

Two-photon transitions of charmonia on the light front

Meijian Li

University of Santiago de Compostela, Spain

In collaborations with: Yang Li, and James P. Vary

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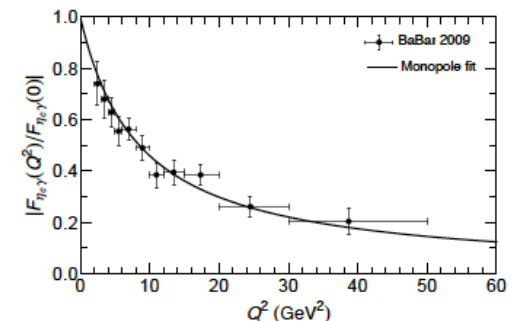
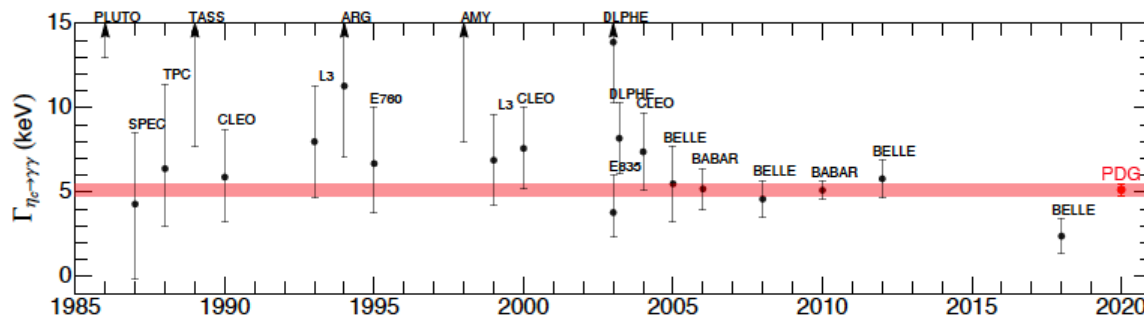
Outline

- **Introduction:** two-photon transition of charmonium
- **Method:** light-front Hamiltonian formalism
- **Results:** transition width and transition form factors
- **Summary**

*Based on: Y. Li, M. Li and J.P. Vary, Phys. Rev. D 105 (2022) 7, L071901;
arXiv:2111.14178 [hep-ph]*

Introduction: two-photon transition

- Charmonium provides an ideal testing ground for various investigations to understand QCD
 - Challenging: relativistic, non-perturbative effects
- The two-photon transition, $H_{c\bar{c}} \rightarrow \gamma^* + \gamma$, provides a clean and important probe to hadron states
- Experimental measurements
 - Diphoton width $\Gamma_{H \rightarrow \gamma\gamma}$: extensive measurements for $\eta_c, \eta_c', \chi_{c0}, \chi_{c2}$
 - Transition form factors $F_{H\gamma}(Q^2) : F_{\eta_{c\gamma}}(Q^2)$ by BABAR 2010; $F_{\chi_{c\gamma}}(Q^2)$ by Belle 2017

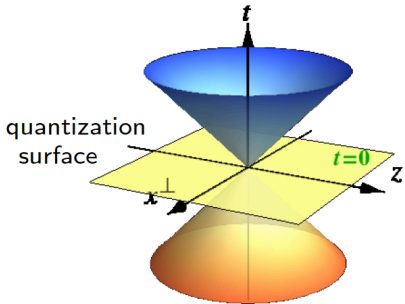
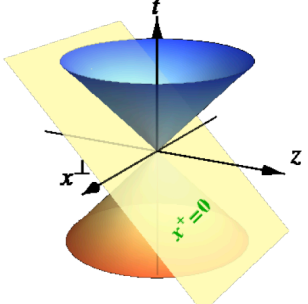
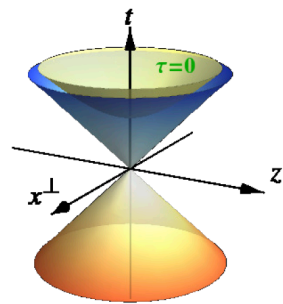


Method: light-front Hamiltonian formalism

Light-front Hamiltonian formalism is a natural framework for addressing relativistic bound-state and scattering problems in QCD

► Light-front quantization

- The quantum field is quantized on the equal light-front time $x^+ = 0$

instant form	front form	point form
time variable $t = x^0$	$x^+ \triangleq x^0 + x^3$	$\tau \triangleq \sqrt{t^2 - \vec{x}^2 - a^2}$
		
Hamiltonian $H = P^0$	$P^- \triangleq P^0 - P^3$	P^μ
kinematical \vec{P}, \vec{J}	$\vec{p}^\perp, P^+, \vec{E}^\perp, E^+, J^-$	\vec{J}, \vec{K}
dynamical \vec{K}, P^0	\vec{F}^\perp, P^-	\vec{P}, P^0
dispersion relation $p^0 = \sqrt{\vec{p}^2 + m^2}$	$p^- = (\vec{p}_\perp^2 + m^2)/p^+$	$p^\mu = mv^\mu \ (v^2 = 1)$

Method: light-front Hamiltonian formalism

► Hamiltonian formalism

- the invariant masses and the boost invariant wavefunctions can be obtained directly by solving the eigenvalue equation

$$(P^+ \hat{P}^- - \vec{P}_\perp^2) |\psi_h(P, j, m_j)\rangle = M_h^2 |\psi_h(P, j, m_j)\rangle$$

- the light-front wavefunction encodes the information of the system, and provides direct access to observables

► Basis representation

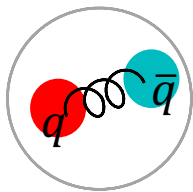
- basis can encode an analytical approximation to the solution
- optimal basis is the key to numerical efficiency

→ Basis Light-Front Quantization (BLFQ)

The charmonium light-front wavefunction by BLFQ

- The charmonium light-front wavefunction is solved using the BLFQ approach in the $|q\bar{q}\rangle$ sector¹,

$$H_{\text{eff}} = \underbrace{\frac{\vec{k}_{\perp}^2 + m_q^2}{x} + \frac{\vec{k}_{\perp}^2 + m_{\bar{q}}^2}{1-x}}_{\text{LF kinetic energy}} + \underbrace{\kappa^4 x(1-x) \vec{r}_{\perp}^2 - \frac{\kappa^4}{(m_q + m_{\bar{q}})^2} \frac{\partial}{\partial x} \left(x(1-x) \frac{\partial}{\partial x} \right)}_{\text{confinement}} + \underbrace{V_g}_{\text{one-gluon exchange}}$$



$$x = p_q^+ / P^+$$

$$\vec{k}_{\perp} = \vec{p}_{q\perp} - x \vec{P}_{\perp}$$

- **Confinement**
Transverse (QCD holography)²
Longitudinal (completes the transverse confinement, and produces desirable distribution amplitudes)
- **One-gluon exchange**

$$V_g = -\frac{4}{3} \frac{4\pi\alpha_s(Q^2)}{Q^2} \bar{u}_{\sigma'} \gamma^{\mu} u_{\sigma} \bar{v}_s \gamma_{\mu} v_{s'}$$

- Basis representation: basis functions are eigenfunctions of H_0

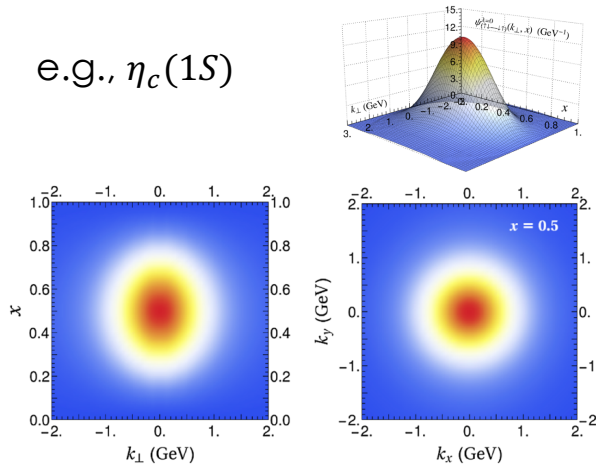
¹ Y. Li, P. Maris, and J. P. Vary, *Phys. Rev. D* **96**, 016022 (2017).

² S. J. Brodsky, G. F. de Teramond, H. G. Dosch, and J. Erlich, *Phys. Rept.* **584**, 1 (2015)

The charmonium light-front wavefunction by BLFQ

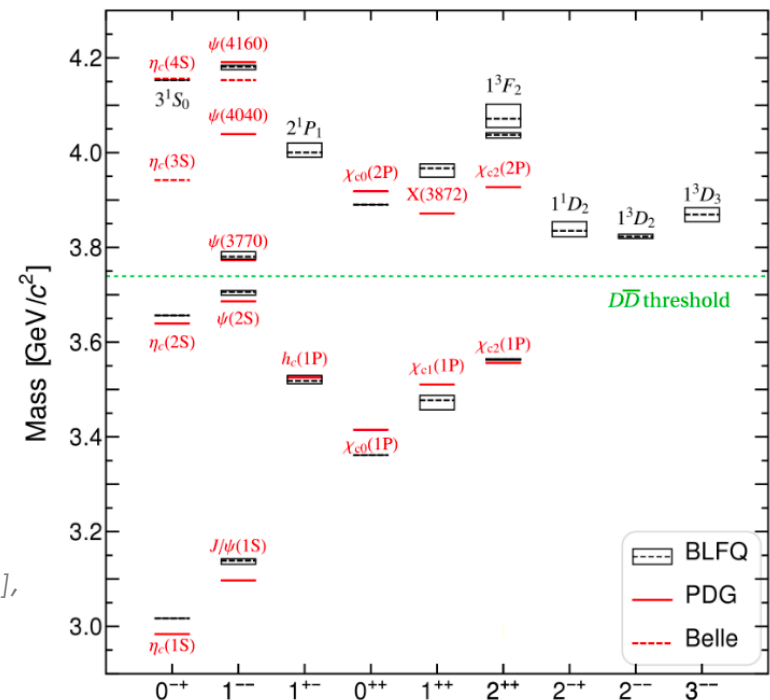
- The charmonium light-front wavefunction is solved using the BLFQ approach in the $|q\bar{q}\rangle$ sector¹,
 - Light-front wavefunctions:
 - Mass spectra:

e.g., $\eta_c(1S)$



- ✓ Access to a variety of observables:
 - Form factors [Li, PRD '18; Mondal, PRD '20],
 - PDFs/GPDs [Lan, PRL '19, PRD '20; Adhikari, PRC '18, '21],
 - radiative transitions [M. Li, PRD '18 & '19],
 - diffractive production [Chen, PLB '17 & PRC '18]

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The two-photon transition

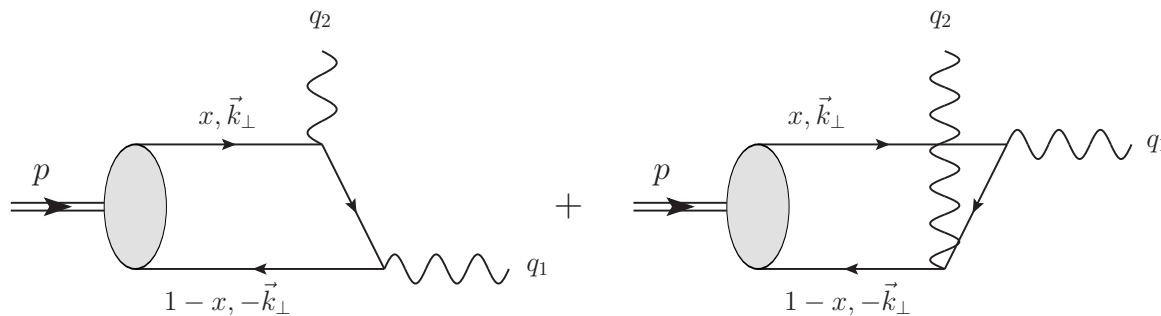
- The amplitude of the hadron-to-two-photon transition, $H_{c\bar{c}}(j, p, \lambda) \rightarrow \gamma^*(q_1, \lambda_1) + \gamma(q_2, \lambda_2)$, is related to the hadron matrix element,

$$\epsilon_\mu^*(q_1, \lambda_1) \epsilon_\nu^*(q_2, \lambda_2) e_\alpha(p, \lambda) \mathcal{M}^{\mu\nu\alpha} =$$

$$H_{\lambda_1, \lambda_2; \lambda}(q_1, q_2) = \epsilon_\nu^*(q_2, \lambda_2) \langle \gamma^*(q_1, \lambda_1) | J^\nu(0) | H(p, \lambda) \rangle$$

The photon light-front wavefunction (by LO perturbation)

The hadron light-front wavefunction (by BLFQ)



Light-cone dominance is manifest in the frame

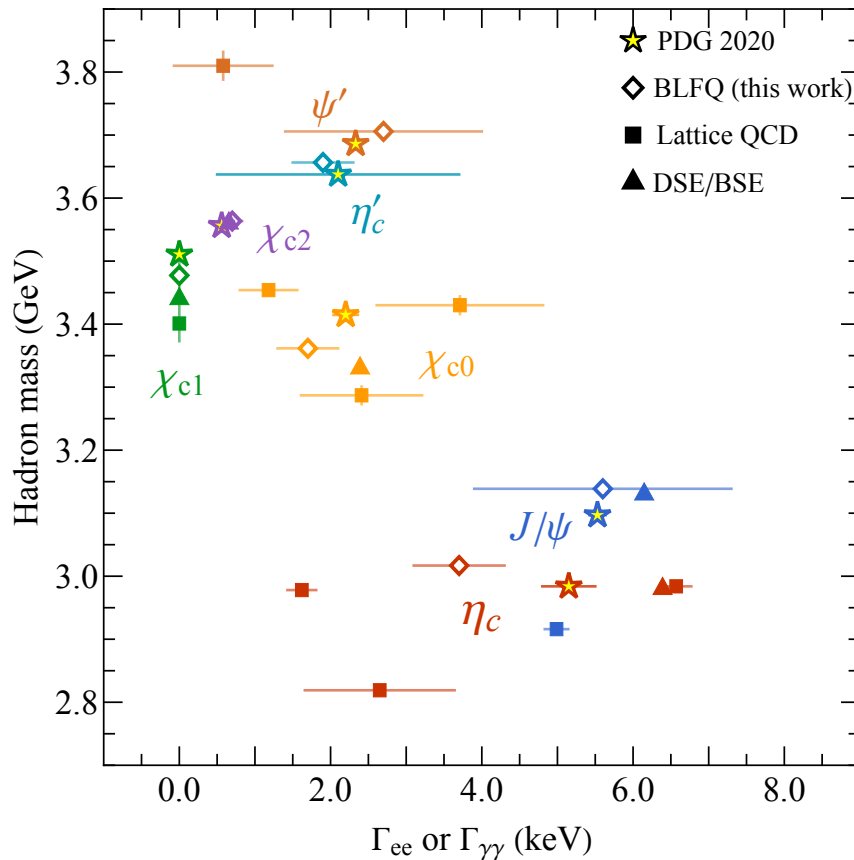
$$q_2^+ = q_1^- = 0$$

- The decay width, can be measured from experiments

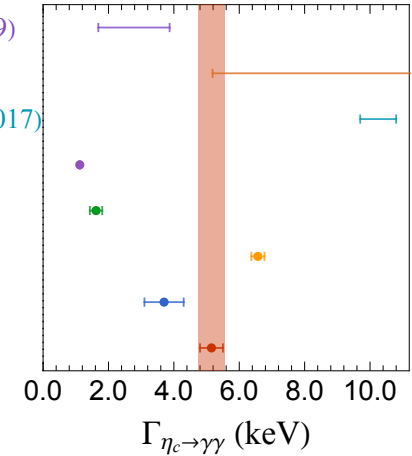
$$\Gamma_{H \rightarrow \gamma\gamma} = \frac{1}{2} \frac{1}{16\pi} \frac{1}{m_H} \frac{1}{2j+1} \sum_{\lambda=-j}^j \sum_{\lambda_1, \lambda_2=\pm 1} |H_{\lambda_1, \lambda_2; \lambda}|^2$$

Results: two-photon decay widths

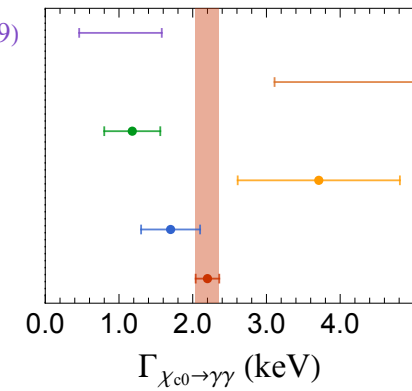
- Without any parameter tuning, our results have reasonable agreement with experimental data¹



NRQM/LF (Babiarz 2019)
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 NNLO NRQCD (Feng 2017)
 Lattice (Chen 2016)
 Lattice (Chen 2020)
 Lattice (Meng 2021)
 BLFQ (this work)
 PDG 2020



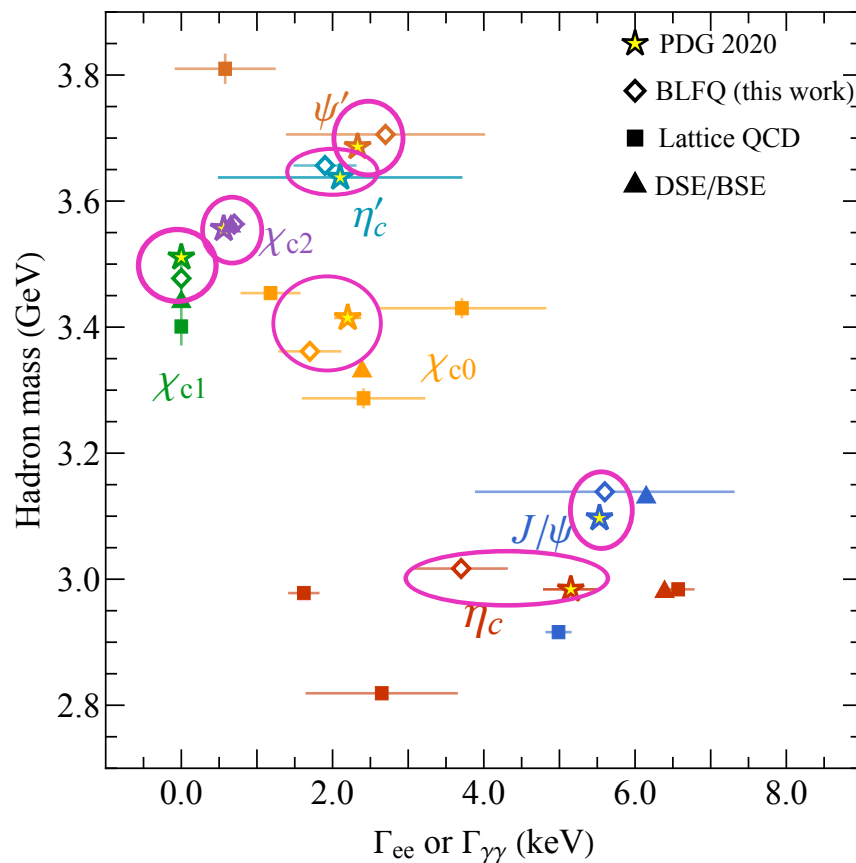
NRQM/LF (Babiarz 2019)
 NRQM (Babiarz 2019)
 Lattice (Chen 2020)
 Lattice (Zou 2021)
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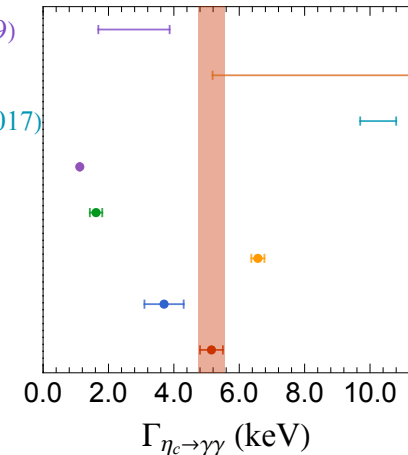
¹ Y. Li, M. Li and J.P. Vary, *Phys. Rev. D* 105 (2022) 7, L071901; arXiv:2111.14178 [hep-ph]

Results: two-photon decay widths

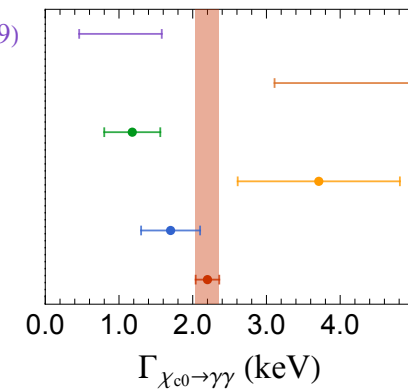
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Results: the transition form factor

(1) pseudoscalar 0^{+-}

- The pseudoscalar transition amplitude is parameterized as

$$\mathcal{M}^{\mu\nu} = 4\pi\alpha_{\text{em}}\epsilon^{\mu\nu\rho\sigma}q_{1\rho}q_{2\sigma}F_{P\gamma\gamma}(q_1^2, q_2^2)$$

- Single-tagged transition form factor: $F_{P\gamma}(Q^2 = -q^2) = F_{P\gamma\gamma}(q^2, 0) = F_{P\gamma\gamma}(0, q^2)$

- The transition form factor in the light-front wavefunction representation reads

$$F_{P\gamma}(Q^2) = 2Q_f^2 \sqrt{2N_c} \int \frac{d^2k_{\perp}}{(2\pi)^3} \int_0^1 \frac{dx}{2\sqrt{x(1-x)}} \frac{\psi_{\uparrow\downarrow-\downarrow\uparrow/P}(x, \vec{k}_{\perp})}{k_{\perp}^2 + m_f^2 + x(1-x)Q^2}$$

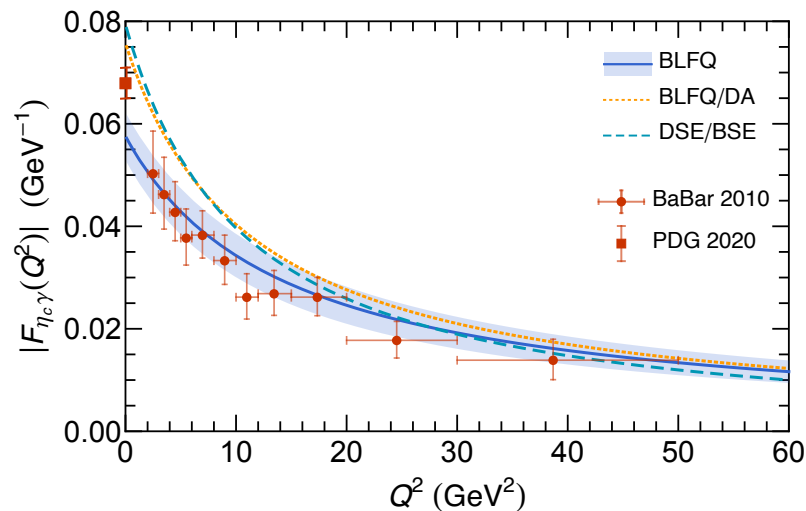
- At large Q , it reduces to the partonic interpretation in terms of light-cone distribution amplitude (LCDA)

$$F_{P\gamma}(Q^2) = \frac{e_f^2 f_P}{Q^2} \int_0^1 dx \frac{\phi_P(x, Q)}{x(1-x)}$$

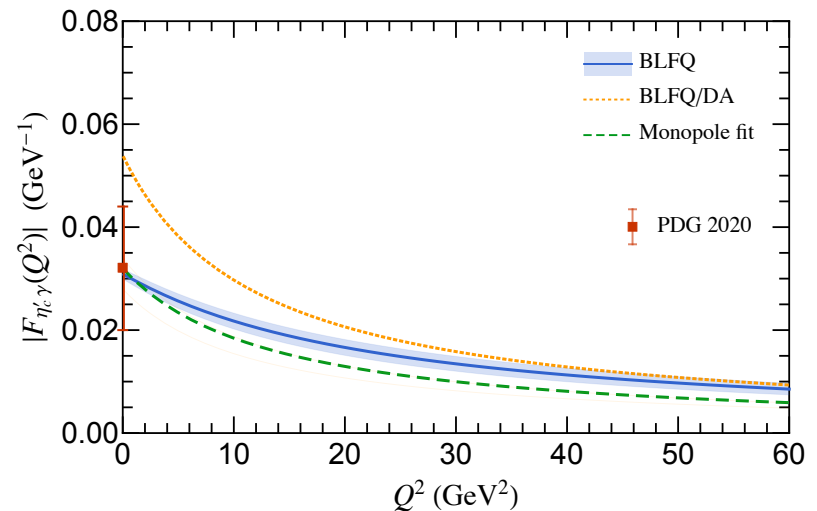
Results: the transition form factor

(1) pseudoscalar 0^{+-}

- η_c



- η_c'



► The calculated transition form factors are in a reasonable agreement with experimental data

- BLFQ (this work), uncertainty is calculated from basis sensitivity
- BLFQ/DA, using the LCDA obtained from the BLFQ wavefunction
- Monopole fit, vector meson dominance model

Results: the transition form factor

(2) scalar 0^{++}

- The transition amplitude is parameterized as

$$\begin{aligned} \mathcal{M}^{\mu\nu}(q_1, q_2) = & 4\pi\alpha_{em} \left\{ [(q_1 \cdot q_2)g^{\mu\nu} - q_2^\mu q_1^\nu] F_{S\gamma\gamma,1}(q_1^2, q_2^2) \right. \\ & \left. + \frac{1}{m_S^2} [q_1^2 q_2^2 g^{\mu\nu} + (q_1 \cdot q_2) q_1^\mu q_2^\nu - q_1^2 q_2^\mu q_2^\nu - q_2^2 q_1^\mu q_1^\nu] F_{S\gamma\gamma,2}(q_1^2, q_2^2) \right\} \end{aligned}$$

- Single-tagged transition form factor: $F_{S\gamma}(Q^2 = -q^2) = F_{S\gamma\gamma,1}(q^2, 0) = F_{S\gamma\gamma,1}(0, q^2)$

- The transition form factor in the light-front wavefunction representation reads

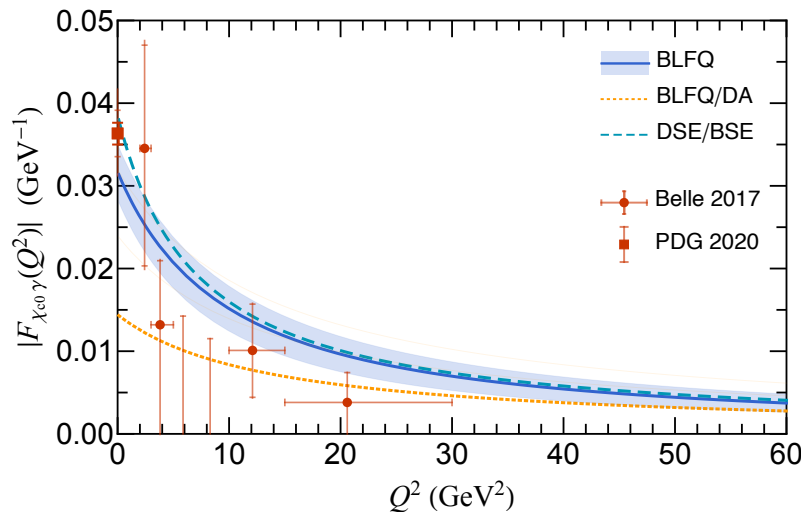
$$\begin{aligned} F_{S\gamma}(Q^2) = & e_f^2 2\sqrt{2N_c} \int_0^1 \frac{dx}{2\sqrt{x(1-x)}} \int \frac{d^2 k_\perp}{(2\pi)^3} \\ & \times \left\{ \psi_{\uparrow\downarrow+\downarrow\uparrow/S}(x, \vec{k}_\perp) \frac{(1-2x)[x(1-x)Q^2 + m_f^2]}{[k_\perp^2 + x(1-x)Q^2 + m_f^2]^2} + \psi_{\uparrow\uparrow/S}(x, \vec{k}_\perp) \frac{\sqrt{2}m_f(k_x + ik_y)}{[k_\perp^2 + x(1-x)Q^2 + m_f^2]^2} \right\} \end{aligned}$$

- At large Q, in terms of distribution amplitude (LCDA)

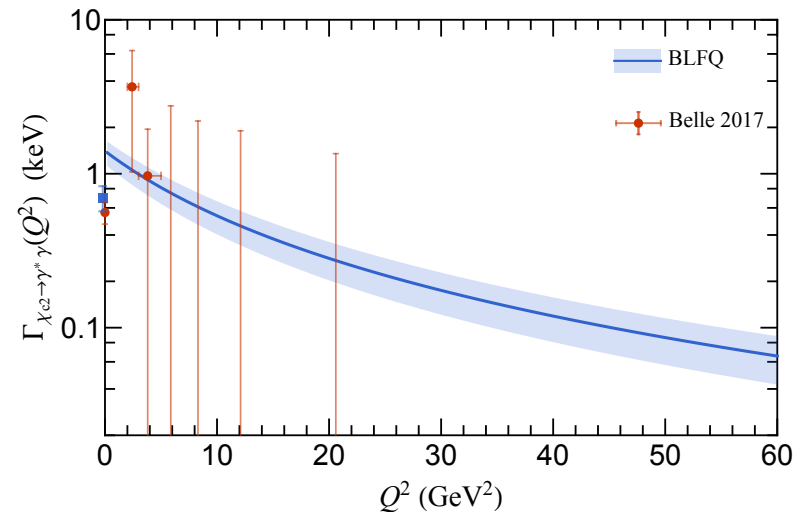
$$F_{S\gamma}(Q^2) = e_f^2 f_S \int_0^1 dx \frac{(1-2x)\phi_S(x, \mu)}{x(1-x)Q^2 + m_f^2}$$

Results: the transition form factor (2) scalar 0^{++} and tensor 2^{++}

- $\chi_{c0} (0^{++})$



- $\chi_{c2} (2^{++})$



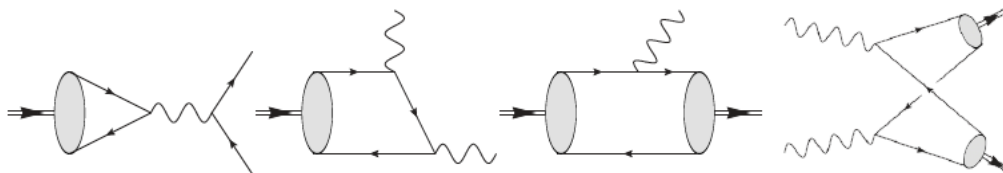
- The calculated transition form factors are in a reasonable agreement with experimental data
 - BLFQ (this work), uncertainty is calculated from basis sensitivity
 - BLFQ/DA, using the LCDA obtained from the BLFQ wavefunction

Summary and outlook

- We investigated the [two-photon transitions](#) of charmonia, η_c , η_c' , χ_{c0} , and χ_{c2} , in the light-front Hamiltonian approach
 - We derived the [formulas](#) of transition form factors in the light-front wavefunction representation
 - Universal for other hadron light-front wavefunctions
 - We computed the decay widths, and the transition form factors, **both** in [good agreements](#) with experimental measurements
 - Reveal relativistic nature of charmonia

Based on: Y. Li, M. Li and J.P. Vary, Phys. Rev. D 105 (2022) 7, L071901; arXiv:2111.14178 [hep-ph]. LFWFs available on Mendeley Data

- Ongoing and future works on radiative transitions
 - A comprehensive study on different leptonic and radiative transitions
 - Extension to bottomonia, heavy-light mesons, and light mesons



Thank you!