

# ATLAS Measurements of Photo-nuclear Processes in Ultra-Peripheral Pb+Pb Collisions

---

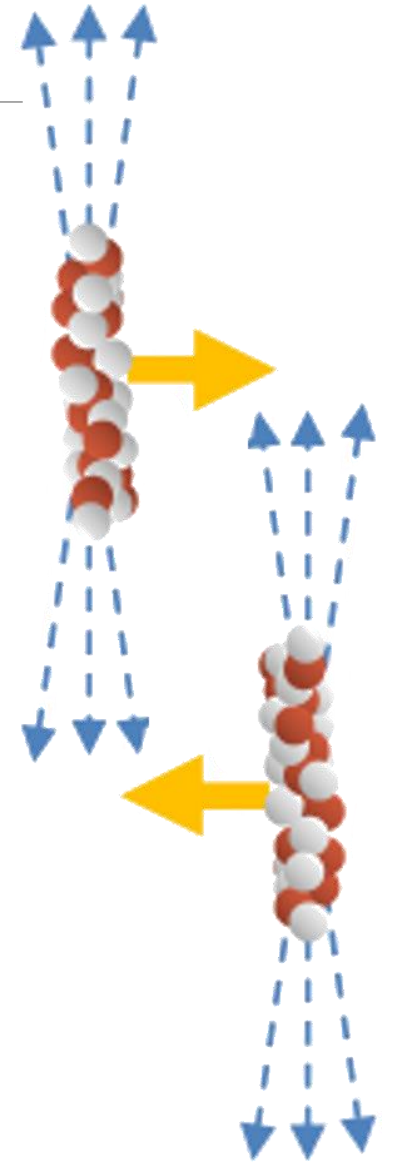
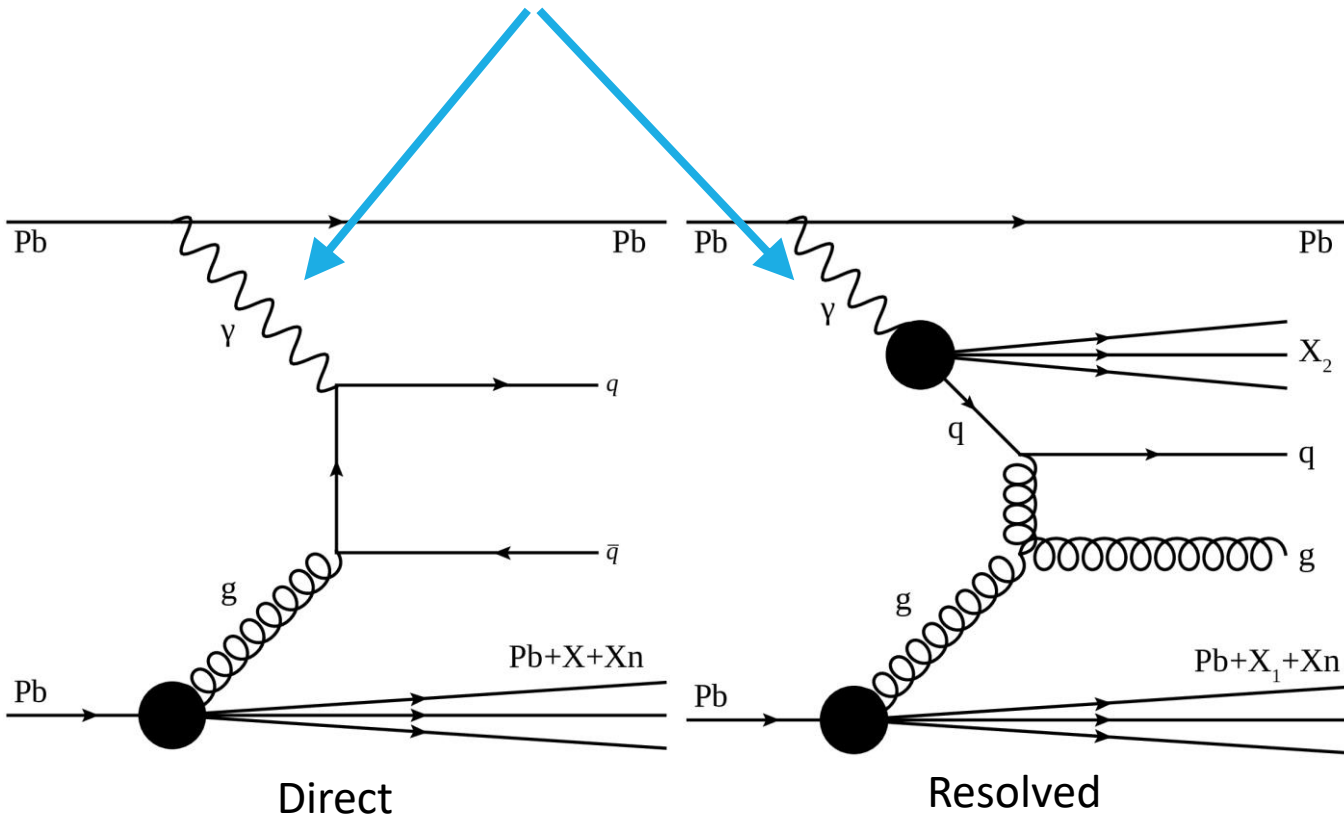
Ben Gilbert

On behalf of the ATLAS Collaboration



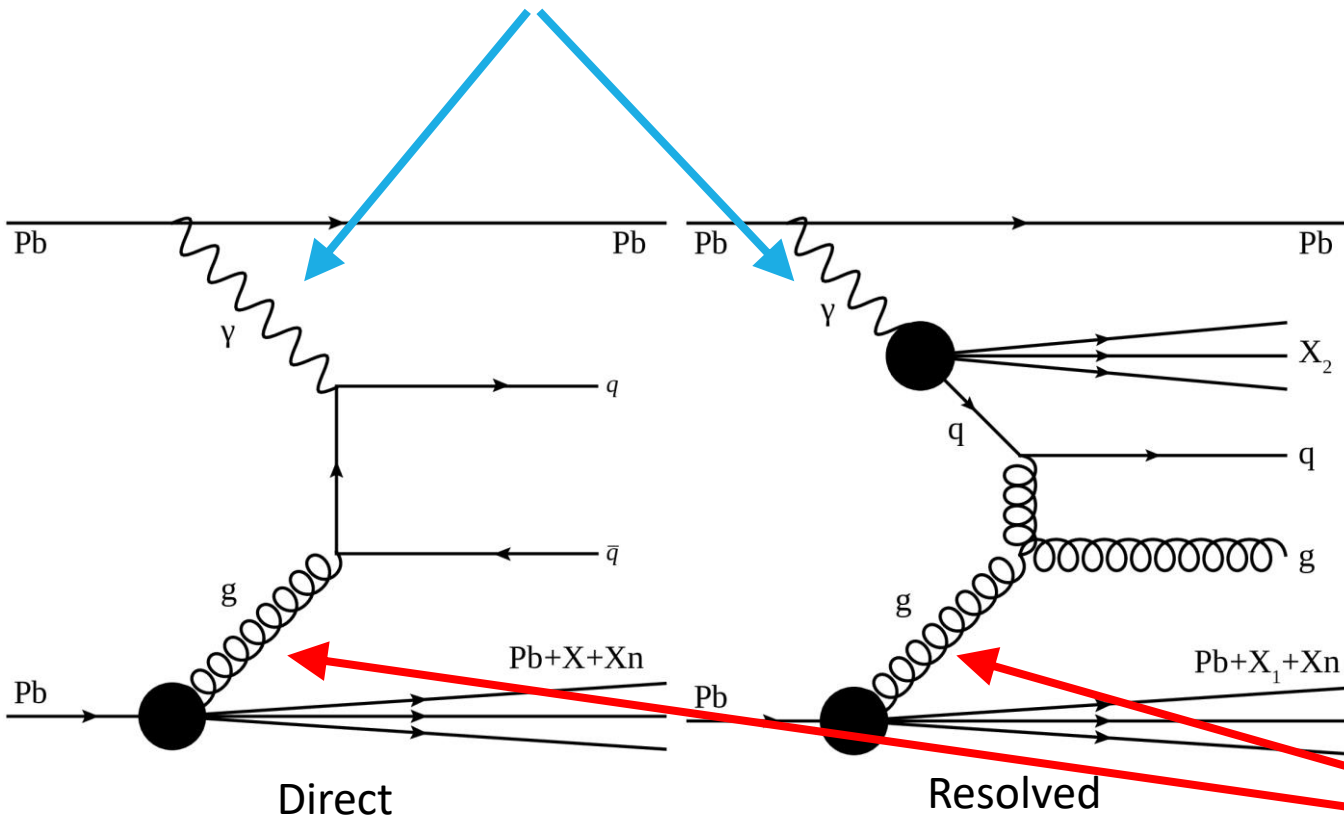
# Introduction: Photo-nuclear Processes

In ultra-relativistic heavy ion collisions, the intense electromagnetic fields of the colliding ions provide a flux of quasi-real photons.

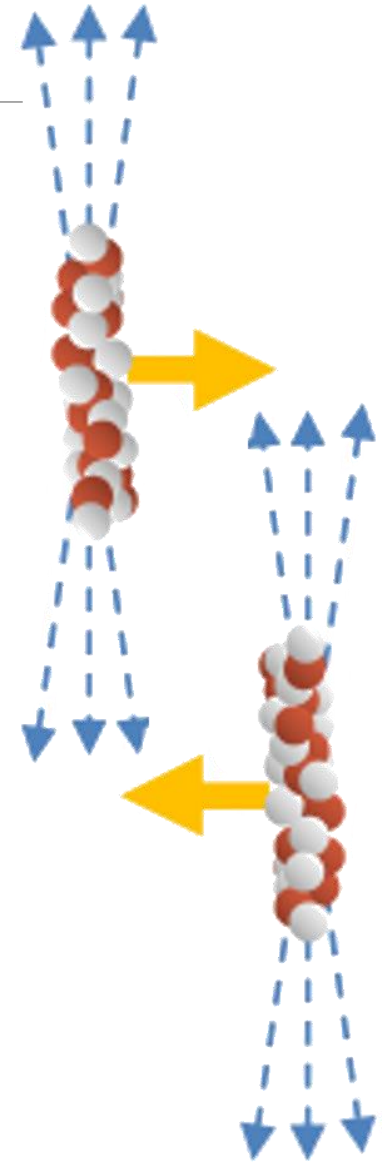


# Introduction: Photo-nuclear Processes

In ultra-relativistic heavy ion collisions, the intense electromagnetic fields of the colliding ions provide a flux of quasi-real photons.



Those photons can scatter off of partons in the other (target) nucleus

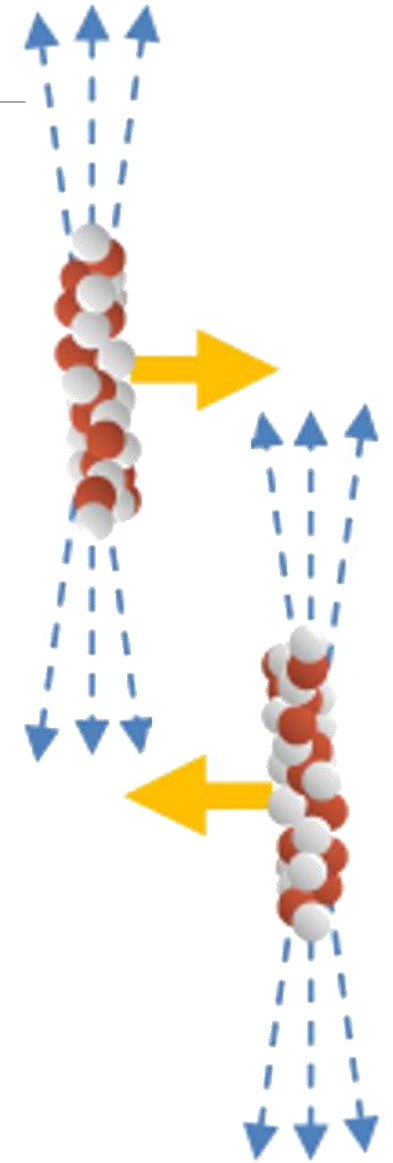
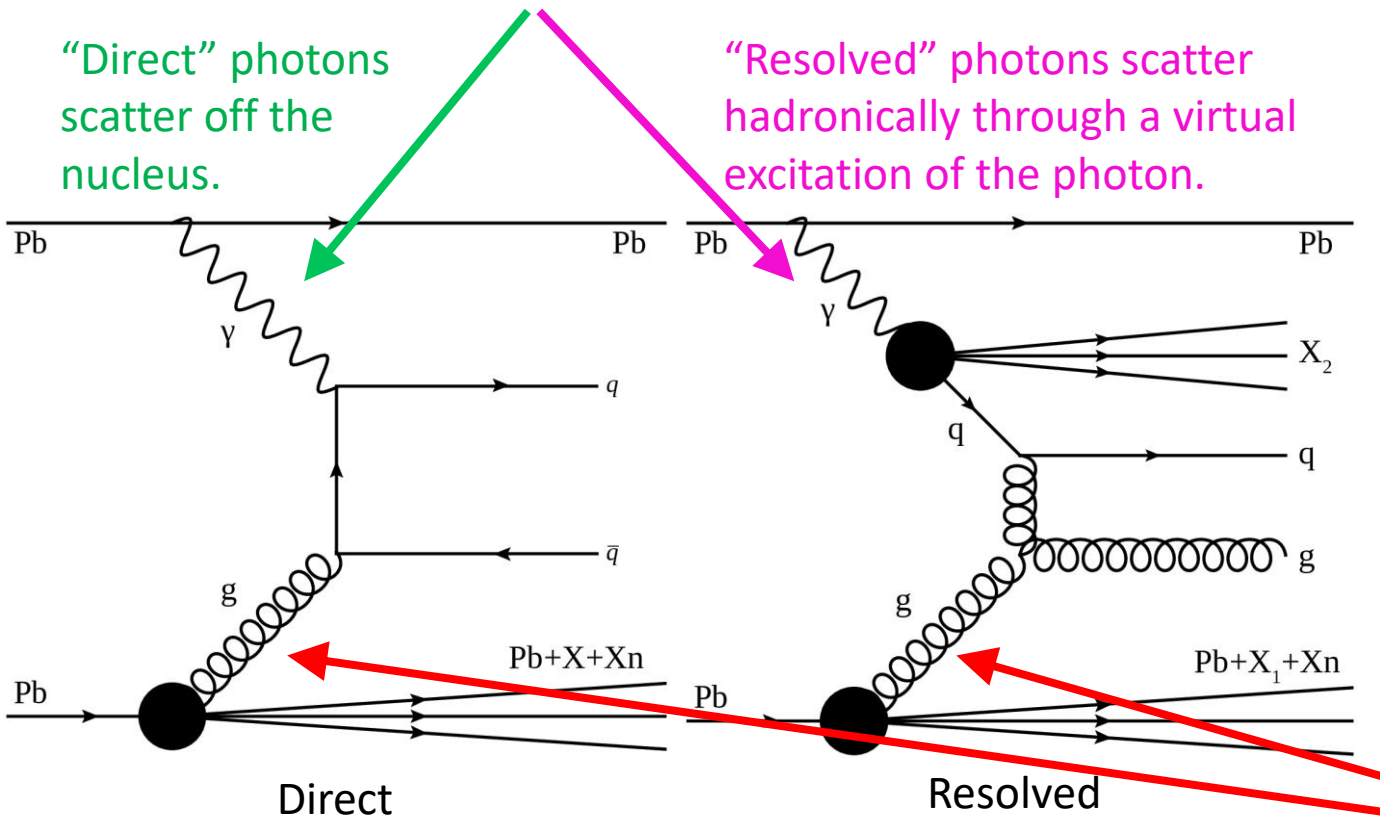


# Introduction: Photo-nuclear Processes

In ultra-relativistic heavy ion collisions, the intense electromagnetic fields of the colliding ions provide a flux of quasi-real photons.

“Direct” photons scatter off the nucleus.

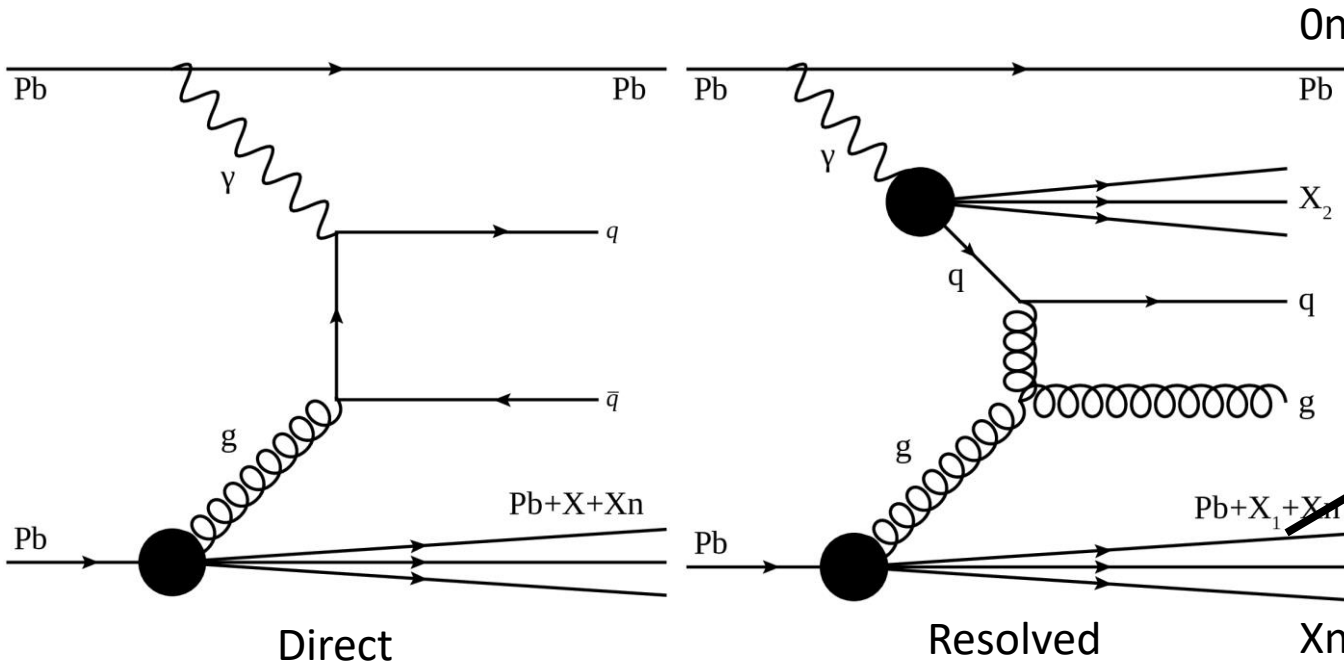
“Resolved” photons scatter hadronically through a virtual excitation of the photon.



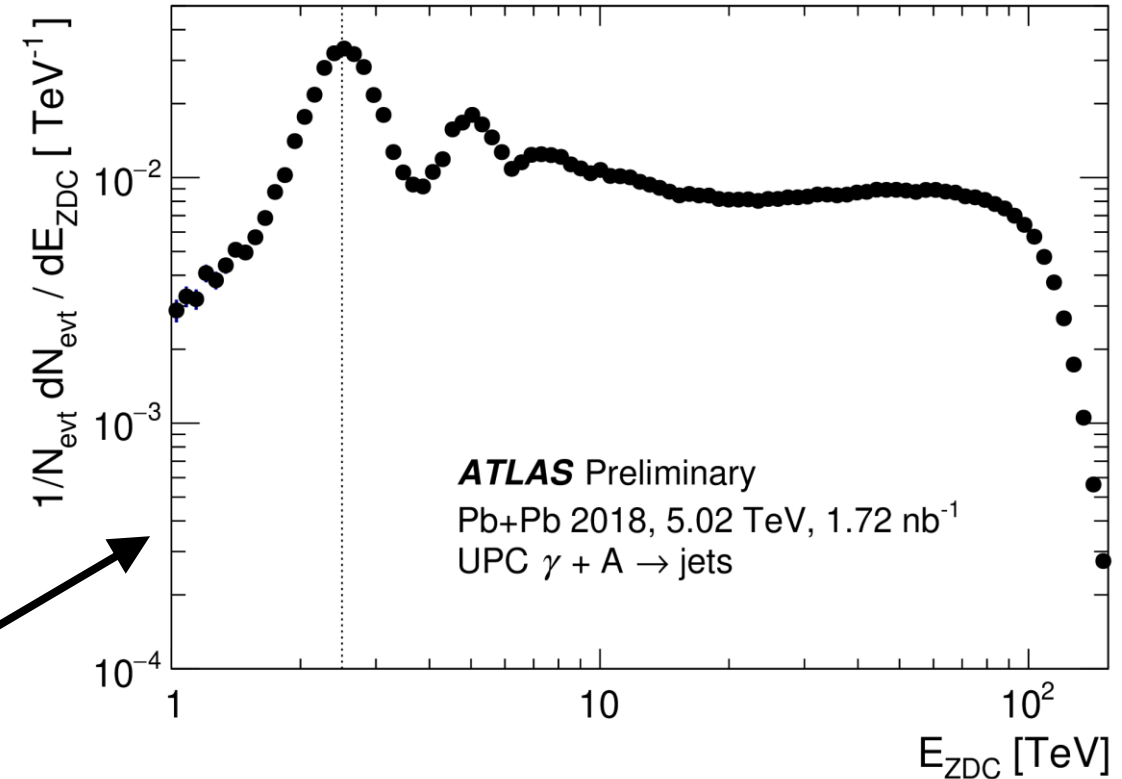
Those photons can scatter off of partons in the other (target) nucleus

# Introduction: Photo-nuclear Processes

$X_n \longleftrightarrow 0_n$

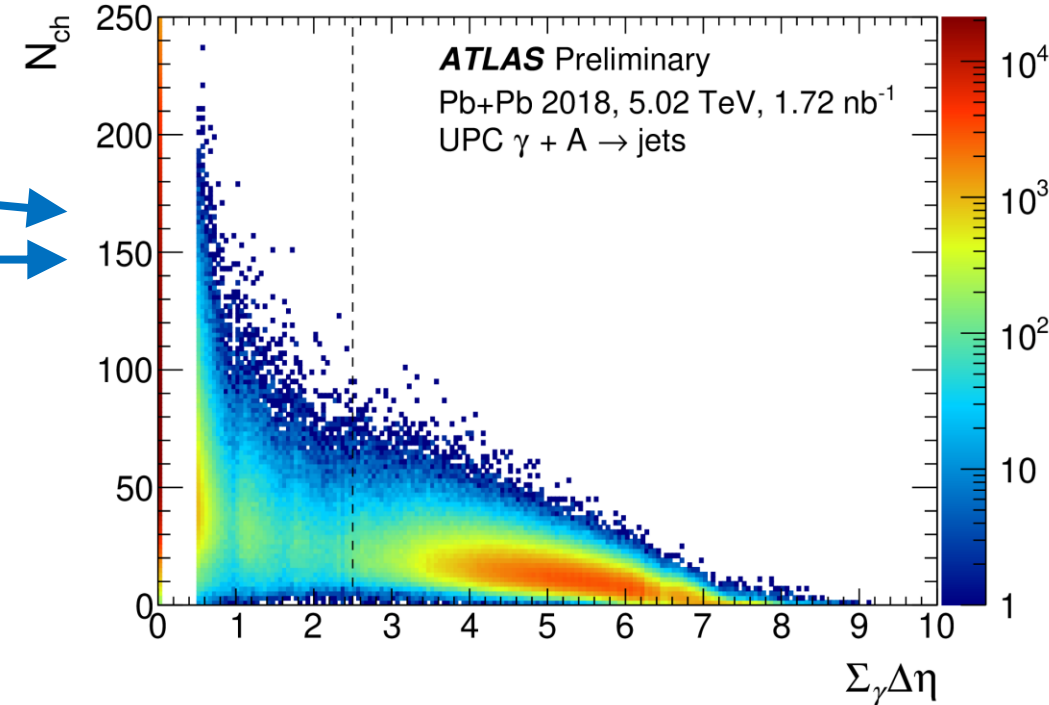
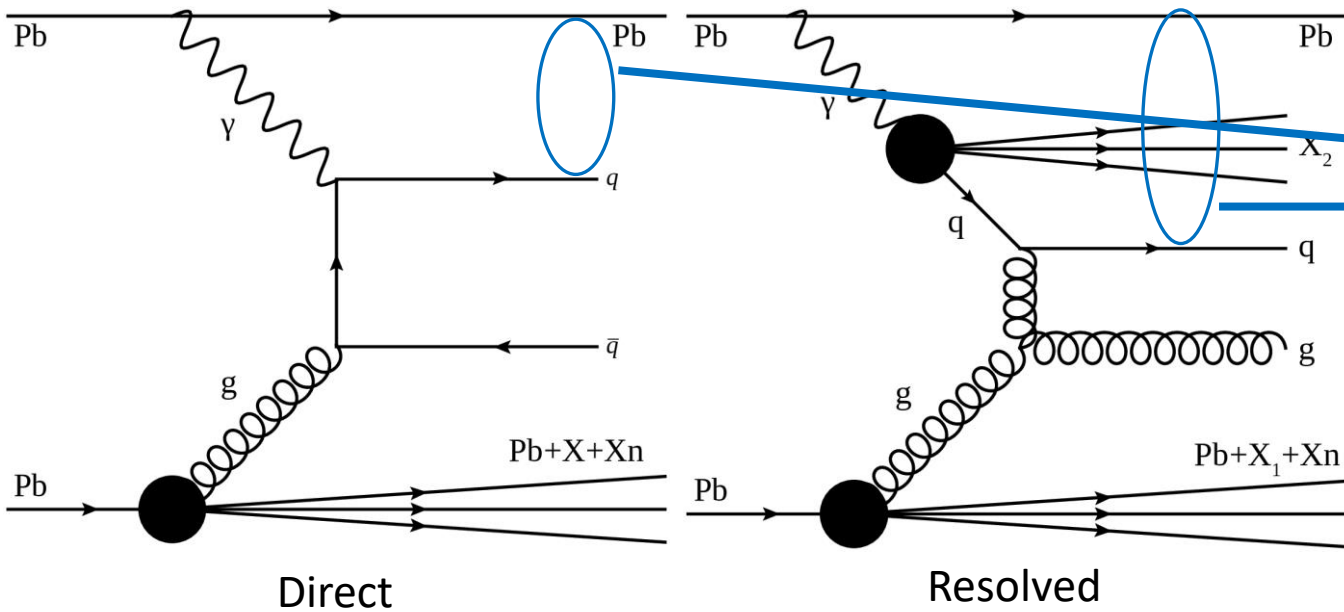
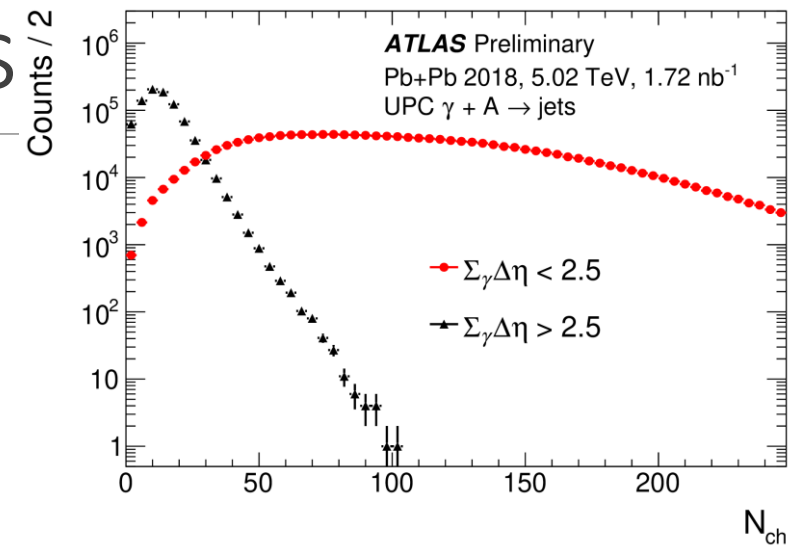


$0_n X_n$  requirement for nuclear breakup in exactly one ATLAS Zero-Degree Calorimeter (ZDC)



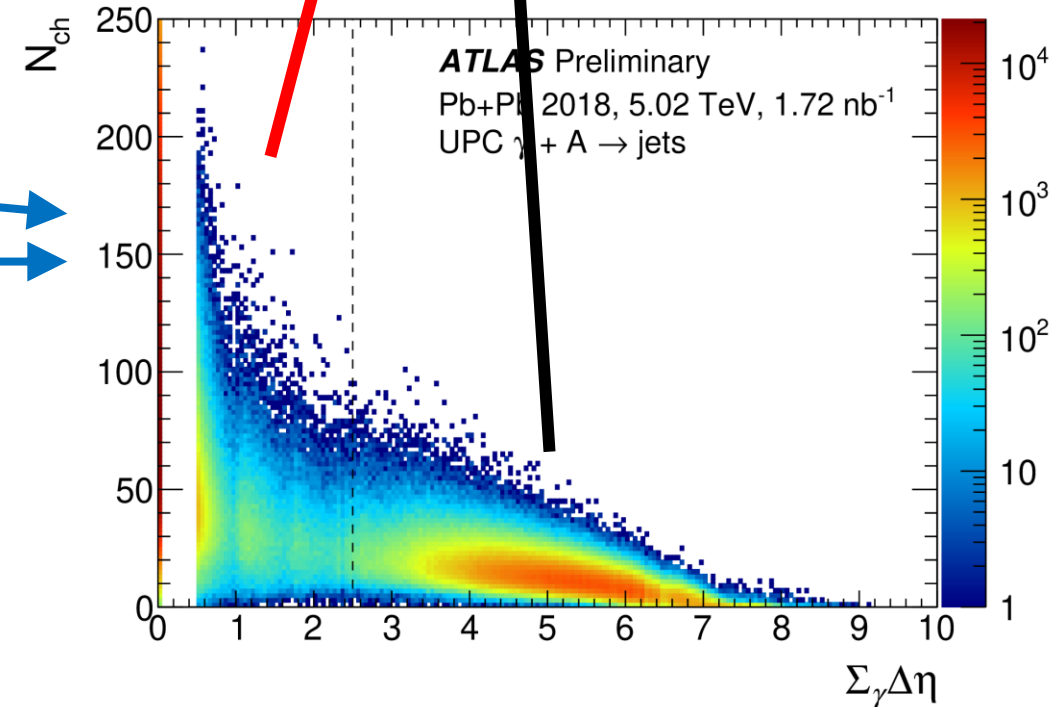
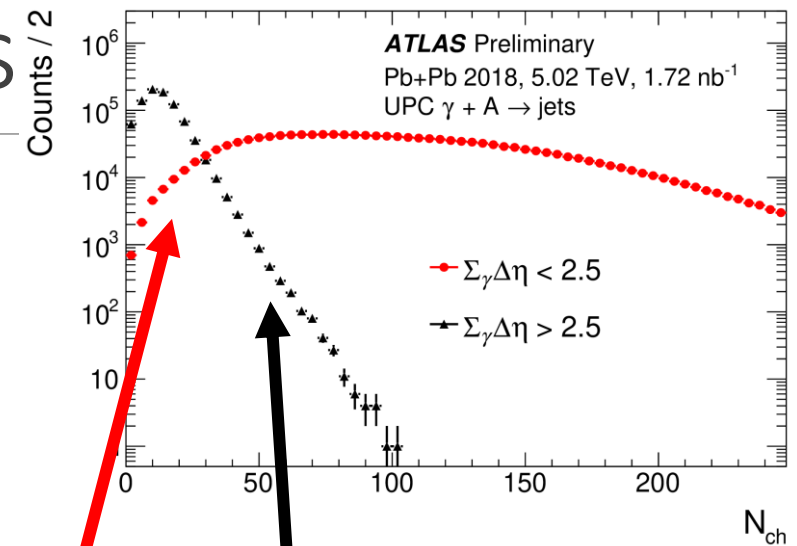
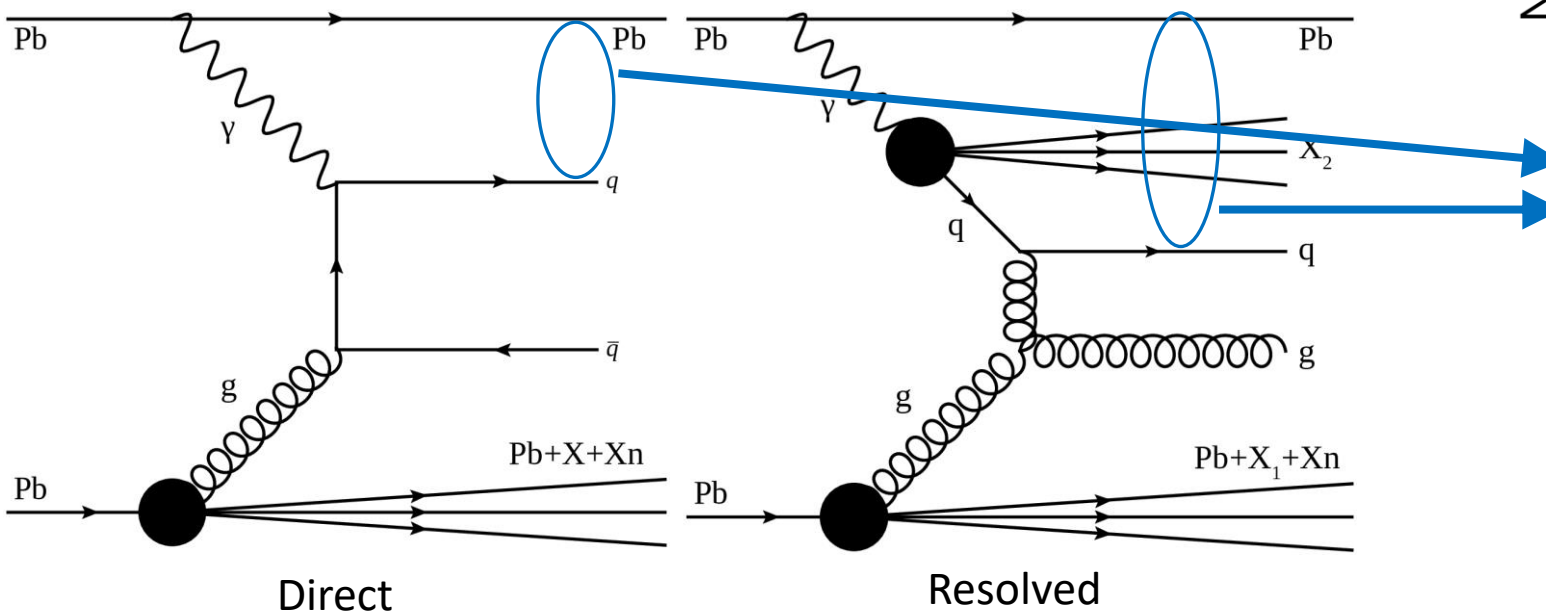
# Introduction: Photo-nuclear Processes

- Large rapidity gaps are required on one side of the detector.
  - This side is dictated by the 0nXn requirement.



# Introduction: Photo-nuclear Processes

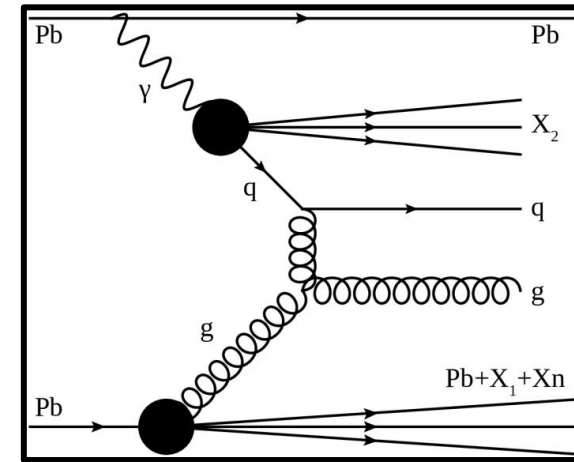
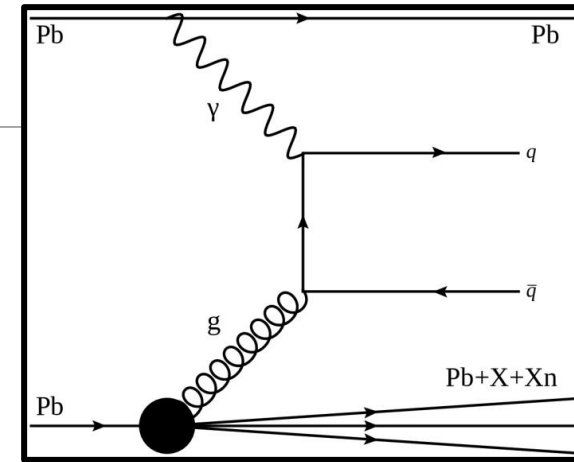
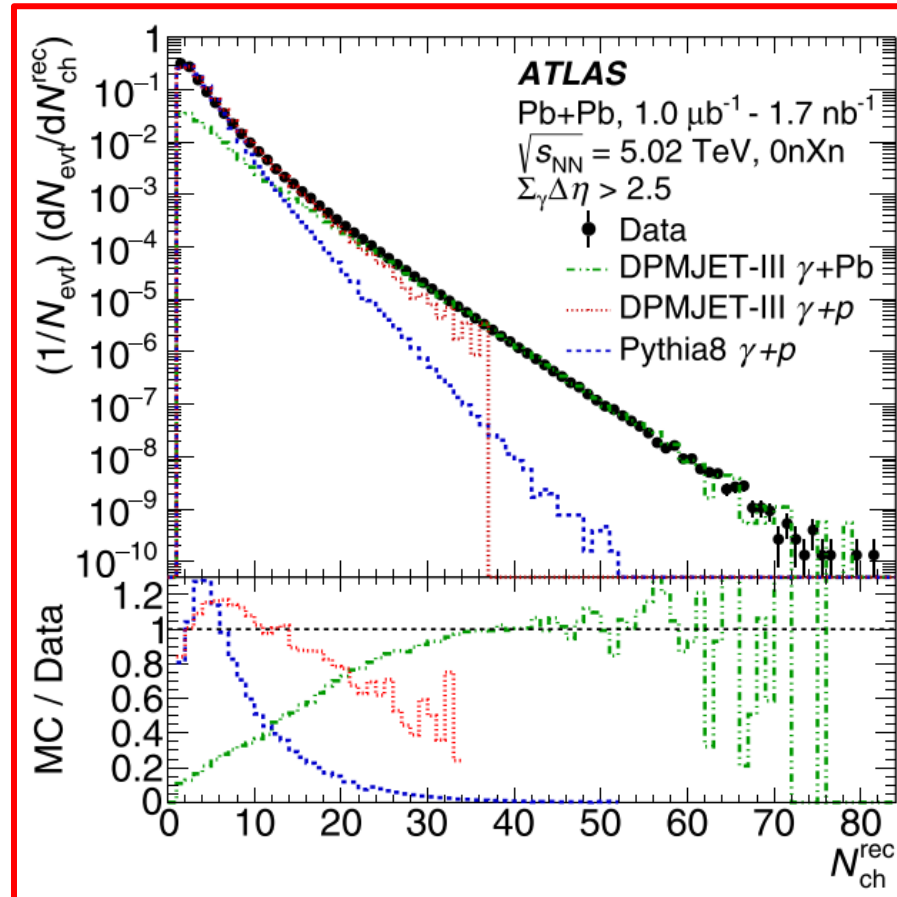
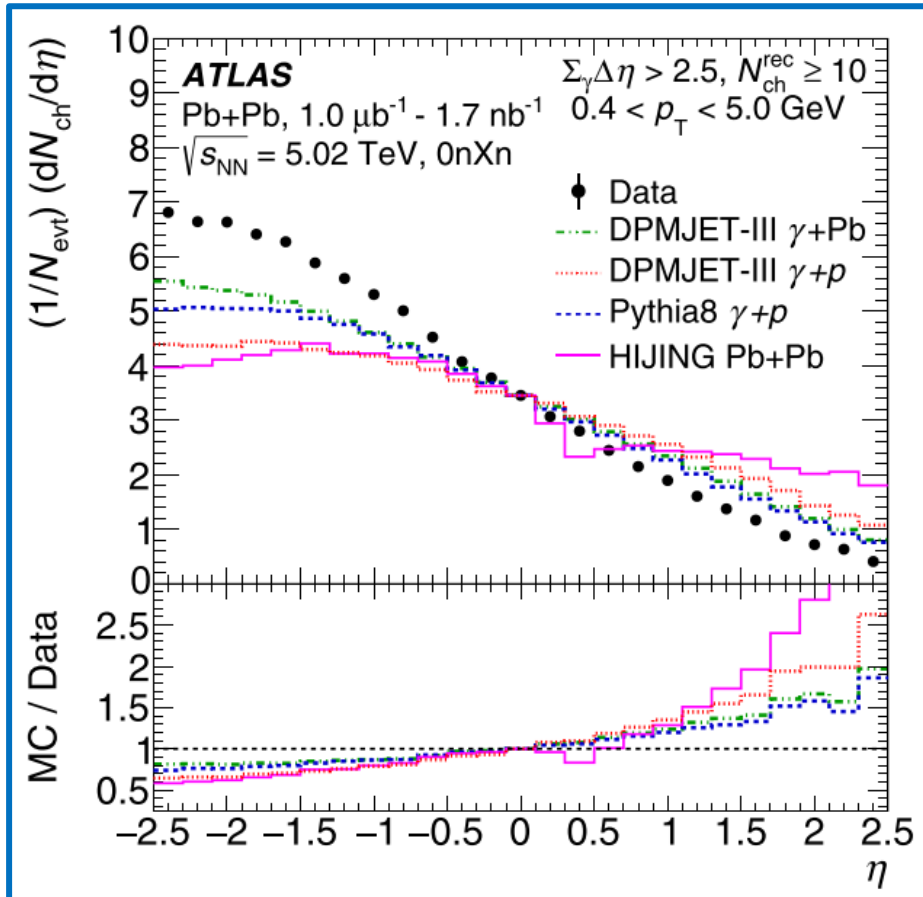
- Large rapidity gaps are required on one side of the detector.
  - This side is dictated by the 0nXn requirement.
- This cut clearly separates out the hadronic multiplicity distribution from the photo-nuclear.





# Photo-nuclear Event Topologies

- Photo-nuclear events are **highly asymmetric** and have a **steeply falling multiplicity** distribution.
  - The  $dN/d\eta$  distribution has not been studied before.
  - None of the MC generators shown can perfectly capture these distributions.





# Photo-nuclear Event Topologies

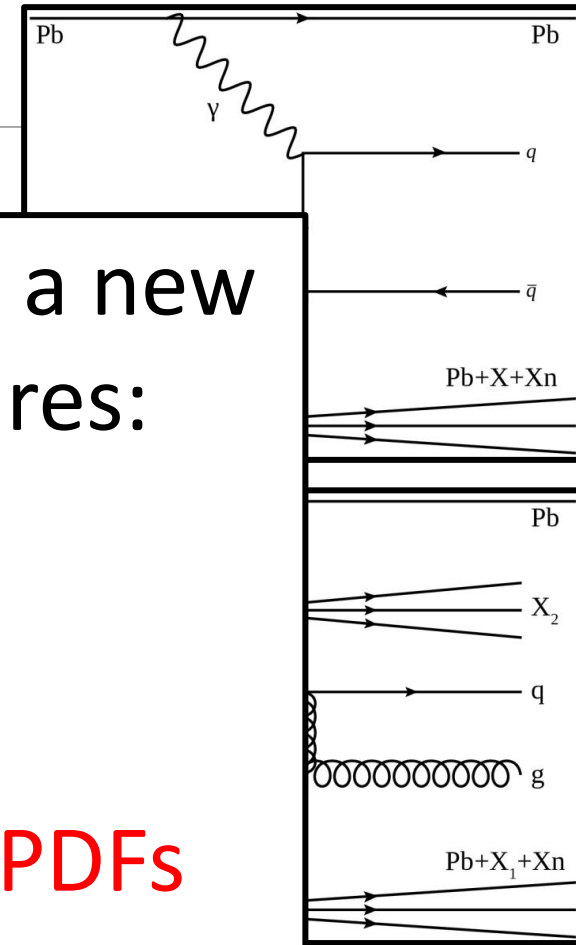
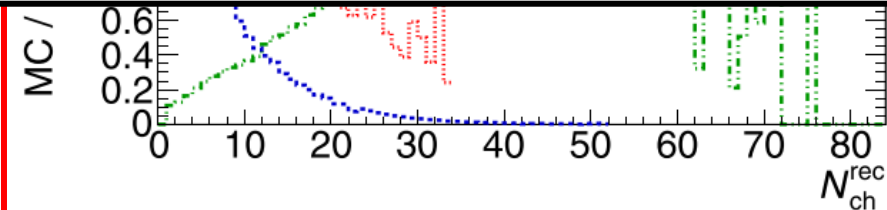
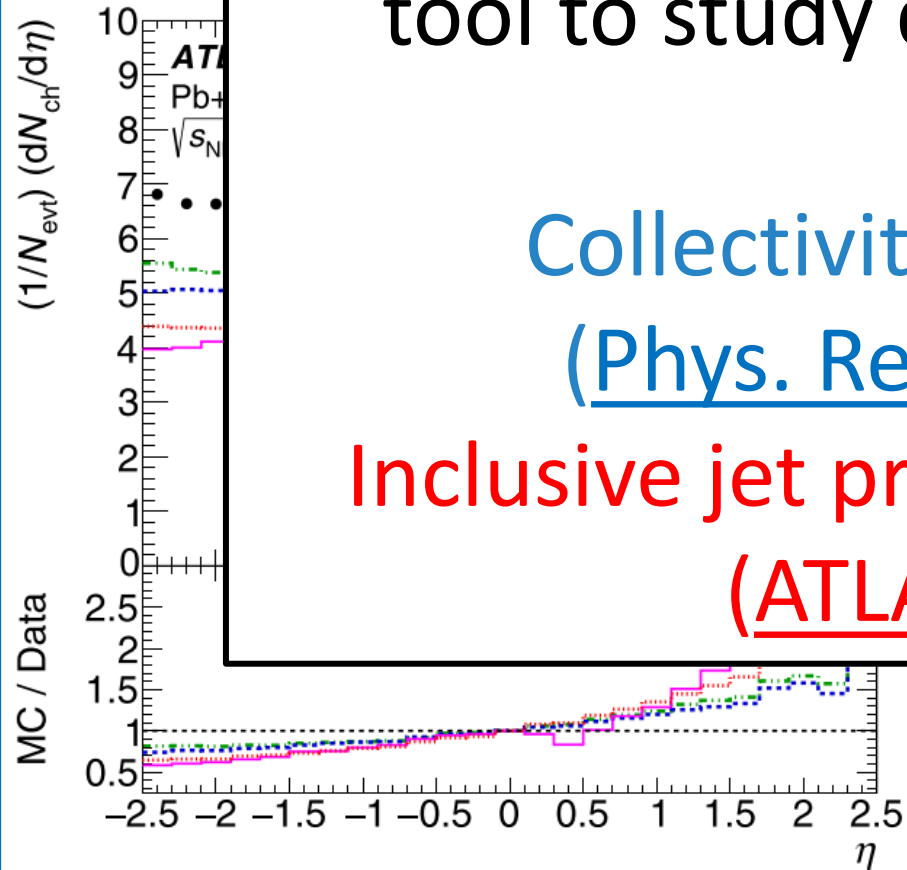
- Photo-nuclear events are **highly asymmetric** and have a **steeply falling multiplicity** distribution.

- The d
- None

Photo-nuclear processes in UPCs provide a new tool to study different physical signatures:

Collectivity in a novel small system  
([Phys. Rev. C 104 \(2021\) 014903](#))

Inclusive jet production as a probe of nPDFs  
([ATLAS-CONF-2022-021](#))



# Constructing Two-Particle Correlations

---

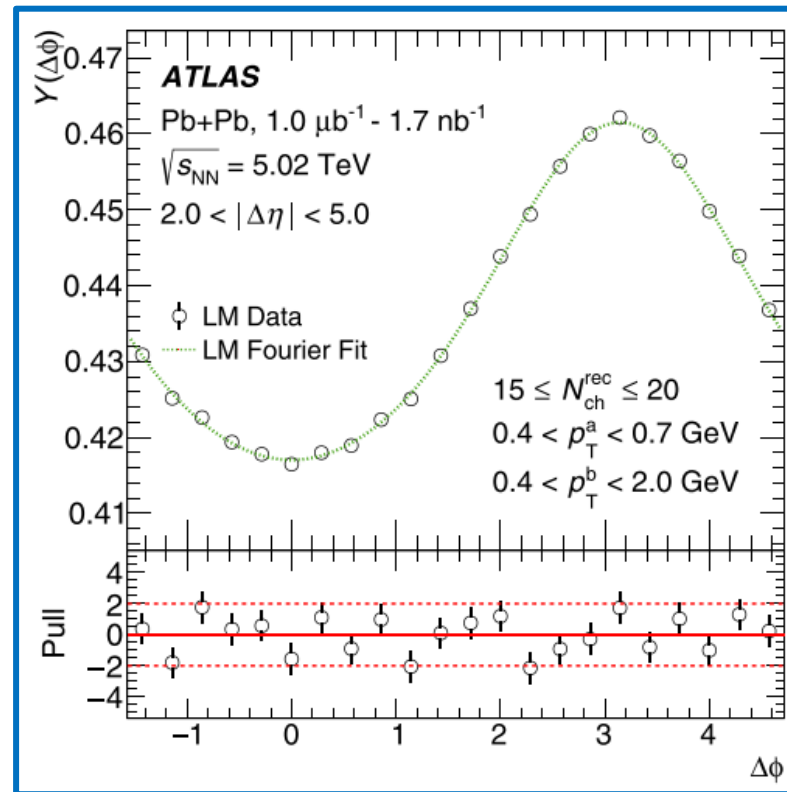
$$Y(\Delta\phi, \Delta\eta) = \frac{1}{N_a} \frac{d^2 N_{\text{pair}}}{d\Delta\phi d\Delta\eta}$$

- Yields are constructed as a function of  $\Delta\eta$  and  $\Delta\phi$  for charged particle tracks.

# Constructing Two-Particle Correlations

$$Y(\Delta\phi, \Delta\eta) = \frac{1}{N_a} \frac{d^2 N_{\text{pair}}}{d\Delta\phi d\Delta\eta} \quad \longrightarrow \quad Y(\Delta\phi) = \int_{|\Delta\eta|=2.0}^{|\Delta\eta|=5.0} Y(\Delta\phi, |\Delta\eta|) d|\Delta\eta|$$

- Yields are constructed as a function of  $\Delta\eta$  and  $\Delta\phi$  for charged particle tracks.
- The yields in  $\Delta\phi$  are determined by integrating over pairs with large  $|\Delta\eta|$ .
- A template for the non-flow subtraction is derived from [low-multiplicity \(LM\) photo-nuclear data](#).



# Constructing Two-Particle Correlations

$$Y(\Delta\phi, \Delta\eta) = \frac{1}{N_a} \frac{d^2 N_{\text{pair}}}{d\Delta\phi d\Delta\eta}$$



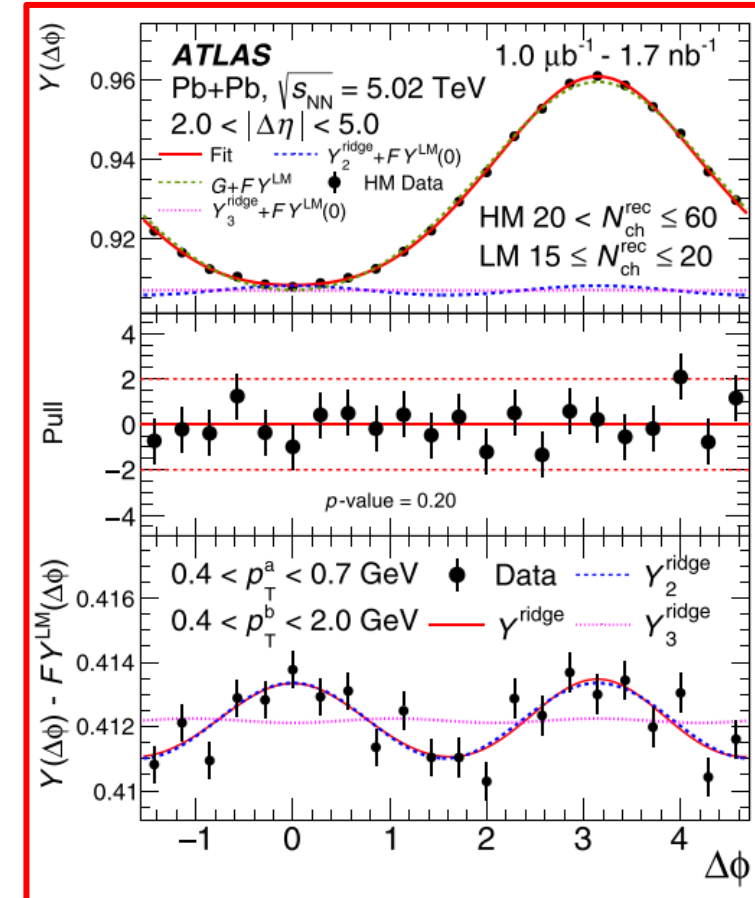
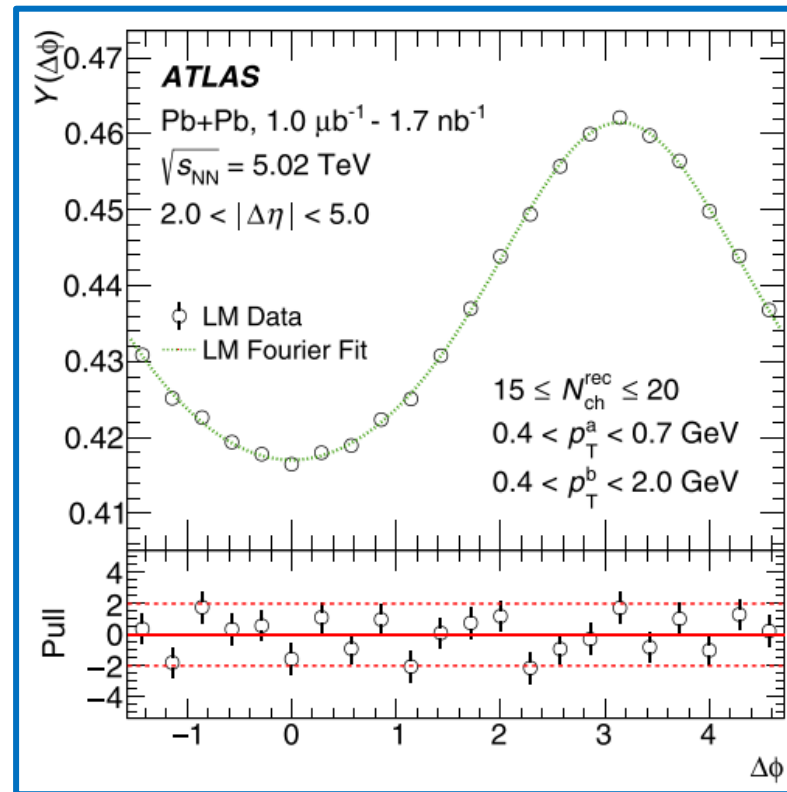
$$Y(\Delta\phi) = \int_{|\Delta\eta|=2.0}^{|\Delta\eta|=5.0} Y(\Delta\phi, |\Delta\eta|) d|\Delta\eta|$$



$$Y^{\text{HM}}(\Delta\phi) = FY^{\text{LM}}(\Delta\phi) + G \left\{ 1 + 2 \sum_{n=2}^4 v_{n,n} \cos(n\Delta\phi) \right\}$$

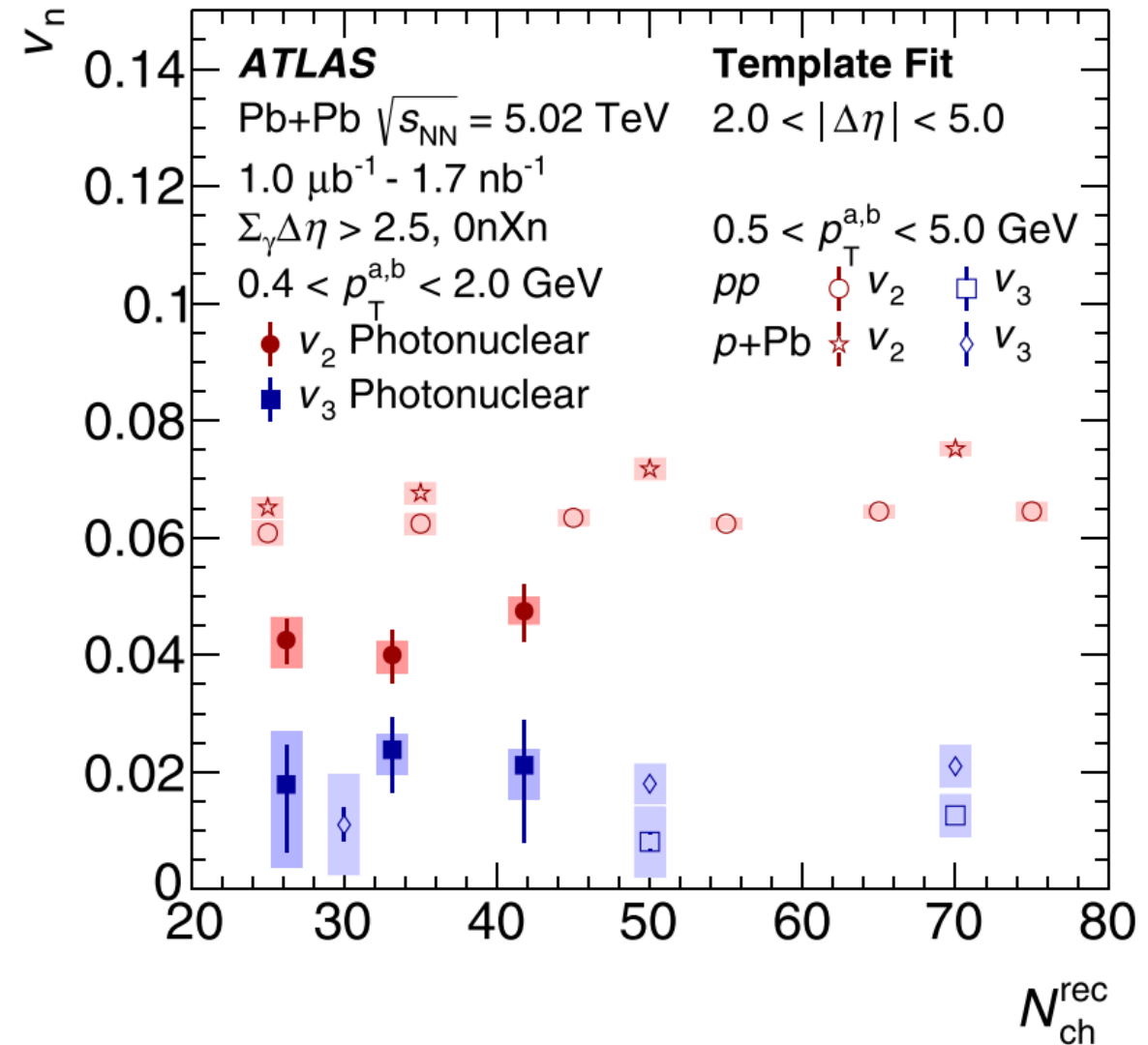
$$= FY^{\text{LM}}(\Delta\phi) + Y^{\text{ridge}}(\Delta\phi)$$

- Yields are constructed as a function of  $\Delta\eta$  and  $\Delta\phi$  for charged particle tracks.
- The yields in  $\Delta\phi$  are determined by integrating over pairs with large  $|\Delta\eta|$ .
- A template for the non-flow subtraction is derived from **low-multiplicity (LM) photo-nuclear data**.
- Template fits to **high-multiplicity (HM) data** allow us to extract the yield not explained by LM non-flow contribution.



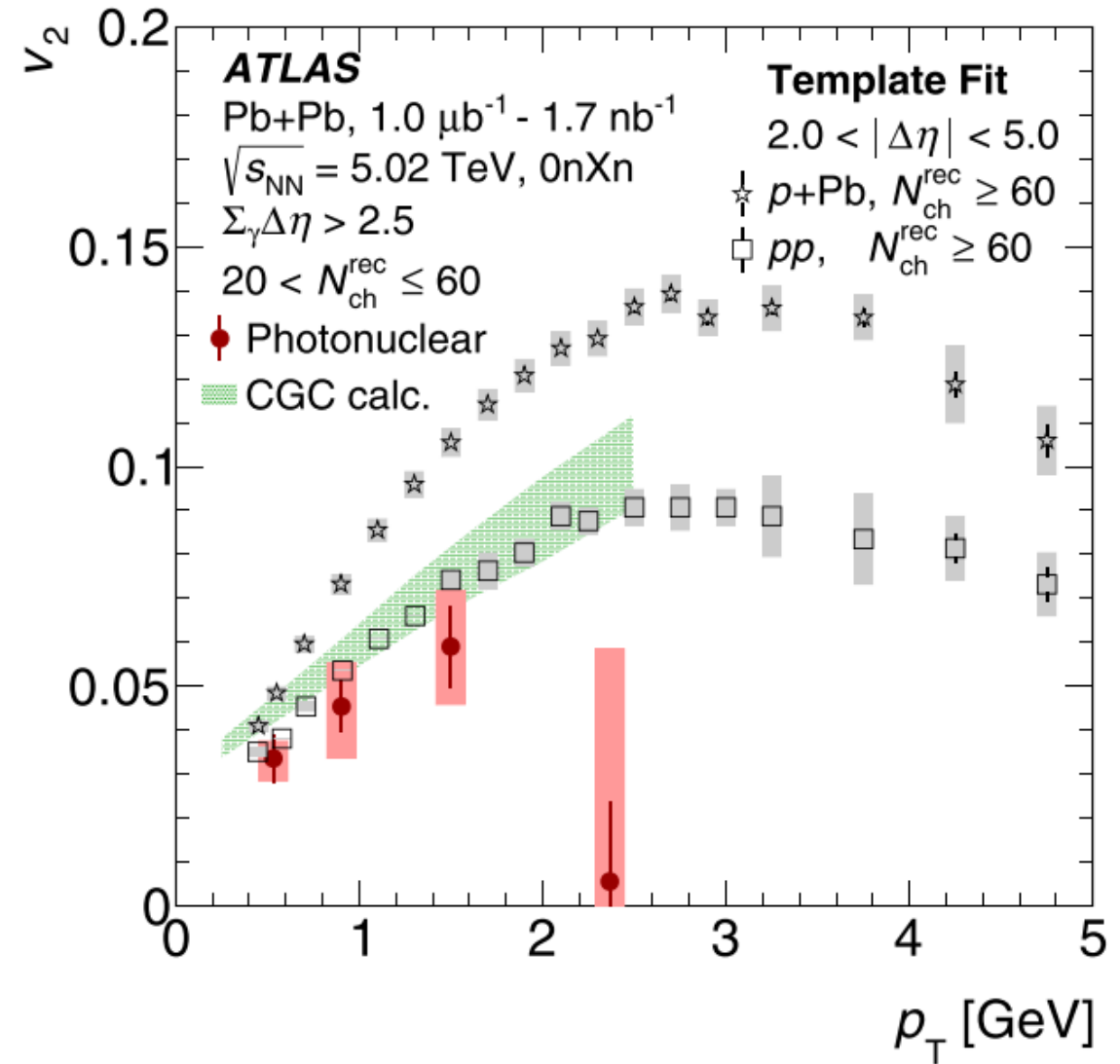
# Multiplicity-Dependence of Correlations

- Applying this procedure for extracting the flow coefficients allows them to be compared to  $pp$  and  $pPb$  with similar multiplicity.
  - The flow signal in photonuclear collisions is **smaller than in  $pp$**  but **clearly significant**.



# $p_T$ -Dependence of Correlations

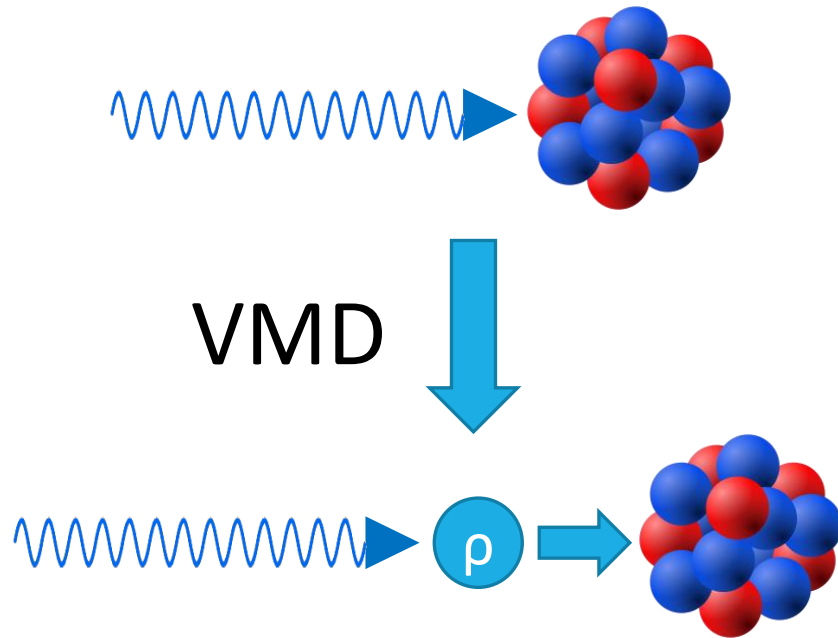
- Applying this procedure for extracting the flow coefficients allows them to be compared to  $pp$  and  $pPb$  with similar multiplicity.
  - The flow signal in photonuclear collisions is **smaller than in  $pp$**  but **clearly significant**.
- The  $v_2$  also demonstrates a  $p_T$ -dependence characteristic of hydrodynamic expansion.



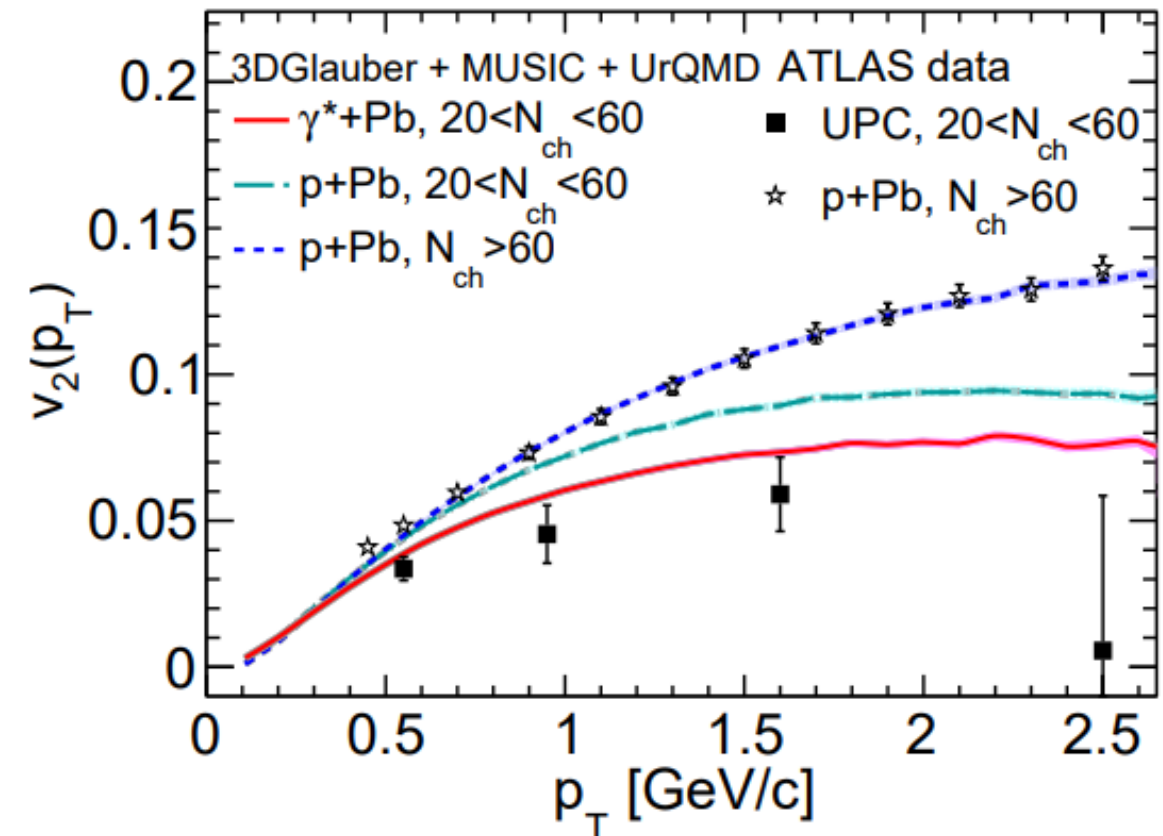


# Interpreting UPC Correlations

- In the Vector Meson Dominance (VMD) model, the resolved photon can be treated as a rho-nucleus collision.
- This model with (3+1)D hydrodynamic evolution (right) can effectively reproduce the ATLAS data for  $v_2(p_T)$ .



Zhao, Shen, and Schenke,  
[nucl-th/2203.06094](https://arxiv.org/abs/nucl-th/2203.06094)



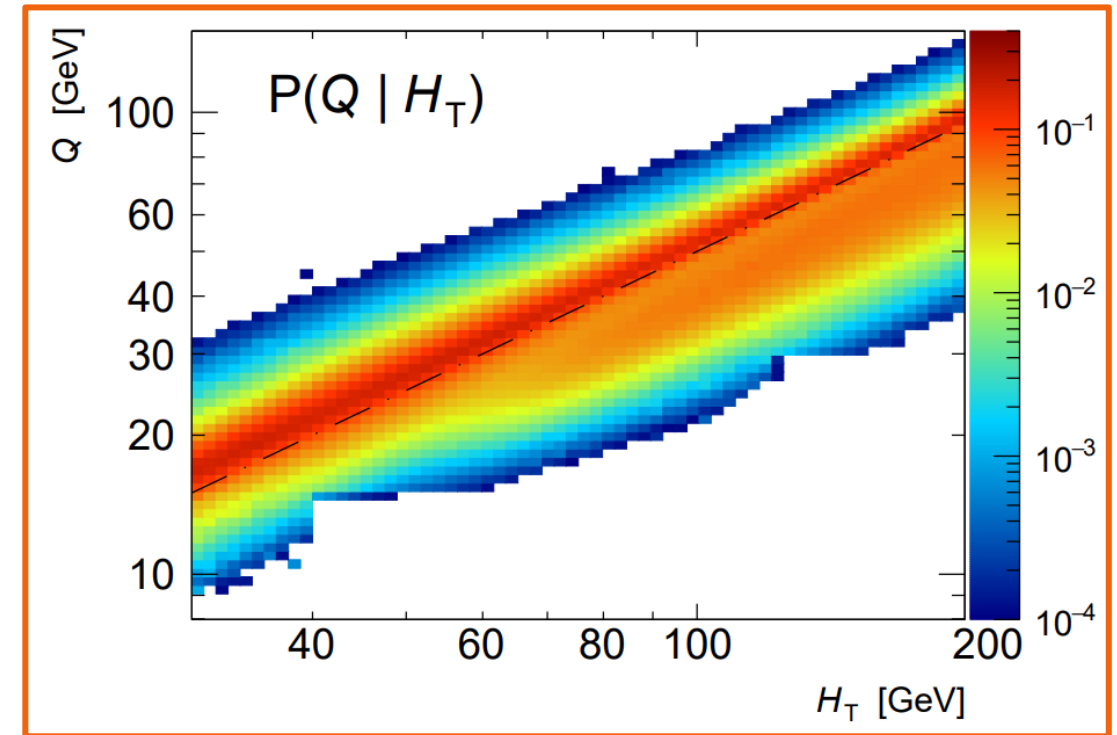
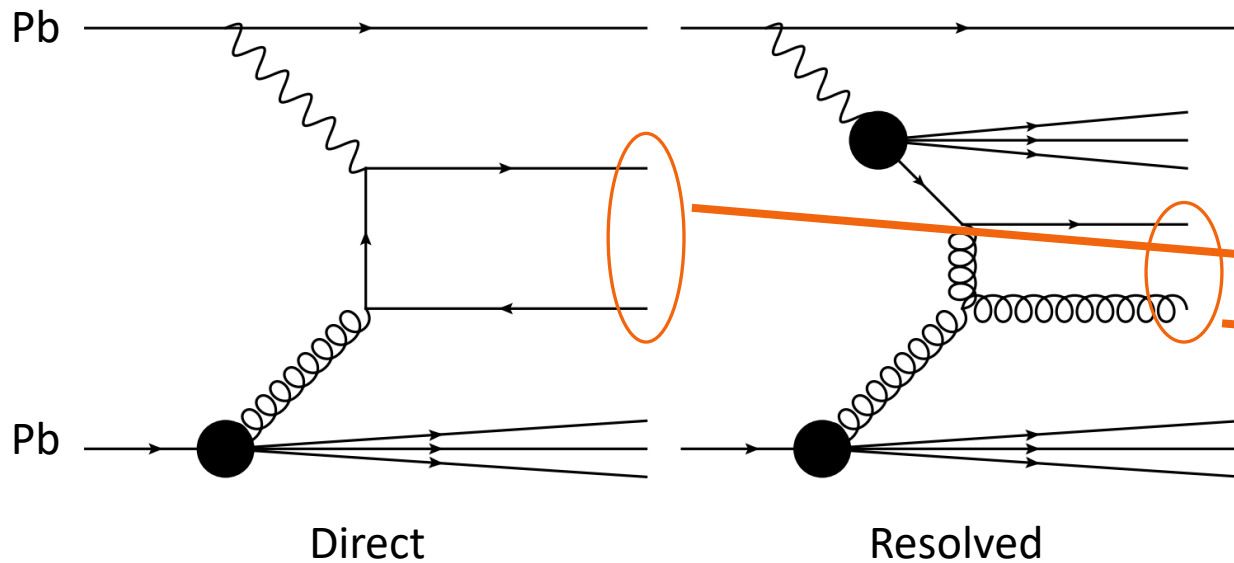
# Photo-nuclear Dijet Events

- If we add a jet selection, we can reconstruct the hard-scattering kinematics.
  - We add a requirement for 2 offline Particle-Flow jets.
  - Jets are required to have  $p_T > 15$  GeV.

$$H_T \equiv \sum_i p_T^i$$

$$x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}}$$

$$z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



Jet kinematics provide access to the hard-scattering kinematics, directly probing nuclear PDF effects.

$$H_T \sim 2Q \quad x_A \sim x$$

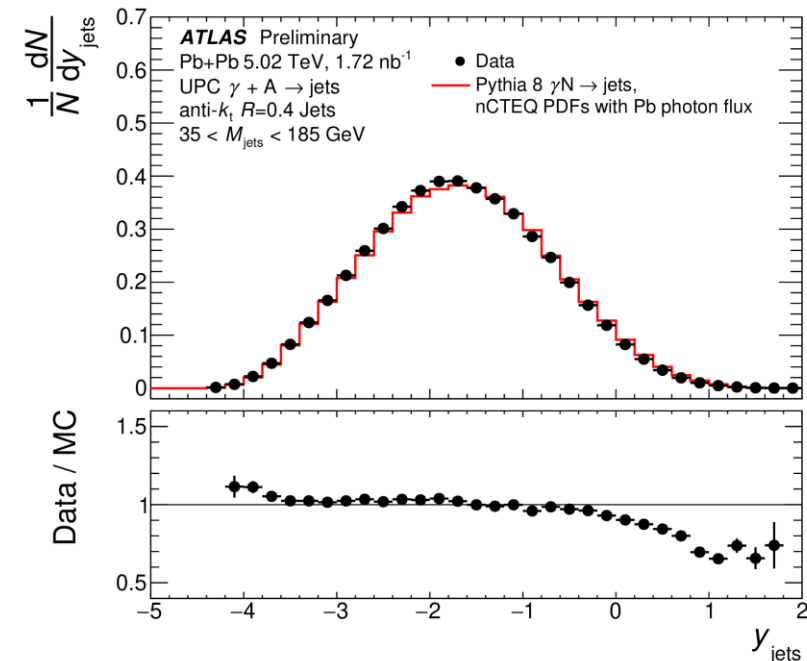
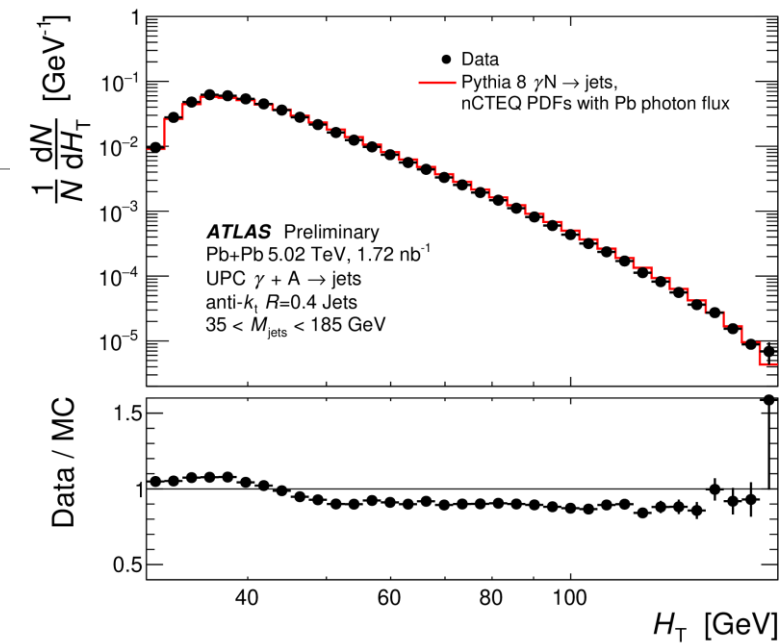
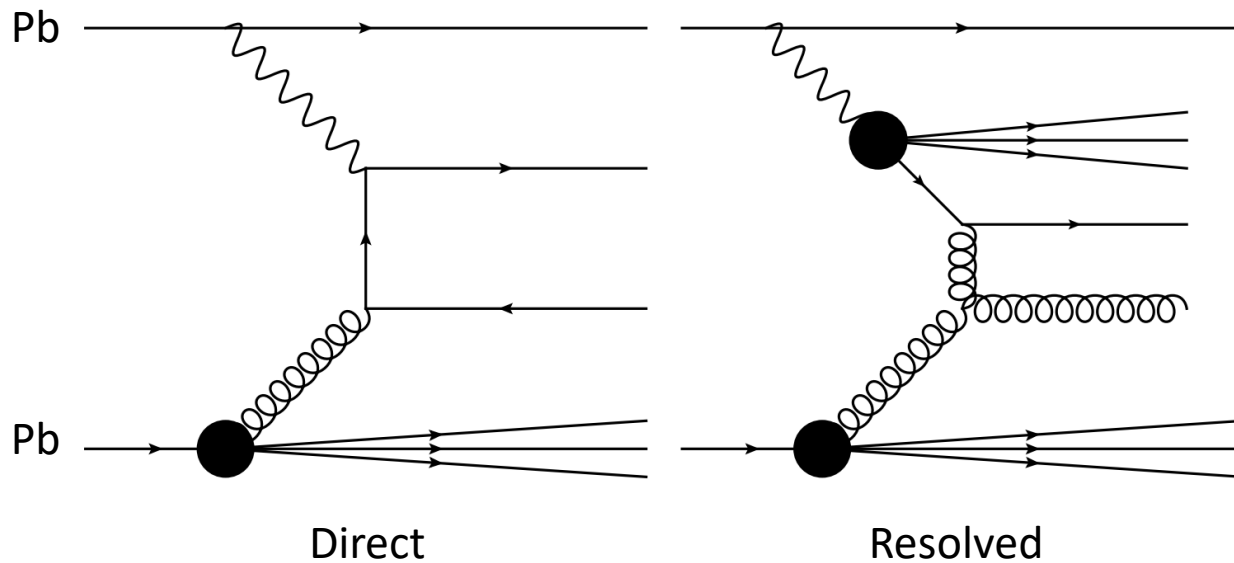
# Photo-nuclear Dijet Events

- If we add a jet selection, we can reconstruct the hard-scattering kinematics.
  - We add a requirement for 2 offline Particle-Flow jets.
  - Jets are required to have  $p_T > 15$  GeV.

$$H_T \equiv \sum_i p_T^i$$

$$x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}}$$

$$z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



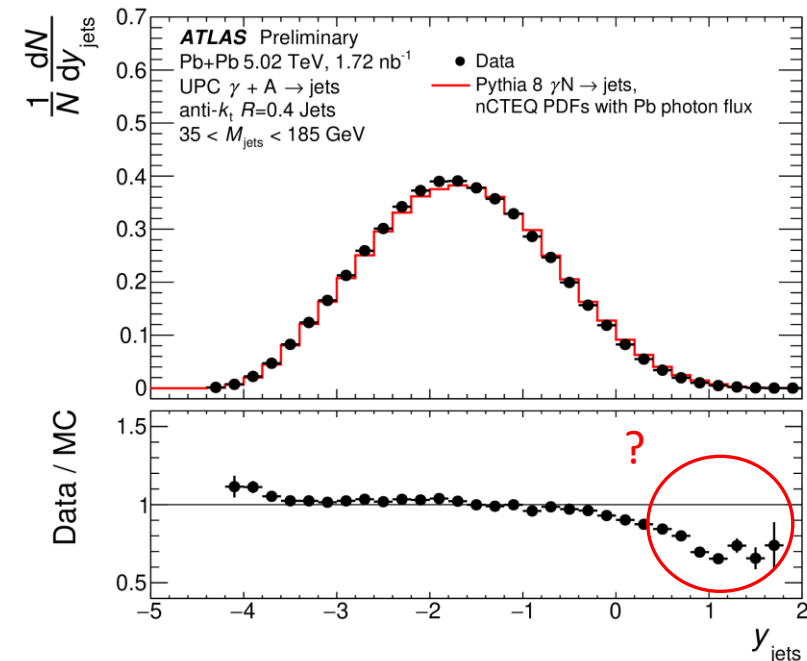
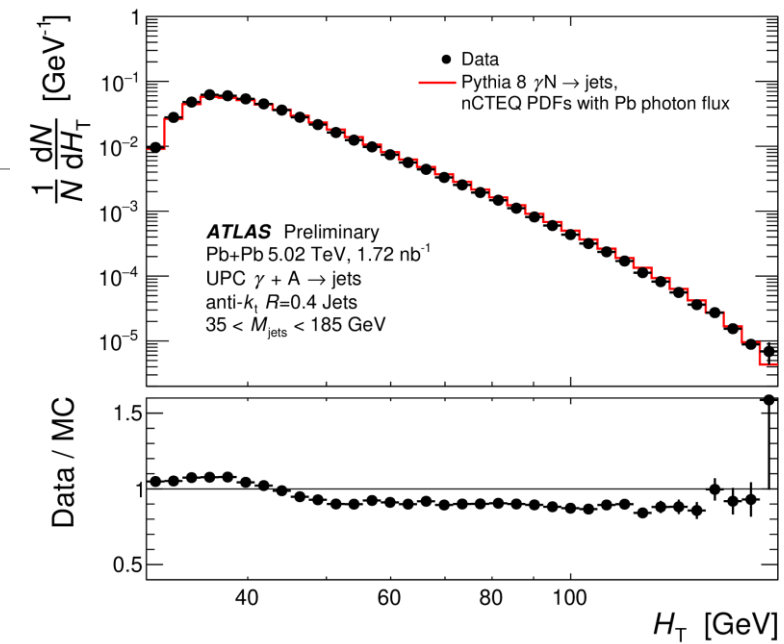
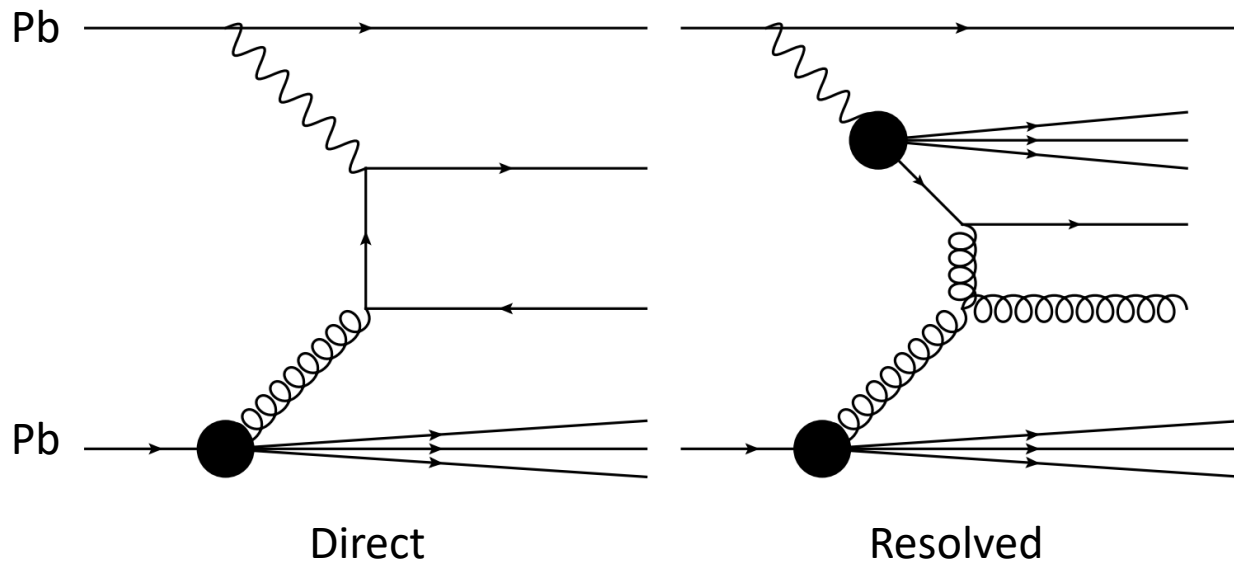
# Photo-nuclear Dijet Events

- If we add a jet selection, we can reconstruct the hard-scattering kinematics.
  - We add a requirement for 2 offline Particle-Flow jets.
  - Jets are required to have  $p_T > 15$  GeV.

$$H_T \equiv \sum_i p_T^i$$

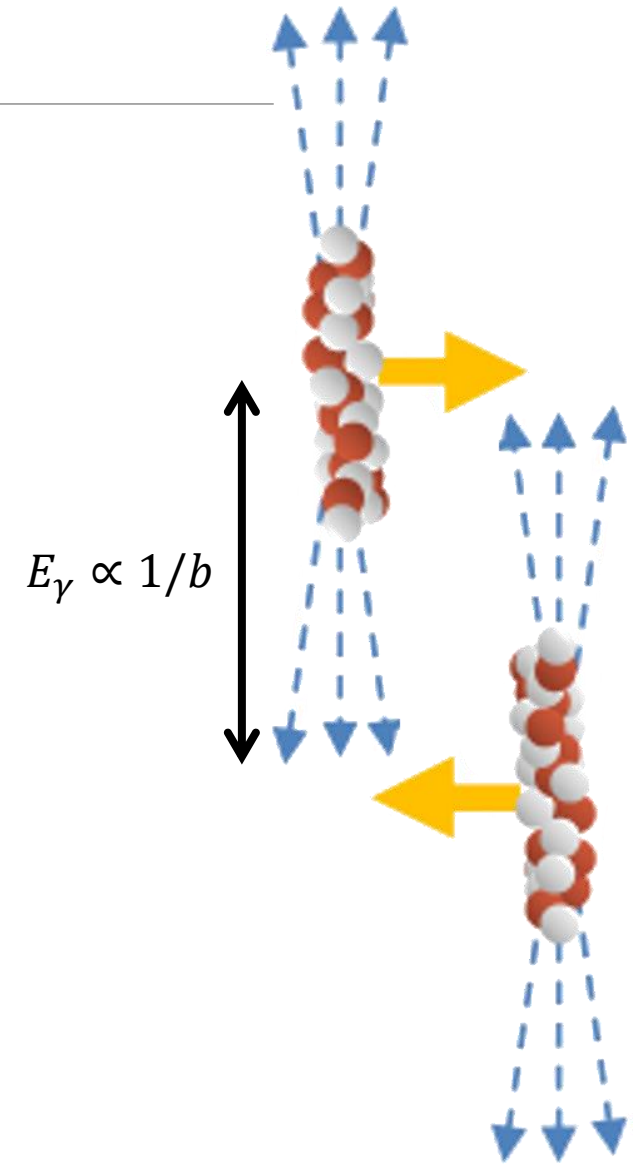
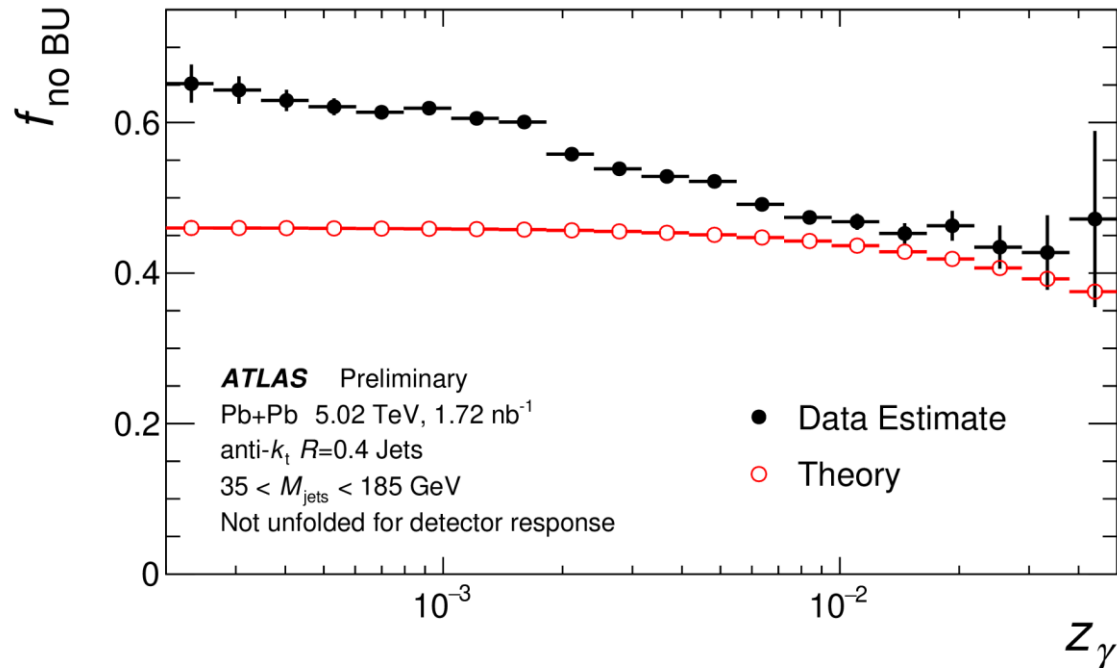
$$x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}}$$

$$z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



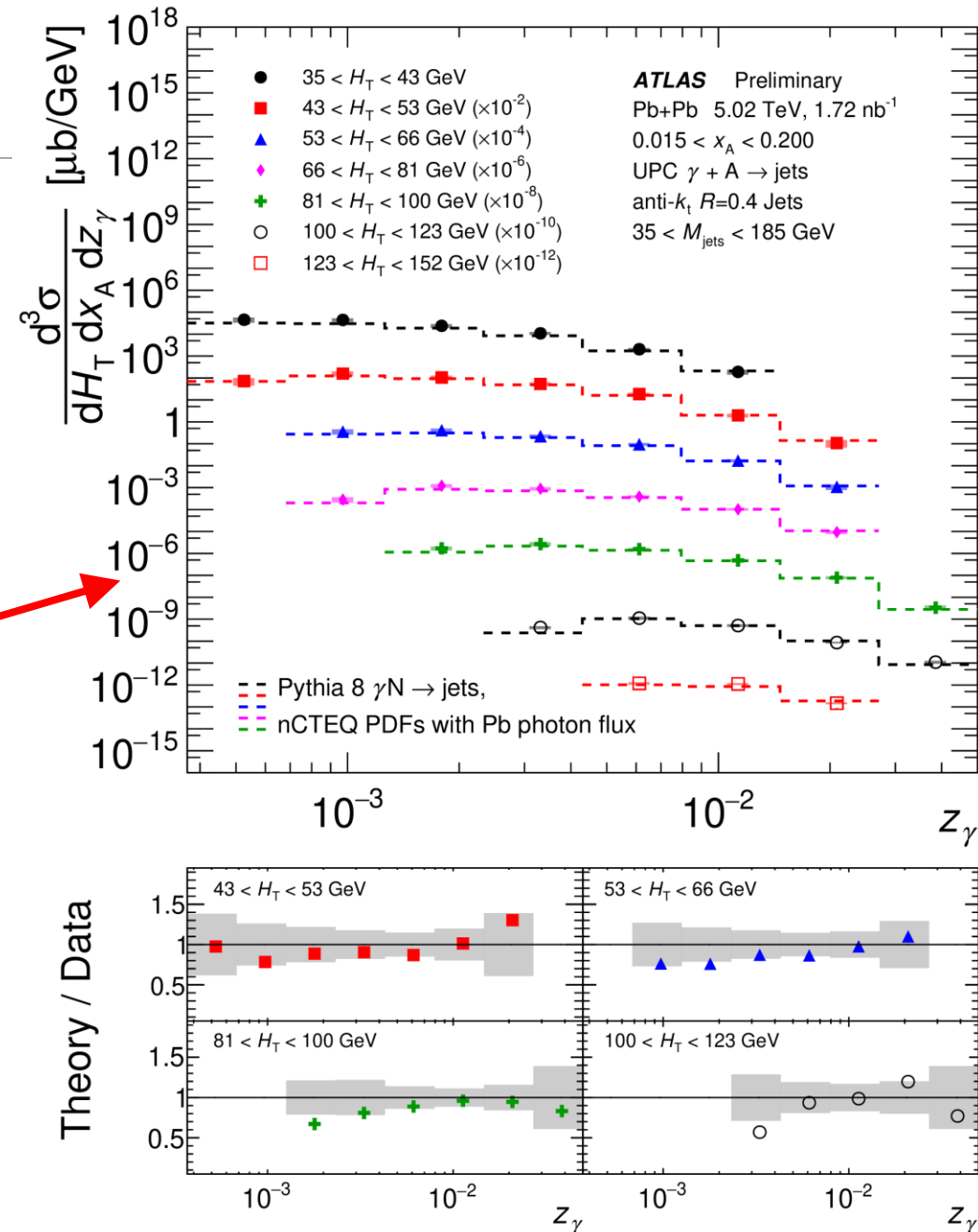
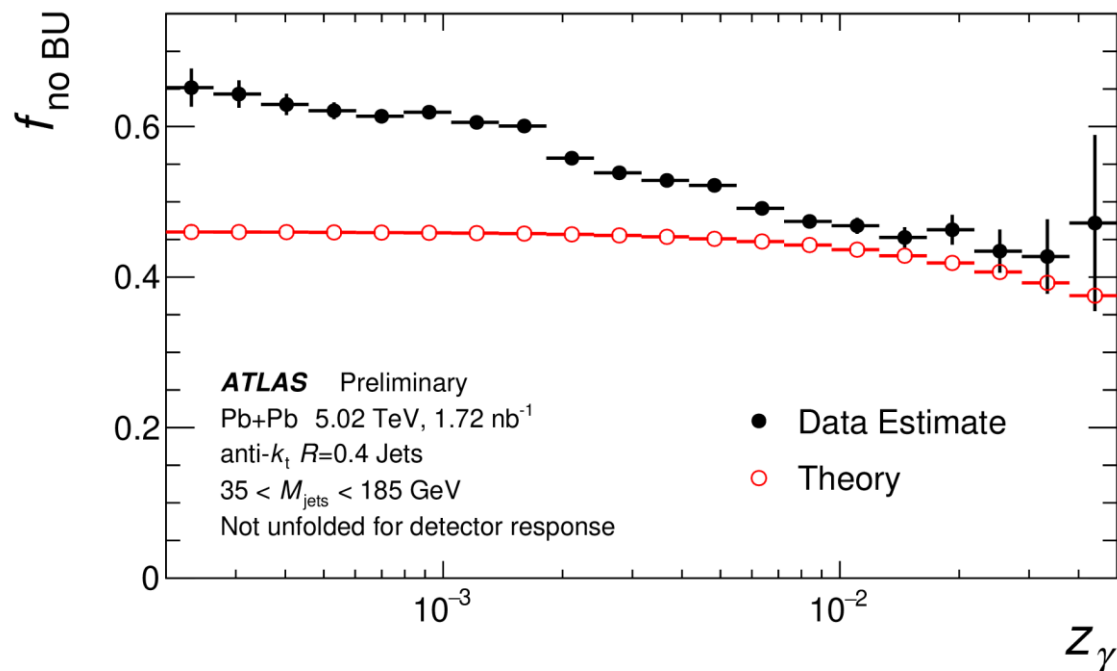
# The Measured Photon Flux

- Breakup of the photon-emitting nucleus is a substantial effect, both in data estimates and theoretical modelling.
  - All model comparisons shown are corrected with the “Theory” points.



# The Measured Photon Flux

- Breakup of the photon-emitting nucleus is a substantial effect, both in data estimates and theoretical modelling.
  - All model comparisons shown are corrected with the “Theory” points.
  - Disagreements appear to arise more at low  $z_\gamma$ , where the breakup model tends to over-correct.



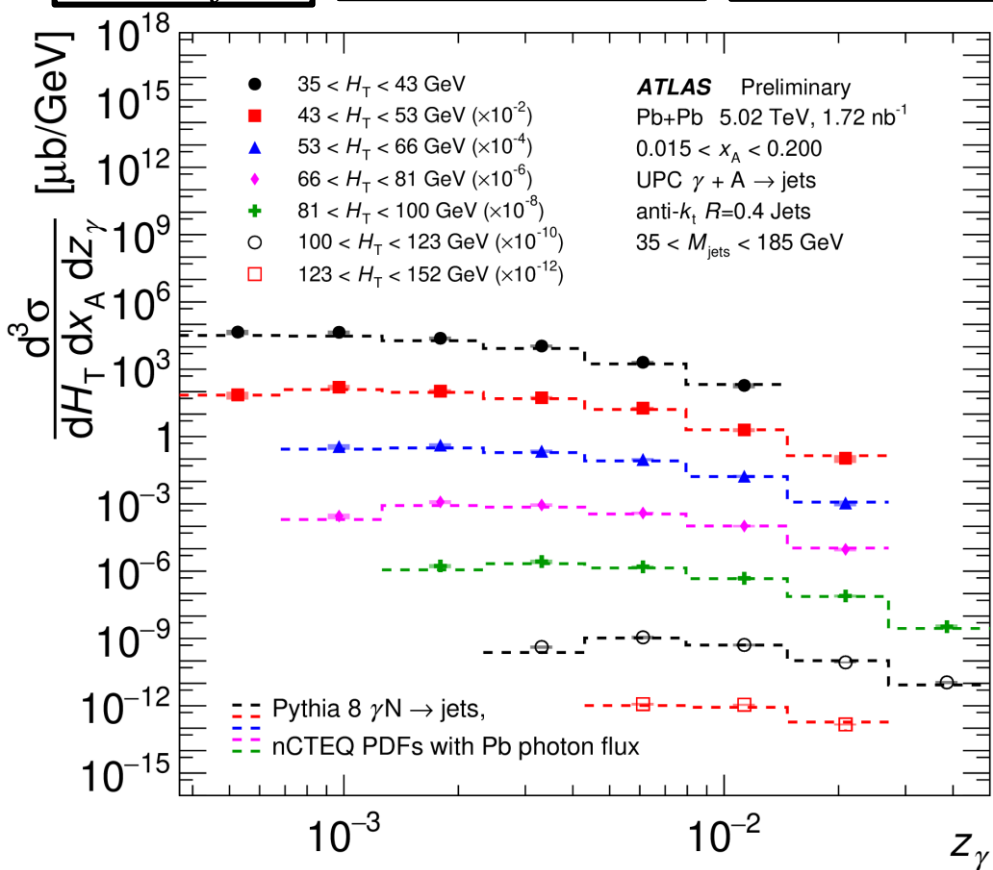


# Scanning in Photon Energy

$$H_T \equiv \sum_i p_T^i$$

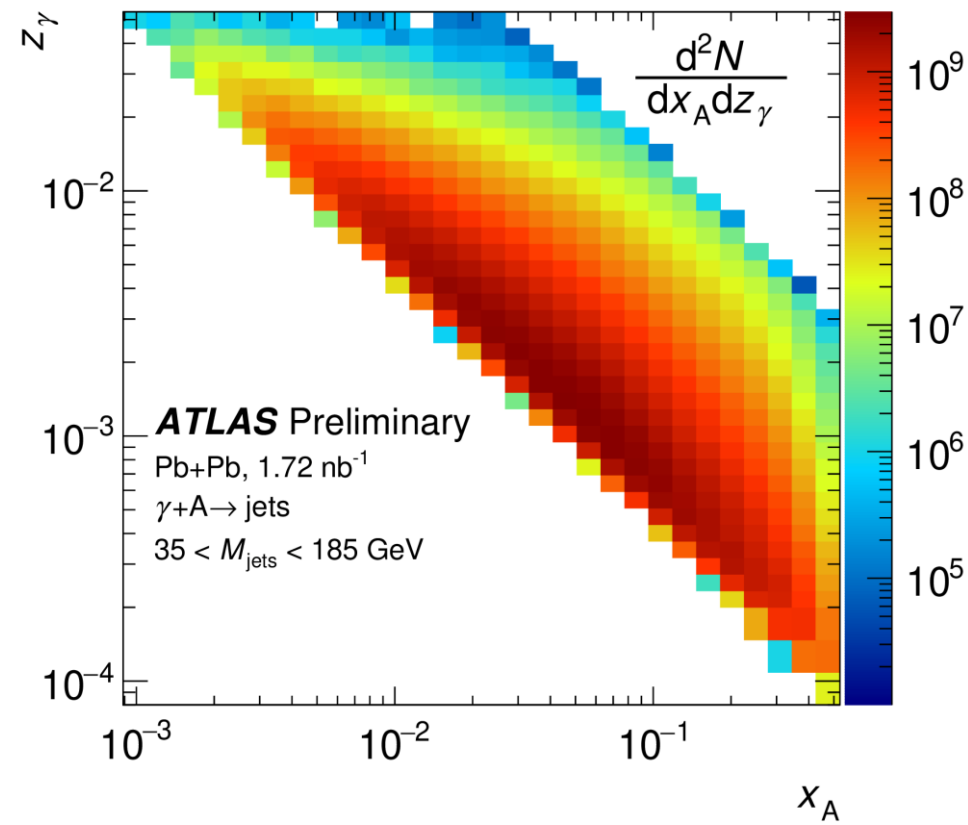
$$x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{S_{NN}}}$$

$$z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{S_{NN}}}$$



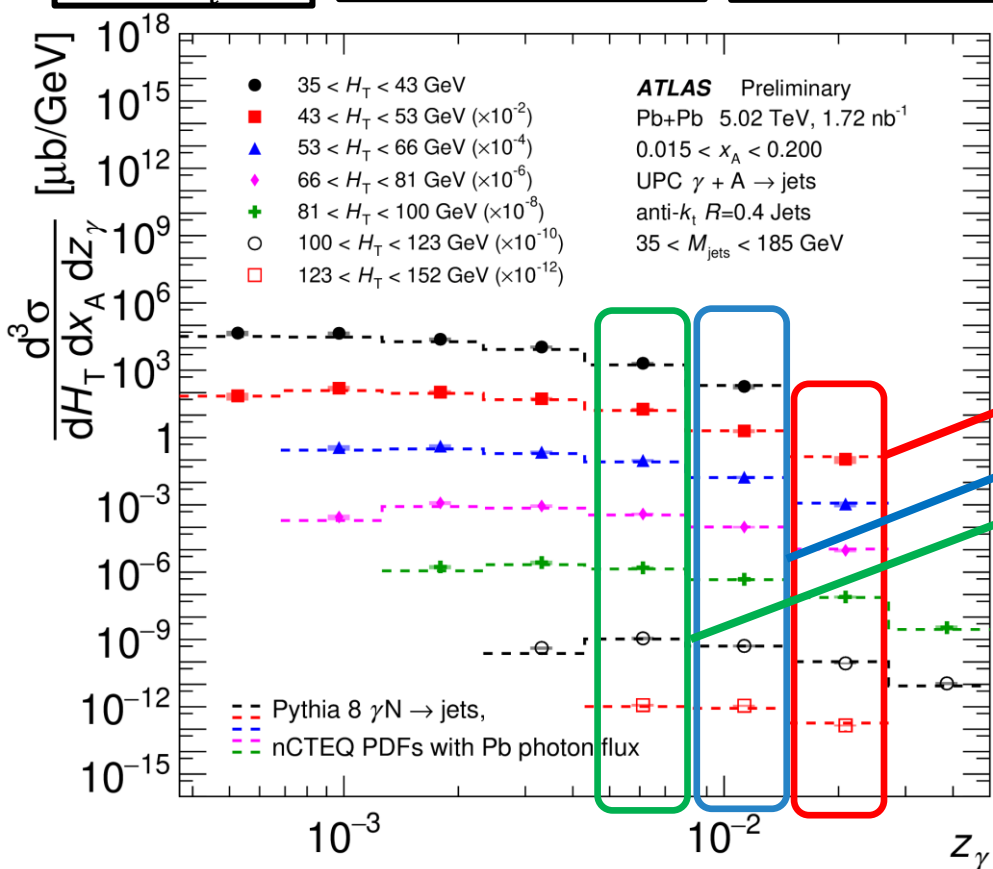
The  $x_A$  distribution has substantial acceptance effects in  $z_\gamma$ .

Selecting a photon energy removes this bias, allowing a more direct measurement of PDF effects.



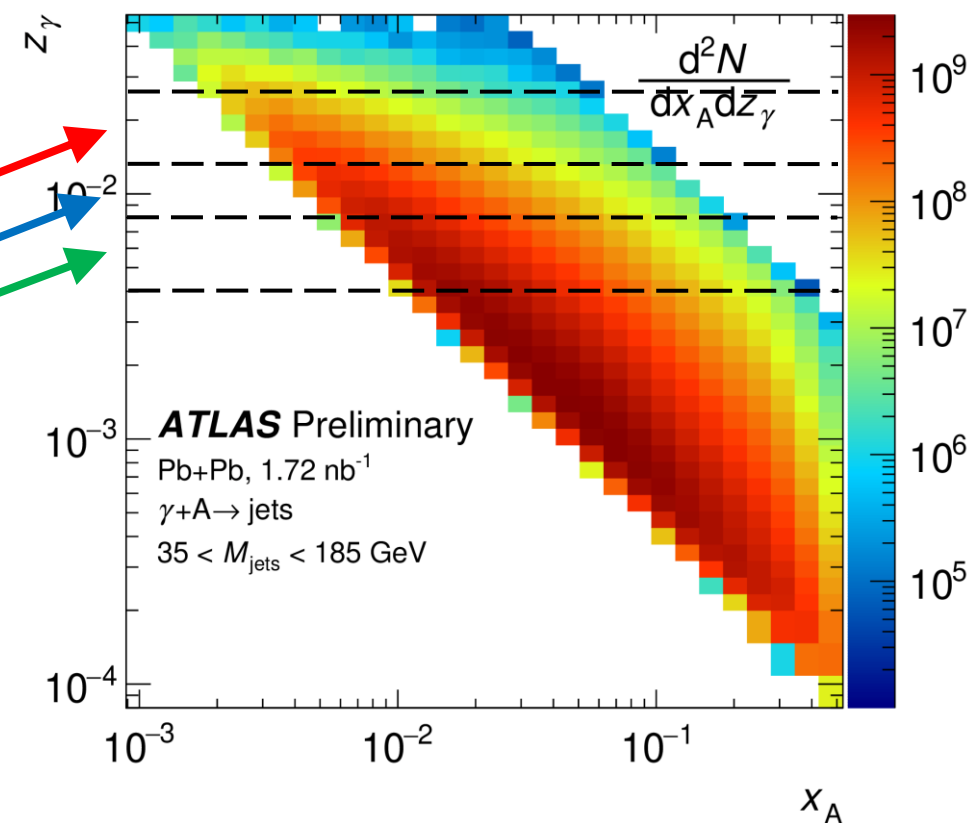
# Scanning in Photon Energy

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



The  $x_A$  distribution has substantial acceptance effects in  $z_\gamma$ .

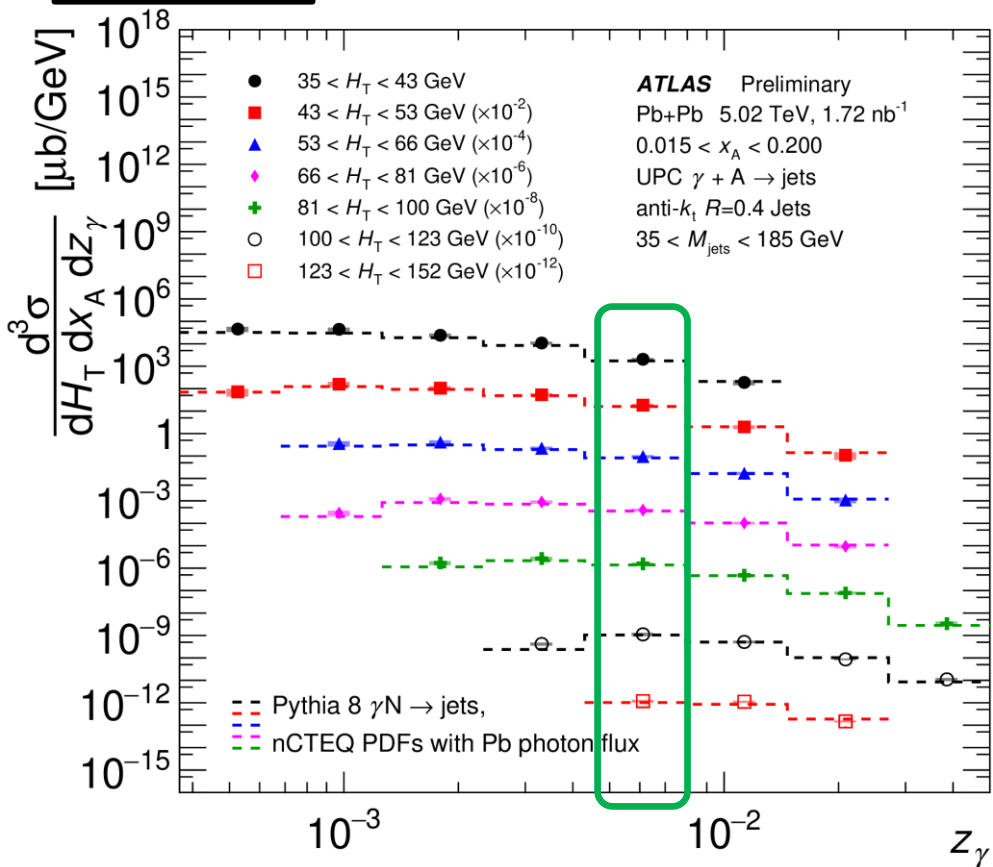
Selecting a photon energy removes this bias, allowing a more direct measurement of PDF effects.



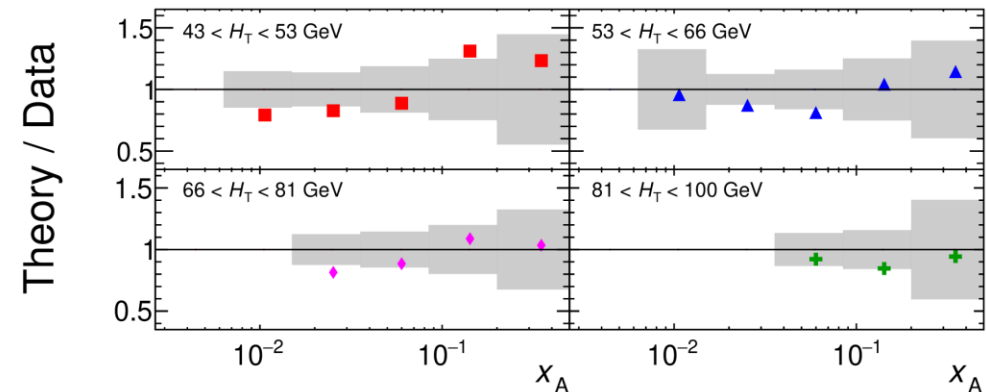
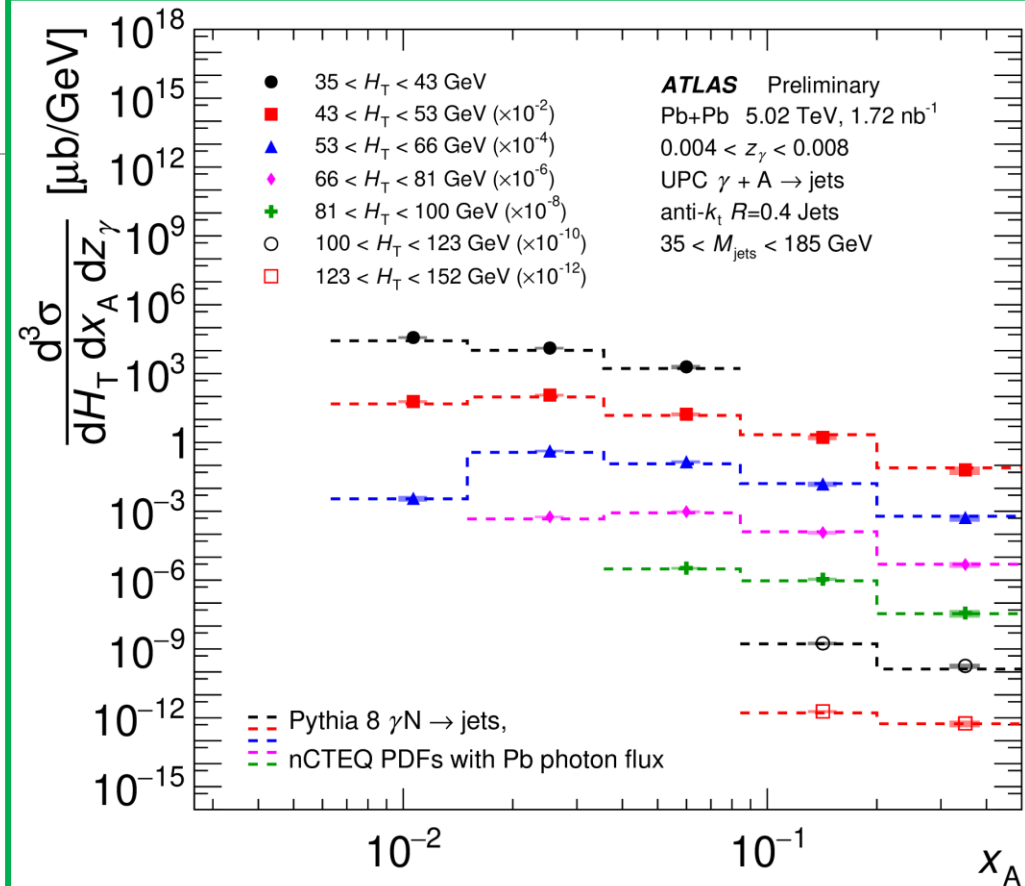
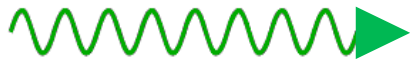
# Measured Cross-Sections

- At intermediate photon energies, we can access higher-x partons.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



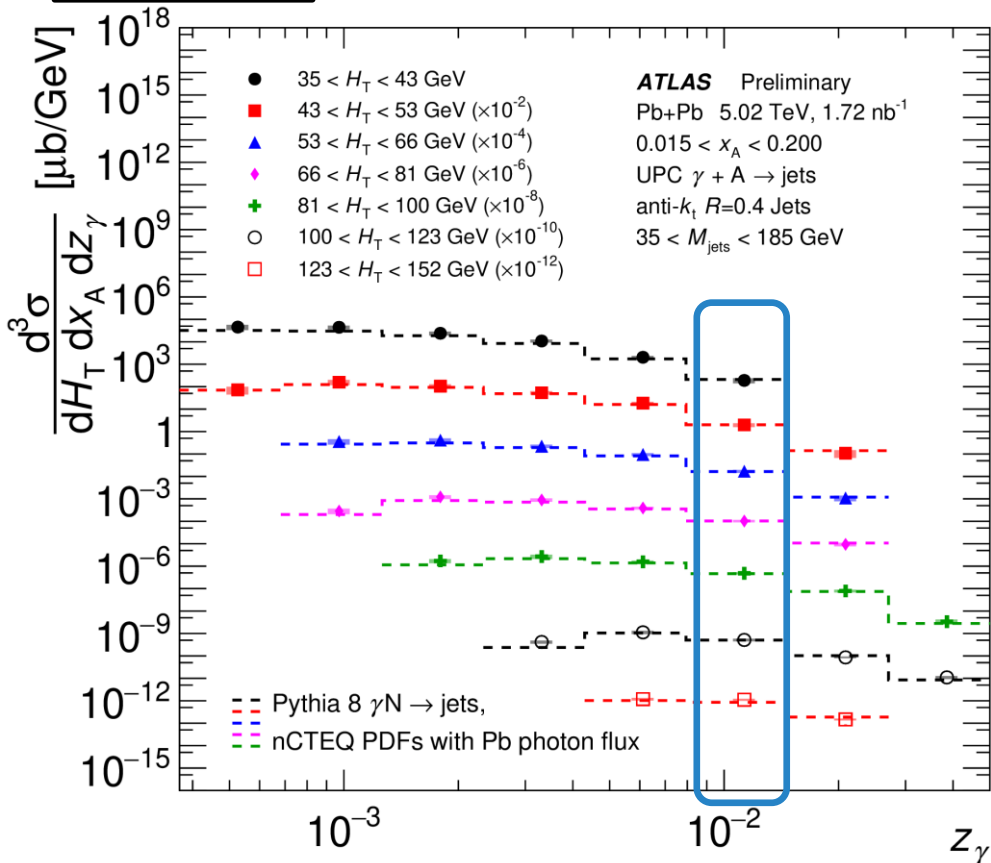
Photon Energy  
0.004 < z<sub>γ</sub> < 0.008



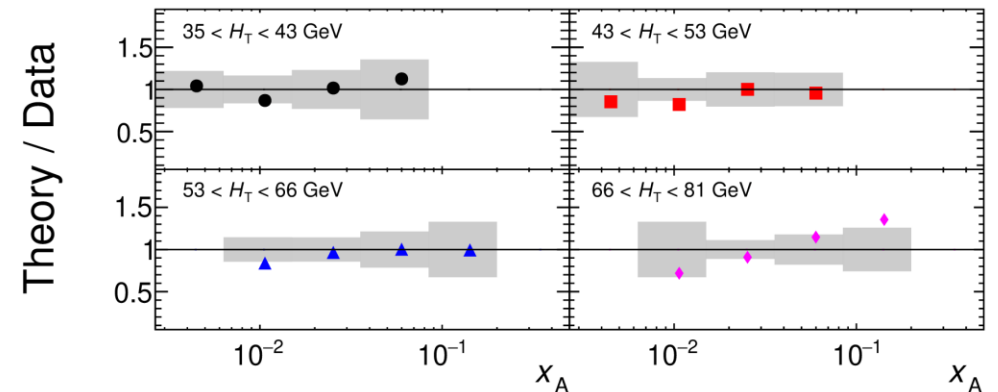
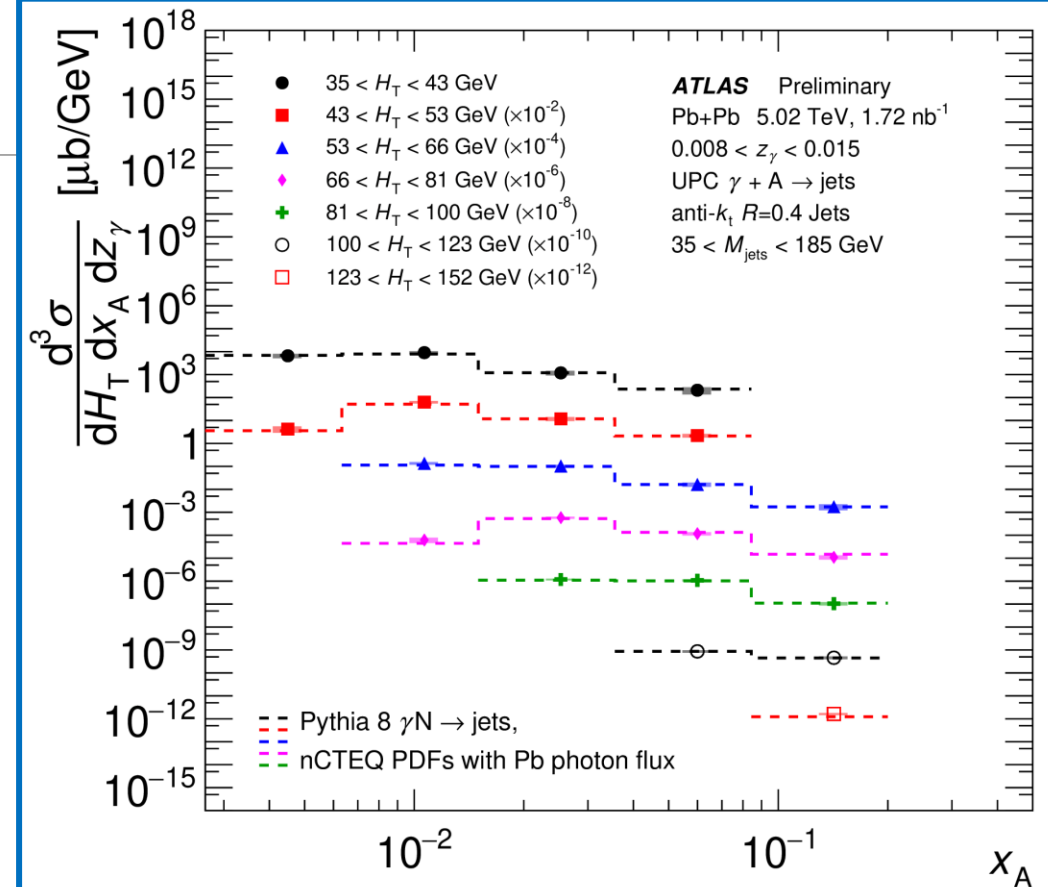
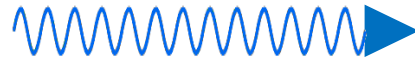
# Measured Cross-Sections

- Going higher in photon energy opens up the low-x shadowing region.
- Results are quite consistent with the theoretical model.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



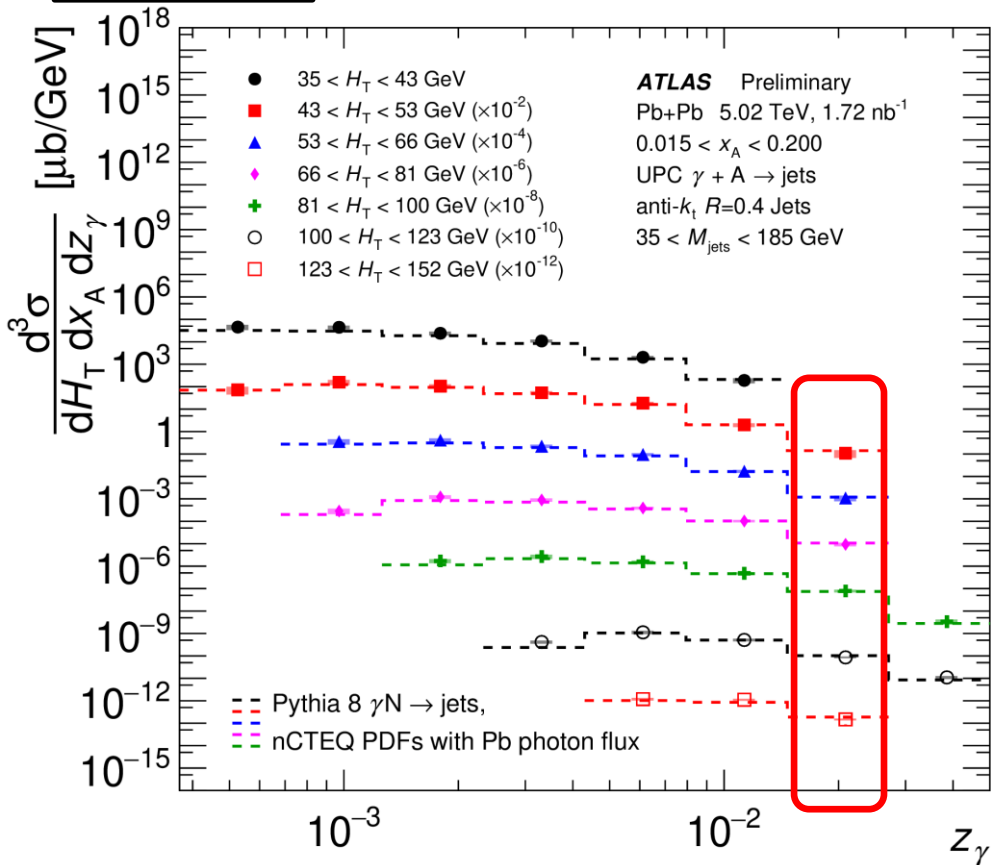
**Photon Energy**  
0.008 < z<sub>γ</sub> < 0.015



# Measured Cross-Sections

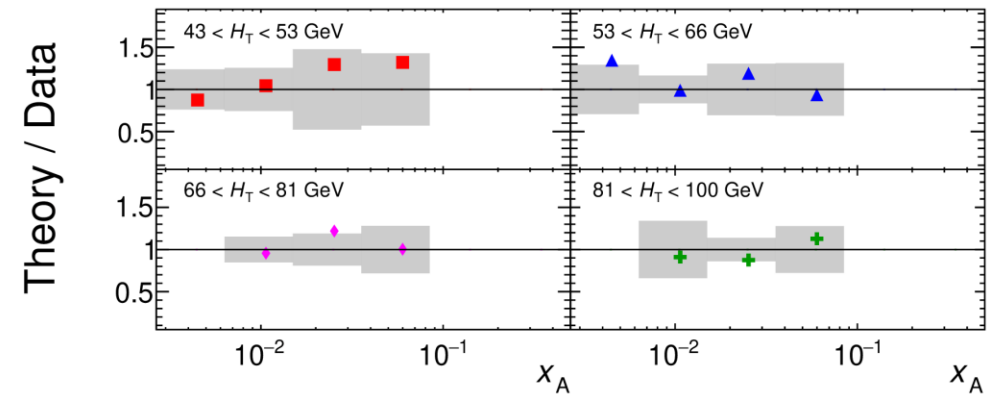
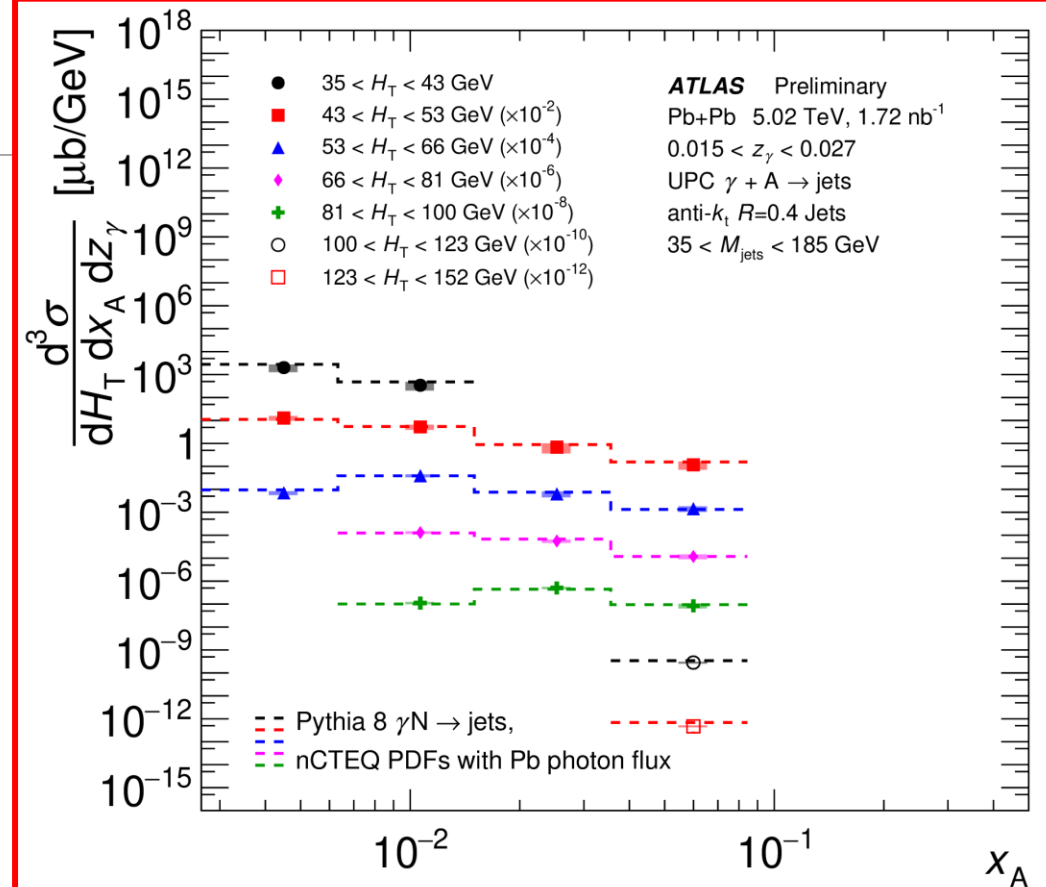
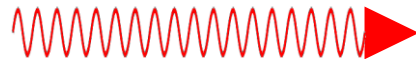
- The highest photon energy allows the most access to low- $x$ .
- Systematic control is more of a challenge near acceptance edges.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



Photon Energy

$0.015 < z_\gamma < 0.027$



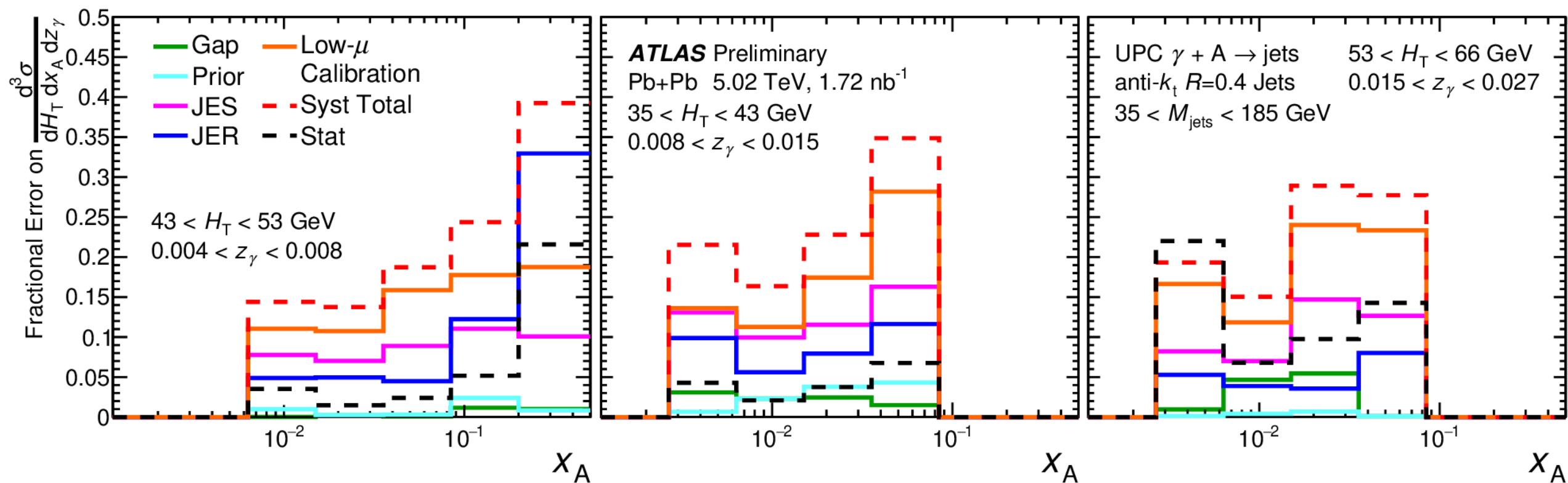
# Systematic Uncertainties on Jet Cross-Sections

Systematics uncertainties are the key limiting factor in our sensitivity to nuclear PDFs.

The jet energy **scale** and **resolution** uncertainties are typically 5-10%.

Control over the **preliminary low- $\mu$  calibration** currently provides the dominant source of uncertainty.

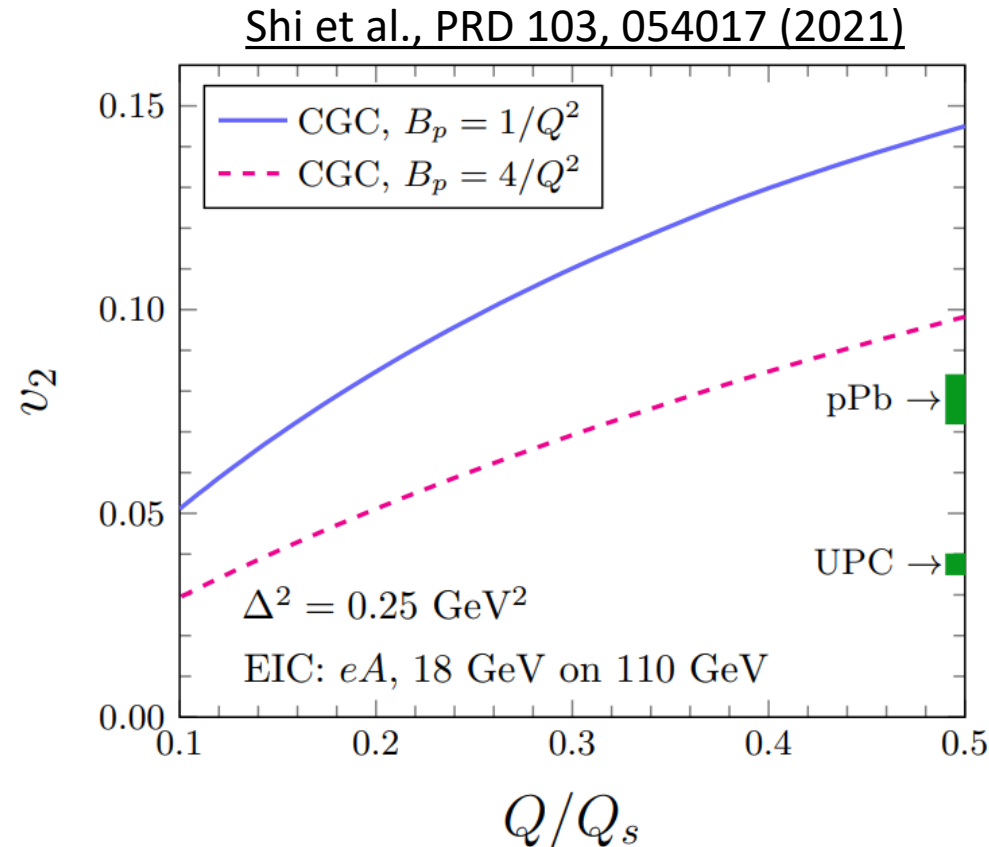
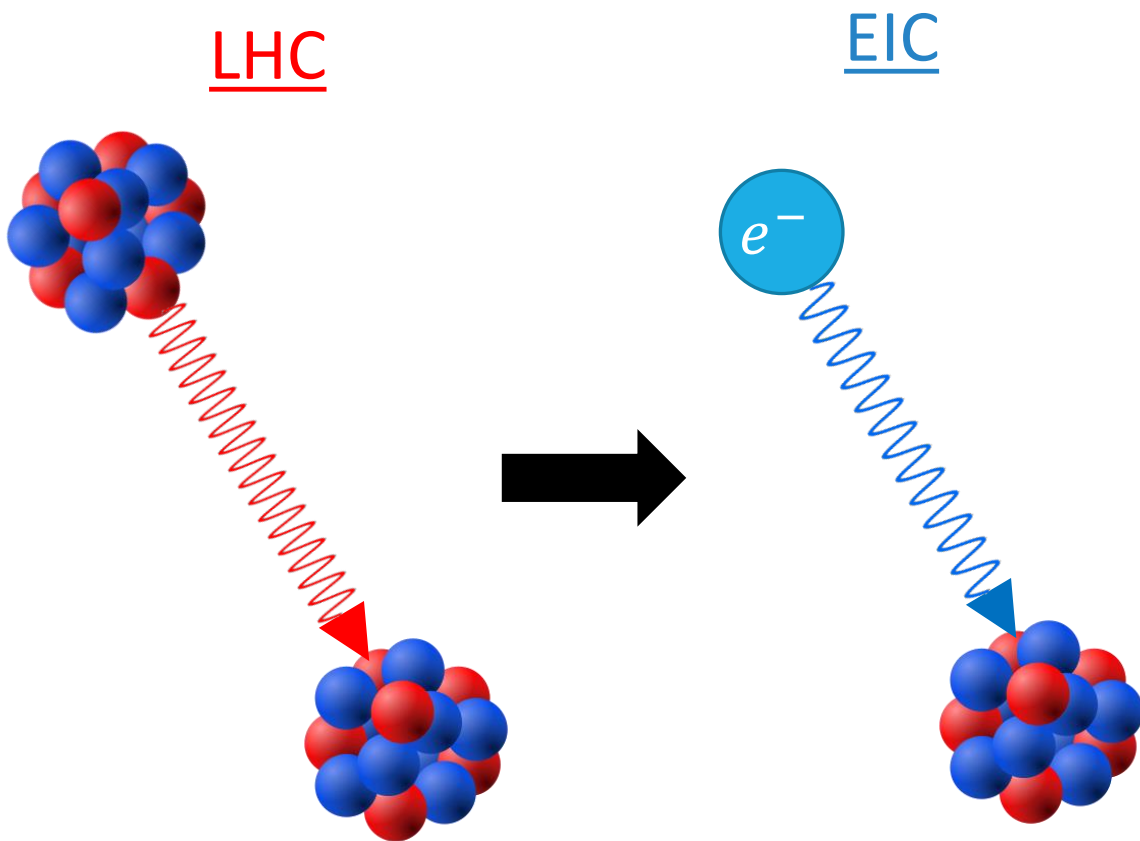
Systematics uncertainties are also evaluated on the **unfolding** and **event selections**.





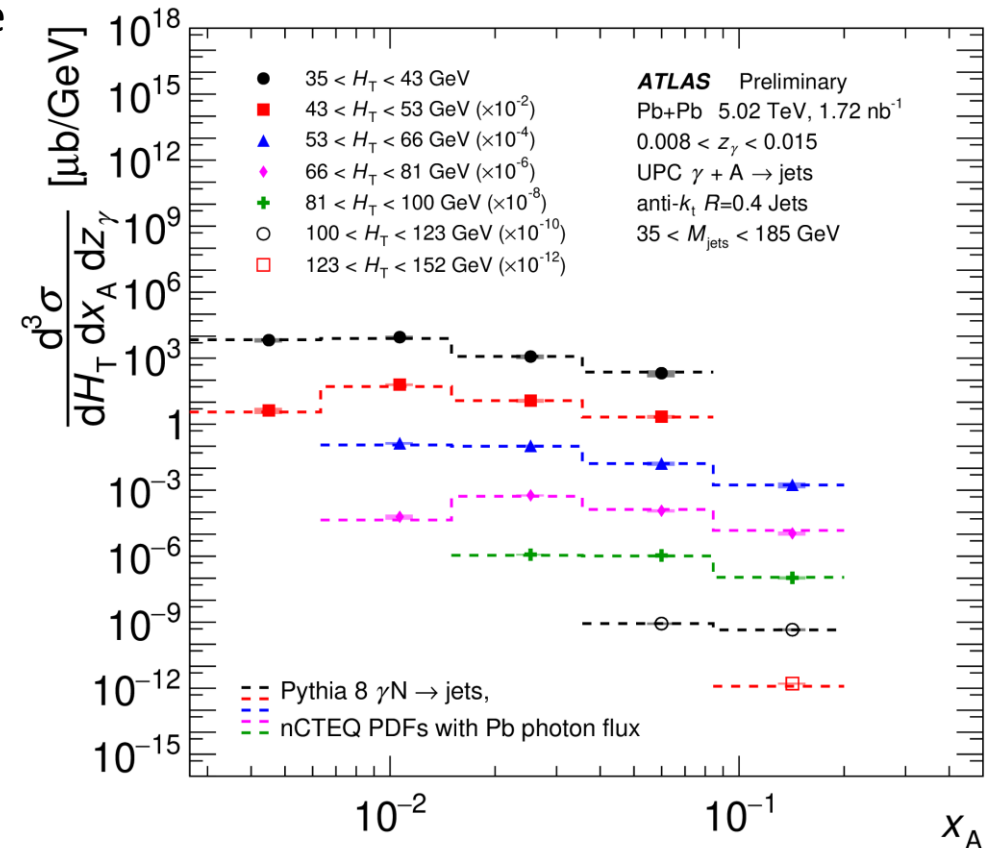
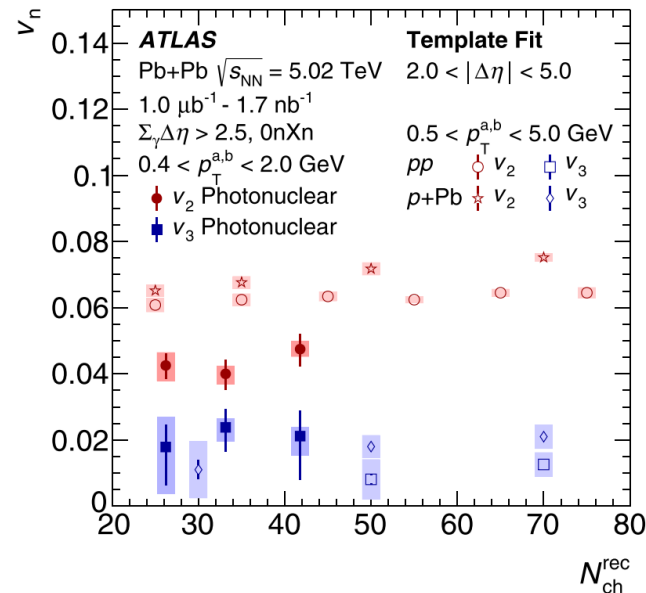
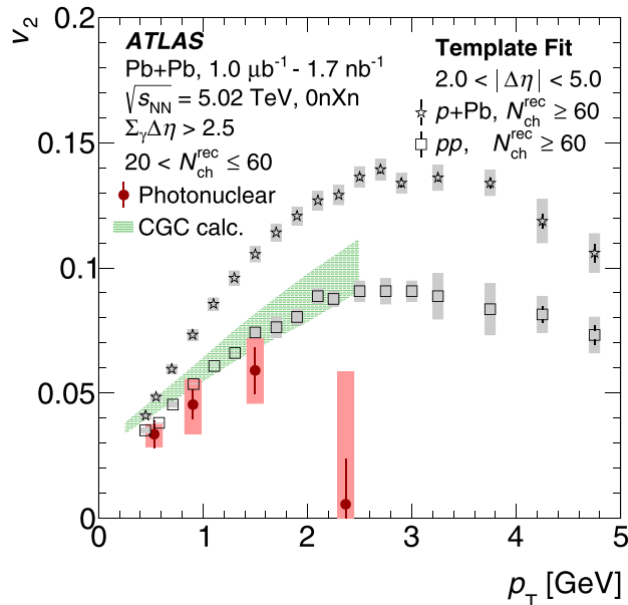
# The Future of Photo-nuclear Measurements

- Photo-nuclear jet production is complimentary to inclusive DIS measurements at the EIC, which will probe lower  $x$  and  $Q^2$ .
  - UPC jet measurements will provide increasingly precise tests of saturation physics before the EIC turns on.
- Photo-nuclear flow at the EIC can be predicted using UPC results by varying the  $Q^2$ , as shown on the right.
  - UPC flow measurements at the LHC will provide a key baseline for precision EIC measurements to come.



# Conclusions

- Photo-nuclear dijet production was measured by ATLAS in 5.02 TeV Pb+Pb collisions with 2018 data.
  - This measurement has been fully unfolded for detector response for the first time.
  - Once final studies of low- $\mu$  jet response in ATLAS can be completed, substantial gains in systematic control can be achieved.
- Studies of 2-particle correlations in photo-nuclear collisions demonstrate collectivity in a novel small system.
  - These results can be effectively modelled as a rho-nucleus collision.
- Both measurements are connected to early physics goals for the EIC.



# Backup

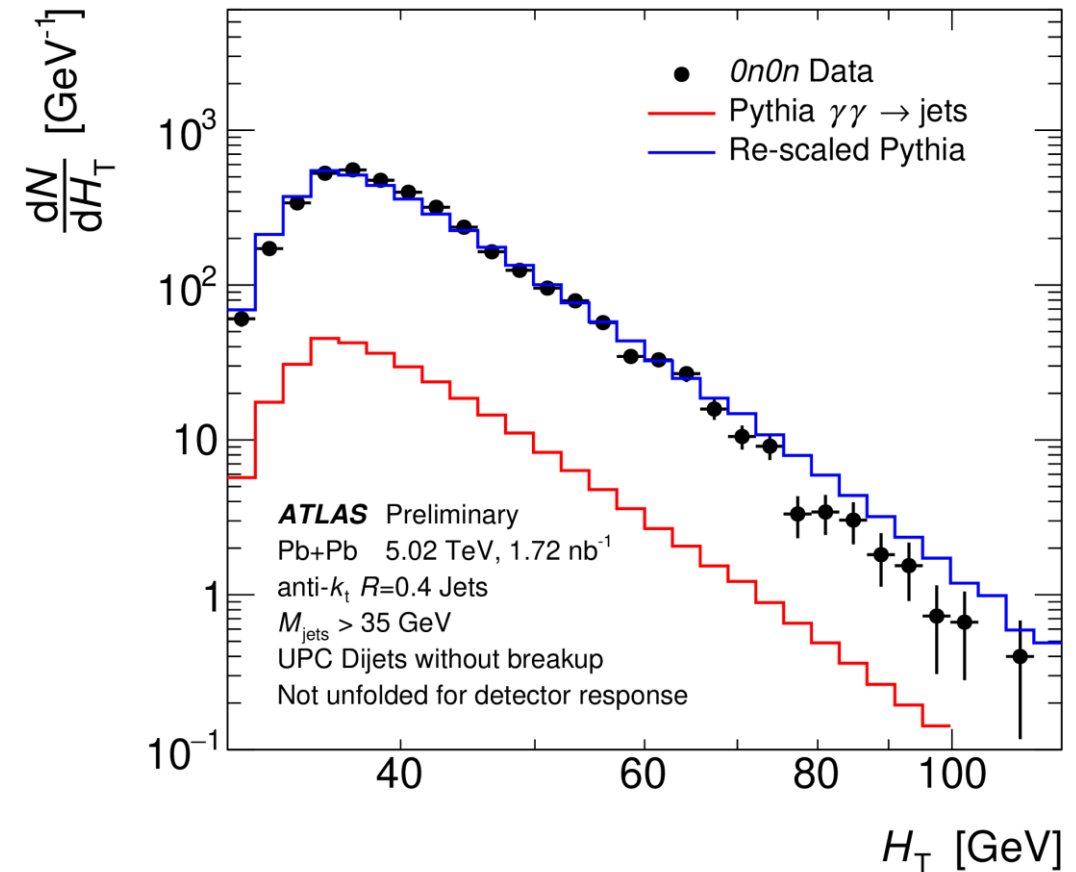
---

# Importance of Neutron Topology: $0n0n$ Events

For the first time, ATLAS has observed dijet production in UPC without nuclear breakup ( $0n0n$ ).

Gaps are required on both sides of the detector:  $\sum \Delta\eta > 2.0$

A factor of 10 more events are observed in data than are predicted from  $\gamma\gamma \rightarrow \text{jets}$ , estimated by Pythia or comparison to  $\gamma\gamma \rightarrow \mu^+\mu^-$  studies.



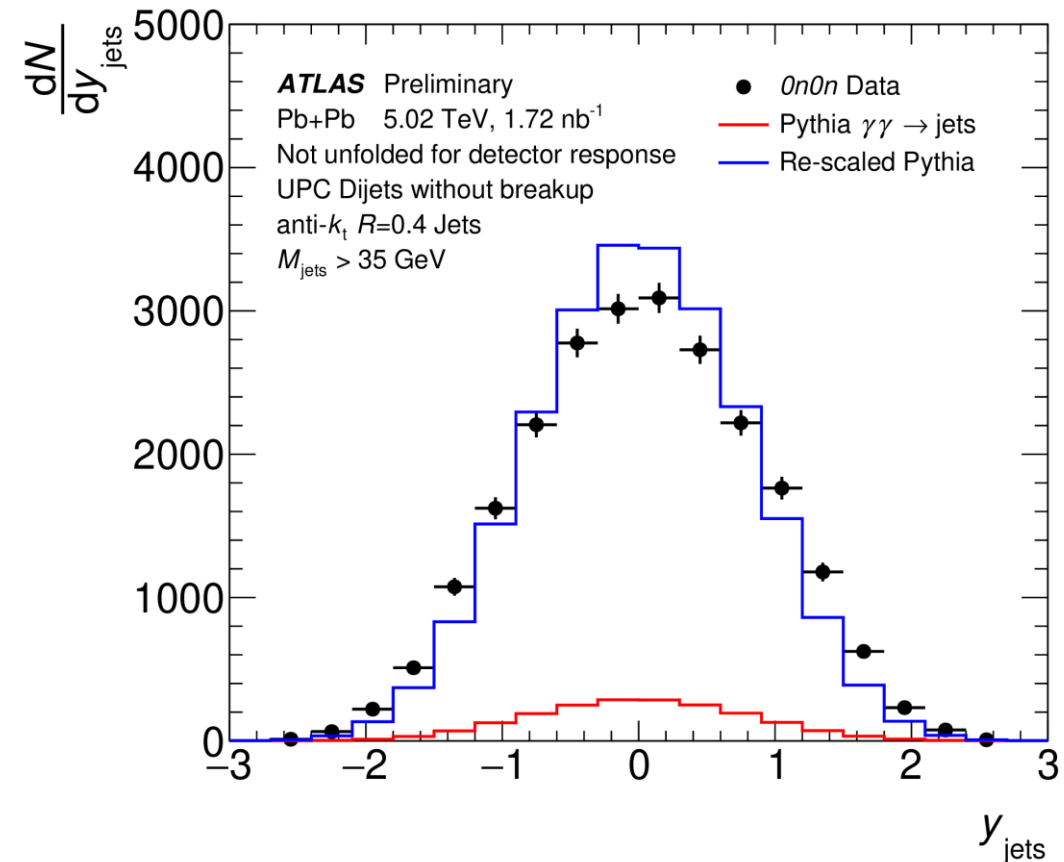
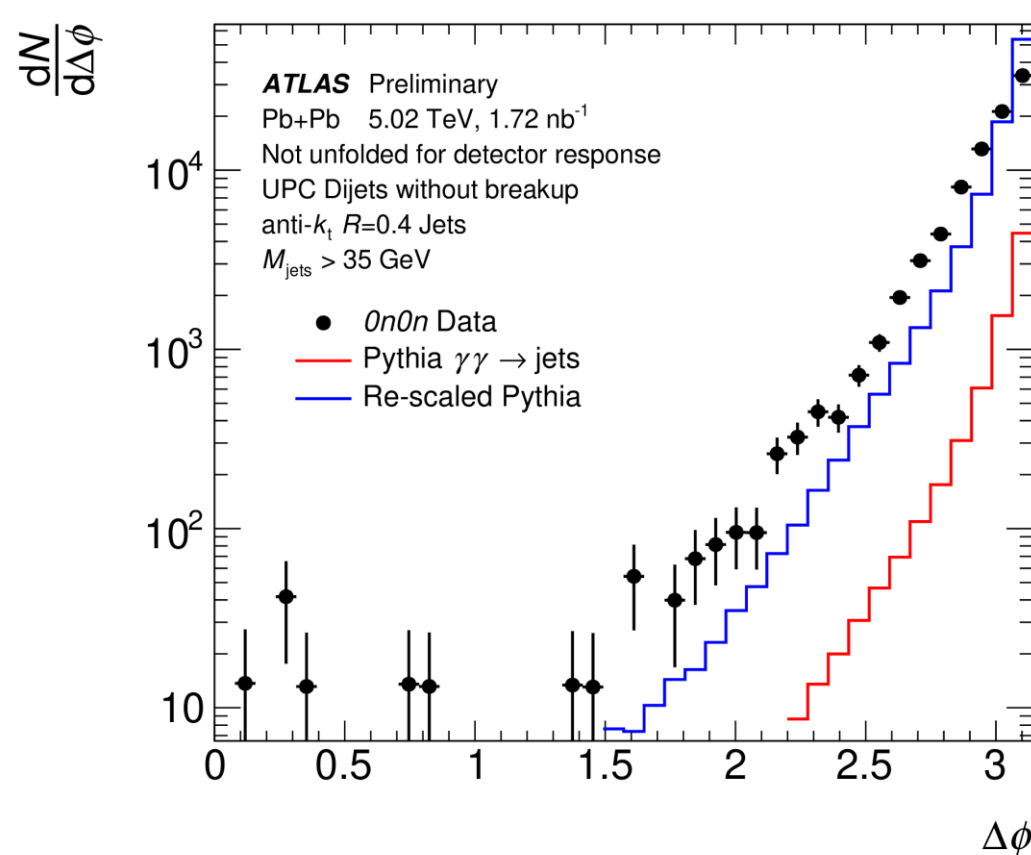
# Importance of Neutron Topology: $0n0n$ Events

For the first time, ATLAS has observed dijet production in UPC without nuclear breakup ( $0n0n$ ).

Gaps are required on both sides of the detector:  $\sum \Delta\eta > 2.0$

A factor of 10 more events are observed in data than are predicted from  $\gamma\gamma \rightarrow \text{jets}$ , estimated by Pythia or comparison to  $\gamma\gamma \rightarrow \mu^+\mu^-$  studies.

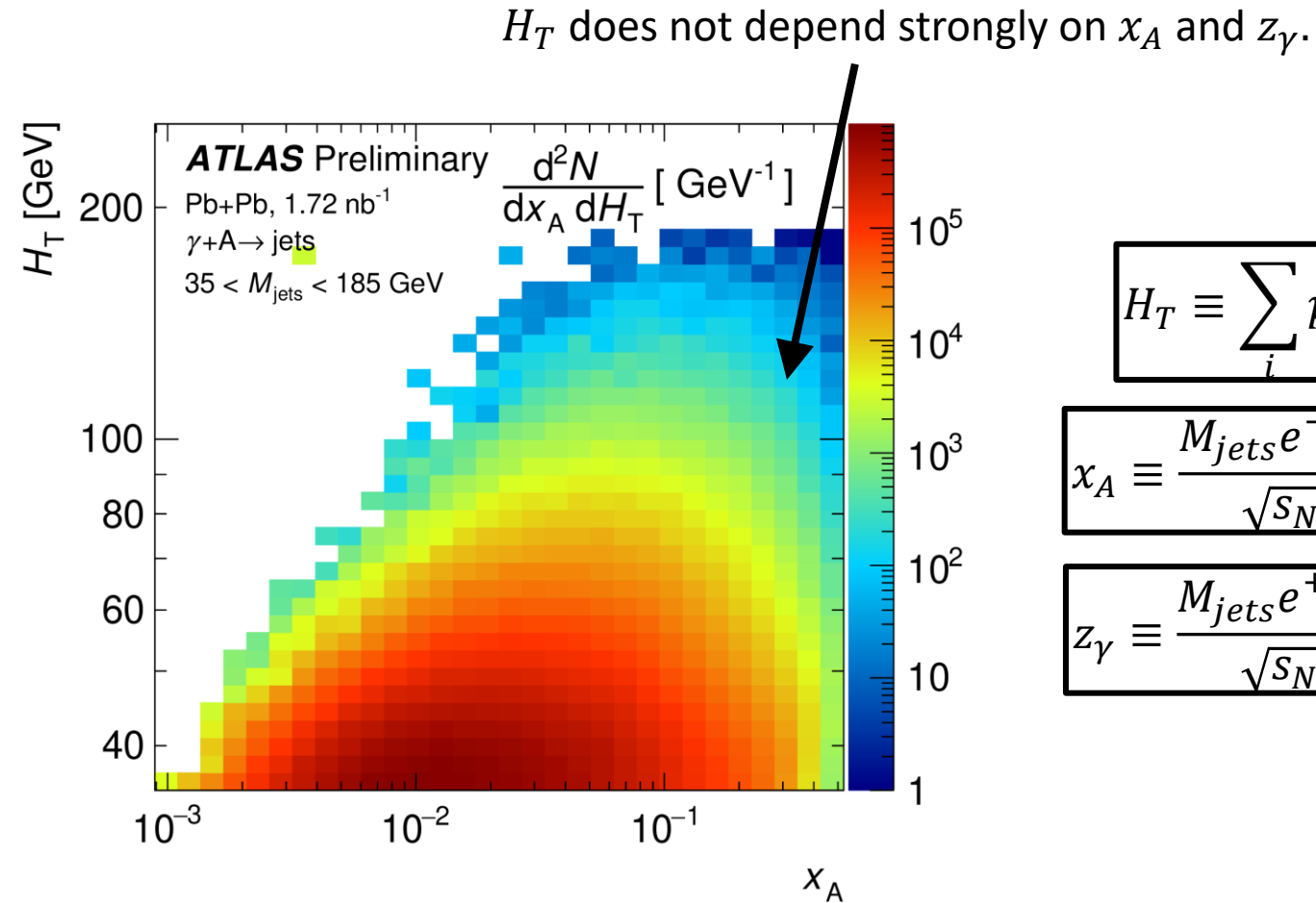
The distribution shapes look distinct from pure  $\gamma\gamma \rightarrow \text{jets}$ .



# Constructing the Jet Cross-Section

$$\frac{d^3\sigma}{dH_T dx_A dz_\gamma} = \frac{1}{\mathcal{L}} \frac{\Delta Y}{\Delta H_T \Delta x_A \Delta z_\gamma}$$

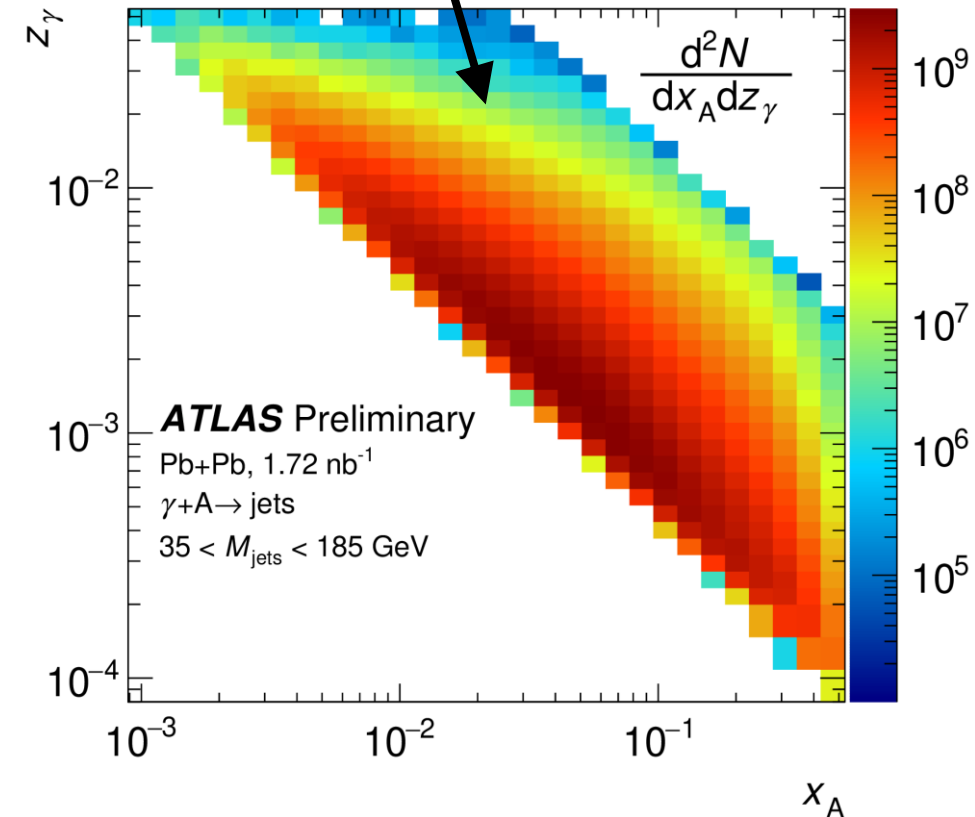
$x_A$  is strongly correlated with  $z_\gamma$ , yielding important acceptance effects in the two variables.



$$H_T \equiv \sum_i p_T^i$$

$$x_A \equiv \frac{M_{\text{jets}} e^{-y_{\text{jets}}}}{\sqrt{s_{NN}}}$$

$$z_\gamma \equiv \frac{M_{\text{jets}} e^{+y_{\text{jets}}}}{\sqrt{s_{NN}}}$$



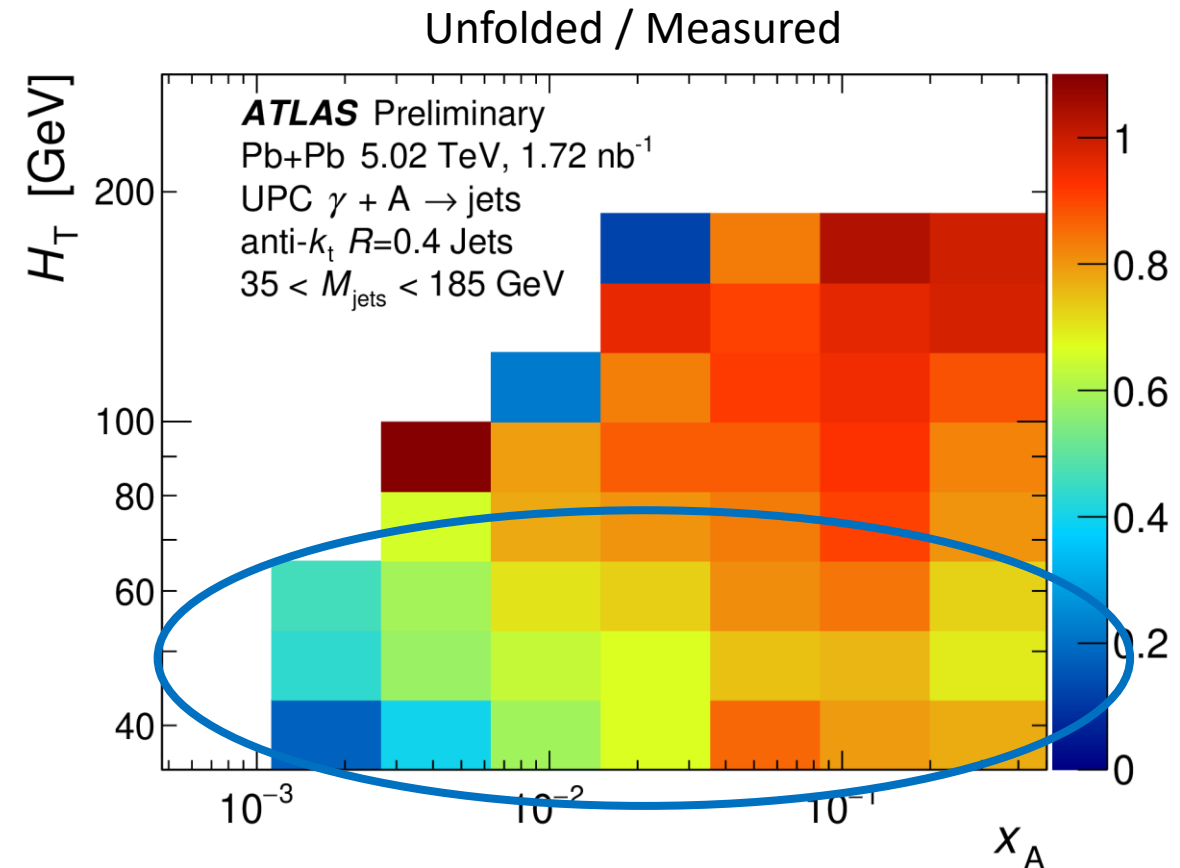
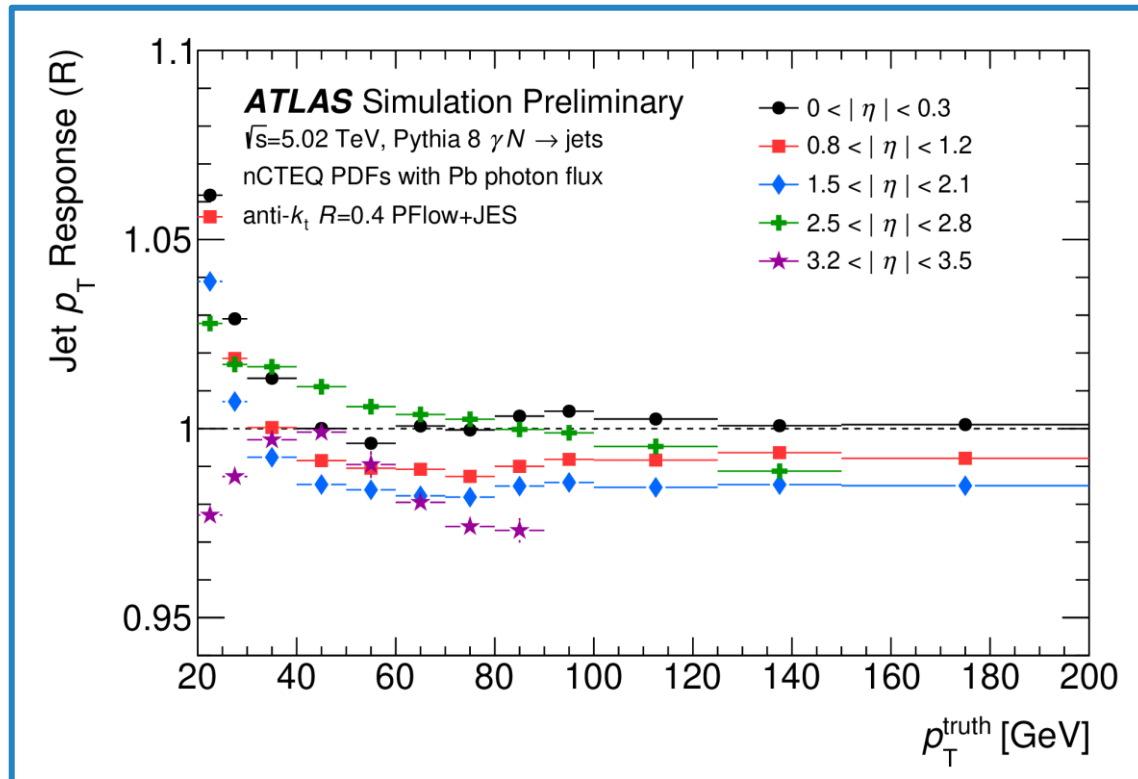


# Unfolding Measured Jet Cross-Sections

The measured cross-sections are then unfolded in 3 dimensions to correct for detector effects.

- Low- $p_T$  flavor effects are the largest correction.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



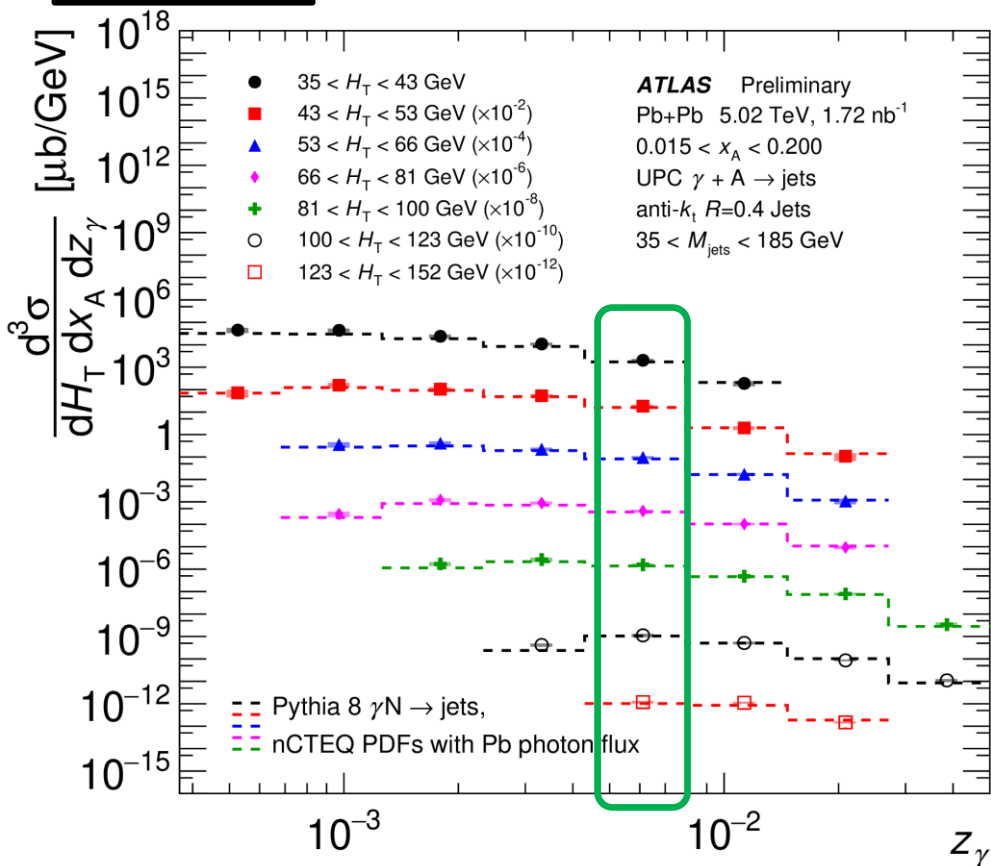
# Measured Jet Cross-Sections

- At intermediate photon energies, we can access higher-x partons.

$$H_T \equiv \sum_i p_T^i$$

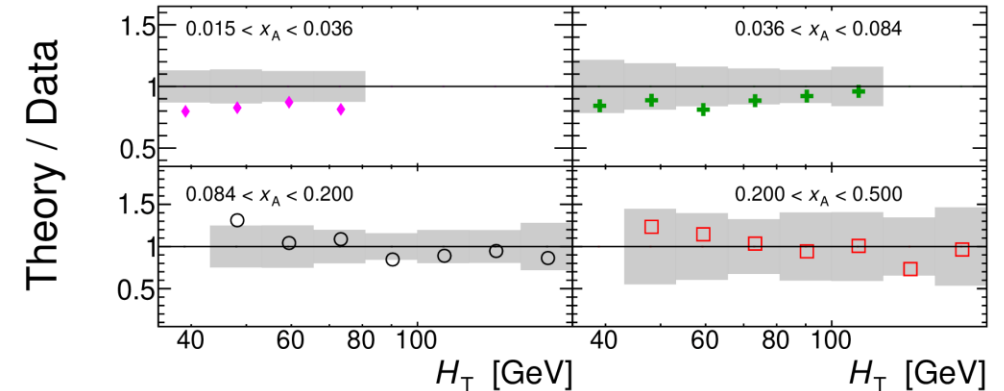
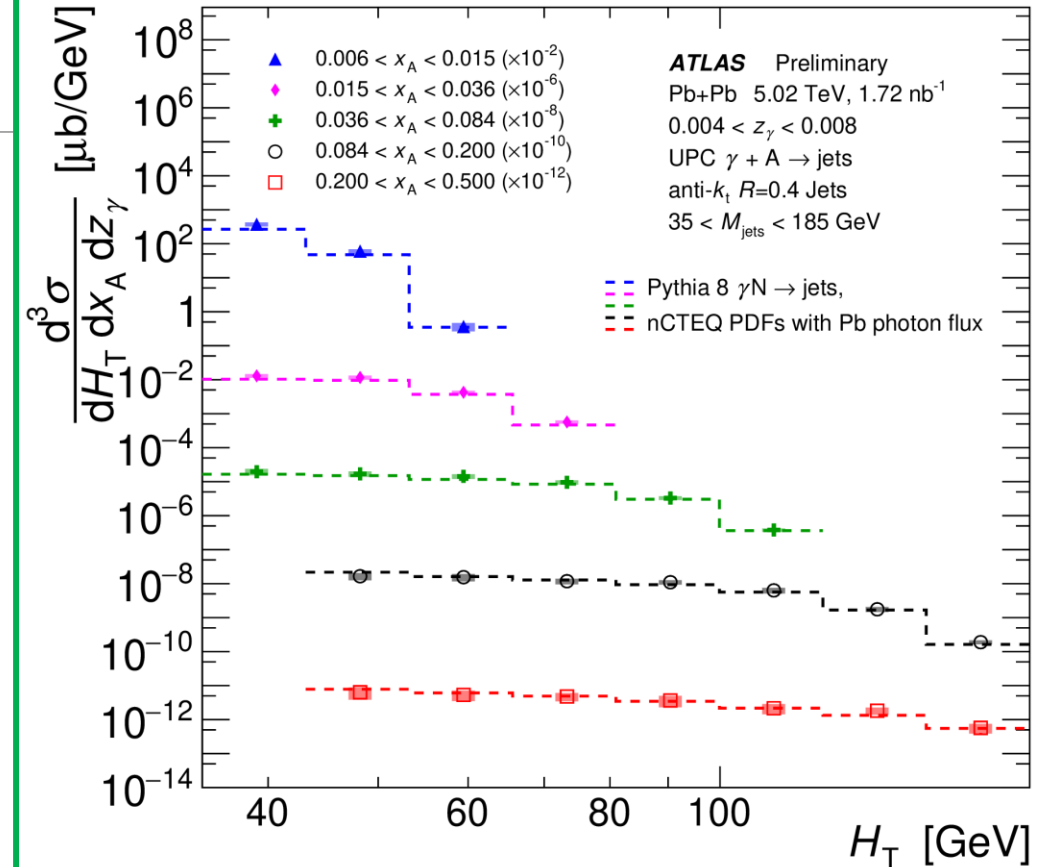
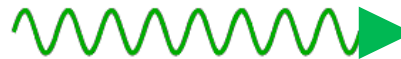
$$x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}}$$

$$z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



Photon Energy

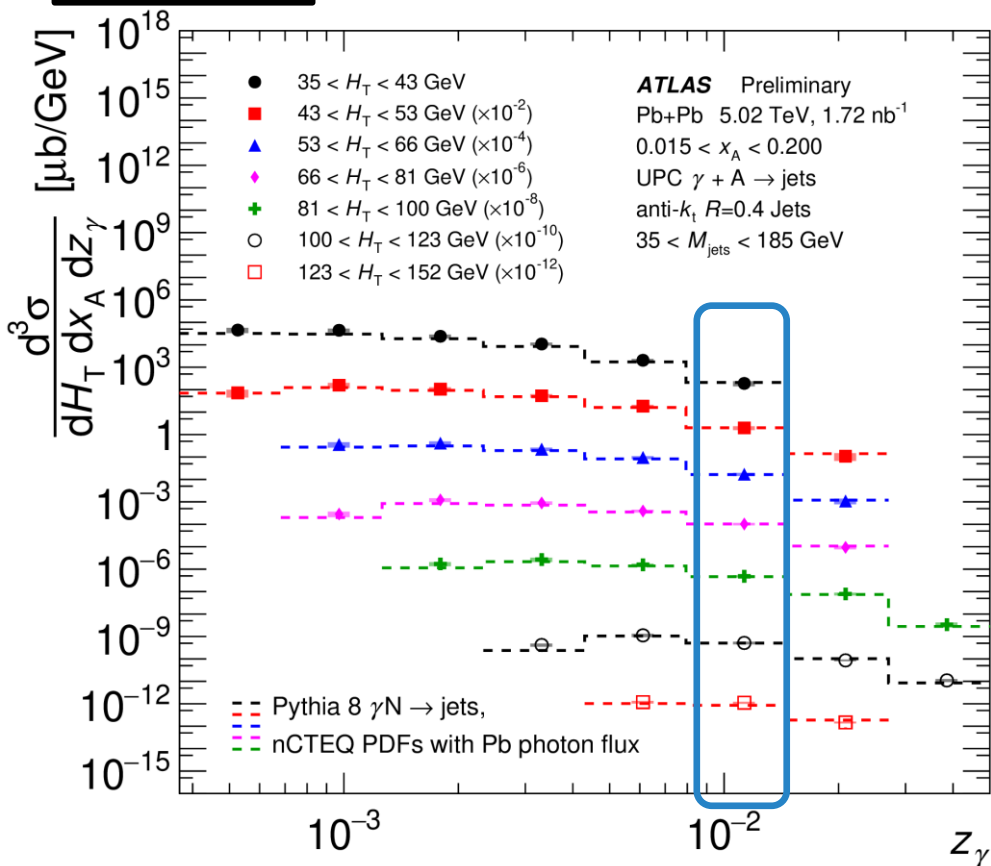
0.004 < z<sub>γ</sub> < 0.008



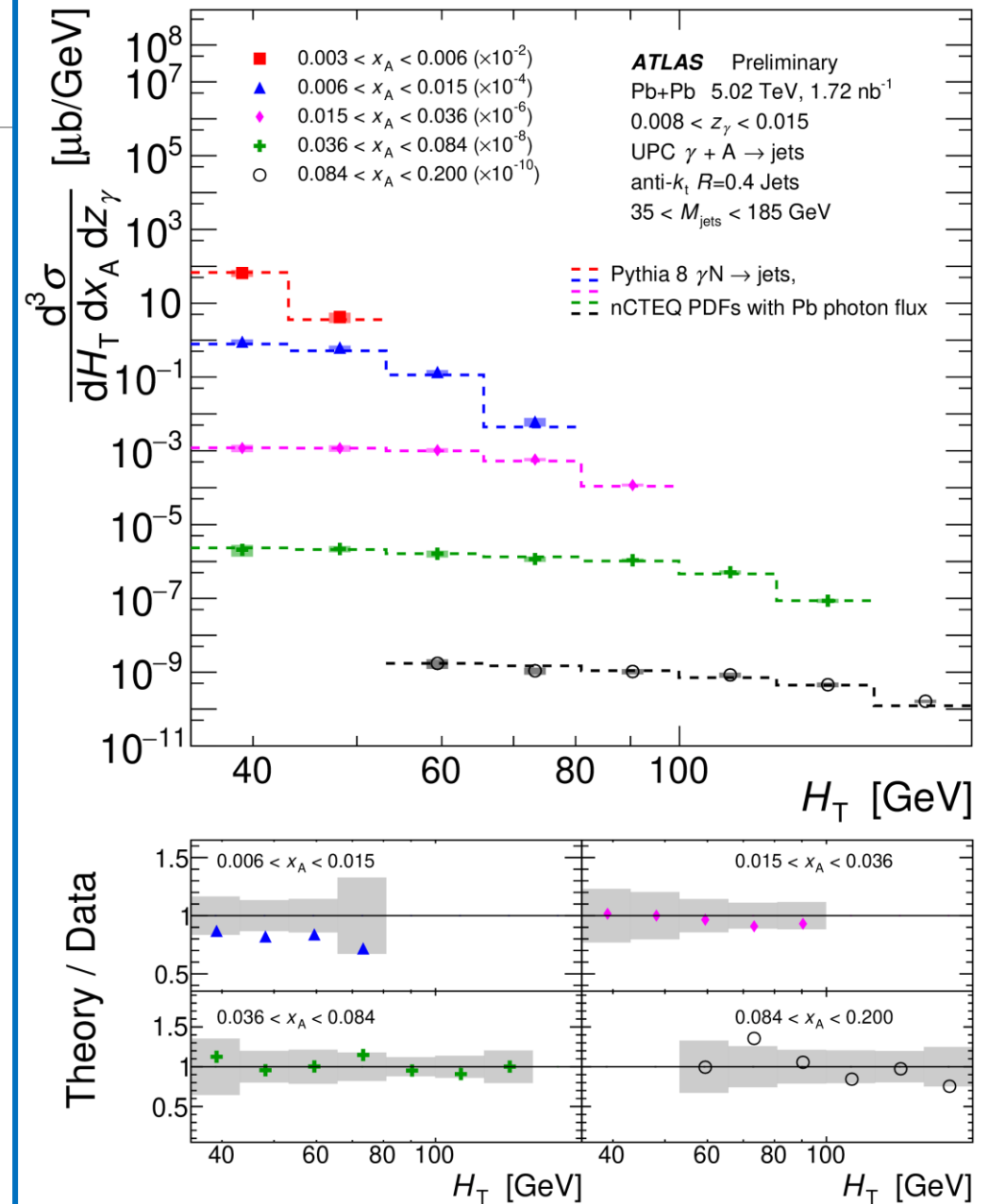
# Measured Cross-Sections

- Going higher in photon energy opens up the low-x shadowing region.
- Results are quite consistent with the theoretical model.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



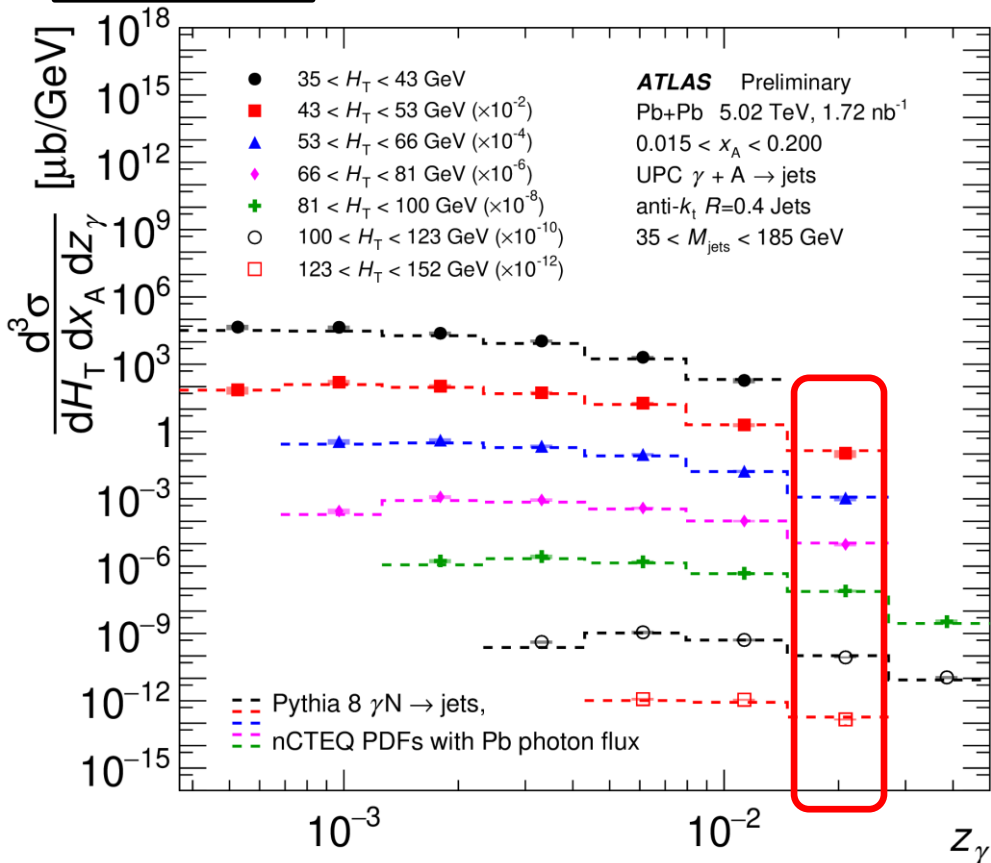
Photon Energy  
0.008 < z<sub>γ</sub> < 0.015  
⚡



# Measured Cross-Sections

- The highest photon energy allows the most access to low-x.
- Systematic control is more of a challenge near acceptance edges.

$$H_T \equiv \sum_i p_T^i \quad x_A \equiv \frac{M_{jets} e^{-y_{jets}}}{\sqrt{s_{NN}}} \quad z_\gamma \equiv \frac{M_{jets} e^{+y_{jets}}}{\sqrt{s_{NN}}}$$



## Photon Energy

$$0.015 < z_\gamma < 0.027$$

