

Probing leptogenesis with gravitational waves

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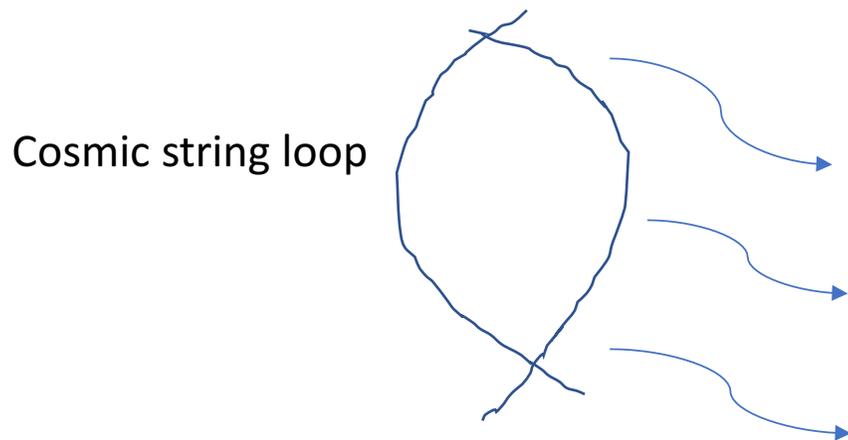
Collab: Pasquale D. Bari, Satyabrata Datta, Ambar Ghosal, Sabir Ramazanov, Federico Urban

Subject of this talk:

How to probe **high scale** leptogenesis with gravitational waves.

High scales => beyond the reach of the collider experiments, therefore a shift of attention to GWs.

Typically the heavy neutrinos which are involved in generating lepton asymmetry in various ways, get masses with a dynamical U(1) gauge symmetry breaking. => associated with cosmic strings

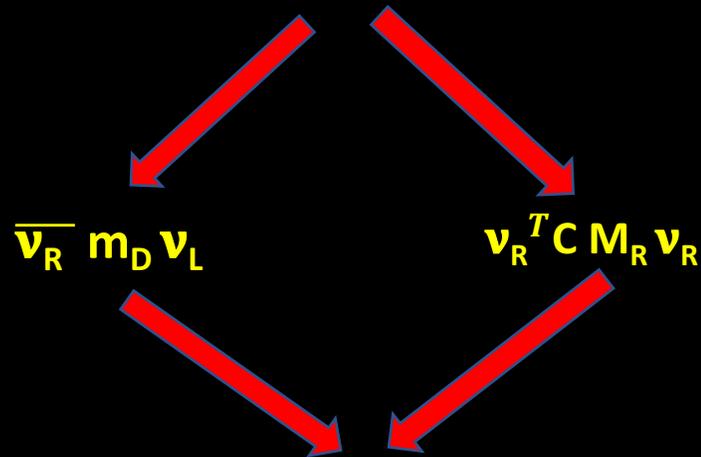


Gravitational waves, Amplitude $\sim [\text{VEV}_{U(1)}]^2$

Heavy mass: $M_R \sim f(\text{VEV}_{U(1)})$

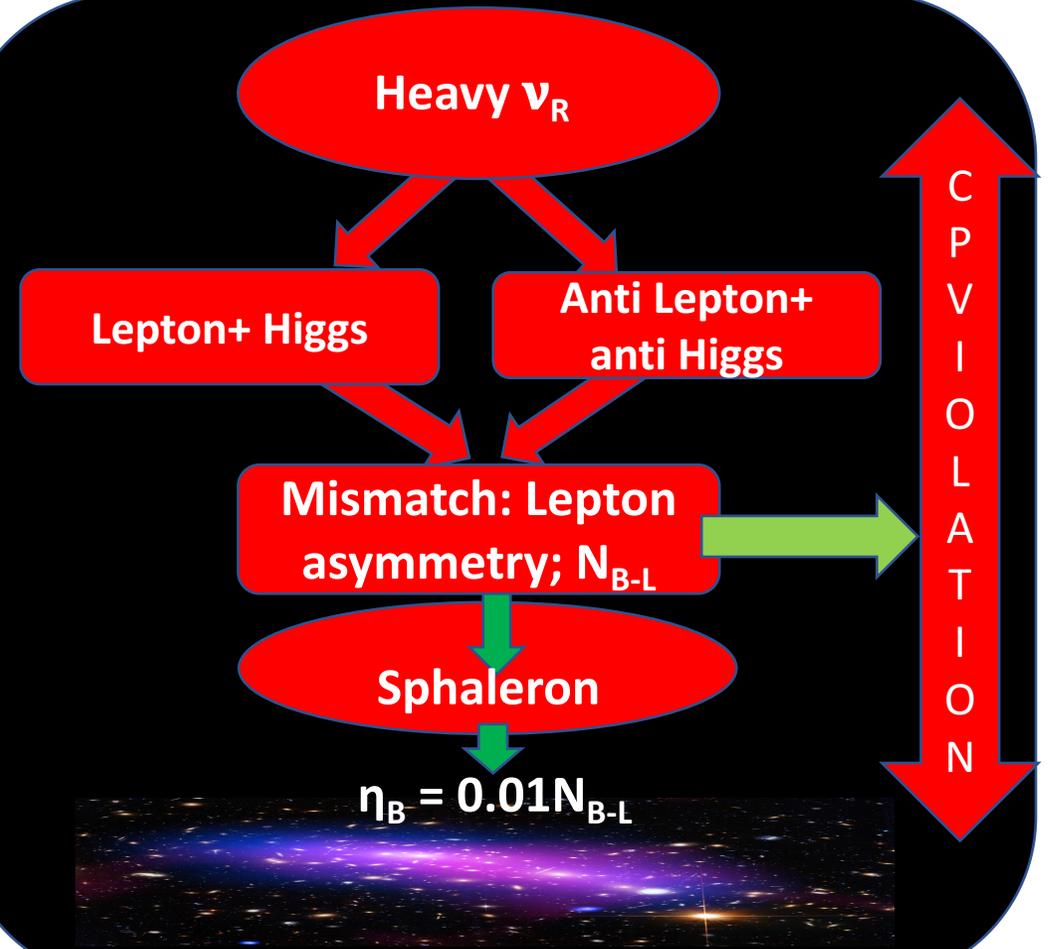
Basic idea to reconcile light neutrino masses and baryogenesis via leptogenesis

Minimal scenario: Introduce two RH neutrino field ν_{Ri}



Type -1 Seesaw : $m_\nu = m_D^T M_R^{-1} m_D$

$M_R \sim 10^{14} \text{ GeV} \Rightarrow m_\nu \approx 0.1 \text{ eV}$
 Light neutrinos are Majorana type



Tests:

Direct: Colliders => low scale leptogenesis (Temp < TeV).

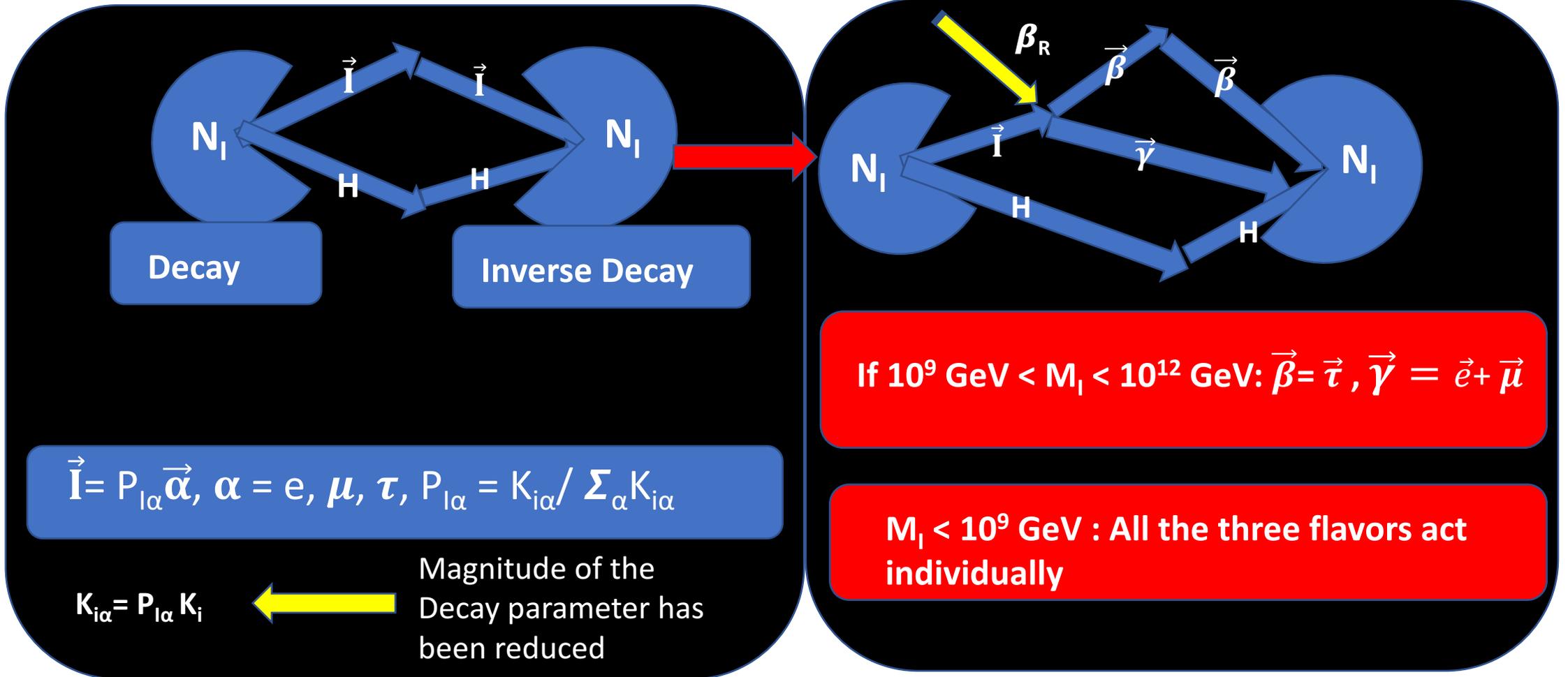
Examples (Not exhaustive):

ARS: (Heavy neutrino oscillation $N_j \leftrightarrow N_i$) (**Akhmedov, Rubakov, Smirnov, 1998**).

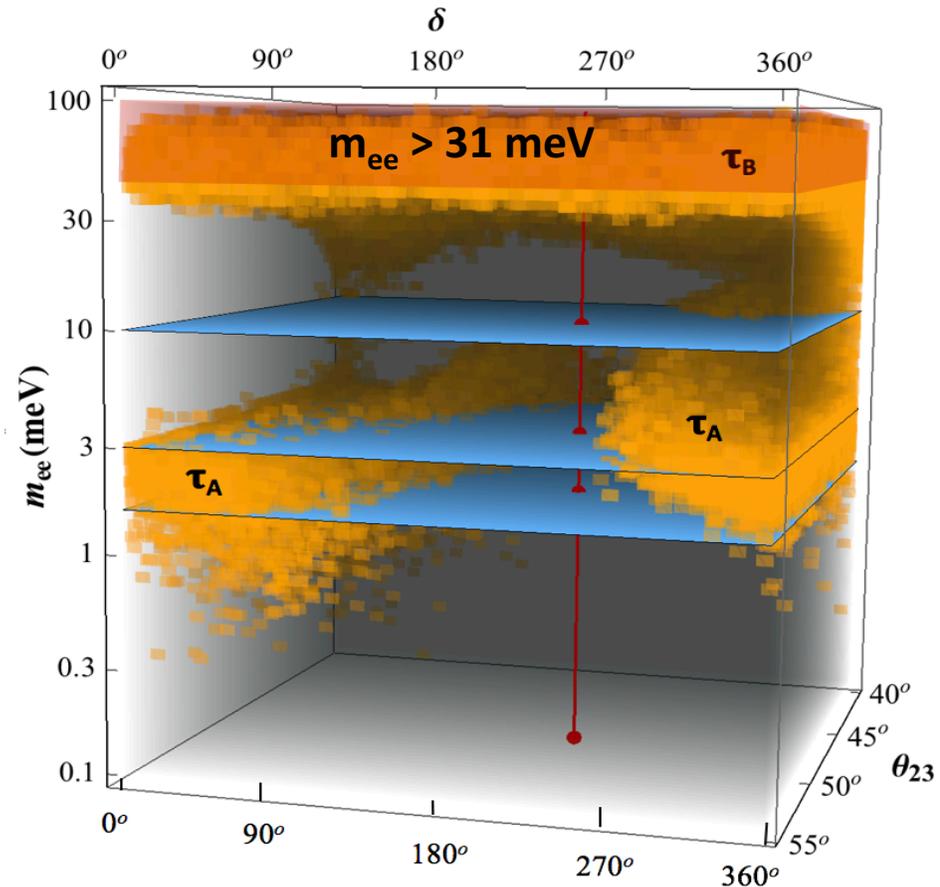
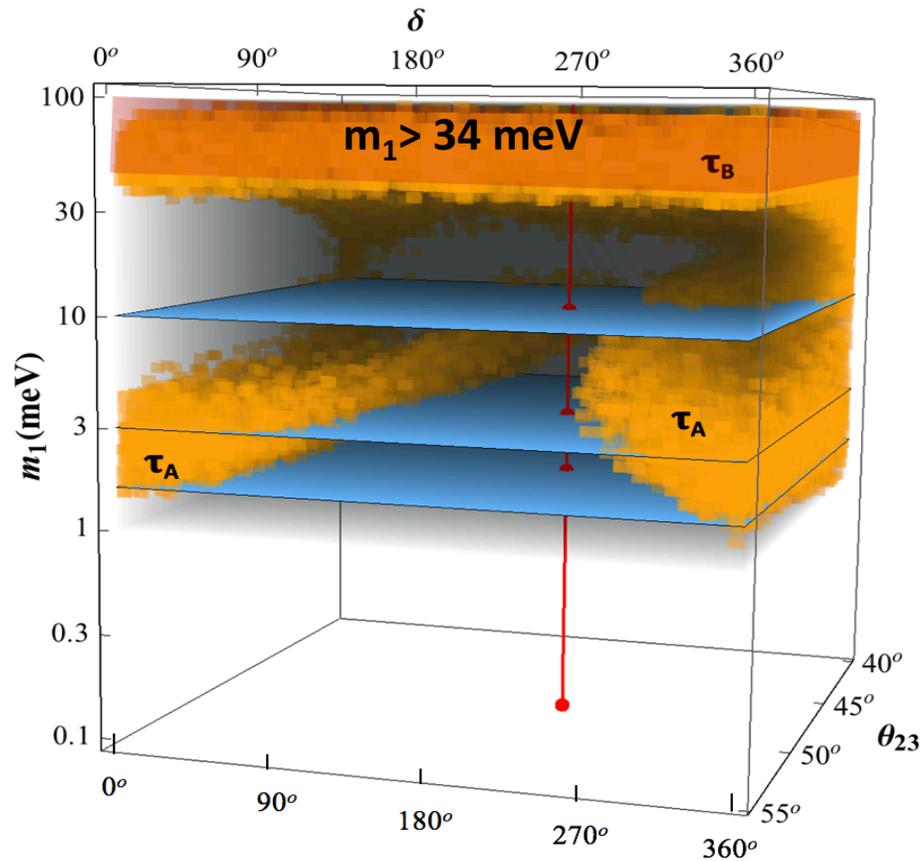
Resonant-leptogenesis (Quasi-degenerate heavy neutrinos enhancement in the CP asymmetry) (**Pfaltsis, Underwood, 2004**).

Higgs decays: (**Hambye, Teresi, 2016**)

Flavour effects : Reduction of N1-washout, Still can not lower the scale significantly but opens up the way of indirect tests in neutrino experiments:



3D Parameter space : Leptogenesis in SO(10) inspired models-A Prototype



KAMLAND2-
ZEN reach:
20 meV

Best fit: ($\theta_{23} \sim 48.6^\circ$, $\delta \sim 221^\circ$)

Di. Bari and Samanta, 2020

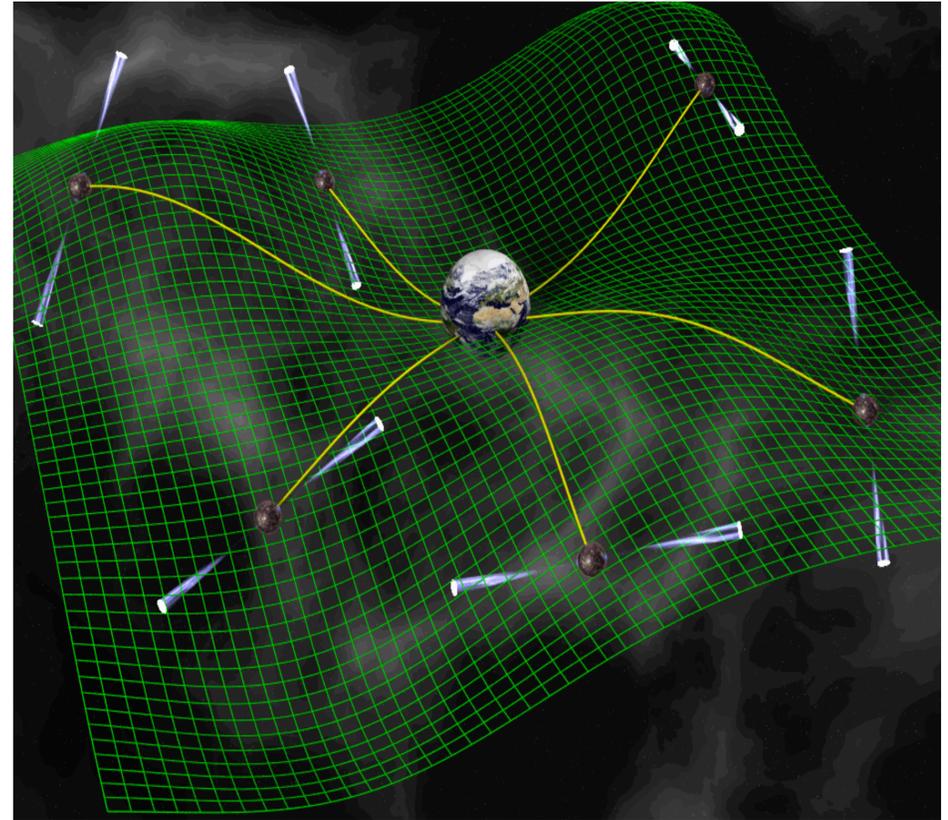
GWs as a probe of leptogenesis , e.g., PTAs

Millisecond pulsars (spins ~100 times a second) produce most stable pulses and are used by the PTAs

When a gravitational wave (a disturbance) passes between the earth and pulsar system, the time of arrival of the signal from the pulsars changes. This induces a change in frequency due to the gravitational wave.

Time residual:

$$R(t) = - \int_0^t \frac{\delta v}{v} dt$$



Pulsar-Timing-Arrays typically work with high amplitude GWs => Could be a Detector of High Scale Symmetry breaking theories, e.g., leptogenesis

Cosmic strings

Cosmological phase transition leads to spontaneous breaking of abelian symmetry
This is associated with topological defect like cosmic strings (CS). (T. Kibble, J. Phys. A 9 (1976)).

They can form close loop and shrink via GW emission.

Still there are debates whether they loose energy via particle radiation or GW emission.

Recent numerical simulation based Nambu-Goto action shows that the dominant emission is GW if the broken symmetry is a local gauge symmetry. (Pillado et al, PRD 2011, T. Vachaspati et al ,PRD, C. Ringeval et al JCAP).

Radiate energy at constant rate $\frac{dE}{dt} = -\Gamma G\mu^2,$ Length dynamics $l(t) = \alpha t_i - \Gamma G\mu(t - t_i)$

G: Newton constant , μ : string tension (~**Square of symmetry breaking scale**), α : loop size (max: 0.1)

Gravitational waves power spectrum and loop number density

Amplitude/energy density $\Omega_{GW}(t_0, f) = \frac{f}{\rho_c} \frac{d\rho_{GW}}{df} = \sum_k \Omega_{GW}^{(k)}(t_0, f).$

Differential energy density $\frac{d\rho_{GW}^{(k)}}{df} = \int_{t_F}^{t_0} \left[\frac{a(\tilde{t})}{a(t_0)} \right]^4 P_{GW}(\tilde{t}, f_k) \frac{dF}{df} d\tilde{t},$

Power spectrum

Amplitude/energy density

$$\Omega_{GW}^{(k)}(t_0, f) = \frac{2kG\mu^2\Gamma_k}{f\rho_c} \int_{t_{osc}}^{t_0} \left[\frac{a(\tilde{t})}{a(t_0)} \right]^5 n\left(\tilde{t}, \frac{2k}{f} \left[\frac{a(\tilde{t})}{a(t_0)} \right]\right) d\tilde{t}.$$

$$P_{GW}(\tilde{t}, f_k) = \frac{2kG\mu^2\Gamma_k}{f_k^2} n(\tilde{t}, f_k) = \frac{2kG\mu^2\Gamma_k}{f^2 \left[\frac{a(t_0)}{a(\tilde{t})} \right]^2} n\left(\tilde{t}, \frac{2k}{f} \left[\frac{a(\tilde{t})}{a(t_0)} \right]\right).$$

Loop number density

μ^2/M_{pl}

Numerical simulation: $n(\tilde{t}, l_k(\tilde{t})) = \frac{0.18}{[l_k(\tilde{t}) + \Gamma G\mu\tilde{t}]^{5/2} \tilde{t}^{3/2}}.$

Cosmic archeology, GW spectral shapes and Leptogenesis

Amplitude sensitivity

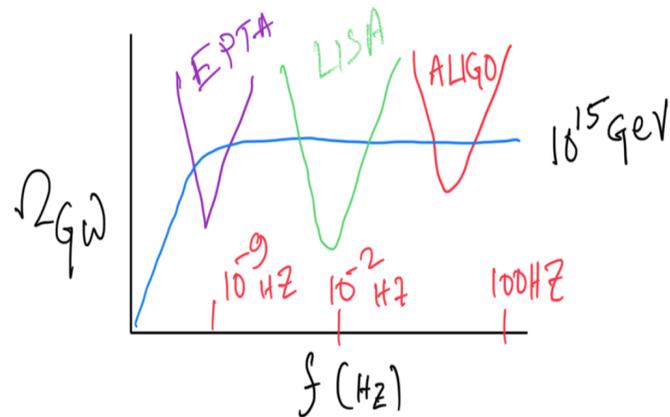
Amplitude + spectral shape sensitivity

Standard Cosmology ($w=1/3$)

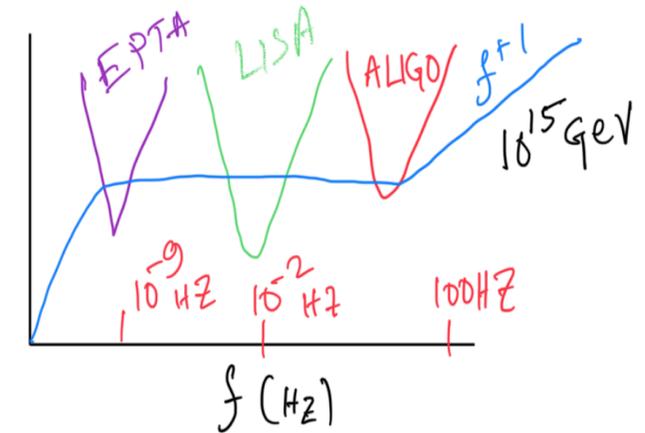
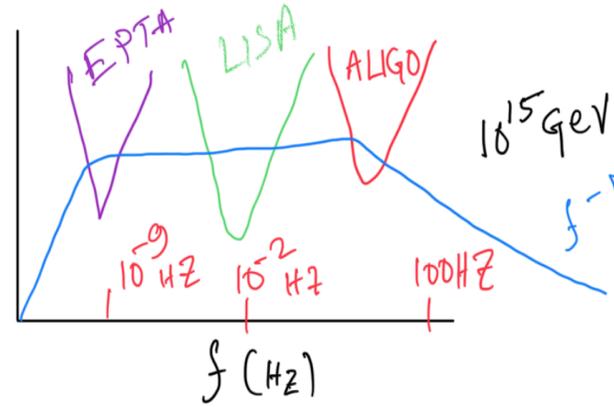
Early Matter domination ($w=0$)

Kination ($w=1$)

Fundamental mode ($k=1$):



A spectral break (f_*, T_*)



Standard Leptogenesis

Murayama et al PRL(2020)
Samanta et al JHEP(2020)

Leptogenesis from Black holes

Samanta et al JCAP(2021)

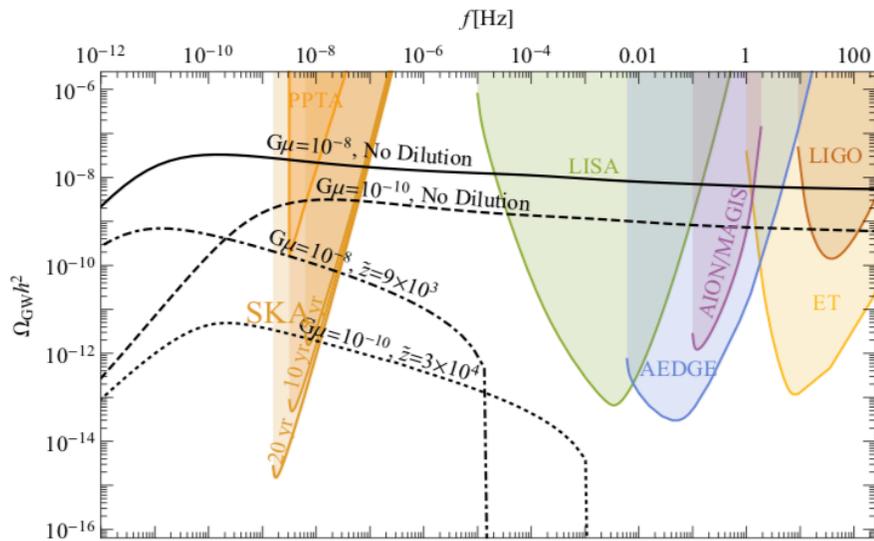
Gravitational Leptogenesis

Samanta et al 2108.08359 (2021)

Some non-standard spectrums

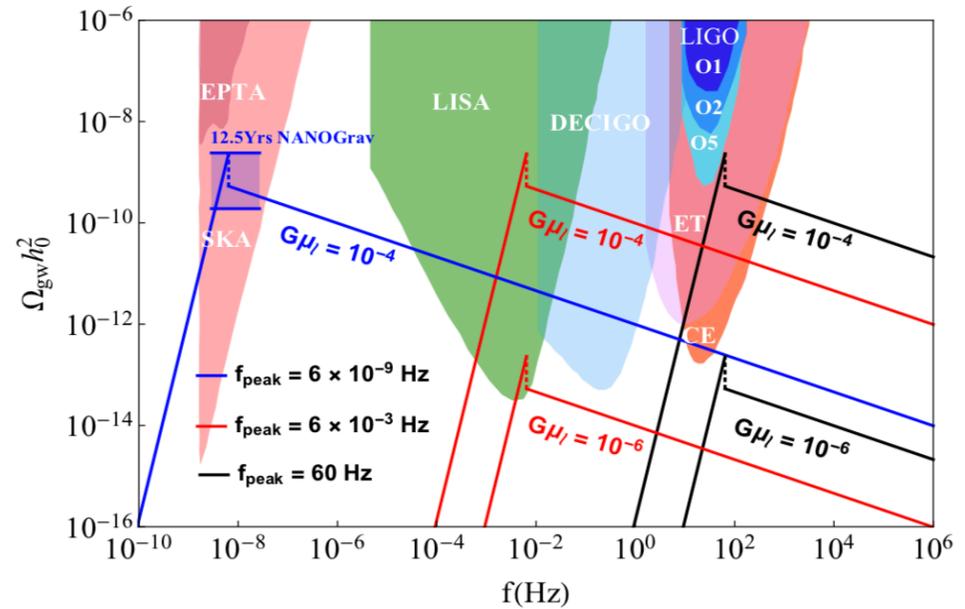
Cosmic strings with inflation dilution

Cui, Lewicki, Morrissey, PRL, 2020



Melting cosmic strings

Emond, Ramazanov, Samanta (2108.05377)



Cosmic archeology, GW spectral shapes and Leptogenesis

G. Shore et al, JHEP 10 025, (2020) **Samanta** et al JHEP 12 (2020) 67

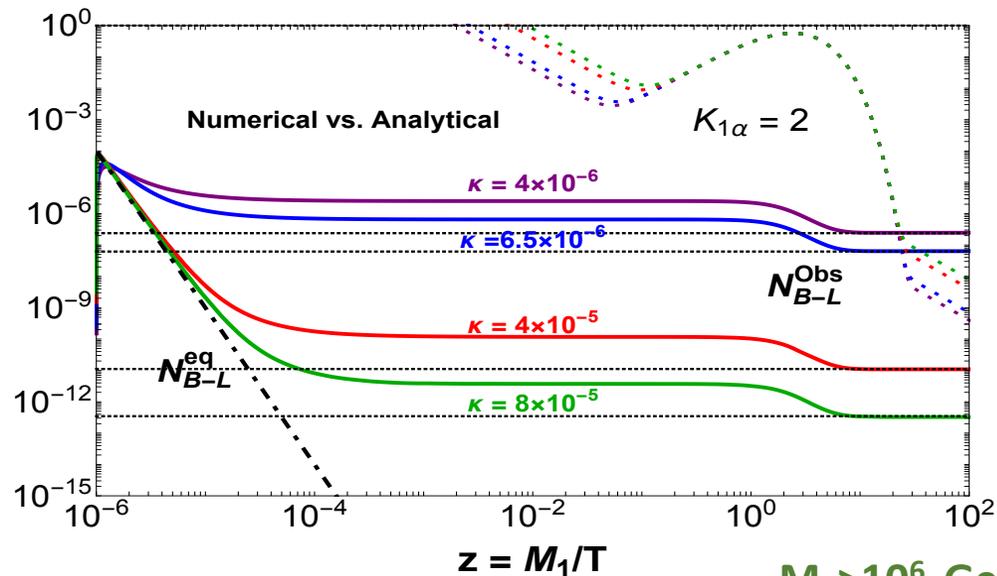
$W = 1, \Omega_{\text{GW}} (f > f_*) \sim f^{+1}$: Gravitational Leptogenesis $\Rightarrow \partial_\mu R_j^\mu / M^2$ generated in seesaw at 2-loops

$$N_{B-L}^{\text{eq}} = \frac{\pi^2 \dot{R}}{36(4\pi)^4} \sum_{j>i} \frac{\text{Im}[k_{ij}^2]}{\zeta(3) T M_i M_j} \left(\frac{M_j^2}{M_i^2}\right) \ln\left(\frac{M_j^2}{M_i^2}\right),$$

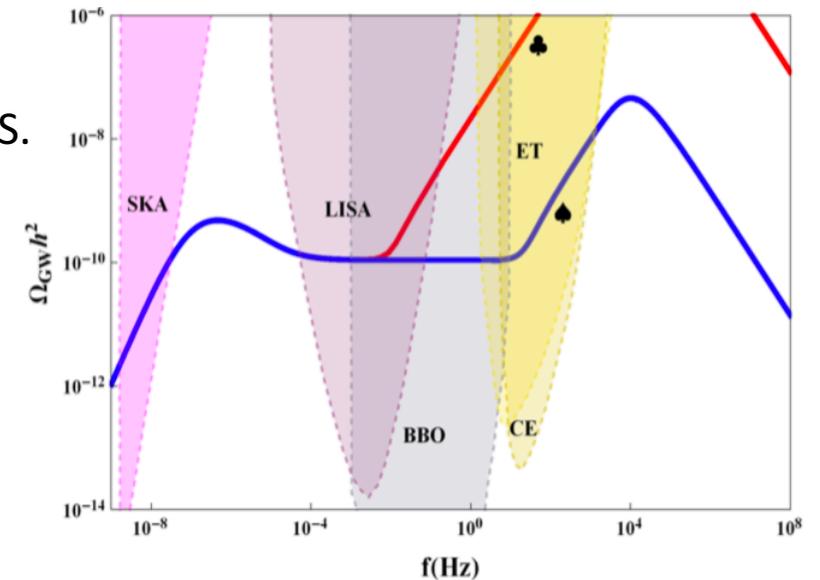
$$\frac{dN_L}{dz} = -2 W_{\Delta L=2} (N_L - N_L^{\text{eq}}),$$

$\sim 1-3w \approx 0$

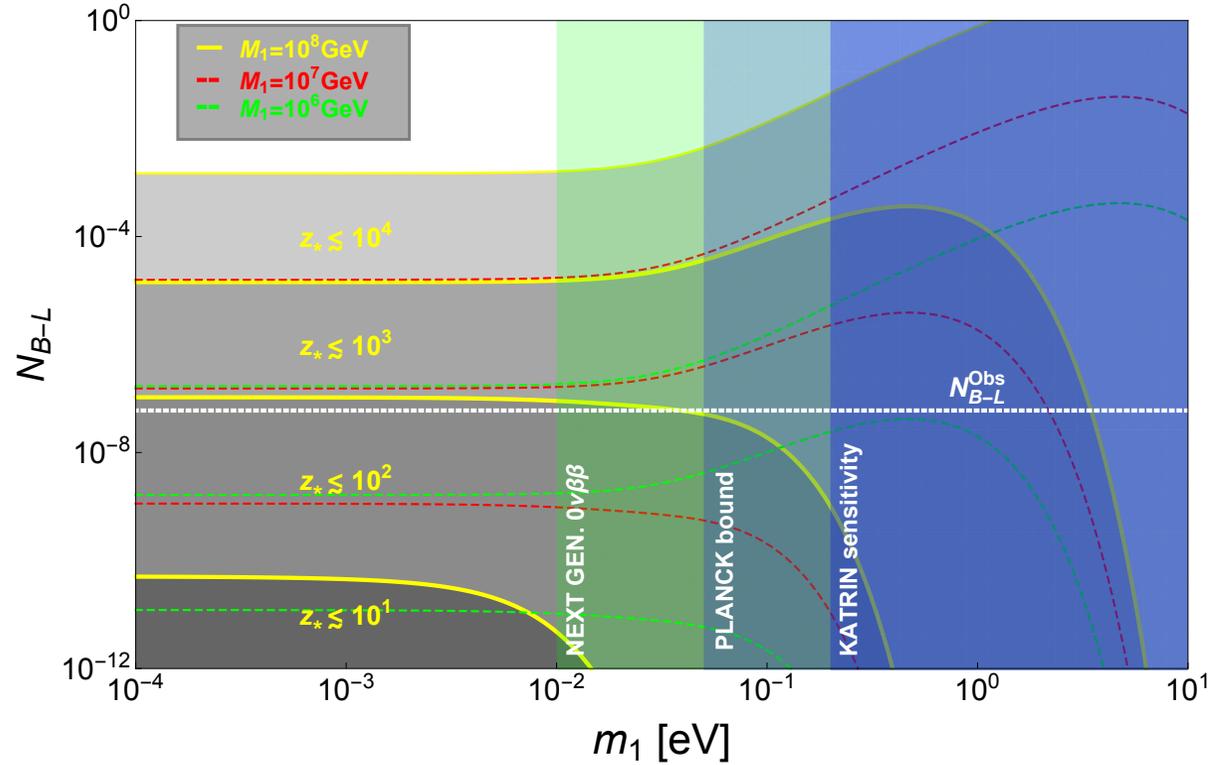
K. Saikawa and S. Shirai, JCAP 05, 035 (2018)



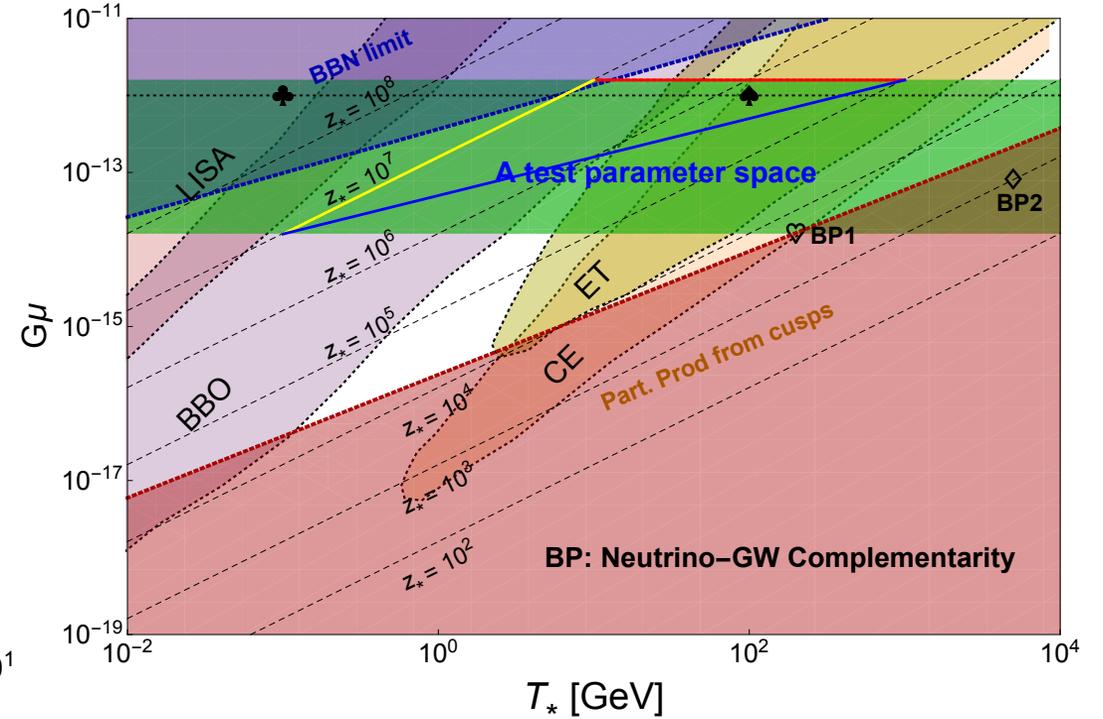
$M_1 > 10^6$ GeV even in 2RH neutrino seesaw model



Neutrino-Gravitational waves complementarity



Large Yukawa coupling => test point is would be side ET



Samanta et al 2108.08359 (2021)

$w = 0, \Omega_{GW} (f > f_*) \sim f^{-1}$ Baryogenesis from ultra-light primordial black holes (Hawking radiation)

Inflation \rightarrow U(1) breaking \rightarrow black hole formation \rightarrow Heavy neutrino emission (HR) \rightarrow baryogenesis

Mass

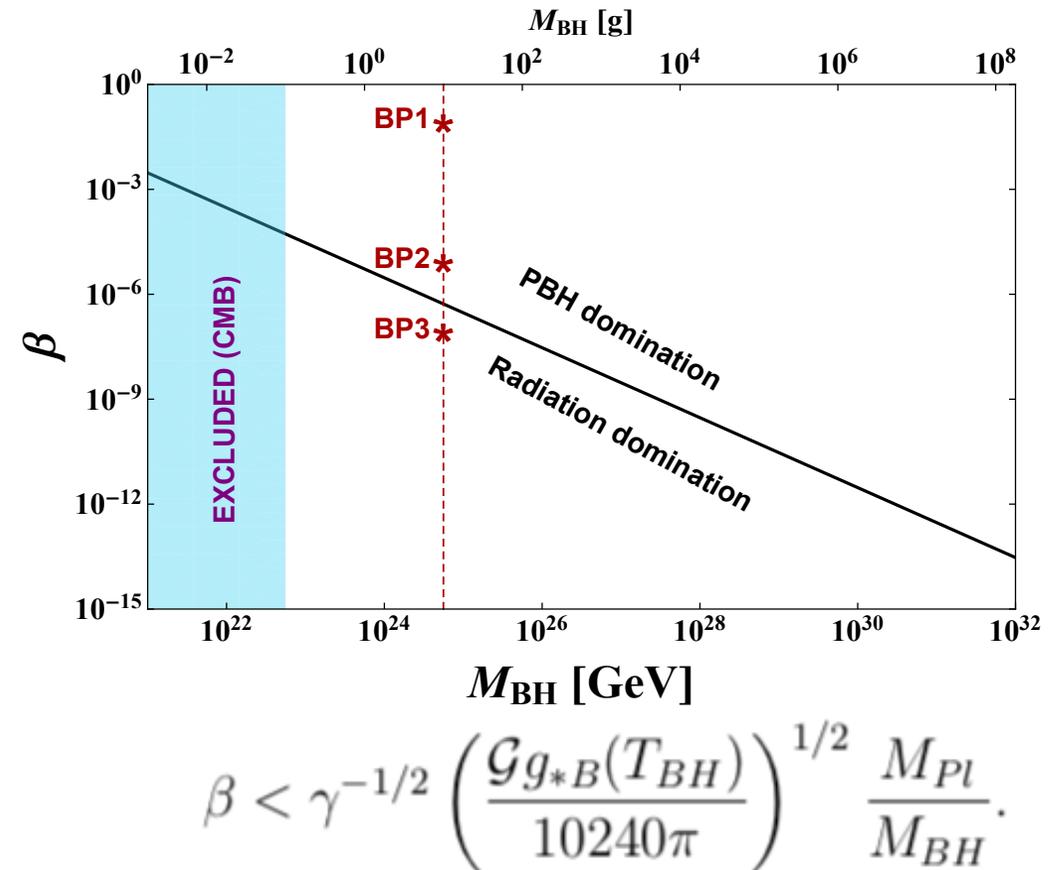
$$M_{BH} = \gamma \frac{4}{3} \pi (H_{Bf}^{-1})^3 \rho_{Bf} \quad \text{with} \quad \rho_{Bf} = \frac{3H_{Bf}^2 M_{Pl}^2}{8\pi}, \quad H_{Bf} = \frac{1}{2t_{Bf}}$$

Mass loss

$$\frac{dM_{BH}}{dt} = - \frac{\mathcal{G} g_{*B}(T_{BH})}{30720\pi} \frac{M_{Pl}^4}{M_{BH}^2}$$

Life-time

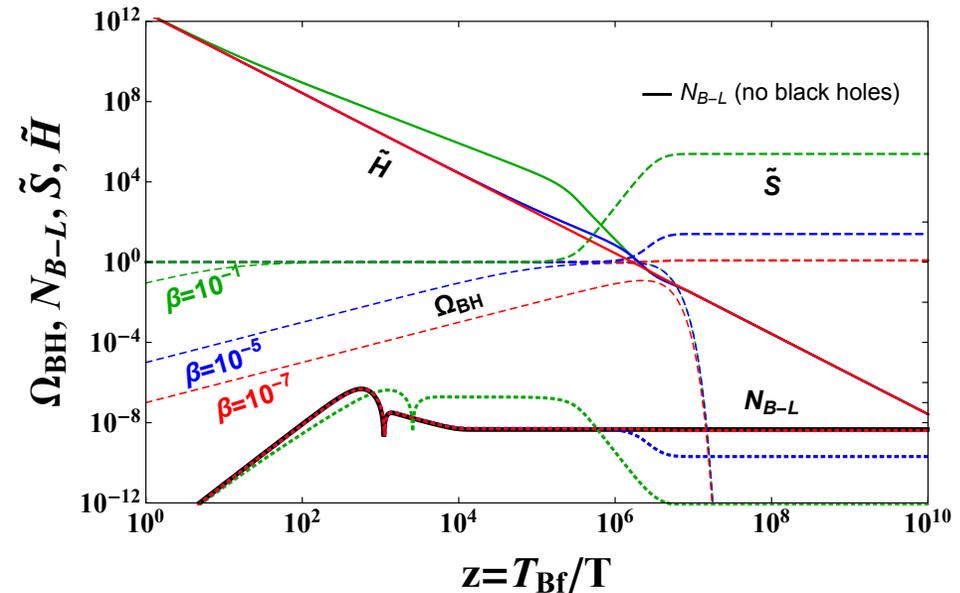
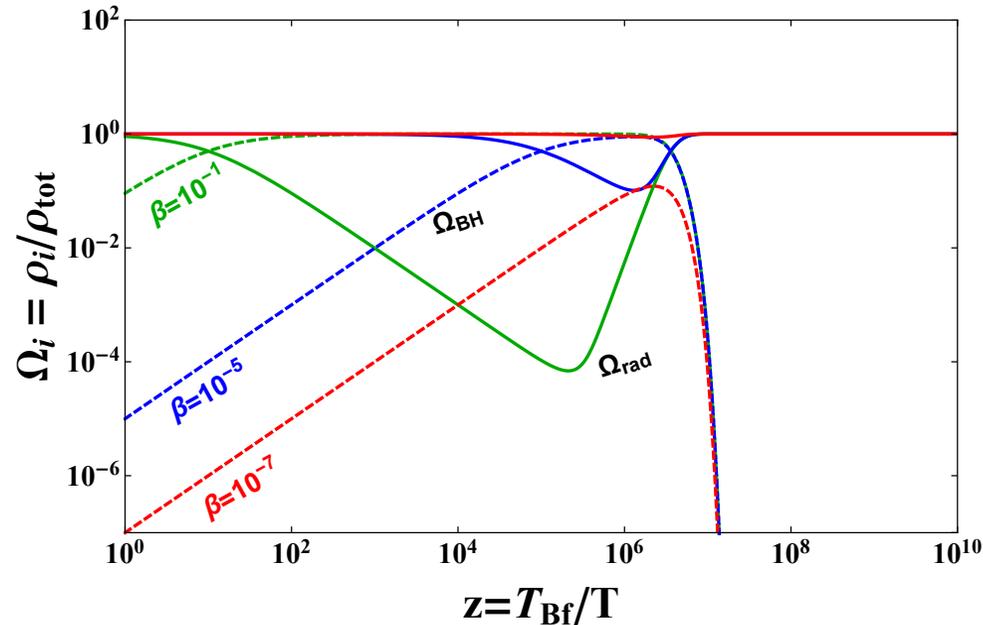
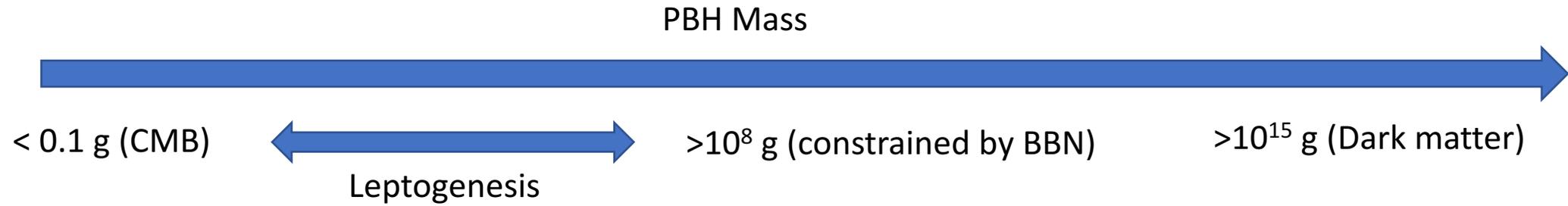
$$\tau = \int_{t_{Bf}}^{t_{ev}} dt = - \int_{M_{BH}}^0 dM_{BH} \frac{30720\pi M_{BH}^2}{\mathcal{G} g_{*B}(T_{BH}) M_{Pl}^4} = \frac{10240\pi M_{BH}^3}{\mathcal{G} g_{*B}(T_{BH}) M_{Pl}^4}$$



$w = 0, \Omega_{\text{GW}} (f > f_*) \sim f^{-1}$ Baryogenesis from ultra-light primordial black holes (Hawking radiation)

Samanta et al JCAP(2021)

Inflation => U(1) breaking => black hole formation => Heavy neutrino emission (HR) => baryogenesis

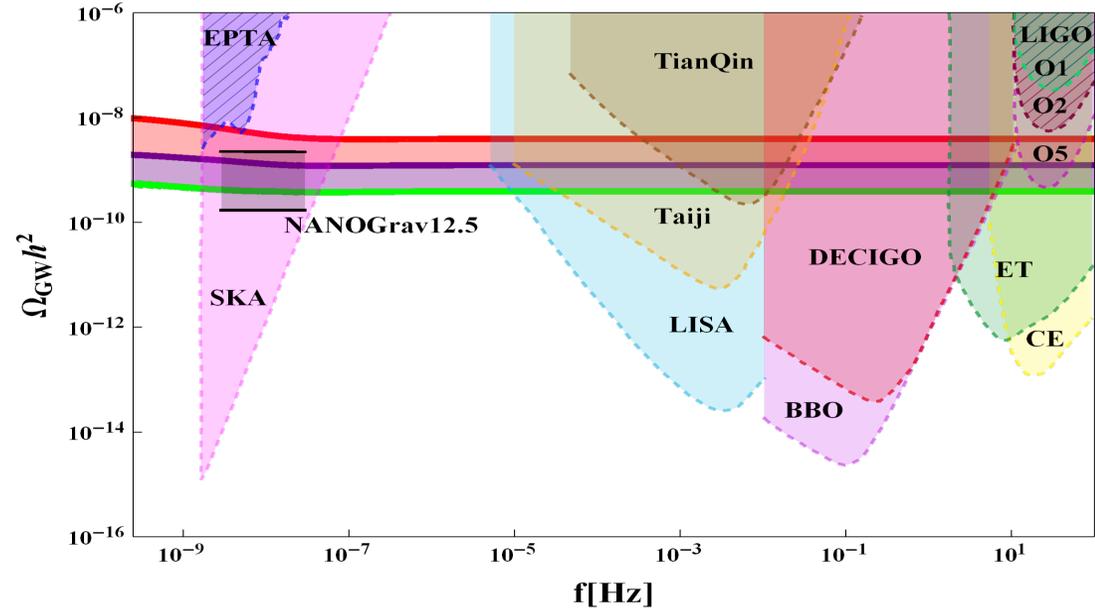
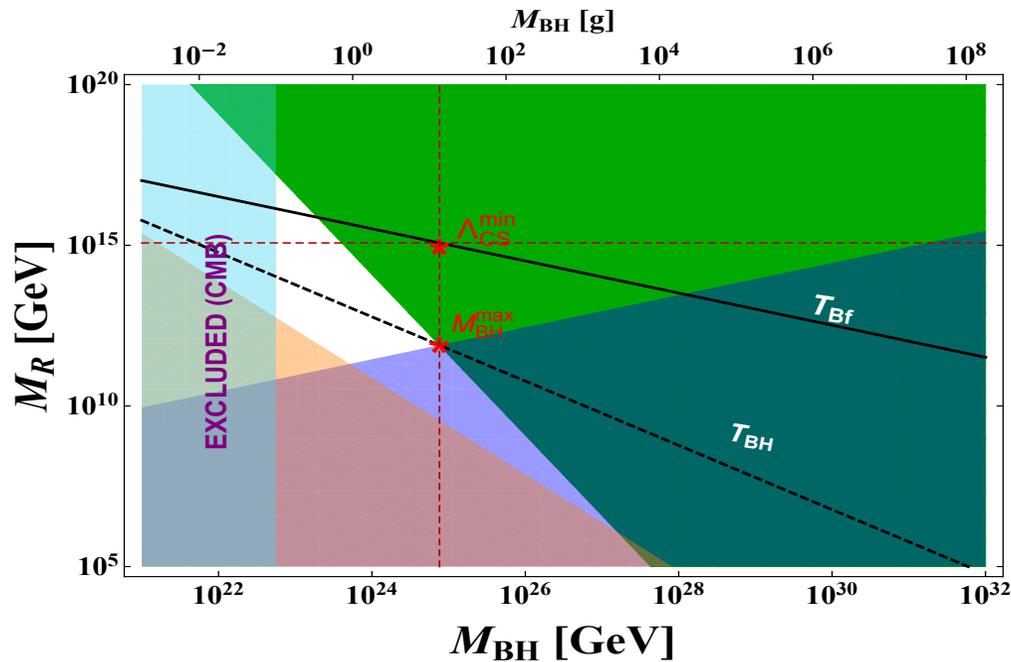


Cosmic archeology, GW spectral shapes and Leptogenesis

Number of radiated particles: $\frac{4\pi}{3} \frac{g_X}{g_{*B}} \left(\frac{M_{BH}}{M_{Pl}}\right)^2$ for $T_{BH} > M_X$, $\frac{1}{48\pi} \frac{g_X}{g_{*B}} \left(\frac{M_{Pl}}{M_X}\right)^2$ for $T_{BH} < M_X$.

Only Ultra light PBHs ($M_{BH} < 13$ g) can produce correct baryon asymmetry

Tension with NANOGrav data, with large loops



Future improvement

We did not consider black hole-string network that could provide spectral distortion.

(**Samanta** et al)

We trying to accommodate **Black** hole generated **Right Hand Neutrino (Black-RHiNOs)** super heavy dark matter in this framework that can be probed with a cosmic ray gravitational waves complementarity.

(**Samanta**, F. Urban et al)

Summary

Gravitational Waves could be potential probe of High scale leptogenesis which otherwise cannot be probed e.g., in colliders.

Towards a clean probe, one may relate the spectral shapes of the GWs to leptogenesis.

I discussed to concrete models for two types of leptogenesis scenarios

Gravitational Leptogenesis (rising GW spectrum) Black hole leptogenesis (falling GW spectrum)