

CS kernel determination from MC (PB-TMDs)

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A. Bermúdez Martínez, A. Vladimirov



Parton Branching (PB) method

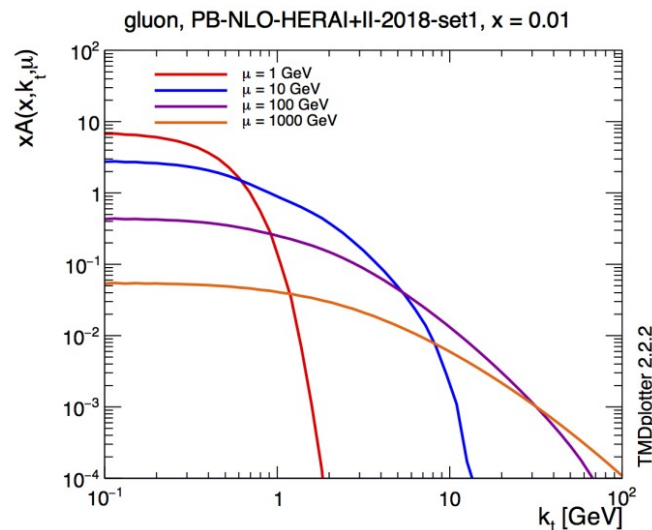
- Evolution of TMDs (and collinear PDFs)
- Resummation of soft gluons at LL and NLL
- Solution valid at LO, NLO and NNLO
- Determination of TMDs from the fully exclusive solution
- **Backward evolution fully determines the TMD shower**

FH et al. [[PLB 772 \(2017\) 446–451](#)]

FH et al. [[JHEP 2018, 70 \(2018\)](#)]

ABM et al. [[PRD 99, 074008 \(2019\)](#)]

→ nation of TMDs from
consistently treats perturbative and non-perturbative
transverse momentum effects



Implemented in the **CASCADE**
generator

[Eur. Phys. J. C 81 \(2021\) 425](#)

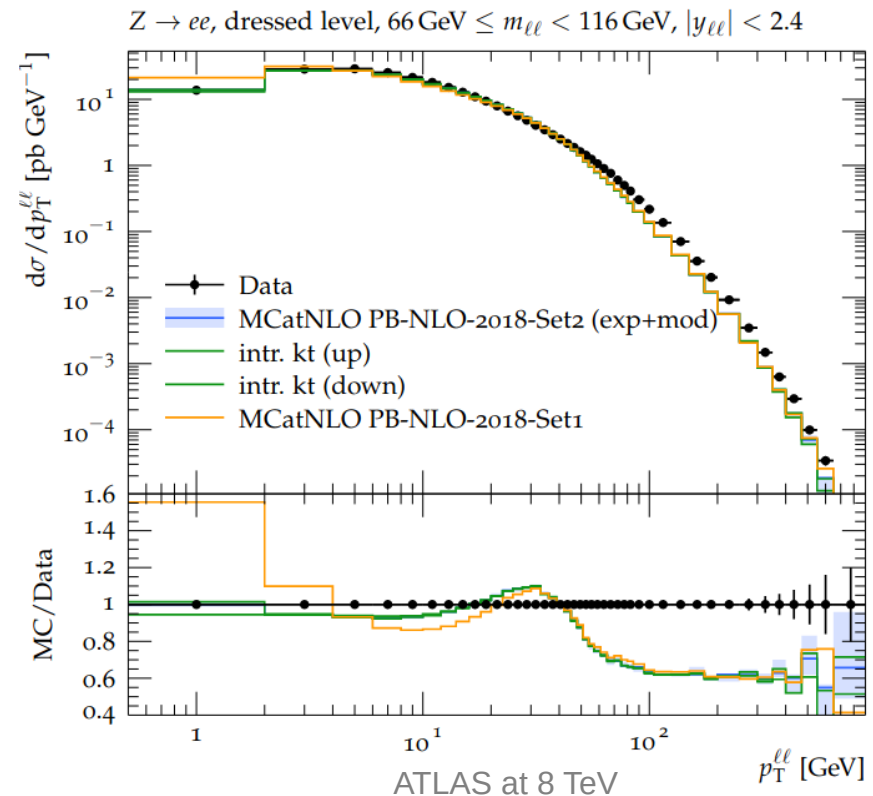
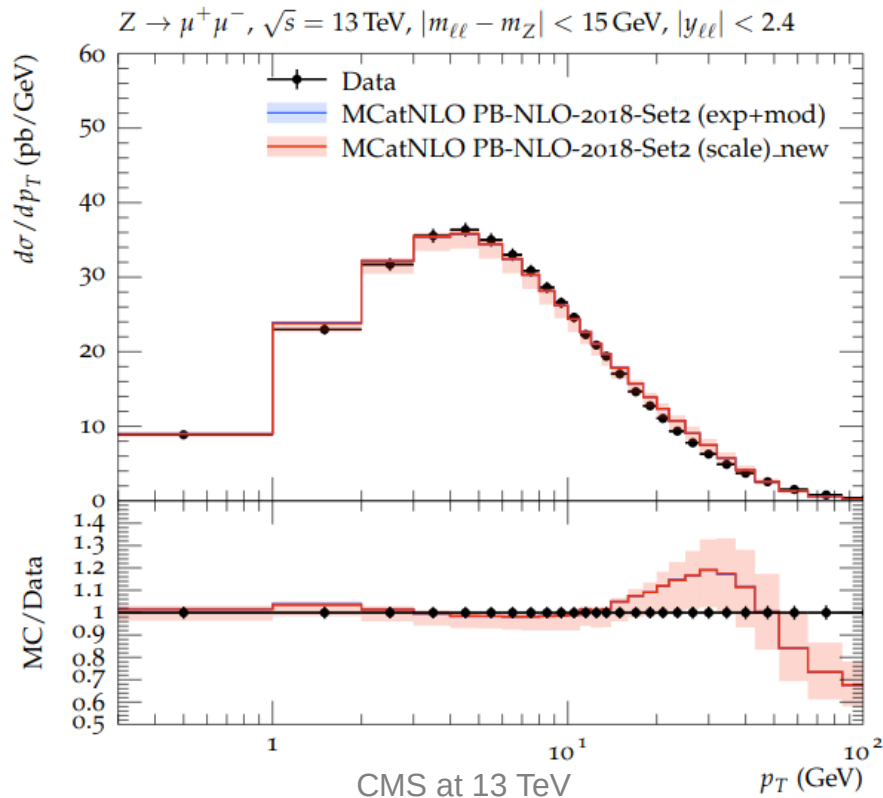
PB-TMDs in high mass DY production

DY pt spectrum

- Combined with MC@NLO
- **Excellent description of DY pT spectrum**
- **Non-perturbative TMD effects not significant at high pT**
- **Missing contributions at high pT**

ABM et al. [[PRD 100, 074027 \(2019\)](#)]

ABM et al. [[EPJC 80, 598 \(2020\)](#)]



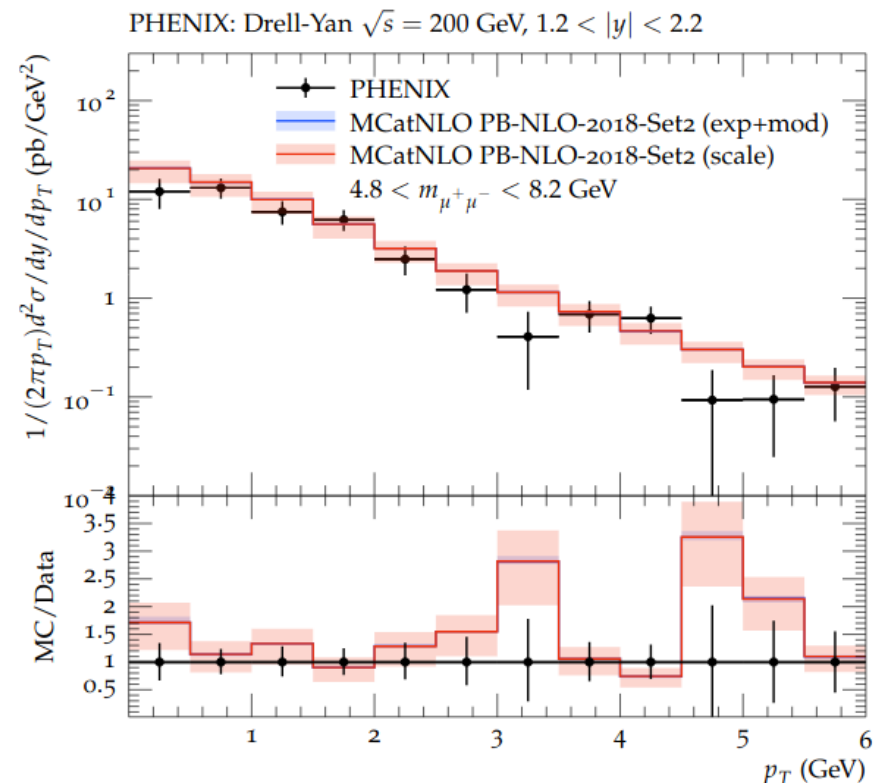
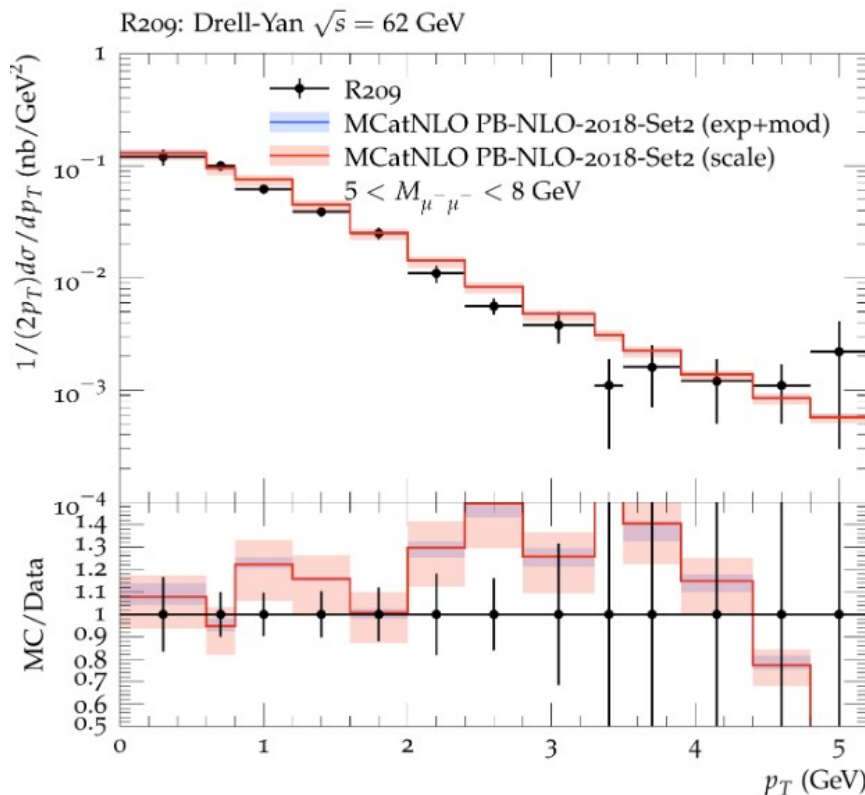
PB-TMDs to low mass DY production

DY pt spectrum

- Combined with MC@NLO
- Excellent description of DY pT spectrum
- **First simultaneous description of both low and high-mass DY pT spectrum**
- **No more low pT crisis** Bacchetta et al. [PRD 100 (2019) 014018]; ABM et al. [EPJC 80, 598 (2020)]

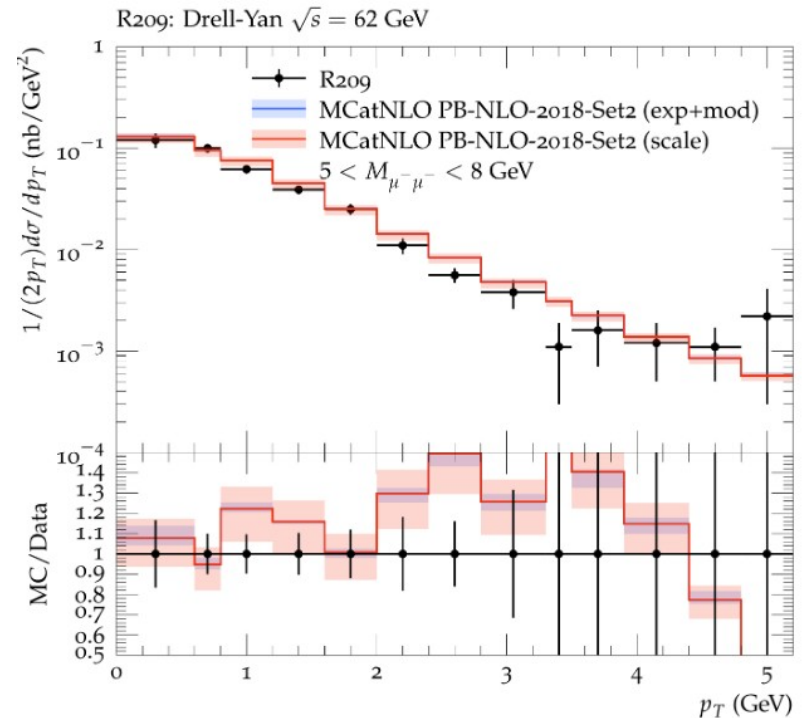
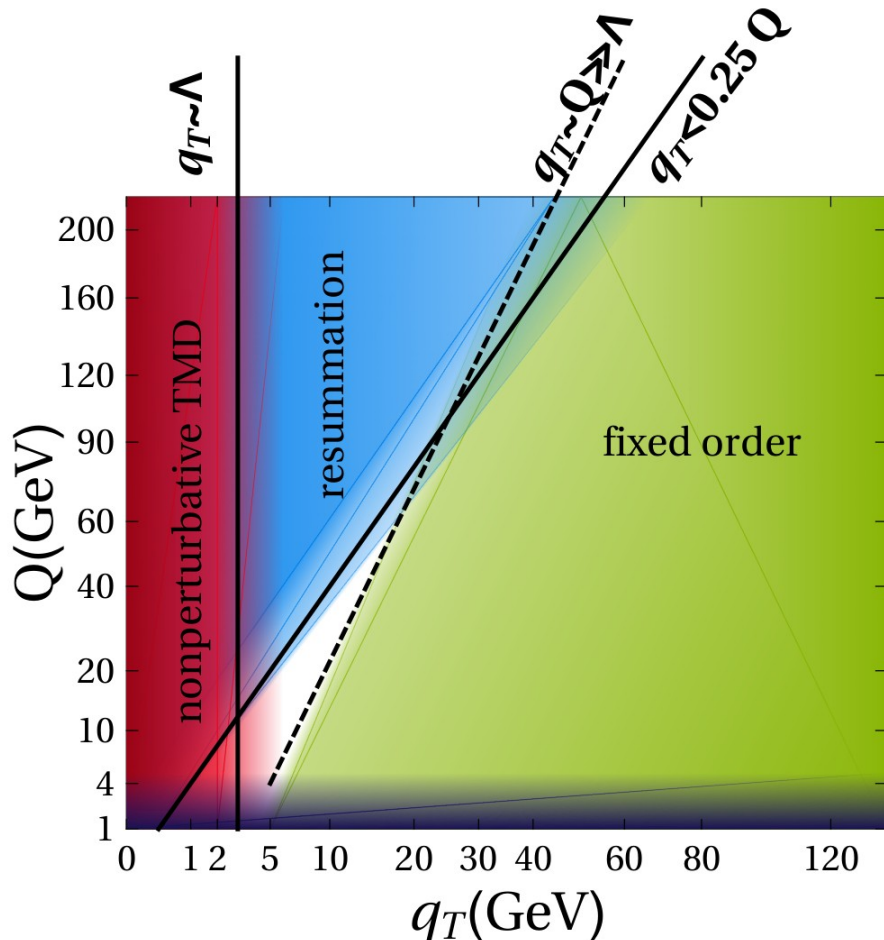
ABM et al. [PRD 100, 074027 (2019)]

ABM et al. [EPJC 80, 598 (2020)]



PB-TMDs to DY production

- Description of the p_T spectrum over a wide DY mass range
- **Valuable non-perturbative information**



Glance at TMD factorization

At small kt :

$$\frac{d\sigma}{dQ^2 dy dq_T^2} = \frac{2\pi}{9} \frac{\alpha_{\text{em}}^2(Q)}{sQ^2} W_{f_1 f_2}(x_1, x_2, Q, q_T),$$

$$W_{f_1 f_2} = |C_V(Q)|^2 \int_0^\infty db b J_0(bq_T) \sum_q e_q^2 f_{q_1}(x_1, b; Q, Q^2) f_{q_2}(x_2, b; Q, Q^2).$$

Notice the dependence on two scales of the parton distributions

Taking the Hankel (Fourier) transform of the xsec:

$$\Sigma(y, Q; b) = \frac{2\pi}{9} \frac{\alpha_{\text{em}}^2(Q)}{sQ^2} |C_V(Q)|^2 \sum_q e_q^2 f_{q_1}(x_1, b; Q, Q^2) f_{q_2}(x_2, b; Q, Q^2).$$

Glance at TMD factorization

Evolving the parton densities to fixed scales:

$$\Sigma(y, Q; b) = \frac{2\pi}{9} \frac{\alpha_{\text{em}}^2(Q)}{sQ^2} |C_V(Q)|^2 \boxed{e^{2\Delta(b; Q \rightarrow (\mu_0, \zeta_0))}} \sum_q e_q^2 f_{q_1}(x_1, b; \mu_0, \zeta_0) f_{q_2}(x_2, b; \mu_0, \zeta_0),$$

Notice the following points:

- The Q dependence of the parton densities transferred to a **residual exponential factor**
- Ratios of the xsec for different Q s do not depend on the collinear parton densities
- **Non-perturbative component** of those ratios would then live only in the exponential

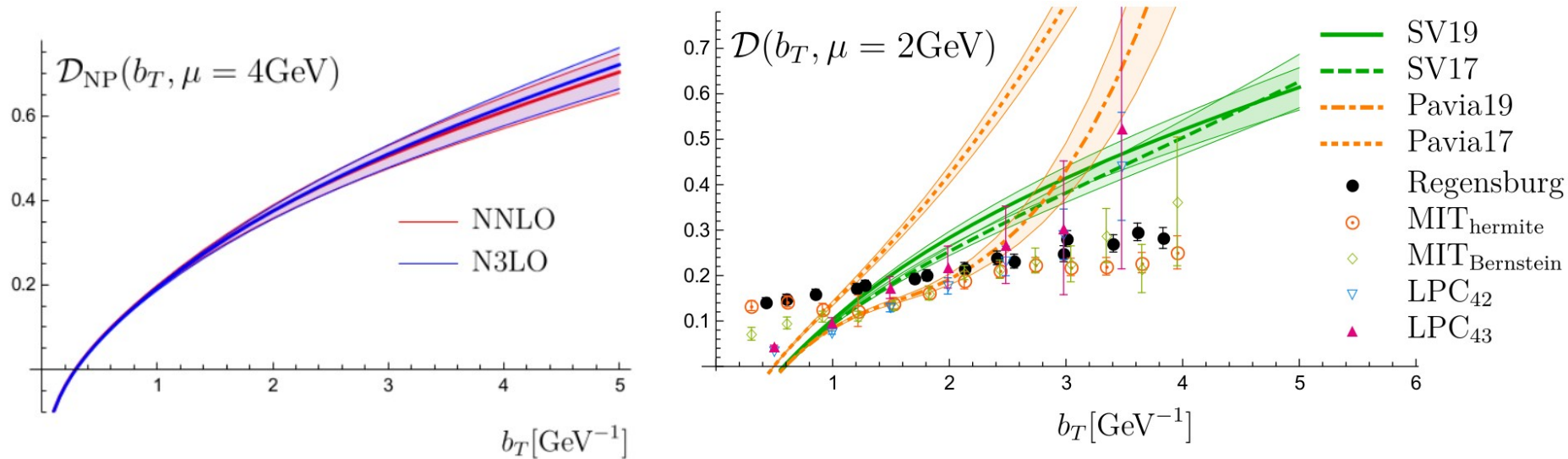
$$\Delta(b; Q \rightarrow (\mu_0, \zeta_0)) = \int_P \left(\gamma_F(\mu, \zeta) \frac{d\mu}{\mu} - \boxed{\mathcal{D}(b, \mu)} \frac{d\zeta}{\zeta} \right),$$

 **Collins-Soper kernel**

CS kernel

- **Self-contained object**
- Perturbative sensitivity at small b
- **Non-perturbative sensitivity at small b**
- Existing lattice calculations

A. Vladimirov, [[Phys. Rev. Lett. 125, 192002](#)]



Let's dive into the new method, applied in MC (PB-TMDs) and extendable to real data!!

CS kernel from Parton Branching

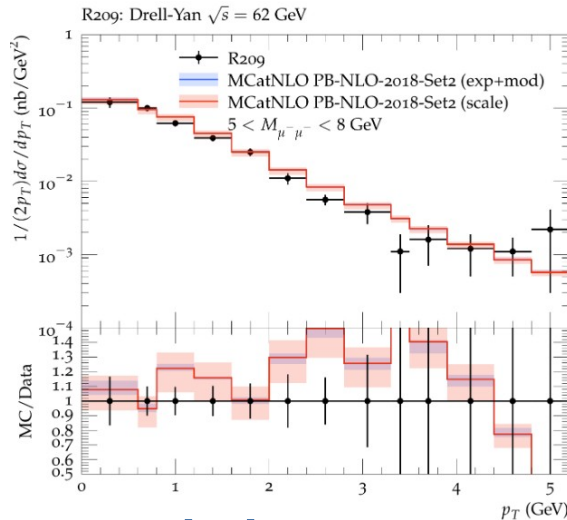
Main formula

$$\frac{\Sigma_1}{\Sigma_2} = \frac{s_2}{s_1} \frac{Q_2^2}{Q_1^2} Z(Q_1, Q_2) e^{2\Delta(b; Q_1 \rightarrow Q_2)}$$

Simulation/
Measurement

Analytically

Get CS kernel
from here



But in b space!

CS determination process

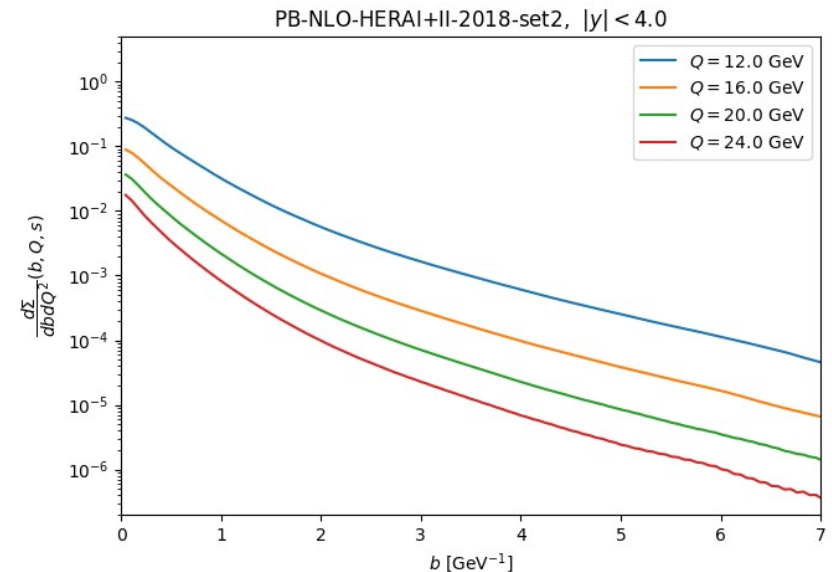
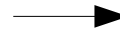
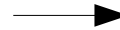
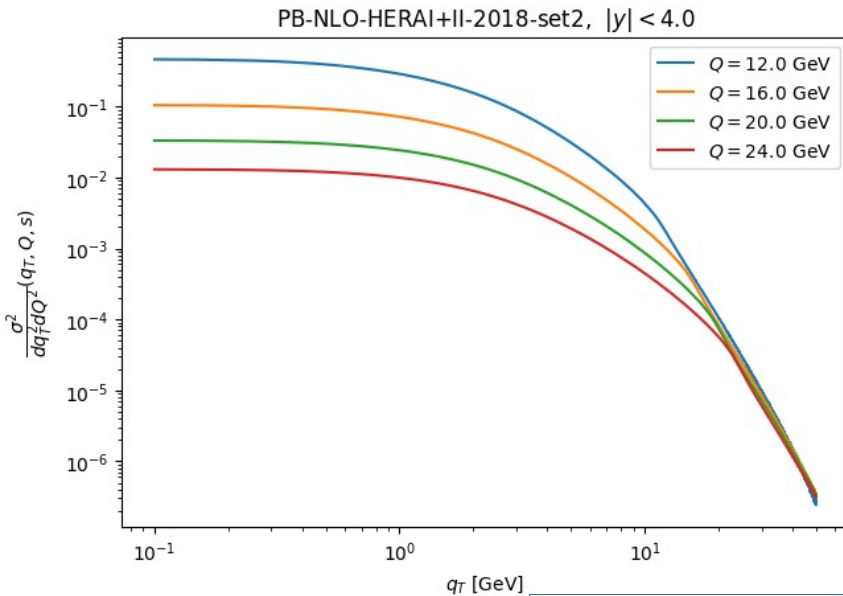
- DY p_T in narrow bins of $Q=\{Q_1, Q_2\}$
- Choose $\{s, y\}$ such that x 's coincide
- Very fine binning in p_T

Get the Fourier transform

CS kernel from Parton Branching

Hankel/Fourier transform of the xsec

$$\Sigma(b, Q, y, s) = \int_0^\infty dq_T q_T J_0(bq_T) \frac{d\sigma}{dQ^2 dy d^2\mathbf{q}_T^2}$$

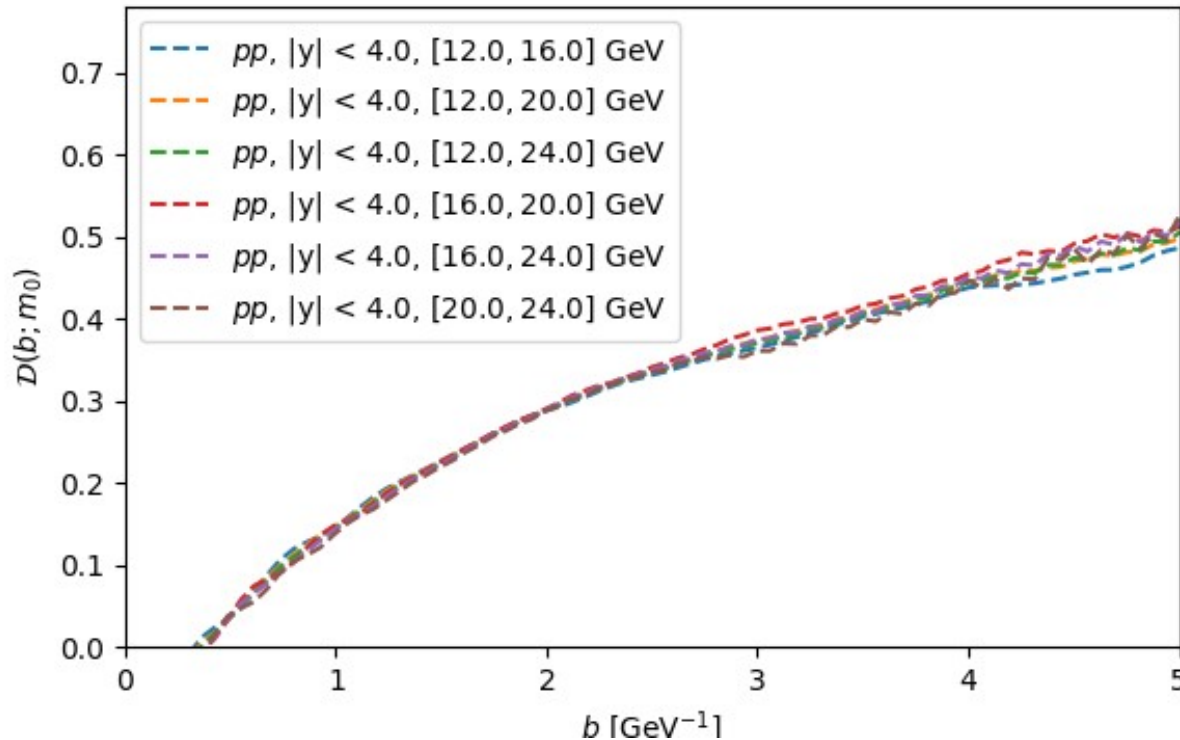


CS determination process

- DY pT in narrow bins of $Q=\{Q1, Q2\}$
- Choose $\{s, y\}$ such that x's coincide
- Very fine binning in pT

CS kernel from Parton Branching

Determine CS kernel



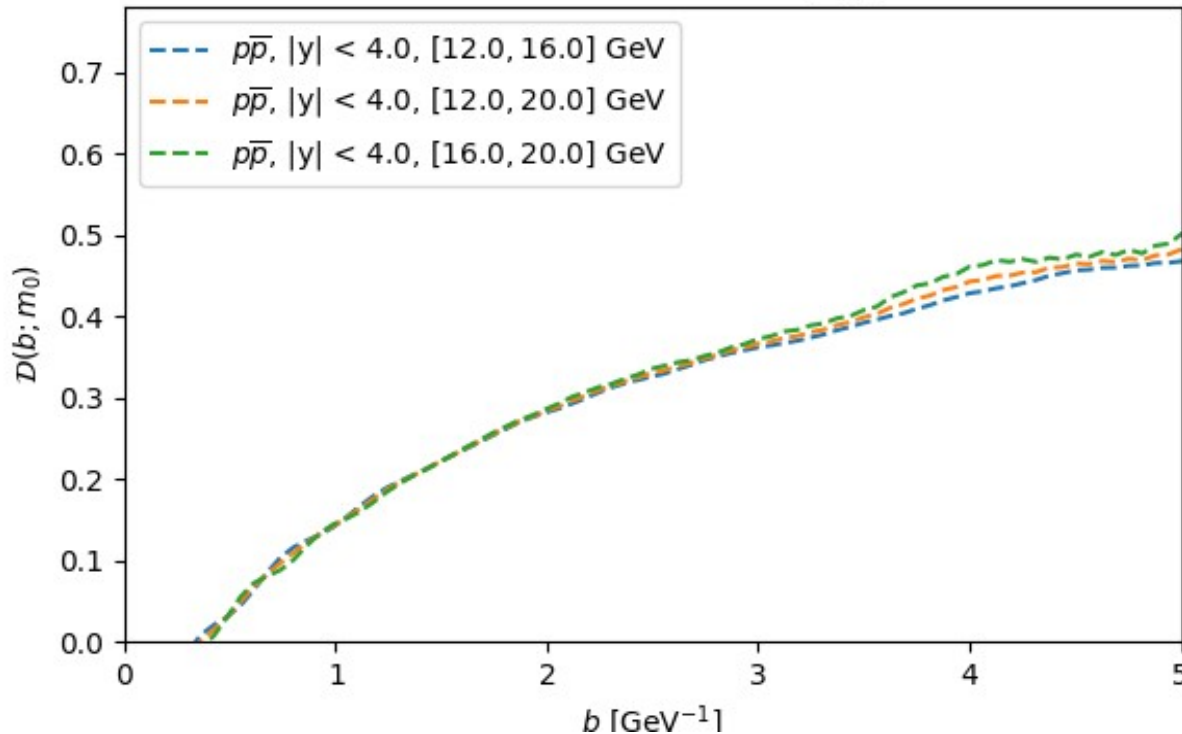
Ideally

- Different Q 's must coincide
- Different processes must coincide
- Should not depend on y range
- Should not depend on collinear PDF input

CS kernel from Parton Branching

Determine CS kernel

PB-NLO-HERAI+II-2018-set2, $|y| < 4.0$



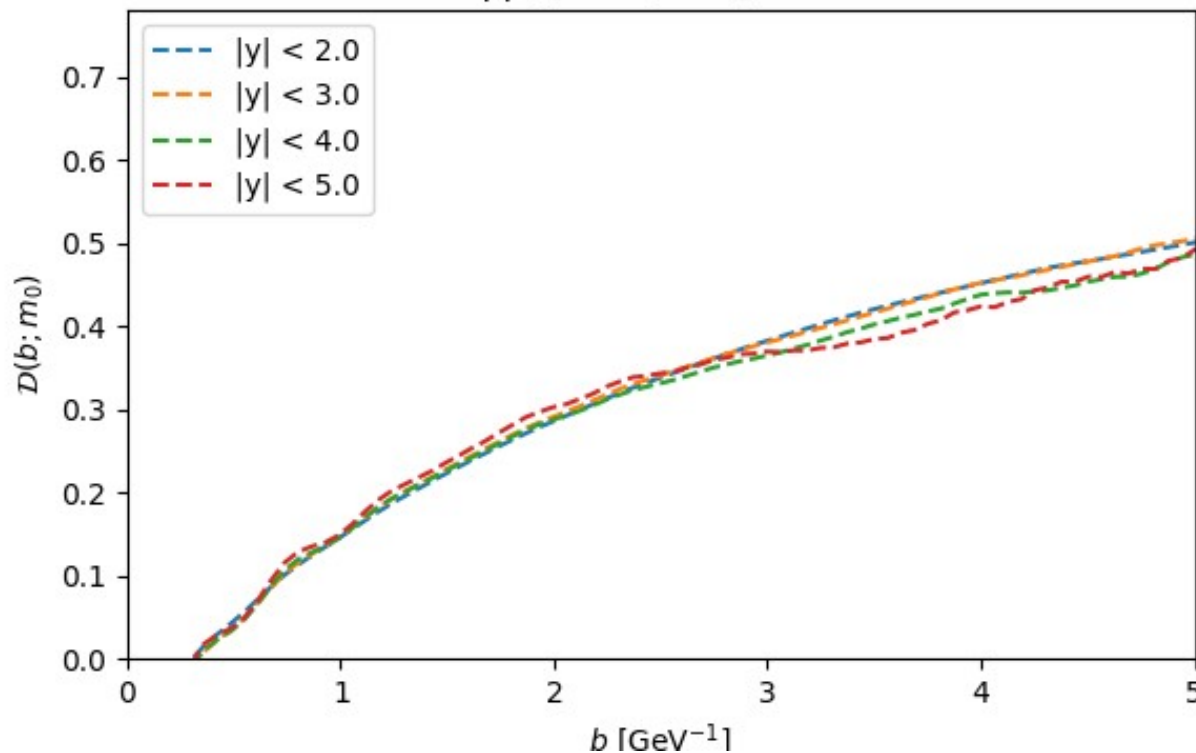
Ideally

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CS kernel from Parton Branching

Determine CS kernel

$pp, [12.0, 16.0]$ GeV



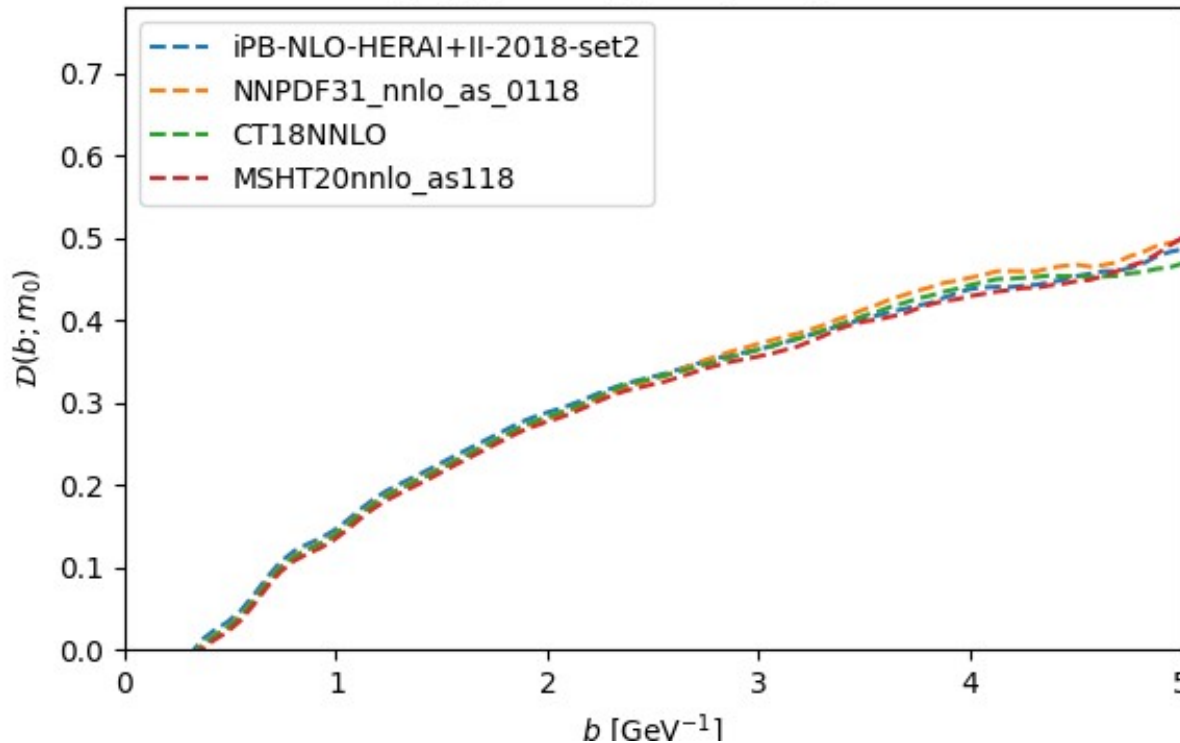
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CS kernel from Parton Branching

Determine CS kernel

$pp, |y| < 4.0, [12.0, 16.0] \text{ GeV}$



Ideally

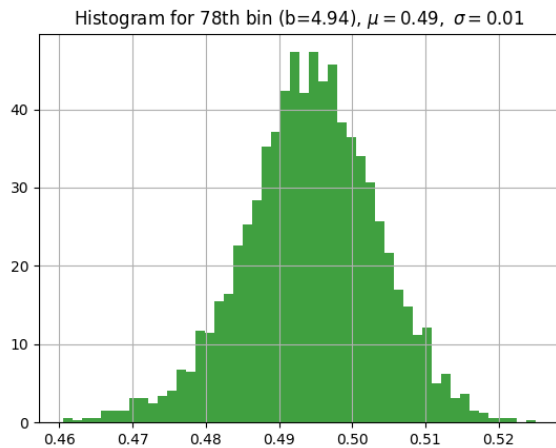
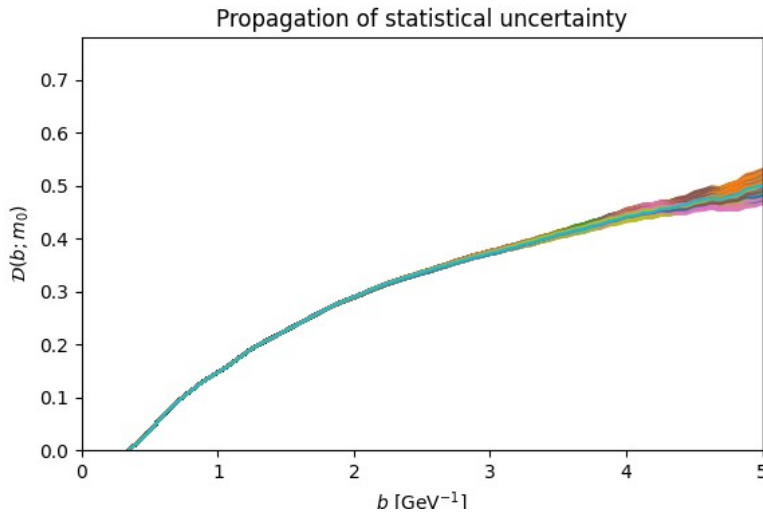
- Different Q's must coincide
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PB passed!!

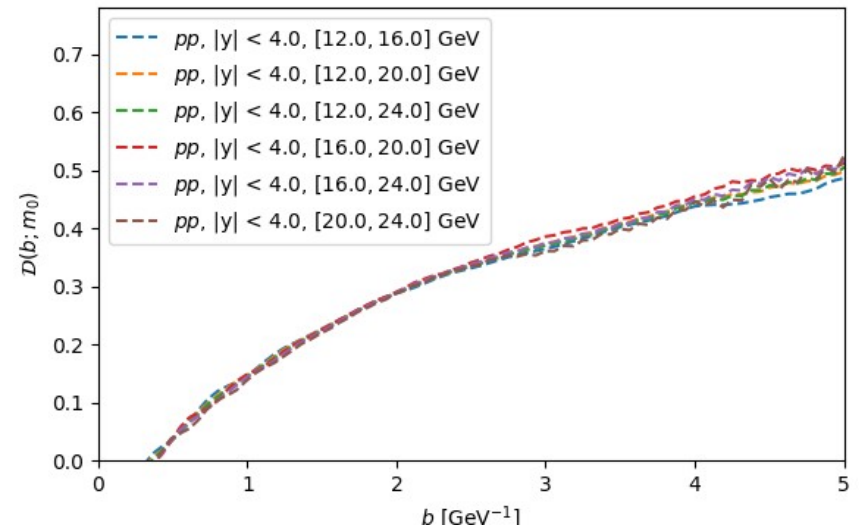
CS kernel from Parton Branching

Uncertainty sources

Statistical propagation

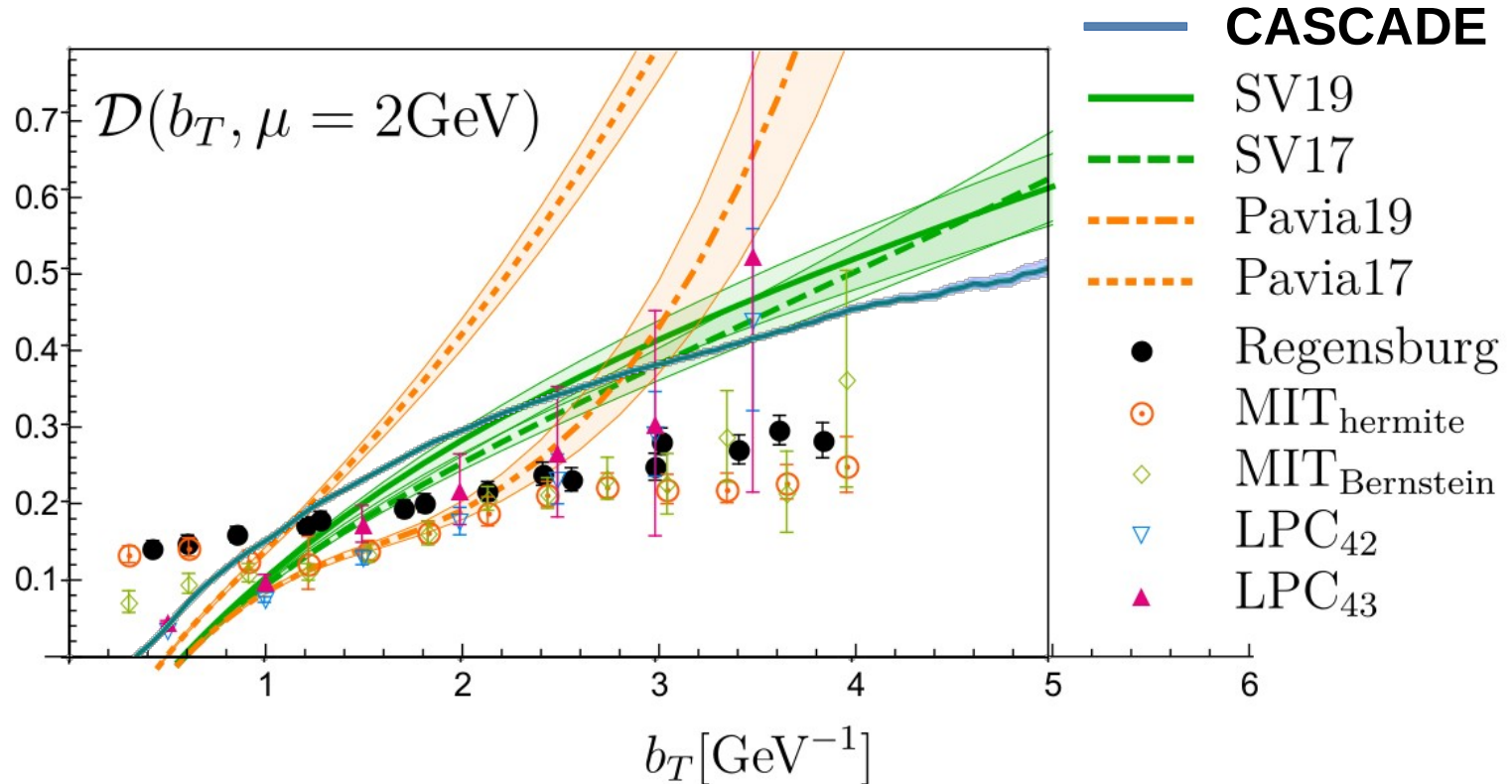


Scale variations



CS kernel from Parton Branching

Determine CS kernel



- There is a general agreement in shape between MC, Lattice, SV19 at large b
- MC shape consistent with perturbative calculation at low b

Conclusions

- CASCADE generator (PB-TMDs) consistent with TMD factorization
- Method to determine CS kernel from MC
- First determination of CS kernel from MC
- MC shape consistent with perturbative calculations at small b
- MC shape consistent with lattice calculations at large b
- Prove of concept → **to be used with real measurements**