

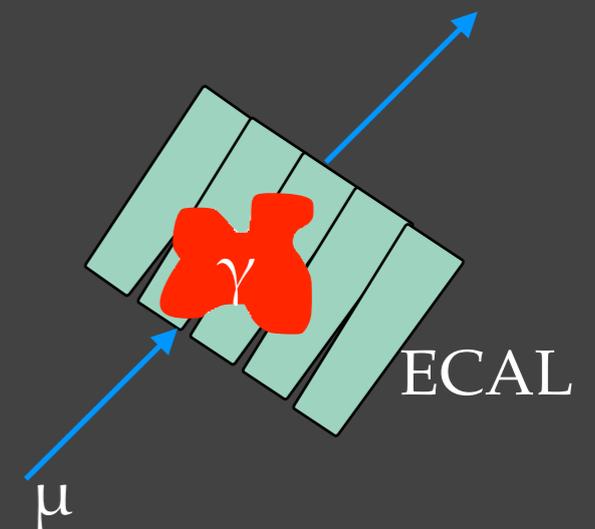
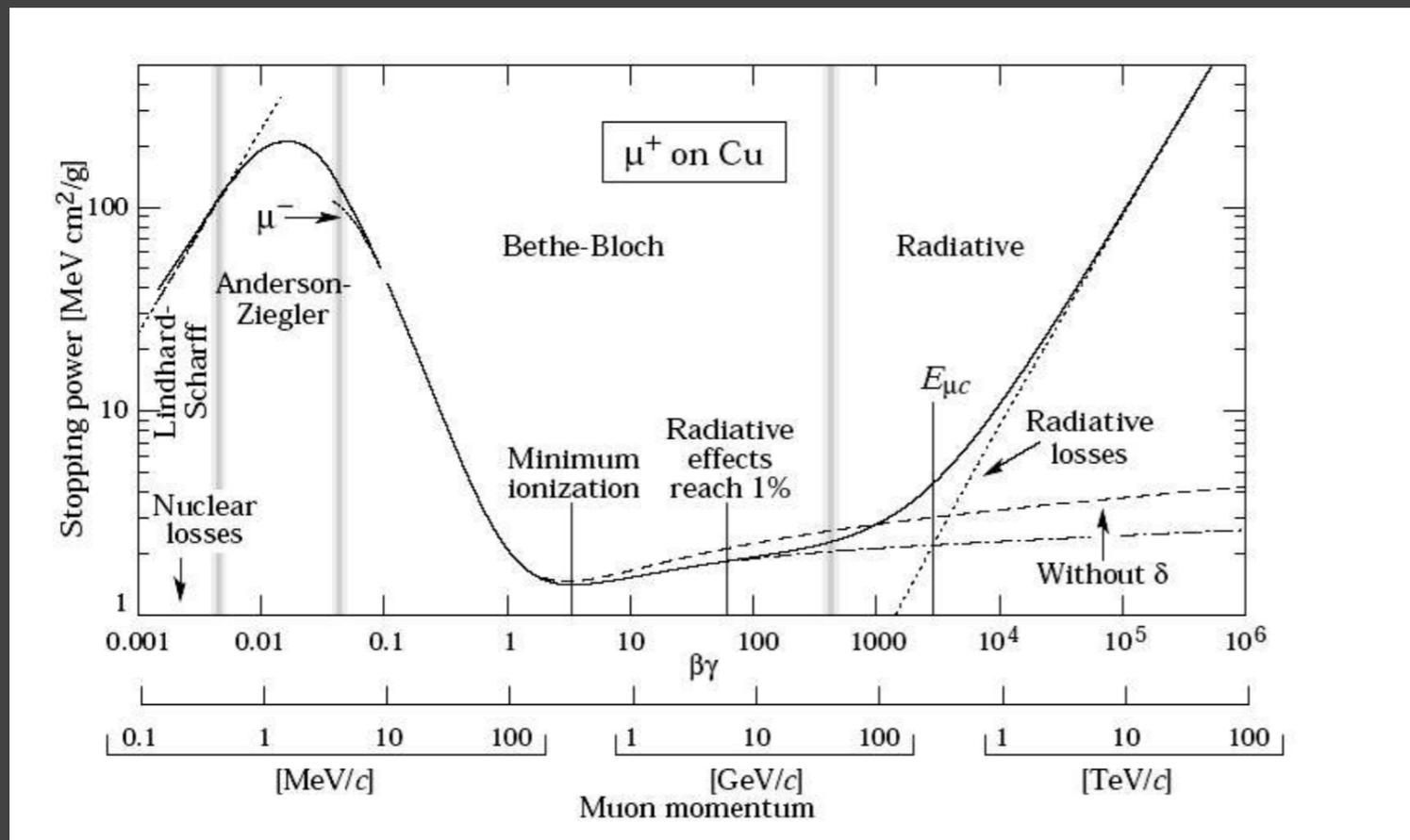
What is beam halo and how do we measure it?

ROGER RUSACK, SHILPI JAIN
UNIVERSITY OF MINNESOTA

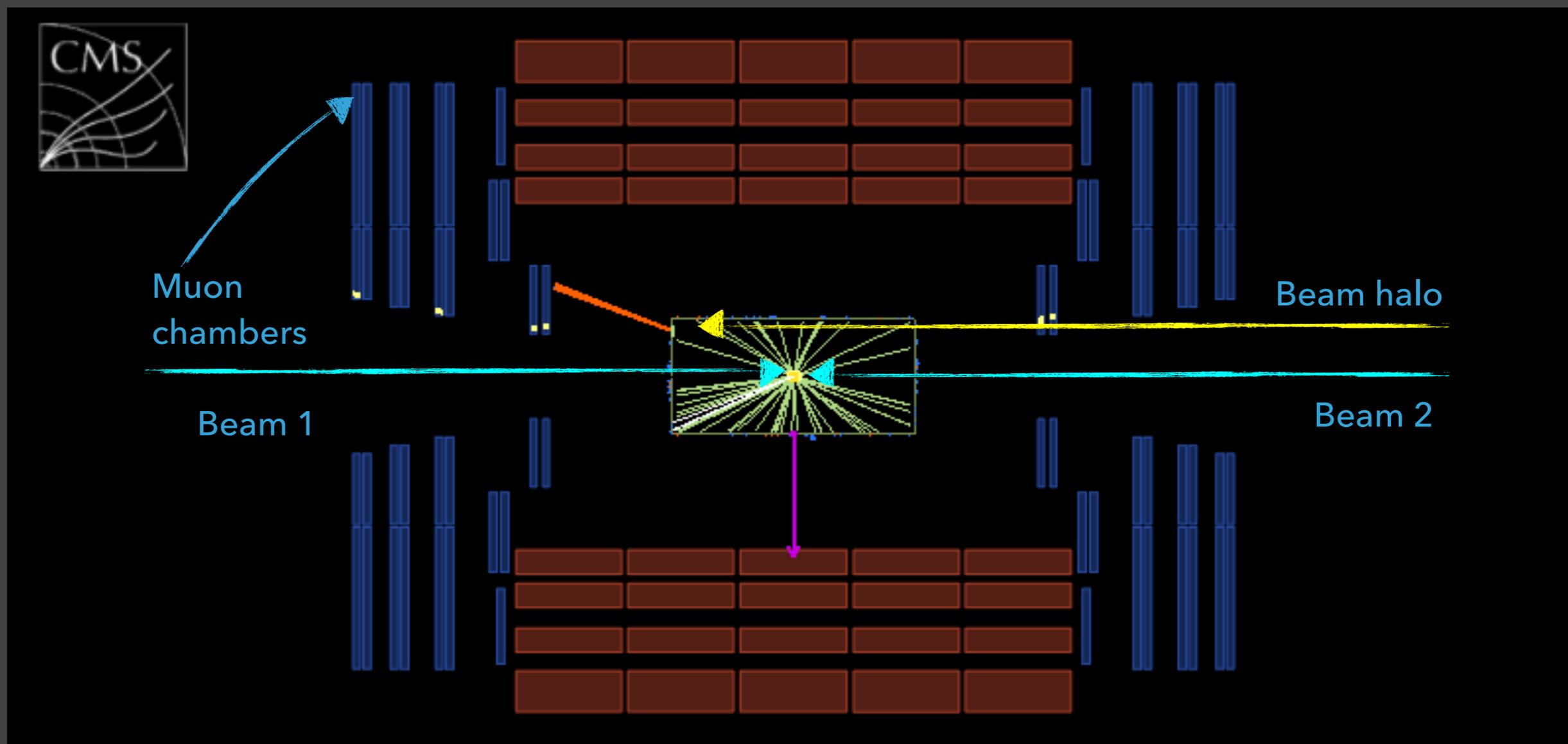
- ▶ There are $\sim 10^{11}$ protons/bunch in LHC beam.
- ▶ The bunches are quite packed - transverse spread is $\sim 15\mu\text{m}$
- ▶ There are inter-particle interactions, beam-gas interactions and synchrotron radiation
- ▶ All of these processes can make protons diffuse to a higher radius
- ▶ Dedicated collimator system is in place in the LHC system to scrape these particles off.
 - ▶ Still some particles escape and form 'beam halo'

Machine induced background

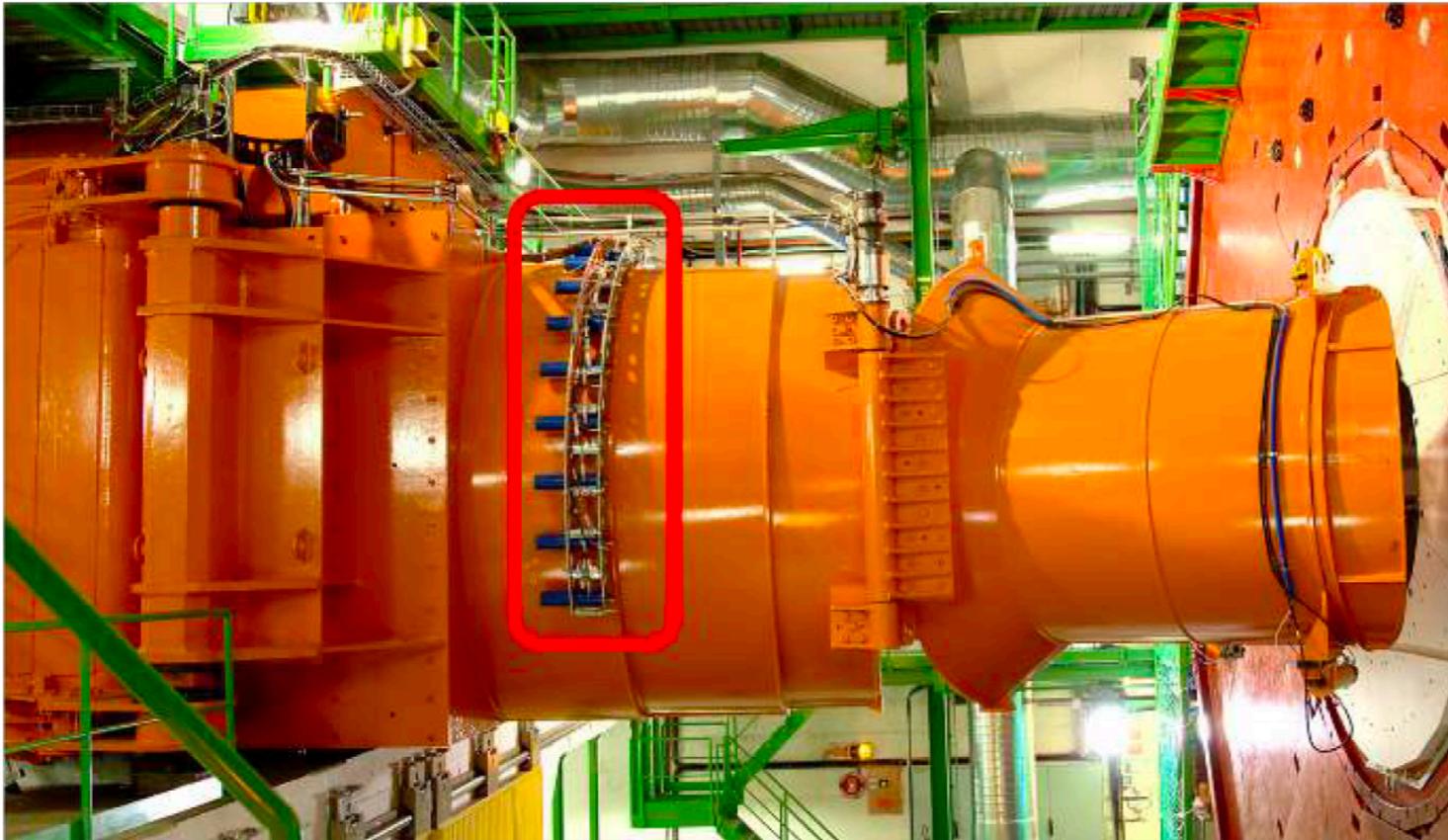
- ▶ Due to collision of these protons with the LHC collimator system, shower particles can be produced
- ▶ These particles travel **parallel** to the beam
- ▶ The ones which reach and are important for the CMS detector are muons
- ▶ These highly energetic muons can interact with the sub-detector giving rise to a high energy deposit in the calorimeter



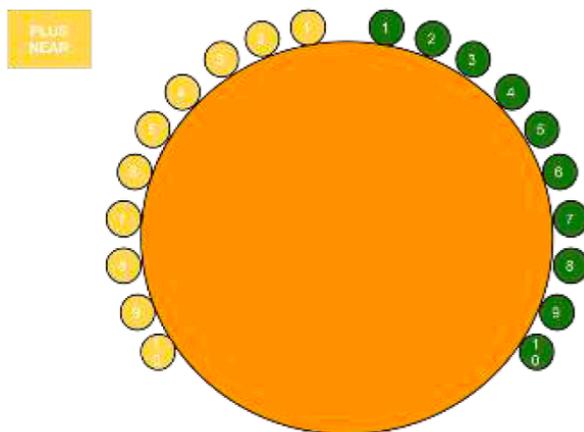
- ▶ A muon with $p_T > 100$ GeV becomes radiative and can do bremsstrahlung dominantly



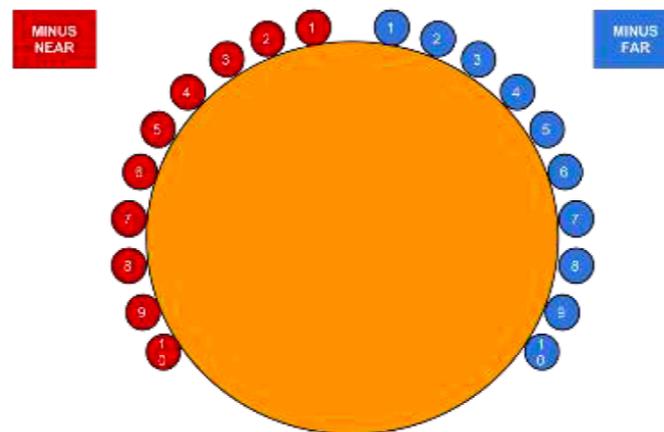
- ▶ A large deposit in ECAL can be seen in line with hits in the muon chambers - typical of beam halo
- ▶ Beam halo is a potential background to many analyses in CMS



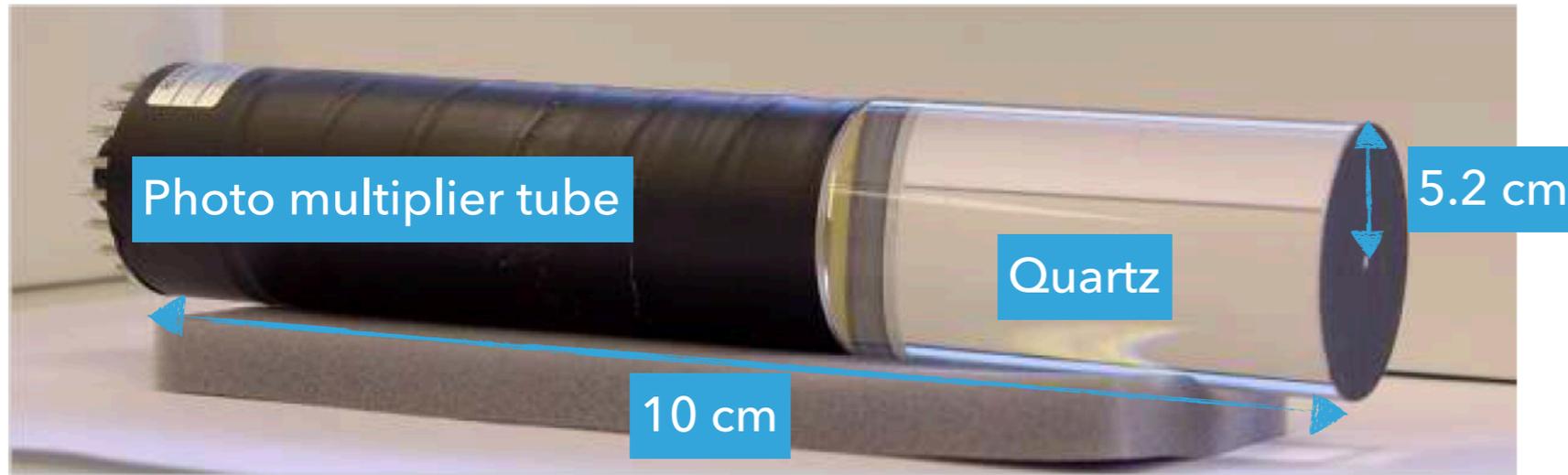
- ▶ To detect such a background, there are beam halo monitors
- ▶ Mounted on rotating shielding
- ▶ 20 on each side of the CMS detector



(a) +Z distribution



(b) -Z distribution

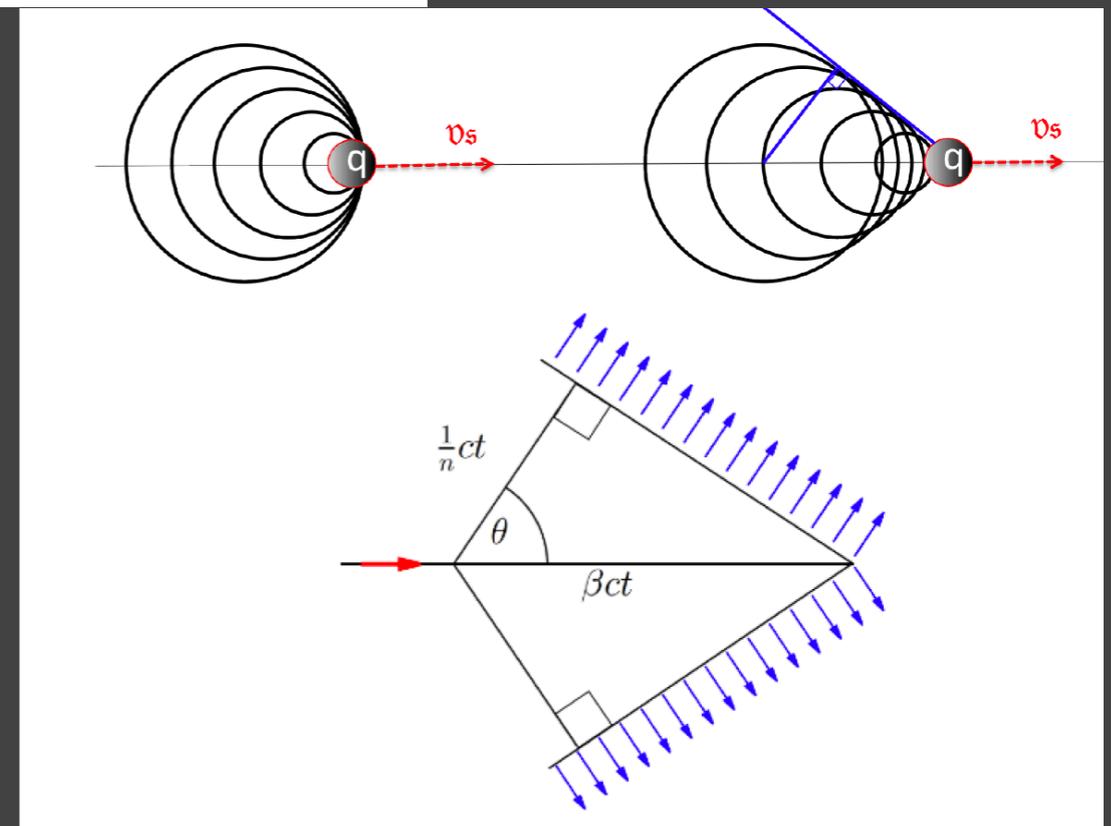


- ▶ Best suited for this purpose: cherenkov light based detector
 - ▶ Suitable material: Fused silica (quartz)
 - ▶ Good transmission in the UV range
- ▶ Photodetector choice: Photo multiplier tube
 - ▶ Sensitive to UV down to 160 nm
 - ▶ High gain and fast rise time

angle at which light is emitted

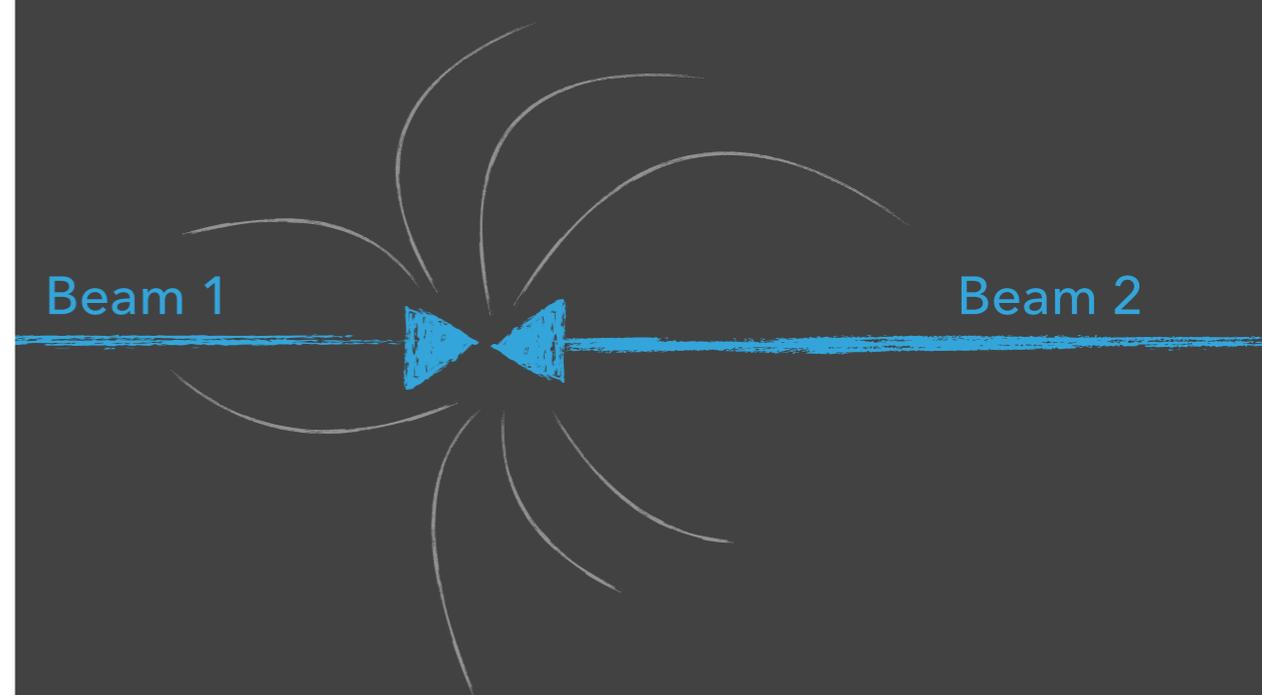
$$\cos \theta = \frac{c}{n\beta c} = \frac{1}{\beta n}$$

n = refractive index
 $\beta = v/c$; v = velocity of the particle

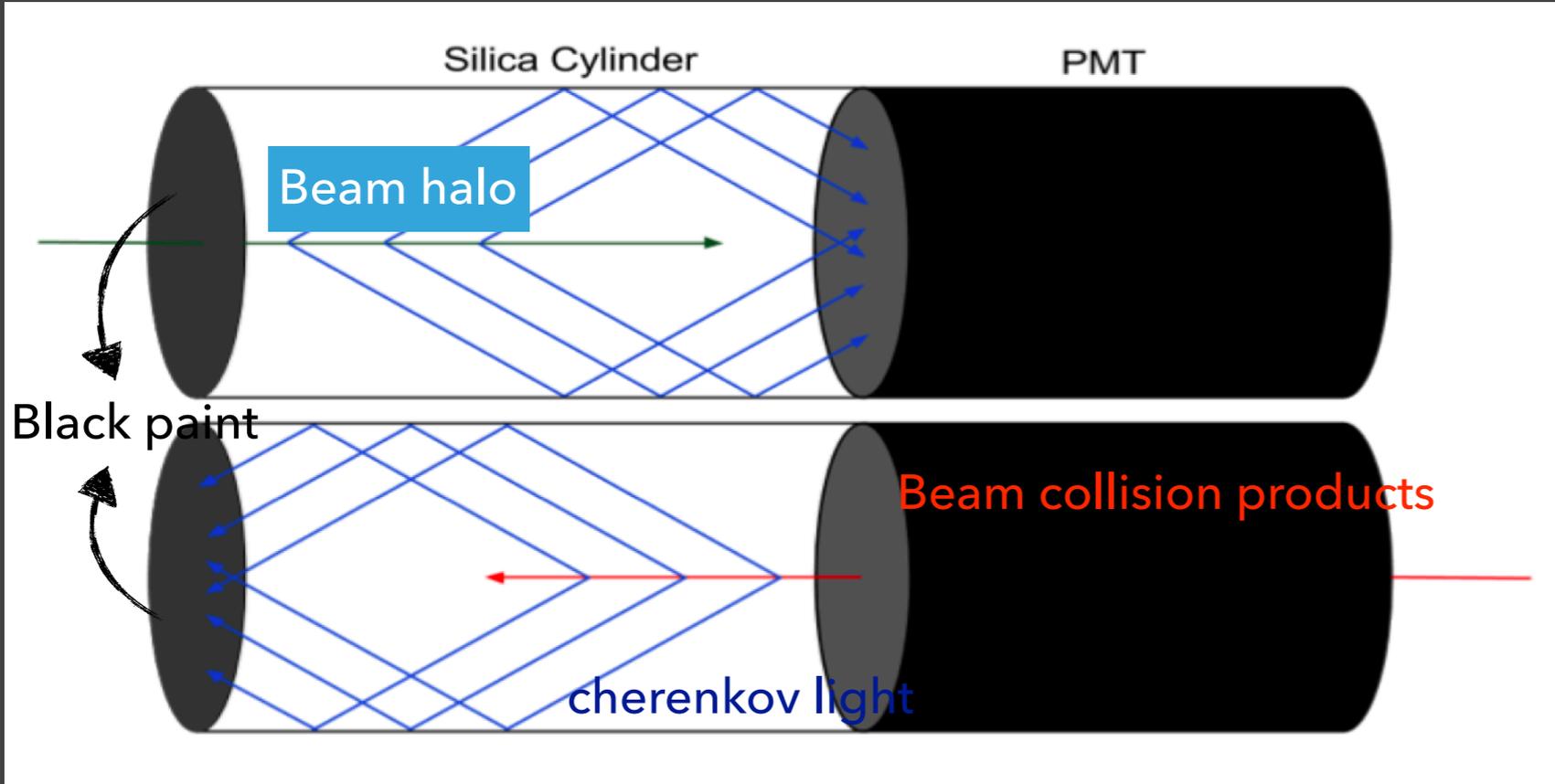


Cherenkov radiation is emitted when particle moves faster than light in that medium - it has directional nature

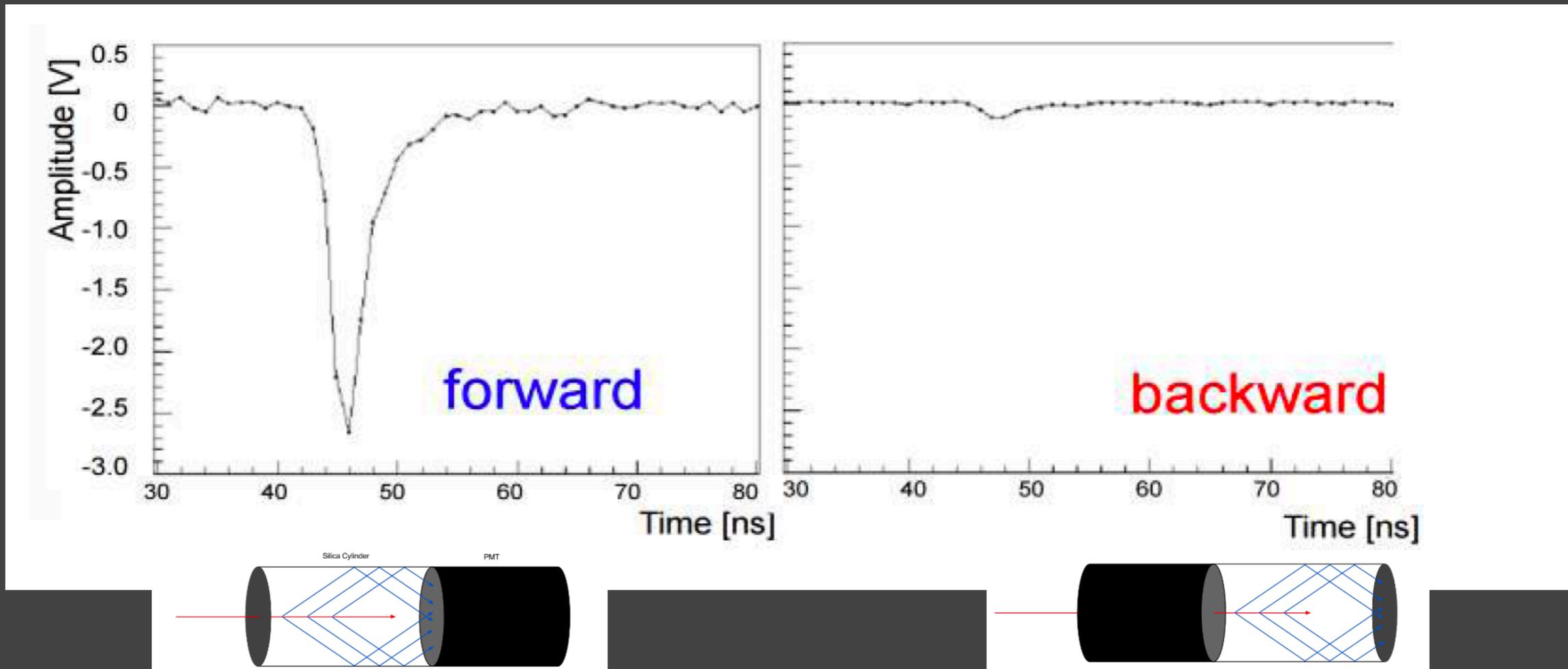
Basic principle



- ▶ Directional property is exploited:
 - ▶ Signal is generated when beam halo crosses the detector
 - ▶ The collision products do not cause any appreciable signal



Signal for forward and backward traveling particles

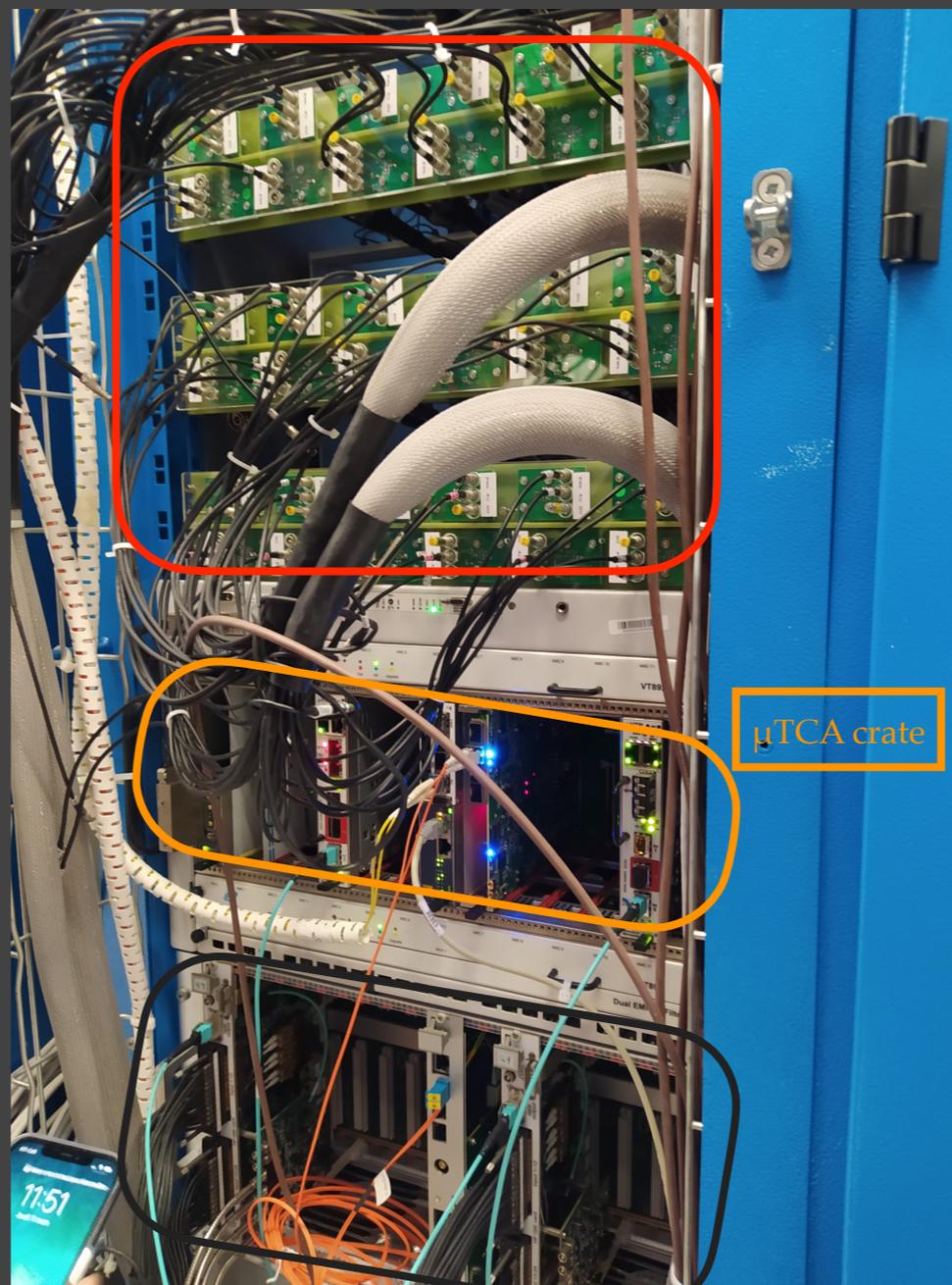


- ▶ Test beam at CERN in 2012 using muons

Readout system in USC (service cavern)



Signal cables (brown)



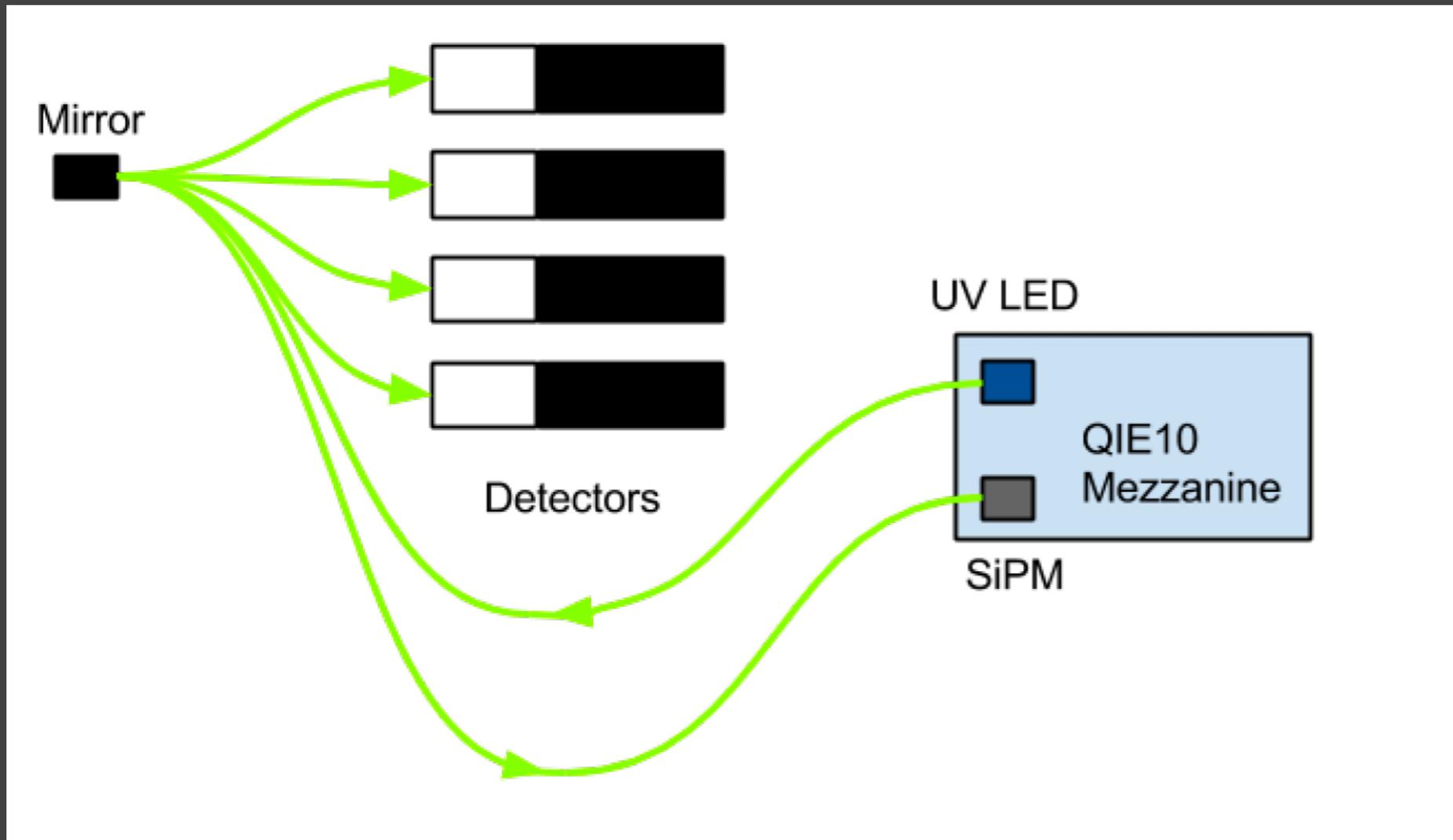
Patch panel

μTCA crate

QIE +
calibration

SF108 rack

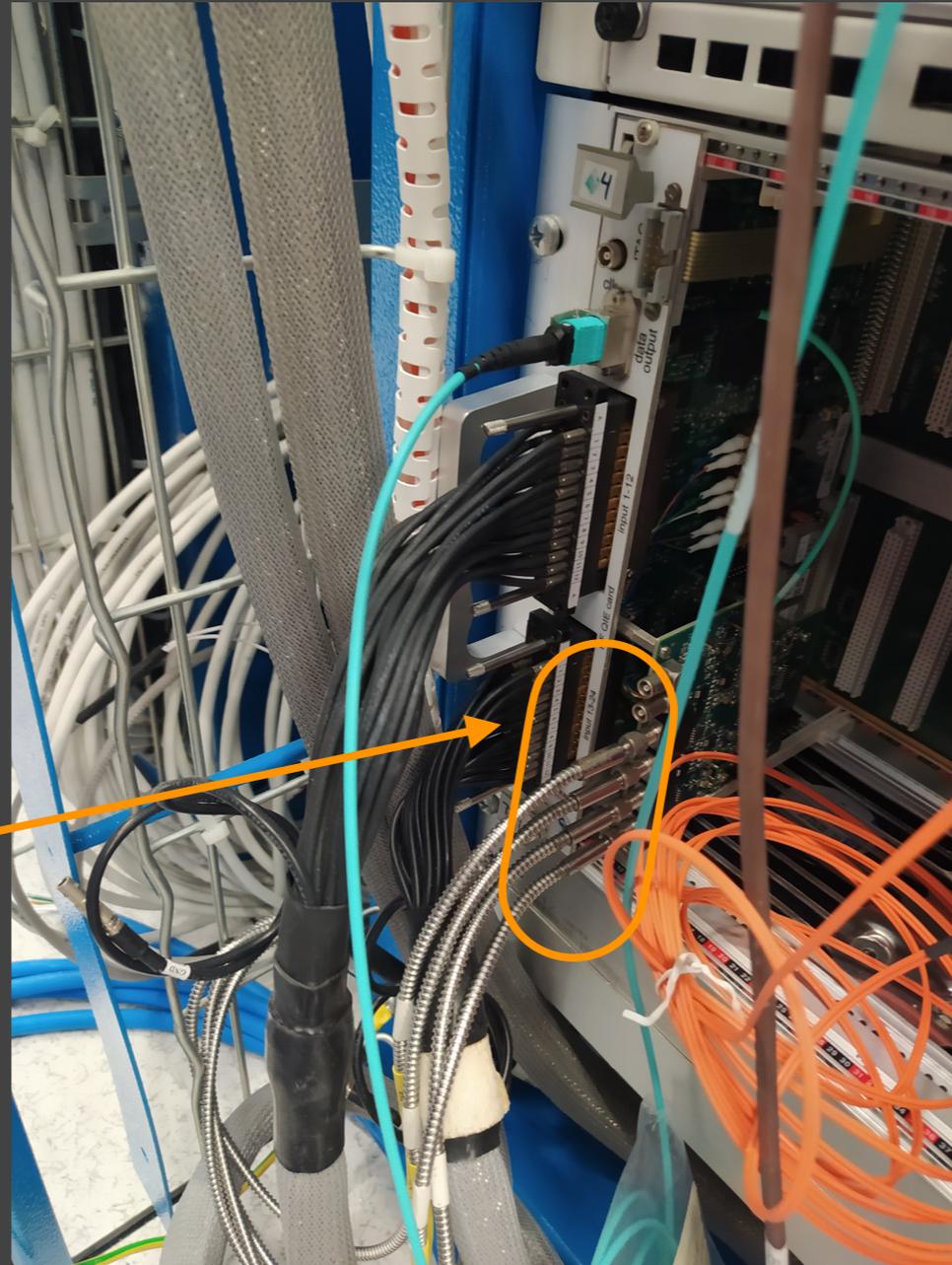
QIE: Integrates and digitizes. It has TDC as well for timing information



- ▶ A pulser circuit can be triggered in empty bunch crossings at fixed intervals to track the performance of the PMTs

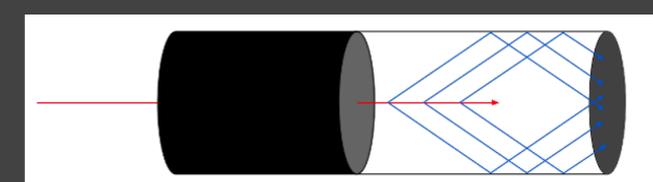
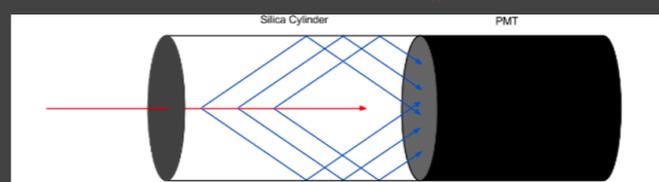
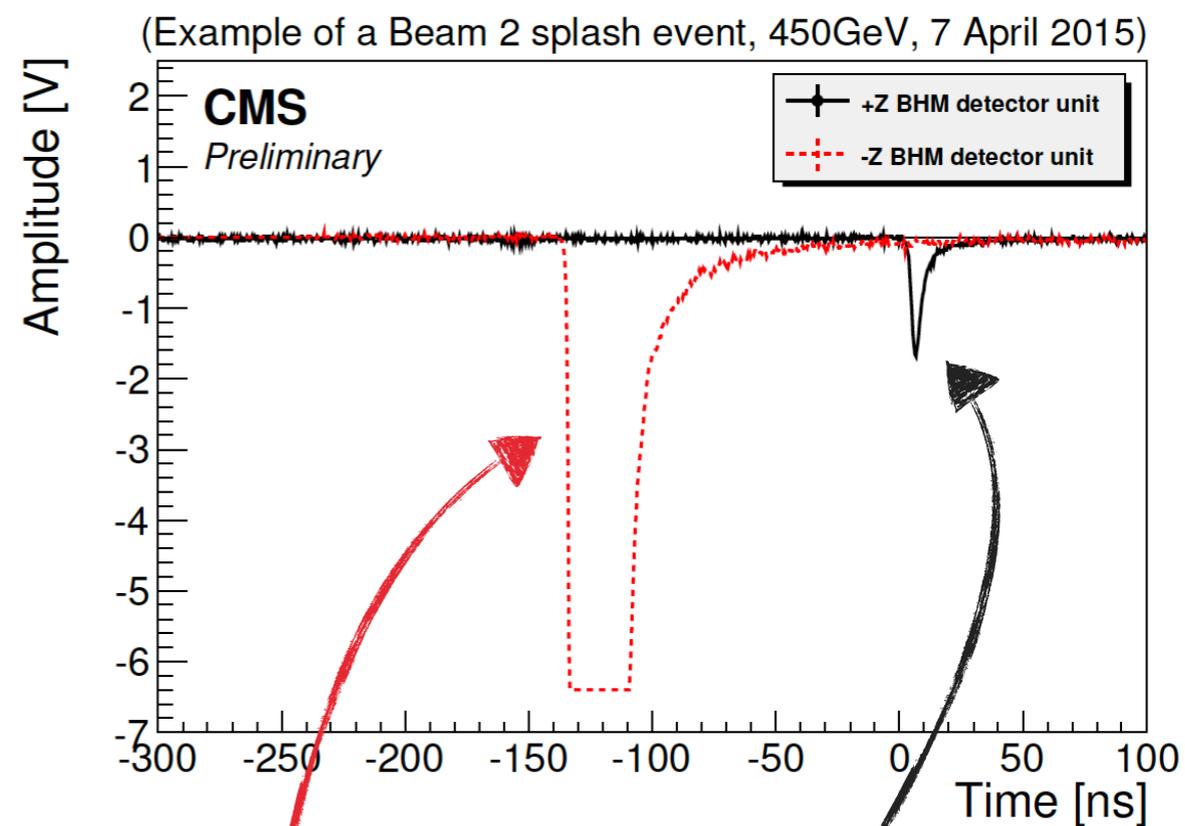
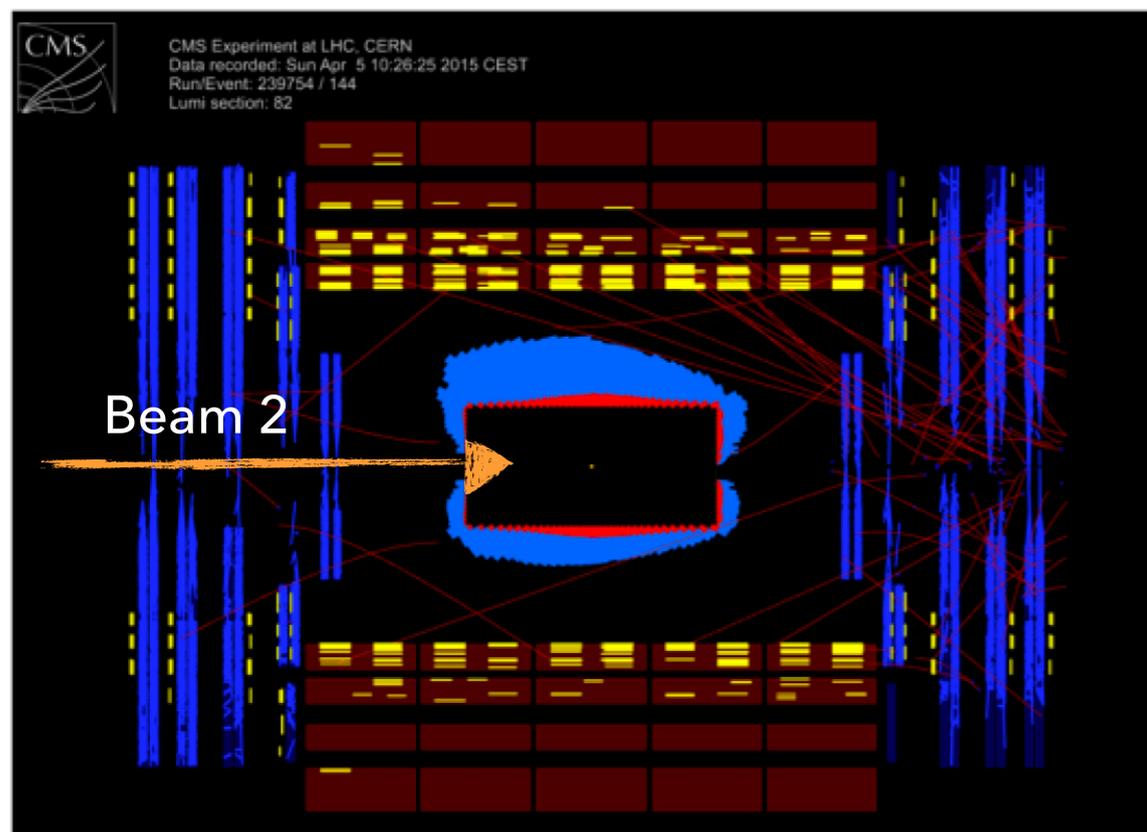


Main crate - SF108



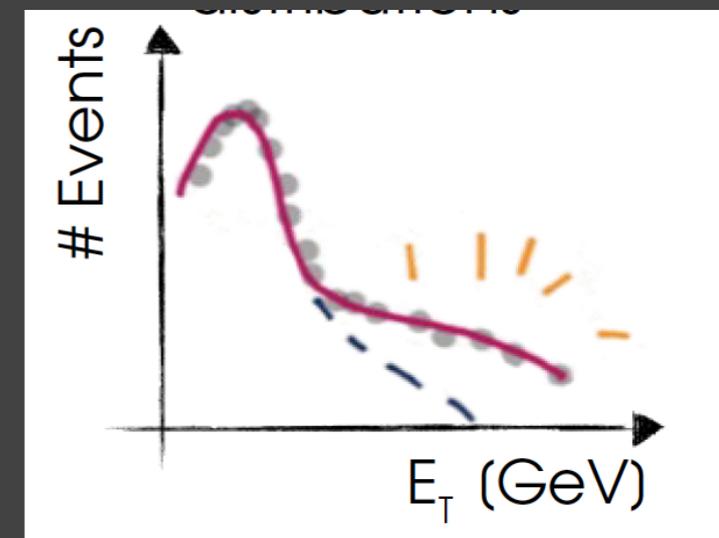
Optical fibers

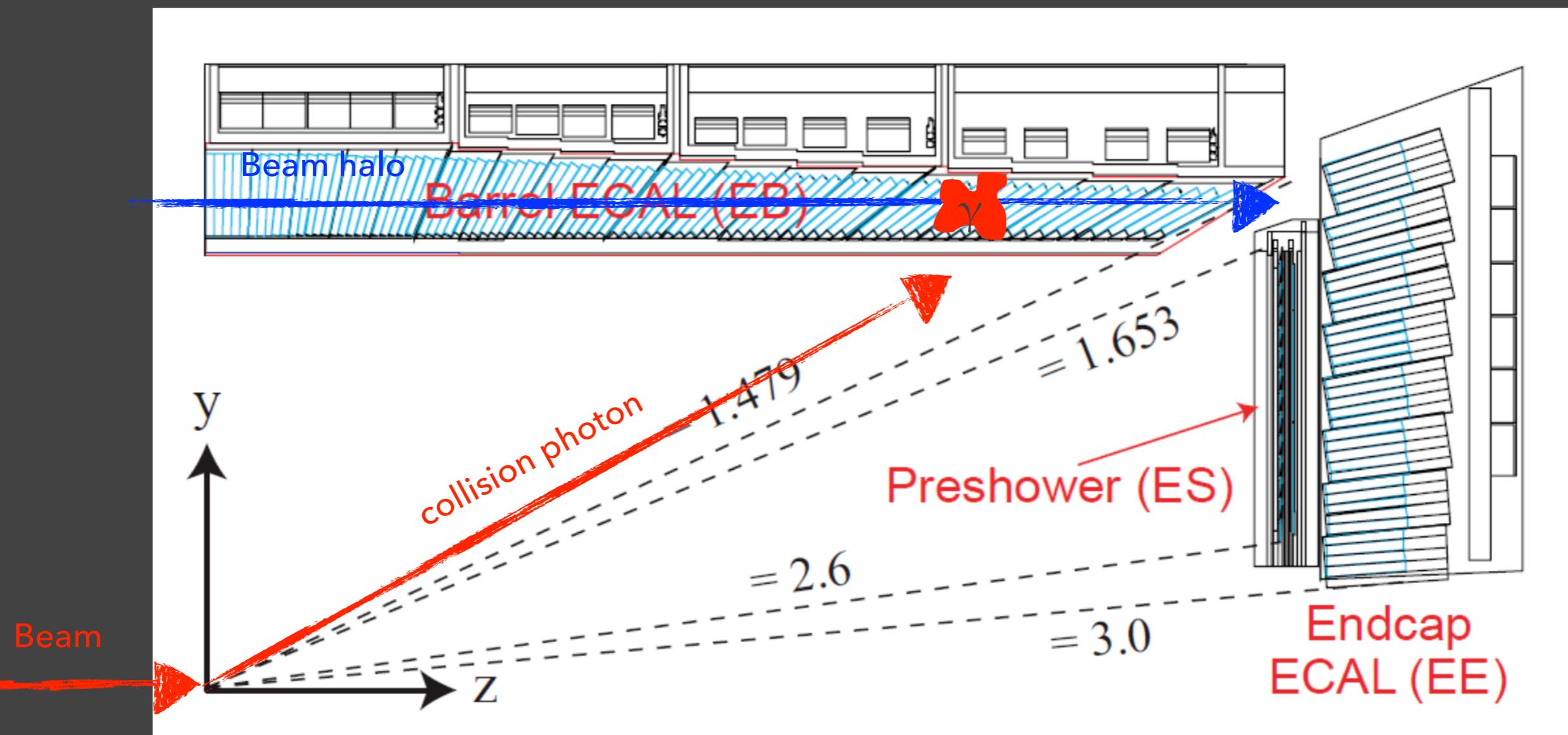
Signals during beam splash in 2015



- ▶ During 2015, during the splash event (i.e. one of the beams directed to the collimators), output from BHM's were recorded.
- ▶ As can be seen, the +Z BHM detector unit does not give any significant signal, where the -Z gives
 - ▶ Demonstration of directional nature of the detector

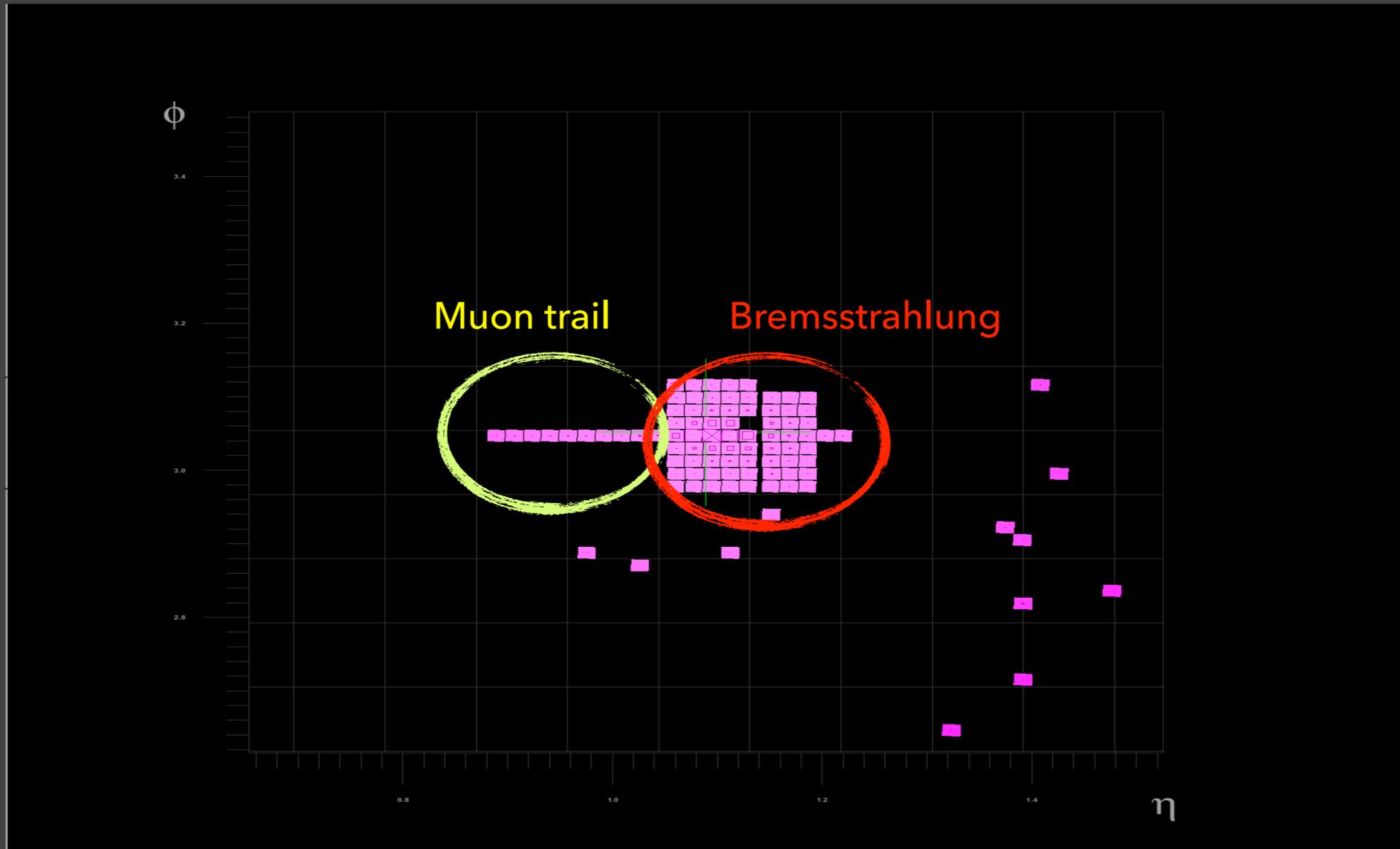
- ▶ The beam halo monitors integrate information over ~ 1.4 s (1 bunch crossing is 25 ns)
 - ▶ In case of any excess of events, the data can be divided into periods of low beam halo and high beam halo
 - ▶ However, with this system, currently there is no per event information available
- ▶ We can still use the CMS calorimeter to build up a per-event decision about the possibility of a beam halo
 - ▶ Different ways to tackle in EB and EE



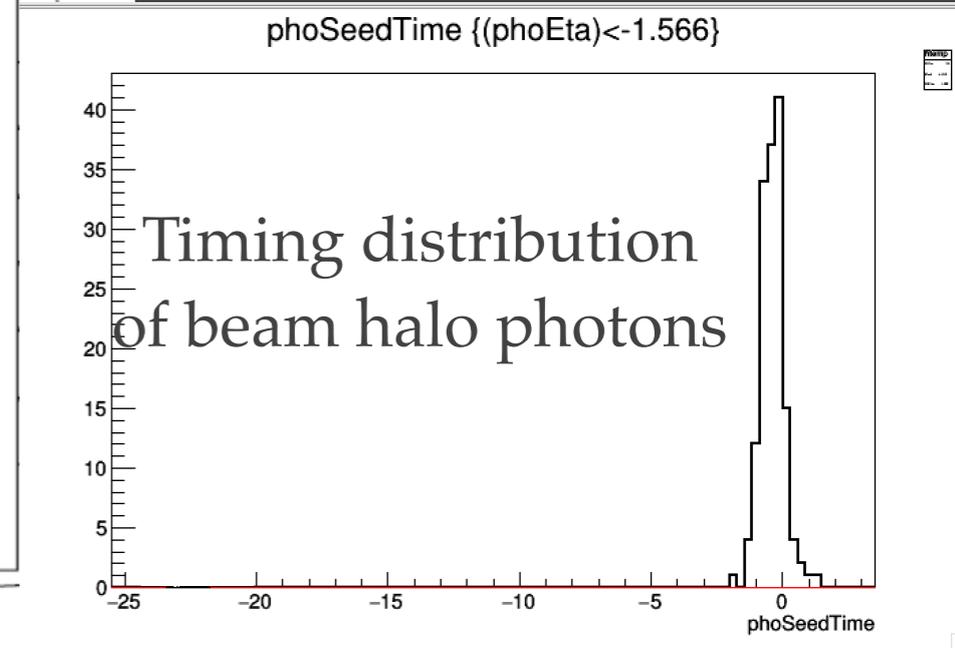
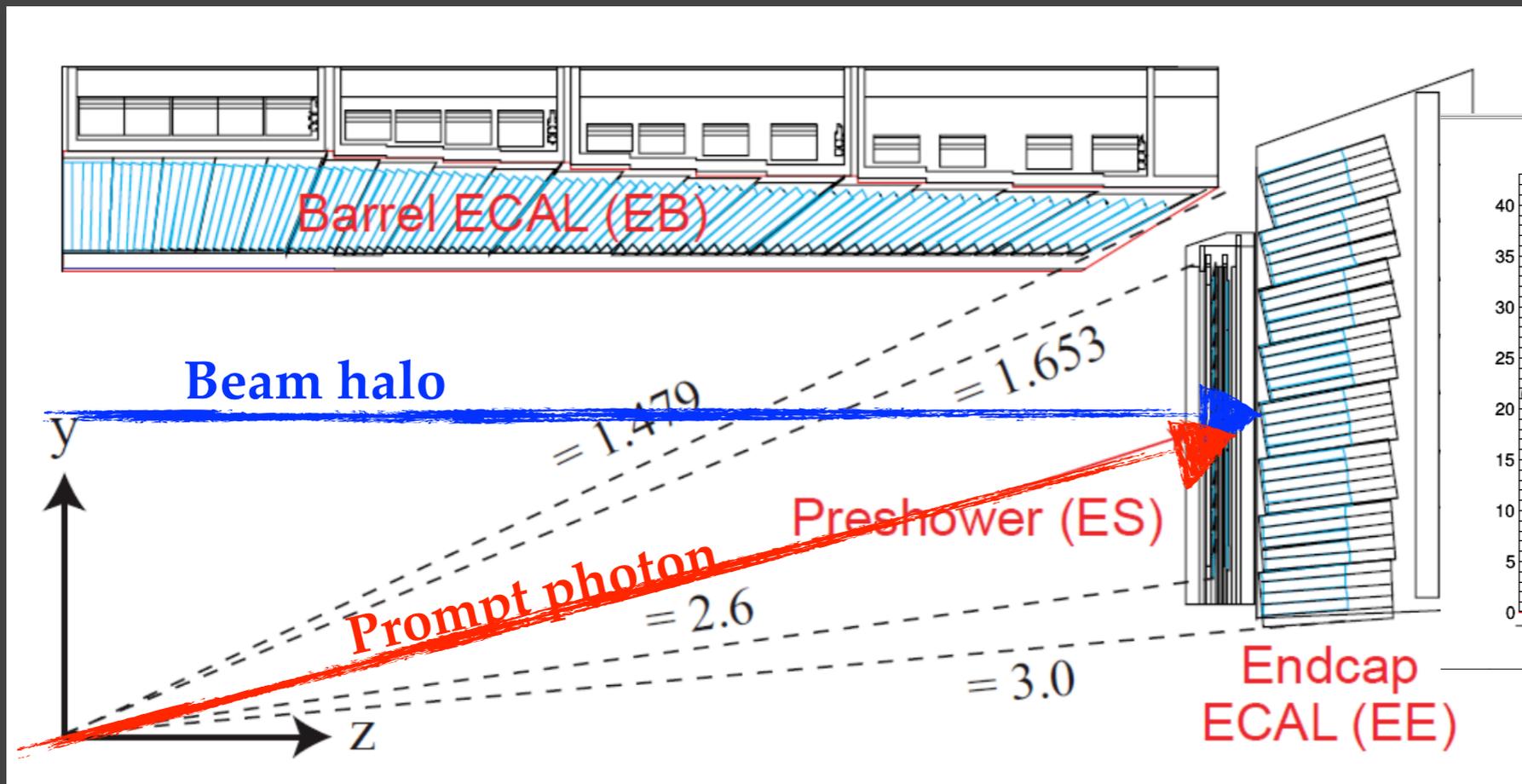


- ▶ The beam halo in the EB arrives earlier than the collision, so timing can be used to handle the background

Inside CMS ECAL Barrel and its removal at the analysis level¹⁵

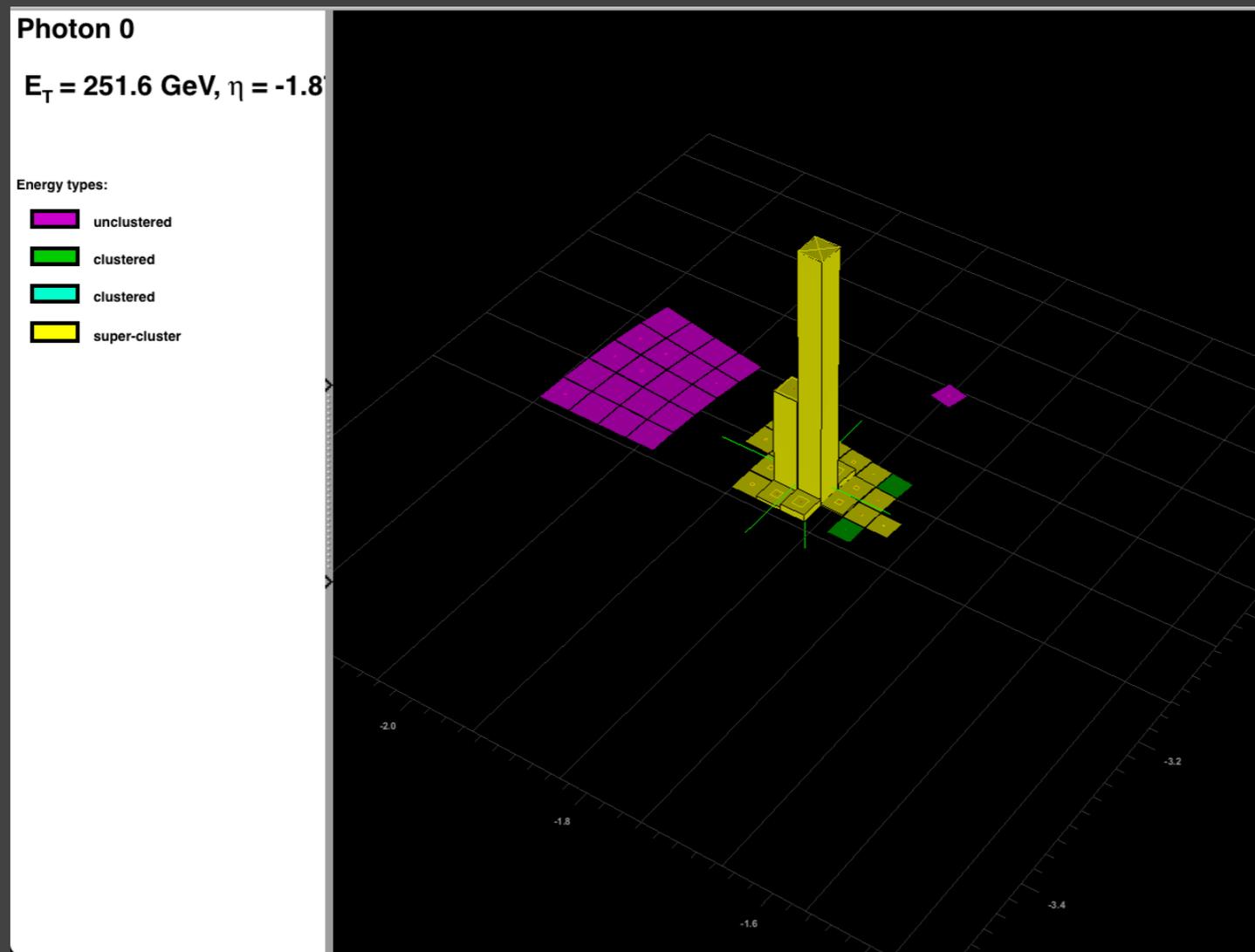


- ▶ By collecting the hits from the muon trail and their total energy, this background can be handled



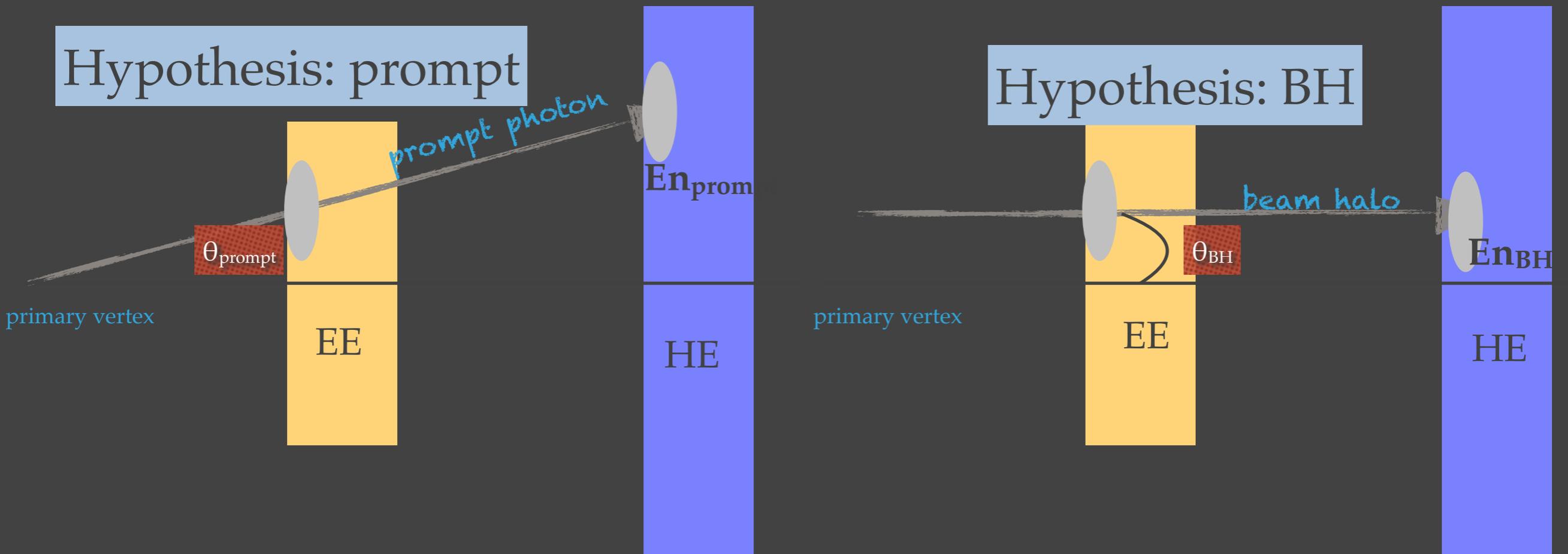
▶ Timing cut:

- ▶ A beam halo coming at a distance of 129./2 will hit at ~0.2 ns earlier compared to a prompt photon (EE is at a distance of 317 cm)
- ▶ Timing distribution of the beam halo photons fall within 3 ns - very similar to that of prompt photons
- ▶ Removes those that interact on the same side as the source, but cannot be used to identify those that interact on the opposite side

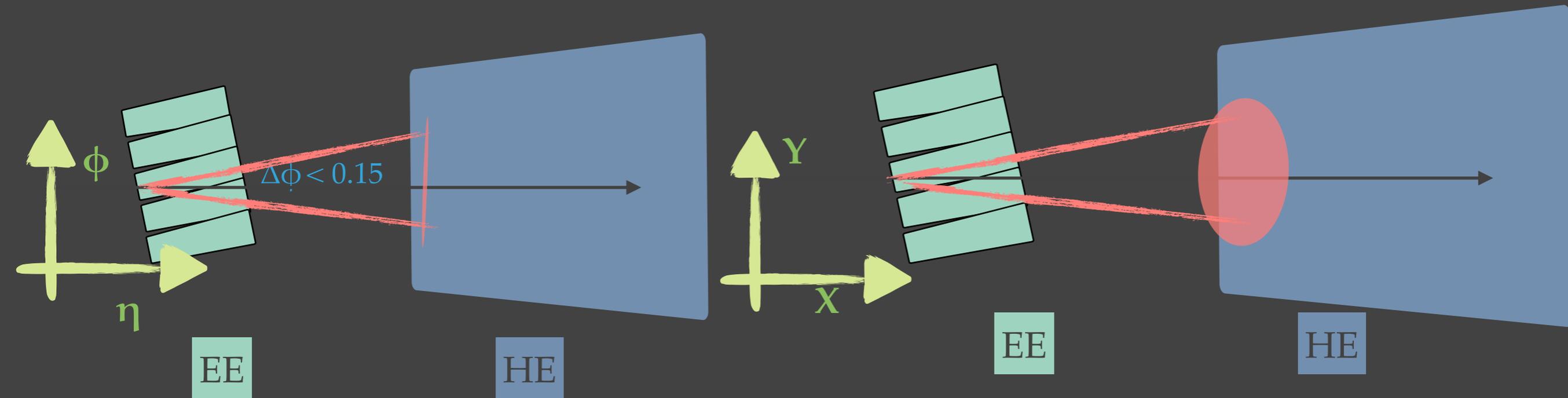


- ▶ No particular muon trail as seen in the EB

- ▶ To tackle this background in the EE, we developed an idea of using the information from various sub-detectors :
 - ▶ ECAL endcap, HCAL endcap and preshower

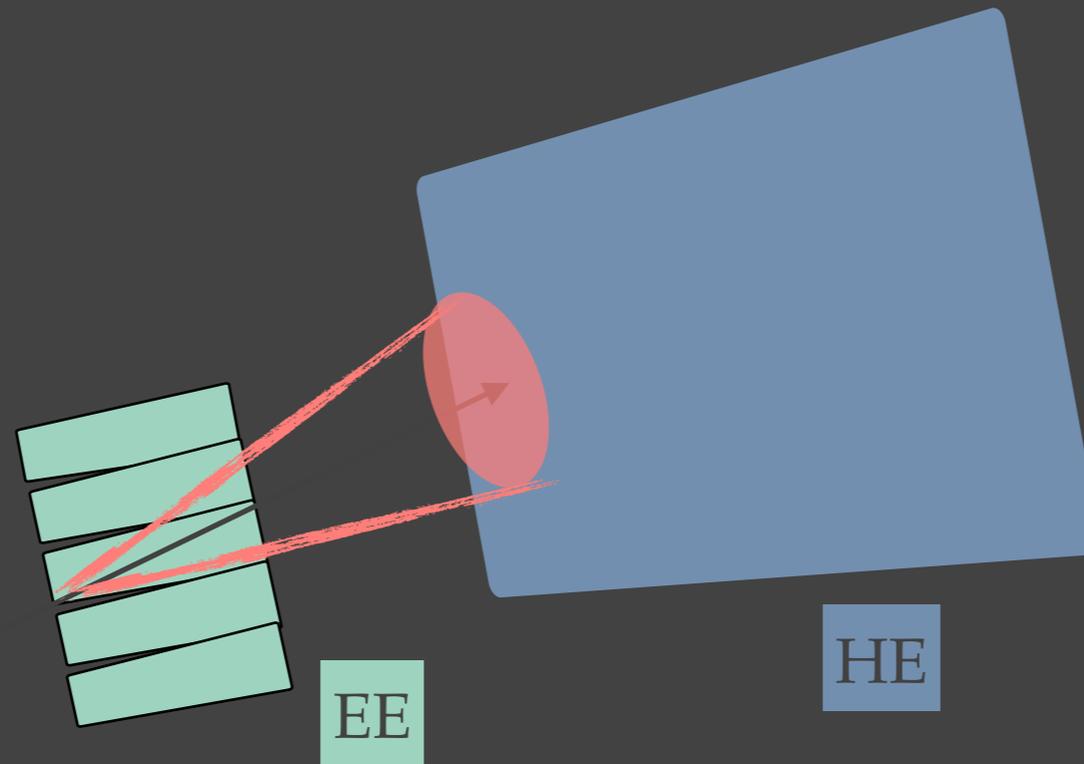


- ▶ A photon in the EE can either be a prompt photon (i.e. coming from the collision vertex) or a beam halo (i.e. traveling parallel to the z axis)
- ▶ Thus two hypothesis are formed:
 - ▶ prompt
 - ▶ BH (beam halo)



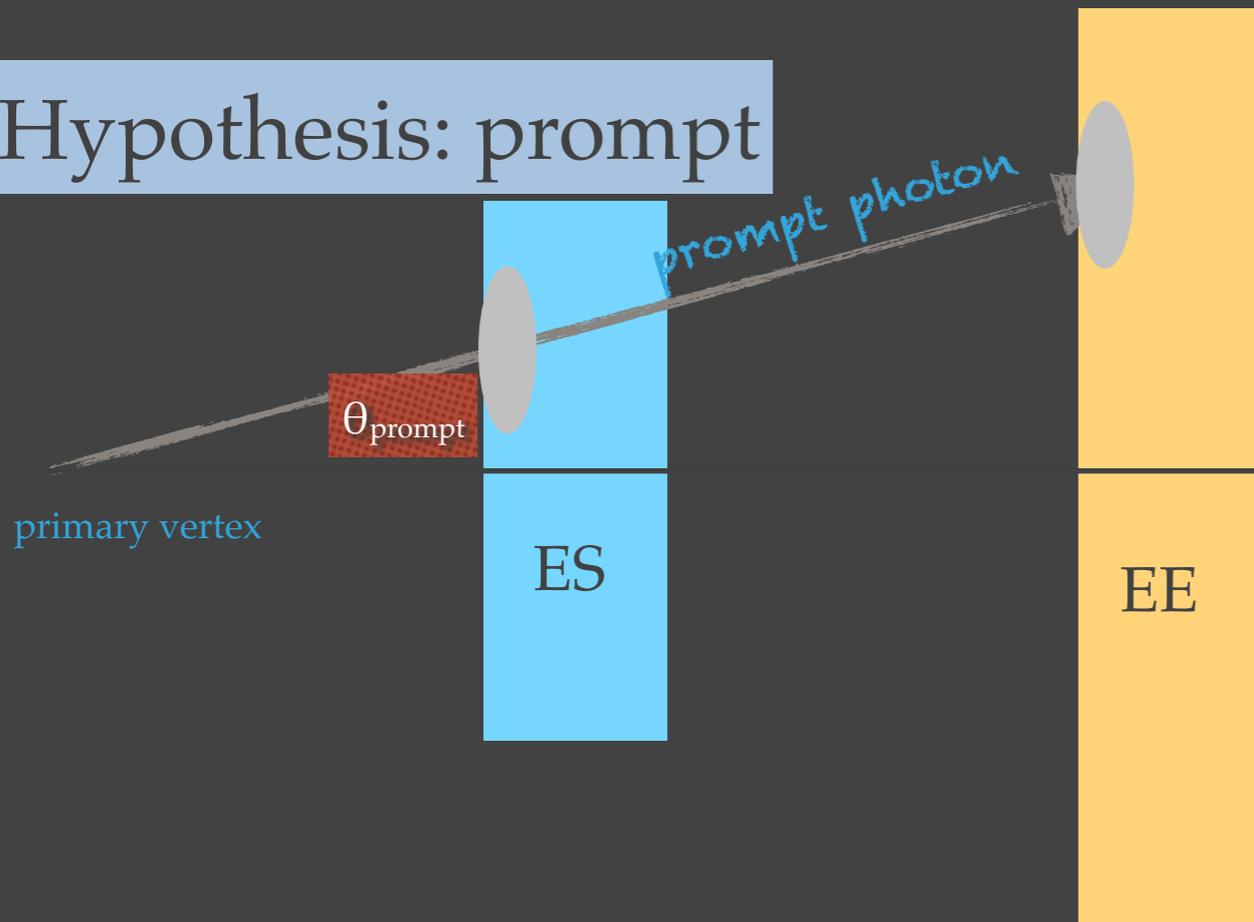
- ▶ Idea is to collect the energy in the HE behind the photon in both the hypothesis.
- ▶ In addition an angle is formed between the line joining the cluster in the EE and the HE
- ▶ In both the hypothesis, a cone is open and hits are collected
 - ▶ For BH hypothesis, the cone is formed in $\Delta\phi$ (since the BH travels parallel)

Angle and energy estimation in 'prompt photon' hypothesis 21

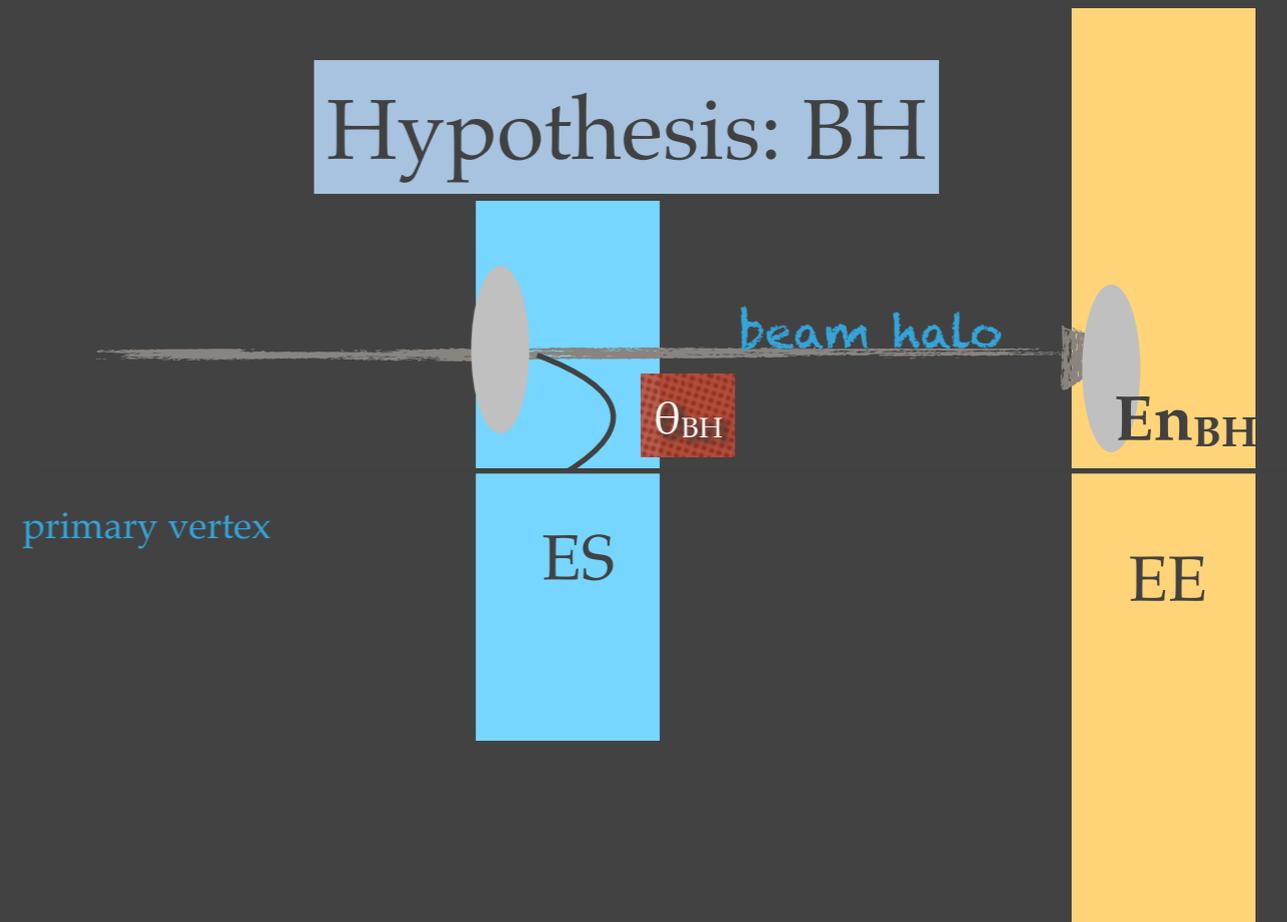


- ▶ Idea is to collect the energy in the HE behind the photon in both the hypothesis
- ▶ In addition an angle is formed between the line joining the cluster in the EE and the HE
- ▶ In both the hypothesis, a cone is open and hits are collected
 - ▶ For BH hypothesis, the cone is formed in $\Delta\varphi$ (since the BH travels parallel)
 - ▶ In case of prompt hypothesis, a cone is opened in $dR = \sqrt{\Delta\eta^2 + \Delta\varphi^2}$

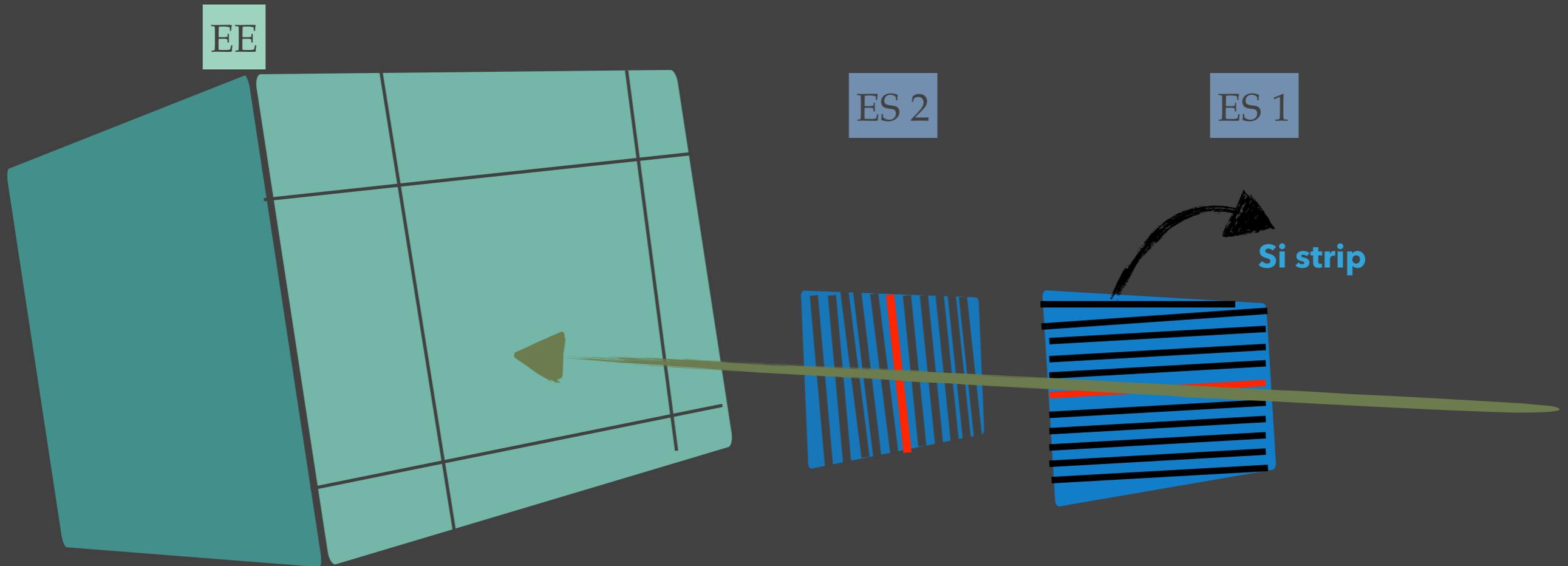
Hypothesis: prompt



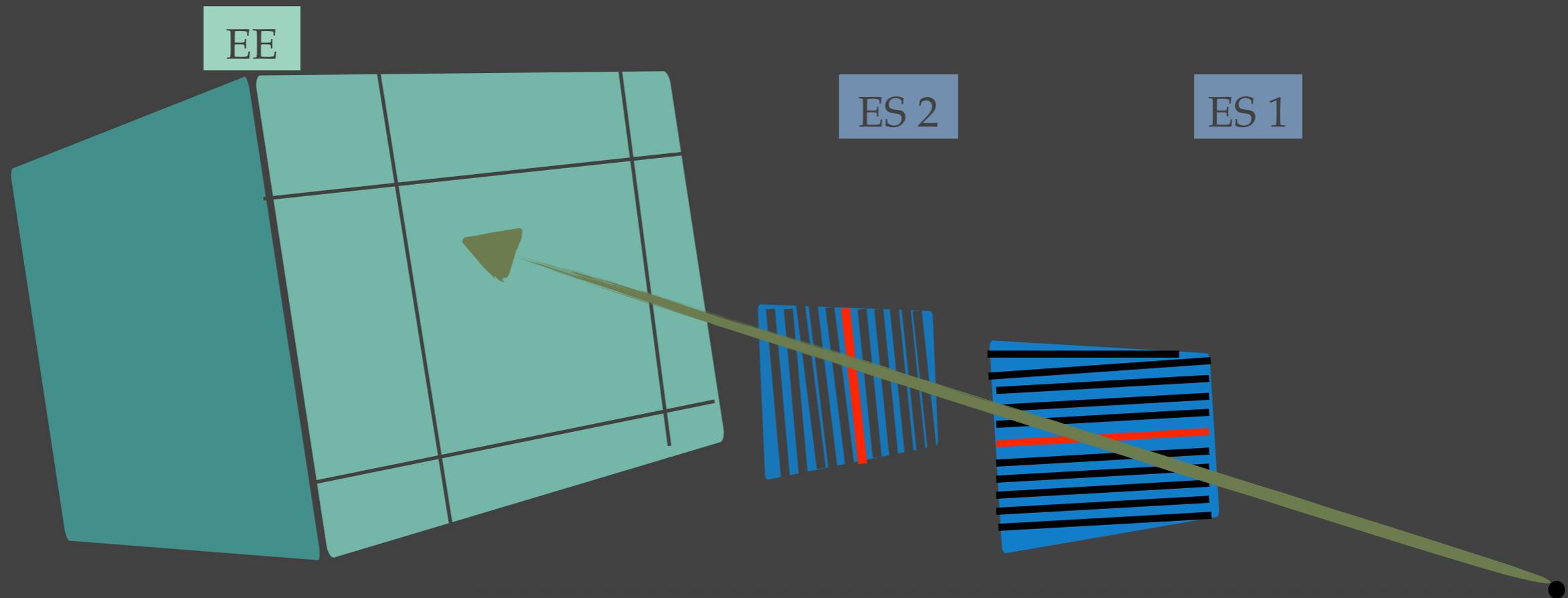
Hypothesis: BH



- ▶ Exactly the same idea of two hypothesis between the preshower and the EE



- ▶ Extrapolate from the position of the photon in the EE to that in the ES assuming the photon traveled parallel to the beam axis
- ▶ Collect the hits in the preshower planes
- ▶ Again form the angle between the line joining the position in the ES and the EE and collect the energy in the ES



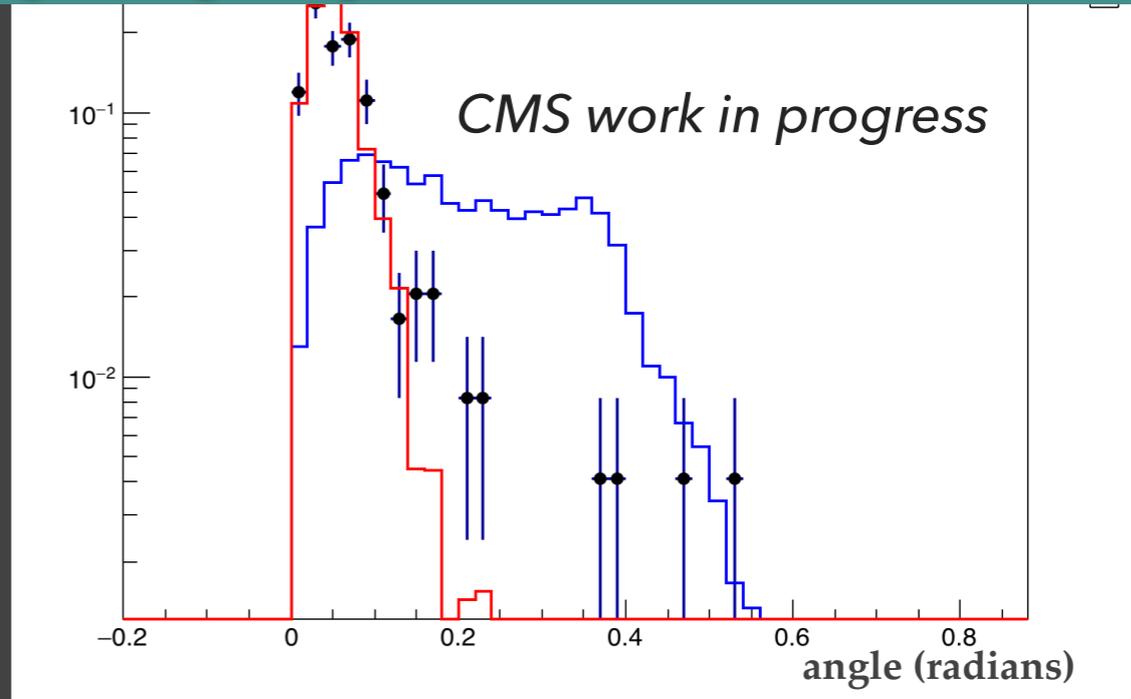
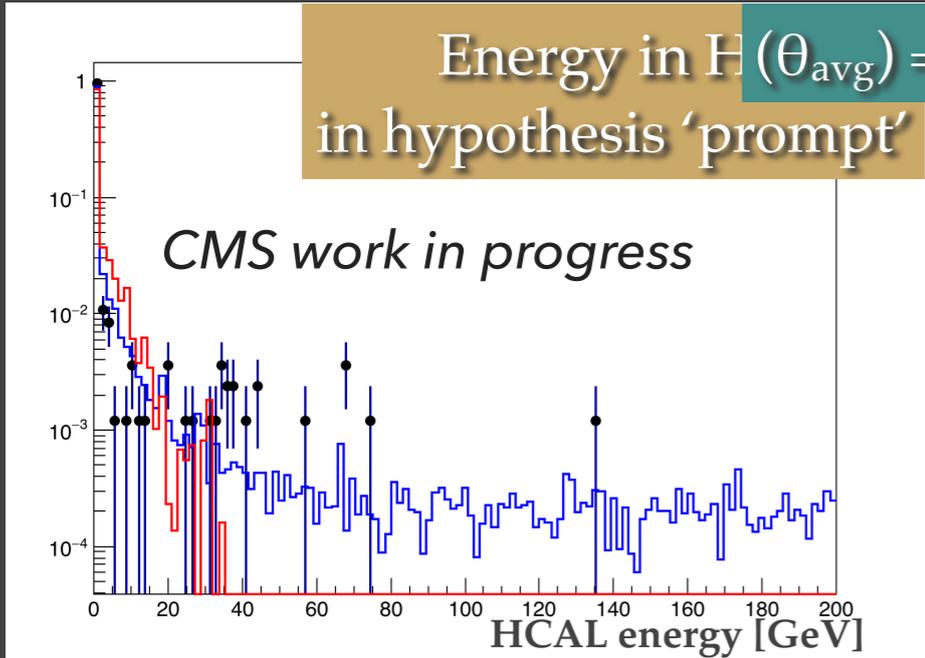
- ▶ Extrapolate from the position of the photon in the EE to that in the ES assuming the photon came from the collision vertex
- ▶ Collect the hits in the preshower planes
- ▶ Again form the angle between the line joining the position in the ES and the EE and collect the energy in the ES

- ▶ Energy in HE/ES in both the hypothesis
- ▶ Energy weighted average angle (average over both the hypothesis to see which one is more favored) the cluster of deposited in the HE/ES with that in the EE wrt the Z axis

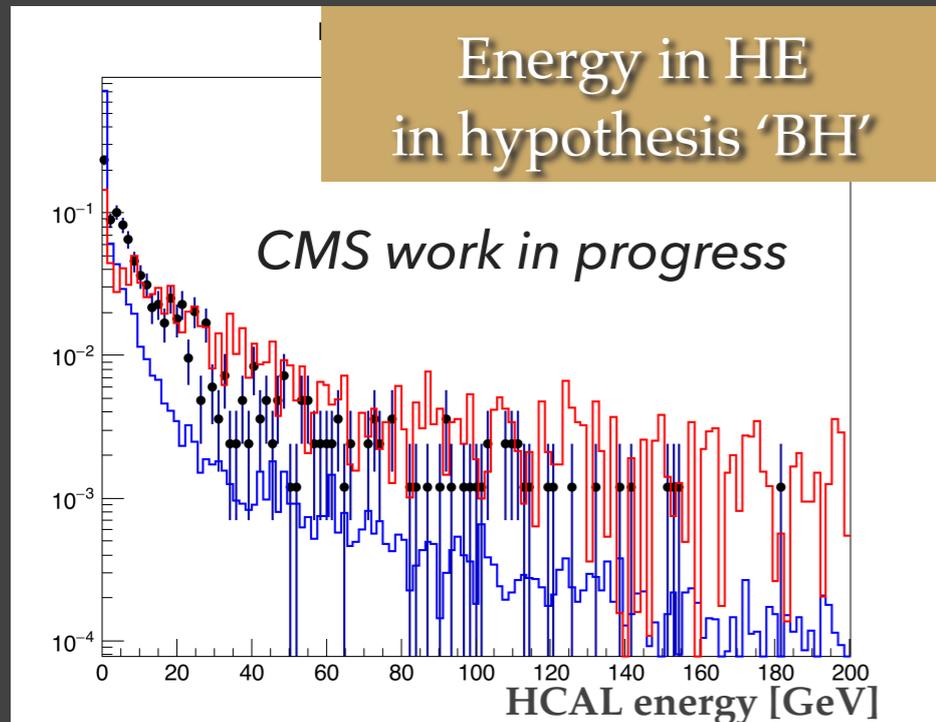
$$(\theta_{\text{avg}}) = (E_{\text{n_prompt}} * \theta_{\text{prompt}} + E_{\text{n_BH}} * \theta_{\text{BH}}) / (E_{\text{n_prompt}} + E_{\text{n_BH}})$$

- ▶ where θ_{prompt} is the angle made by the photon in hypothesis 'prompt' and θ_{BH} is the angle in hypothesis 'BH'. $E_{\text{n_prompt}}$ is the energy in the HE/ES in the 'prompt' hypothesis and $E_{\text{n_BH}}$ is the energy in the hypothesis 'BH'
- ▶ Beam being parallel to the Z axis would tend to have angles ~ 0

Energy in H (θ_{avg}) = $(E_{n_prompt} * \theta_{prompt} + E_{n_BH} * \theta_{BH}) / (E_{n_BH} + E_{n_prompt})$
 in hypothesis 'prompt'



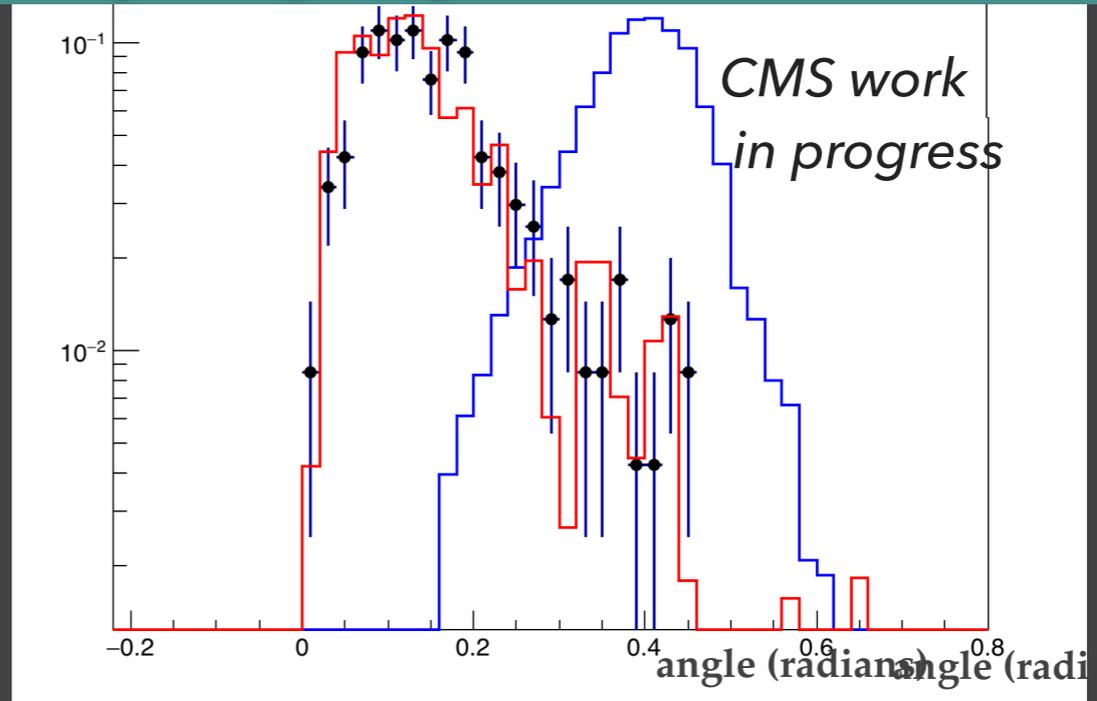
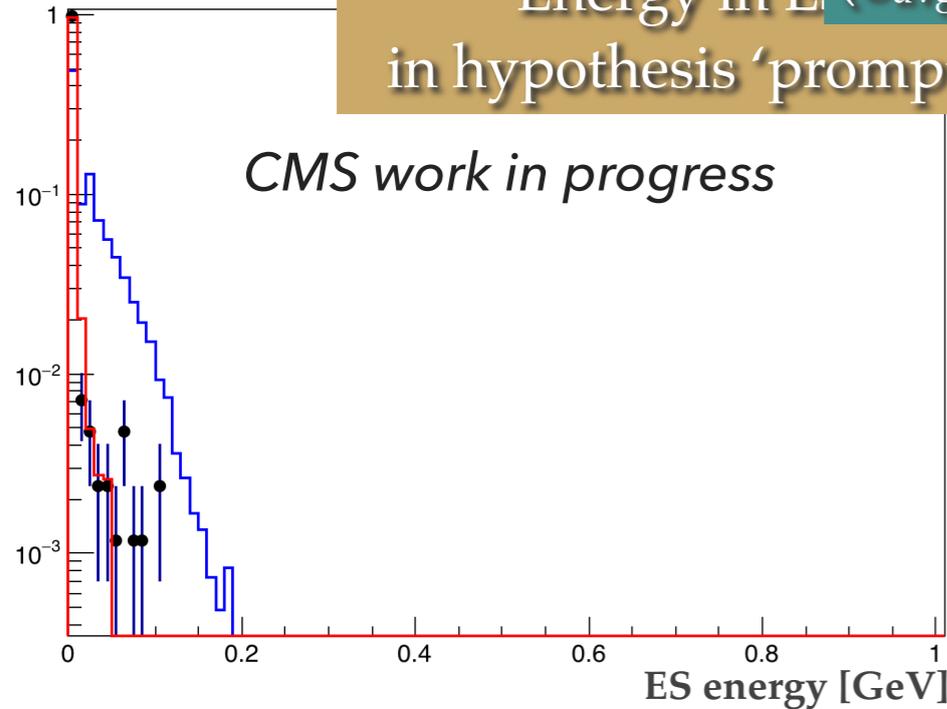
Energy in HE
 in hypothesis 'BH'



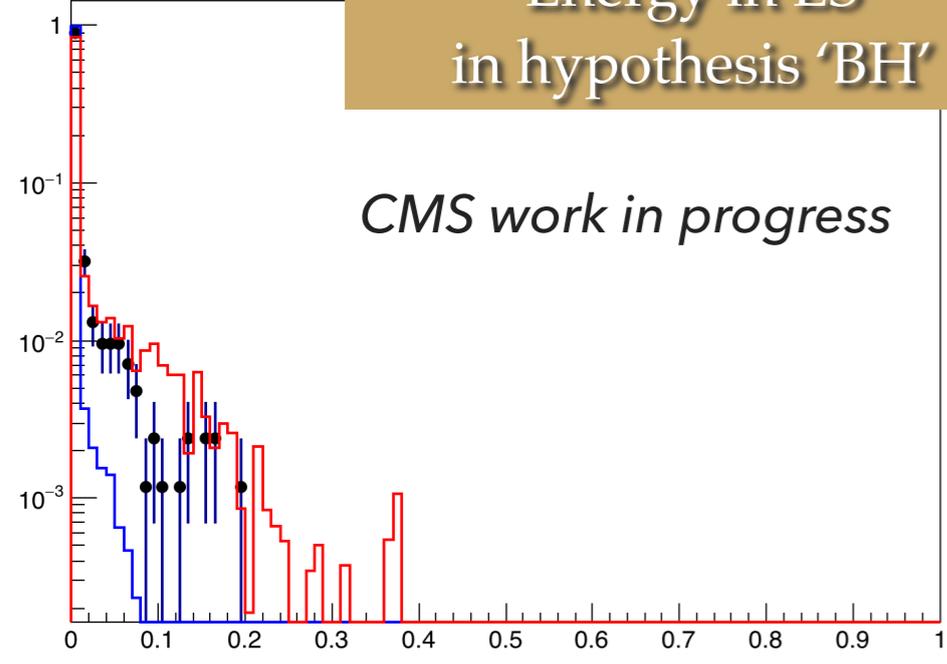
— Data BH
 — MC BH
 — MC prompt

- ▶ As expected beam halo in both data and MC peak ~ 0 whereas the prompt photons are away from 0

Energy in ES $(\theta_{avg}) = (E_{n_prompt} * \theta_{prompt} + E_{n_BH} * \theta_{BH}) / (E_{n_BH} + E_{n_prompt})$
 in hypothesis 'prompt'



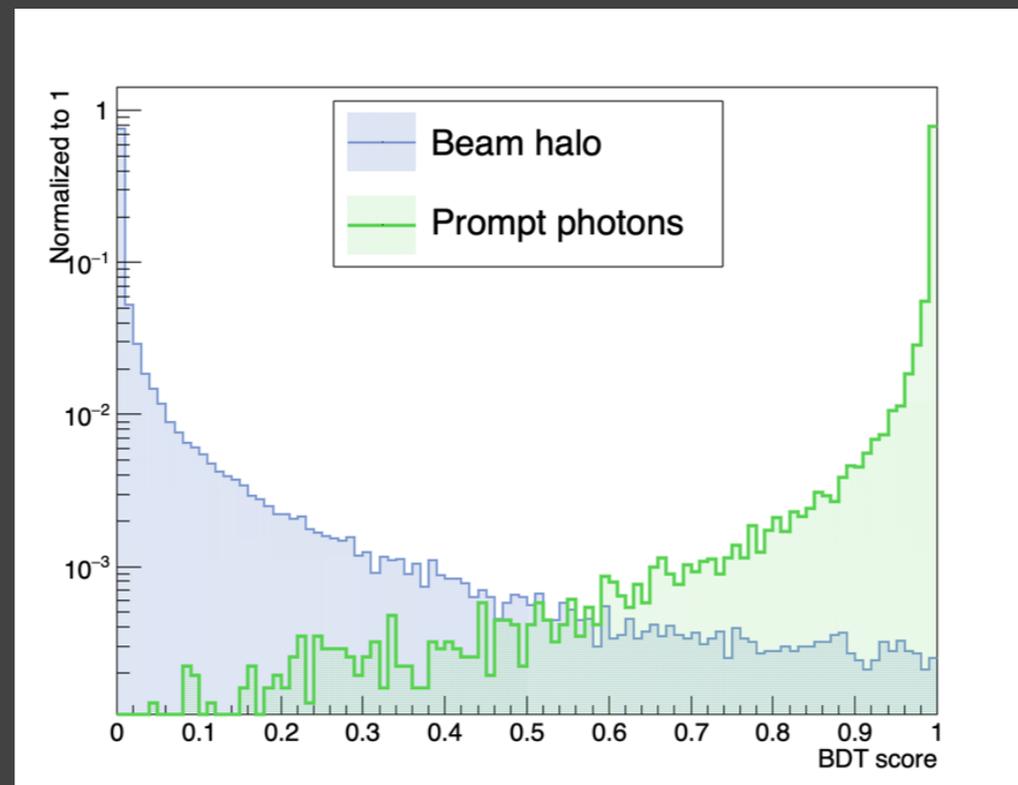
Energy in ES
 in hypothesis 'BH'



- Data BH
- MC BH
- MC prompt

▶ As expected beam halo in both data and MC peak ~0 whereas the prompt photons are away from 0

- ▶ The variables thus formed + a few more shower shape variables in the EE are fed to an xgboost
- ▶ The output of xgboost is used to reject the beam halo in the analysis
- ▶ This tagger is being used in the mono-photon analysis
 - ▶ Extends the acceptance of this analysis to EE



- ▶ Currently this detector is under testing and preparation phase
- ▶ Since quartz and PMTs are radiation hard detectors, plan for run 3 is to continue using the current detectors with an upgrade to the firmware
- ▶ For HL-LHC, there is a possibility to readout the BHM into L1 trigger system and provide an independent measurement of the beam halo for the CMS searches
- ▶ References: [1](#), [2](#), [3](#)

BACKUP

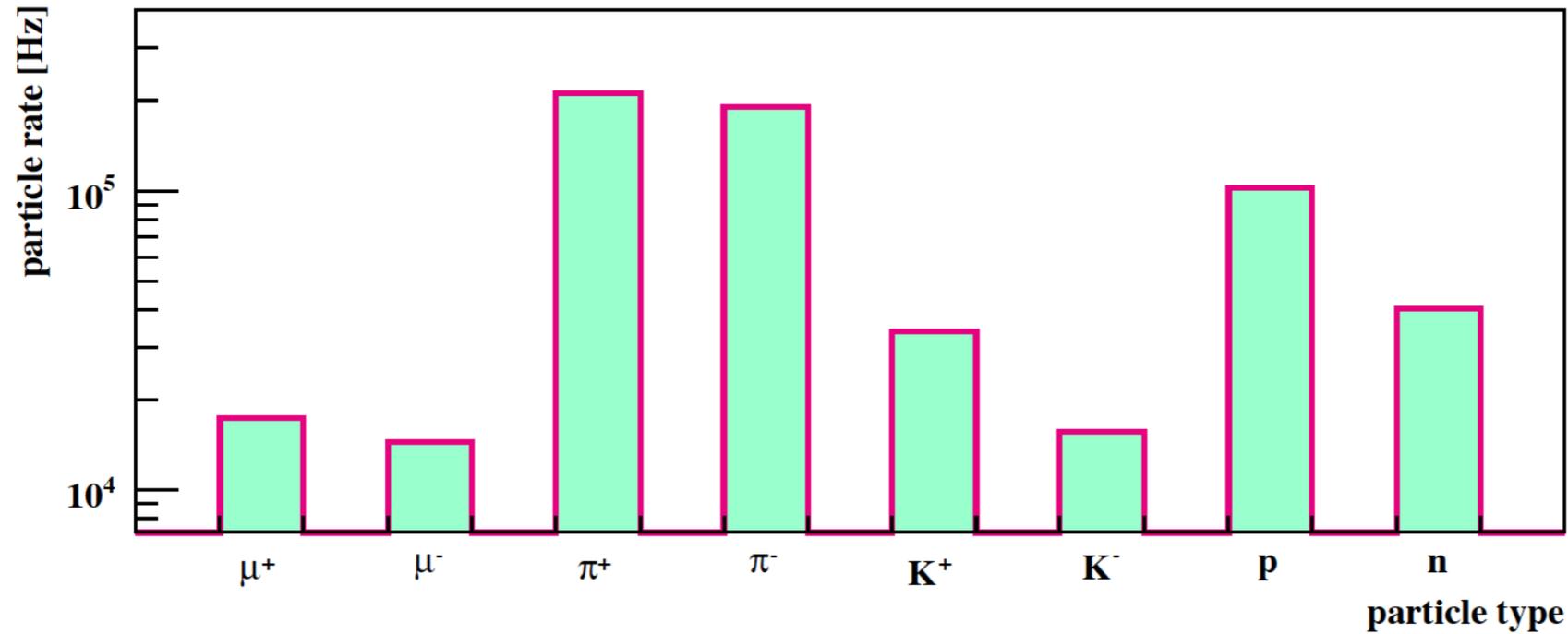


Figure 3: Rates of the different generated particles at the scoring planes at $z = \pm 23$ m.

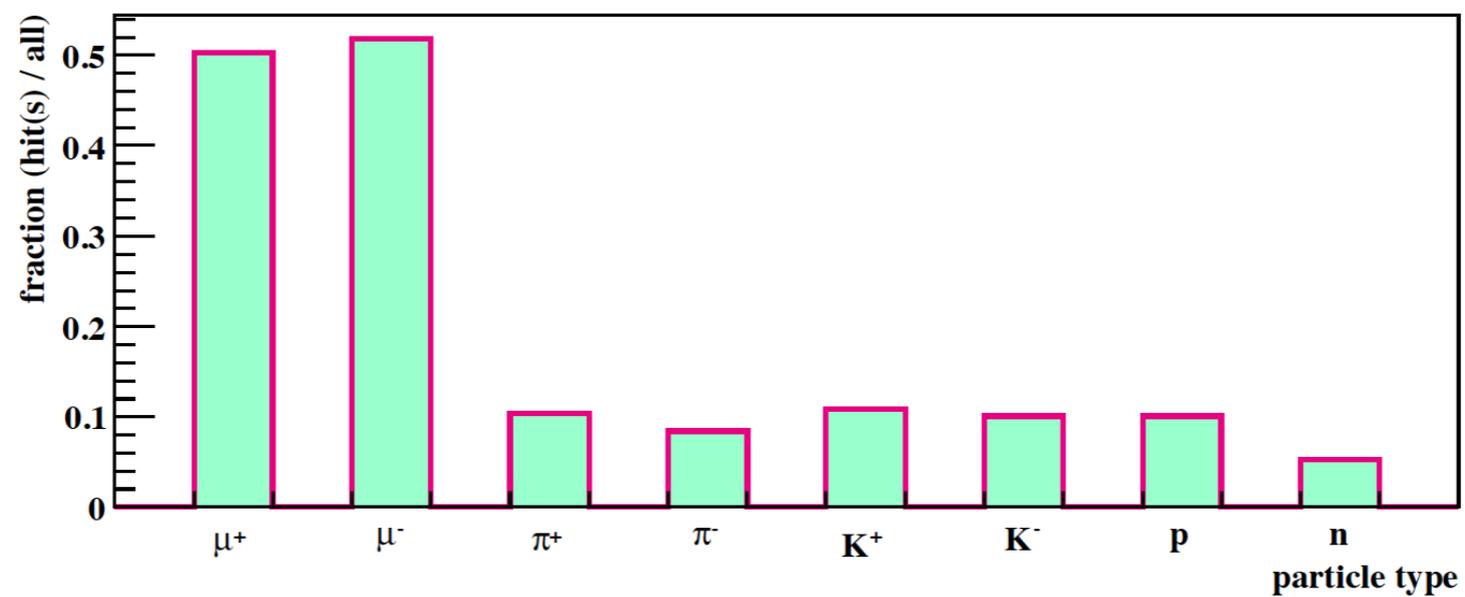
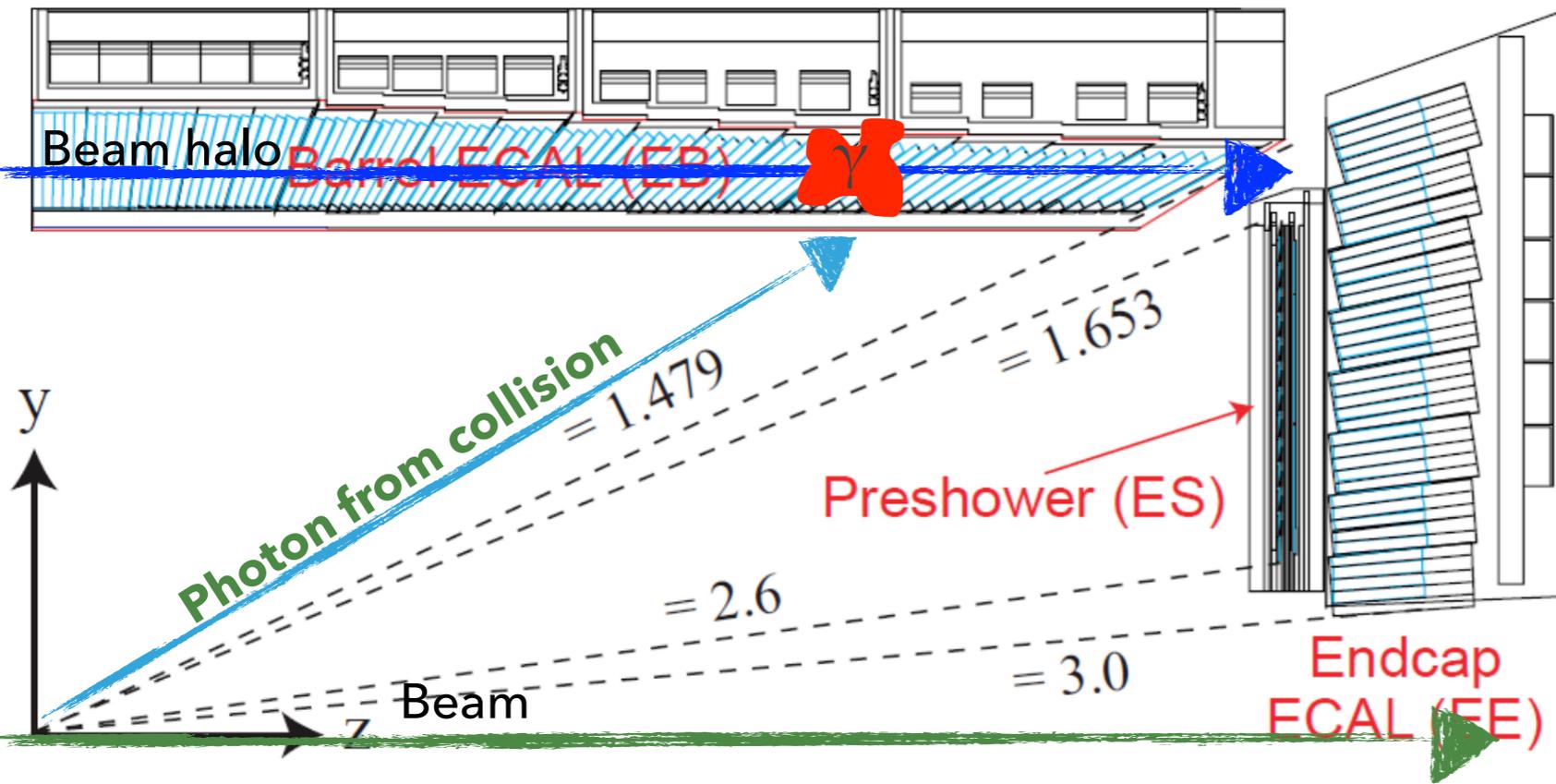
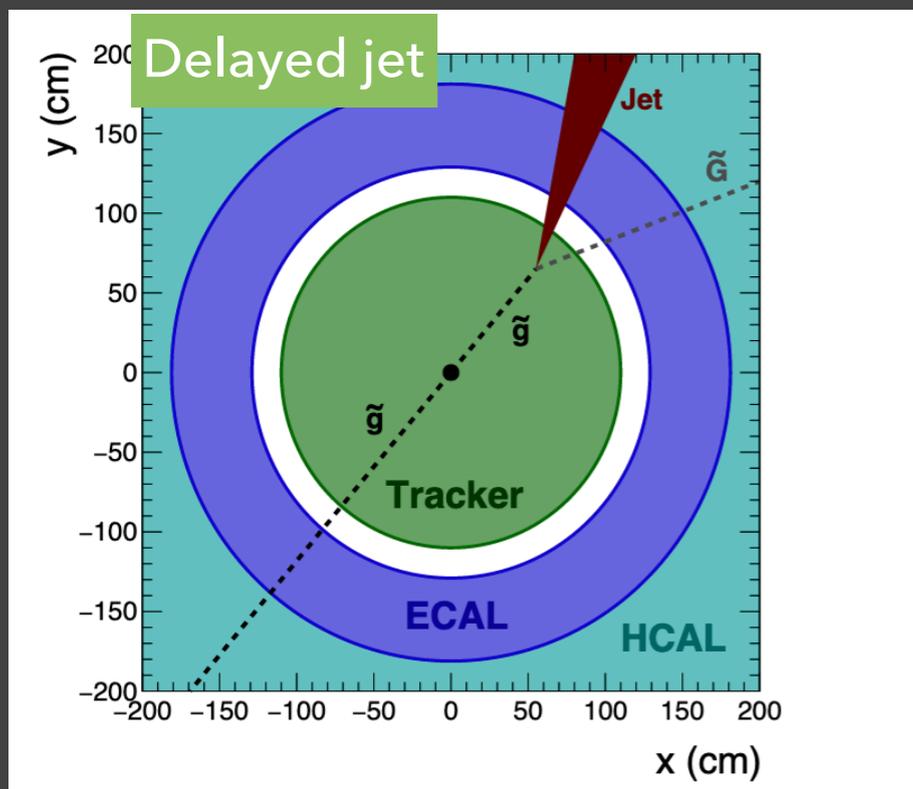


Figure 4: Fraction of generated particles that cause at least one hit in the CMS detector. Whereas this about 50% for muons it is only about 10% for hadrons.

Beam halo as background to several searches



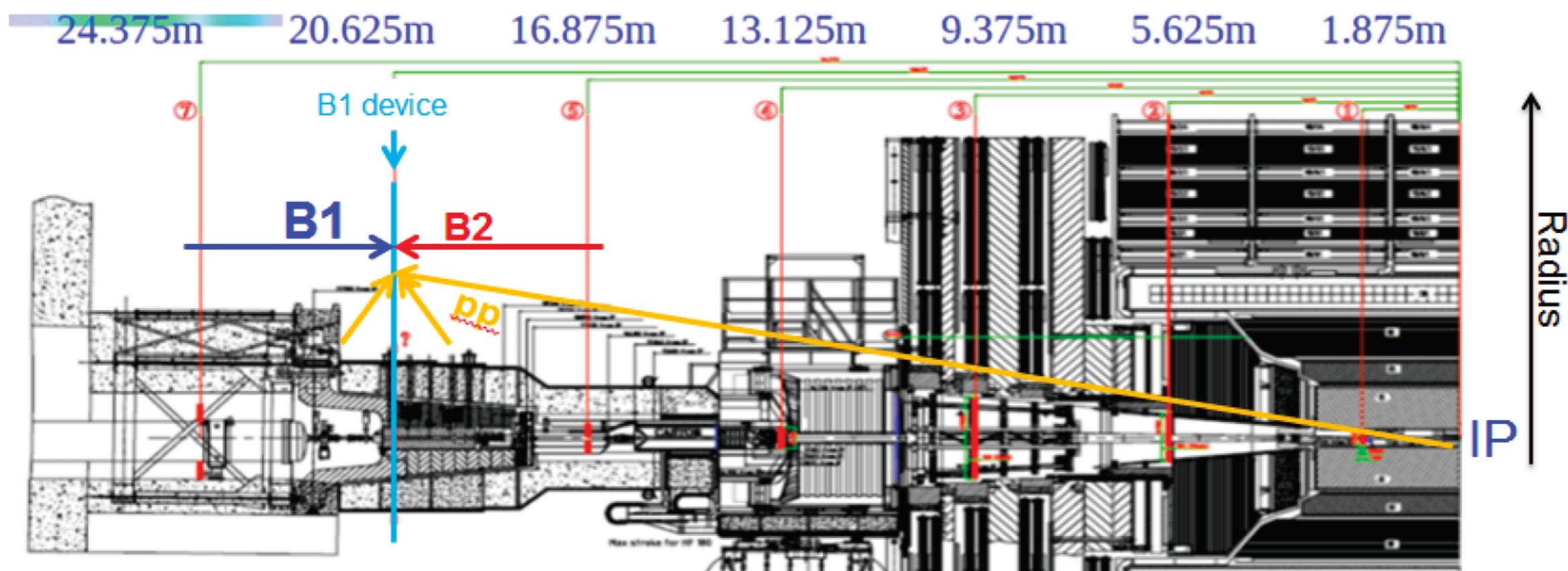
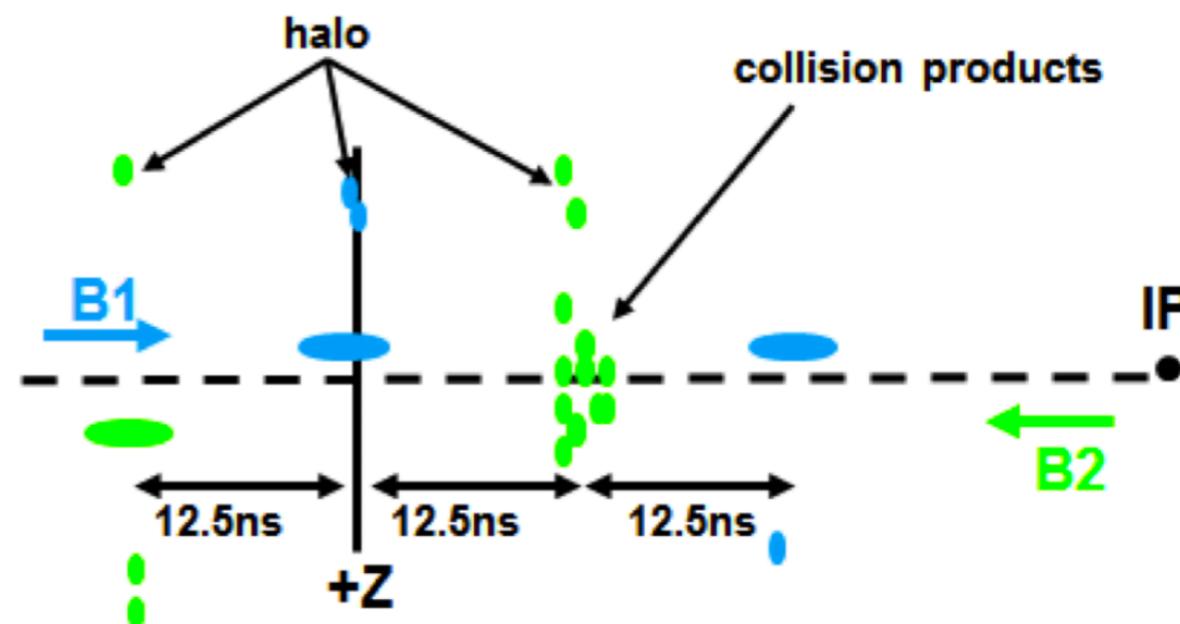
- ▶ The beam halo **arrives at a different time (in the ECAL barrel) compared to the collision products**
- ▶ Usually to reject such background, timing of the photon is required to be within some window consistent with collisions.

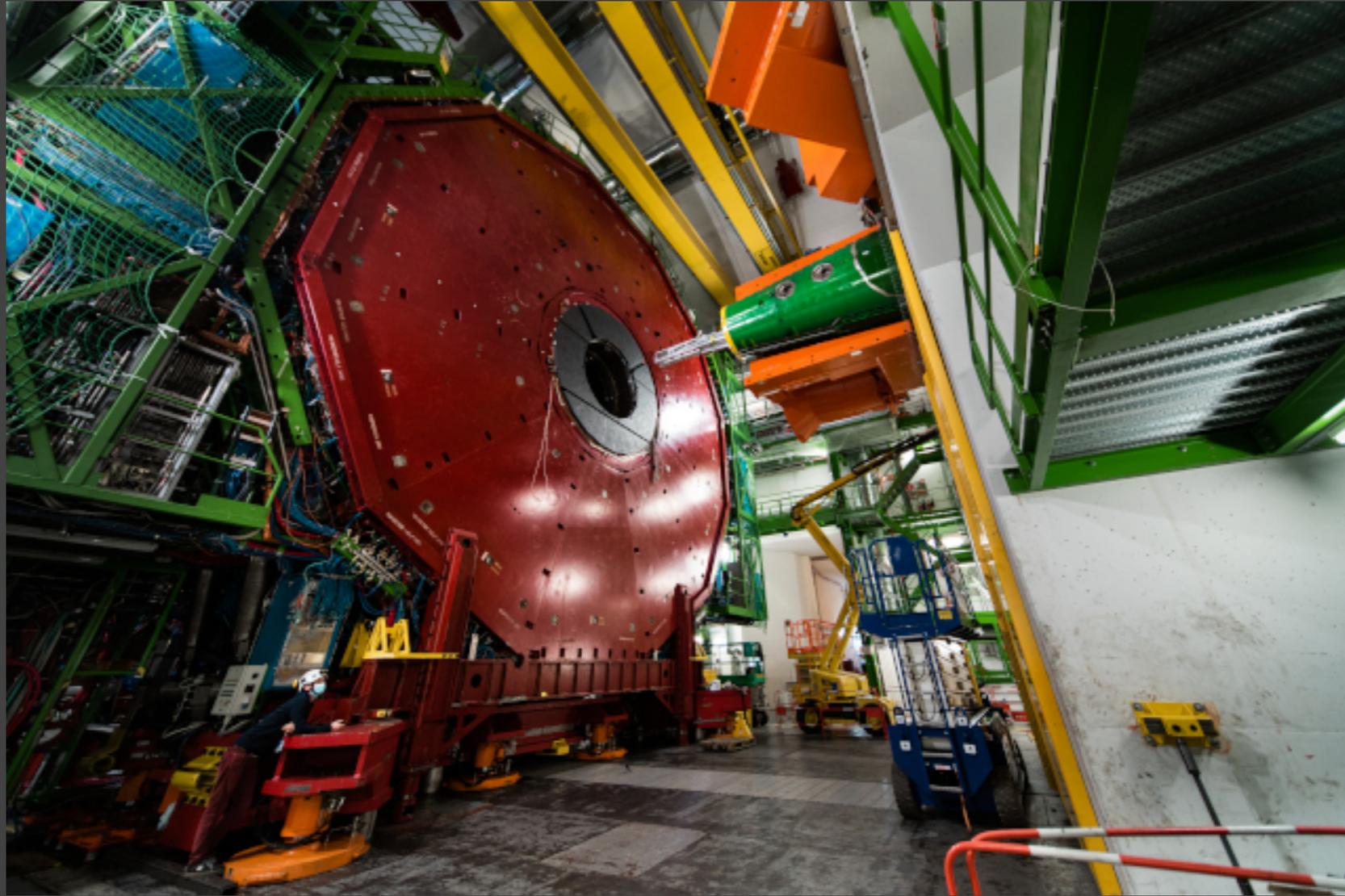


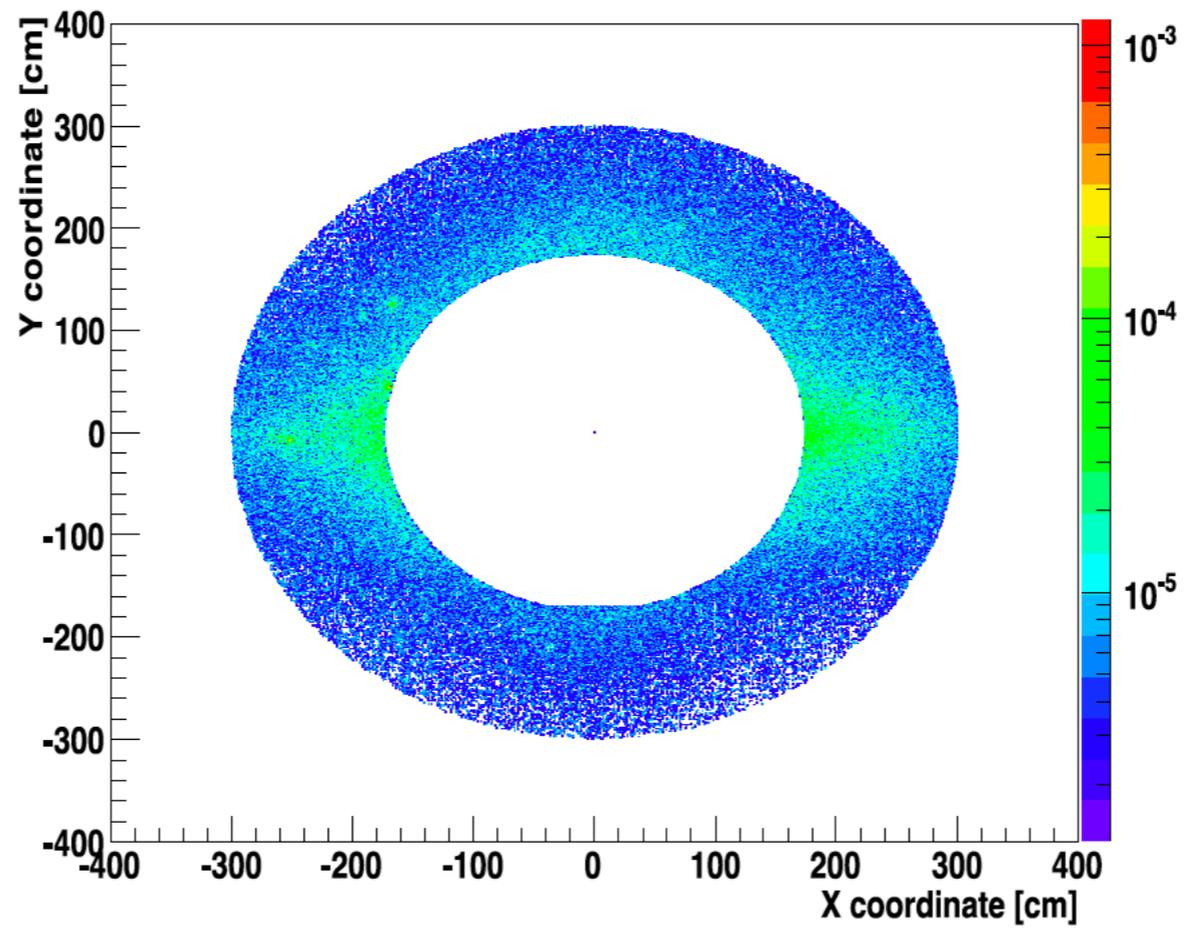
- ▶ However, in searches like:
 - ▶ Delayed jet: Pair of long-lived gluinos travel ~ 1 m such that the decayed product (gluon)
 - ▶ Jet is delayed in time (matches with the time of beam halo) and is only reconstructed only in the calorimeter
 - ▶ Beam halo can be huge

Choice of placement – golden locations

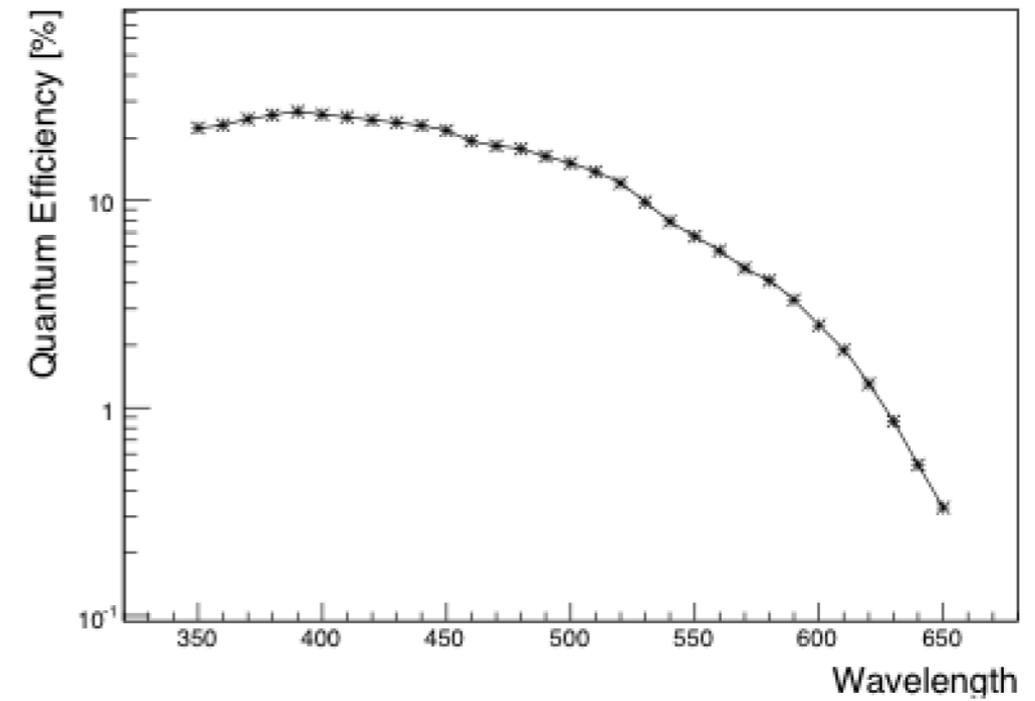
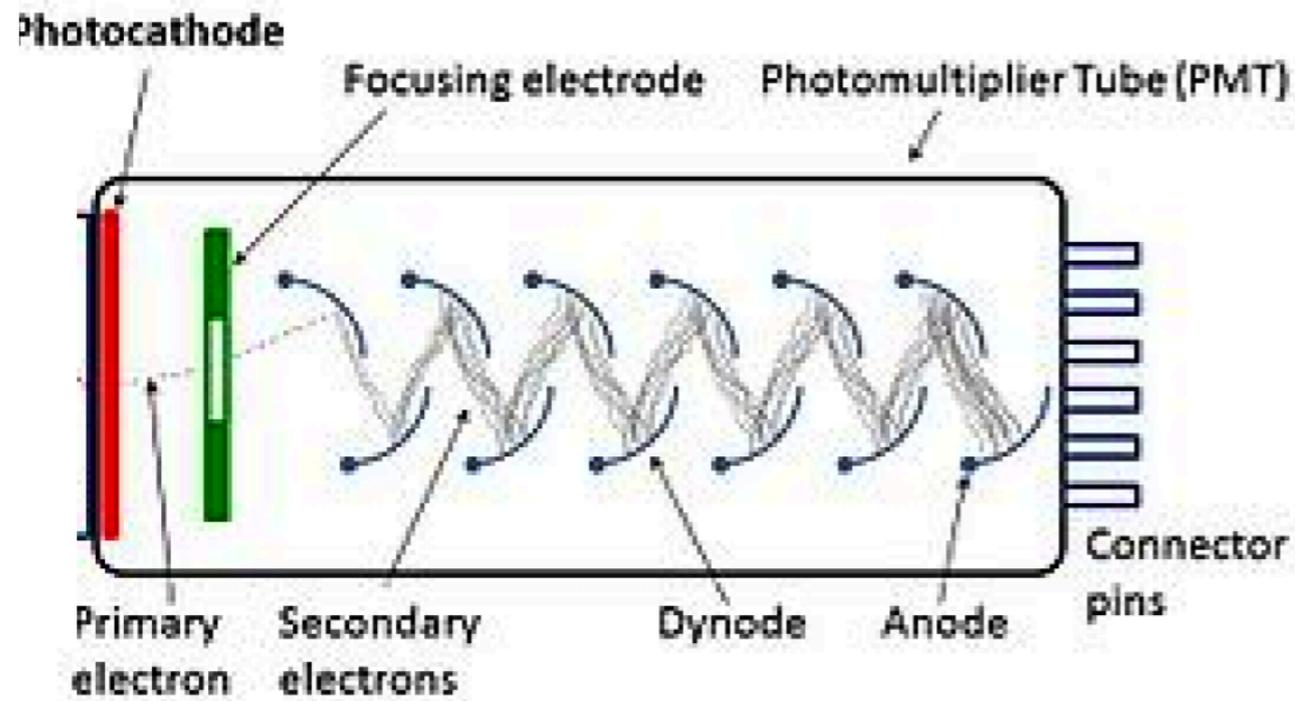
- ▶ Bunch spacing of 25 ns imposes a maximum time separation of 12.5 ns between the incoming beam halo and the outgoing beam (and its halo) and the collision products
- ▶ 7 golden locations which correspond to this maximum time
- ▶ Final choice is at a distance 20.625 m away : less harsh environment, less impact of the CMS magnetic field and least effected by the magnetic field



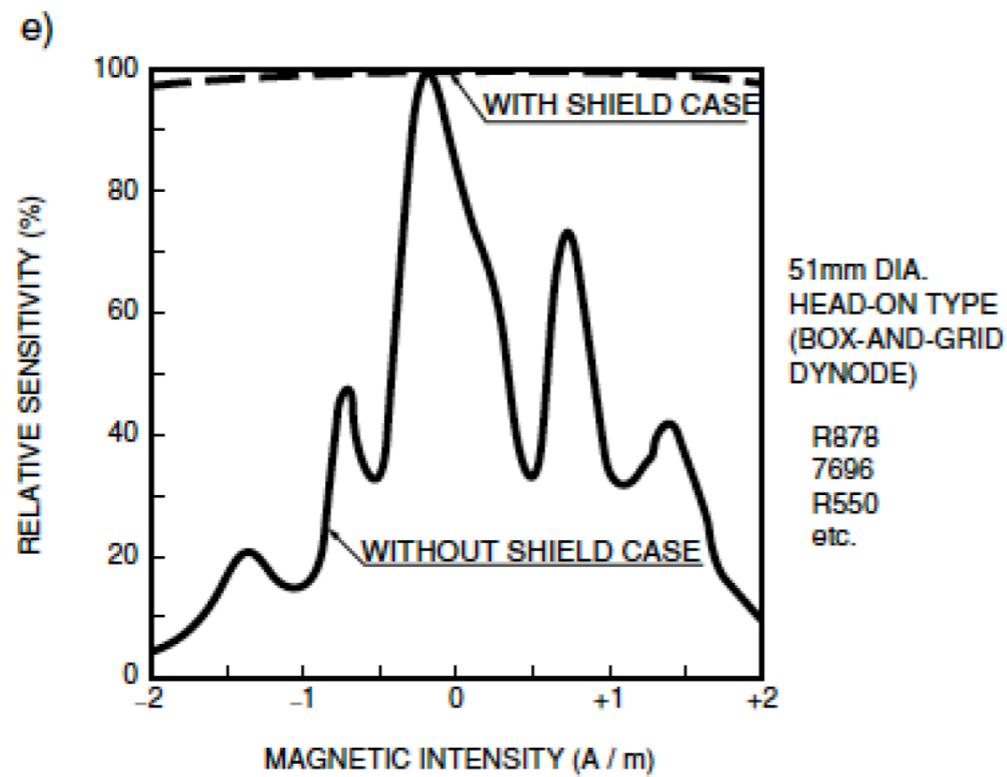




Photomultiplier tube



~25% at 400 nm



THBV3_0544EA

Figure 6.1: Magnetic Characteristics of 51 mm head-on photomultiplier tube as the one used for the BHM detector unit [Hama].

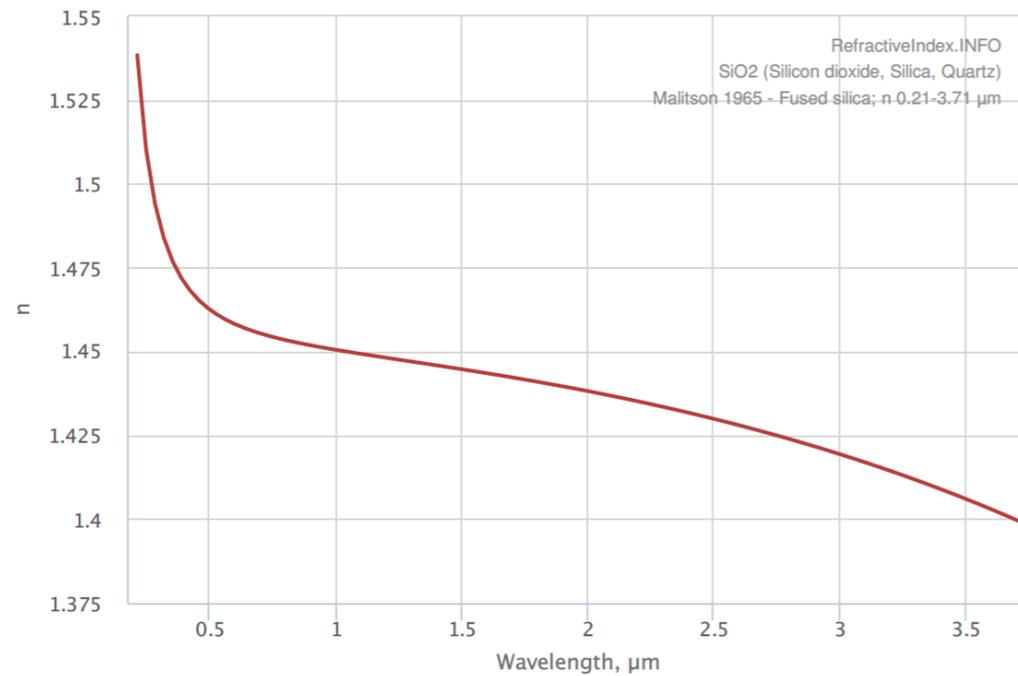


Figure 3.3: The index of refraction of quartz as a function of the wavelength [Pol].

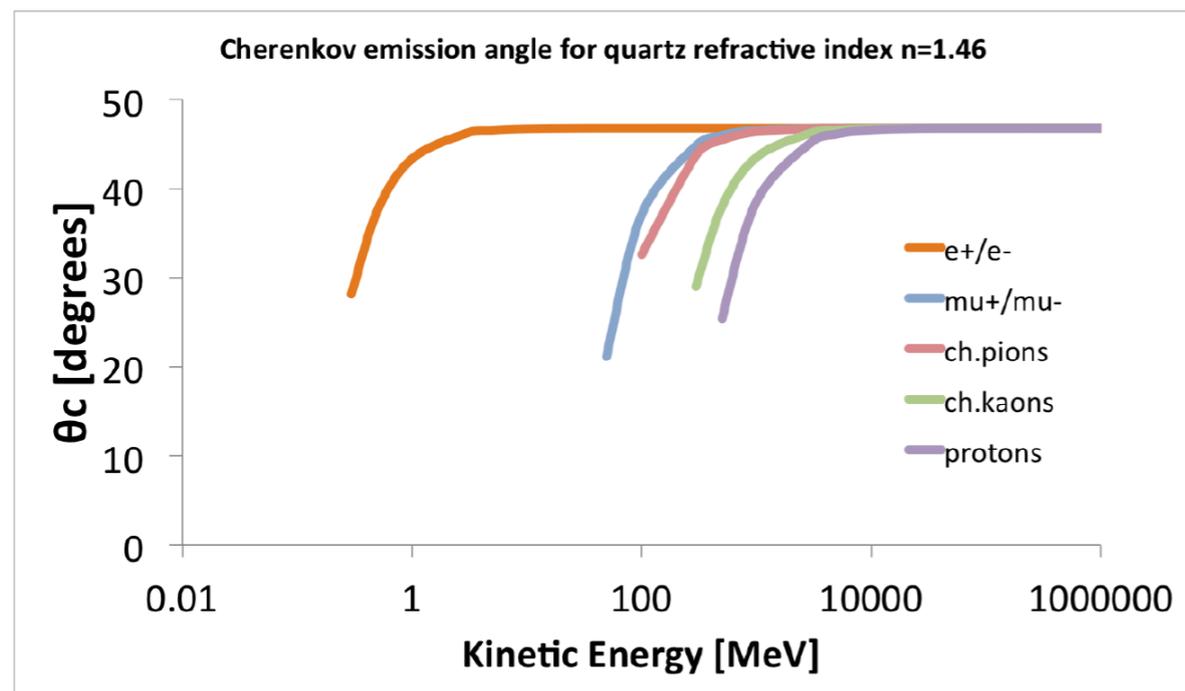
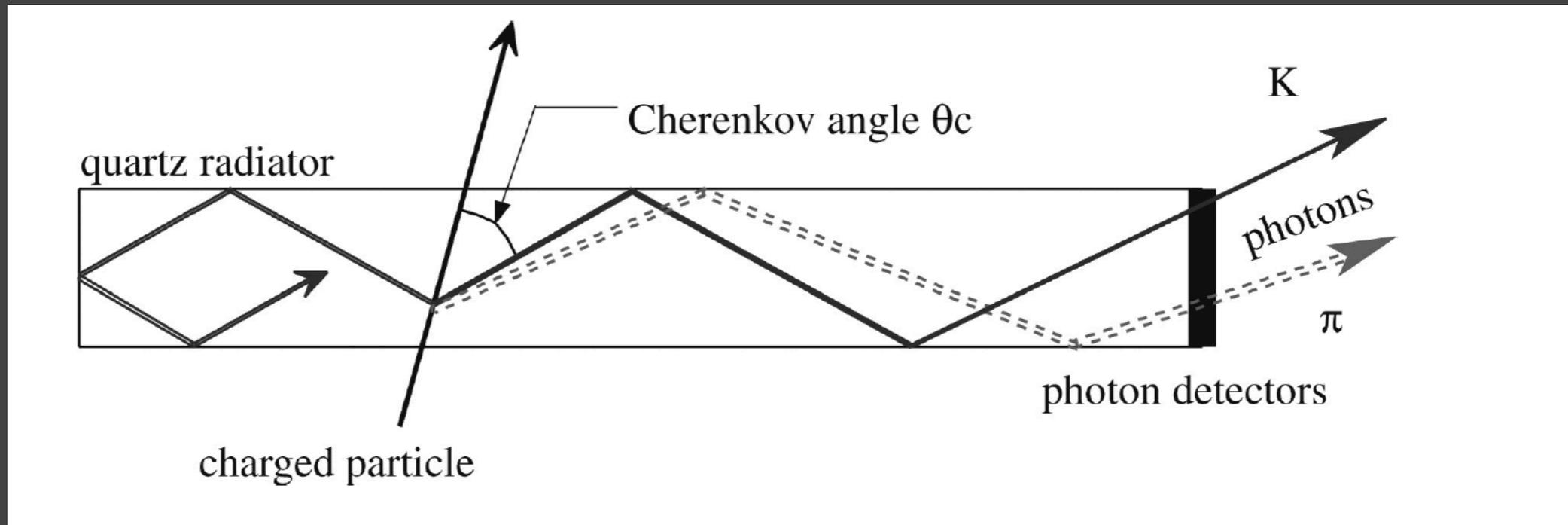


Figure 3.4: Cherenkov angles as a function of momentum for different particle species for $n=1.46$.



$$\theta_0(\lambda) = \arcsin \left(\frac{n_1(\lambda)}{n_2(\lambda)} \right) = \arcsin \left(\frac{1}{n_2(\lambda)} \right) = \arcsin \left(\frac{1}{1.46} \right) \approx 43^\circ$$

λ (nm)	n	$\beta_{threshold}$	θ_c for $\beta=1$
200	1.55	0.645	49.5°
300	1.485	0.673	48.0°
400	1.47	0.680	47.2°
500	1.465	0.683	47°
600	1.46	0.685	46.8°

Table 3.1: Index of refraction of quartz n, threshold velocity $\beta_{threshold}$ and angle of emission of Cherenkov photon $\theta_c(\beta = 1)$ for various wavelengths.

particle type	Rest mass (MeV)	E (MeV) for $\beta = 0.65$	E (MeV) for $\beta = 0.99$
e^+, e^-	0.511	0.161	0.190
μ^+, μ^-	105.7	33.38	39.37
protons	938.2	296.38	349.58
charged kaons	493.6	155.93	183.92
charged pions	139.6	44.09	52.00

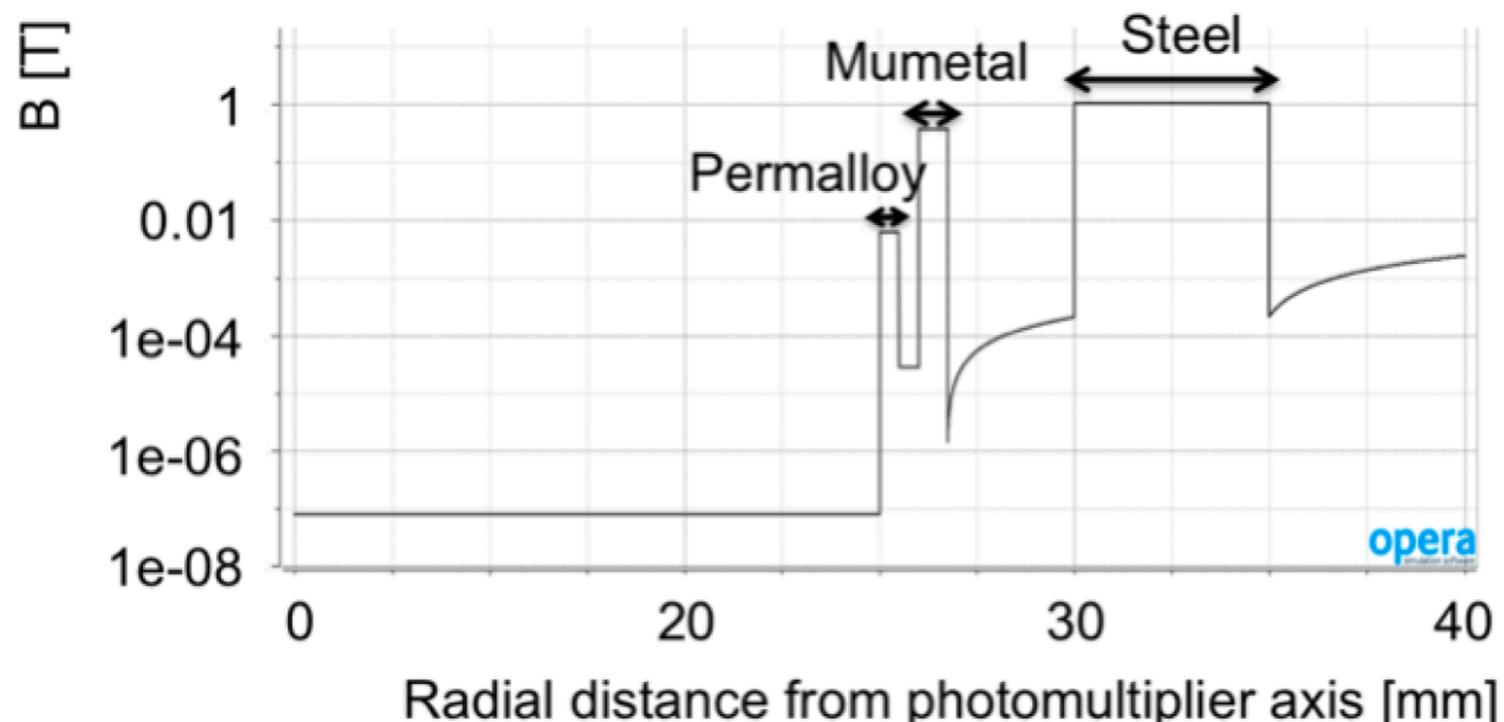
Table 3.2: Total energy for $\beta=0.65$ and $\beta=0.99$ for various charged particles.

Characteristic	Value
Acceptance per beam	300 cm^2
Time resolution	ns
Particle sensitivity	Direction-sensitive response to charged particles with a rejection power of $O(10^3)$ Insensitive to γ 's and thermal neutrons
Radiation hardness	100 krad
Magnetic field tolerance	18 mT with angle of 20° to beam axis
Logistical issues	Cost effective Robust mechanics Limited weight

Table 4.2: The summary of the BHM design requirements.



- ▶ PMT has increased sensitivity to the magnetic field
- ▶ To take care of this, it is shielded with a specific casing
- ▶ Shielding consists of 3 highly permeable materials which reduce the field to a very low level:



- ▶ permalloy
- ▶ mu-metal
- ▶ Steel

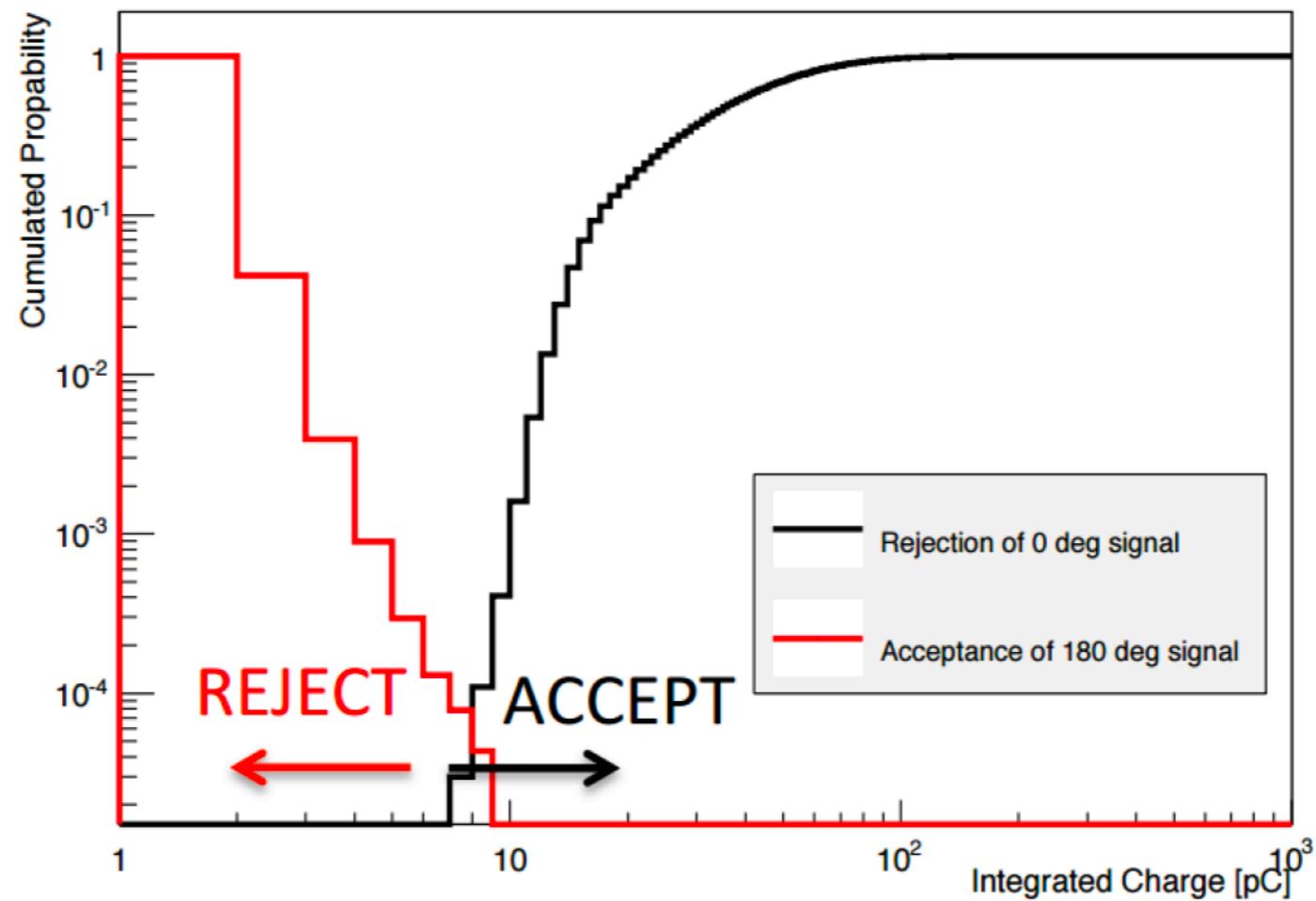
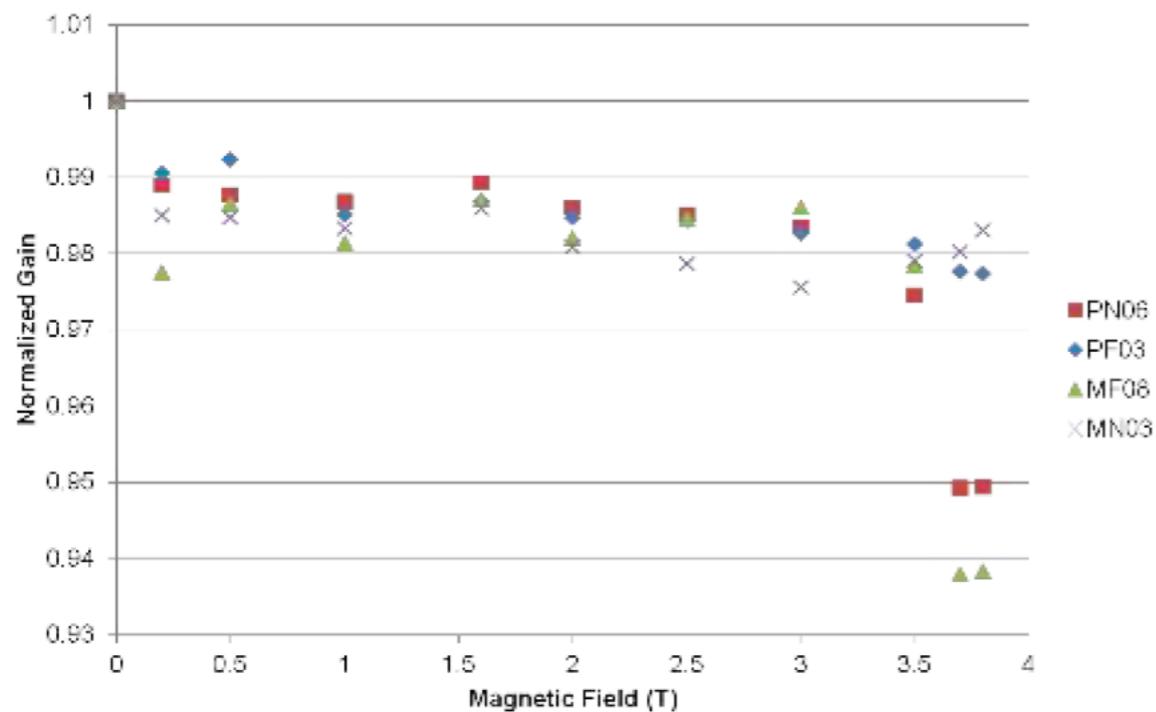
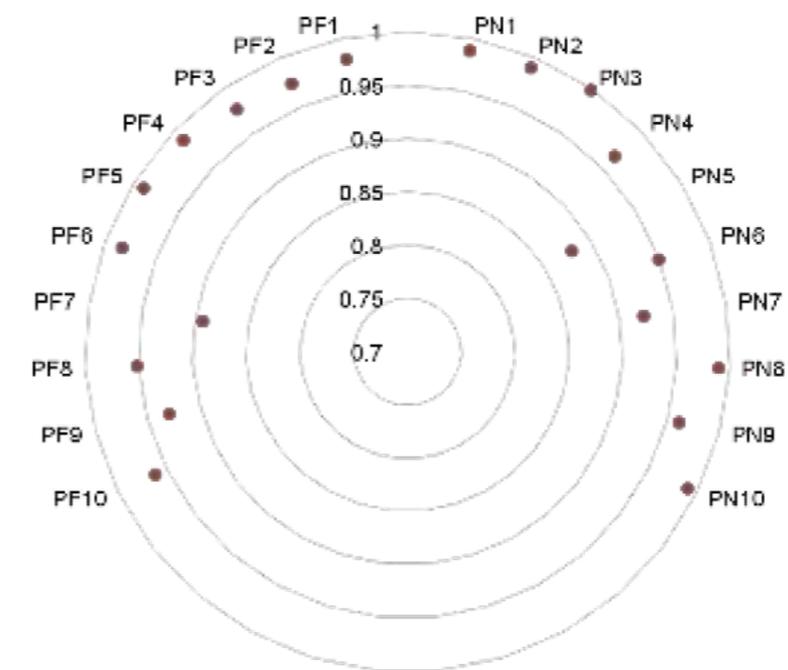


Figure 3.8: Cumulative integral of the normalized distributions of the signal charge measured for electrons with angles of 0° and 180° .



(a) PMT gain vs. Magnetic field



(b) % decrease in PMT gain at 3.8 T

Figure 4.7: *Top:* Normalized PMT gain versus magnetic field strength, shown for four sample units. *Bottom:* Percentage decrease in PMT gain at maximum field of 3.8T, shown for all units measuring beam 1.

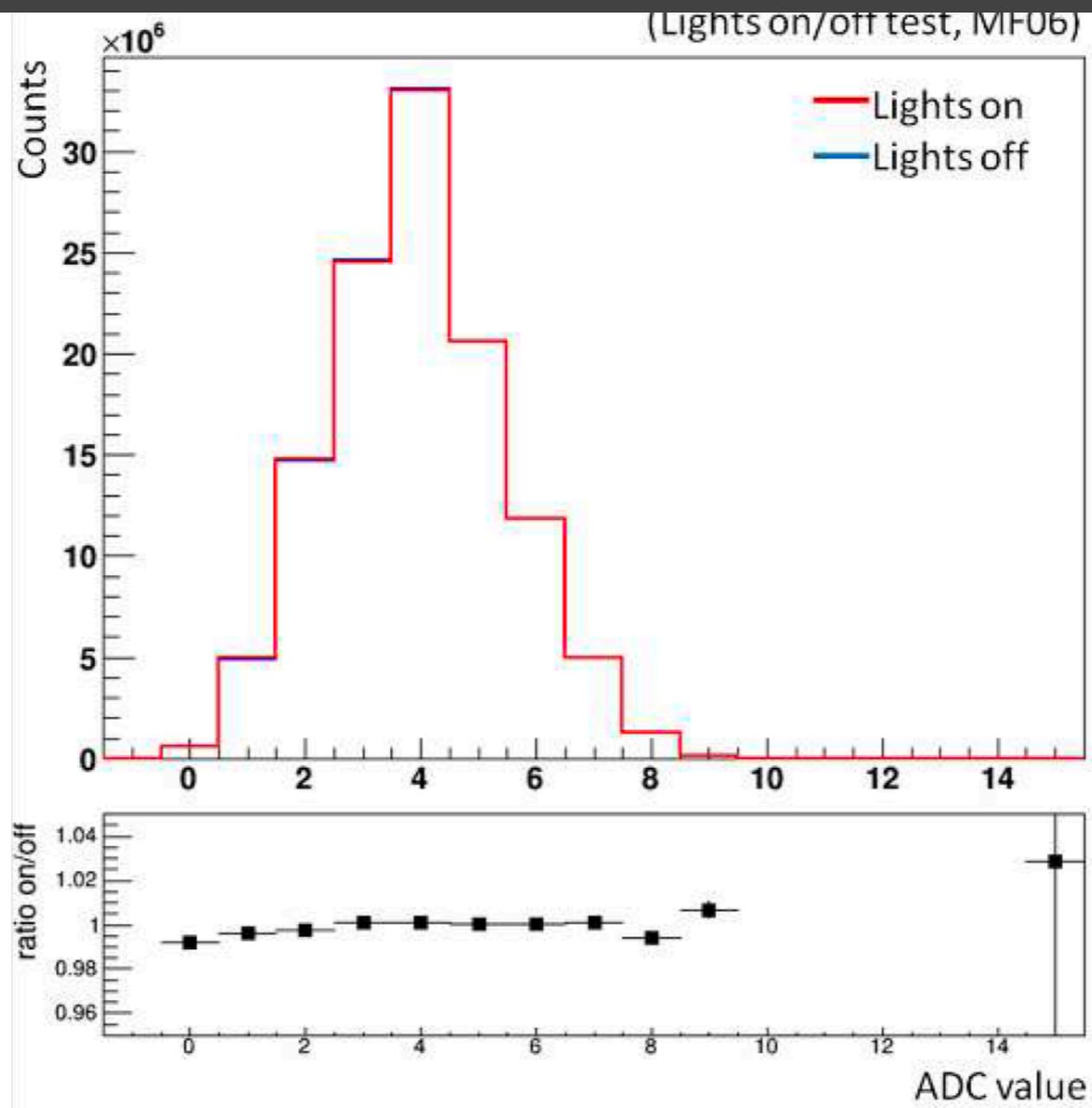


Figure 4.8: Detector units showed no change in response, regardless of light in the CMS cavern.

Location	Distance from IP [m]	MIB flux [particles/ cm^2]/s	Ratio of MIB flux/Collision products flux
4	13.125	(0.108±0.022)- (0.868± 0.343)	2×10^{-3} - 8×10^{-2}
5	16.875	(0.135±0.017)- (1.556± 0.335)	4×10^{-4} - 2×10^{-3}
6	20.625	(0.113±0.016)- (1.883±0.843)	4×10^{-4} - 4×10^{-3}

Table 4.1: Comparison of the ratio and the absolute flux in the candidate locations.

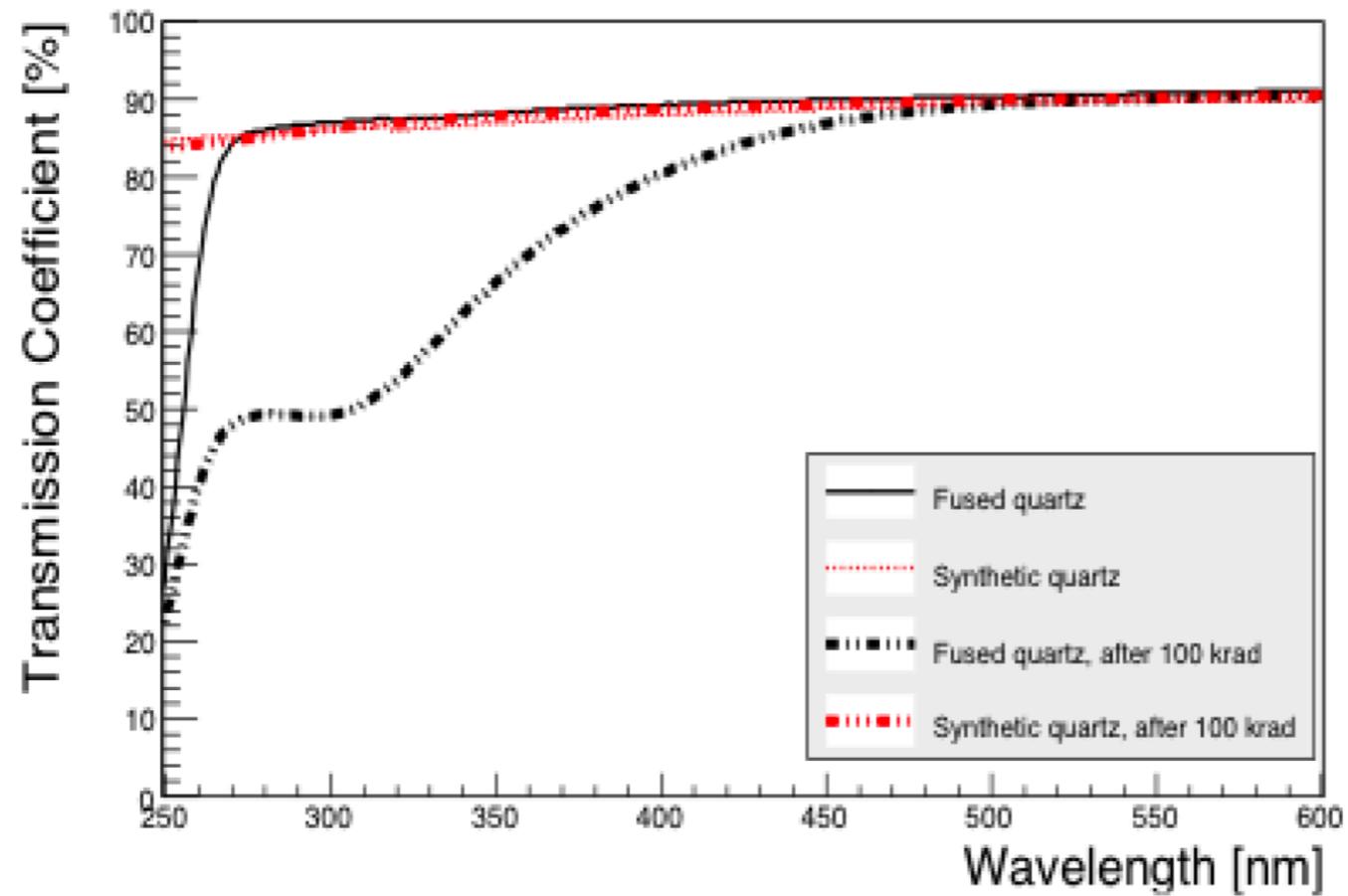


Figure 5.2: The transmission coefficients for natural and synthetic fused quartz as a function of the wavelength before and after irradiation with a dose of 100 krad.

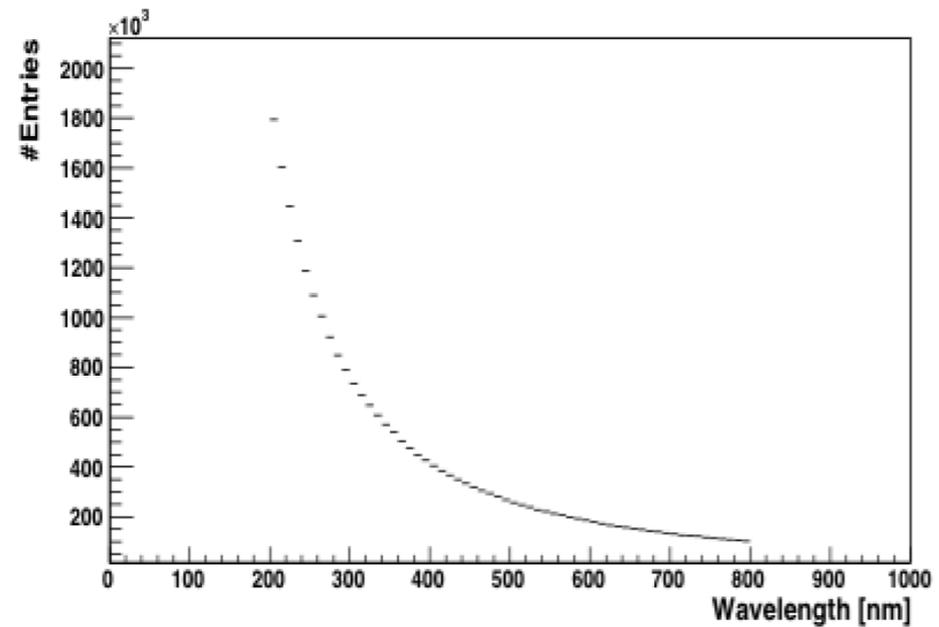


Figure 5.4: The wavelength of the light produced, as simulated when a muon of 4 GeV crosses 10 cm long quartz radiator, entering from the center of the front face of the bar.

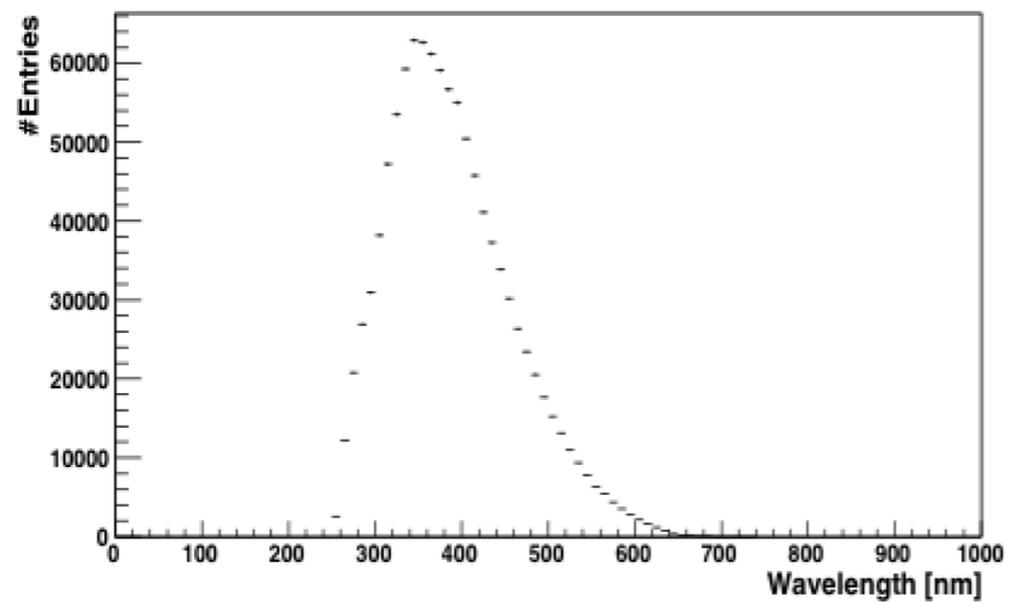


Figure 5.5: The wavelength of the light detected by the photocathode, as simulated when a muon of 4 GeV crosses 10 cm long quartz radiator, entering from the center of the front face of the bar.

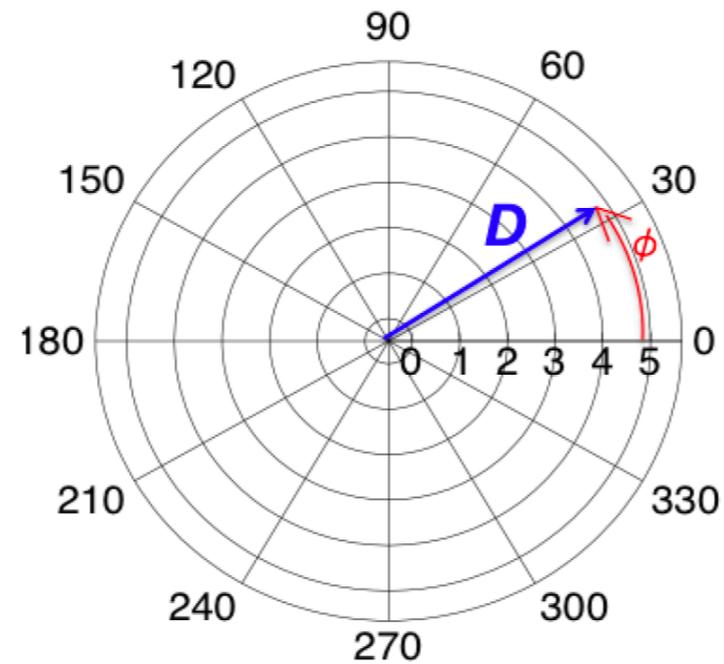


Figure 5.17: A diagram illustrating on a polar plot. The directional gain, D , at a particular impact angle corresponds to the radial length. The rings manifest the radial scale. Impact angles from 0° to 90° correspond to the forward direction (i.e. towards the photocathode), while angles from 90° to 180° correspond to the backward direction (i.e. towards the black painted side of the radiator).

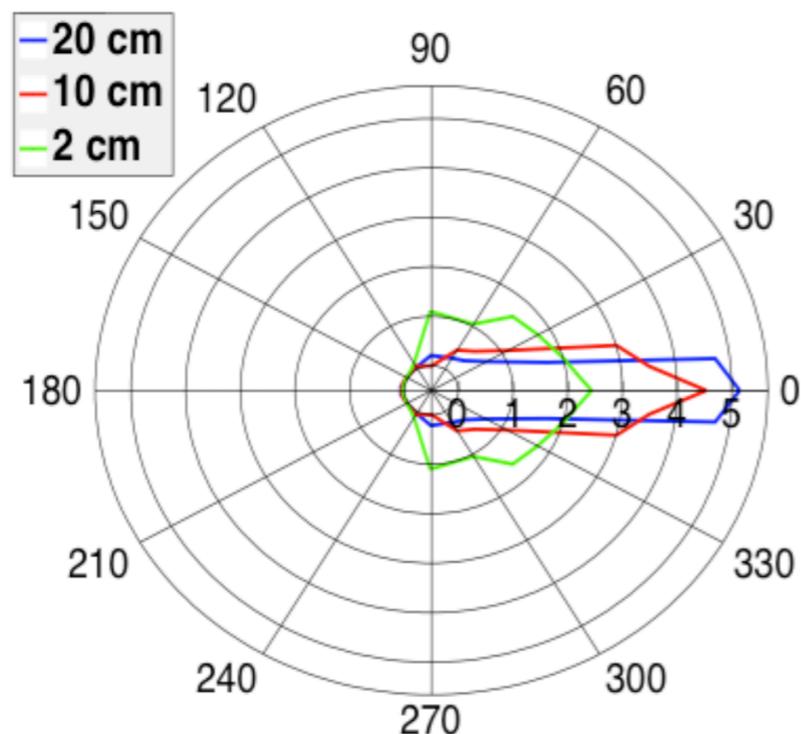
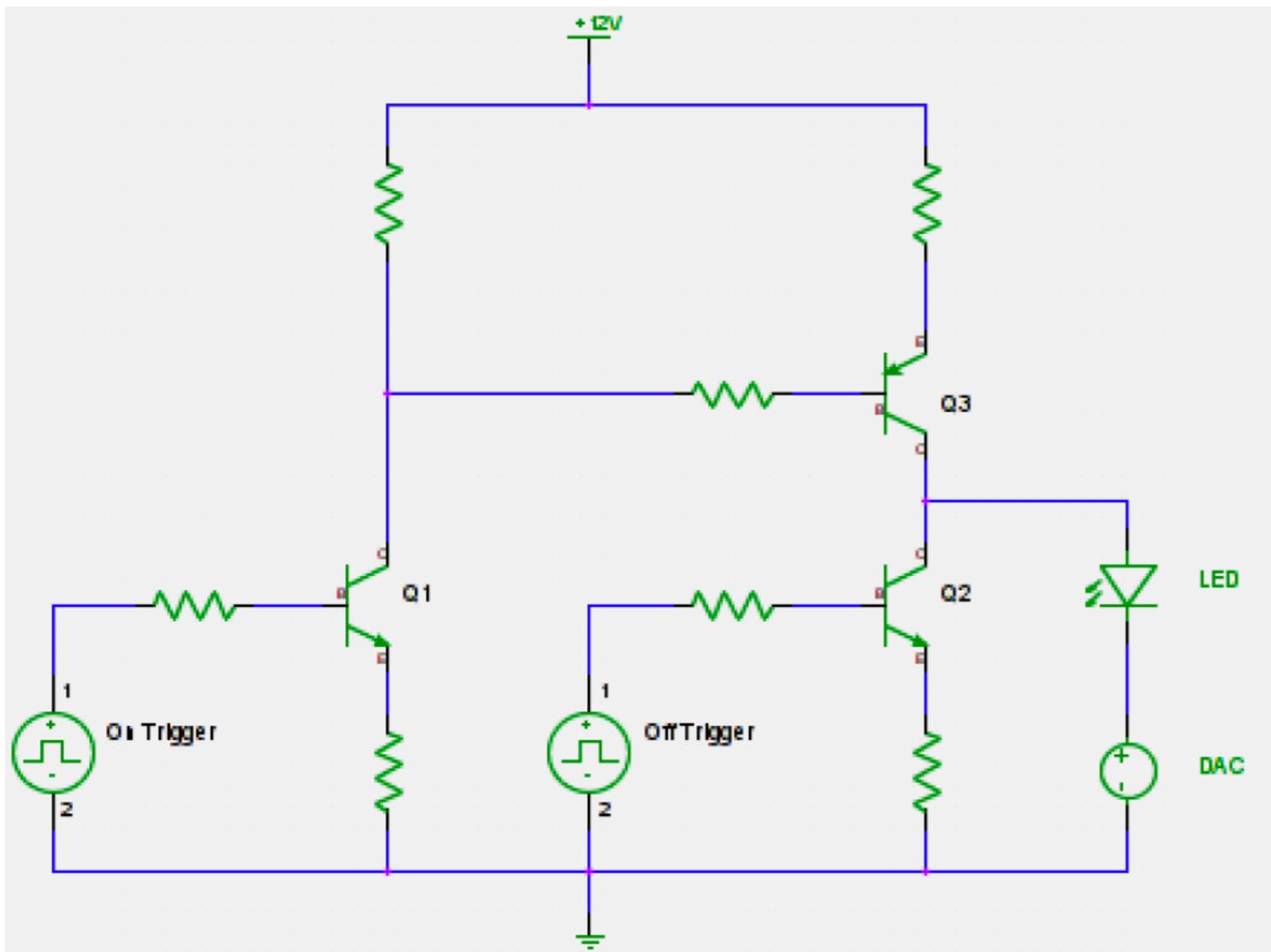


Figure 5.18: The directional gain of three radiators of length 2 cm, 10 cm and 20 cm as function of the impact angle for 4 GeV muons.

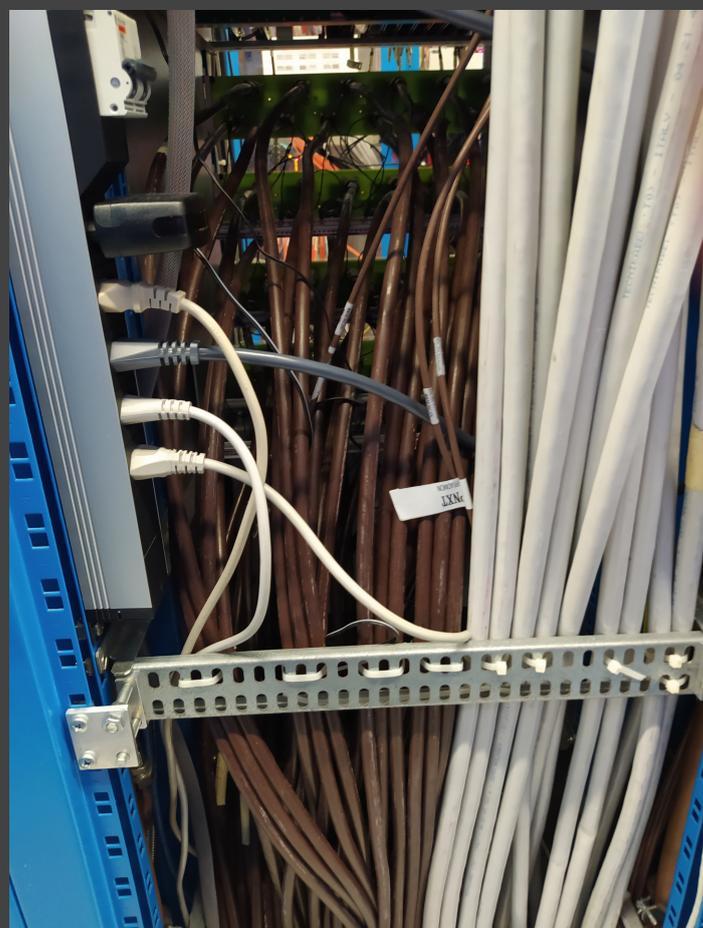
- ▶ Average number of photoelectrons = average value of photoelectrons integrated over ϕ

$$D(l, \phi) = \frac{\text{Number of photoelectrons}(l, \phi)}{\text{Average number of photoelectrons}(l)}$$

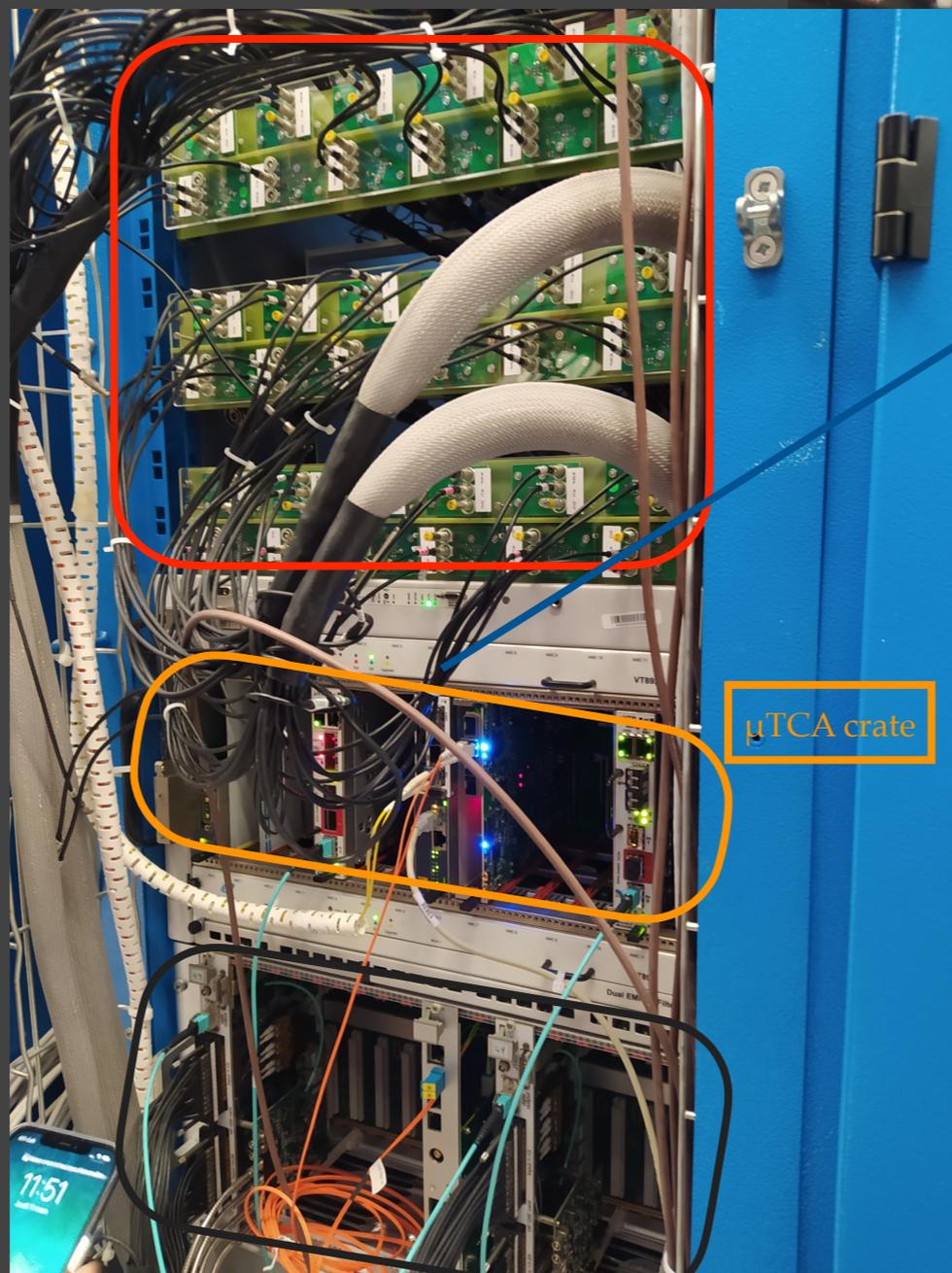
Pulser circuit



Readout system in USC (service cavern)



Signal cables (brown)

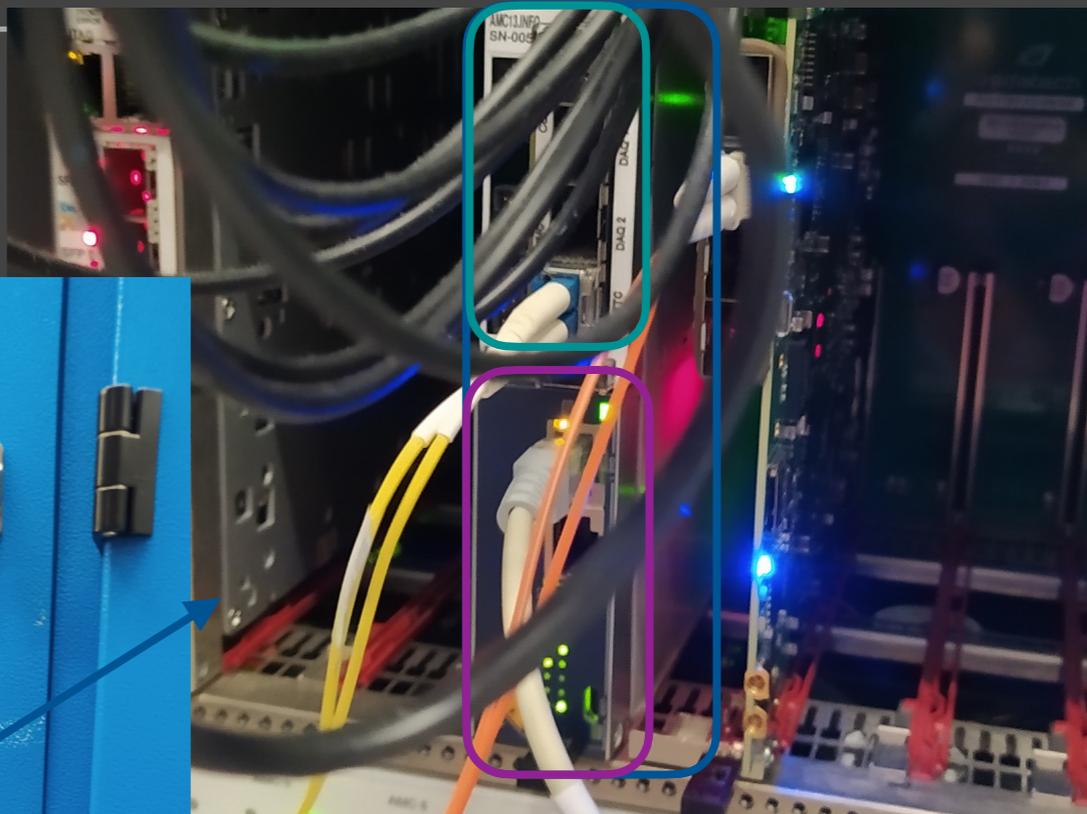


Patch panel

μTCA crate

QIE +
calibration

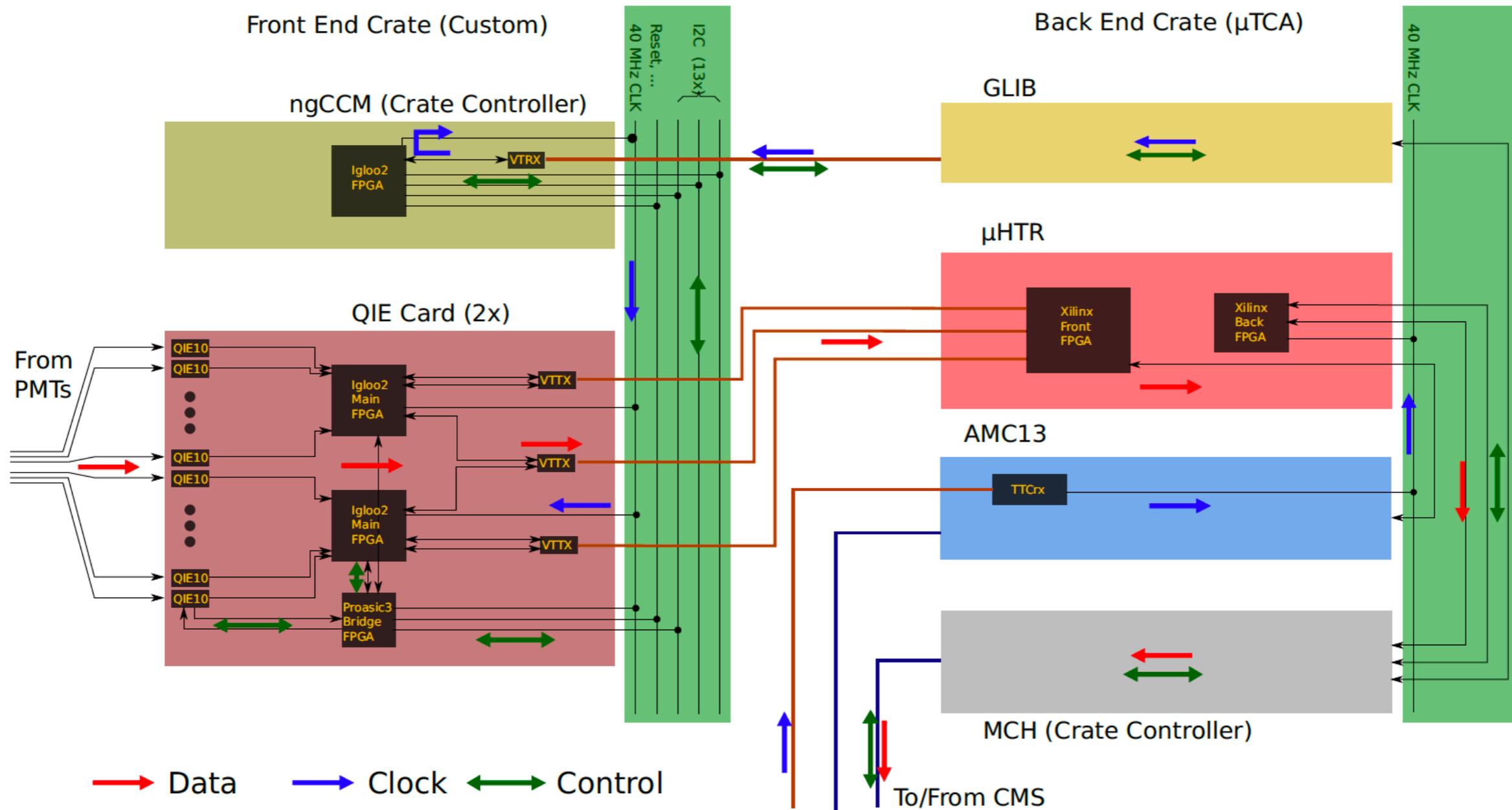
SF108 rack



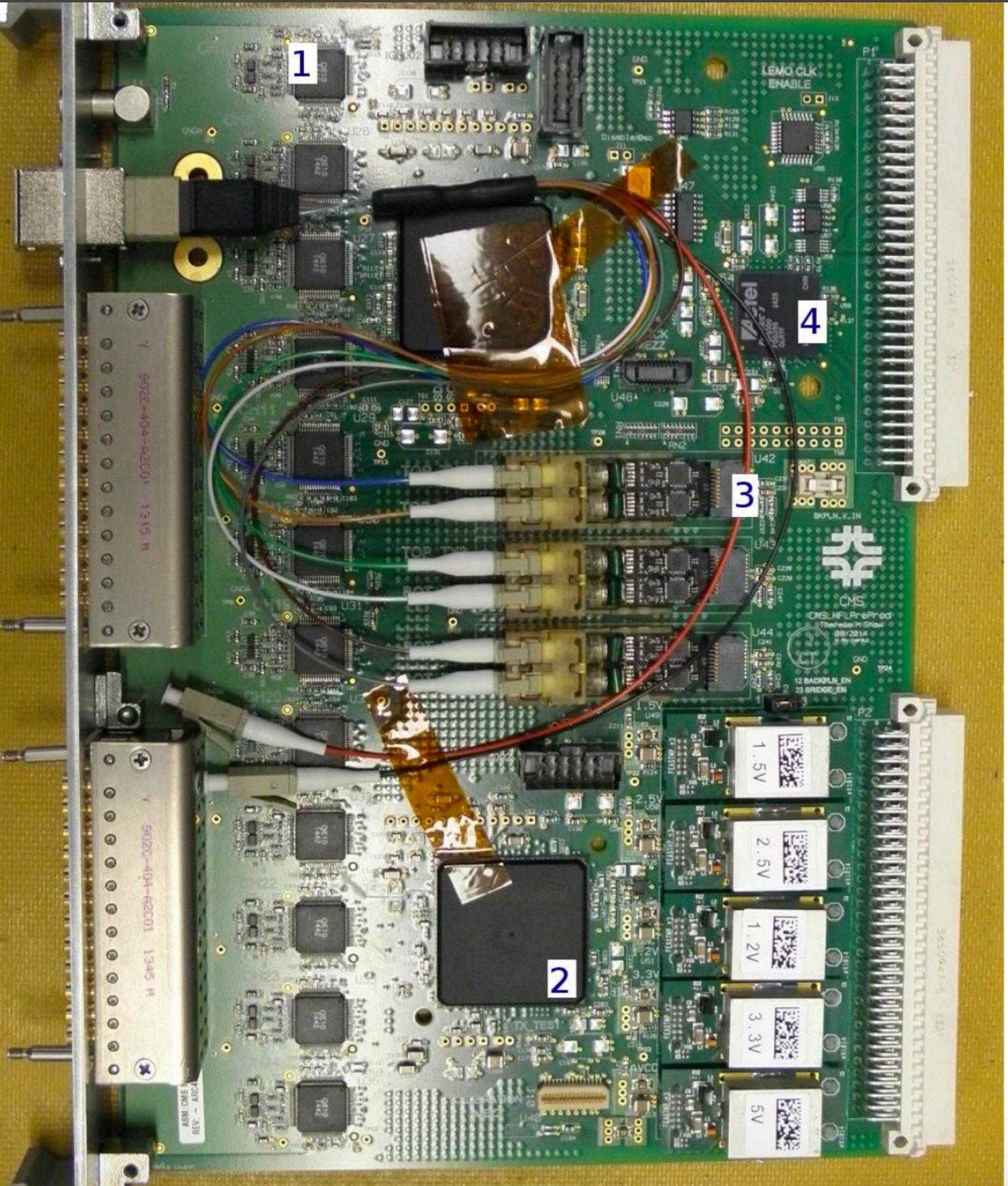
MCH
(μTCA controller
hub)

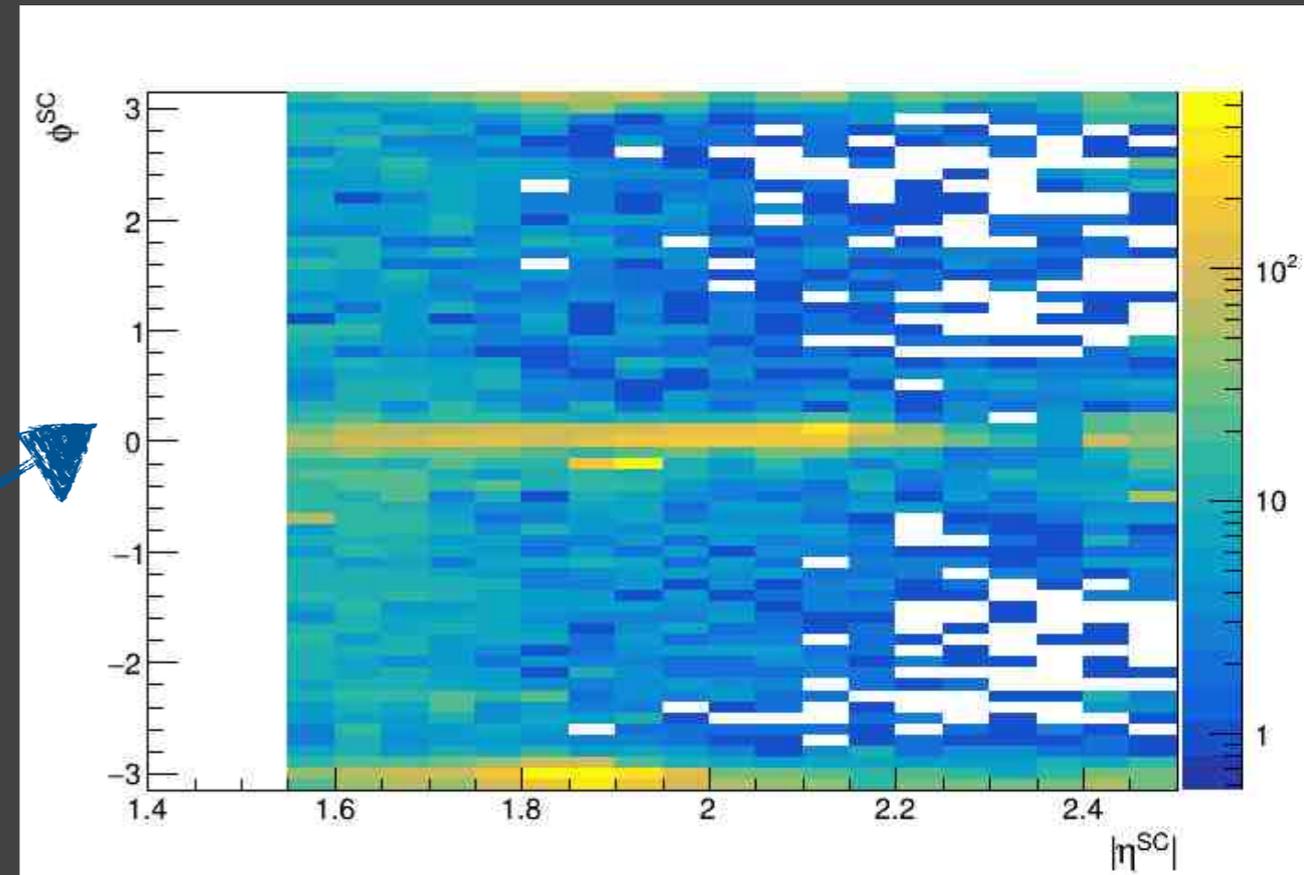
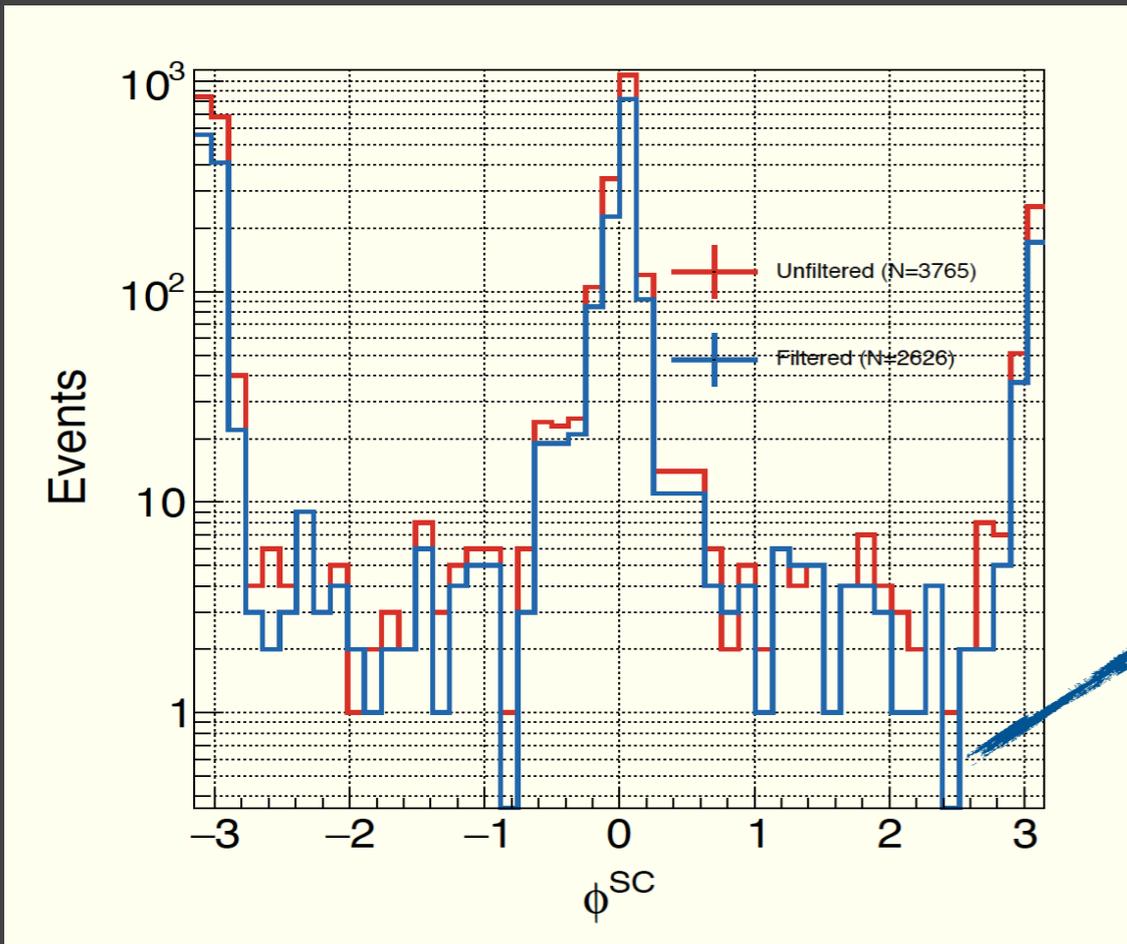
AMC13
(Advanced Mezzanine
card - provides clk and TTC to the
system)

QIE: Integrates and digitizes. It has TDC as well for timing information

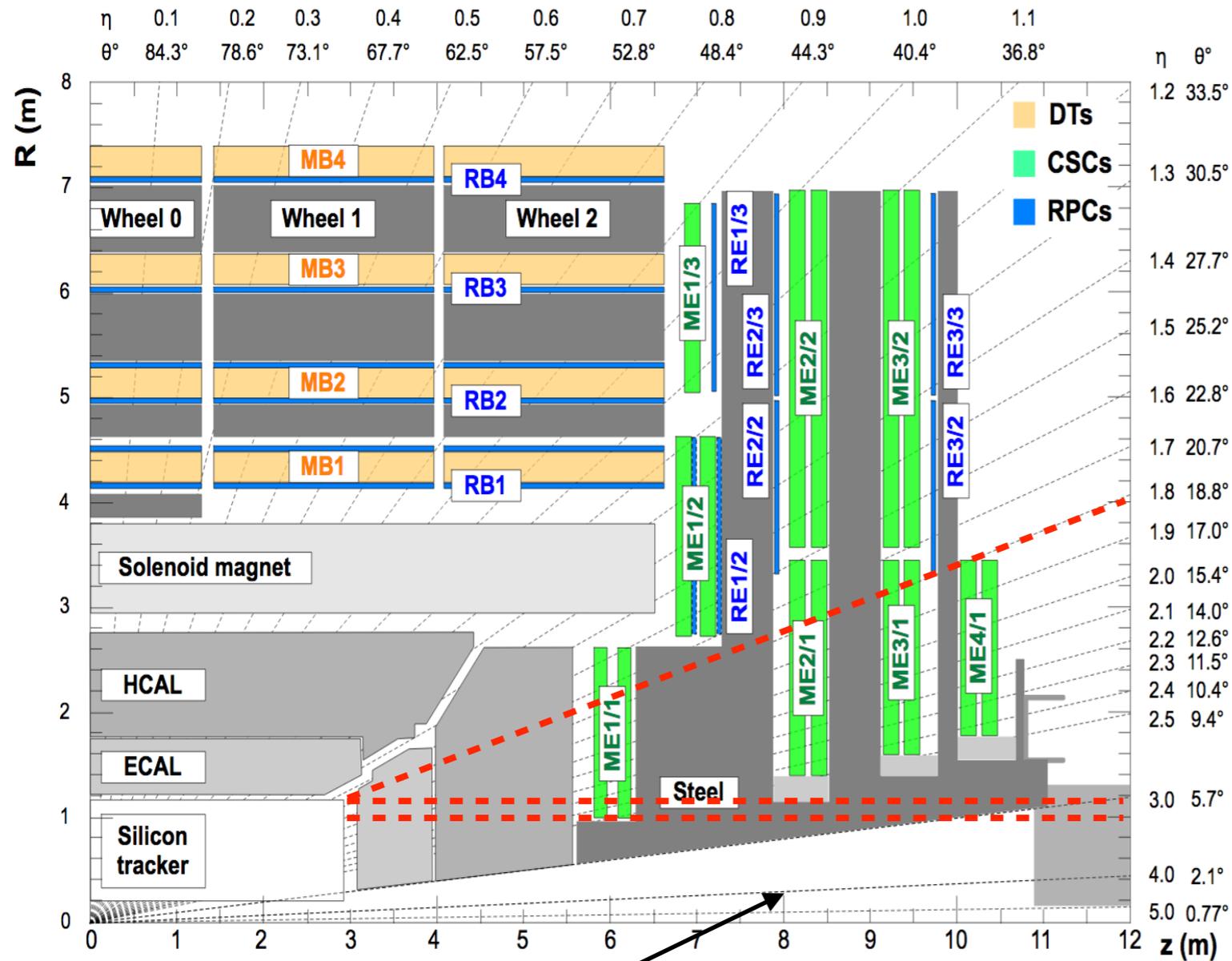


QIE readout board





- ▶ Most of the beam halo peaks $\sim \eta > 1.8$



- ▶ Muon chambers (CSCs) are mostly used to tag muons from the beam halo.
- ▶ However, a large region is not covered by the CSCs

Since beam halo travels parallel:
 All of the EE above $\eta = 1.8$ is not covered
 by the CSCs