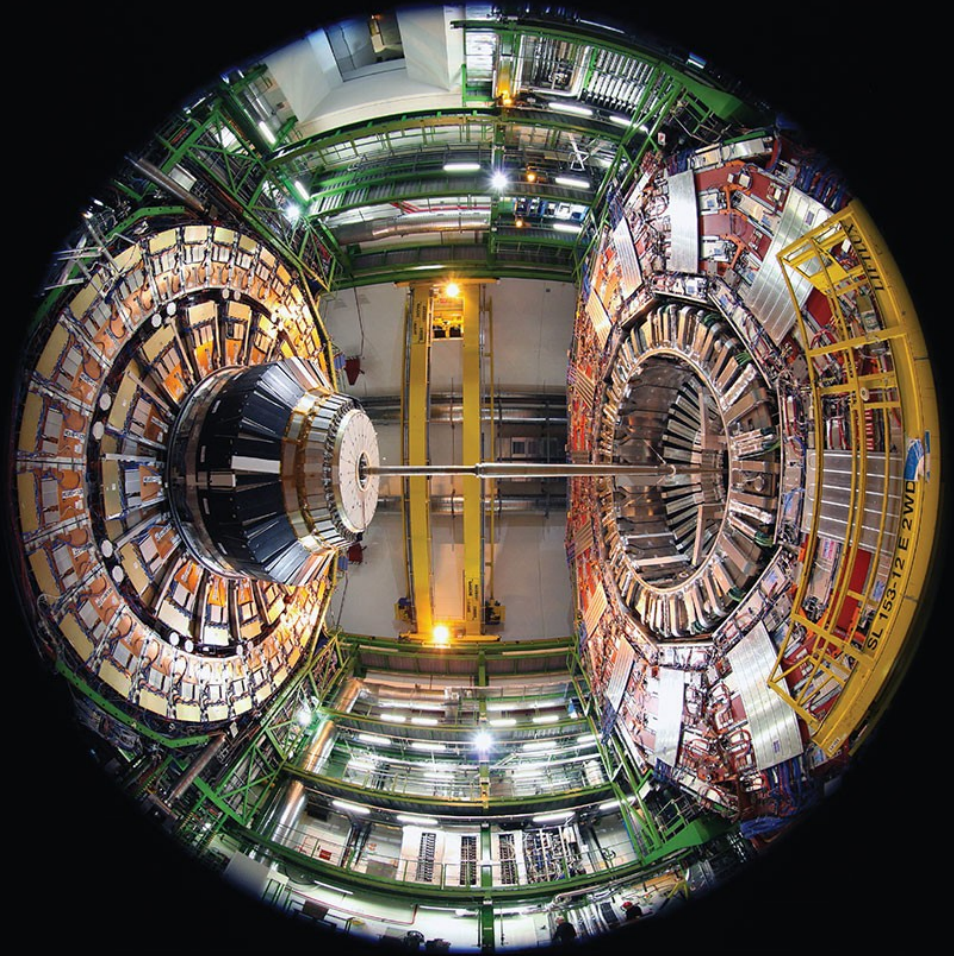


Calorimetry in high energy physics *Part 4*

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Lecture plan

What is calorimeter in HEP?

Electromagnetic and hadronic showers

Detection techniques

**Calorimeter
response & resolution**

Closing remarks

Energy resolution

- Calorimeter's energy resolution is determined by **fluctuations**
- Input energy $E \propto N$ **number of secondary particles**
 - Poisson distribution of $N \rightarrow \sigma(E)/E \propto \sqrt{N}/N \propto 1/\sqrt{E}$
 - Although in reality only a fraction can be detected (threshold effects)
- Other types of fluctuations
 - Signal quantum fluctuations (e.g. photoelectron statistics)
 - Sampling fraction
 - Shower leakage
 - Instrumental effects (electronic noise, light attenuation, non-uniformity, etc.)
 - Hadronic-specific fluctuations (EM fraction, invisible energy)

Sampling fluctuations

■ Two aspects in sampling fluctuations

■ **Sampling fraction:** fraction of energy deposited in active material **by a mip**

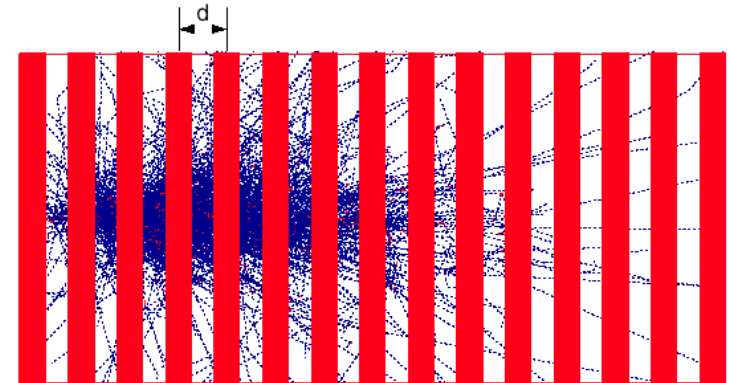
$$f_{\text{samp}} = \frac{E_{\text{loss in active}}^{\text{mip}}}{E_{\text{total loss}}^{\text{mip}}}$$

$$f_{\text{samp}} = \frac{d_{\text{active}} \times (dE/dx)_{\text{active}}^{\text{mip}}}{d_{\text{absorber}} \times (dE/dx)_{\text{absorber}}^{\text{mip}} + d_{\text{active}} \times (dE/dx)_{\text{active}}^{\text{mip}}}$$

○ Lower sampling → less particles collected → larger fluctuations

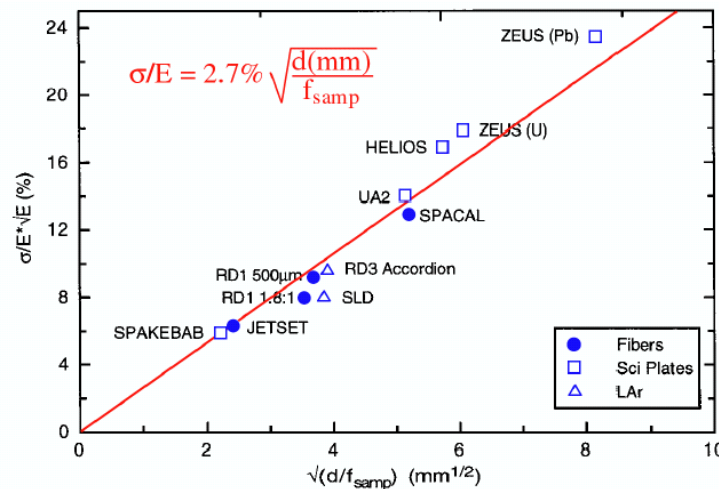
■ **Active layer thickness**

- Large fraction of low energy electrons (< 1MeV) produced in absorber
- Traveling a small distance in active material



Thicker active layer → worse resolution
Lower sampling → worse resolution

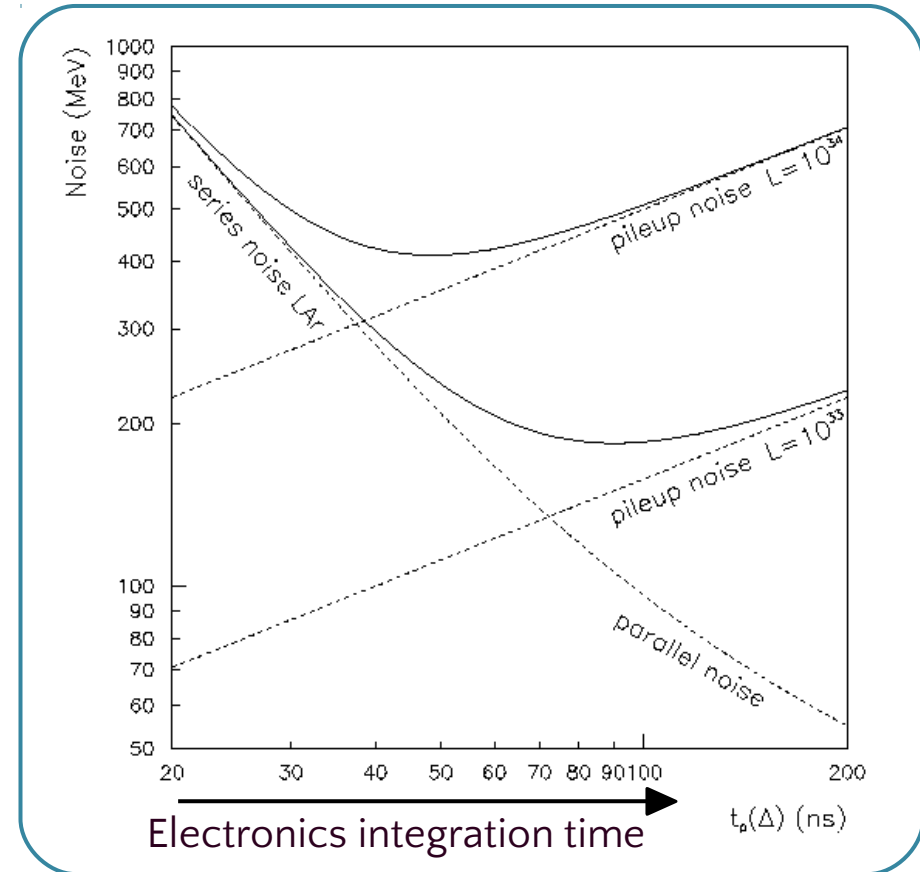
$$\frac{\sigma(E)}{E} \propto \sqrt{\frac{d_{\text{active}}}{f_{\text{samp}}}} \frac{1}{\sqrt{E}}$$



Noise

- Noise fluctuations are constant in energy
 - → Impact resolution in $1/E$ (mainly low energy)
- Usually comes from the **electronics readout** system
- But at hadron colliders
 - Contributions from **pile-up interactions**
 - = fluctuations due to multiple low energy collisions

Electronic noise vs pile-up noise
(ATLAS LAr calorimeter)



Leakage

- Energy from secondary particles escaping measurement

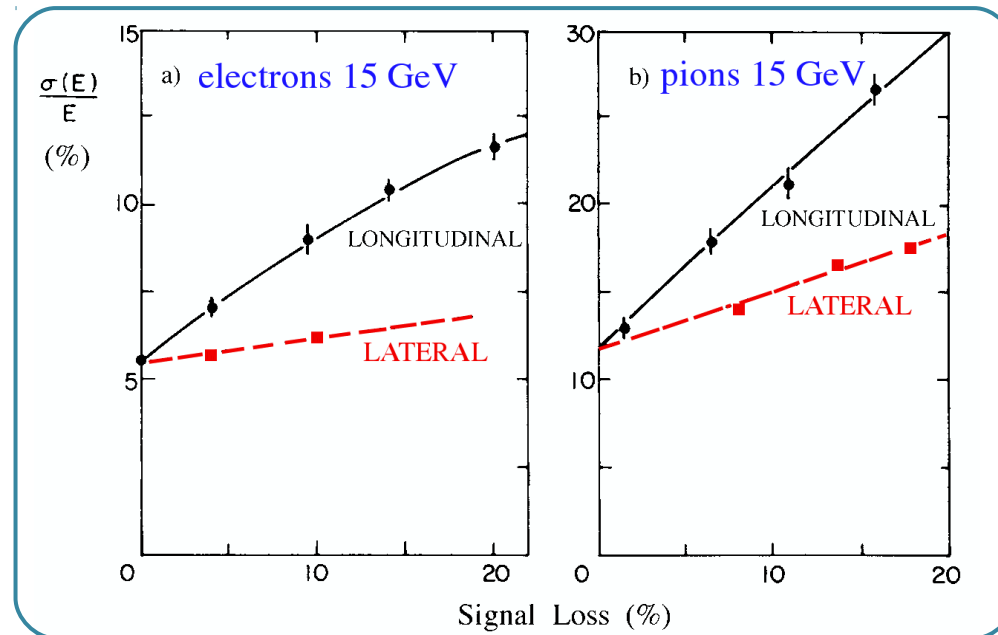
- Non-Poissonian fluctuations

- **Longitudinal** leakage (rear of the detector)

- A detector is never infinitely deep
- Dangerous since increases as $\log(E)$
- Alleviated if calorimeter “sufficiently” deep

- **Lateral** leakage

- One tends to limit the lateral size over which the signal is integrated
- Need to limit integration of channels with low S/N
- Need to limit integration of nearby showers



Contribution of leakage fluctuations to energy resolution

Energy resolution: parametrization

$$\frac{\sigma(E)}{E} = \frac{\overset{\text{Stochastic}}{S}}{\sqrt{E}} \oplus \frac{\underset{\text{Noise}}{N}}{E} \oplus \overset{\text{Constant}}{C} \quad \oplus = \text{quadratic sum}$$

■ Stochastic term

- Everything with a **Poisson-like** statistics
- Intrinsic particle fluctuations, sampling, quantum fluctuations

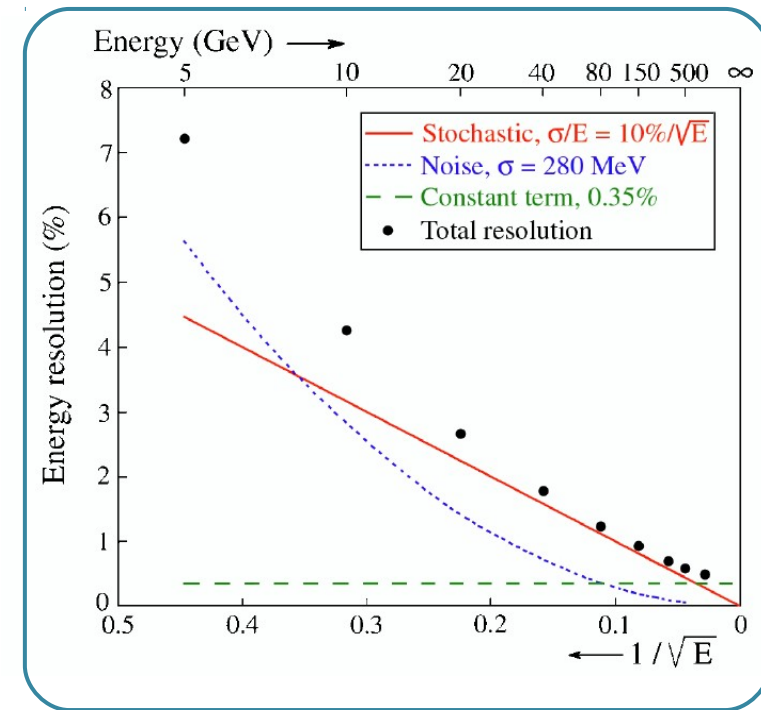
■ Noise term

- Internal (e.g., electronics) and external (e.g. pile-up) noise

■ Constant term

- Fluctuations due to **leakage**
- Imperfections in construction, **non-uniformity**
 - Local variations of temperature, light attenuation, material thicknesses, etc.

Energy resolution in EM ATLAS barrel calorimeter



Homogeneous vs sampling calorimeters

Table 33.8: Resolution of typical electromagnetic calorimeters. E is in GeV.

Technology (Experiment)	Depth	Energy resolution	Date
NaI(Tl) (Crystal Ball)	$20X_0$	$2.7\%/E^{1/4}$	1983
$\text{Bi}_4\text{Ge}_3\text{O}_{12}$ (BGO) (L3)	$22X_0$	$2\%/\sqrt{E} \oplus 0.7\%$	1993
CsI (KTeV)	$27X_0$	$2\%/\sqrt{E} \oplus 0.45\%$	1996
CsI(Tl) (BaBar)	$16\text{--}18X_0$	$2.3\%/E^{1/4} \oplus 1.4\%$	1999
CsI(Tl) (BELLE)	$16X_0$	1.7% for $E_\gamma > 3.5$ GeV	1998
PbWO_4 (PWO) (CMS)	$25X_0$	$3\%/\sqrt{E} \oplus 0.5\% \oplus 0.2/E$	1997
Lead glass (OPAL)	$20.5X_0$	$5\%/\sqrt{E}$	1990
Liquid Kr (NA48)	$27X_0$	$3.2\%/\sqrt{E} \oplus 0.42\% \oplus 0.09/E$	1998
Scintillator/depleted U (ZEUS)	$20\text{--}30X_0$	$18\%/\sqrt{E}$	1988
Scintillator/Pb (CDF)	$18X_0$	$13.5\%/\sqrt{E}$	1988
Scintillator fiber/Pb spaghetti (KLOE)	$15X_0$	$5.7\%/\sqrt{E} \oplus 0.6\%$	1995
Liquid Ar/Pb (NA31)	$27X_0$	$7.5\%/\sqrt{E} \oplus 0.5\% \oplus 0.1/E$	1988
Liquid Ar/Pb (SLD)	$21X_0$	$8\%/\sqrt{E}$	1993
Liquid Ar/Pb (H1)	$20\text{--}30X_0$	$12\%/\sqrt{E} \oplus 1\%$	1998
Liquid Ar/depl. U (DØ)	$20.5X_0$	$16\%/\sqrt{E} \oplus 0.3\% \oplus 0.3/E$	1993
Liquid Ar/Pb accordion (ATLAS)	$25X_0$	$10\%/\sqrt{E} \oplus 0.4\% \oplus 0.3/E$	1996

Homogeneous

Sampling

Fluctuations in hadron showers

- Same types of fluctuations as in EM showers +

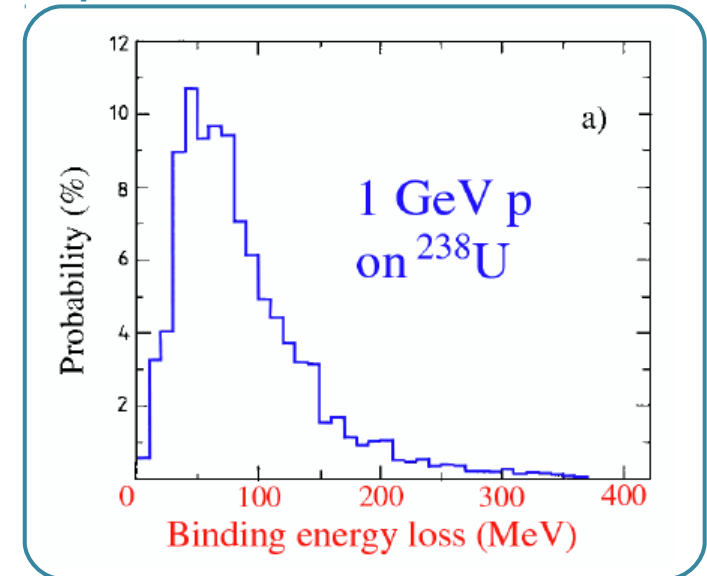
- **Fluctuations in visible energy**

- Fluctuations in losses due to nuclear binding energy
- Note: Correlation with the number of neutrons produced in spallation reactions

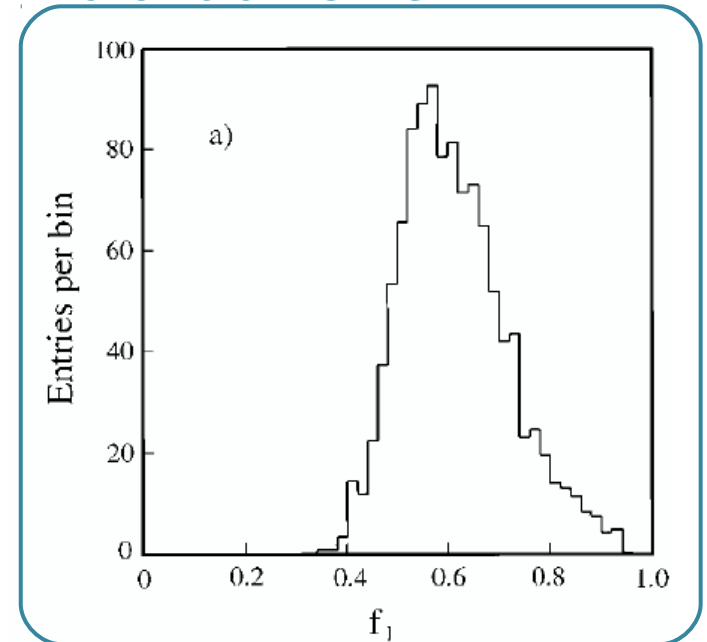
- **Fluctuation in the EM shower fraction**

- Dominating effect in most hadron calorimeters, where $e/h \neq 1$
- Due to the irreversibility of π_0 production \rightarrow asymmetry in EM fraction distribution
- Ideally need to measure the EM fraction for each shower

Binding energy loss for 1 GeV proton in Uranium



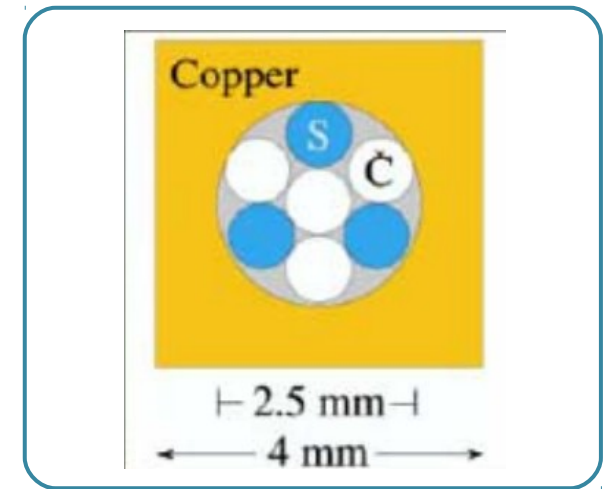
EM fraction of 150 GeV π showers in SPACAL



Measurement of the EM fraction with Dual Readout

- For non-compensating calorimeters
 - Can improve the resolution by **measuring the EM fraction**
 - Largest source of fluctuations
- Can be done with **Dual Readout**
 - DREAM prototype exploring this idea
- Combination of **quartz fibers** and scintillator fibers
 - Quartz fibers only sensitive to EM component
 - Collecting **Cerenkov light emitted by electrons**
- Allows to measure separately the EM and hadronic components of the shower

Conceptual design of a dual readout calorimeter



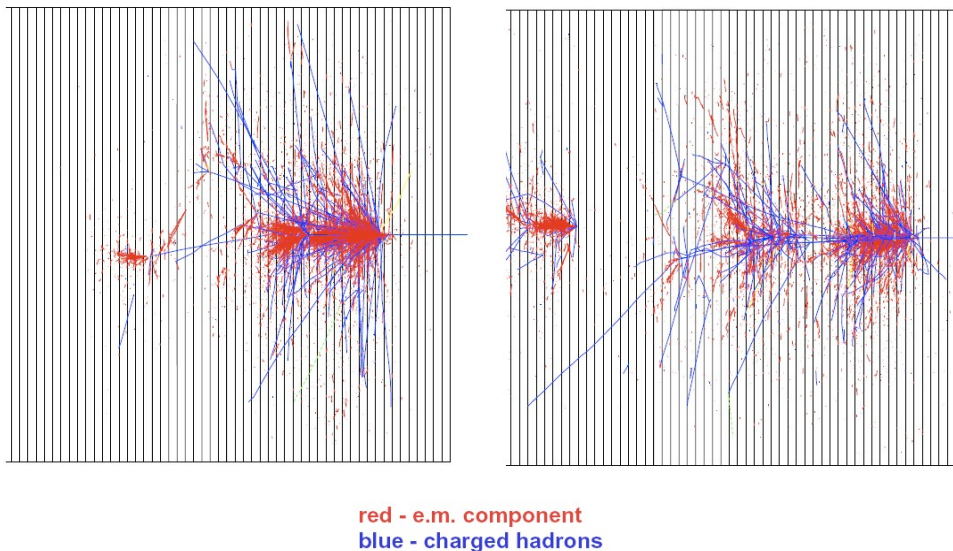
DREAM prototype



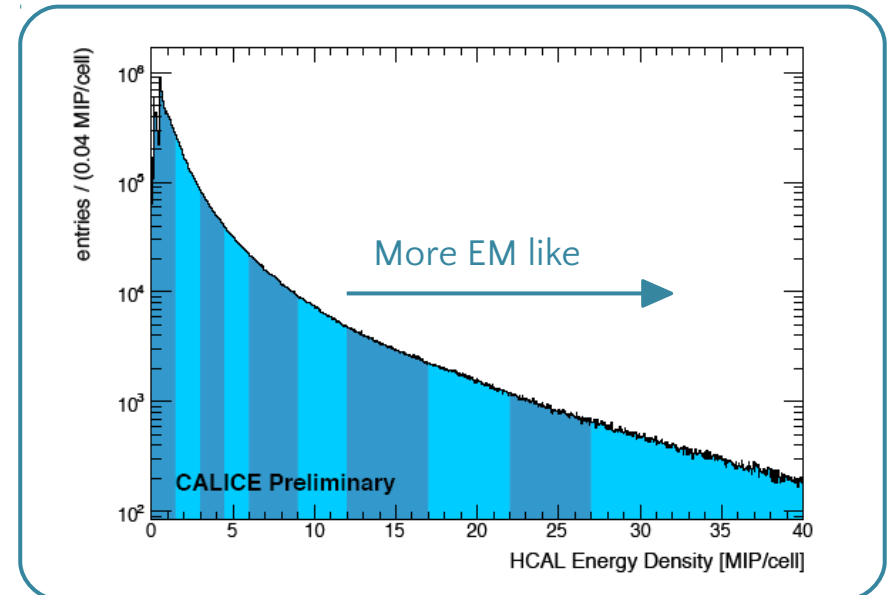
Measurement of the EM fraction indirectly

- One can also **infer** the shower components indirectly
 - And apply different calibrations according to the type of energy deposit (EM or hadronic)
- General idea: EM showers are narrow and dense while hadronic showers are more diffuse
 - → Apply weights according to energy density
- Requires calorimeters with **high-granularity** and **fine segmentation**

EM and hadronic components in showers



Cell energy density



Lecture plan

What is calorimeter in HEP?

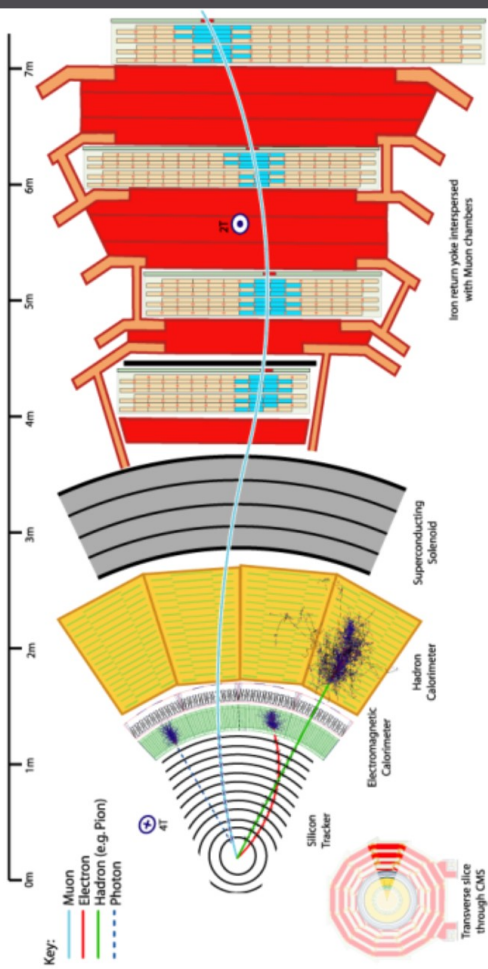
Electromagnetic and hadronic
showers

Detection techniques

Calorimeter
response & resolution

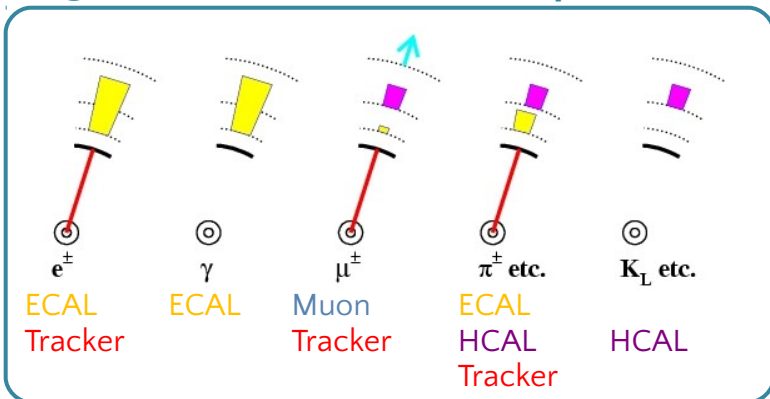
Closing remarks

Closing remarks (1/2)



- In the end a calorimeter is used for
 - Measuring particle energies as precisely as possible (needs a **linear response and a good resolution**)
 - **Identifying particles/showers**, in particular electrons/photons and hadrons
- Many parameters can be optimized
 - Type of calorimeter, material, segmentation, granularity, etc.
- The perfect calorimeter doesn't exist (yet)
 - But one can **combine calorimeter measurements with information from other subdetectors** (e.g. trackers)
 - And make use of their complementary measurements
 - Strategy used in **"Particle Flow"** reconstruction algorithms

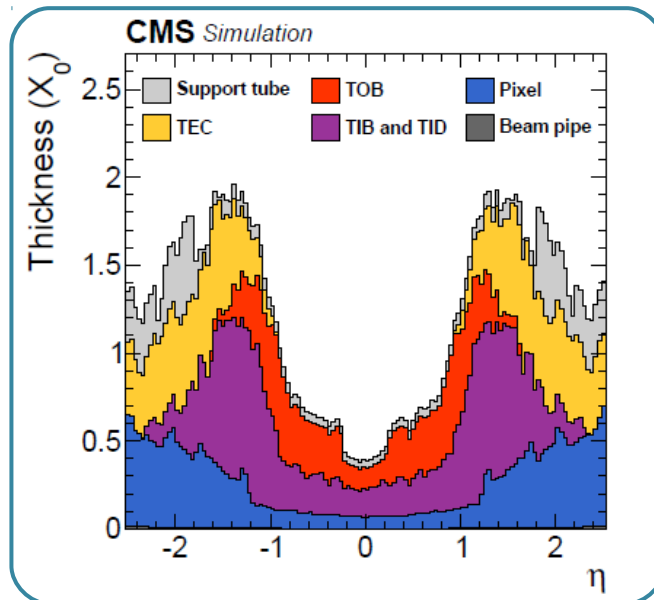
Signatures from different particles



Closing remarks (2/2)

- Real conditions of a full detector in its environment are harsh
 - Very **high number of calorimeter channels** (100k to several millions)
 - **Magnetic field** (impact on photodetectors, electronics, mechanics)
 - **Material** in front of the calorimeter (mechanical structure, other sub-detectors)
 - **Radiations**, Pile-up (in-time and out-of-time)
 - ...
- Degrades performance compared to standalone devices or test beams
- Needs to be taken into account when designing/optimizing a calorimeter

Material in front of the CMS calorimeters



Tracks from multiple simultaneous interactions



Exercises

- In the next two slides are a few exercises related to the resolution of calorimeters
- The solutions can be found on the ESIPAP Indico page

Energy resolution – Sampling term

- We consider a sampling calorimeter using Lead as absorber and Plastic Scintillator as active material, with the following properties:
 - 5 mm thick lead plates
 - 3 mm thick scintillator tiles
 - A resolution of $16\% / \sqrt{E}$ (sampling term)
- Compute the sampling fraction of this calorimeter
- What is the sampling fraction that would be required to get a sampling term of $13\% / \sqrt{E}$?
- Consequently, what is the lead plate thickness that would be required, if we use the same scintillator thickness as before (3mm)?

material	Z	A	ρ	dE/dx	λ_0	X_0	R_M	ϵ
			[g/cm ³]	[MeV/cm]	[cm]	[cm]	[cm]	[MeV]
Al	13	27.0	2.70	4.37	37.2	8.9	4.68	39.3
Liq. Ar	18	40.0	1.40	2.11	80.9	14.0		29.8
Fe	26	55.9	7.87	11.6	17.1	1.76	1.77	20.5
Cu	29	63.5	8.96	12.9	14.8	1.43	1.60	18.7
W	74	183.9	19.3	22.6	10.3	0.35	0.92	7.9
Pb	82	207.2	11.35	12.8	18.5	0.56	1.60	7.2
U	92	238.0	18.95	20.7	12.0	0.32	1.00	6.6
NaI			3.67	4.84	41.3	2.59		12.4
Plastic scintillator			1.032	2.03	68.5	42.9		87.1

Energy resolution – Comparison of two EM calorimeters

- We are comparing the resolutions of the ATLAS and CMS EM calorimeters, as measured in test beams:

CMS

$$\frac{\sigma(E)}{E} = \frac{0.03}{\sqrt{E(\text{GeV})}} \oplus \frac{0.3}{E(\text{GeV})} \oplus 0.005$$

ATLAS

$$\frac{\sigma(E)}{E} = \frac{0.1}{\sqrt{E(\text{GeV})}} \oplus \frac{0.3}{E(\text{GeV})} \oplus 0.007$$

- Fill the following table for both calorimeters. And comment these numbers.

	10 GeV	1 TeV
Stochastic [%]		
Noise [%]		
Constant [%]		
$\sigma(E) / E$ [%]		