Phenomenological study of quarkonium suppression and the impact of the energy gap between singlets and octets

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Based on ...

 Escobedo and Ferreiro, Simple model to include initial-state and hot-medium effects in the computation of quarkonium nuclear modification factor, Phys. Rev. D 105 (2022) 1, 1.

Outline

- Introduction
- 2 Initial-effect model
- The initial temperature
- 4 Computation of R_{AA}
- Conclusions

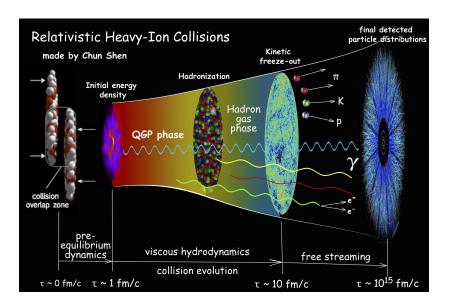
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- Quarkonium is a good probe of the QGP. Not only it is formed in the initial instants of the collision but also it was predicted by Matsui and Satz that quarkonium melts in the QGP. Therefore, by measuring quarkonium suppression we can infer the properties of the medium.
- Moreover, a heavy ion is not equivalent to an uncorrelated ensemble of protons. This can also modify the probability of quarkonium formation. Initial state effects.

A heavy ion collision



R_{AA} , the nuclear modification factor

In the Glauber model, the nucleus is seen as a collection of uncorrelated protons moving eikonally on the longitudinal direction.

$$R_{AA} = \frac{N_{HQ}^{AA}}{N_{col}N_{HQ}}$$

where

- N_{HQ}^{AA} is how much quarkonium is produced in a heavy ion collision.
- N_{col} is the number of nucleon-nucleon collisions in that heavy ion collision.
- \bullet N_{HQ} is how much quarkonium produce in a proton-proton collision.

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R_{AA} is not one if

- There are medium effects.
- There are initial effects, i.e. if a heavy ion collision is not equivalent to N_{col} proton-proton collisions.

Hot medium effects

Phenomena that modify quarkonium population in a QGP

- Screening of chromoelectric fields at large distances. Inhibits quarkonium formation if quarkonium's size is larger than screening length.
- Medium induced decay width.
- Recombination. A heavy quark-antiquark pair can meet inside of the medium and form a new bound state. Important for charmonium but sub-leading effect for bottomonium.

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Survival probability

If we ignore recombination, we can encode medium effects in a survival probability S.

• Given S as a function of the initial temperature. Can we compute R_{AA} ?

Motivation

To develop a framework to easily compute R_{AA} in the cases in which the survival probability is given by a simple analytical formula.

Questions we wish to answer:

- What is the initial temperature as a function of x_{\perp} ?
- How does the quarkonium production probability depends on x_{\perp} ?
- Compute R_{AA} for a given S.

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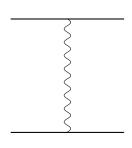
 The Glauber model can be understood as the assumption that nucleons only interact by Pomeron exchange.

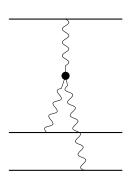
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- In the parton language, we can relate the Pomeron with the exchange of a pair of gluons.
- Corrections beyond the Glauber model can be implemented by including some non-trivial interaction between Pomerons.
- In our case, we use a model that includes a triple Pomeron vertex.

Pomeron exchanges





Some definitions

The thickness function is the density of nucleons in the transverse plane

$$T_A(x,y) = \int_{-\infty}^{\infty} \rho(x,y,z)$$

In the Glauber model, the density of collisions at a given point is proportional to

$$n_{col}(x_{\perp}, b) = T_A\left(x_{\perp} + \frac{b}{2}\right) T_B\left(x_{\perp} - \frac{b}{2}\right)$$

where b is the impact parameter. Note the unconventional choice, we use a reference system in which $x_{\perp}=0$ corresponds to the center of the overlapping region, not to the center of one of the nuclei.

Some definitions II

A participant is a nucleon that collides at least once. The density of participants at a given point in the transverse plane is

$$n_{part}(x_{\perp}, b) = T_A \left(x_{\perp} + \frac{b}{2} \right) \left(1 - \left(1 - \frac{T_B \left(x_{\perp} - \frac{b}{2} \right) \sigma}{B} \right)^B \right)$$
$$+ T_B \left(x_{\perp} - \frac{b}{2} \right) \left(1 - \left(1 - \frac{T_A \left(x_{\perp} + \frac{b}{2} \right) \sigma}{A} \right)^A \right)$$

 R_{AA} is often given as a function of the total number of participants N_{part} . More participants implies more central collisions.

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In this model, the probability to create a particle at a given point in the transverse plane is proportional to

$$S^{sh}(\mathbf{s}, \mathbf{b}) = \frac{T_A\left(\mathbf{s} + \frac{\mathbf{b}}{2}\right)}{1 + A F(y, p_T)T_A\left(\mathbf{s} + \frac{\mathbf{b}}{2}\right)} \frac{T_B\left(\mathbf{s} - \frac{\mathbf{b}}{2}\right)}{1 + B F(-y, p_T)T_B\left(\mathbf{s} - \frac{\mathbf{b}}{2}\right)}$$

Note that it depends on both rapidity and transverse momentum.

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- Therefore:

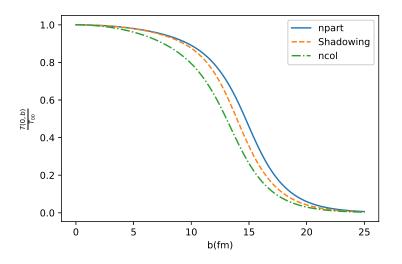
$$T_0(x_{\perp},b) = T_{00} \left(\frac{S_{\pi}^{sh}(x_{\perp},b)}{S_{\pi}^{sh}(0,0)} \right)^{1/4}$$

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• As a comparison, we can also compute the temperature that we obtain changing S_{π}^{sh} to n_{part} or n_{col} .

Temperature in the center of the overlapping region



Temperature by quarkonium

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$$\langle T_{HQ}(b) \rangle = T_{00} \frac{\int d^2 s S_{HQ}^{sh}(\mathbf{s}, \mathbf{b}) \left(\frac{S_{\pi}^{sh}(\mathbf{s}, \mathbf{b})}{S_{\pi}^{sh}(\mathbf{0}, \mathbf{0})}\right)^{1/4}}{\int d^2 s S_{HQ}^{sh}(\mathbf{s}, \mathbf{b})}$$

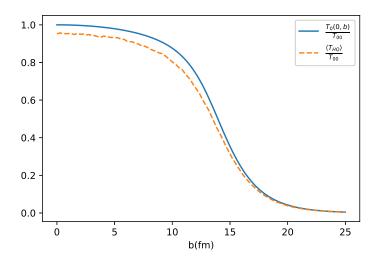
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 We might ask if the temperature seen by a quarkonium state is close to the temperature at the center of the plateau.

$\langle T_{HQ} \rangle$ versus $T_0(0,b)$



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- A pNRQCD inspired model.
- A model that takes into account the finite energy gap between singlets and octets.

Computation of R_{AA}

$$R_{AB}(b) = \frac{N_{HQ}^{AB}(b)}{N_{HQ}^{pp} T_{AB}(b)}$$

where

$$N_{HQ}^{AB}(b) = N_{HQ}^{pp} \int d^2s S_{HQ}^{sh}(\mathbf{s}, \mathbf{b}) S_{med}(\mathbf{s}, \mathbf{b})$$

and

$$T_{AB}(b) = \int d^2s T_A\left(\mathbf{s} + \frac{\mathbf{b}}{2}\right) T_B\left(\mathbf{s} - \frac{\mathbf{b}}{2}\right)$$

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- If $S_{med} = 1$ there are no medium effects. R_{AA}^{CNM} .
- If $S_{HQ}^{sh} = T_{AB}$ we ignore the triple Pomeron vertex. Original Glauber model. R_{AA}^{T} .

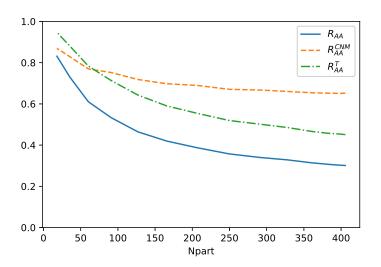
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ight)^{3t_0\Gamma(\mathbf{s}, \mathbf{b}, t_0)} & T_0(\mathbf{s}, \mathbf{b}) \geq T_f \ 1 & T_0(\mathbf{s}, \mathbf{b}) < T_f \end{cases}$$



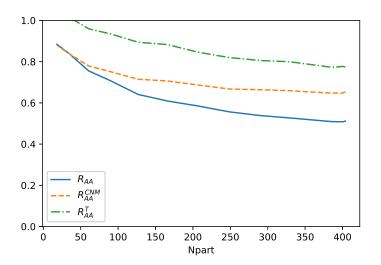
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$$S_{med}(\mathbf{s}, \mathbf{b}) = \begin{cases} e^{-\frac{3aT_0(\mathbf{s}, \mathbf{b})^3t_0}{b^2} \left(e^{-\frac{b}{T_0(\mathbf{s}, \mathbf{b})} \left(1 + \frac{b}{T_0(\mathbf{s}, \mathbf{b})} \right) - e^{-\frac{b}{T_f}} \left(1 + \frac{b}{T_f} \right) \right)} & T_0(\mathbf{s}, \mathbf{b}) > T_f \\ 1 & T_0(\mathbf{s}, \mathbf{b}) \le T_f \end{cases}$$



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- Qualitatively, the suppression due to initial effects is less dependent on centrality.
- We have developed a simple framework to include initial state effects given S_{med} .