

Signatures of gluon saturation from structure-function measurements

arXiv:2203.05846

DIS2022 WG6

Nestor Armesto ¹ Tuomas Lappi ^{2 3} Heikki Mäntysaari ^{2 3}
Hannu Paukkunen ^{2 3} Mirja Tevio ^{2 3}

¹Instituto Galego de Física de Altas Enerxías IGFAE, Universidade de Santiago de Compostela, 15782 Santiago de Compostela, Galicia-Spain

²Department of Physics, University of Jyväskylä, P.O. Box 35, 40014 University of Jyväskylä, Finland

³Helsinki Institute of Physics, P.O. Box 64, 00014 University of Helsinki, Finland

3.5.2022

- Color Glass Condensate (CGC) framework describes non-linear effects (gluon saturation)
 - ▶ Bjorken- x dependence from Balitsky-Kovchegov (BK) evolution equation

- Color Glass Condensate (CGC) framework describes non-linear effects (gluon saturation)
 - ▶ Bjorken- x dependence from Balitsky-Kovchegov (BK) evolution equation
- In collinear factorization framework the Q^2 evolution comes from Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) evolution equations

- Color Glass Condensate (CGC) framework describes non-linear effects (gluon saturation)
 - ▶ Bjorken- x dependence from Balitsky-Kovchegov (BK) evolution equation
- In collinear factorization framework the Q^2 evolution comes from Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) evolution equations
- To see saturation effects on experimental data we have to distinguish the genuine difference between DGLAP and BK dynamics

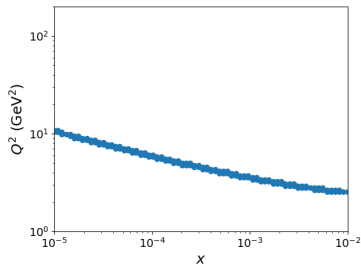
- Color Glass Condensate (CGC) framework describes non-linear effects (gluon saturation)
 - ▶ Bjorken- x dependence from Balitsky-Kovchegov (BK) evolution equation
- In collinear factorization framework the Q^2 evolution comes from Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) evolution equations
- To see saturation effects on experimental data we have to distinguish the genuine difference between DGLAP and BK dynamics
- Both frameworks require input which are fitted to the same experimental data
→ The results do not deviate dramatically and the distinguishing DGLAP/BK evolution is difficult

Our method to see difference in DGLAP/BK

- 1 We want to be as independent as possible of initial condition parametrization

Our method to see difference in DGLAP/BK

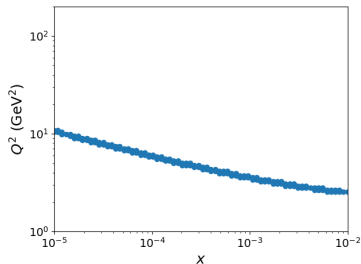
- 1 We want to be as independent as possible of initial condition parametrization
- 2 We "force" collinear factorization and CGC $F_{2,L}$ to agree in a line in (x, Q^2) plane



Matching line in (x, Q^2) plane

Our method to see difference in DGLAP/BK

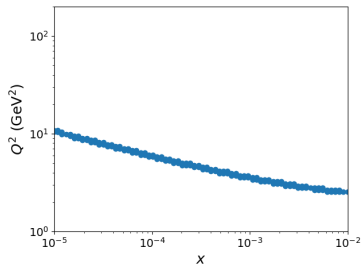
- 1 We want to be as independent as possible of initial condition parametrization
- 2 We "force" collinear factorization and CGC $F_{2,L}$ to agree in a line in (x, Q^2) plane
- 3 Differences between the two frameworks outside the chosen line quantify signatures of gluon saturation



Matching line in (x, Q^2) plane

Our method to see difference in DGLAP/BK

- 1 We want to be as independent as possible of initial condition parametrization
- 2 We "force" collinear factorization and CGC $F_{2,L}$ to agree in a line in (x, Q^2) plane
- 3 Differences between the two frameworks outside the chosen line quantify signatures of gluon saturation
- 4 With differences we can approximate the accuracy of $F_{2,L}$ saturation measurements in EIC and LHeC/FCC-he



Matching line in (x, Q^2) plane

$F_{2,L}$ with collinear factorization vs. CGC

Collinear factorization:

- Collinear factorization $F_{2,L}$ using APFEL [1] and LHAPDF [2] libraries
- NNPDF31_nlo_as_0118_1000 as proton PDF set
- nNNPDF20_nlo_as_0118_Au197 as nuclear PDF set
- Both PDF sets have 1000 Monte Carlo replicas

Color Glass Condensate (CGC):

- Dipole picture $F_{2,L}$ fitted to HERA data
- Leading order total photon-nucleus cross sections
- Running coupling BK evolution ¹

- We match collinear factorization $F_{2,L}$ to corresponding CGC structure functions in a line in (x, Q^2) plane

¹T. Lappi and H. Mäntysaari. “Single inclusive particle production at high energy from HERA data to proton-nucleus collisions”. In: *Phys. Rev. D* 88 (2013), p. 114020. [arXiv: 1309.6963 \[hep-ph\]](#)

Bayesian reweighting method [4, 5]:

For each PDF replica f_k we define

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(\mathcal{O}_i - \mathcal{O}_i[f_k])^2}{(\delta_{\text{BK}} \mathcal{O}_i)^2}$$

Bayesian reweighting method [4, 5]:

For each PDF replica f_k we define

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(\mathcal{O}_i - \mathcal{O}_i[f_k])^2}{(\delta_{\text{BK}} \mathcal{O}_i)^2}$$

and so called **Giele-Keller** weights
[6]

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}},$$

Bayesian reweighting method [4, 5]:

For each PDF replica f_k we define

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(\mathcal{O}_i - \mathcal{O}_i[f_k])^2}{(\delta_{\text{BK}} \mathcal{O}_i)^2}$$

and so called **Giele-Keller** weights
[6]

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}},$$

which always sum up to unity,

$$\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k = 1.$$

Bayesian reweighting method [4, 5]:

For each PDF replica f_k we define

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(\mathcal{O}_i - \mathcal{O}_i[f_k])^2}{(\delta_{\text{BK}} \mathcal{O}_i)^2}$$

and so called **Giele-Keller** weights
[6]

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}},$$

which always sum up to unity,

$$\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k = 1.$$

Giele-Keller weights favor replicas
with $\chi^2 \approx 0$.

Bayesian reweighting method [4, 5]:

For each PDF replica f_k we define

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(\mathcal{O}_i - \mathcal{O}_i[f_k])^2}{(\delta_{\text{BK}} \mathcal{O}_i)^2}$$

and so called **Giele-Keller** weights [6]

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}},$$

which always sum up to unity,

$$\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k = 1.$$

Giele-Keller weights favor replicas with $\chi^2 \approx 0$.

Then we define reweighted observables as

$$\mathcal{O}^{\text{Rew}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k \mathcal{O}[f_k]$$

We also construct a PDF set matched to BK in (x, Q^2) line (Back up)

Fixing matching parameters

- We want to match the reweighted values to BK values as closely as possible
 - ▶ Finite number of replicas (1000) prevent the absolute match
- Effective number of replicas [4, 7]

$$N_{\text{eff}} = \exp \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k \ln \left(\frac{N_{\text{rep}}}{\omega_k} \right)$$

gives an approximation on how many PDF replicas have significant weight

- We adjust δ_{BK} in χ_k^2 in order to fix $N_{\text{eff}} \approx 10$

$$\chi_k^2 = \sum_{i=1}^{N_{\text{data}}} \frac{(y_i - y_i[f_k])^2}{(\delta_{\text{BK}} y_i)^2}$$

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}}$$

$$\mathcal{O}^{\text{Rew}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \omega_k \mathcal{O}[f_k]$$

Choosing the matching line

- We want to do the matching in a common region of validity for both frameworks:

Choosing the matching line

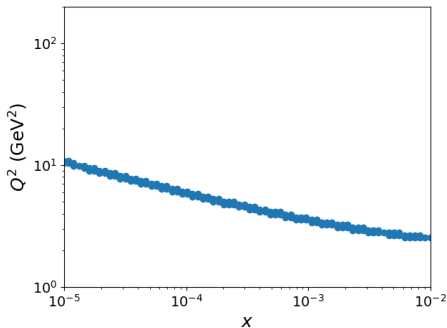
- We want to do the matching in a common region of validity for both frameworks:
 - ▶ In a region $Q^2 \gg Q_s^2$ where saturation effects are moderate

Choosing the matching line

- We want to do the matching in a common region of validity for both frameworks:
 - ▶ In a region $Q^2 \gg Q_s^2$ where saturation effects are moderate
 - ▶ With enough small $\alpha_s \log(Q^2)$ so that DGLAP evolution dynamics is reliable

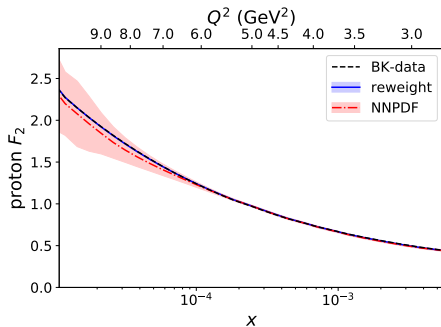
Choosing the matching line

- We want to do the matching in a common region of validity for both frameworks:
 - ▶ In a region $Q^2 \gg Q_s^2$ where saturation effects are moderate
 - ▶ With enough small $\alpha_s \log(Q^2)$ so that DGLAP evolution dynamics is reliable
- **We choose to do the matching on points $Q^2(x) \approx 10 \times Q_s^2(x)$**

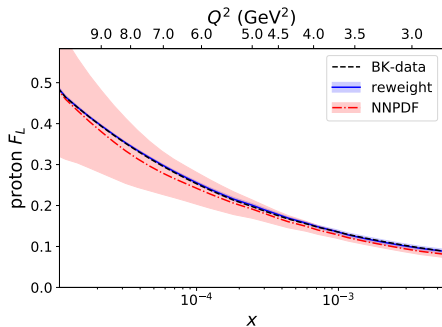


Matching line in (x, Q^2) plane

Proton matching



(a) F_2

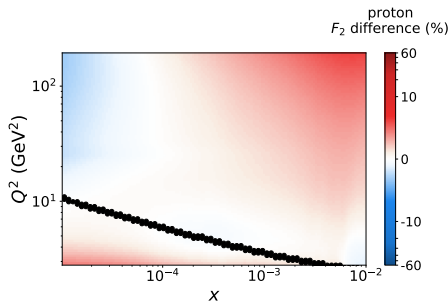


(b) F_L

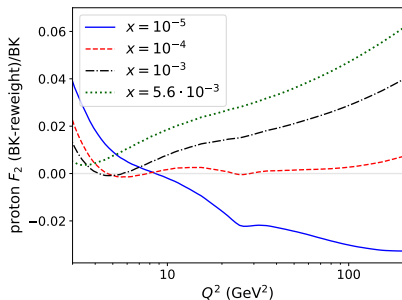
The structure functions for proton as a function of x at $Q^2 \approx 10Q_s^2(x)$

- Separate matching for proton F_2 and F_L are both almost perfect

Relative difference of proton F_2^{Rew} to F_2^{BK}



(a) F_2

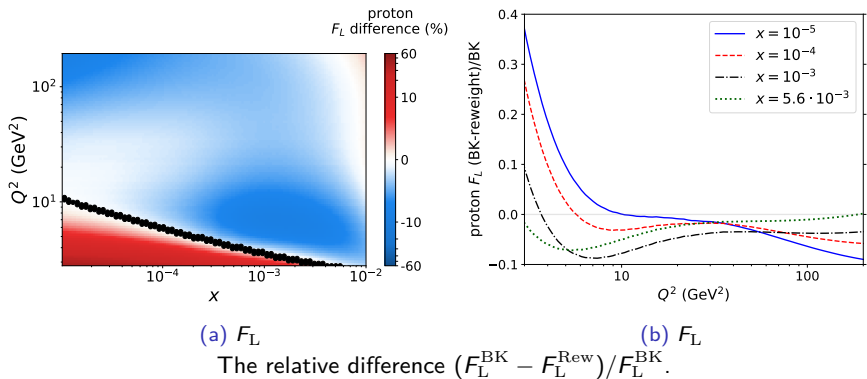


(b) F_2

Relative difference $(F_2^{\text{BK}} - F_2^{\text{Rew}})/F_2^{\text{BK}}$

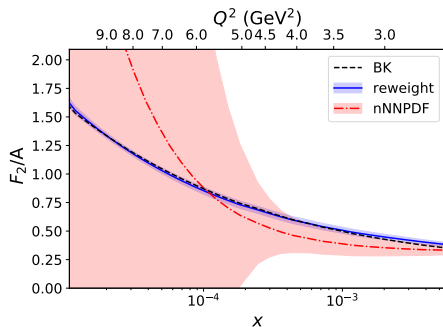
- For proton F_2 the relative difference is only a few percent
- Generically slower x dependence in BK evolution

Relative difference of proton F_L^{Rew} to F_L^{BK}

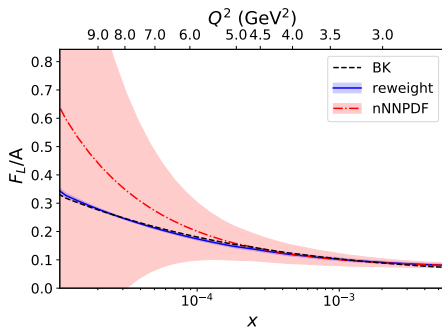


- For proton F_L the relative difference is:
 - ▶ $\lesssim 10\%$ for $x = 10^{-3} \dots 5.6 \times 10^{-3}$ (EIC)
 - ▶ $\lesssim 40\%$ for $x = 10^{-5} \dots 10^{-4}$ (LHeC/FCC-he)
- F_L is much more sensitive to saturation than F_2

Nuclear matching



(a) F_2

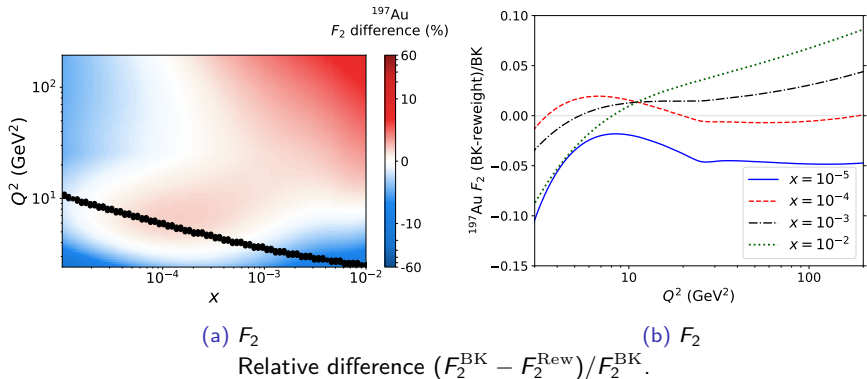


(b) F_L

The structure functions for ^{197}Au as a function of x at $Q^2 \approx 10Q_s^2(x)$.

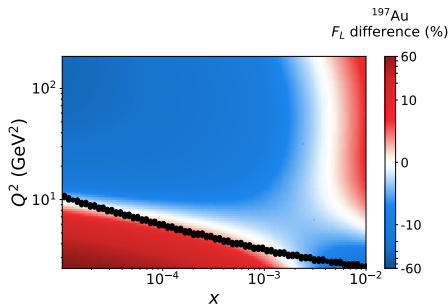
- Nuclear reweight is not as successful as for proton since there are not enough Monte Carlo replicas to get a precise match

Relative difference of nuclear F_2 to F_2^{BK}

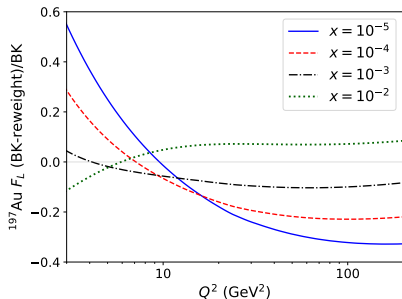


- For nuclear F_2 the relative difference is $\lesssim 10\%$
- The relative difference is much larger than in the proton case
 - ▶ It is expected since saturation effects are stronger in nuclei

Relative difference of nuclear F_L^{Rew} to F_L^{BK}



(a) F_L



(b) F_L

The relative difference $(F_L^{\text{BK}} - F_L^{\text{Rew}})/F_L^{\text{BK}}$.

For nuclear F_L the relative difference is:

- $\lesssim 15\%$ for $x = 10^{-3} \dots 10^{-2}$ (EIC)
- $\lesssim 60\%$ for $x = 10^{-5} \dots 10^{-4}$ (LHeC/FCC-he)

- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$

Summary

- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$
- The deviation outside the matching line describes signatures of saturation

- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$
- The deviation outside the matching line describes signatures of saturation
- In order to see saturation in protons in EIC
 - ▶ F_L the measurements have to be $\mathcal{O}(10\%)$
 - ▶ F_2 the measurements have to be $\mathcal{O}(1\%)$

- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$
- The deviation outside the matching line describes signatures of saturation
- In order to see saturation in protons in EIC
 - ▶ F_L the measurements have to be $\mathcal{O}(10\%)$
 - ▶ F_2 the measurements have to be $\mathcal{O}(1\%)$
- In LHeC/FCC-he the differences are a few times larger

Summary

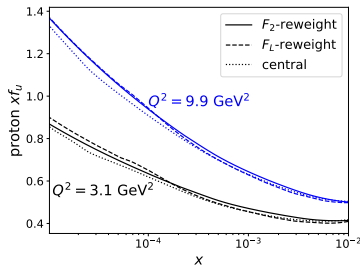
- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$
- The deviation outside the matching line describes signatures of saturation
- In order to see saturation in protons in EIC
 - ▶ F_L the measurements have to be $\mathcal{O}(10\%)$
 - ▶ F_2 the measurements have to be $\mathcal{O}(1\%)$
- In LHeC/FCC-he the differences are a few times larger
- Saturation is stronger in heavy nuclei than in proton

Summary

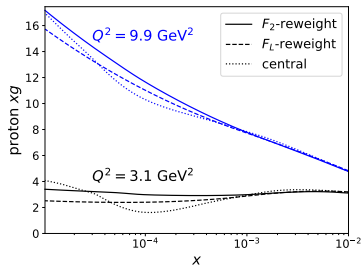
- With Bayesian reweighting we match proton/nuclear structure functions to corresponding BK values in a line $Q^2 \approx 10 \times Q_s^2(x)$
- The deviation outside the matching line describes signatures of saturation
- In order to see saturation in protons in EIC
 - ▶ F_L the measurements have to be $\mathcal{O}(10\%)$
 - ▶ F_2 the measurements have to be $\mathcal{O}(1\%)$
- In LHeC/FCC-he the differences are a few times larger
- Saturation is stronger in heavy nuclei than in proton
- F_L is more sensitive to saturation than F_2

- [1] Valerio Bertone, Stefano Carrazza, and Juan Rojo. “APFEL: A PDF Evolution Library with QED corrections”. In: *Comput. Phys. Commun.* 185 (2014), pp. 1647–1668. DOI: 10.1016/j.cpc.2014.03.007. arXiv: 1310.1394 [hep-ph].
- [2] Andy Buckley et al. “LHAPDF6: parton density access in the LHC precision era”. In: *Eur. Phys. J. C* 75 (2015), p. 132. DOI: 10.1140/epjc/s10052-015-3318-8. arXiv: 1412.7420 [hep-ph].
- [3] T. Lappi and H. Mäntysaari. “Single inclusive particle production at high energy from HERA data to proton-nucleus collisions”. In: *Phys. Rev. D* 88 (2013), p. 114020. arXiv: 1309.6963 [hep-ph].
- [4] Richard D. Ball et al. “Reweighting NNPDFs: the W lepton asymmetry”. In: *Nucl. Phys. B* 849 (2011). [Erratum: *Nucl.Phys.B* 854, 926–927 (2012), Erratum: *Nucl.Phys.B* 855, 927–928 (2012)], pp. 112–143. DOI: 10.1016/j.nuclphysb.2011.03.017. arXiv: 1012.0836 [hep-ph].
- [5] Richard D. Ball et al. “Reweighting and Unweighting of Parton Distributions and the LHC W lepton asymmetry data”. In: *Nucl. Phys. B* 855 (2012), pp. 608–638. DOI: 10.1016/j.nuclphysb.2011.10.018. arXiv: 1108.1758 [hep-ph].
- [6] Walter T. Giele and Stephane Keller. “Implications of hadron collider observables on parton distribution function uncertainties”. In: *Phys. Rev. D* 58 (1998), p. 094023. DOI: 10.1103/PhysRevD.58.094023. arXiv: hep-ph/9803393.
- [7] Hannu Paukkunen and Pia Zurita. “PDF reweighting in the Hessian matrix approach”. In: *JHEP* 12 (2014), p. 100. DOI: 10.1007/JHEP12(2014)100. arXiv: 1402.6623 [hep-ph].

Back up: Weighted proton PDFs



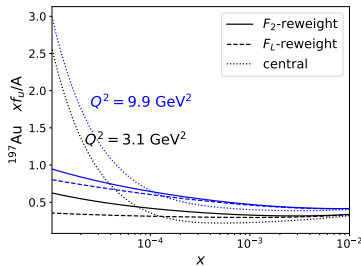
(a) Proton up quark



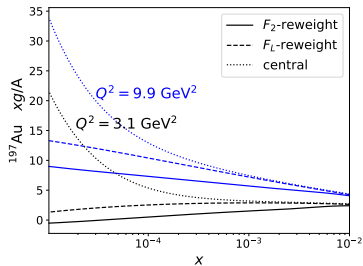
(b) Proton gluon

- Reweighting has slightly stronger effect on gluon distribution than on up quark
- Moderate effects expected since NNPDF3.1 PDFs are fitted to same HERA data as BK boundary conditions

Back up: Weighted nuclear PDFs



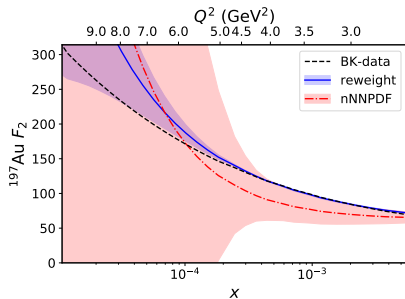
(c) Nuclear up quark



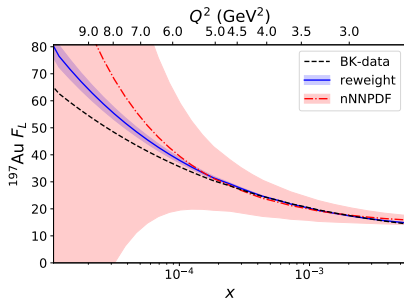
(d) Nuclear gluon

- Nuclear PDFs are affected more than proton PDFs
- Reweighting has stronger effect on gluon distribution than on up quark

Back up: Reweight with smaller x region



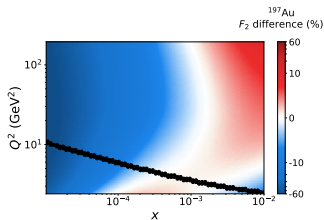
(a) F_2



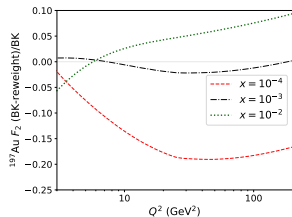
(b) F_L

Nuclear reweight in region $x = 10^{-4} \dots 10^{-2}$

Back up: Reweight with smaller x region

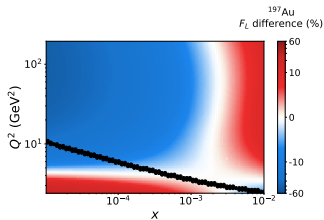


(a) F_2

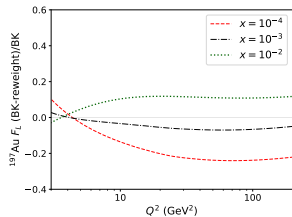


(b) F_2

The relative difference $(F_2^{\text{BK}} - F_2^{\text{Rew}})/F_2^{\text{BK}}$ with nuclear reweight in region $x = 10^{-4} \dots 10^{-2}$.



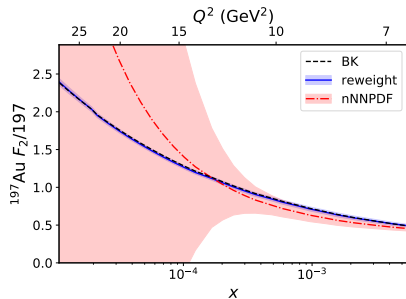
(a) F_L



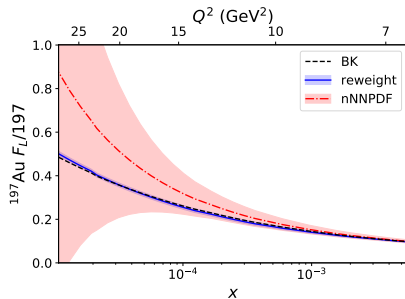
(b) F_L

The relative difference $(F_L^{\text{BK}} - F_L^{\text{Rew}})/F_L^{\text{BK}}$ with nuclear reweight in region $x = 10^{-4} \dots 10^{-2}$.

Back up: Reweight in line $Q^2(x) \approx 27 \times Q_s^2(x)$



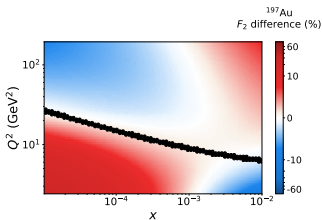
(a) F_2



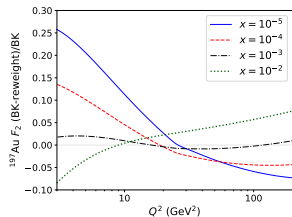
(b) F_L

Nuclear reweight in line $Q^2(x) \approx 27 \times Q_s^2(x)$.

Back up: Reweight in line $Q^2(x) \approx 27 \times Q_s^2(x)$

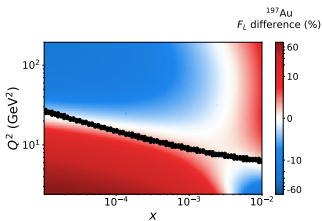


(a) F_2

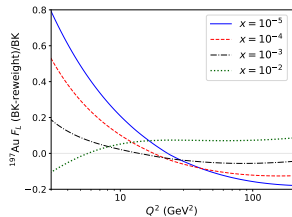


(b) F_2

The relative difference $(F_2^{\text{BK}} - F_2^{\text{Rew}})/F_2^{\text{BK}}$ with nuclear reweight in line $Q^2(x) \approx 27 \times Q_s^2(x)$.



(a) F_L



(b) F_L

The relative difference $(F_L^{\text{BK}} - F_L^{\text{Rew}})/F_L^{\text{BK}}$ with nuclear reweight in line $Q^2(x) \approx 27 \times Q_s^2(x)$.

Giele-Keller weights which favor replicas with $\chi^2/N_{\text{data}} \approx 0$

$$\omega_k = \frac{e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_k^2}}$$

Weights used with experimental data favor replicas with $\chi^2/N_{\text{data}} \approx 1$

$$\omega_k = \frac{(\chi_k^2)^{(N_{\text{data}}-1)/2} e^{-\frac{1}{2}\chi_k^2}}{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} (\chi_k^2)^{(N_{\text{data}}-1)/2} e^{-\frac{1}{2}\chi_k^2}}$$