

INVISIBLE TRACES OF CONFORMAL SYMMETRY BREAKING

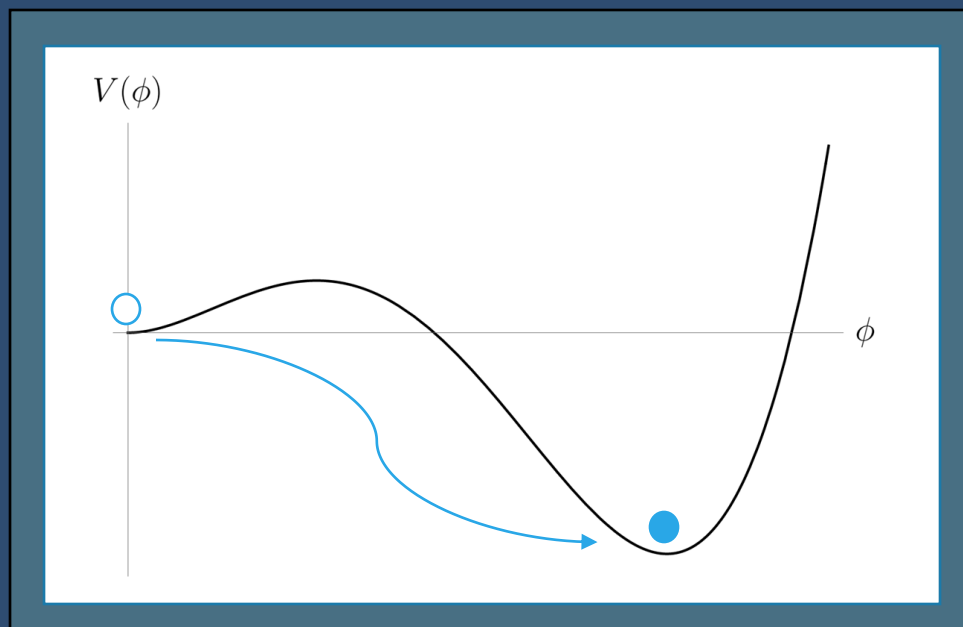
Maciej Kierkla

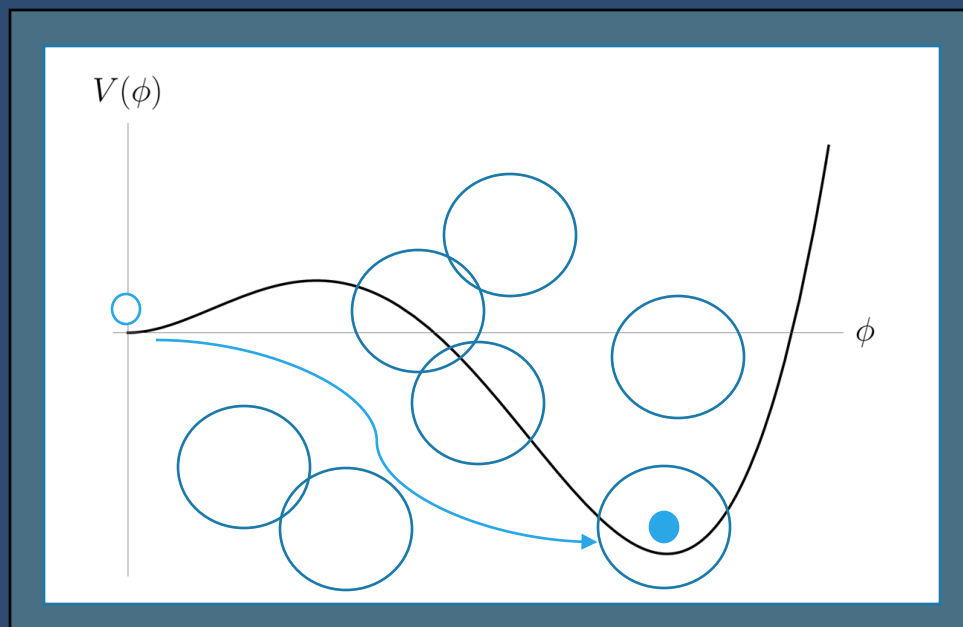
Faculty of Physics, University of Warsaw,

in collaboration with Alexandros Karam, Bogumiła Świeżewska.



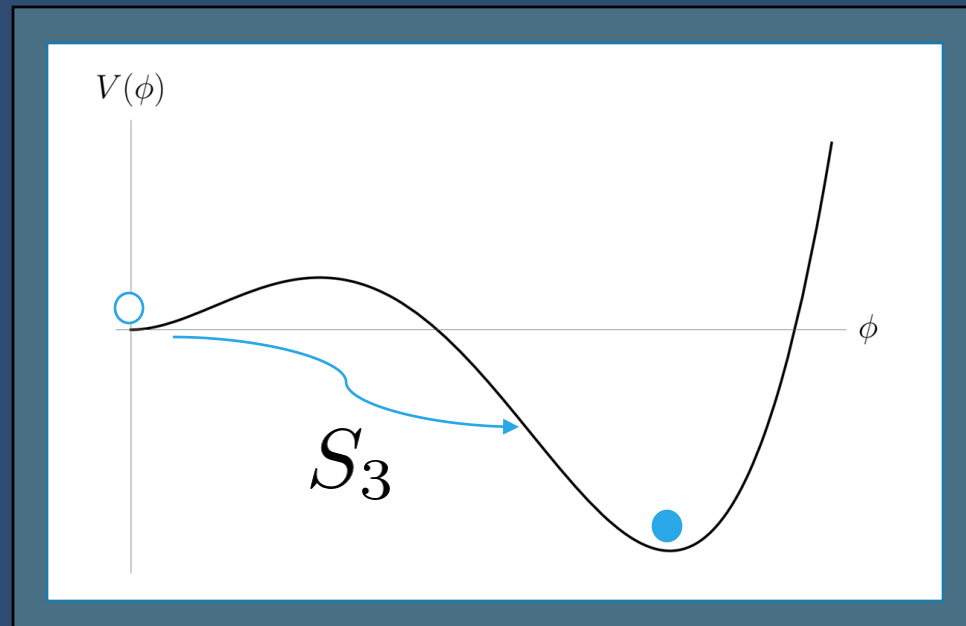






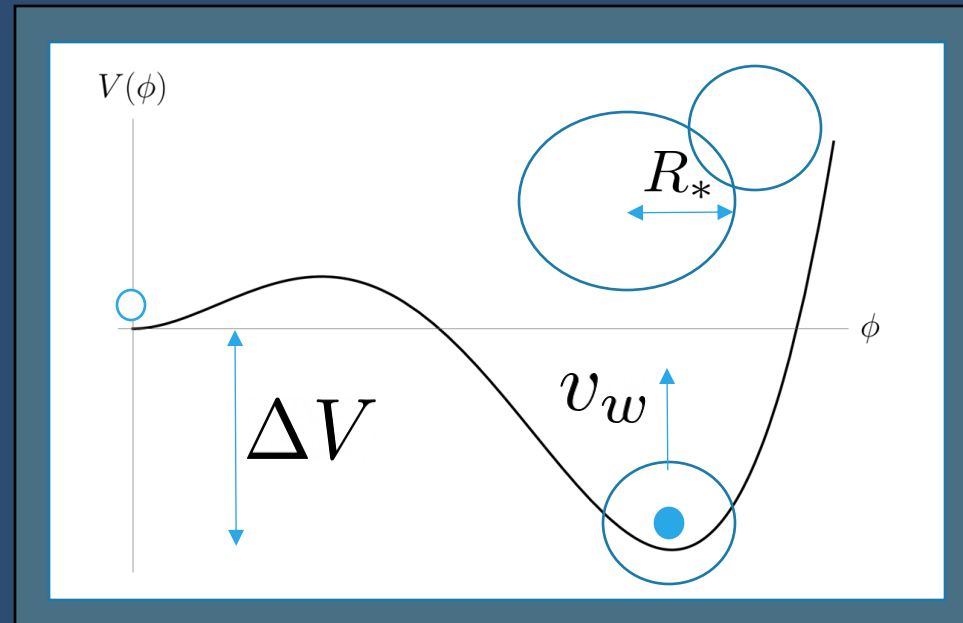
Thermal tunnelling

- Decay rate is given as: $\Gamma(T) \simeq T^4 \left(\frac{S_3}{2\pi T} \right)^{3/2} e^{-S_3/T}$
- The action of the field configuration $S_3 = 4\pi \int r^2 dr \left[\frac{1}{2} \left(\frac{d\varphi}{dr} \right)^2 + V_{\text{eff}}(\varphi, T) \right]$



Cosmological phase transitions

- Strength of the transition $\alpha \sim \frac{\Delta V}{\rho_{\text{rad}}(T_p)}$
- Average bubble radius R_* or time scale $\beta \sim R_*^{-1}$
- Bubble wall velocity v_w



Model



Phase Transition
Properties



GW power
spectrum
 $\Omega_{\text{GW}}h^2$

$SU(2)_cSM$



$$V_{\text{tree}} = \frac{1}{4} (\lambda_1 h^4 + \lambda_2 h^2 \phi^2 + \lambda_3 \phi^4)$$

Classical
conformal symmetry

All masses generated
via Coleman-
Weinberg mechanism

Hierarchy
problem alleviated



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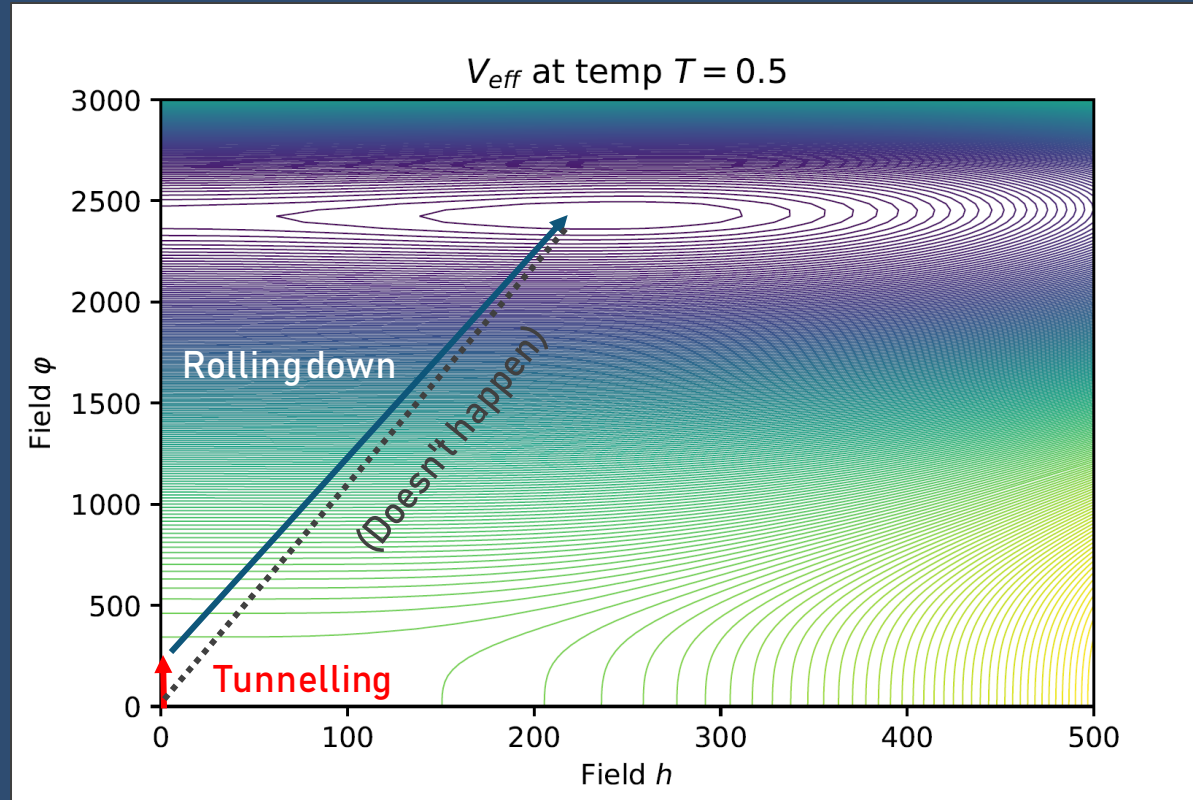
Perturbative and stable
up to the Planck scale

Vector DM candidate,
gauge boson X

Exhibits supercooling

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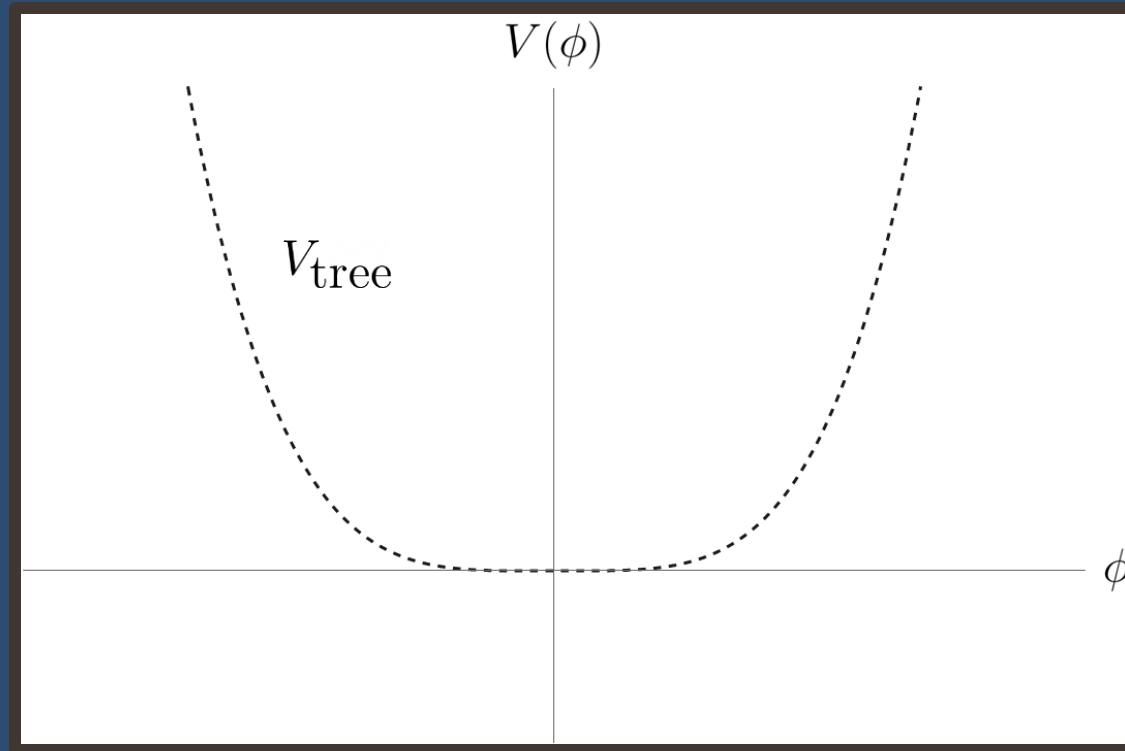
Tunneling scenario in SU(2)cSM



Tunnelling occurs only in the new scalar direction!

Coleman-Weinberg Mechanism

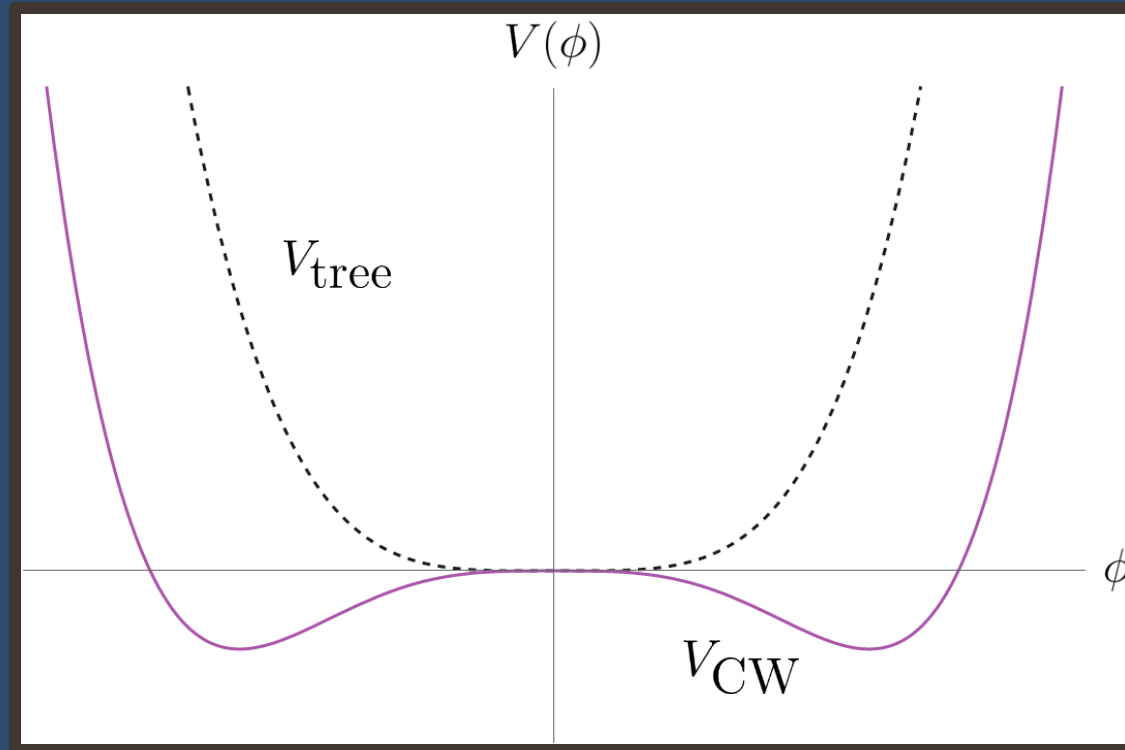
Example: massless
scalar electrodynamics.



$$V_{CW} = \underbrace{\frac{\lambda}{4!} \phi^4}_{V_{\text{tree}}}$$

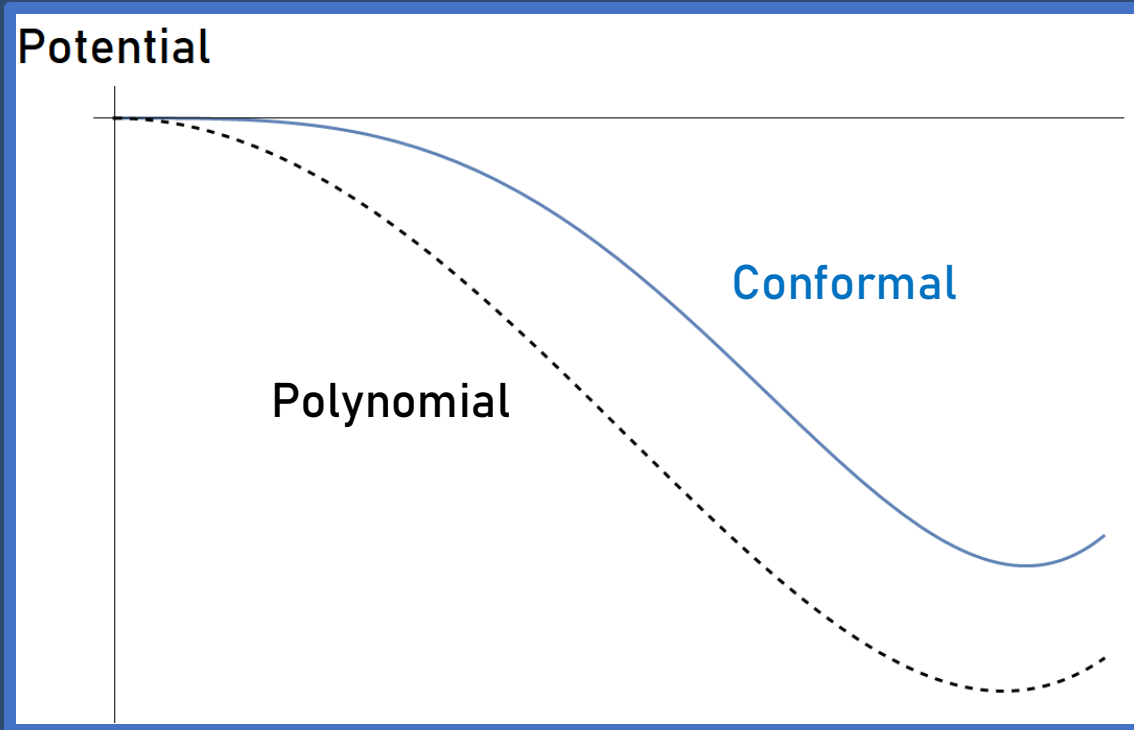
Coleman-Weinberg Mechanism

Example: massless scalar electrodynamics.



$$V_{CW} = \underbrace{\frac{\lambda}{4!} \phi^4}_{V_{\text{tree}}} + \underbrace{\frac{\lambda^4}{64\pi^2} \phi_c^4 \left(\ln \frac{\phi_c^2}{\mu^2} - \frac{3}{2} \right)}_{\text{"scalar" correction}} + \underbrace{\frac{3e^4}{64\pi^2} \phi_c^4 \left(\ln \frac{\phi_c^2}{\mu^2} - \frac{5}{6} \right)}_{\text{"boson" correction}}$$

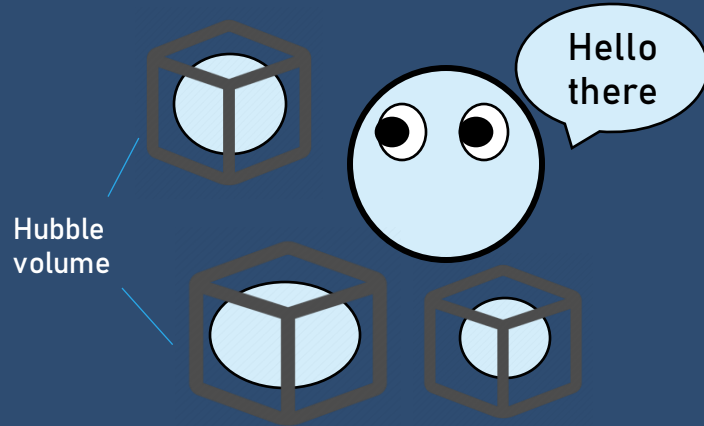
Introducing: supercooling



Features:

- phase transition happens at temperatures significantly below EW scale,
- thermally produced barrier lasts till $T=0$,
- Induces strong Gravitational Wave signal.

- Nucleation temperature



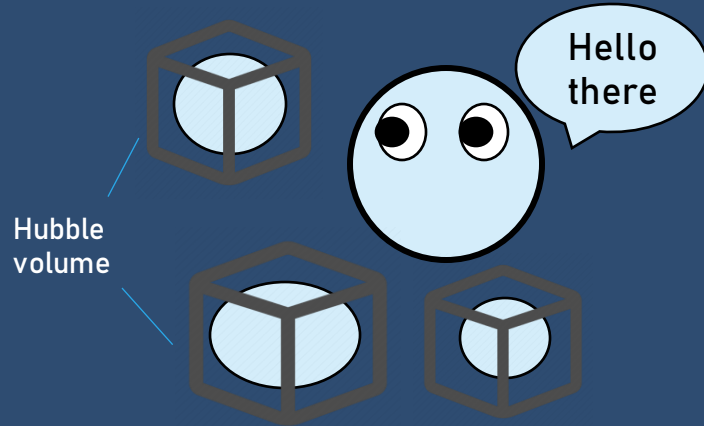
We calculate:
$$N(T_n) = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

One can also use an approximation:

But not this one
$$\frac{S_3}{T} \simeq 140$$

$$\frac{\Gamma(T_n)}{H(T_n)^4} \simeq 1 \Rightarrow \frac{S_3}{T_n} = 4 \log \left(\frac{T_n}{H(T_n)} \right)$$

- Nucleation temperature

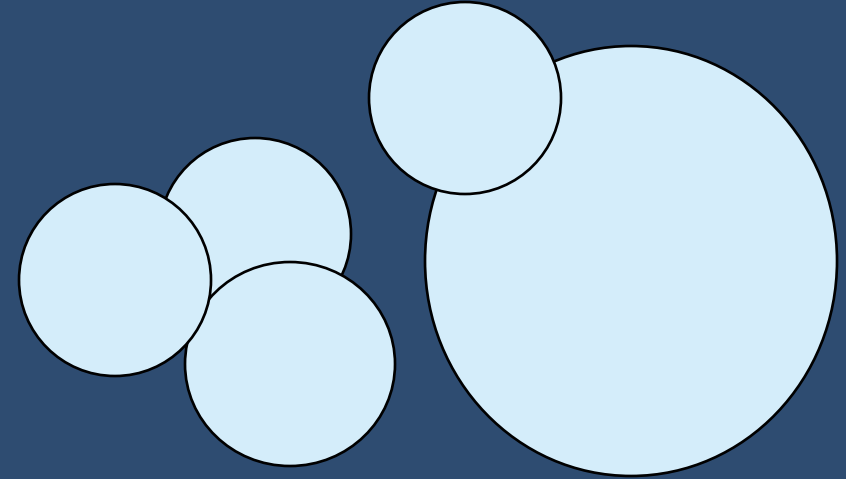


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- Percolation temperature



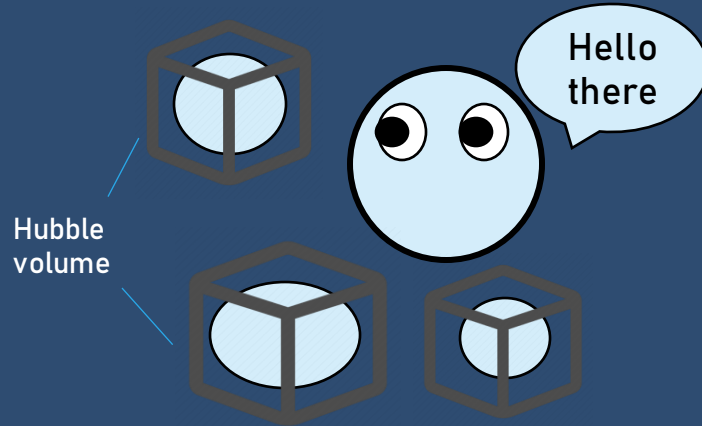
Probability of point still in false vacuum is $P = e^{-I(T)}$, where

$I(T)$ is the volume converted into true vacuum

Then we solve for condition:

$$I(T_p) \simeq 0.34$$

- Nucleation temperature



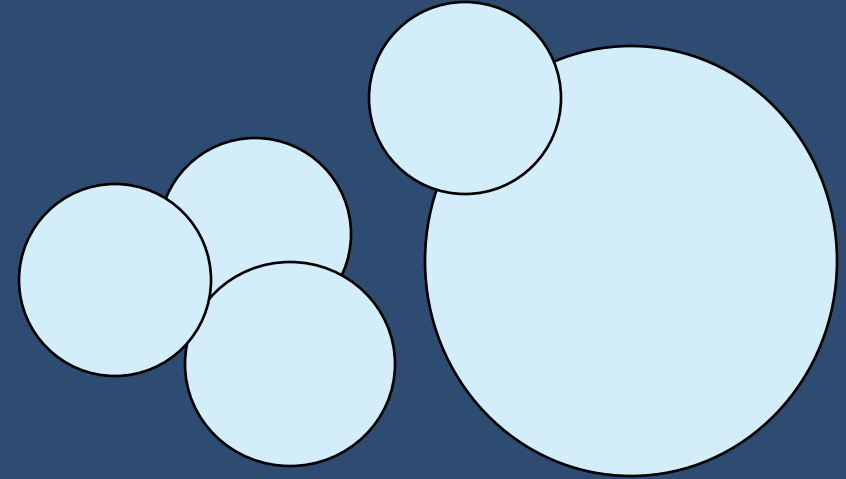
Are not equal
in models with
supercooling

We calculate:
$$N(T_n) = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

One can also use an approximation:

$$\frac{\Gamma(T_n)}{H(T_n)^4} \simeq 1 \Rightarrow \frac{S_3}{T_n} = 4 \log \left(\frac{T_n}{H(T_n)} \right)$$

- Percolation temperature



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$$\Omega_{\text{GW}} = \Omega_{\text{collisions}} + \Omega_{\text{sound waves}} + \Omega_{\text{turbulence}}$$

How do we know which source dominates?

Efficiency factors:

$$\kappa_{col} = \frac{E_{wall}}{E_V}$$

$$\kappa_{sw} \sim 1 - \kappa_{col}$$

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Where the energy goes?

There is a lot of friction

Bubble expansion accelerates

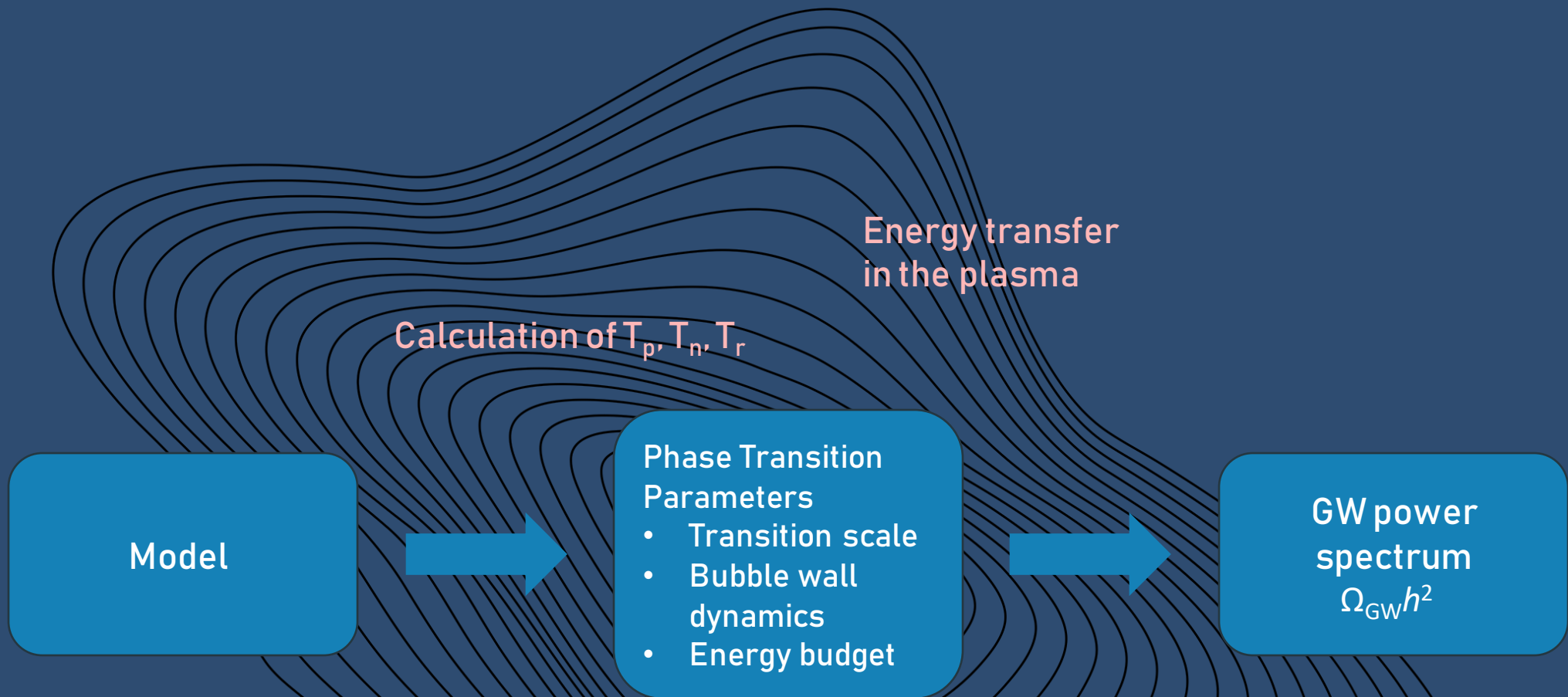
Energy is dissipated in the surrounding plasma

Energy goes to the bubble's wall

And the main GW source is...

Sound waves + Turbulences

Bubble collisions



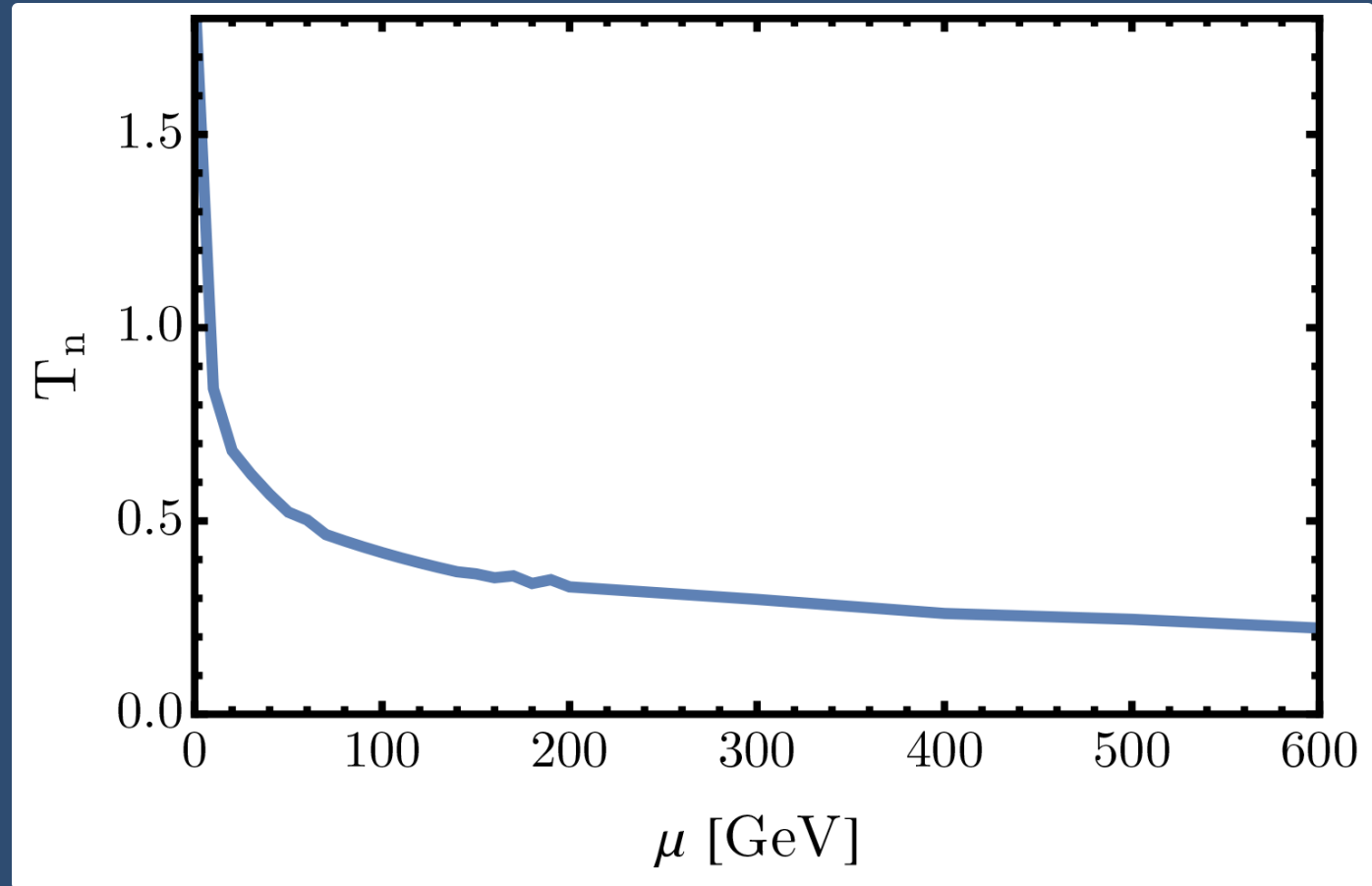
Calculation of T_p, T_n, T_r

Energy transfer
in the plasma

RG scale dependence

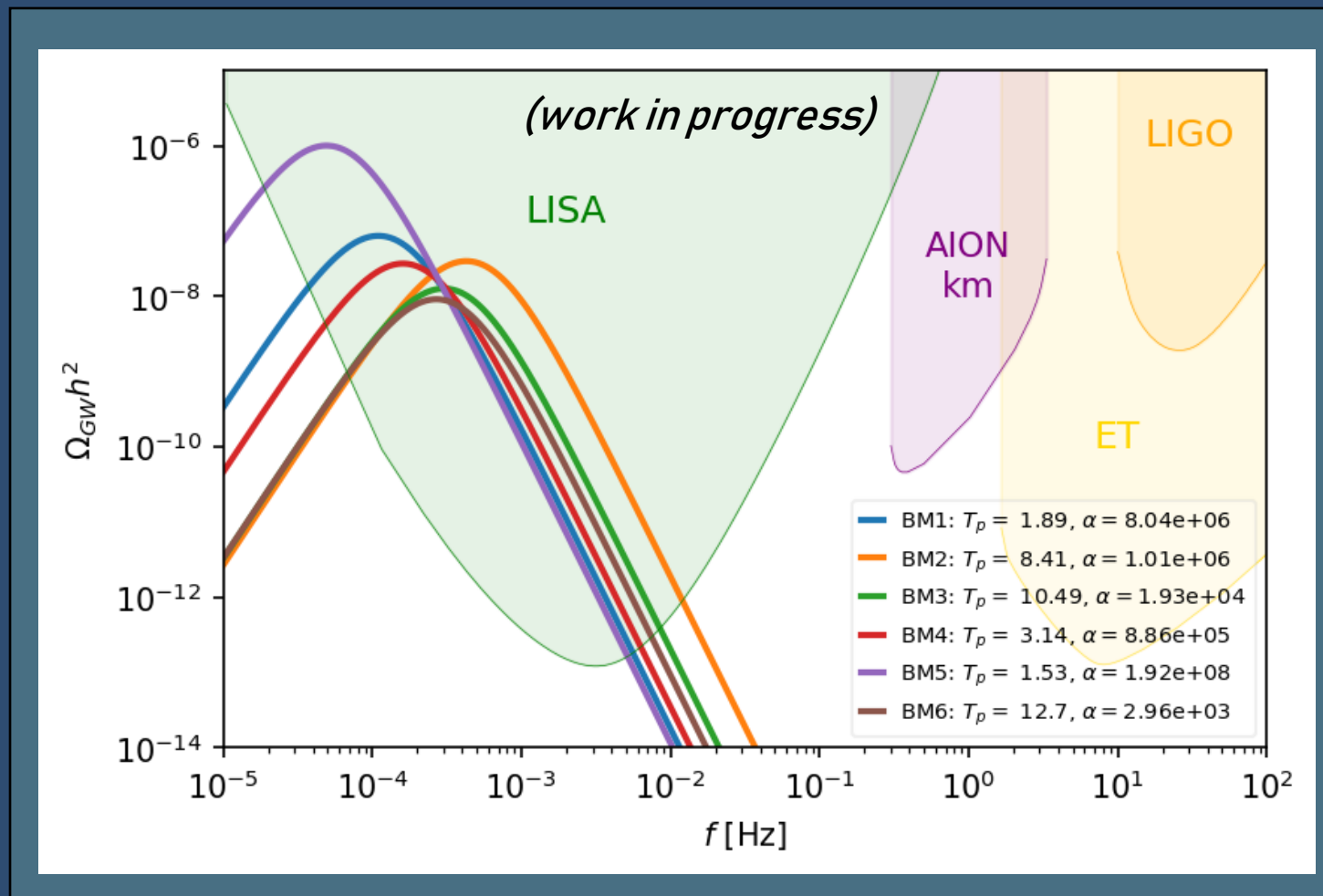
Efficiency factors i.e
inclusion of all
possible sources

- RG scale dependence



This affects other parameters and therefore the resulting spectrum.

Gravitational Waves spectra in SU(2)cSM



Goal: provide accurate predictions for LISA.

Thank you



NARODOWE CENTRUM NAUKI