

# Short-range correlations and the phenomenological determination of nuclear PDFs

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# Introduction

1. One of the standard ways of parametrizing nuclear PDFs (nPDFs) is by extending the proton PDF parametrizations to account for  $A$ -dependence.
2. E.g. in the nCTEQ group:
  - ▶ *PDF of nucleus* ( $A$  - mass,  $Z$  - charge,  $N$  - number of neutrons)

$$f_i^{(A,Z)}(x, Q) = \frac{Z}{A} f_i^{p/A}(x, Q) + \frac{N}{A} f_i^{n/A}(x, Q)$$

- ▶ bound proton PDFs are parametrized

$$x f_i^{p/A}(x, Q_0) = x^{c_1} (1-x)^{c_2} P(x, \{c_k\})$$

- ▶ bound neutron PDFs are constructed assuming *isospin symmetry*
- ▶  $A$ -dependence

$$c_k \rightarrow c_k(A) \equiv p_k + a_k \left(1 - A^{-b_k}\right)$$

## 3. Sum rules

$$\int_0^1 dx f_{u_v}^{p/A}(x, Q) = 2, \quad \int_0^1 dx f_{d_v}^{p/A}(x, Q) = 1, \quad \int_0^1 dx \sum_i x f_i^{p/A}(x, Q) = 1.$$

# Introduction

Let's consider a parametrization inspired by nuclear dynamics, one option are models based on: **Short Range Correlations** (SRC).

- ▶ SRC pairs can have isospin  $I = 0, 1$ , possible configurations:  $(pn)$ ,  $(pp)$ ,  $(nn)$
- ▶ Partonic content of SRC pairs could be expressed as a convolution of distributions of a parton inside a nucleon and a nucleon inside a pair, then the distribution of the full nucleus:

$$f_i^A = \frac{Z}{A} \left[ (1 - [z_A^{(pp)} + z_A^{(pn)}]) f_{i/p} + z_A^{(pp)} f_{\text{SRC}}^{p/(pp)} \otimes f_{i/p} + z_A^{(pn)} f_{\text{SRC}}^{p/(pn)} \otimes f_{i/p} \right] \\ + \frac{N}{A} \left[ (1 - [n_A^{(nn)} + n_A^{(pn)}]) f_{i/n} + n_A^{(nn)} f_{\text{SRC}}^{n/(nn)} \otimes f_{i/n} + n_A^{(pn)} f_{\text{SRC}}^{n/(pn)} \otimes f_{i/n} \right]$$

- ▶ For phenomenological purpose we can simplify it assuming:

$$f_{i/p}^{\text{SRC}} \equiv [f_{\text{SRC}}^{p/(pp)} + f_{\text{SRC}}^{p/(pn)}] \otimes f_{i/p} \qquad z_A \equiv z_A^{(pp)} + z_A^{(pn)} \\ f_{i/n}^{\text{SRC}} \equiv [f_{\text{SRC}}^{n/(nn)} + f_{\text{SRC}}^{n/(pn)}] \otimes f_{i/n} \qquad n_A \equiv n_A^{(nn)} + n_A^{(pn)}$$

- ▶ As a consequence we will be able to determine only total number of paired neutrons and protons.

# SRC parametrization

Our phenomenological SRC inspired parametrization takes form:

$$f_i^A(x, Q) = \frac{Z}{A} \left[ (1 - z_A) f_{i/p}(x, Q) + z_A f_{i/p}^{\text{SRC}}(x, Q) \right] \\ + \frac{N}{A} \left[ (1 - n_A) f_{i/n}(x, Q) + n_A f_{i/n}^{\text{SRC}}(x, Q) \right]$$

with  $f_{i/p}(f_{i/n})$  being the free proton (neutron) PDFs and  $f_{i/p}^{\text{SRC}}(f_{i/n}^{\text{SRC}})$  the effective SRC proton (neutron) distributions.

The full nPDF  $f_i^A$  need to fulfill:

1. DGLAP evolution.
2. Momentum and number sum rules:

$$\int_0^1 dx x f_i^A(x, Q) = 1, \quad \int_0^1 dx f_{u_v}^A(x, Q) = \frac{A + Z}{A}, \quad \int_0^1 dx f_{d_v}^A(x, Q) = \frac{A + N}{A}.$$

We assume that both  $f_{i/n}$  and  $f_{i/n}^{\text{SRC}}$  can be determined using isospin symmetry. We also restrict  $f_{i/p}^{\text{SRC}}(f_{i/n}^{\text{SRC}})$  (and  $f_i^A$ ) to be define on  $x \in (0, 1)$ , then  $f_{i/p}^{\text{SRC}}(f_{i/n}^{\text{SRC}})$ :

- ▶ fulfill DGLAP evolution equation,
- ▶ obey the same sum rules as free proton (neutron) distributions.

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For the purpose of global analysis we:

- ▶ fix the free proton PDFs to the nCTEQ15 proton,
- ▶ parametrize the SRC PDFs as:

$$x f_i^{p/A}(x, Q_0) = x^{c_1} (1 - x)^{c_2} e^{c_3 x} (1 + e^{c_4 x})^{c_5}$$

Free parameters:

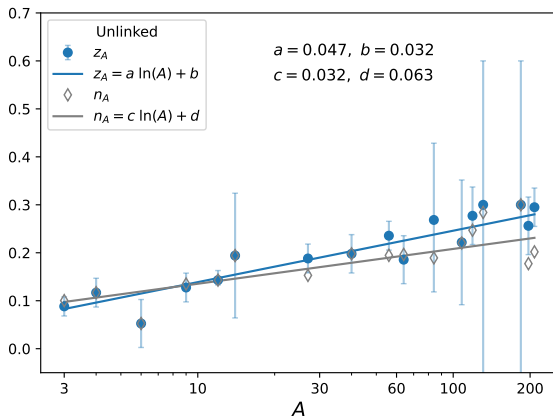
- ▶  $x$ -shape: set of  $\{c_k\}$  parameters for each flavour (total of 21),
- ▶  $A$ -dependence: pairs of  $(z_A, n_A)$  parameters which are independent for each nuclei (instead we could use nuclear model to constrain them).

Used data:

- ▶ all DIS & DY data used in the nCTEQ15 analysis [[PRD 93, 085037 \(2016\)](#)],
- ▶ high- $x$  DIS data from JLAB which we used in the nCTEQ15hix analysis [[PRD 103, 114015 \(2021\)](#)],
- ▶  $p$ Pb data for  $W/Z$  production from the LHC used in the nCTEQ15WZ analysis [[EPJC 80, 968 \(2020\)](#)].

## Results: $A$ -dependence of the $(z_A, n_A)$ parameters

With such a setup we are able to obtain a very good fit to data with  $\chi^2/N_{\text{DOF}} = 0.8$ .

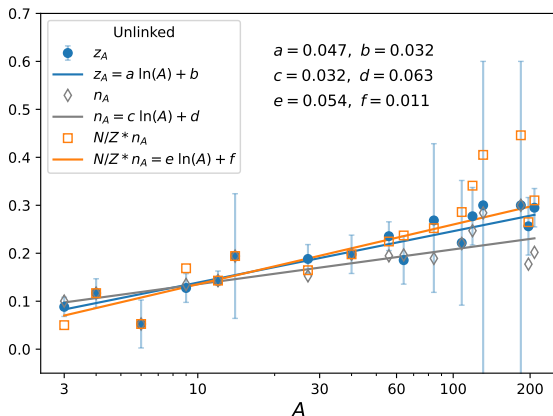


The number of protons and neutrons in SRC pairs is approximately equal, e.g.

- ▶  $^{197}_{79}\text{Au}$  ( $z_A=0.256, n_A=0.178$ ):  $79 \times z_A \simeq 20.2$  protons and  $118 \times n_A \simeq 21.0$  neutrons.
- ▶  $^{208}_{82}\text{Pb}$  ( $z_A=0.295, n_A=0.202$ ):  $82 \times z_A \simeq 24.2$  protons and  $126 \times n_A \simeq 25.5$  neutrons.

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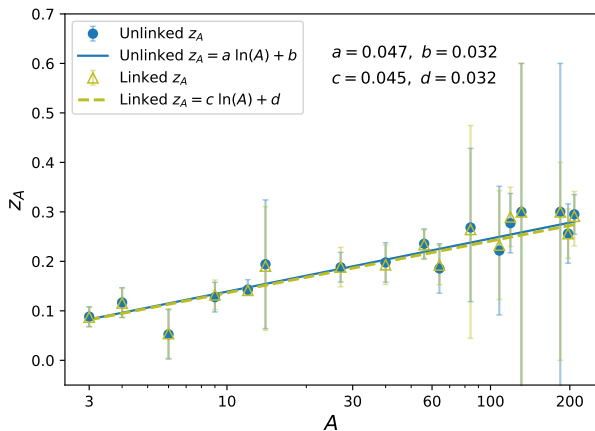
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- ▶ Correcting for the access of neutrons we obtained a very comparable numbers of protons and neutrons bounded in the SRC pairs.
- ▶ This is consistent with the hypothesis that the SRC pairs are dominantly proton-neutron combinations.
- ▶ We can use this observation to restrict number of fit parameters by linking  $n_A = (Z/N)z_A$ .



## Results - Linked fit $n_A = (Z/N)z_A$



- The obtained  $z_A$  values are nearly the same as for the **Unlinked** fit.
- Fit quality is very comparable  $\chi^2/N_{\text{DOF}} = 0.82$  (vs  $\chi^2/N_{\text{DOF}} = 0.8$ ).

# Results

- ▶ In order to judge the obtain results in context of nPDFs it is useful to compare them with nPDFs obtained using standard approach.
- ▶ We performed a “standard” fit using the same data and nCTEQ15-like parametrization.

$\chi^2/N_{\text{data}}$	DIS	DY	$W/Z$	JLab	$\chi^2_{\text{tot}}$	$\frac{\chi^2_{\text{tot}}}{N_{\text{DOF}}}$
<b>Reference</b>	0.85	0.97	0.88	0.72	1408	0.85
<b>SRC Unlinked</b>	0.84	0.75	1.11	0.41	1300	0.80
<b>SRC Linked</b>	0.85	0.84	1.14	0.49	1350	0.82

**Table:** The **Reference** fit has 19 shape and 3  $W/Z$  normalization parameters. The **Unlinked** and **Linked SRC** fits have 21 shape, 3  $W/Z$  normalization, and 30 and 19 SRC parameters, respectively. There are 1684 data points after cuts.

- ▶ Better overall quality of the SRC fits.

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- ▶ Especially better description of the (precise) high- $x$  JLAB data.

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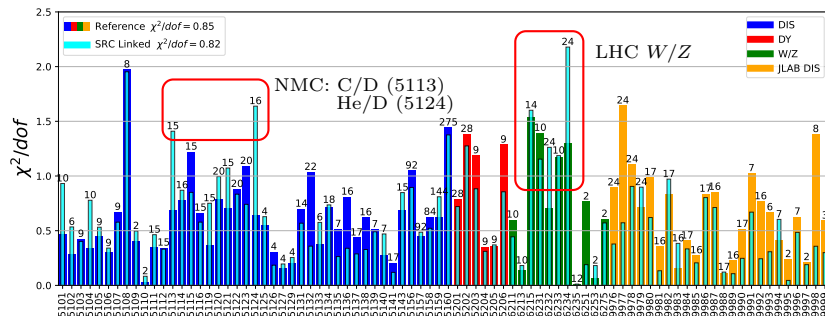
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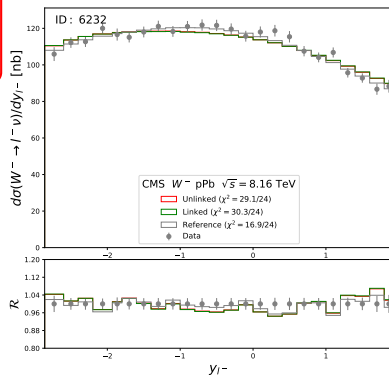
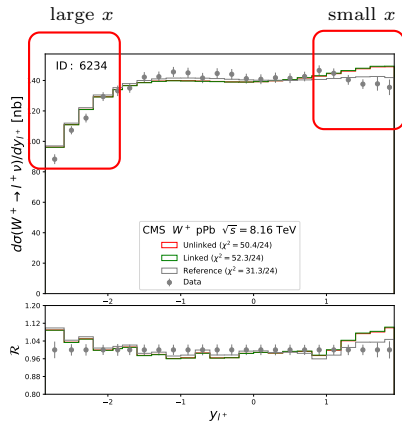
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- ▶ Worse description of the *W/Z* data from LHC - lowest available *x* values.

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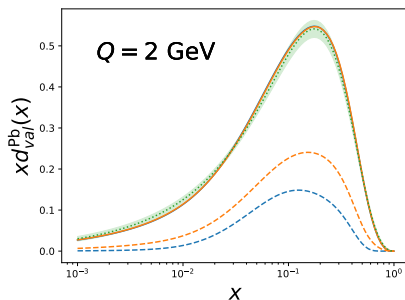
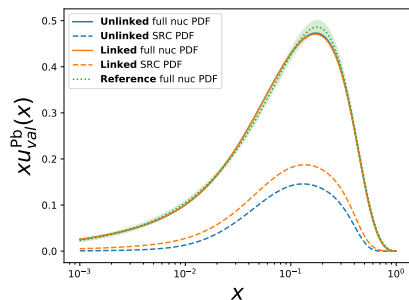


- ▶ Better overall quality of the SRC fits.
- ▶ Especially better description of the (precise) high- $x$  JLAB data.
- ▶ Worse description of the  $W/Z$  data from LHC - lowest available  $x$  values.
- ▶ For most of the experiments we observe decrease in  $\chi^2$ ,  
 exceptions: 5113 - NMC DIS for C/D, 5124 - NMC DIS for He/D, 6234 & 6232 - CMS  $W^\pm$  from Run II.

# Results



## Results: PDFs



- ▶ nPDFs obtained from SRC fits lie within the error bands of the **Reference** fit.
- ▶ The SRC components of the full nPDFs are in the range 20% to 30% – in agreement with the  $\{z_A, n_A\}$  values.

- ▶ The simple SRC-based picture of nPDFs leads to comparable or better data description than the traditional nPDF parameterization.
- ▶ The  $\{z_A, n_A\}$  parameters increase logarithmically in a uniform pattern with the nuclear  $A$  values from  $\sim 10\%$  to  $\sim 30\%$ .
- ▶ For the **Unlinked** fit the obtained values of  $\{z_A, n_A\}$  suggest approximately equal number of protons and neutrons in the SRC pairings. While this is consistent with other observations suggesting the SRC pairs are predominantly  $pn$  further analysis is required to confirm it.
- ▶ Even when the  $\{z_A, n_A\}$  parameters are constrained in the **Linked** fit, we obtain a very good fit to the data, yielding lower  $\chi^2$  than in the **Reference** fit. This should be tested with more data and/or can be used to further constrain the used parametrization.
- ▶ The improvement of the JLab data at the expense of the  $W/Z$  data suggests this approach may be more appropriate for larger  $x$  than for smaller  $x$ .
- ▶ It is notable that all the above results, obtained from purely data driven fits, seem to support the SRC-based description of nuclei.