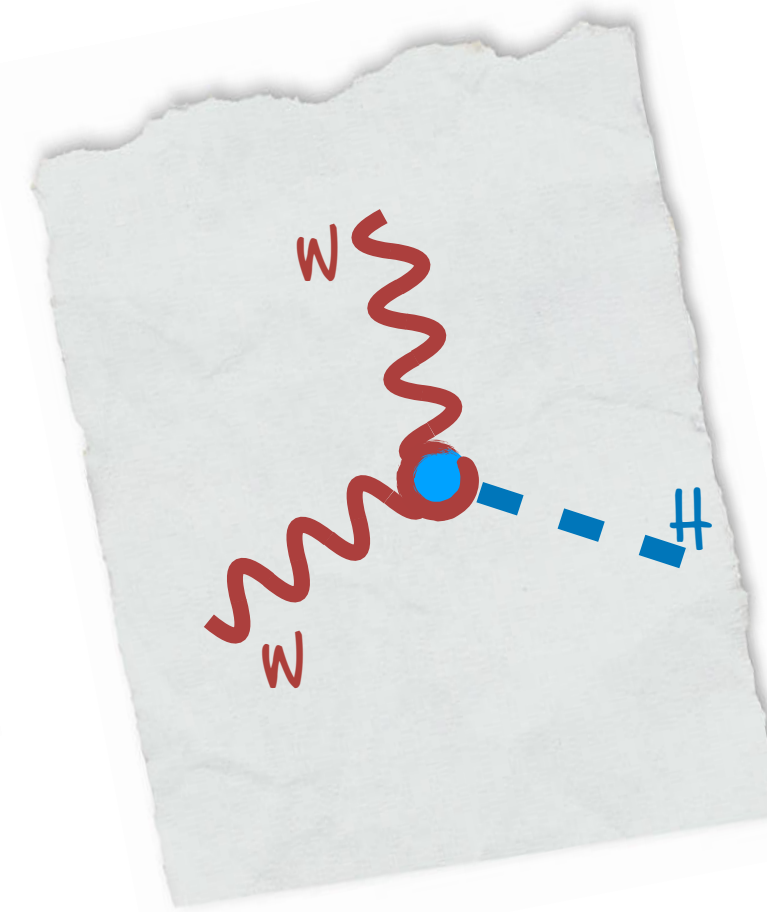
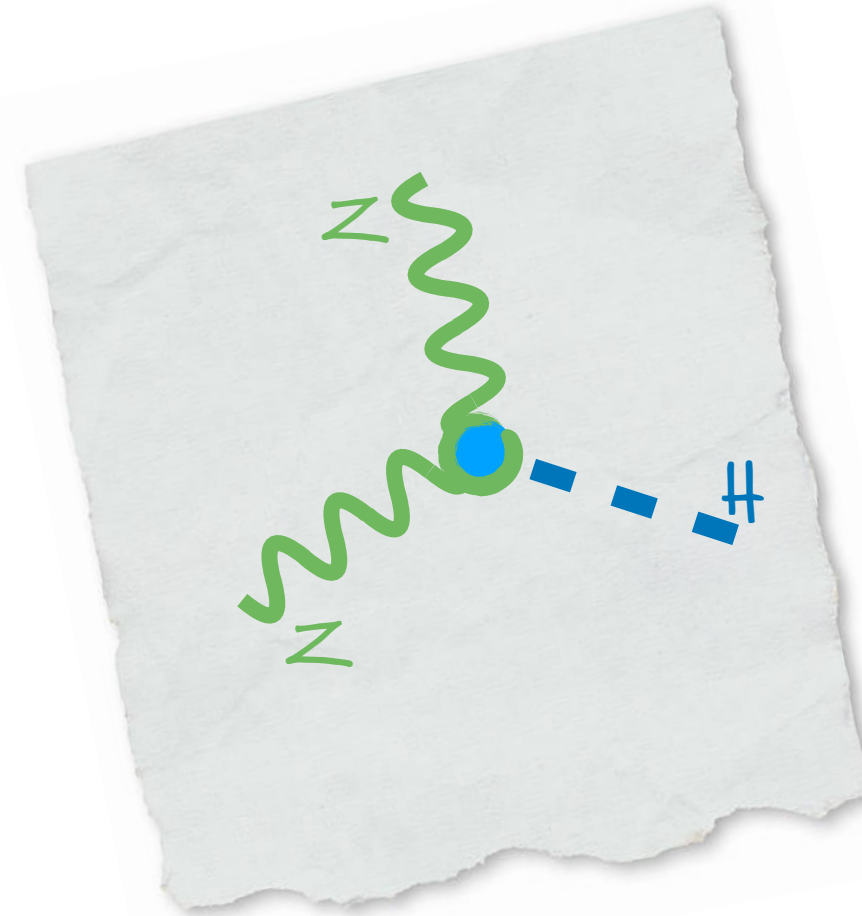
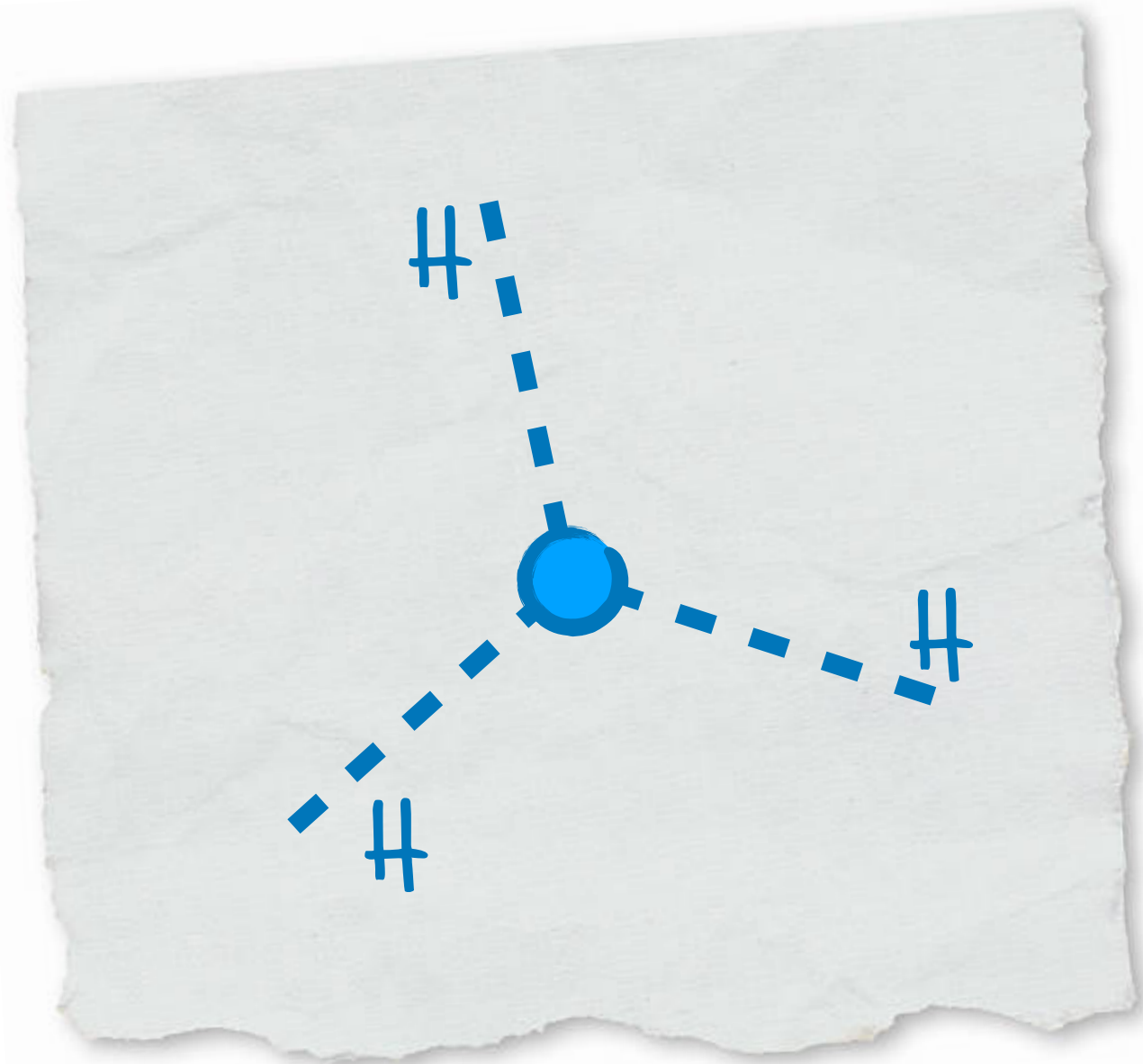


Higgs boson self-interactions

... and Higgs-gauge-boson interactions



Roberto Salerno
Cesare Cazzaniga

Roy Lemmon



Science & Technology Facilities Council
Daresbury Laboratory



Science and
Technology
Facilities Council

Foreword

These studies have been done in a contest of a 4 months M2 internship. The goal was to put in place a full analysis chain from the generation of the samples through the selection of candidates to the statistical analysis.

There are many caveats (results are on the optimistic side)

- The centrally produced samples are (yet) not used
- Not all the systematics uncertainties are included
- Only main backgrounds are considered so less selection cuts included leading to higher signal efficiency

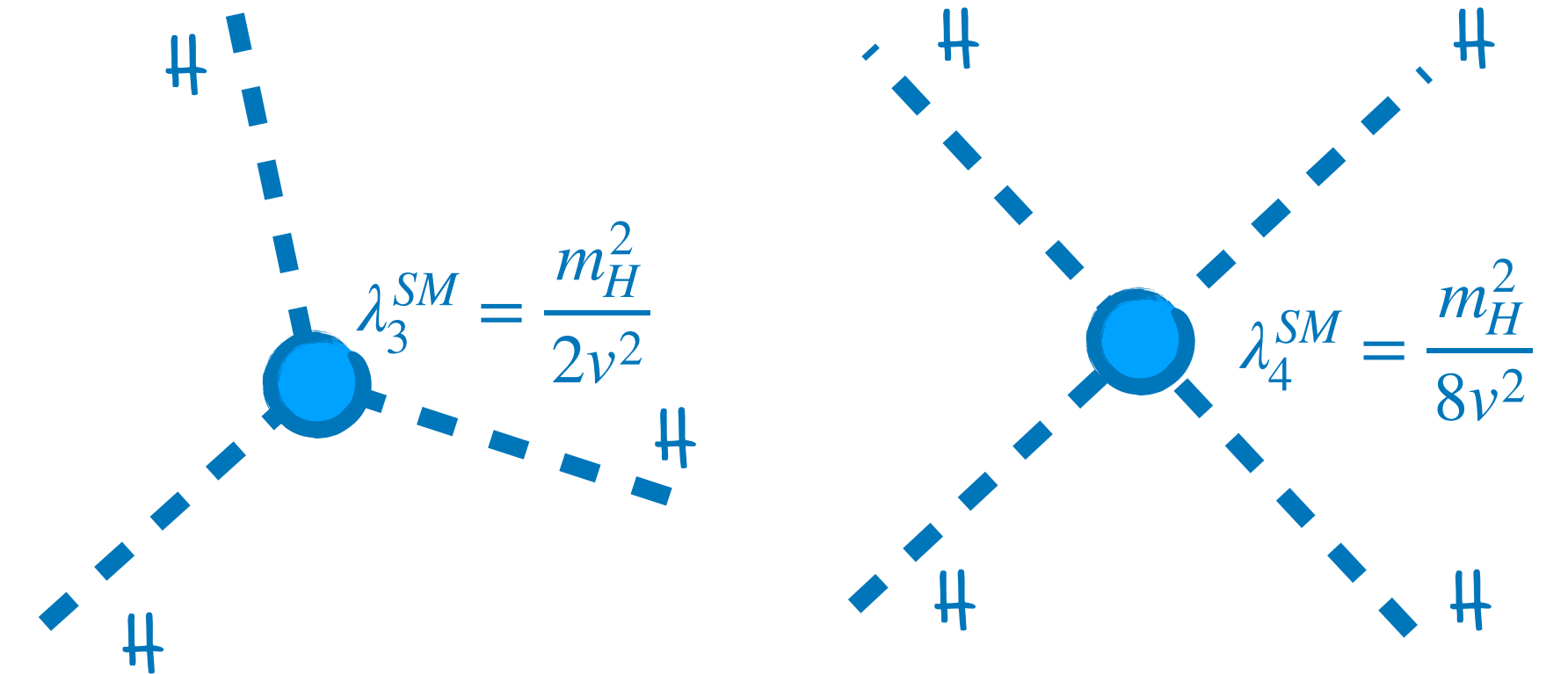


Does the Higgs boson interact with itself?

A self-interacting Higgs (as SM predicts) would be unlike anything yet seen in nature. All other interactions change particle identity.

The Higgs boson cubic (λ_3^{SM}) and quartic (λ_4^{SM}) couplings are the keys to check the EWSB. The Higgs boson potential is :

$$\mathcal{L} \subset -\frac{m_h^2}{2}h^2 - \lambda_3^{SM}vh^3 - \lambda_4^{SM}h^4$$

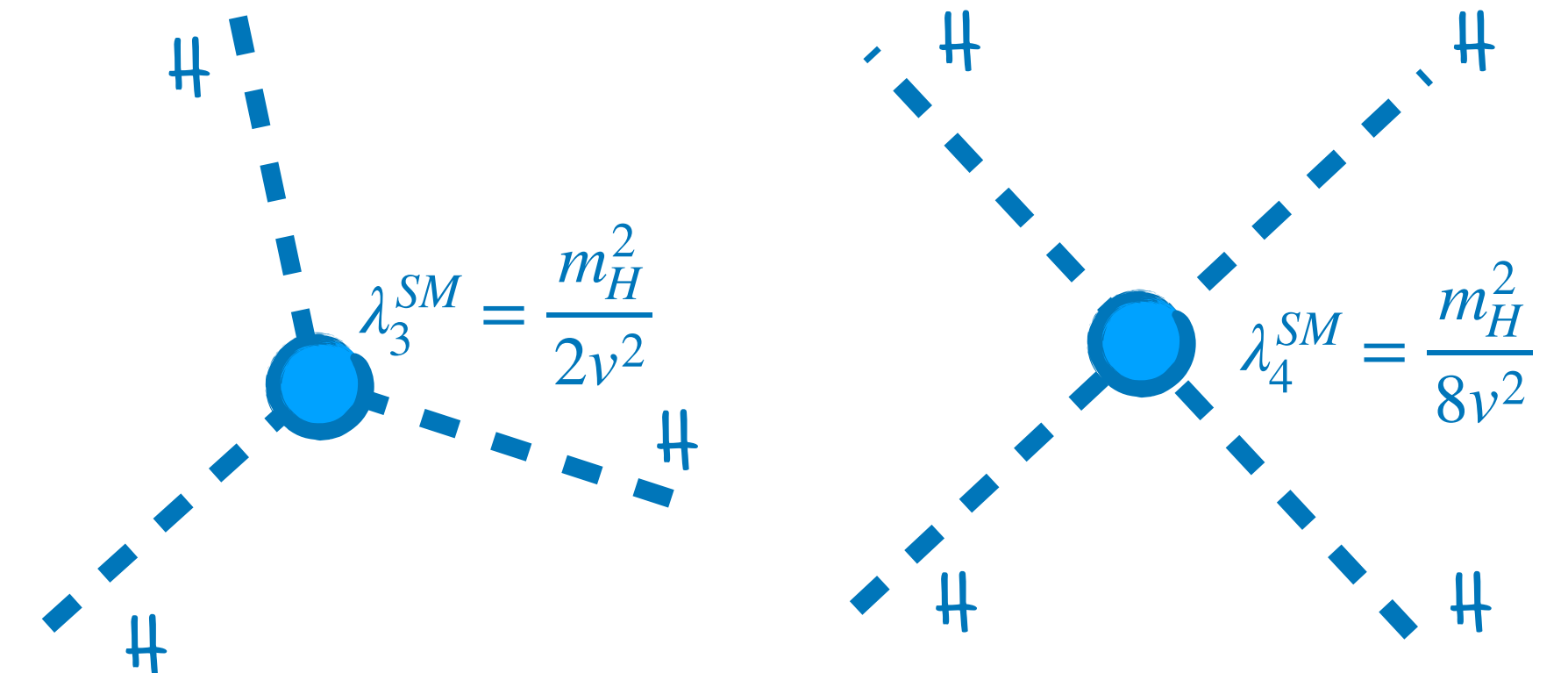


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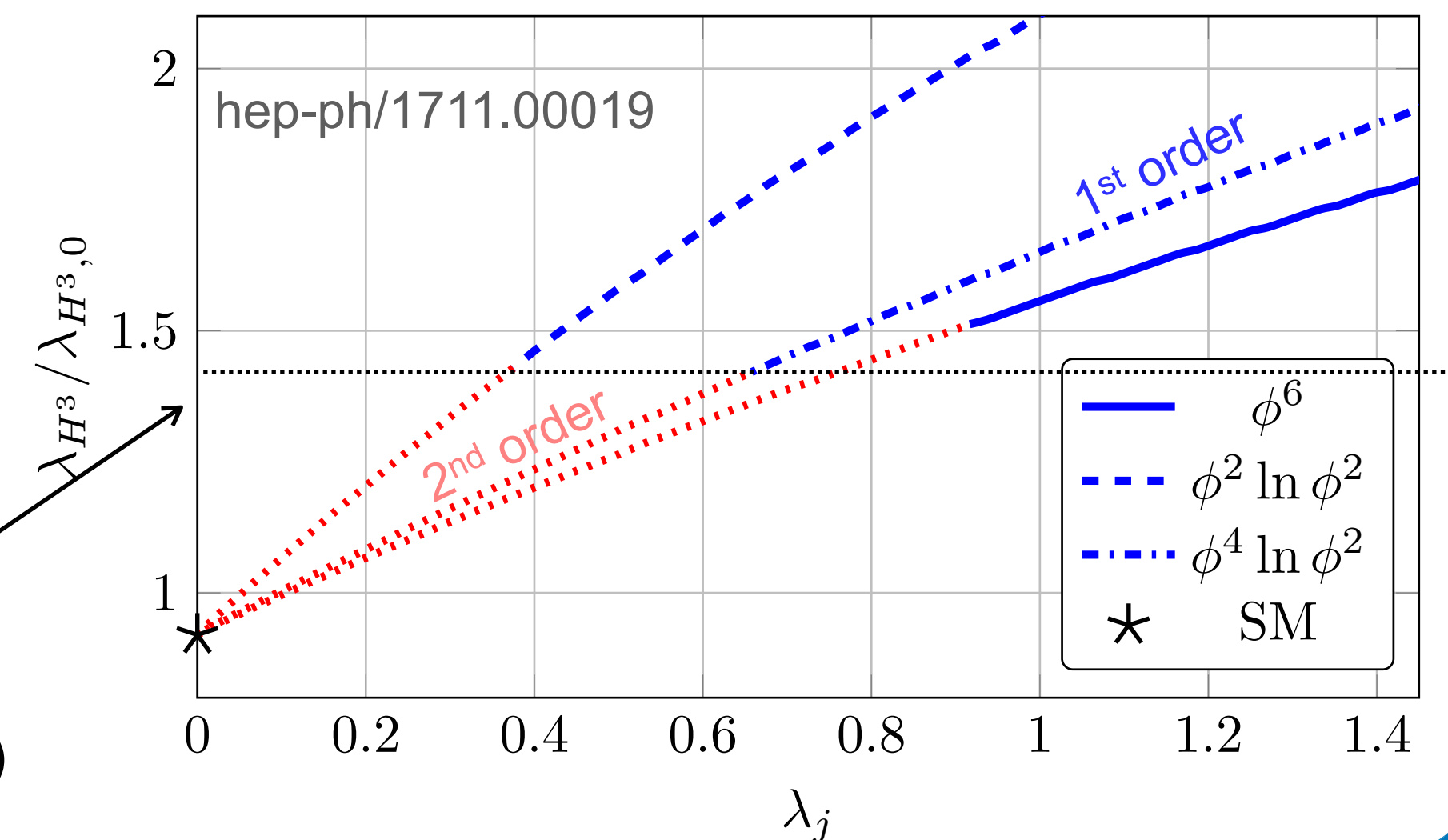
$$\mathcal{L} \subset -\frac{m_h^2}{2}h^2 - \lambda_3^{SM}vh^3 - \lambda_4^{SM}h^4$$



Link with the cosmology

Deviations from SM Higgs boson self-coupling cause a modified potential that allows **first-order electroweak phase transition** and hence an explanation of the observed matter vs anti-matter asymmetry!

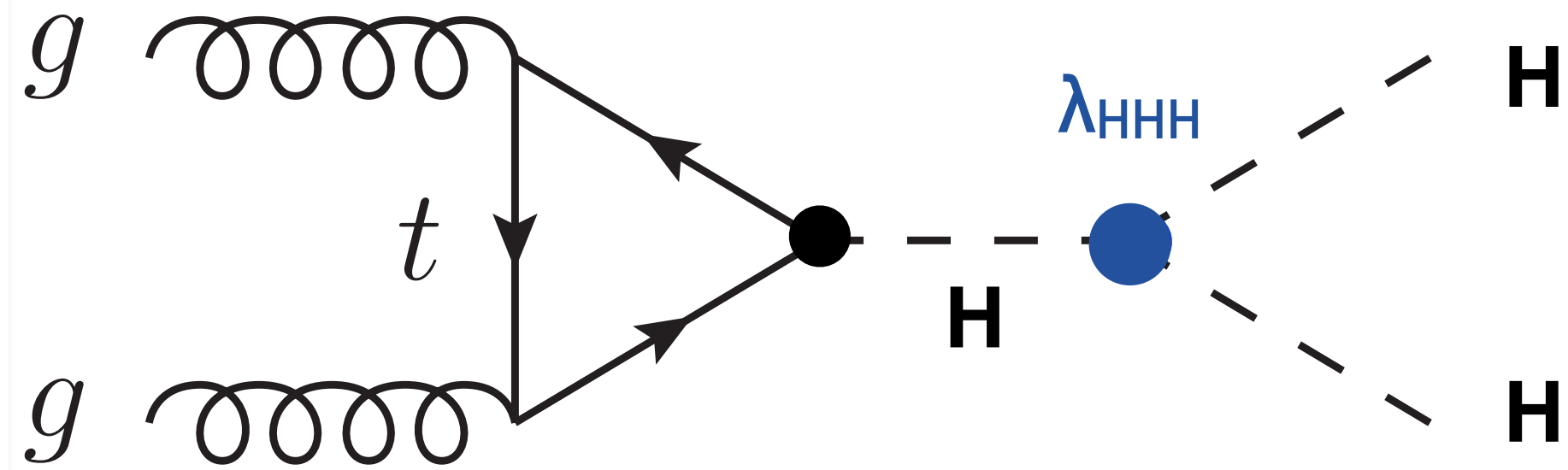
We need to probe size of modification down to 1.4, the expected uncertainty of the measurement should be $\mathcal{O}(10\%)$



The Higgs boson self-coupling before FCC-ee

Direct measurements @ HL-LHC

Profiting of on-shell production of two Higgs bosons

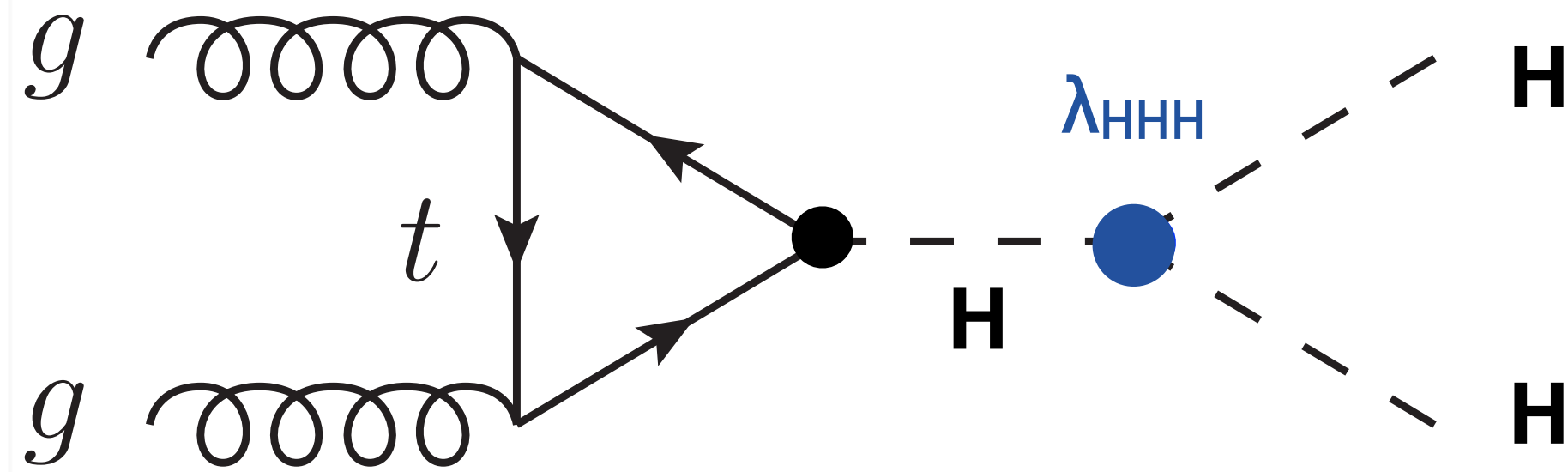





- 👍 Theoretically clean
- 👎 Very rare process ($\sigma \sim 40 \text{ fb}$ @ 14 TeV)
- 👎 Experimentally challenging

The Higgs boson self-coupling before FCC-ee

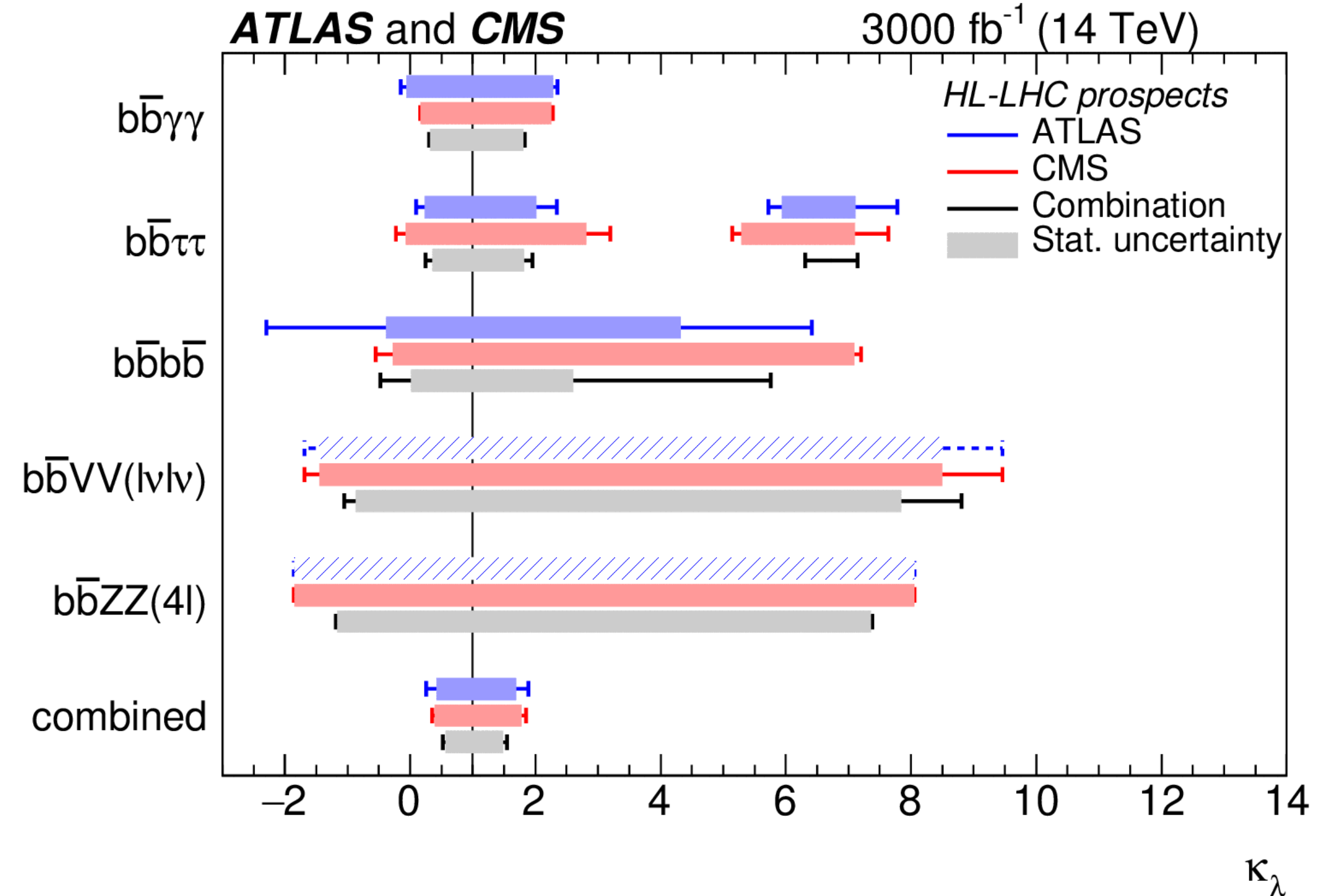
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Profiting of on-shell production of two Higgs bosons



-  Theoretically clean
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Expected results

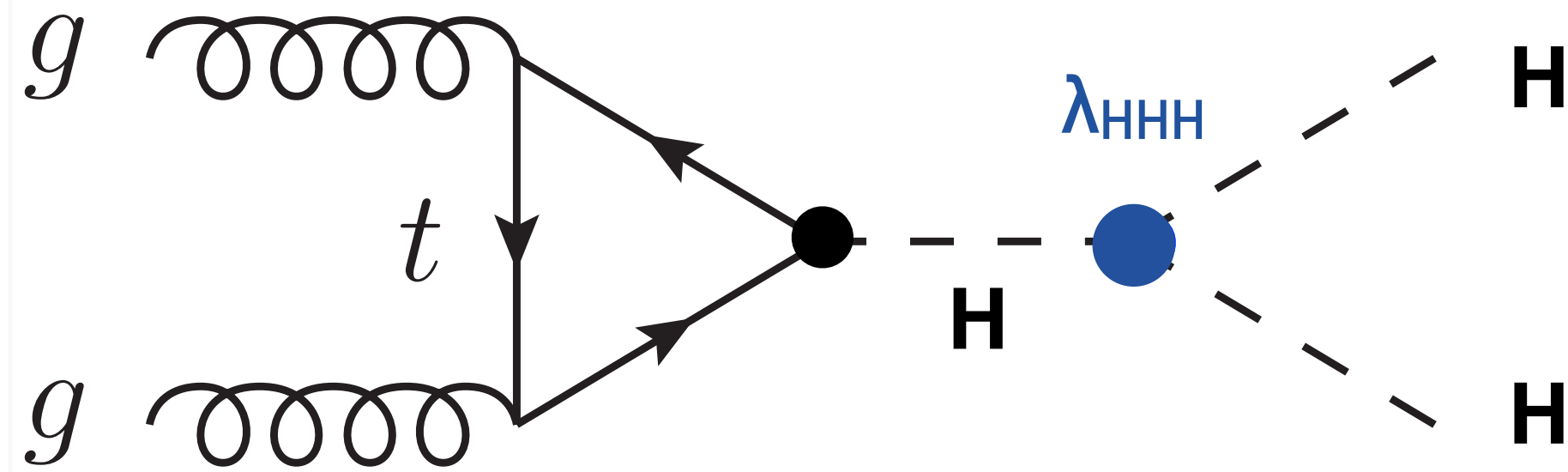





arXiv:1902.00134

The Higgs boson self-coupling before FCC-ee

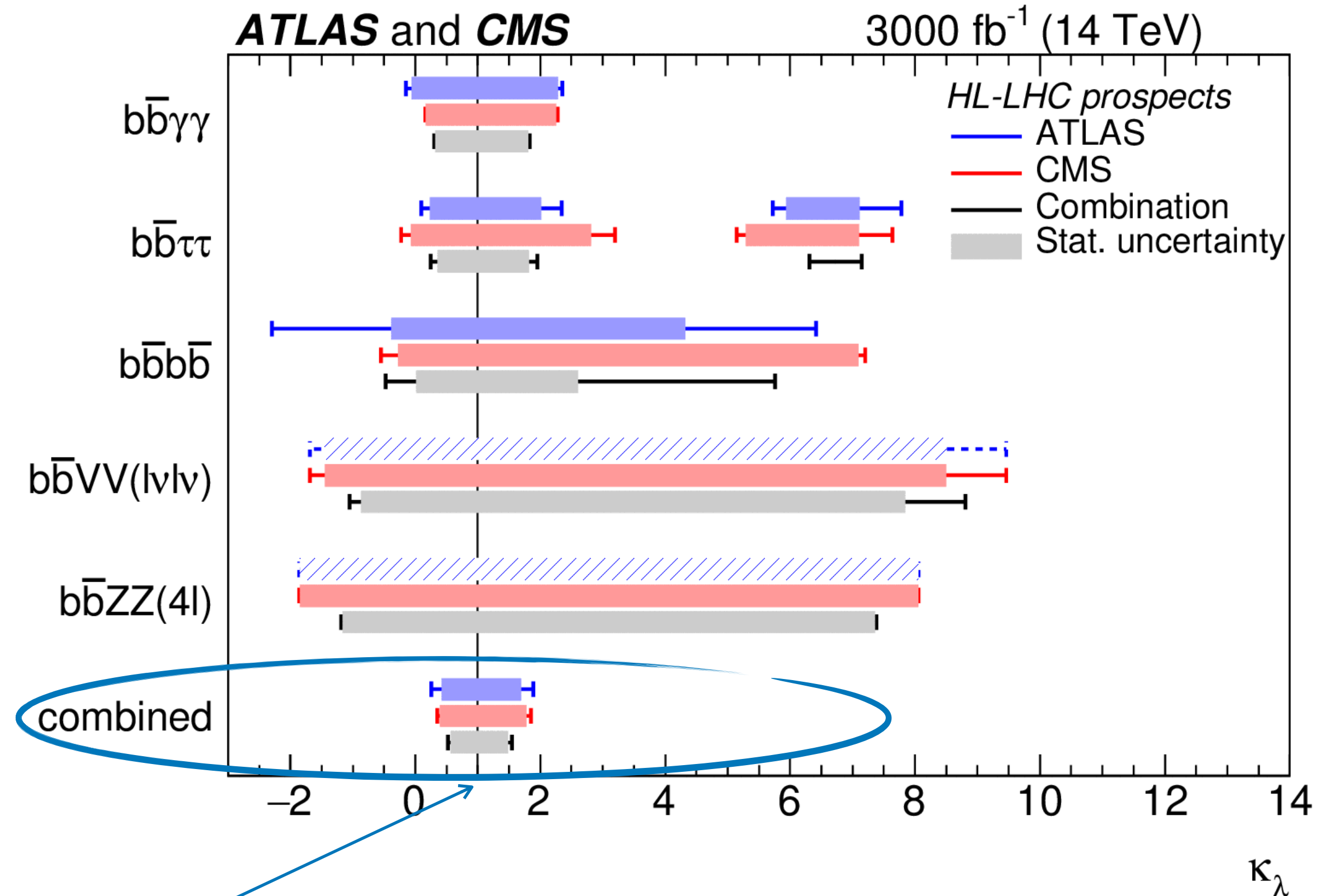
Direct measurements @ HL-LHC

Profiting of on-shell production of two Higgs bosons



-  Theoretically clean
-  Very rare process ($\sigma \sim 40 \text{ fb}$ @ 14 TeV)
-  Experimentally challenging

Expected results



arXiv:1902.00134

$\mathcal{O}(50\%)$ precision
2 experiments and various channels

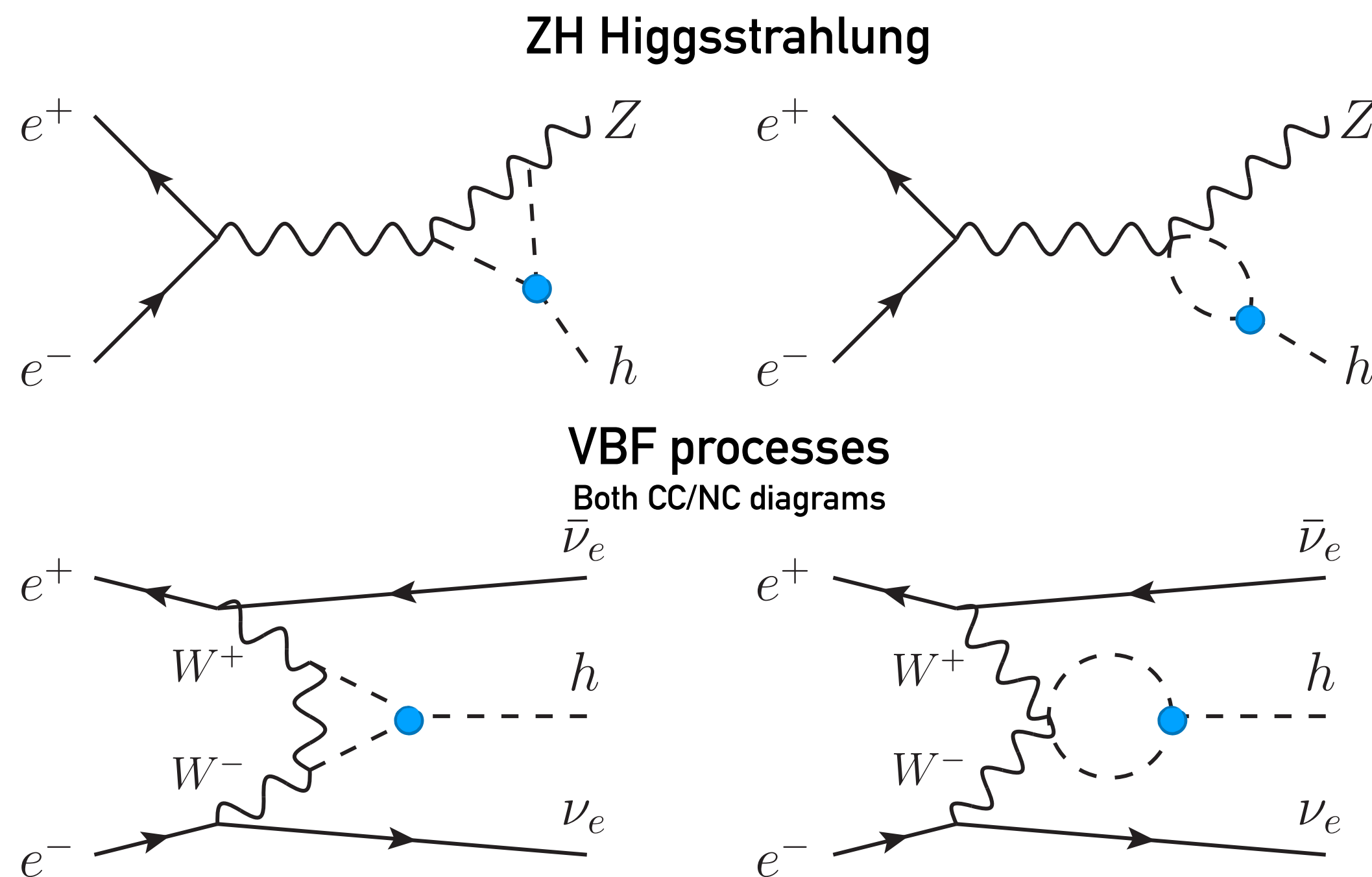
Higher-order corrections to single-Higgs processes

λ_{HHH} does not enter single-Higgs processes at LO but it affects both Higgs boson production and decay at NLO.

Higher-order corrections to single-Higgs processes

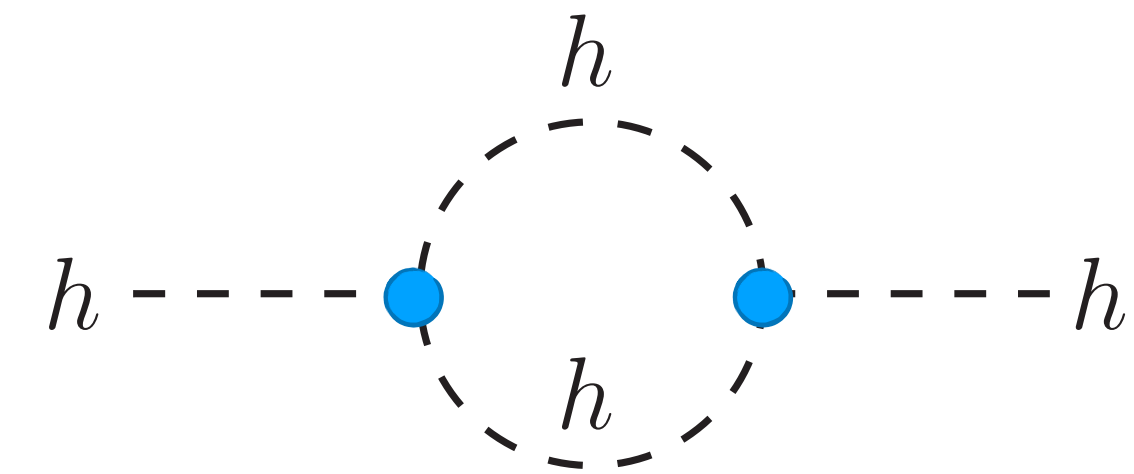
λ_{HHH} does not enter single-Higgs processes at LO but it affects both Higgs boson production and decay at NLO.

Vertex corrections (linear in k_λ)



Higgs boson self-energy (quadratic in k_λ)


Universal modifications via wave function renormalisation



λ_3 effect

The NLO corrections to an observable Σ

$$\Sigma_{\text{NLO}} = \boxed{Z_H} \Sigma_{\text{LO}} (1 + \kappa_\lambda \boxed{C_1})$$



Universal coefficient
from wave function

Process and kinematic
dependent coefficient

λ_3 effect

The NLO corrections to an observable Σ

$$\Sigma_{\text{NLO}} = \boxed{Z_H} \Sigma_{\text{LO}} (1 + \kappa_\lambda \boxed{C_1})$$

Universal coefficient
from wave function

Process and kinematic
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Decay modes

C_1^Γ [%]	$\gamma\gamma$	ZZ	WW	$f\bar{f}$	gg
on-shell H	0.49	0.83	0.73	0	0.66

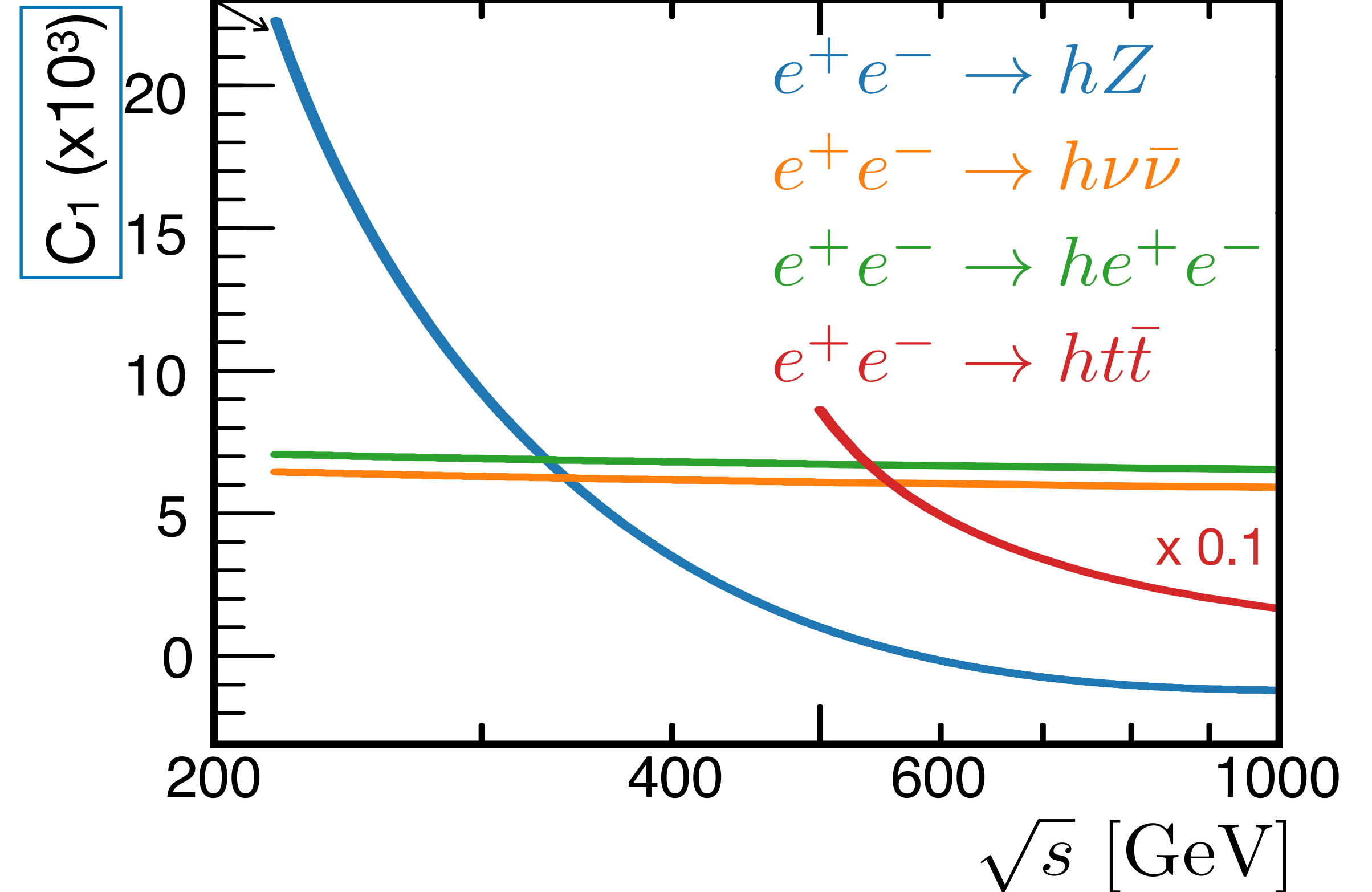
$$\boxed{C_1^{\Gamma_{\text{tot}}}} \equiv \sum_j \text{BR}^{\text{SM}}(j) C_1^\Gamma(j) \sim 2.3 \times 10^{-3}$$

Close to SM predictions the BRs
are sensitive to modified λ_{HHH}

$$\delta \text{BR}_{\lambda_3}(i) = \frac{(\kappa_\lambda - 1)(C_1^\Gamma(i) - C_1^{\Gamma_{\text{tot}}})}{1 + (\kappa_\lambda - 1)C_1^{\Gamma_{\text{tot}}}}$$

Production mechanisms

0.22 at threshold



Negligible impact of a modified λ_{HHH} on the angular asymmetries.
At 1-loop C_1 for HZ, WW-boson fusion and ZZ-boson fusion are independent of the beam polarization.

λ_3 effect

The NLO corrections to an observable Σ

$$\Sigma_{\text{NLO}} = \boxed{Z_H} \Sigma_{\text{LO}} (1 + \kappa_\lambda \boxed{C_1})$$

\downarrow Universal coefficient from wave function \downarrow Process and kinematic dependent coefficient

Decay modes

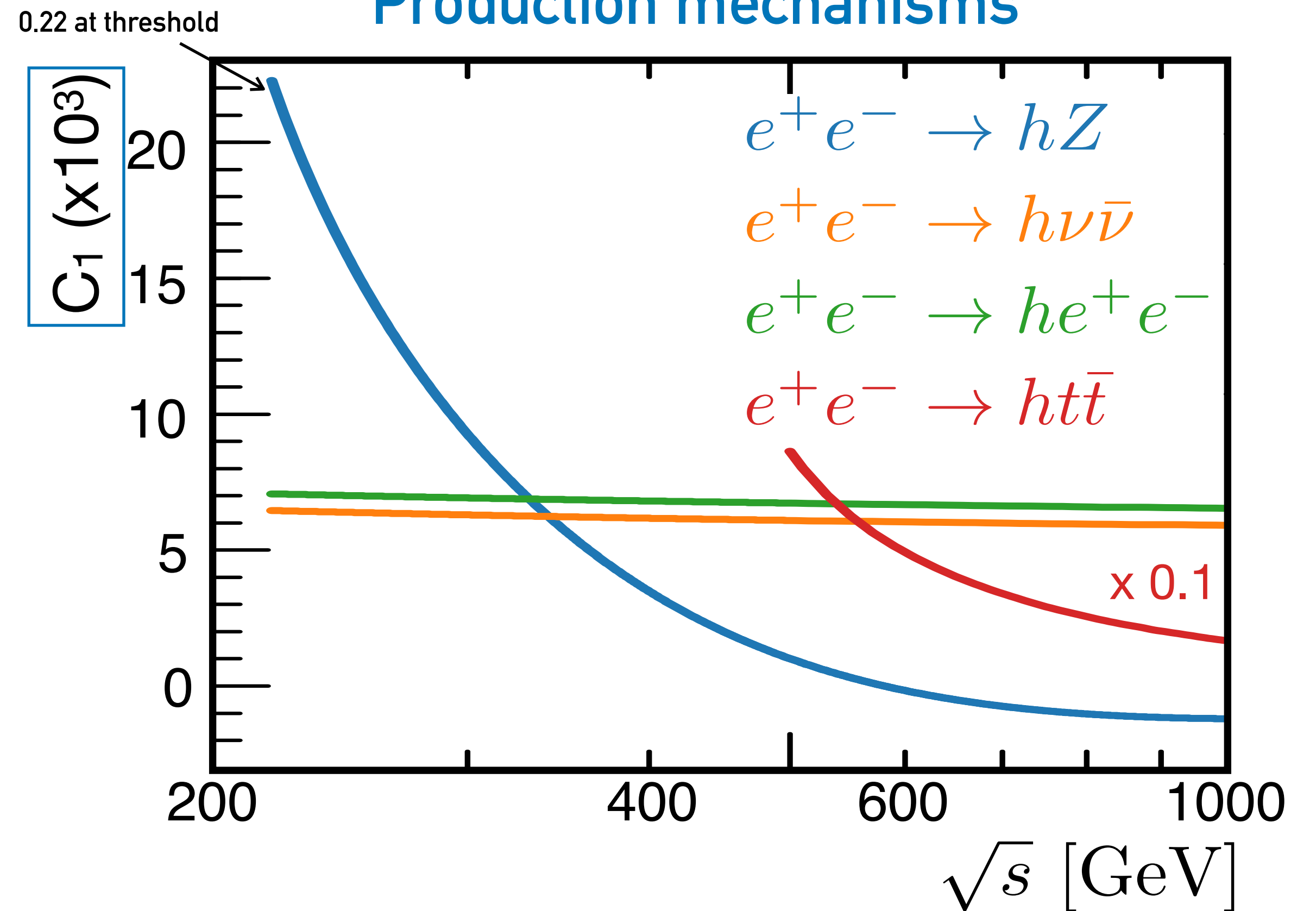
C_1^Γ [%]	$\gamma\gamma$	ZZ	WW	$f\bar{f}$	gg
on-shell H	0.49	0.83	0.73	0	0.66

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Close to SM predictions the BRs are sensitive to modified λ_{HHH}

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Production mechanisms



Negligible impact of a modified λ_{HHH} on the angular asymmetries. At 1-loop C_1 for HZ, WW-boson fusion and ZZ-boson fusion are independent of the beam polarization.

Great opportunity from the combination of inclusive or exclusive analyses and various production modes at two collision energies

Inclusive analyses @240 GeV

Exploited various Z decays, using the recoil techniques

Z($\mu\mu$)H

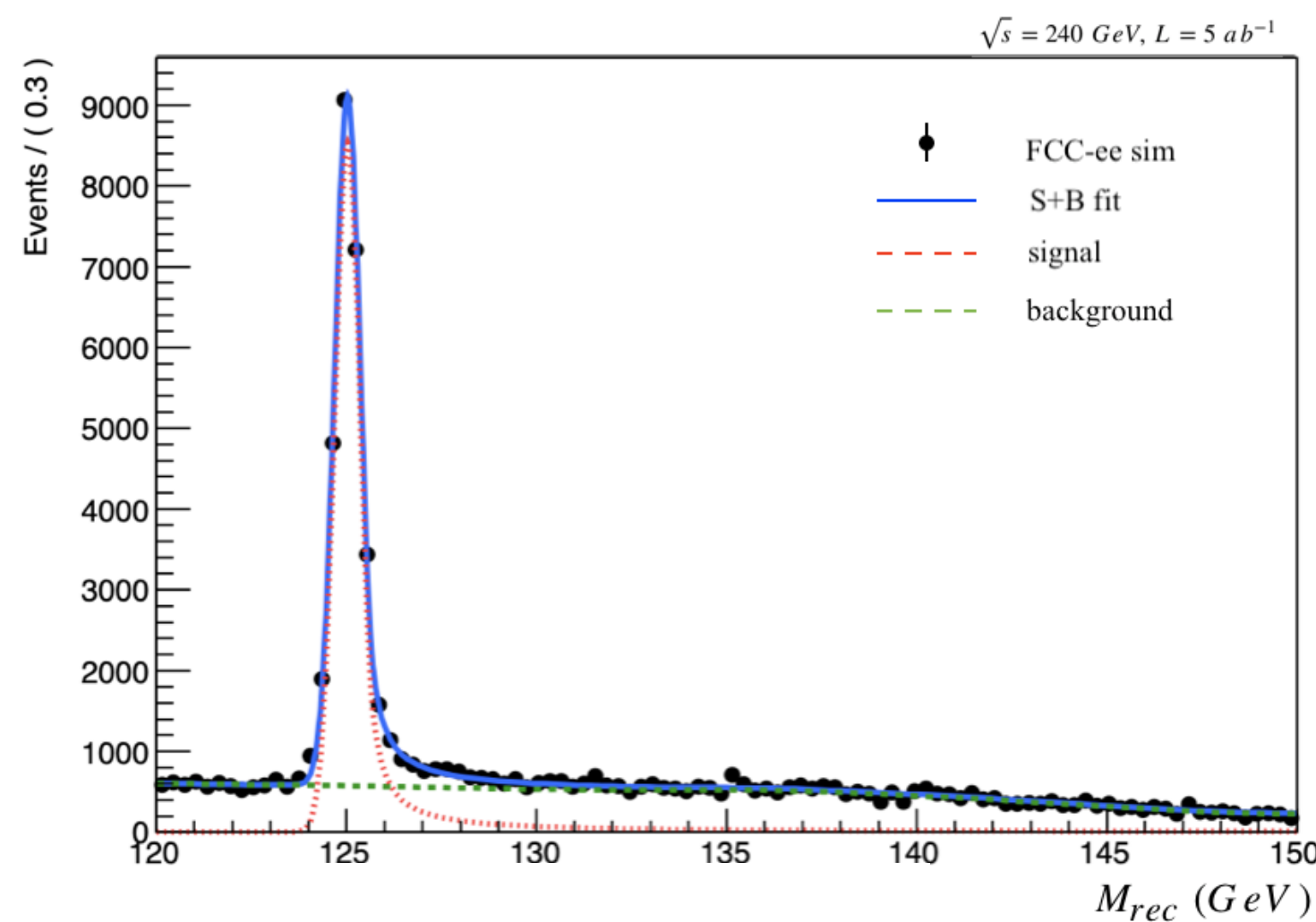
- ▶ $\mu^+\mu^-$ with $p_{T\mu 1} > 20$ GeV, $p_{T\mu 2} > 5$ GeV
- ▶ Minimum $|M_{\mu^+\mu^-} - M_Z|$
- ▶ $80 < M_{\mu^+\mu^-} < 100$ GeV

Z(ee)H

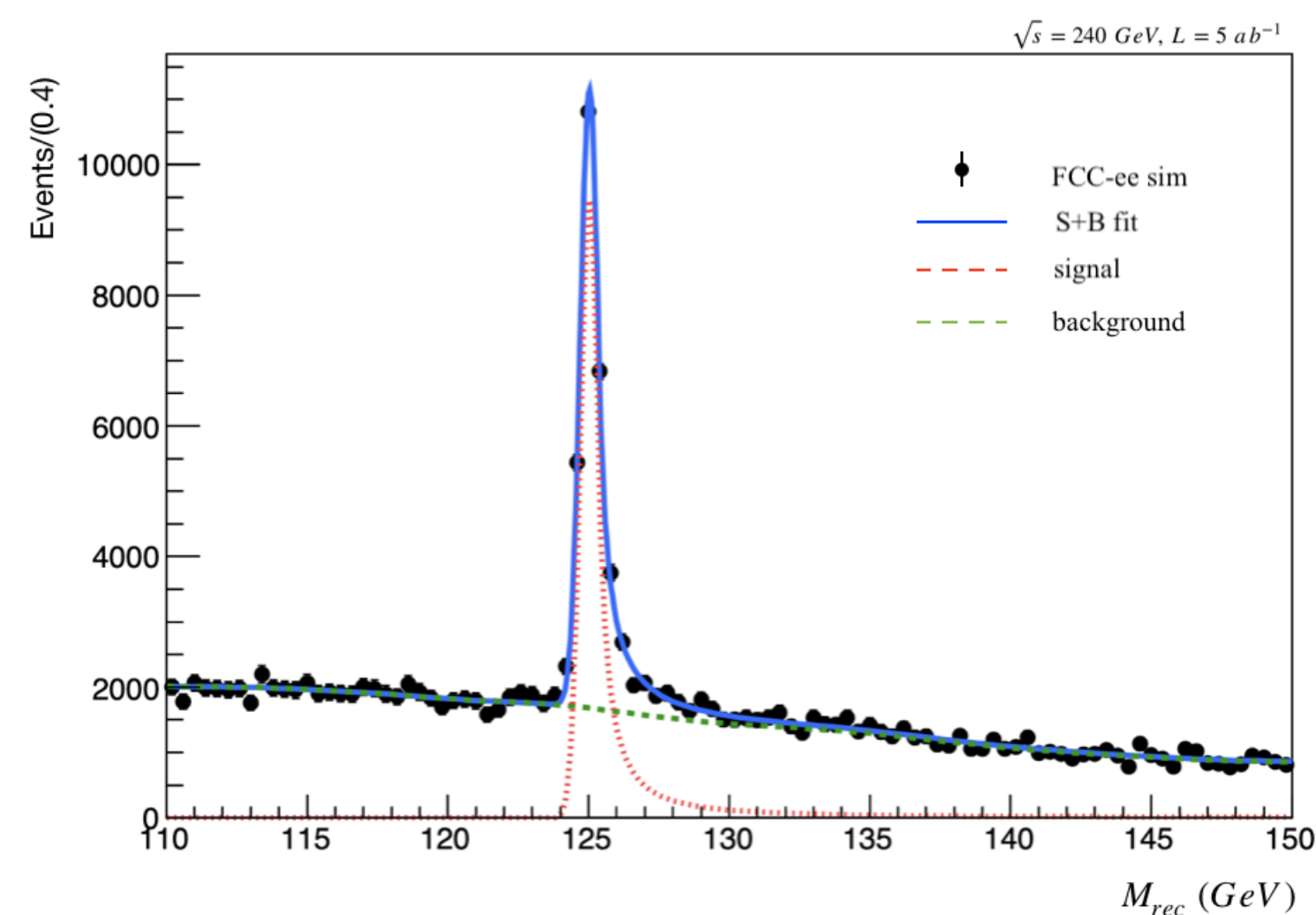
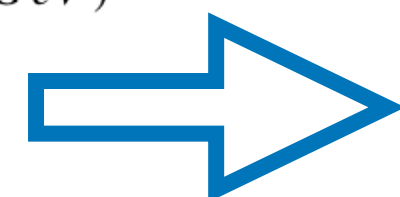
- ▶ e^+e^- with $p_{Te1} > 10$ GeV, $p_{Te2} > 5$ GeV
- ▶ Minimum $|M_{e^+e^-} - M_Z|$
- ▶ $60 < M_{e^+e^-} < 120$ GeV

Z(bb)H

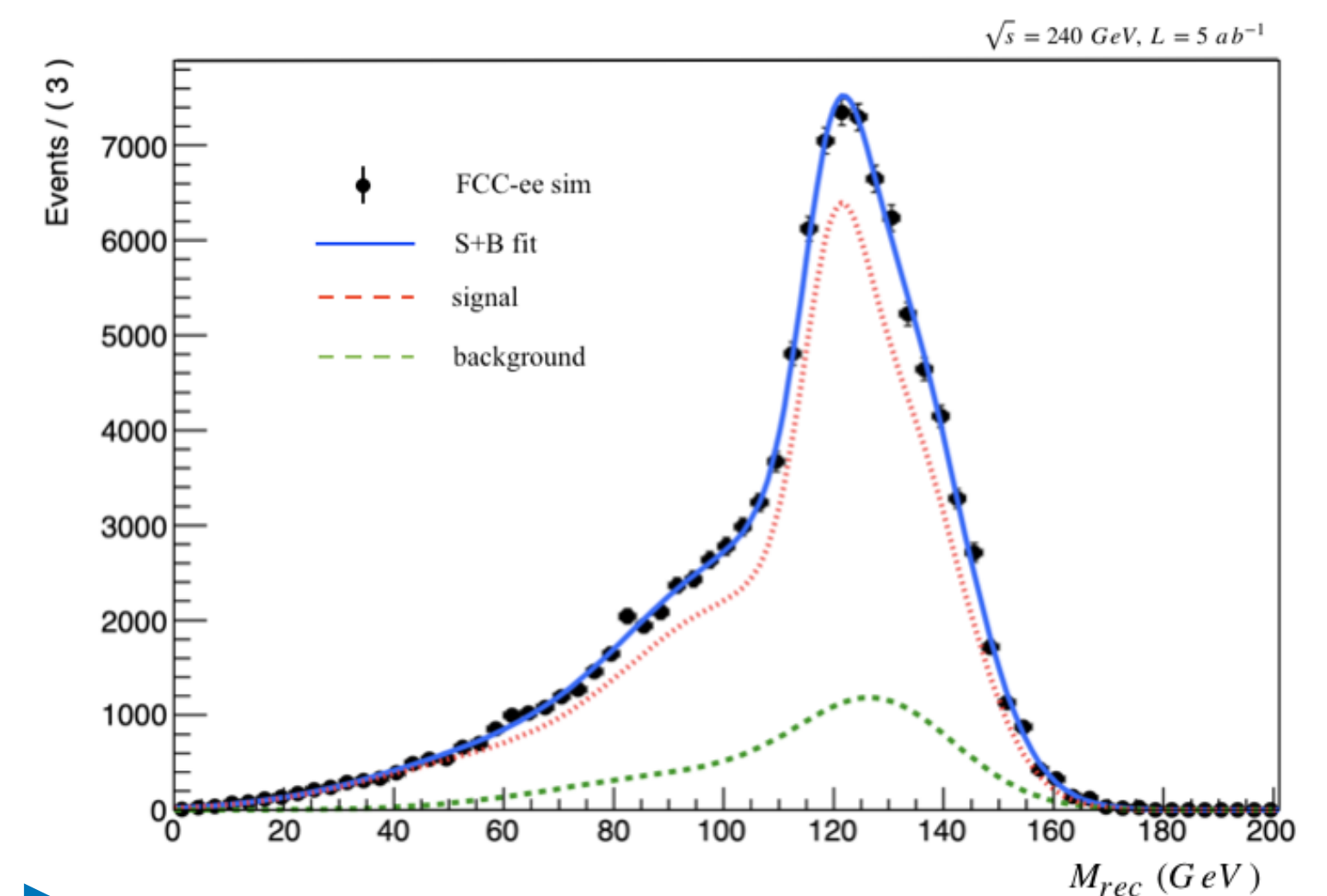
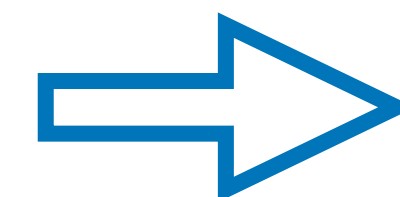
- ▶ ≥ 2 b-jets + $p_{Tjj} > 60$ GeV
- ▶ $M_{jj} > 45$ GeV
- ▶ $H_T > 10$ GeV
- ▶ BDT (17 variables)



Golden channel



Larger background



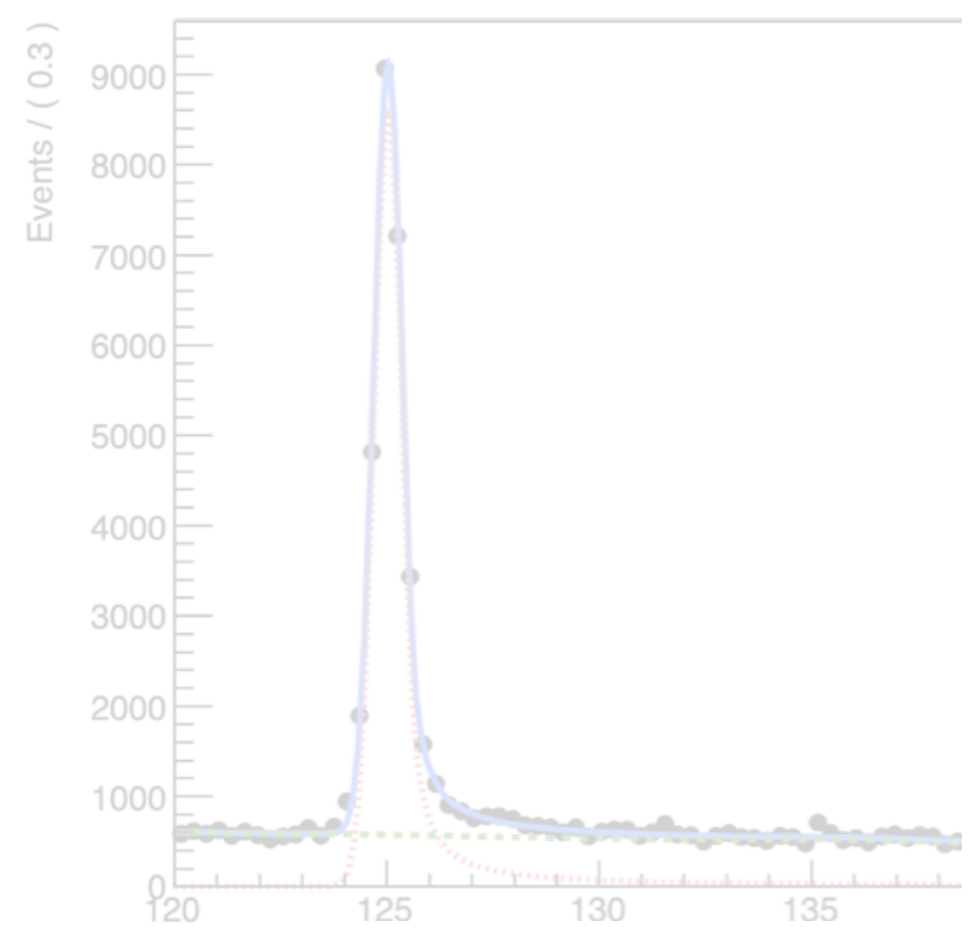
Broader resonant peak

Inclusive analyses @240 GeV

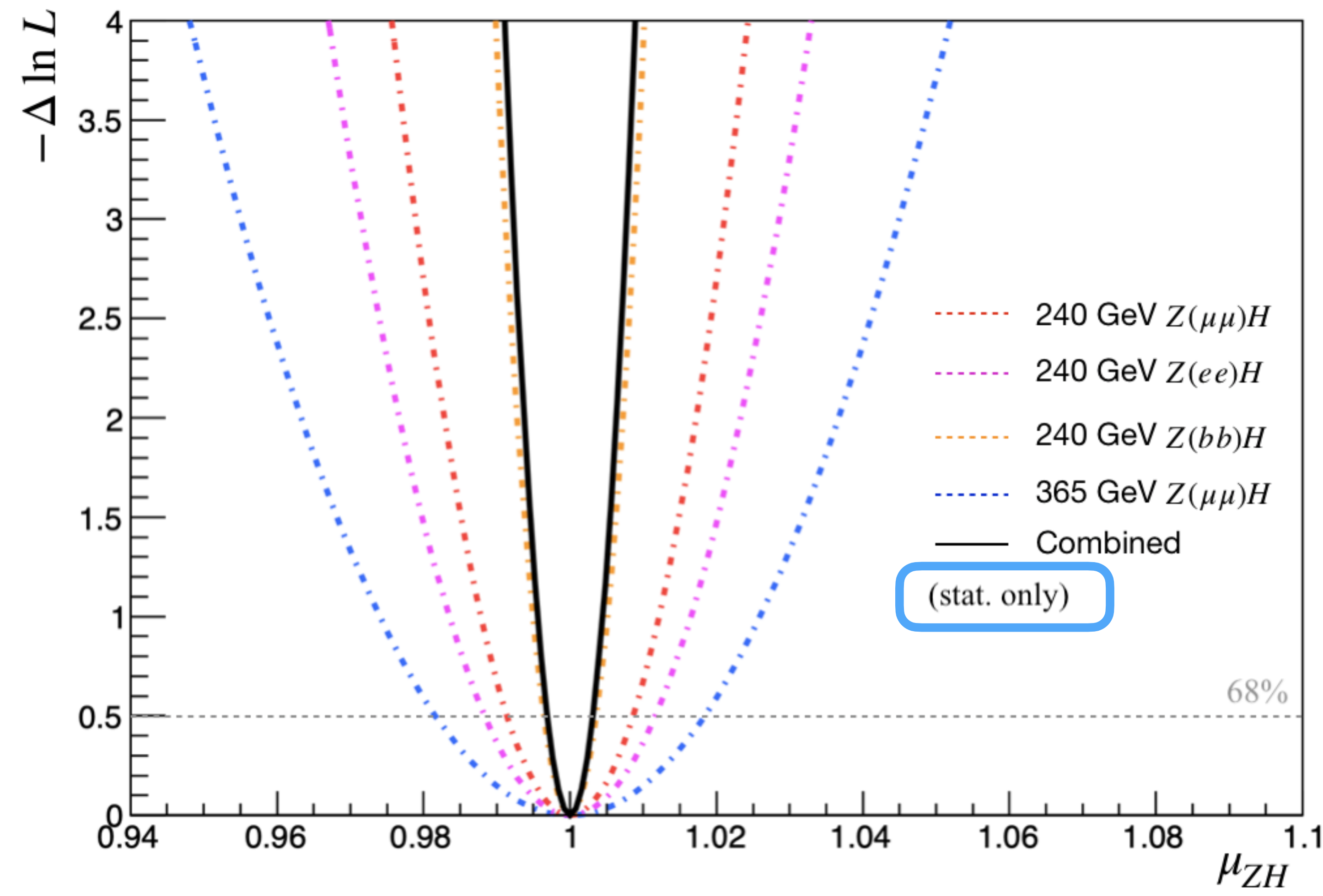
Exploited various Z decays, using the recoil techniques

Z($\mu\mu$)H

- ▶ $\mu^+\mu^-$ with $p_{T\mu 1} > 20$
- ▶ Minimum $|M_{\mu^+\mu^-} - M_{ZH}|$
- ▶ $80 < M_{\mu^+\mu^-} < 100$ GeV

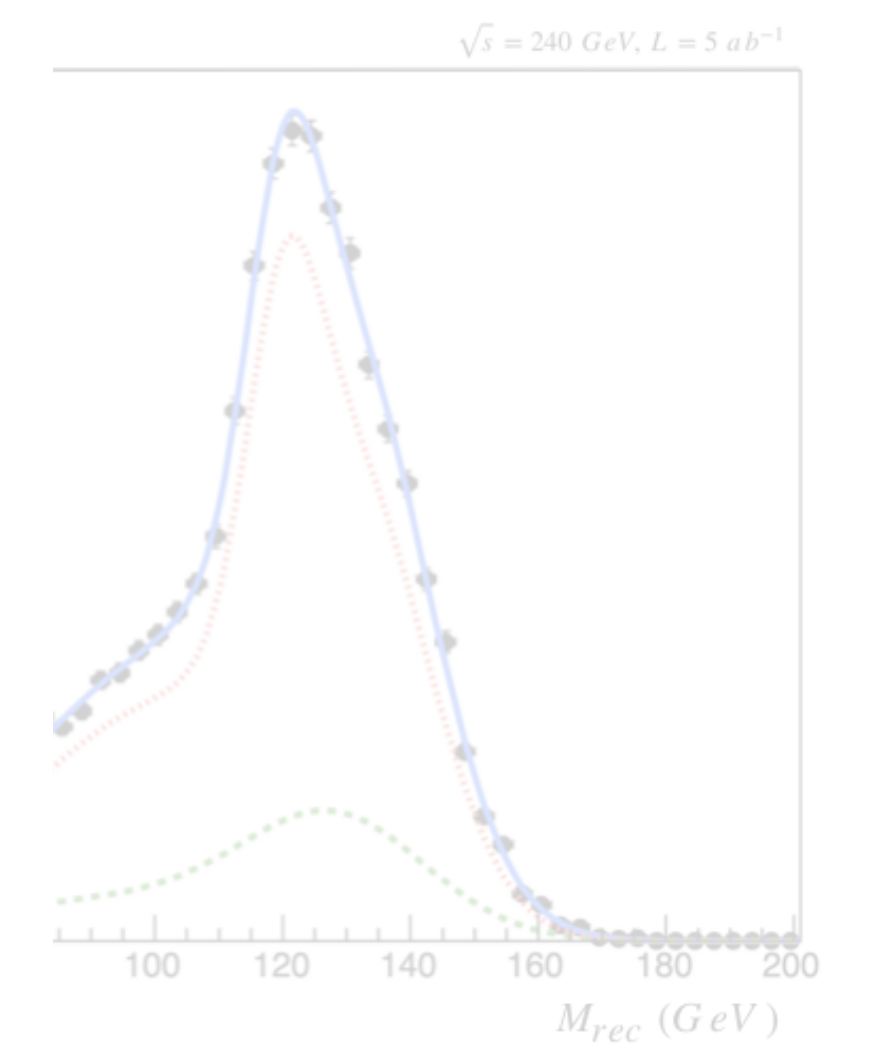


Golden char



Z(bb)H

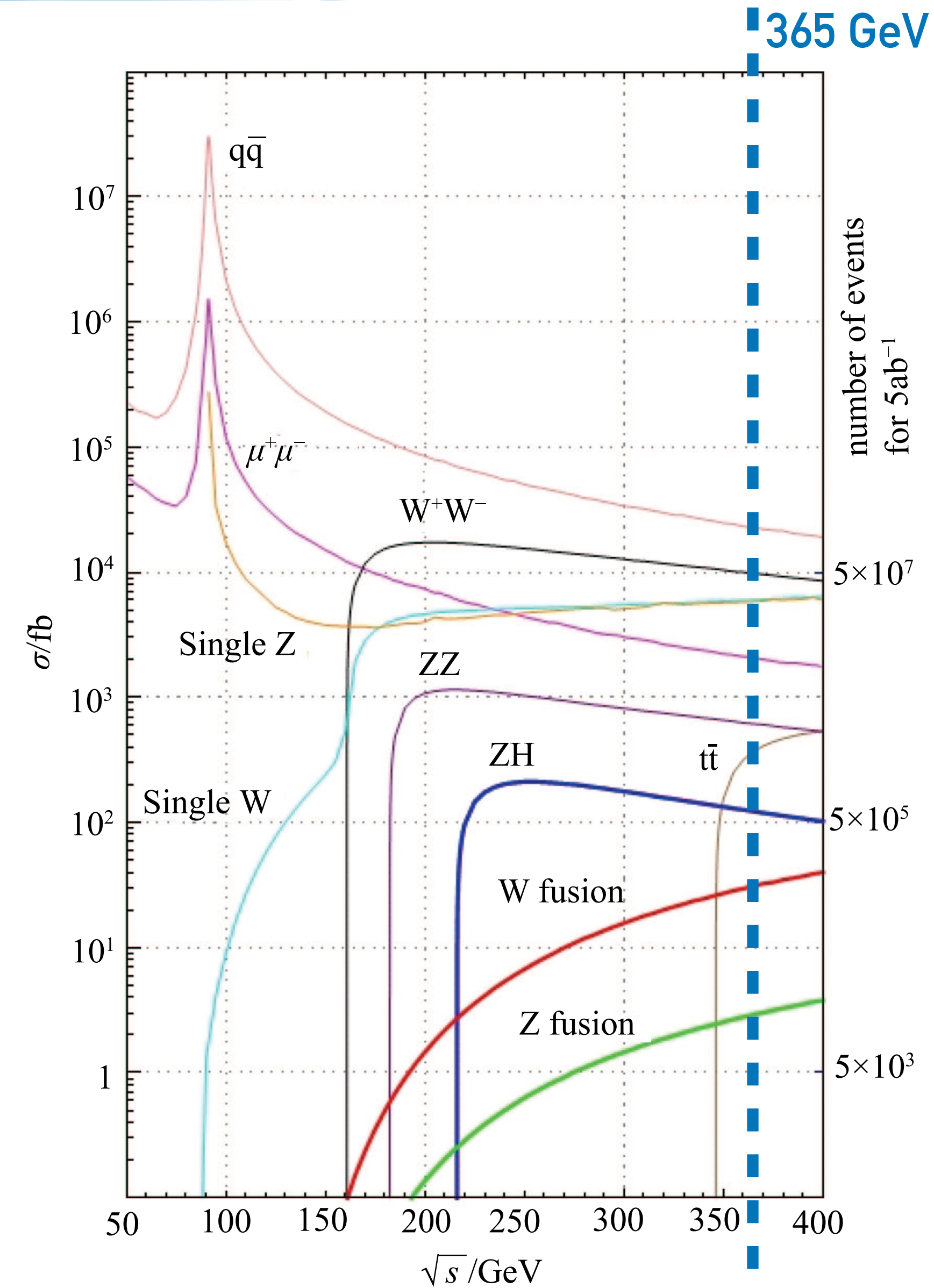
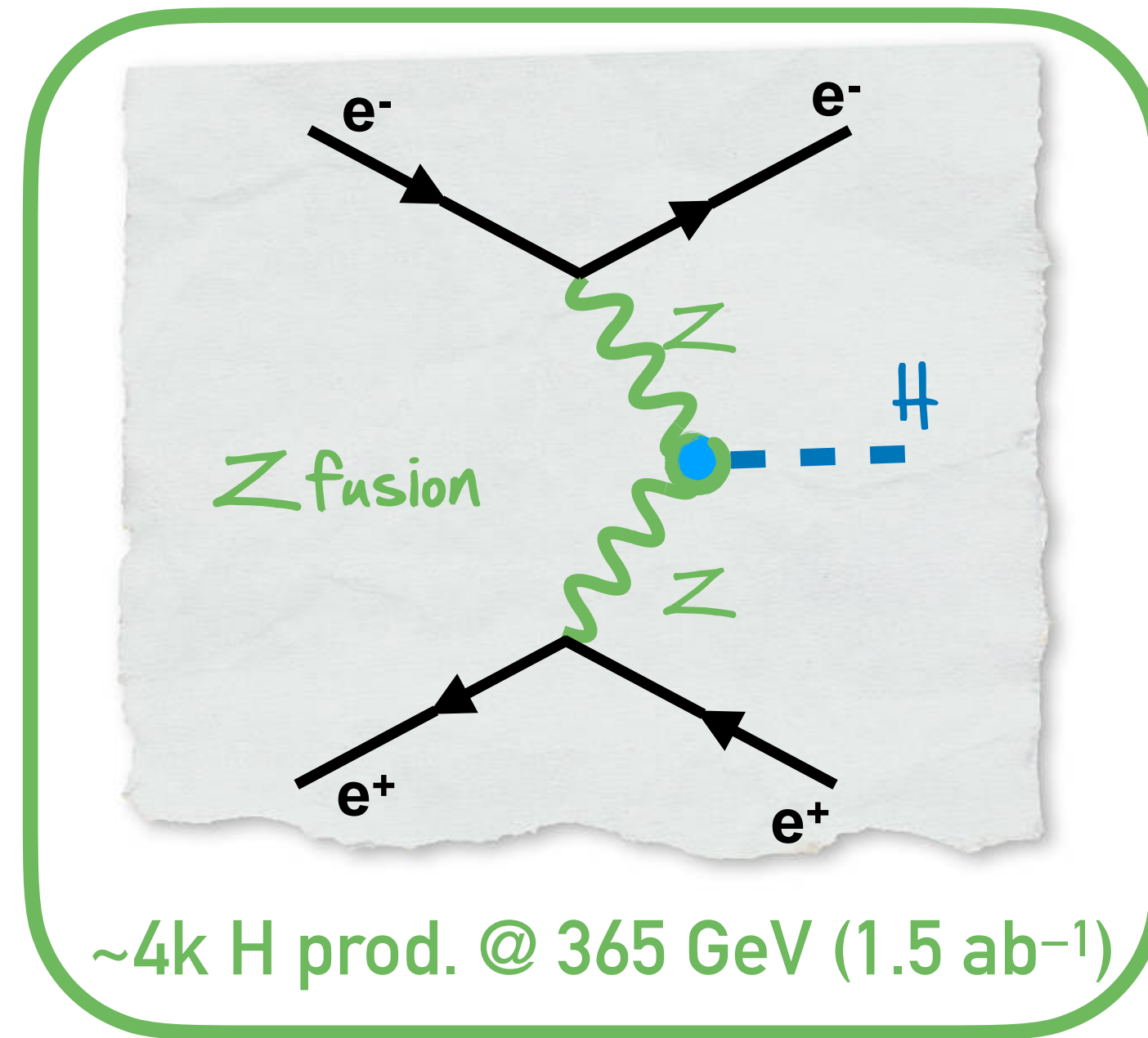
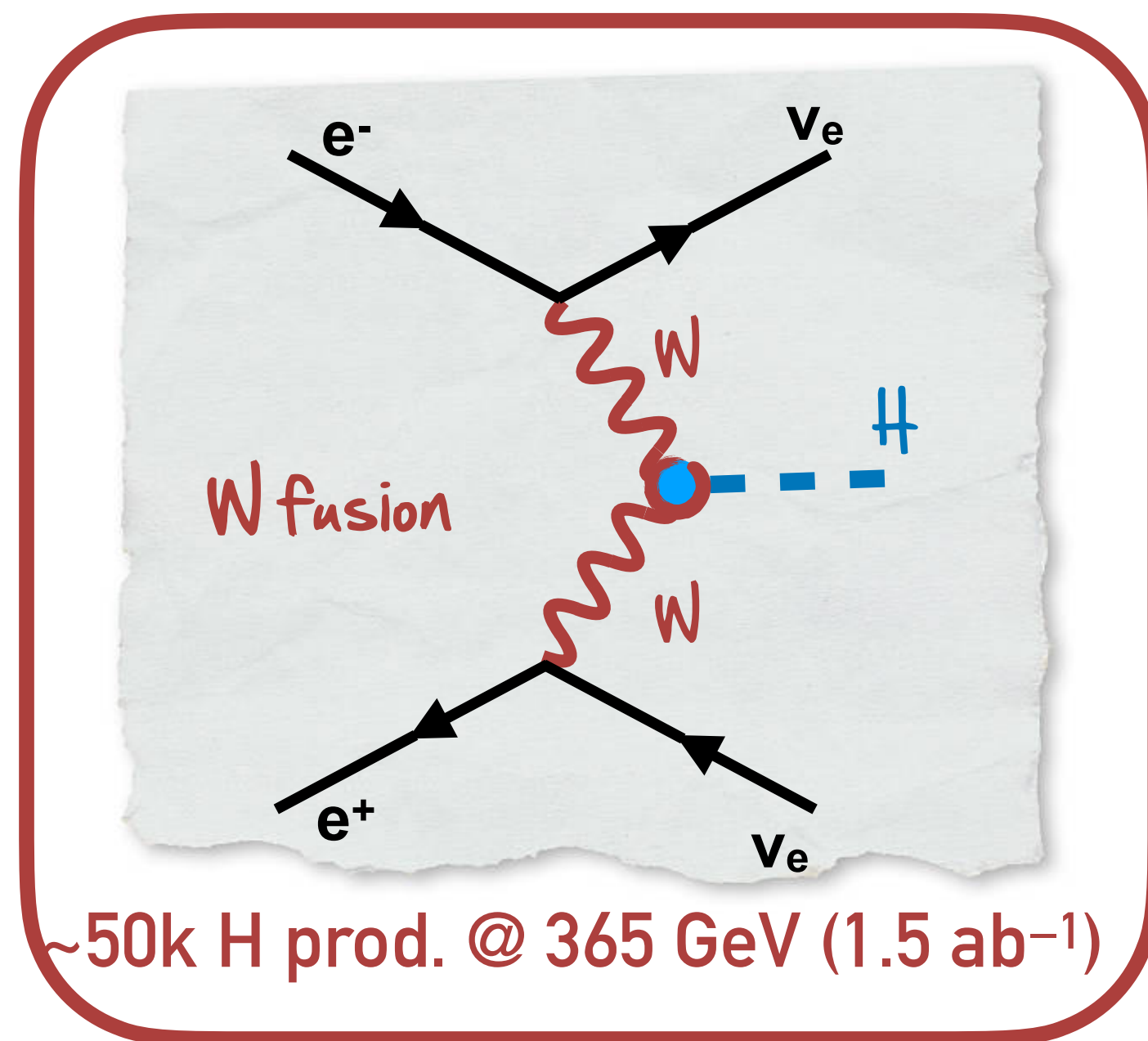
- ▶ $ets + p_{Tjj} > 60$ GeV
- ▶ 5 GeV
- ▶ 1 GeV
- ▶ 7 variables)



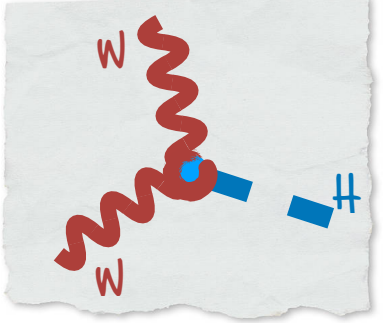
resonant peak

VBF production @365 GeV

The production channel starts to become relevant due to the logarithmic raise $\sim \ln^2(s/M_V^2)$ of the t-channel exchange of vector bosons

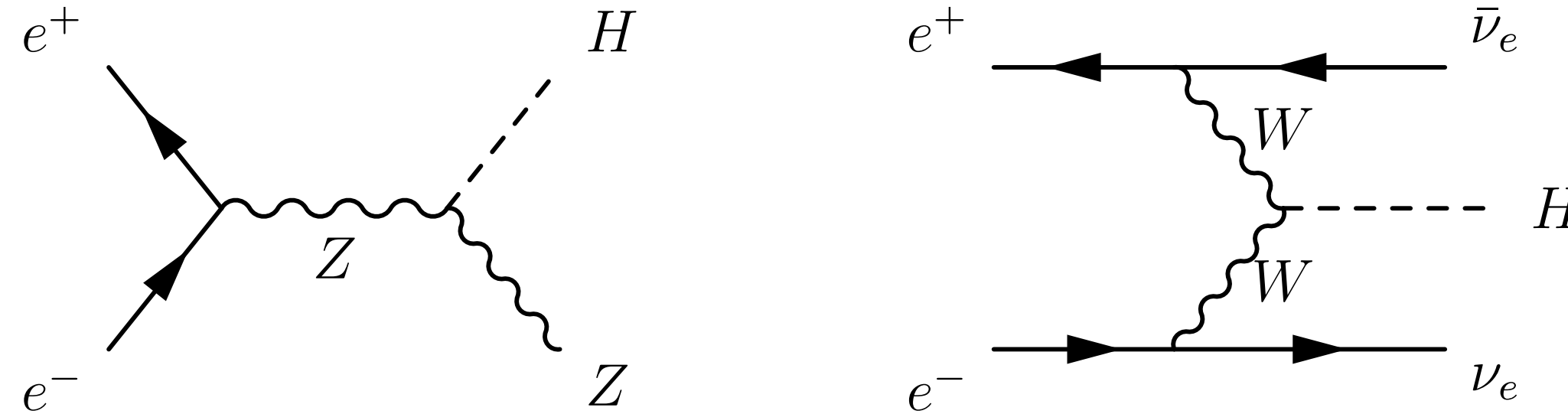


Production dominated by the W fusion because of larger charged currents



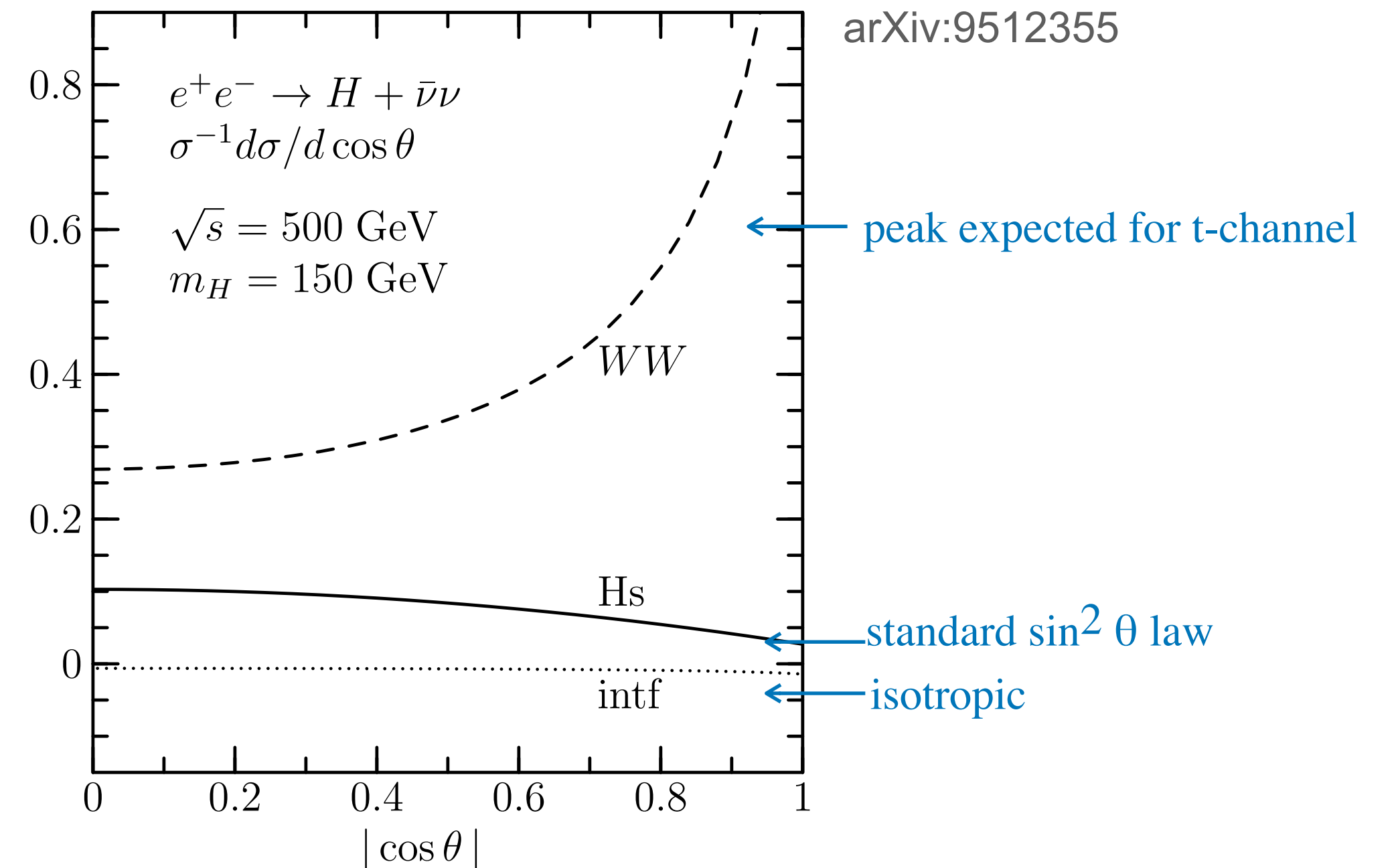
WW-boson fusion : $ee \rightarrow \nu\nu H(bb)$

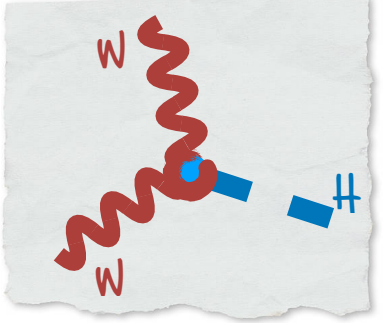
Two production mechanisms contribute



For $\nu_e \bar{\nu}_e$ decays of the Z boson, the two production amplitudes interfere
 Positive interference term of the same size as their individual cross sections

Need to exploit angular distribution
 to separate the processes





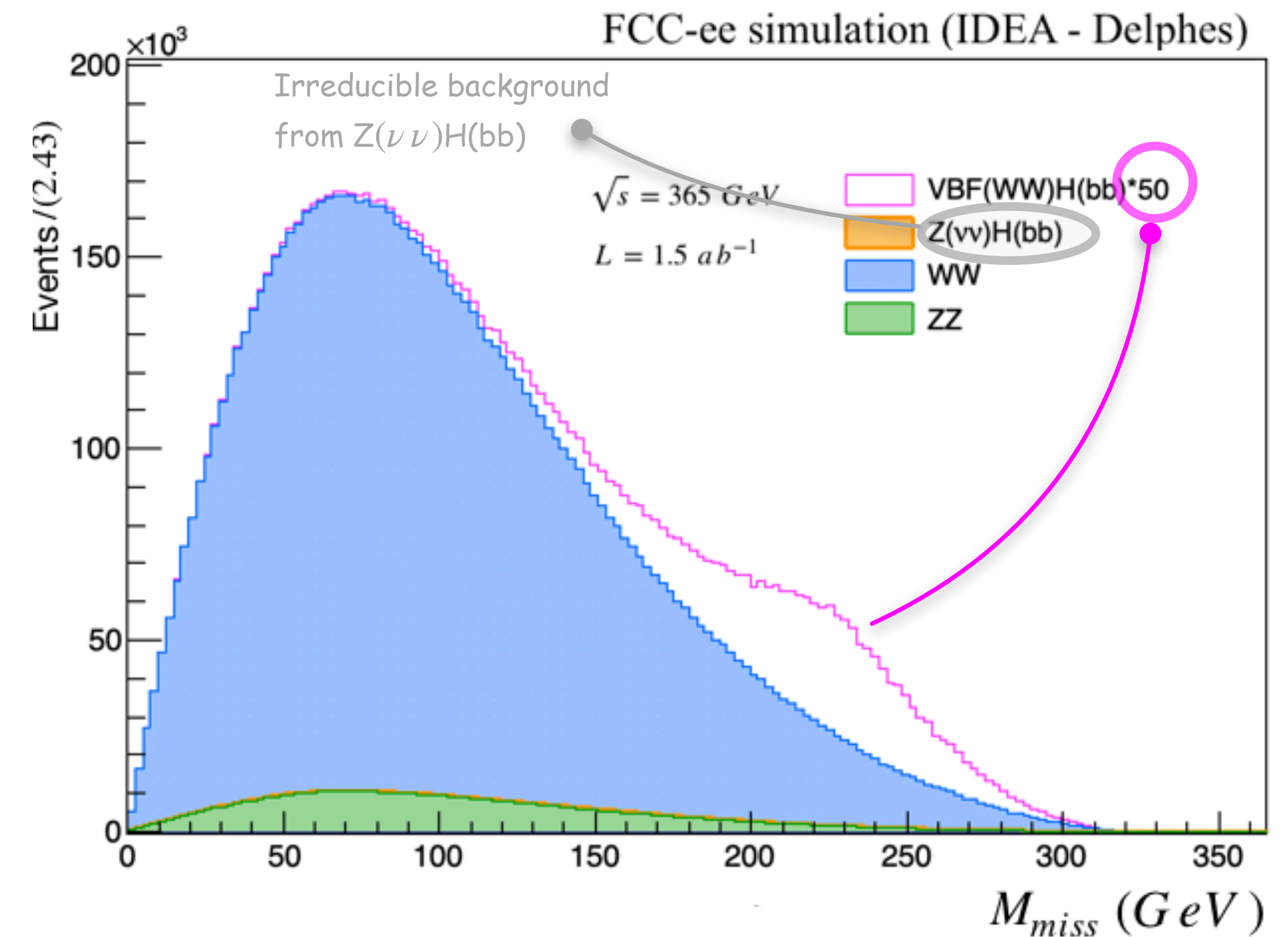
WW-boson fusion : $ee \rightarrow \nu\nu H(bb)$

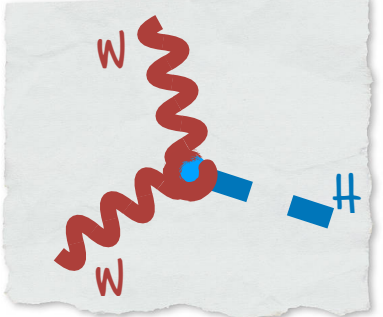
1. Preselection cuts

- 2 b-jets, $|\eta_{jj}| < 3$
- $H_T > 10$ GeV
- $MET > 10$ GeV

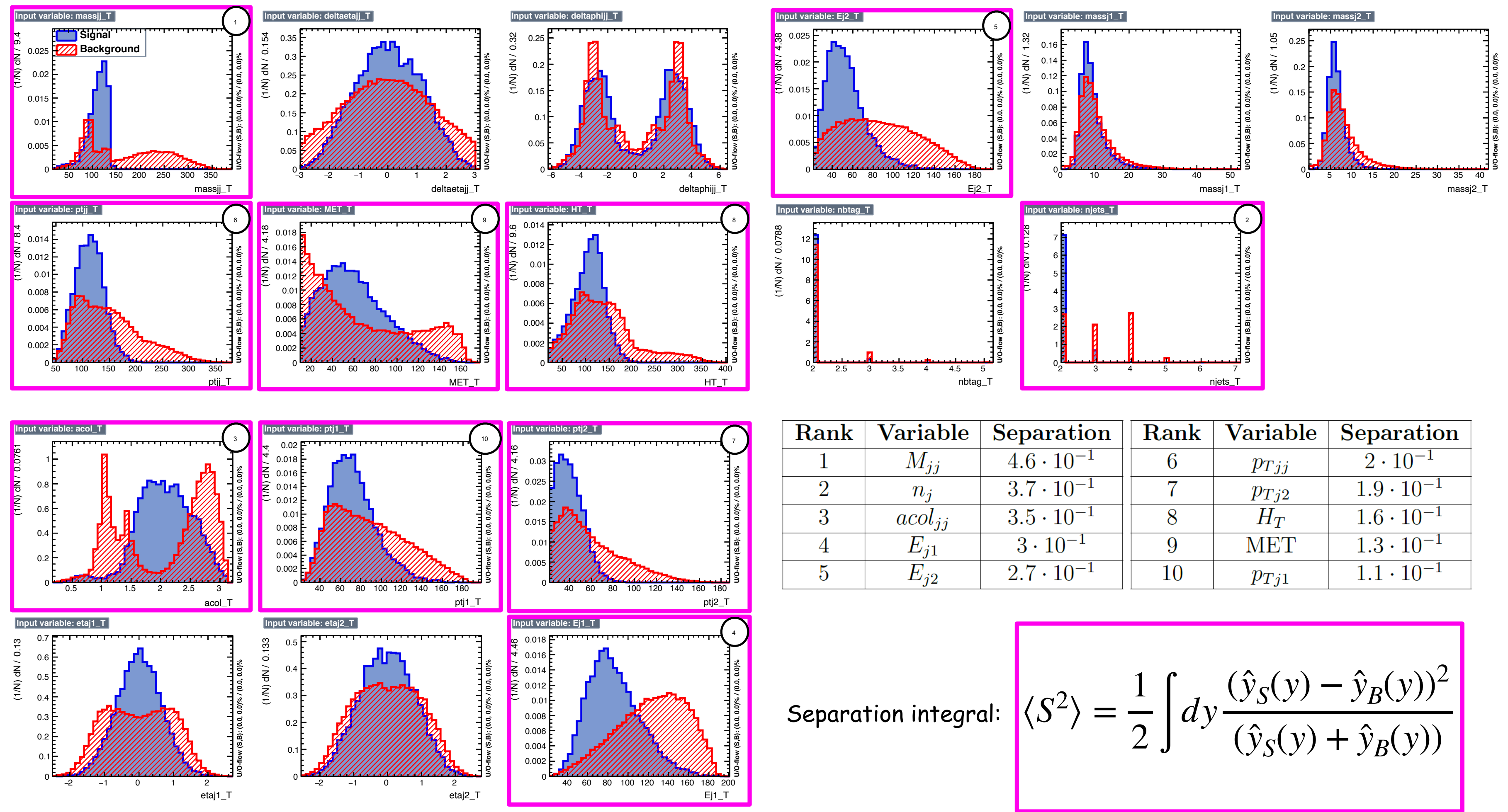
2. Adaptive BDT to reduce the backgrounds

- 17 input variables
- trained with a 20k sig. and 100k back. events
- 800 trees, min. node size of 1%, a max. depth of 3

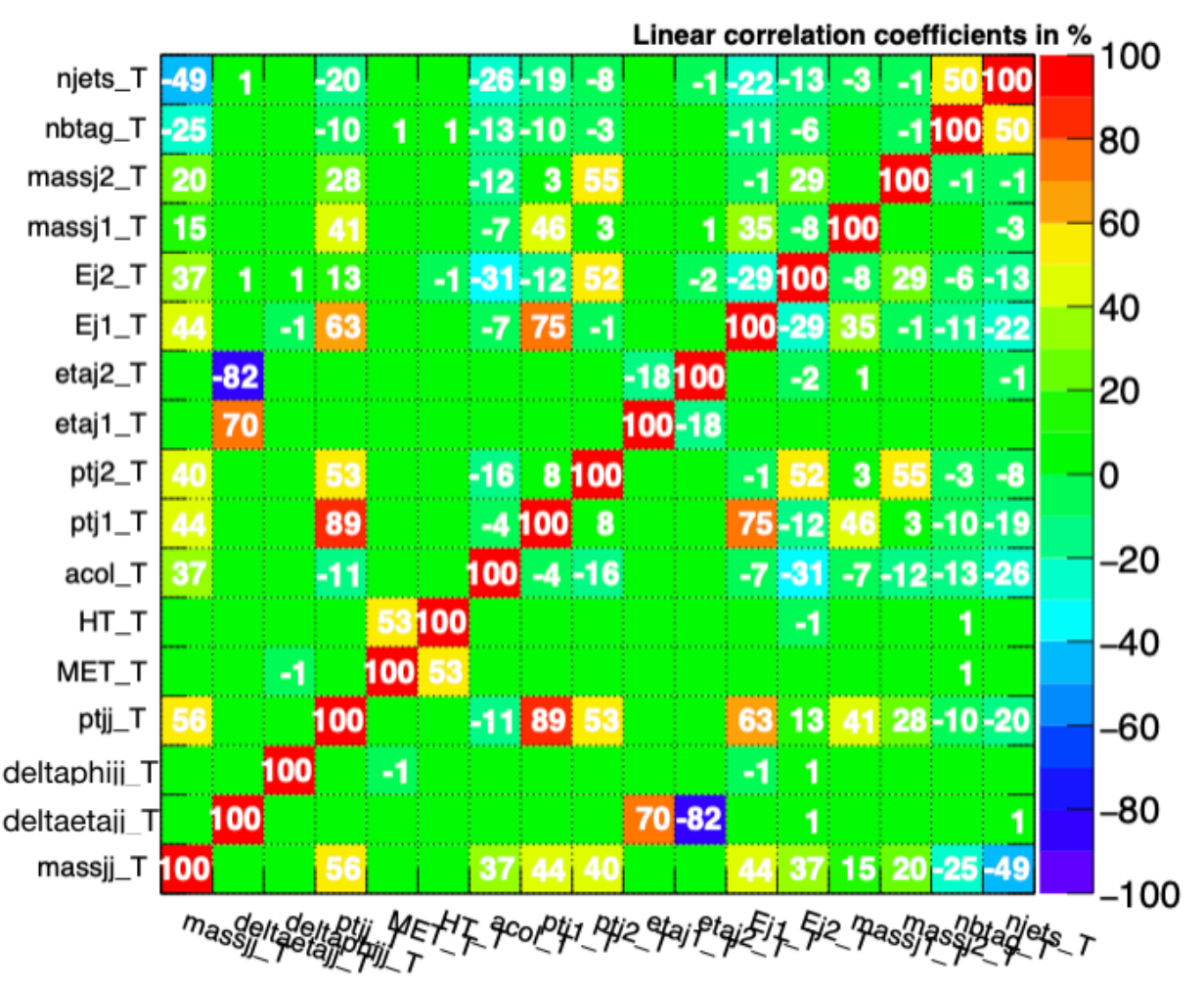




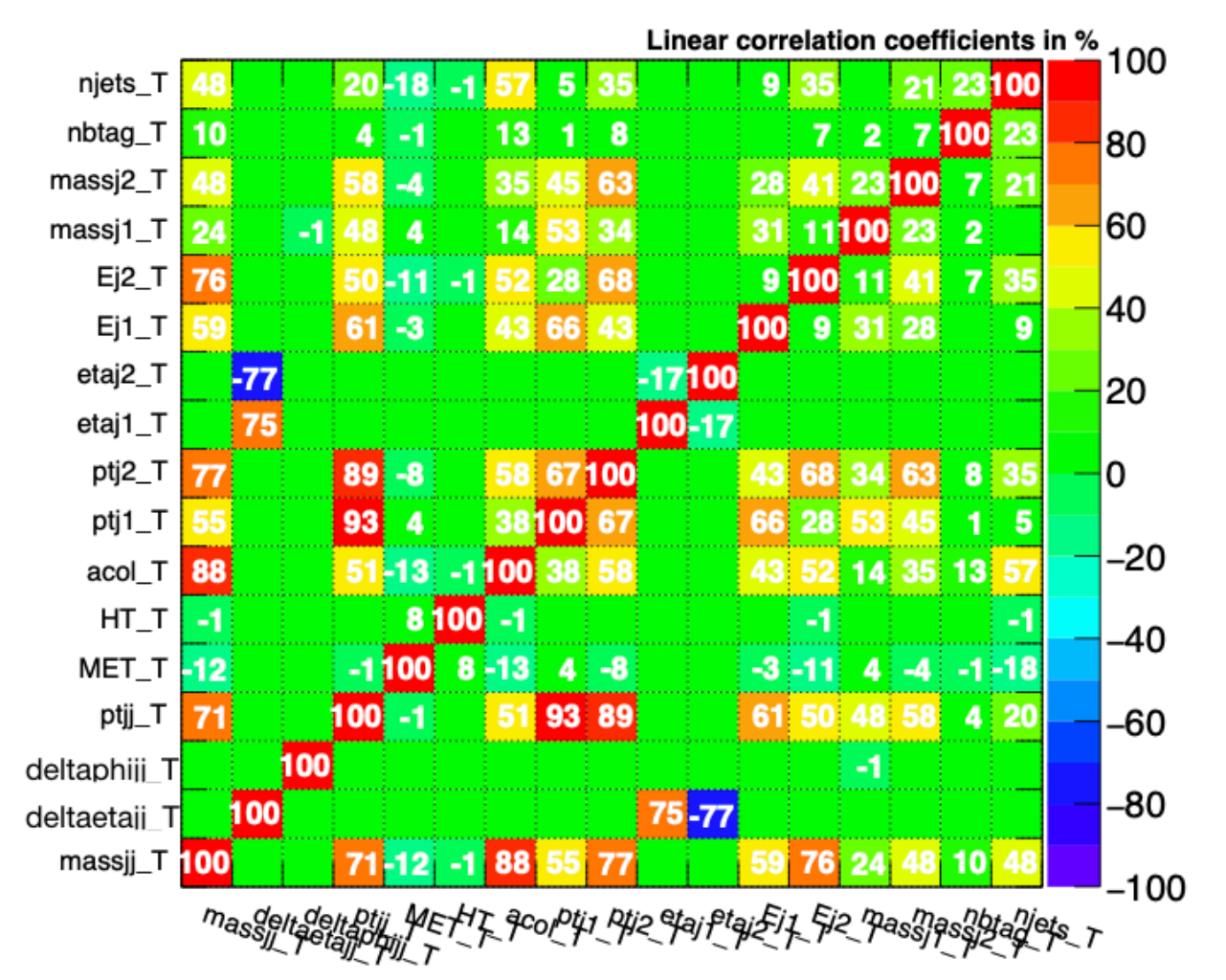
BDT variables and correlations



Correlation Matrix (signal)



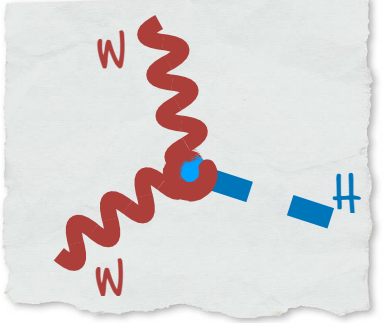
Correlation Matrix (background)



Rank	Variable	Separation	Rank	Variable	Separation
1	M_{jj}	$4.6 \cdot 10^{-1}$	6	p_{Tjj}	$2 \cdot 10^{-1}$
2	n_j	$3.7 \cdot 10^{-1}$	7	p_{Tj2}	$1.9 \cdot 10^{-1}$
3	$acol_{jj}$	$3.5 \cdot 10^{-1}$	8	H_T	$1.6 \cdot 10^{-1}$
4	E_{j1}	$3 \cdot 10^{-1}$	9	MET	$1.3 \cdot 10^{-1}$
5	E_{j2}	$2.7 \cdot 10^{-1}$	10	p_{Tj1}	$1.1 \cdot 10^{-1}$

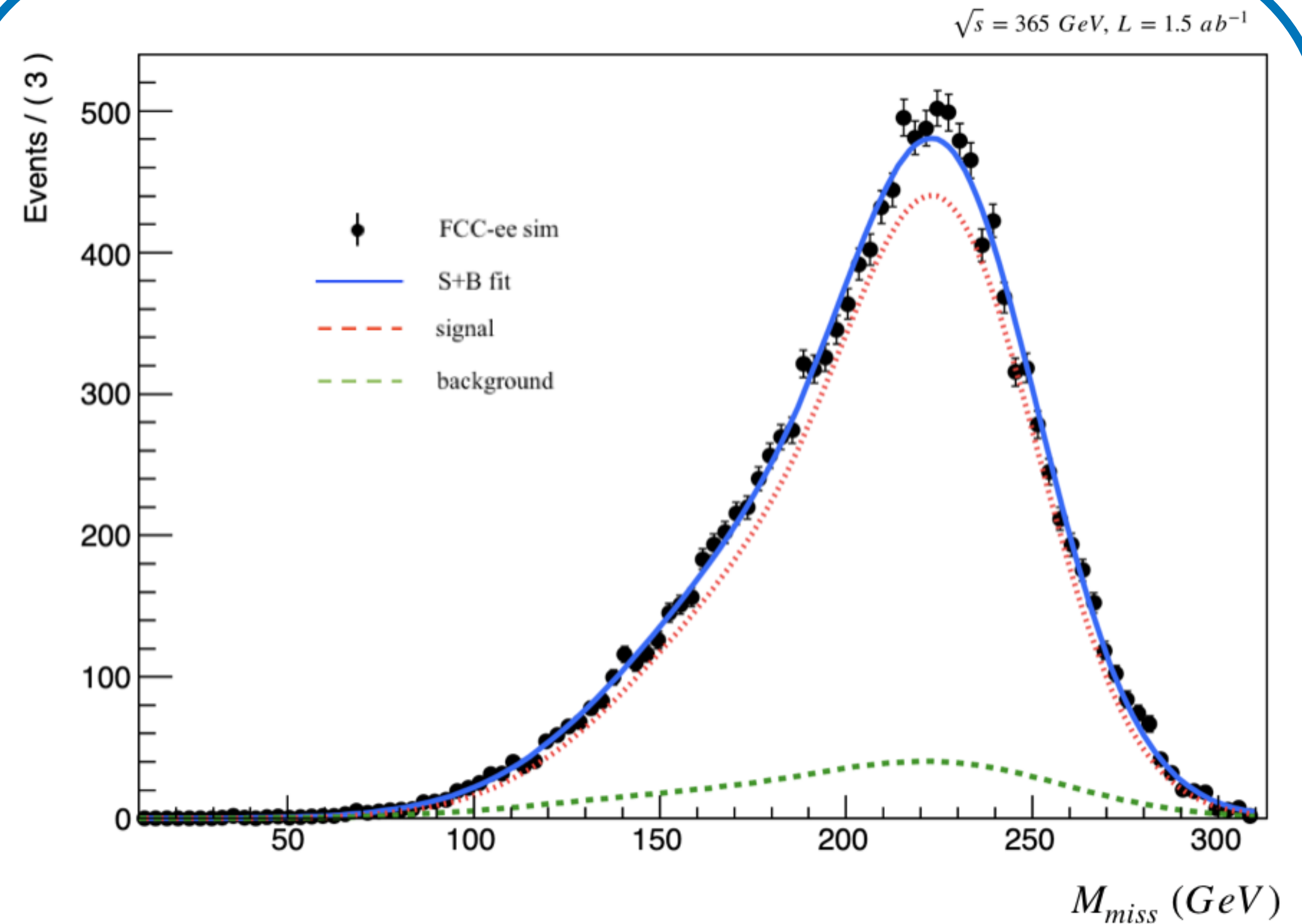
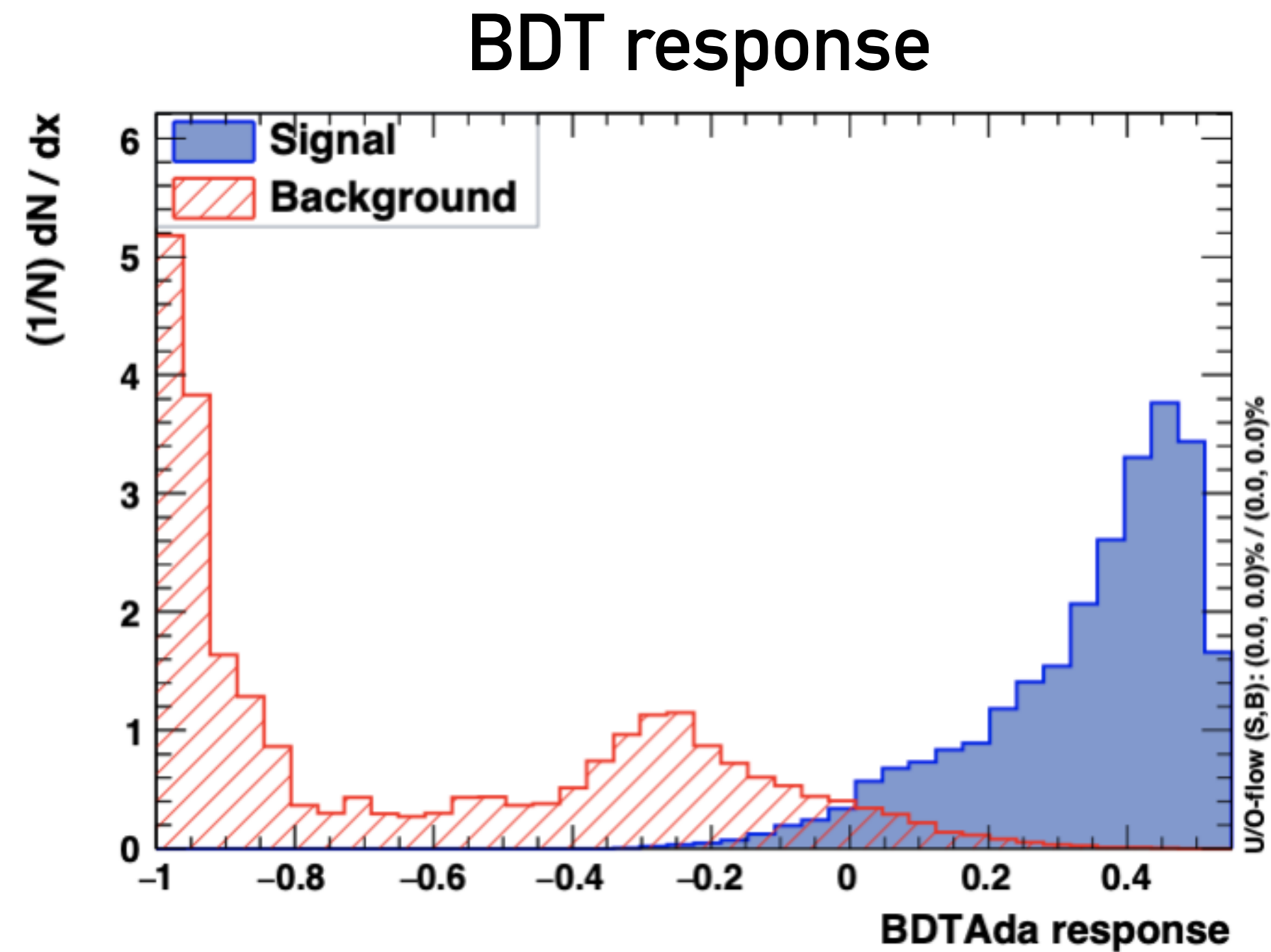
Separation integral:

$$\langle S^2 \rangle = \frac{1}{2} \int dy \frac{(\hat{y}_S(y) - \hat{y}_B(y))^2}{(\hat{y}_S(y) + \hat{y}_B(y))}$$



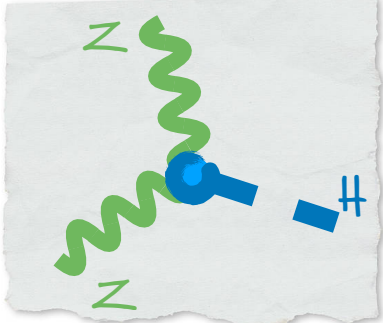
Final discrimination variable

The missing mass after preselection and BDT cuts



Signal (convolution of 2 Gaussians)
 Irreducible background mainly from $Z(\nu\nu)H(bb)$

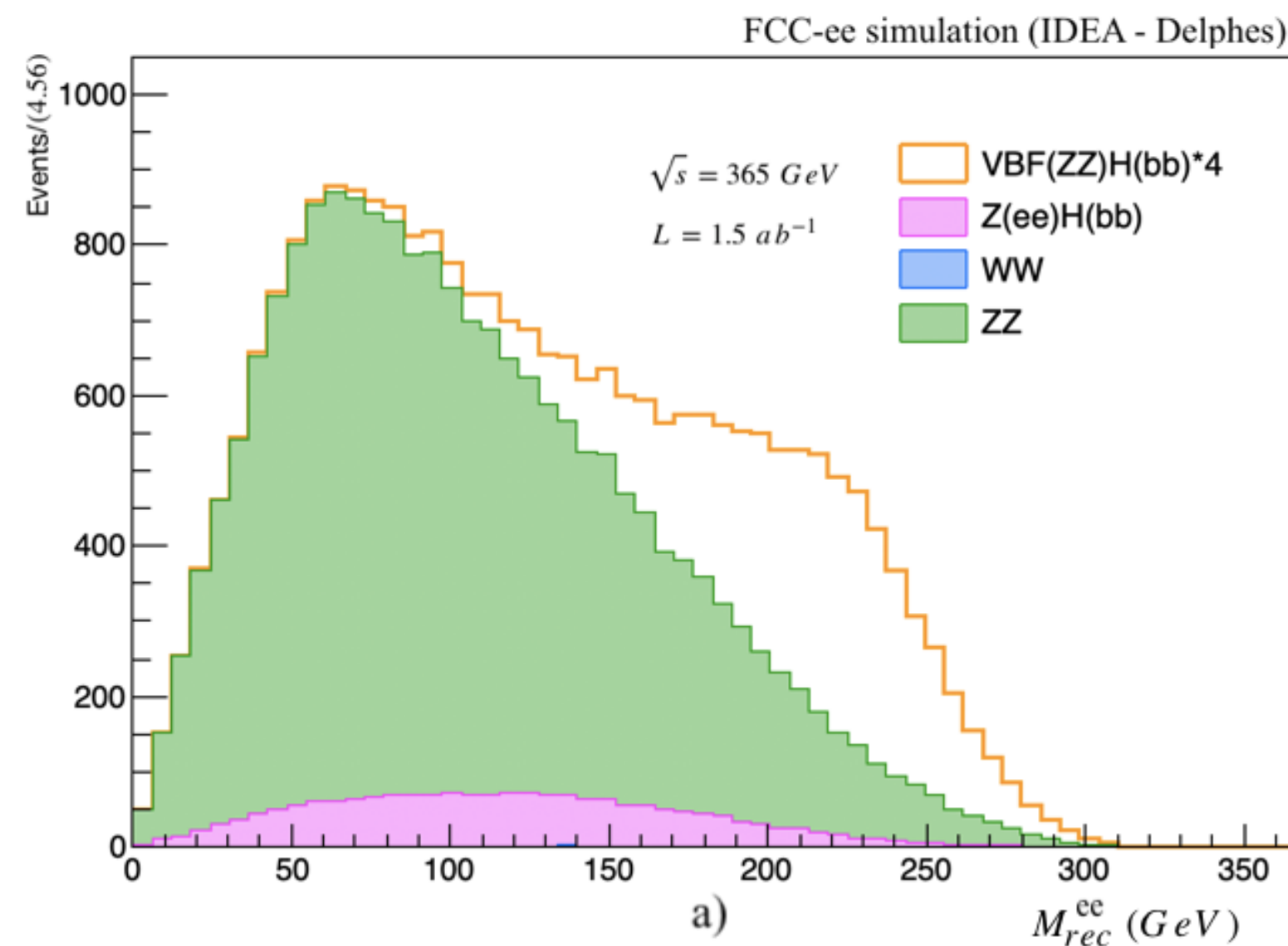
MC samples	$\nu_e\bar{\nu}_e H(bb)$	$Z(\nu\bar{\nu})H(bb)$	WW	ZZ
Number of events (normalized)	$3.05 \cdot 10^4$	$2.06 \cdot 10^4$	$1.61 \cdot 10^7$	$9.49 \cdot 10^5$
$n_{bj} \geq 2, \Delta\eta < 3, HT > 20, MET > 10 \text{ GeV}$	47%	48%	0.09%	5.5%
BDTAda response ≥ 0.12	42 %	3.4 %	0.002 %	0.06 %



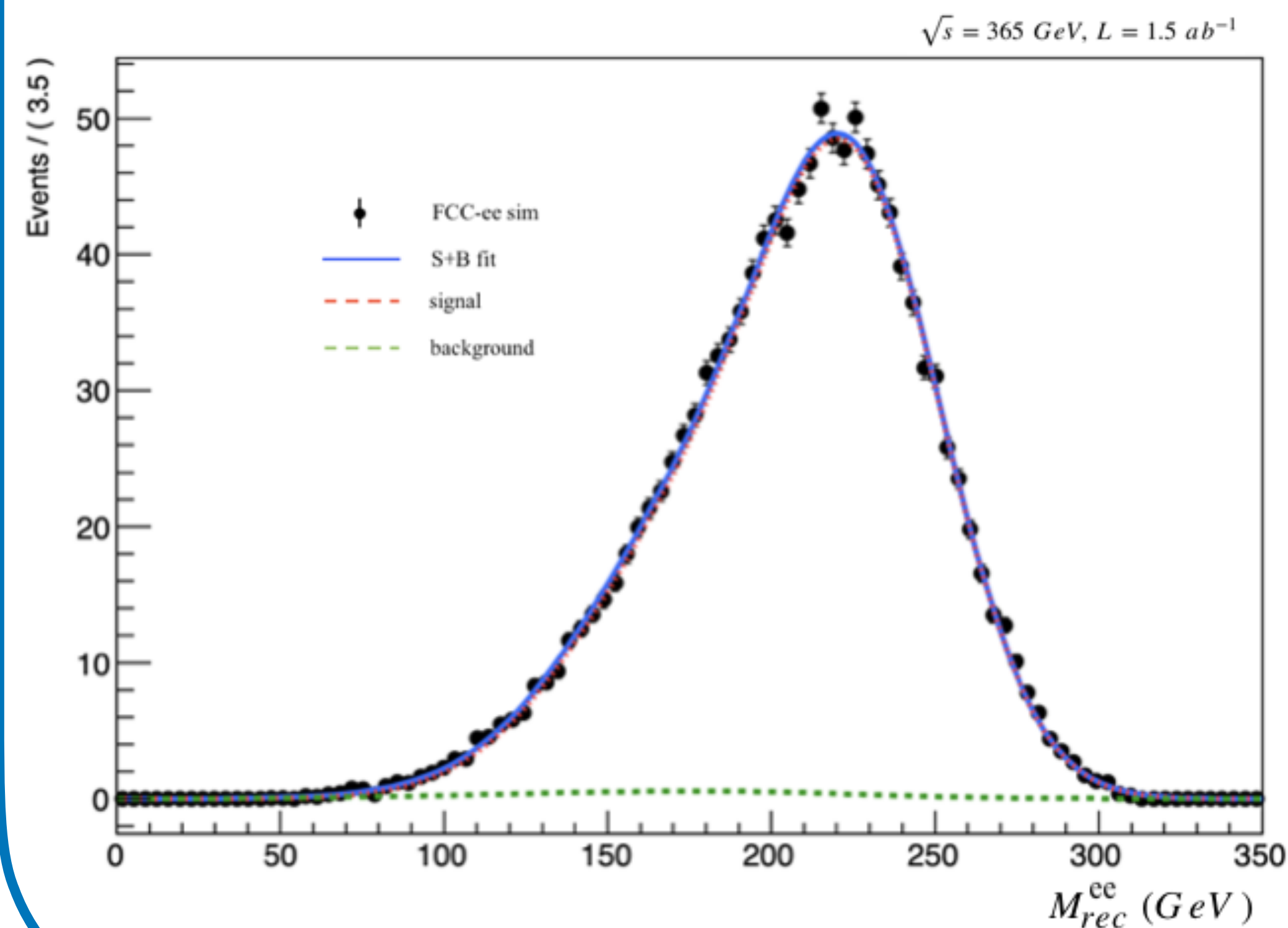
ZZ-boson fusion : $ee \rightarrow eeH(bb)$

1. Preselection cuts

- 2 jets + 2 electrons
- $m_{ee} > 80 \text{ GeV}$
- $\text{MET} > 10 \text{ GeV}$

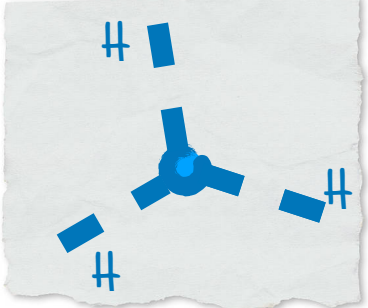


Fit to the recoil mass spectrum in after the BDT



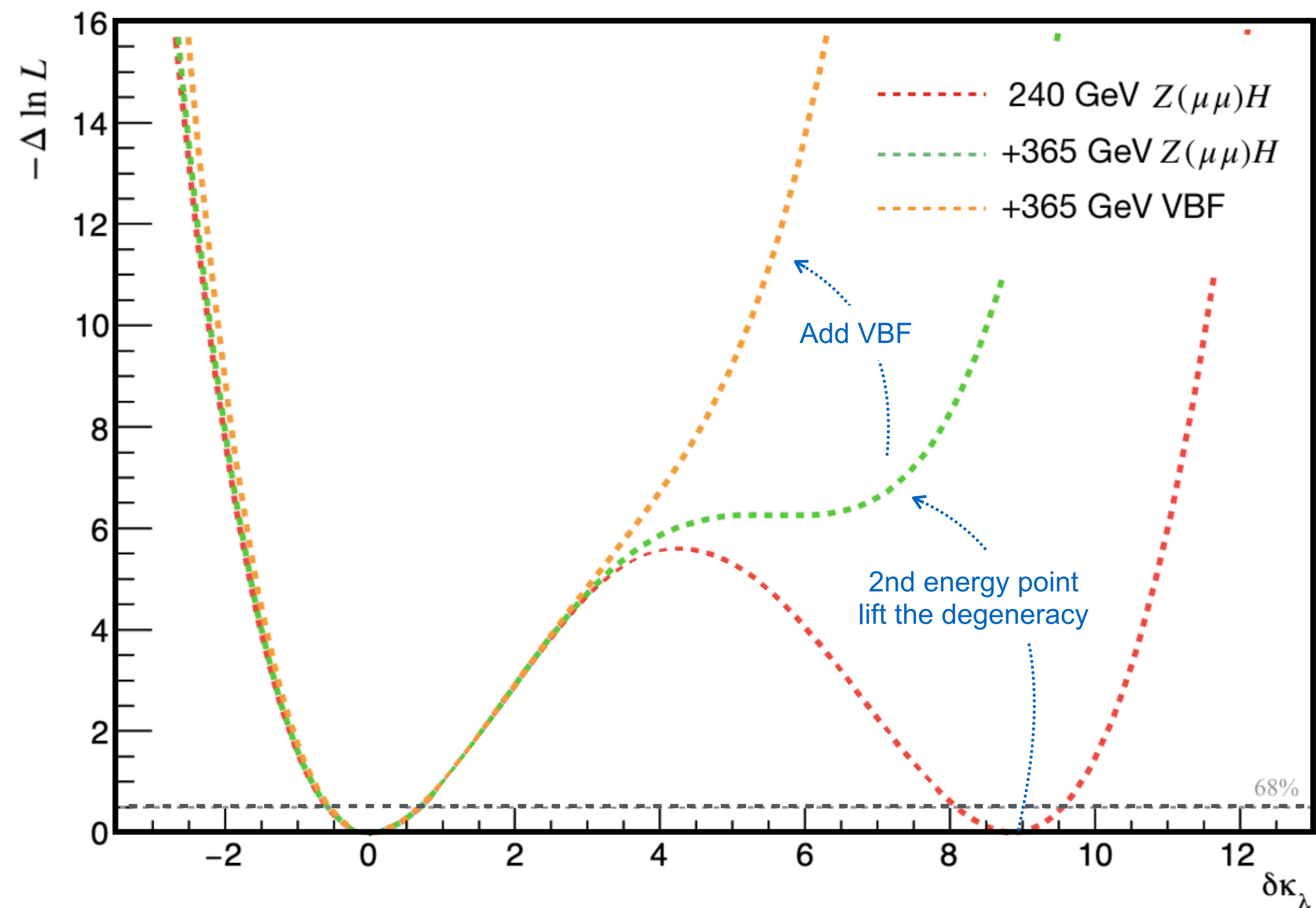
2. BDT to further reduce the backgrounds

Rank	Variable	Separation	Rank	Variable	Separation
1	$M_{e^+e^-}$	$9.1 \cdot 10^{-1}$	5	M_{jj}	$3.8 \cdot 10^{-1}$
2	$acol_{e^+e^-}$	$7.1 \cdot 10^{-1}$	6	η_{e2}	$2.4 \cdot 10^{-1}$
3	$acol_{jj}$	$7 \cdot 10^{-1}$	7	E_{j1}	$2.1 \cdot 10^{-1}$
4	n_{bj}	$4.6 \cdot 10^{-1}$	8	η_{j1}	$1.4 \cdot 10^{-1}$



Putting all together

1D fit with only $\delta\kappa_\lambda$ floating



The secondary minimum easily excluded adding a 2nd energy point

Thoughts on “detector requirements”

Statistics is the essence of the Higgs boson self-coupling studies

Access to various production mechanisms and two energies points are great opportunities

We need to include

- > hadronic Z decays in inclusive analyses
- > highest H BR channel ($H \rightarrow bb$) in WW-fusion and ZZ-fusion channels
- > exploit angular distribution(s) to better separate HZ and VBF channels

it means

- > Efficient flavour tagging
- > Optimal jet angular and energy resolutions

Conclusions

Preliminary results based on the work done during a M2 internship have been shown.

The analysis chain has been put in place to measure the Higgs boson self-coupling from higher-order corrections to single-Higgs processes

The analyses are going to be redone (improved selection, adding systematics, ...) using the centrally produced samples within the FCCAnalyses framework.