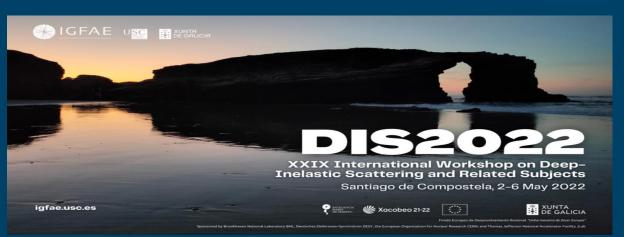
PB TMD fits at NLO with dynamical resolution scale

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- Recap of Parton Branching method
- Fixed and Dynamical soft-gluon resolution scale z_M
- Fits with fixed z_M at NLO
- Fits with dynamical z_M at NLO

Merged talk:

Z+b jet production in 4FL and 5FL

TMD evolution in the PB formalism:

$$\begin{split} \widetilde{A_a}\big(x,k_\perp,\mu^2\big) = \\ \widetilde{A_a}\big(x,k_\perp,\mu_0^2\big) \, \Delta_a\big(\mu^2\big) + \sum_b \int \frac{d^2\mu_\perp'}{\pi{\mu'}^2} \, \Theta\big(\mu^2 - {\mu'}^2\big) \, \Theta\big({\mu'}^2 - \mu_0^2\big) \\ \times \frac{\Delta_a\big(\mu^2\big)}{\Delta_a\big({\mu'}^2\big)} \int_x^{\textbf{Z}_{\textbf{M}}} dz \, P_{ab}^R(\,z,\alpha_s(\textbf{\textit{q}}_\perp)) \\ \times \, \widetilde{A_b}\left(\frac{x}{z},k_\perp + (1-z)\mu',{\mu'}^2\right) \, dz \, P_{ab}^R(\,z,\alpha_s(\textbf{\textit{q}}_\perp)) \end{split}$$

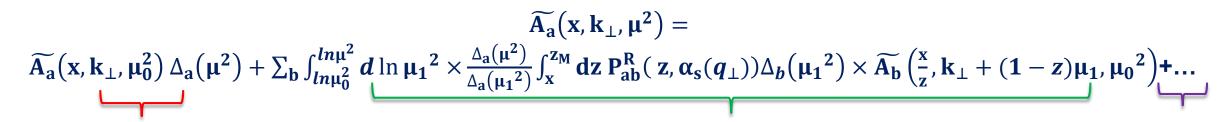
- $ightharpoonup z_M$: Resolvable branching : $z < z_M$ Non-resolvable branching : $z > z_M$
- **Splitting functions:** $P_{ab}^{R}(z)$: The real emission parts of the DGLAP splitting function: Probability that a branching will happen
- $x_bp^+, k_{t,b}$ b> Sudakov form factor: $\Delta_a = \exp(-\int_{\ln \mu_0^2}^{\ln \mu^2} d(\ln \mu'^2) \sum_b \int_0^{z_M} dz \, z \, P_{ba}^R(\alpha_s, z))$

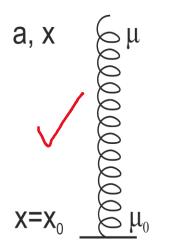
The probability of an evolution without any resolvable branching

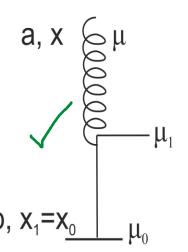
 $x_a p^+, k_{t,a}$ a $z = x_a/x_b$ c $q_{t,c}
ightarrow \mu$ μ is evolution scale

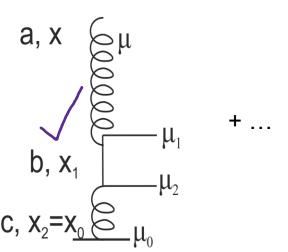
At every step kinematics can be calculated!

Iterative form of the PB evolution equation:









Solvable by MC iterative technique:

- generated μ_1^2 : if $\mu_1^2 > \mu^2$ stop, otherwise splitting,
- generated the next scale μ_2^2 : if $\mu_2^2 > \mu^2$ stop, otherwise splitting,
- ...

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Color coherence phenomena:

Angular ordering of the soft gluon emissions

$$\begin{aligned} \Theta_{i+1} &> \Theta_i \\ |q_{\perp,i}| &= (1-z_i) |E_i| \sin \Theta_i \end{aligned}$$

Associating " $|E_i| \sin \Theta_i$ " with μ'

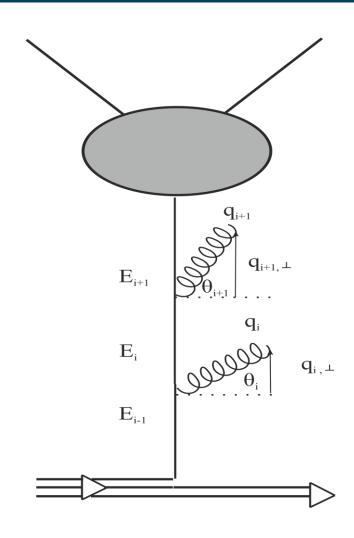
$$q_{\perp,i}^2 = (1 - z_i)^2 \, \mu_i^{\prime 2}$$

• The **argument of** α_s should be q_{\perp}^2

$$\alpha_s(q_\perp^2) = \alpha_s((1-z)^2\mu'^2)$$

• resolvable & non-resolvable \rightarrow condition on $\min q_{\perp,i}^2 \rightarrow z_M$

$$z_{M} = 1 - \left(\frac{q_{0}}{\mu'}\right)$$



Fixed and dynamical resolution scale

- \triangleright Fixed z_M :
- μ independent

$$\mathbf{z_M} = \mathbf{1} - \boldsymbol{\epsilon}$$

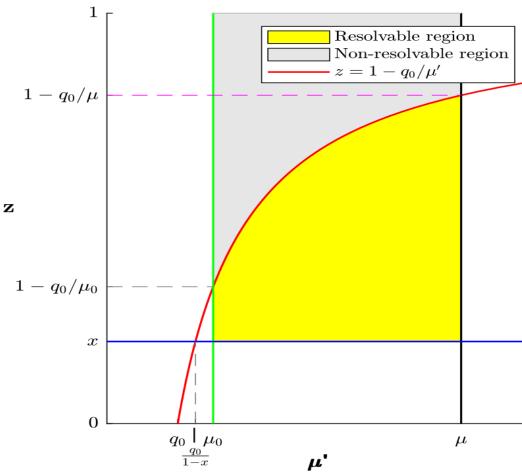
where ϵ is small: 10^{-3} , 10^{-4} , 10^{-5} ,...

Dynamical Resolution scale in Angular Ordering:

$$\mathbf{z}_{\mathbf{M}} = \mathbf{1} - \left(\frac{\mathbf{q}_{\mathbf{0}}}{\mathbf{\mu}'}\right)$$

where q_0 is smallest emitted transverse momentum for resolvable partons

- Sudakov form factor Δ_a : non- resolvable region
- Splitting functions P_{ab}^R : resolvable region

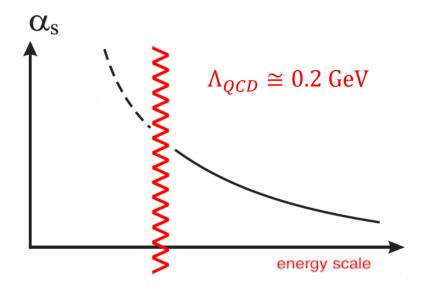


Dynamical resolution scale

The Condition on q₀ of

$$z_M = 1 - \left(\frac{q_0}{\mu'}\right)$$

- Scale of strong coupling: $\alpha_s(q_\perp^2) = \alpha_s((1-z)^2\mu'^2)$
- Lowest scale in α_s corresponds to minimal q_{\perp}
- $q_{\perp,min} = q_0 \& q_0 > \Lambda_{QCD} =>$ we stay in the weak coupling region!



PB TMD fits at NLO with fixed zmax

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The Past PB TMD fits at NLO calculation using angular ordering : fixed z_M

"NLO DIS Matrix Element (ME) and NLO evolution kernel"

Associating the evolution scale with some physical interpretation:

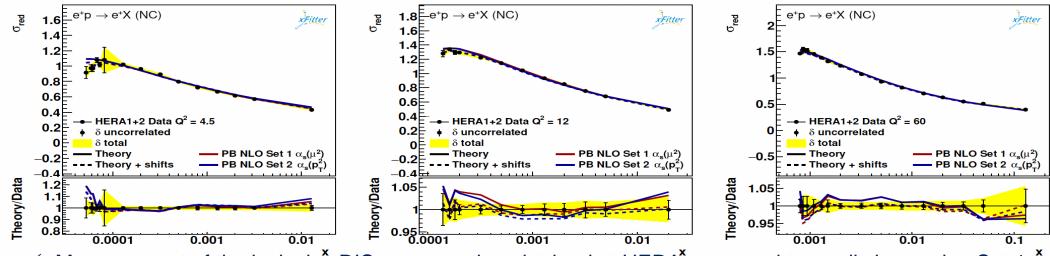
Two scenario
$$\begin{cases} \bullet & \text{Set 1} & \longrightarrow & \alpha_S(\mu'^2) \\ \bullet & \text{Set 2} & \longrightarrow & \alpha_S(q_\perp^2) = \alpha_S((1-z)^2\mu'^2) \end{cases}$$

■ The resulting TMD parton densities, PB-NLO-2018-set1 and PB-NLO-2018-set2 are available in TMDLIB2: The European Physical Journal C 81.8 (2021): 1-10

Data set: HERA 1+ 2 inclusive DIS data

PB NLO Set1 $lpha_{ m s}(\mu_i^2)$			
	χ^2	d.o.f	$\chi^2/{\rm d.o.f}$
$\mu_0^2 = 1.9 \text{GeV}^2$	1363.37	1131	1.21
PB NLO Set 2 $lpha_{ m s}(q_{ti}^2)$			
	χ^2	d.o.f	$\chi^2/{ m d.o.f}$
$\mu_0^2 = 1.4 \mathrm{GeV}^2$	1369.80	1131	1.21

.Phys Rev. D 99, 074008 (2019),



✓ Measurement of the inclusive DIS cross section obtained at HERA*compared to predictions using Set 1 and Set 2

PB TMD fits at NLO with dynamical zmax

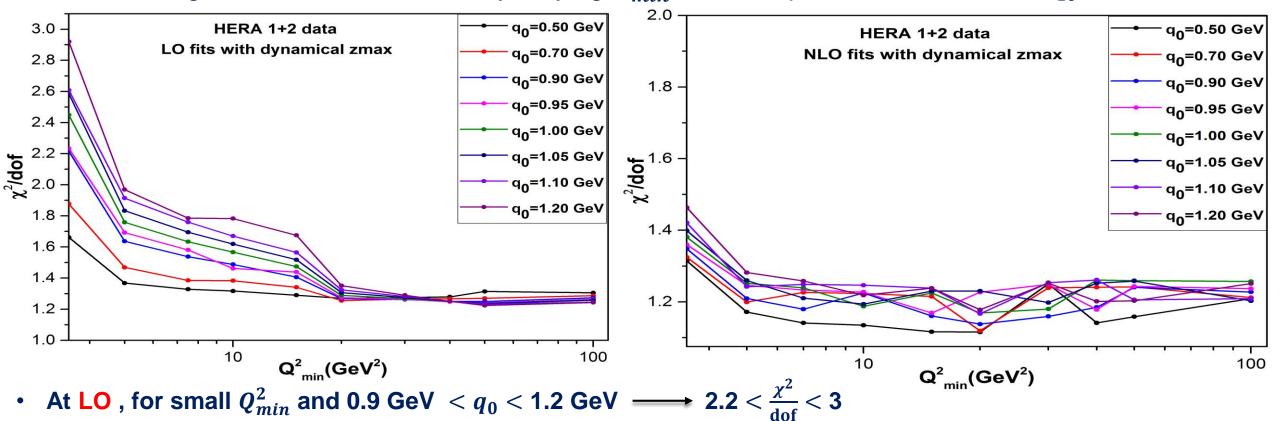
New study

From fixed resolution scale to dynamical resolution scale

AT NLO, for small Q_{min}^2 and all values of q_0 , we have better fits with good

New fits with dynamical zmax at LO and NLO with HERA 1 + 2 Data set: Using Aritter arXiv:1709.01151v1

 \checkmark Performing different fits, each time by varying Q_{min}^2 and on top of that with different q_0 values

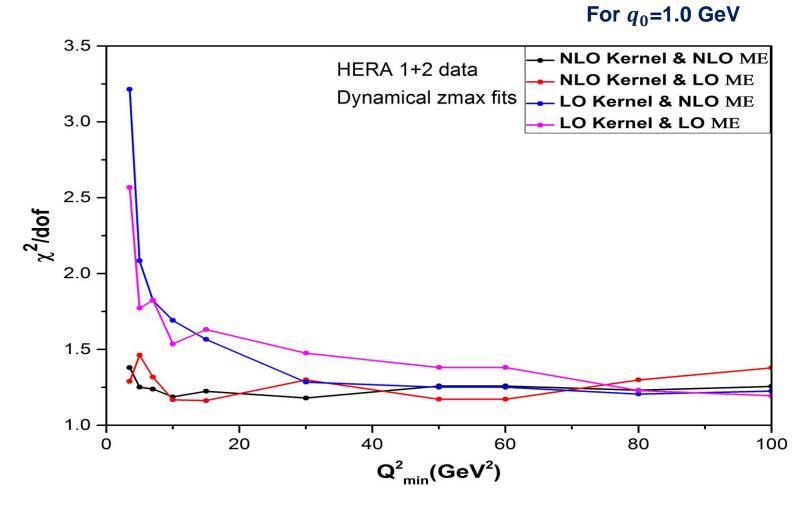


Does the difference between LO and NLO come from the kernels? or ME?!...

4 states for this purpose:

- 1. Fitting with NLO kernel & NLO ME
- 2. Fitting with NLO kernel & LO ME
- 3. Fitting with LO kernel & LO ME
- 4. Fitting with LO kernel & NLO ME

The difference is dominated by the kernel not ME..!



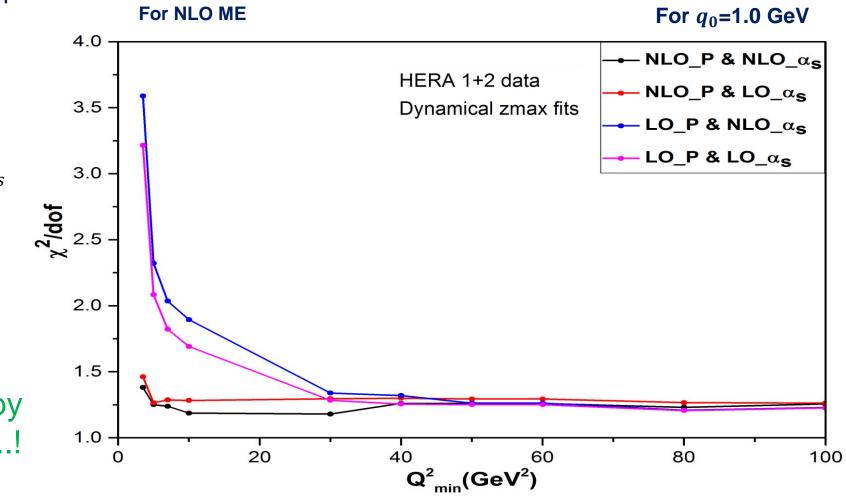
Which part of the kernel is responsible?

$$P_{ab}(z,\mu^2)$$
? or α_s ?

4 states for this purpose:

- 1. Fitting with NLO P_{ab} & NLO α_s
- 2. Fitting with NLO P_{ab} & LO α_s
- 3. Fitting with LO P_{ab} & LO α_s
- 4. Fitting with LO P_{ab} & NLO α_s

The difference is dominated by the splitting functions not α_s ..!



Which part of the splitting functions is responsible for the difference between LO and NLO?

> For high values of q_0 (e.g, [1.0 Gev, 1.2 Gev]) or low values of $z_M = 1 - \left(\frac{q_0}{\mu'}\right)$, LO and NLO have different behavior.

The first piece for checking is
$$\frac{1}{z}$$

- In the NLO, all the splitting functions have pieces with (1/z) term : $P_{ab}(z,\mu^2) \sim P_{qq}(1/z,\mu^2)$, $P_{qq}(1/z,\mu^2)$, $P_{qq}(1/z,\mu^2)$, $P_{qq}(1/z,\mu^2)$
- In the LO, just the splitting functions with "gluon" in the final state have (1/z) piece:

$$P_{gg}(z,\mu^2) = \frac{1}{1-z} + \frac{1}{z} - 2 + z(1-z),$$

$$P_{gq}(z,\mu^2) = \frac{1+(1-z)^2}{z}$$

• And the splitting functions with "quark" in the final state don't have (1/z) piece:

$$P_{qq}(z, \mu^2) = \frac{2}{1-z} - 1 - z,$$

$$P_{qq}(z, \mu^2) = z^2 + (1-z)^2$$

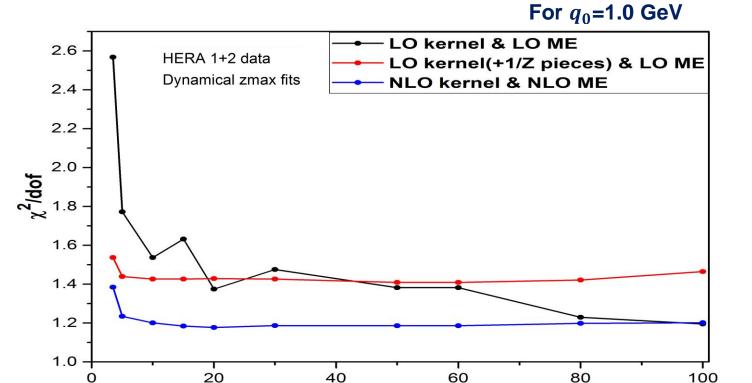
➤ Is the lack of (1/z) piece in LO splitting function with quark in the final state responsible for this difference?

Does the difference come from 1/z piece of NLO splitting function?

For better understanding: "We added to the LO splitting functions(P_{qq} , P_{qq}) the 1/z pieces of NLO"

✓
$$P_{qq}(z, \mu^2) = \frac{2}{1-z} - 1 - z + (\frac{1}{z})$$
 pieces of P_{qq} NLO

$$\checkmark P_{qg}(z,\mu^2) = z^2 + (1-z)^2 + (\frac{1}{z}) \text{ pieces of } P_{qg} \text{ NLO}$$



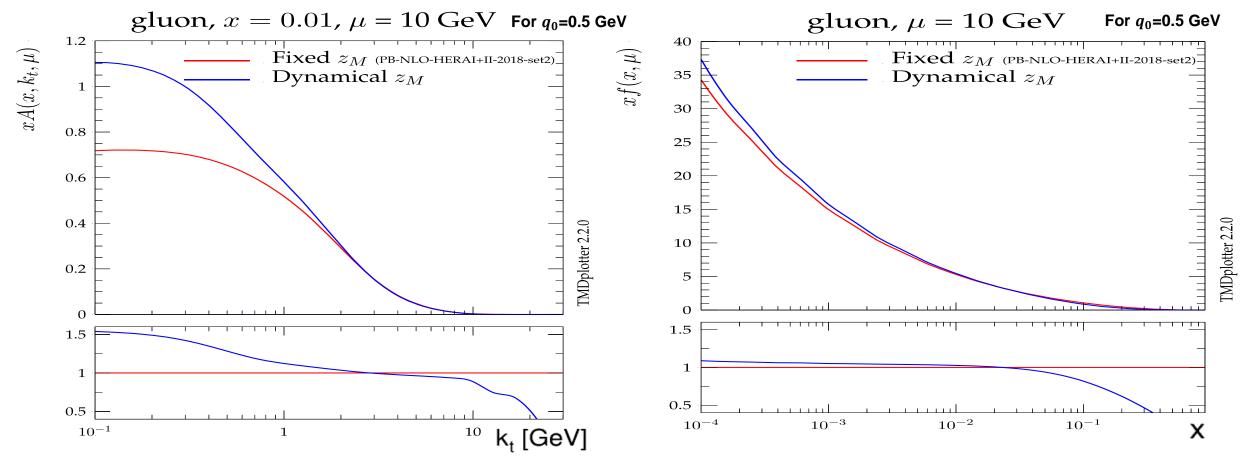
 $Q_{min}^2(GeV^2)$

- ✓ In NLO we have an extra (1/z) pieces in the quark channels compared with LO which is responsible for this difference!
- ✓ With this piece we are describing data well! Amount of $\frac{\chi^2}{\text{dof}}$ is reasonably good!

** For PB-TMD fit with dynamical zmax we obtain a reasonably good $\frac{\chi^2}{dof}$ at NLO! **

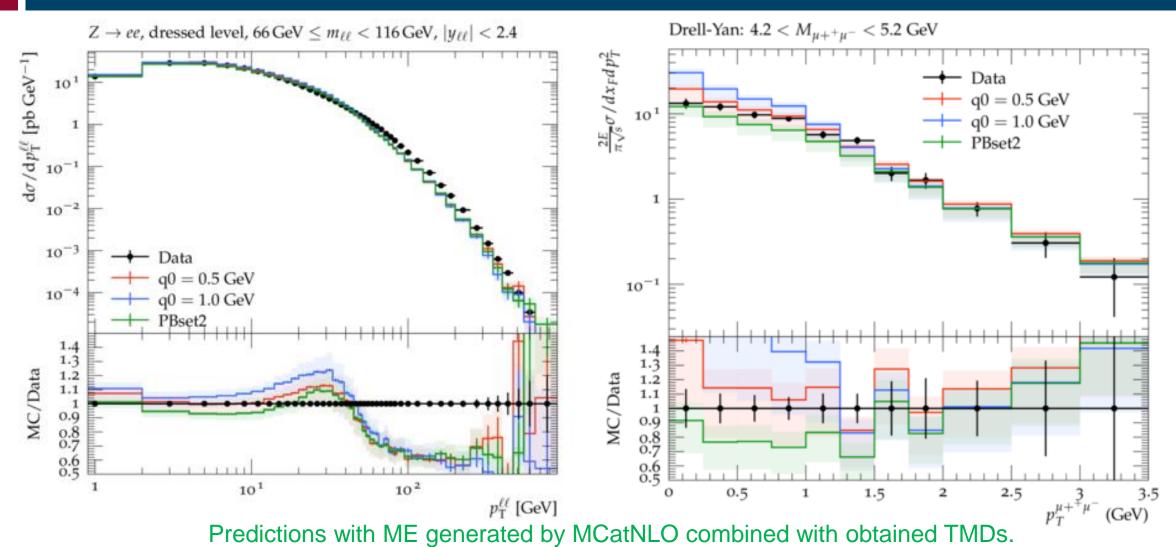
How does dynamical zmax affect the fitted TMD (iTMD)?

Set 2: fixed zmax & $\alpha_s(q_{\perp}^2) = \alpha_s((1-z)^2\mu'^2)$



 \checkmark The dynamical zmax fit implies an effect not only in the k_T dependence but also in the x dependence!

The predictions in dynamical zmax frame

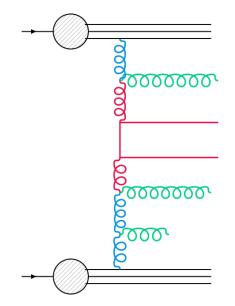


The merged talk: Z+b jet production in 4FL and 5FL

5 FLNS

- Full coupled evolution with all flavours & α(M_znf=5)=0.118
- HERAPDF parametrization form
- Using full HERAI+II inclusive DIS data
- chi2/dof=1.21

[Phys. Rev. D 99 (2019) no. 7, 074008]



4 FLNS

- The same functional form and data as 5FL-parameters re-fitted
- $M_b \to \infty \& \alpha (M_7^{nf=4}) = 0.1128$
- chi2/dof=1.25

[arXiv:2106.09791]

Matrix elements from MC@NLO (HERWIG6 subtraction)

- 5FLVNS : Z + one parton process
- 4FLVNS: Z + bb process

PDFs: TMDs (4FL & 5FL)

- 5FLVNS: b-quark is treated as a light quark
- 4FLVNS: no b-quark in the parton density

Parton shower following TMDs for intial state

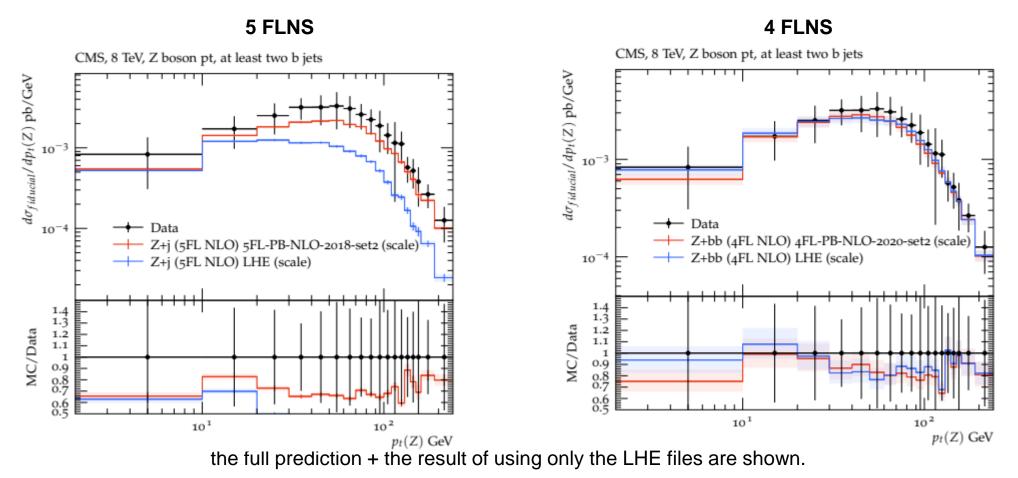
5FL & 4FL PB-TMDs included in the Cascade3

Eur. Phys. J. C 81 (2021) 425

$Z+b\overline{b}$ as a function of $p_t(Z)$

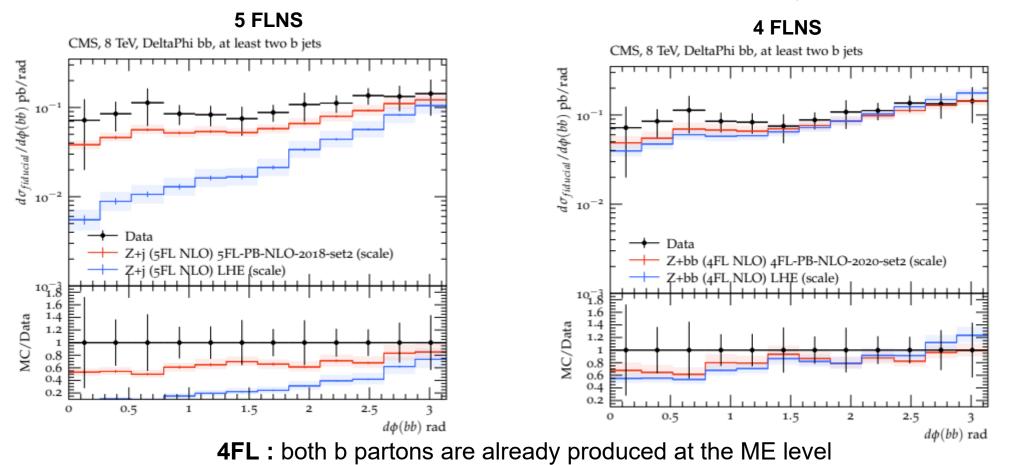
Differential cross section for $Z + b\bar{b}$ as a function of $p_t(Z)$ as measured by CMS collaboration

[Eur. Phys. J. C 77 (2017) 751]



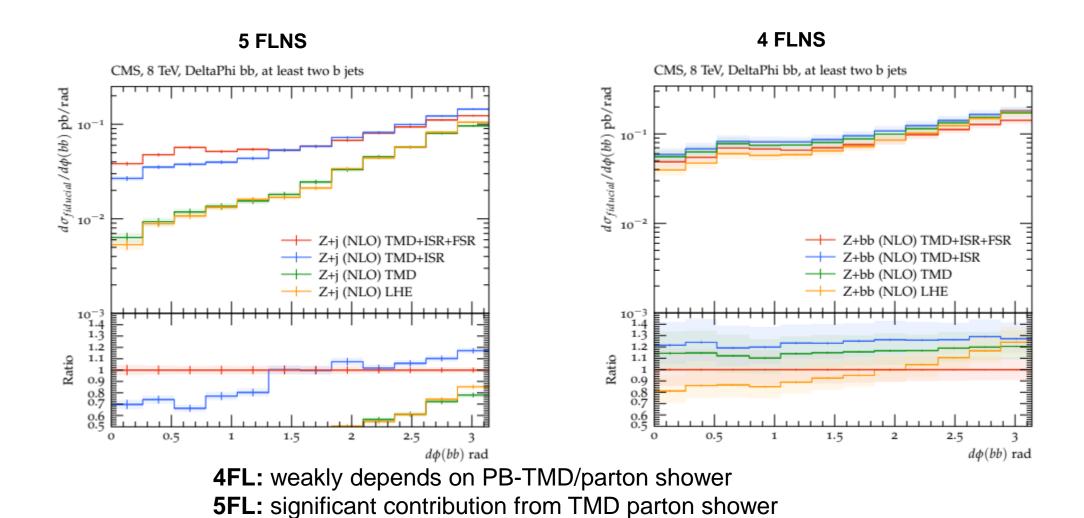
Z+bb as a function of $\Delta \phi$ (b \overline{b})

Differential cross section for Z +b \bar{b} as a function of $\Delta \phi$ (b \bar{b}) as measured by CMS collaboration



5FL: bb must be simulated in the parton shower.

Breakdown of the different contributions to quantify their roles



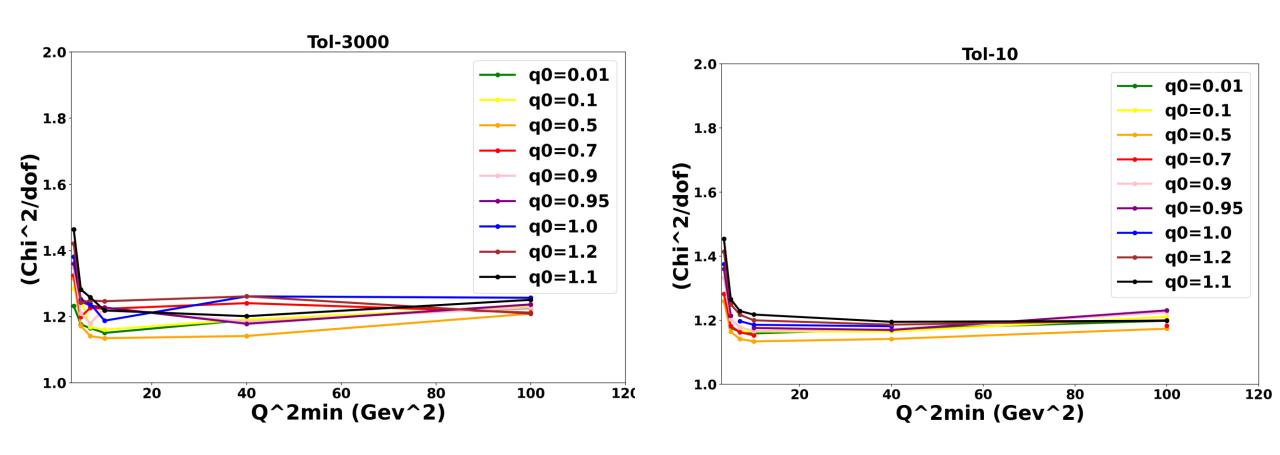
- > PB TMD fits at NLO with dynamical zmax for the first time!
- > For PB-TMD fit to HERA data with dynamical zmax, we obtain a reasonably good $\frac{\chi^2}{dof}$ at NLO!
- > The difference between LO and NLO fits is mostly due to (1/z) pieces in quark channel in NLO splitting functions!
- \succ The dynamical zmax impacts both the k_T dependence and the x dependence of the fitted parton distribution!
- > The next step: Using the PB TMD with dynamical zmax in phenomenology of LHC and lower energy colliders!
- ➤ The 4FL and 5FL PB-TMD distributions used to calculate Z + bb production
 - •Good agreement with measurements obtained by the CMS collaboration
 - •The evolution of the PB-TMD parton densities as well as in the PB-TMD parton shower is checked.

Thank you ...

BACK UP ...



Back up...



PB TMD fits at NLO with dynamical zmax:

