#### **Asymptotic Freedom** & **High Derivative Gauge Theories**

#### **Manuel Asorey**





Departamento de Física Teórica **Universidad** Zaragoza

WORKSHOP ON SM AND BEYOND Corfu, August 2021

- SM is an effective field theory
- UV completion of Standard Model may involve higher derivatives (s = 0, s = 1/2, s = 1)

- SM is an effective field theory
- UV completion of Standard Model may involve higher derivatives (s = 0, s = 1/2, s = 1)
- HD give rise to renormalizable quantum gravity theories
- Inflationary models involving high derivative theories provide the best fits of the scalar/tensor relations

- SM is an effective field theory
- UV completion of Standard Model may involve higher derivatives (s = 0, s = 1/2, s = 1)
- HD give rise to renormalizable quantum gravity theories
- Inflationary models involving high derivative theories provide the best fits of the scalar/tensor relations
- Problems with ghosts, causality and unitarity

#### **High Derivative Gauge Theories**

$$S = \frac{1}{4g^2} \int d^4x F^a_{\mu\nu} F^{\mu\nu a} + \frac{1}{4g^2 \Lambda^{2n}} \int d^4x F^a_{\mu\nu} \Delta^n F^{\mu\nu a},$$

where

$$\Delta = d_A^* d_A + d_A d_A^*$$

is Hodge-covariant Laplacian operator

$$\Delta^{\nu b}_{\mu a} = -D^2 \delta^{\nu}_{\mu} \delta^{b}_{a} + 2f^{b}_{\ ca} F_{\mu}^{\nu c}$$

Instantons are minima in each topological sector

#### **Perturbation Theory**

#### **Gauge fixing [***α***-gauge**]

$$S_{\alpha} = \frac{\alpha}{2g^2 \Lambda^{2n}} \int d^4 x \, \partial^{\mu} A^a_{\mu} (-\partial^{\sigma} \partial_{\sigma})^n \partial^{\nu} A^a_{\nu}$$

#### **One-loop divergences**

$$\Gamma^{ab}_{\mu\nu}(p) = -c_n \frac{C_2(G)}{16\pi^2\epsilon} i\,\delta^{ab}\left(p^2\eta_{\mu\nu} - p_\mu p_\nu\right)$$

with

$$c_n = \frac{29}{3} - 23n + 5n^2$$
  $n \ge 2$ ,

$$c_1 = -\frac{43}{3}, \quad c_0 = \alpha - \frac{13}{3}$$

#### **Perturbation Theory**

#### **One-loop renormalization**

$$S_{\text{count}} = c_n \frac{C_2(G)}{128\pi^2} \left(\frac{2}{\varepsilon} + \log \frac{\Lambda_{\text{QCD}}^2}{\Lambda^2}\right) F^a_{\mu\nu} F^{\mu\nu a},$$

 $\beta$ -function of the coupling constant

$$\beta_n = c_n \frac{g^3 C_2(G)}{32\pi^2} \quad n \ge 2$$

Asymptotic freedom only for n = 0, 1, 2, 3, 4

$$\widetilde{\mathbf{c}}_{0} = -\frac{22}{3}, \mathbf{c}_{1} = -\frac{43}{3}, \mathbf{c}_{2} = -\frac{49}{3}, \mathbf{c}_{3} = -\frac{43}{3}, \mathbf{c}_{4} = -\frac{7}{3}, \mathbf{c}_{5} = \frac{59}{3}$$

#### **Perturbation Theory**

#### **One-loop form factor**

$$\Gamma^{ab}_{\mu\nu}(p) = -\frac{C_2(G)}{32\pi^2} i\delta^{ab} \left(p^2 \eta_{\mu\nu} - p_{\mu}p_{\nu}\right) \Pi(p^2)$$

with

$$\Pi(p^2) = \left(b_n \log \frac{p^2 + \Lambda^2}{\Lambda^2} + c_0 \log \frac{p^2}{\Lambda_{\rm QCD}^2}\right),\,$$

$$b_n = 14 - \alpha - 23n + 5n^2$$
 for  $n \ge 2$   
 $b_0 = 0$ ,  $b_1 = -10 - \alpha$ 

## **Scaling Regimes**

There are two different asymptotic regimes with two different beta functions:

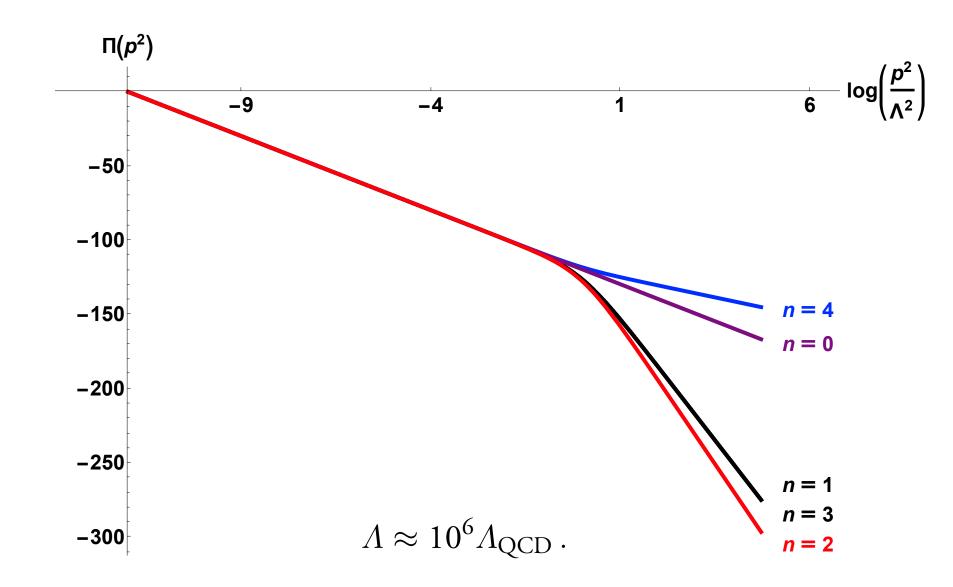
• UV regime  $p \gg \Lambda$ 

$$\beta_{\rm UV} = c_n \frac{g^3 C_2(G)}{32\pi^2}$$

• **IR regime**  $\Lambda_{\text{QCD}}$ 

$$\beta_{\rm IR} = -\frac{22}{3} \frac{g^3 C_2(G)}{32\pi^2}$$

#### **Two-point form factor**



#### Generalization

Replace the Hodge-covariant Laplacian operator by a generalized Laplacian

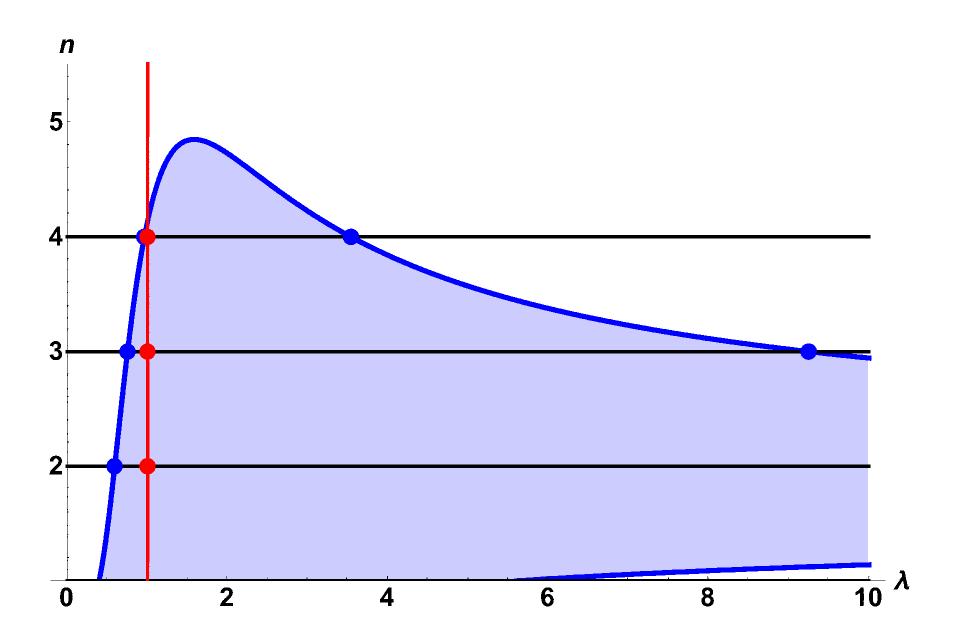
 $\Delta \Rightarrow {}^{\lambda}\Delta$ 

$${}^{\lambda}\Delta^{\nu b}_{\mu a} = -\delta^{b}_{a}\,\delta^{\nu}_{\mu}\,D^{2} + 2\,\lambda f^{b}_{\ ca}\,F_{\mu}{}^{\nu c}$$

**One-loop** β-function coefficients

$$c_{n} = -\frac{7}{3} + 5n + 4n^{2} - (4 + 10n + 4n^{2}) \lambda + (16 - 18n + 5n^{2}) \lambda^{2}$$
$$n \ge 2$$
$$c_{1} = \frac{38}{3} - 18\lambda - 9\lambda^{2}$$

## **Two-point form factor**



## **UV Finite theories**

The range of asymptotically free theories is more restrictive for  $\lambda \neq 1$ 

For some values of  $\lambda \neq 1$  the theory is finite

$$n = 1 \quad \lambda_1 = -2.55 \quad \lambda_2 = 0.55$$
  

$$n = 2 \quad \lambda = 0.59$$
  

$$n = 3 \quad \lambda_1 = 0.75 \quad \lambda_2 = 9.25$$
  

$$n = 4 \quad \lambda_1 = 0.96 \quad \lambda_2 = 3.54$$

The theory is free of UV divergences

No instanton solutions  $\Rightarrow$  new QCD potentials for axions

## **Higher Derivative Ghosts**

- BRST symmetry is preserved
- High derivative terms are not renormalized
- $\beta_{\Lambda} = -\frac{1}{n}\beta_{g}$
- Effective ghost masses run to infinite if  $\beta_{\sigma} \neq 0$
- Unitarity and Causality might be recovered in UV

M.A., F. Falceto and L. Rachwal JHEP 05(2021)075

 Asymptotic freedom is a very stringent condition on UV completion of gauge theories

- Asymptotic freedom is a very stringent condition on UV completion of gauge theories
- Only theories with less than eight extra derivatives can be UV asymptotically free

- Asymptotic freedom is a very stringent condition on UV completion of gauge theories
- Only theories with less than eight extra derivatives can be UV asymptotically free
- The effective masses of extra ghosts fields run to infinity under RG flow

- Asymptotic freedom is a very stringent condition on UV completion of gauge theories
- Only theories with less than eight extra derivatives can be UV asymptotically free
- The effective masses of extra ghosts fields run to infinity under RG flow
- Unitarity and Causality might be recovered

- Asymptotic freedom is a very stringent condition on UV completion of gauge theories
- Only theories with less than eight extra derivatives can be UV asymptotically free
- The effective masses of extra ghosts fields run to infinity under RG flow
- Unitarity and Causality might be recovered
- Implications for high derivative theories of quantum gravity

- Asymptotic freedom is a very stringent condition on UV completion of gauge theories
- Only theories with less than eight extra derivatives can be UV asymptotically free
- The effective masses of extra ghosts fields run to infinity under RG flow
- Unitarity and Causality might be recovered
- Implications for high derivative theories of quantum gravity
- Are there similar phenomena in gravity?