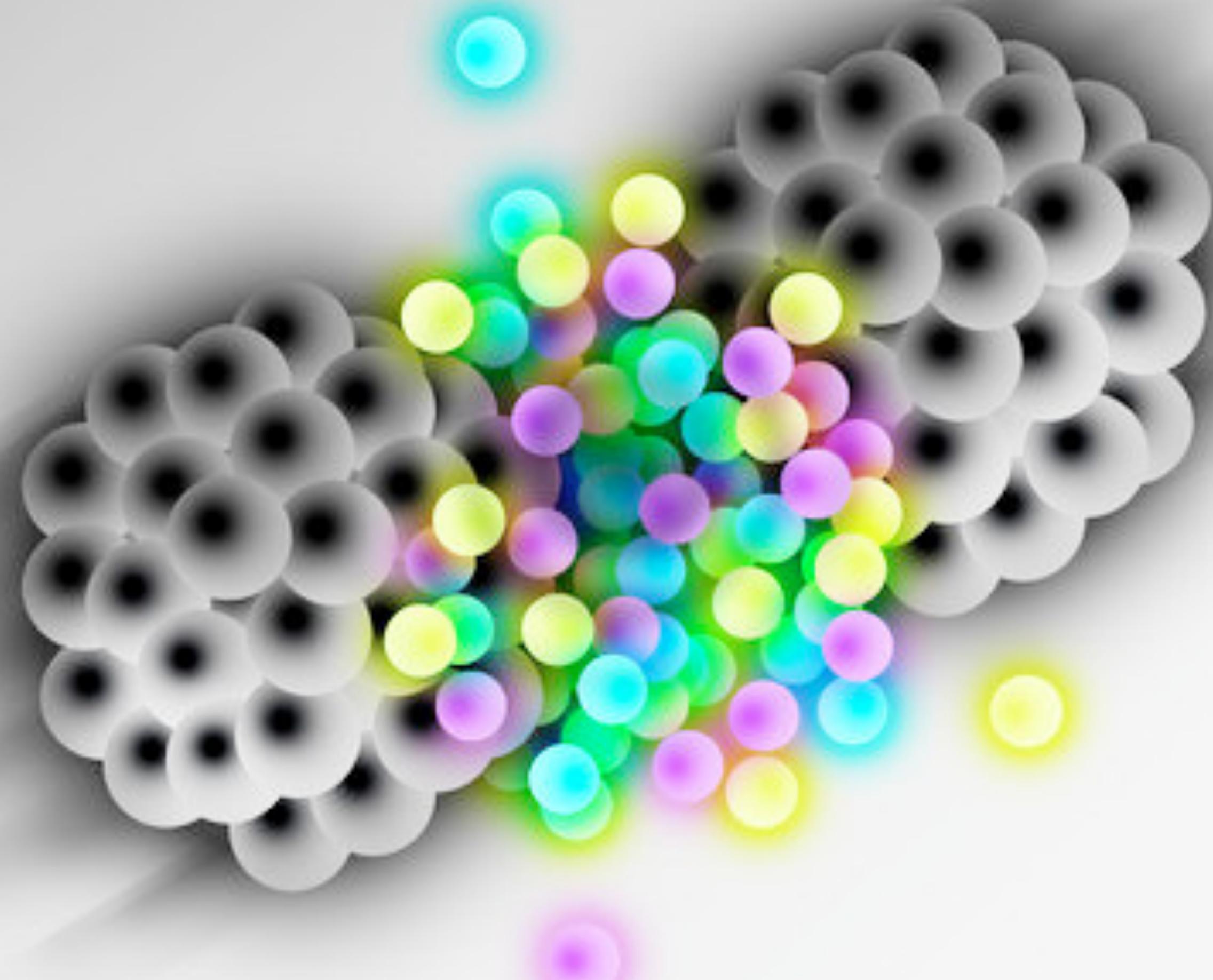


New winds in heavy-ion physics



Liliana Apolinário

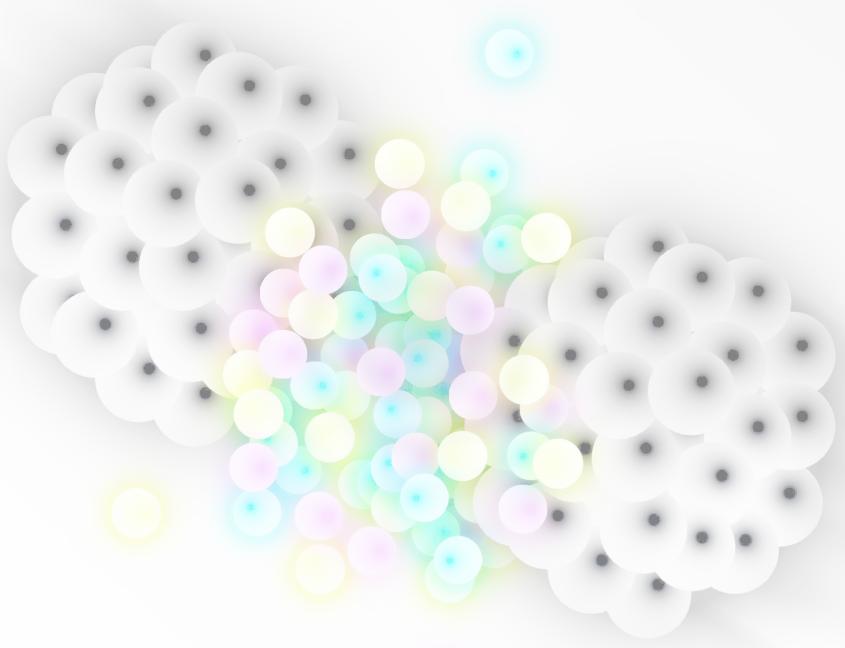


TÉCNICO
LISBOA

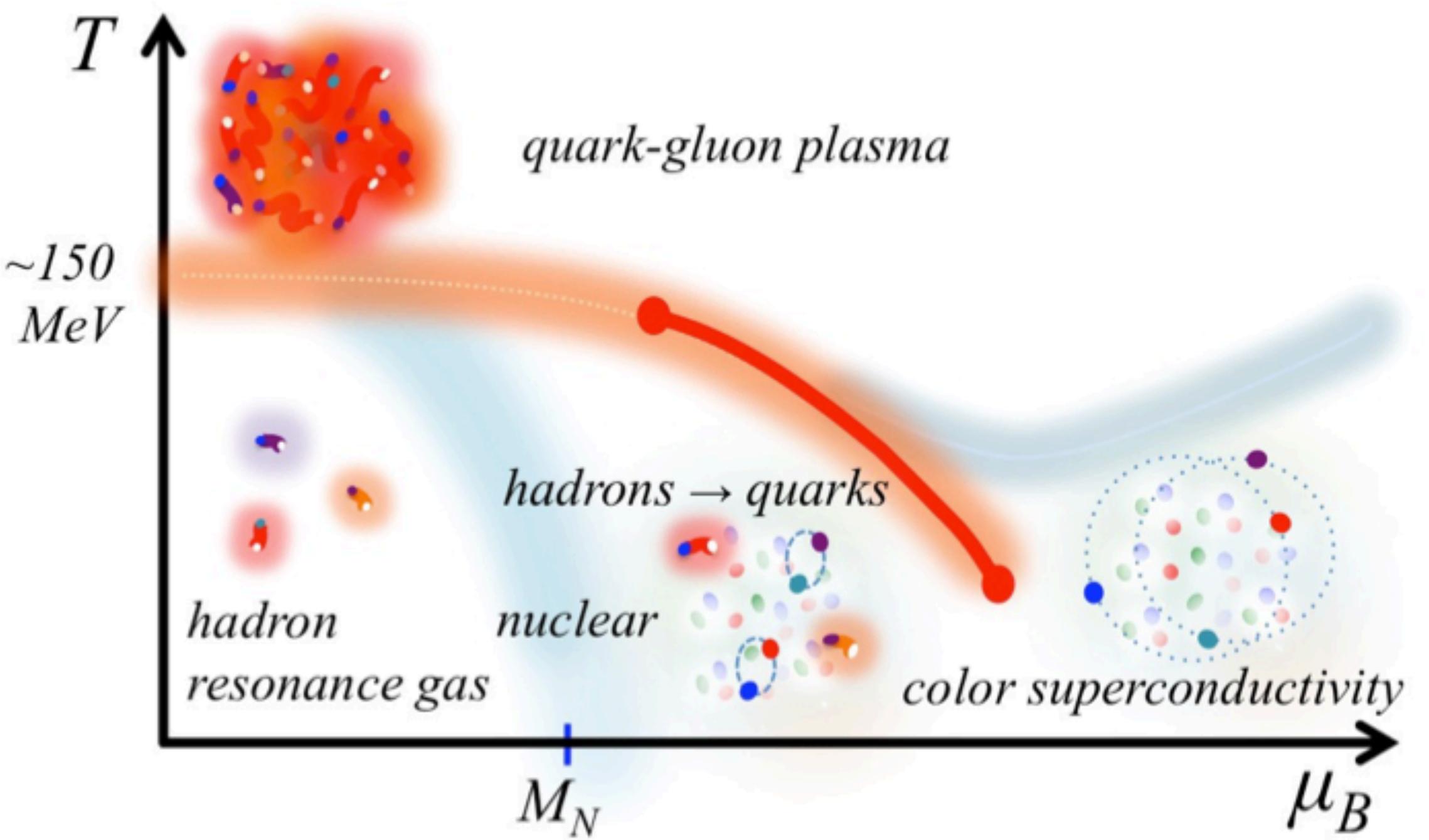
DIS2022

Monday, May 2nd

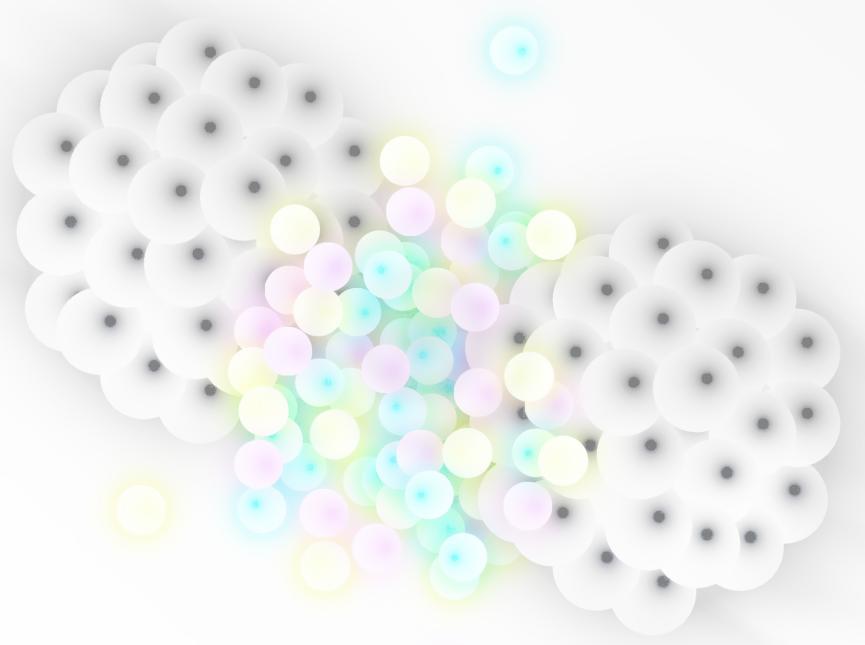
Heavy-Ion Collisions



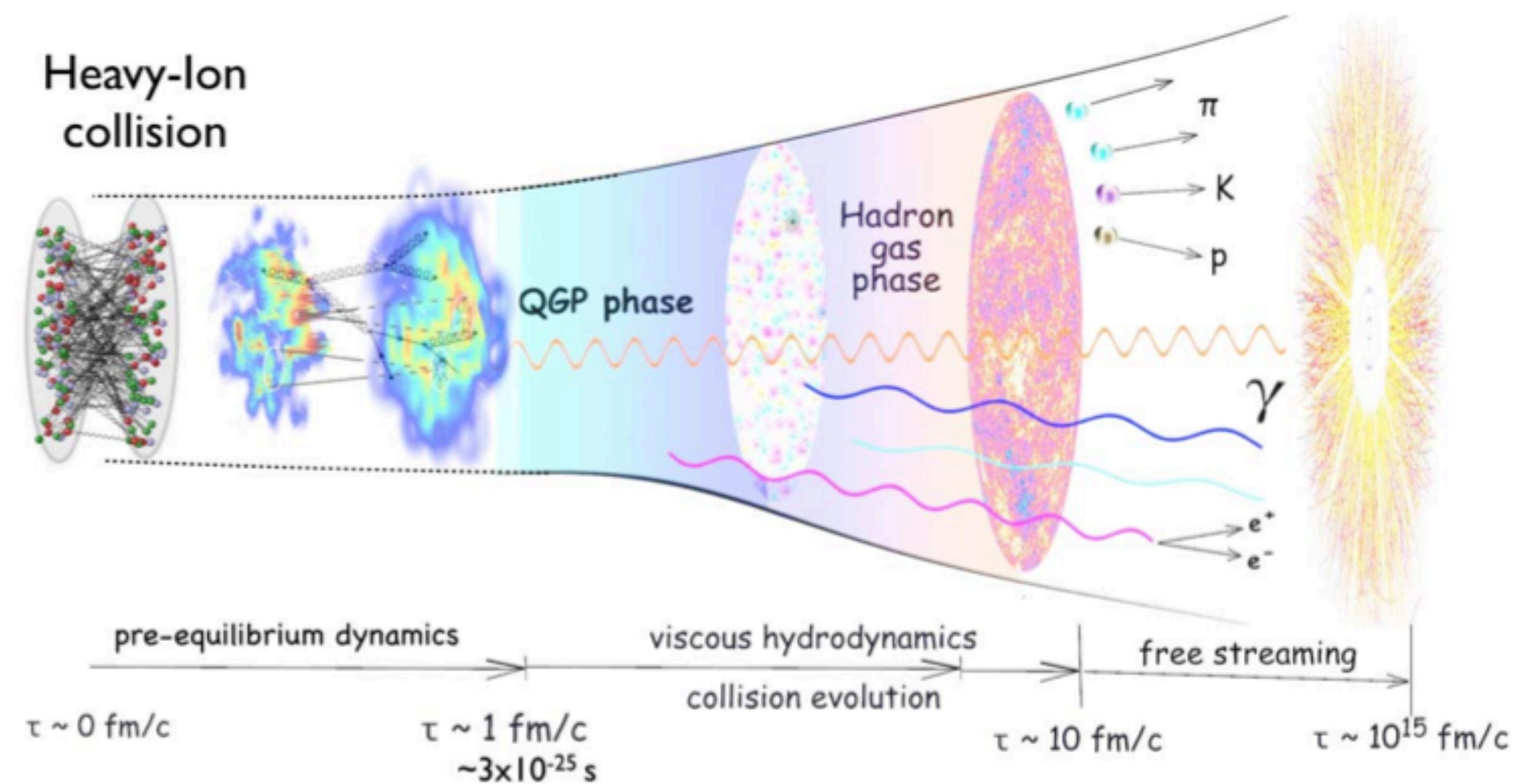
- Heavy-ion collision:
 - Probe the QCD phase diagram
 - Understand the QCD fundamental interactions
 - Collectivity from a gauge-field theory?
 - Tools used to study created matter shared with nearby physics fields research
 - QGP vs colliding nuclei?



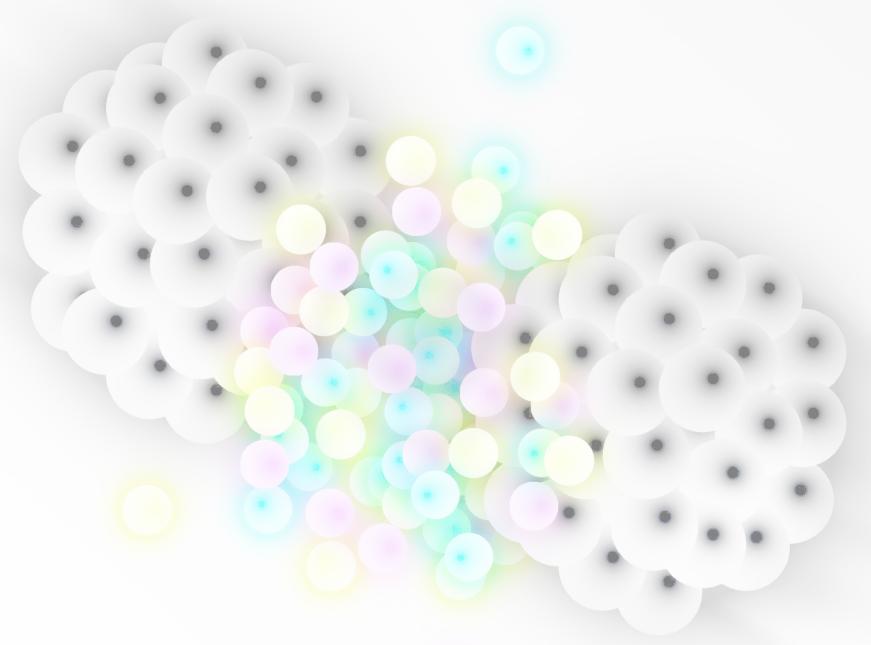
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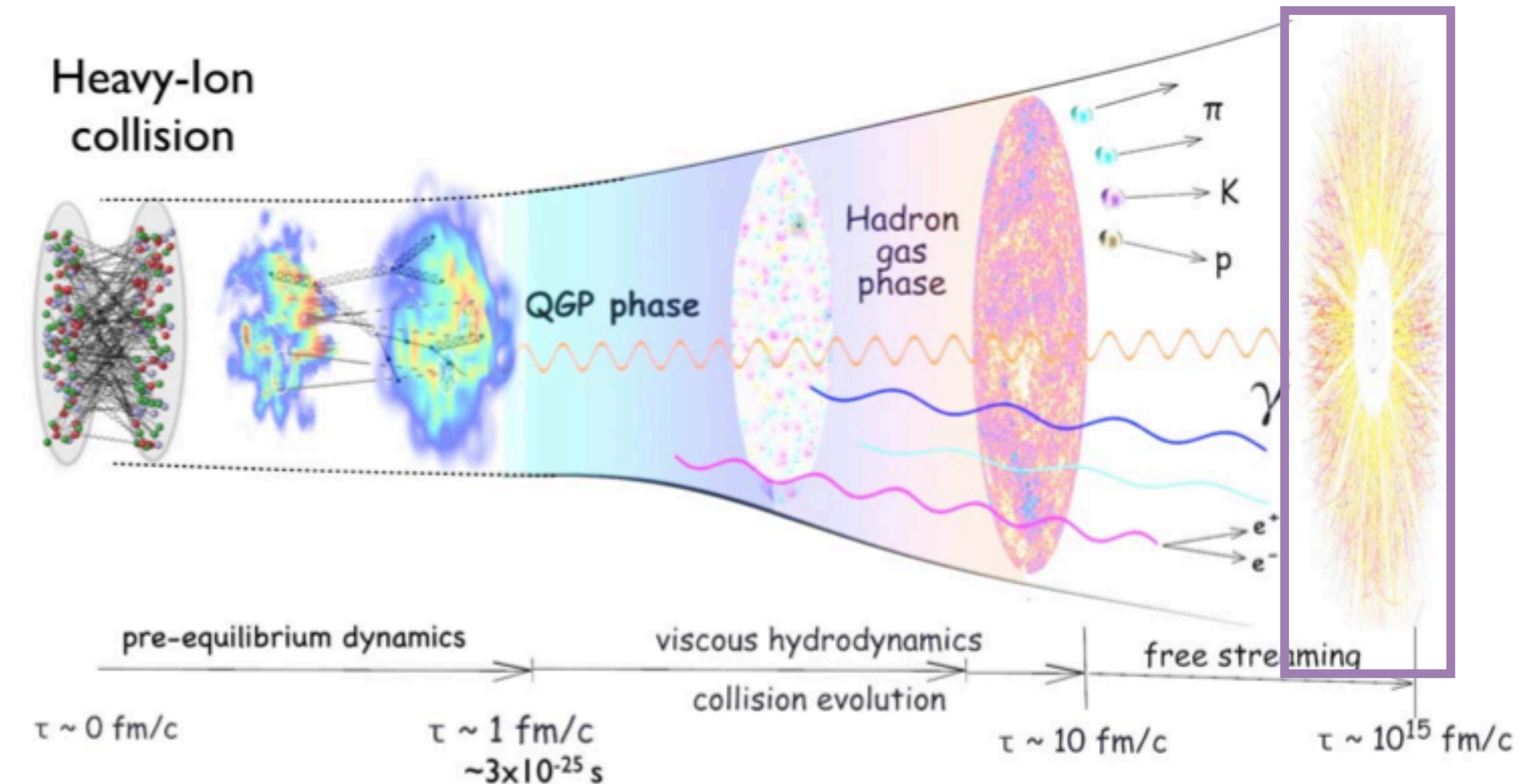
- Different QGP probes will access different wavelengths:



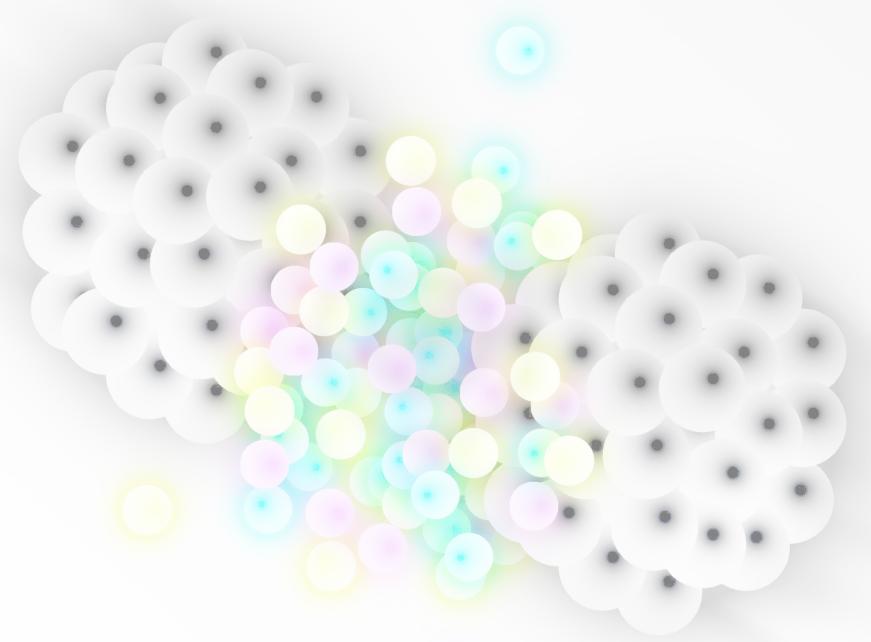
Heavy-Ion Collisions



- Different QGP probes will access different wavelengths:
 - Soft probes (bulk of the collision): low momentum particles - hydrodynamic based description

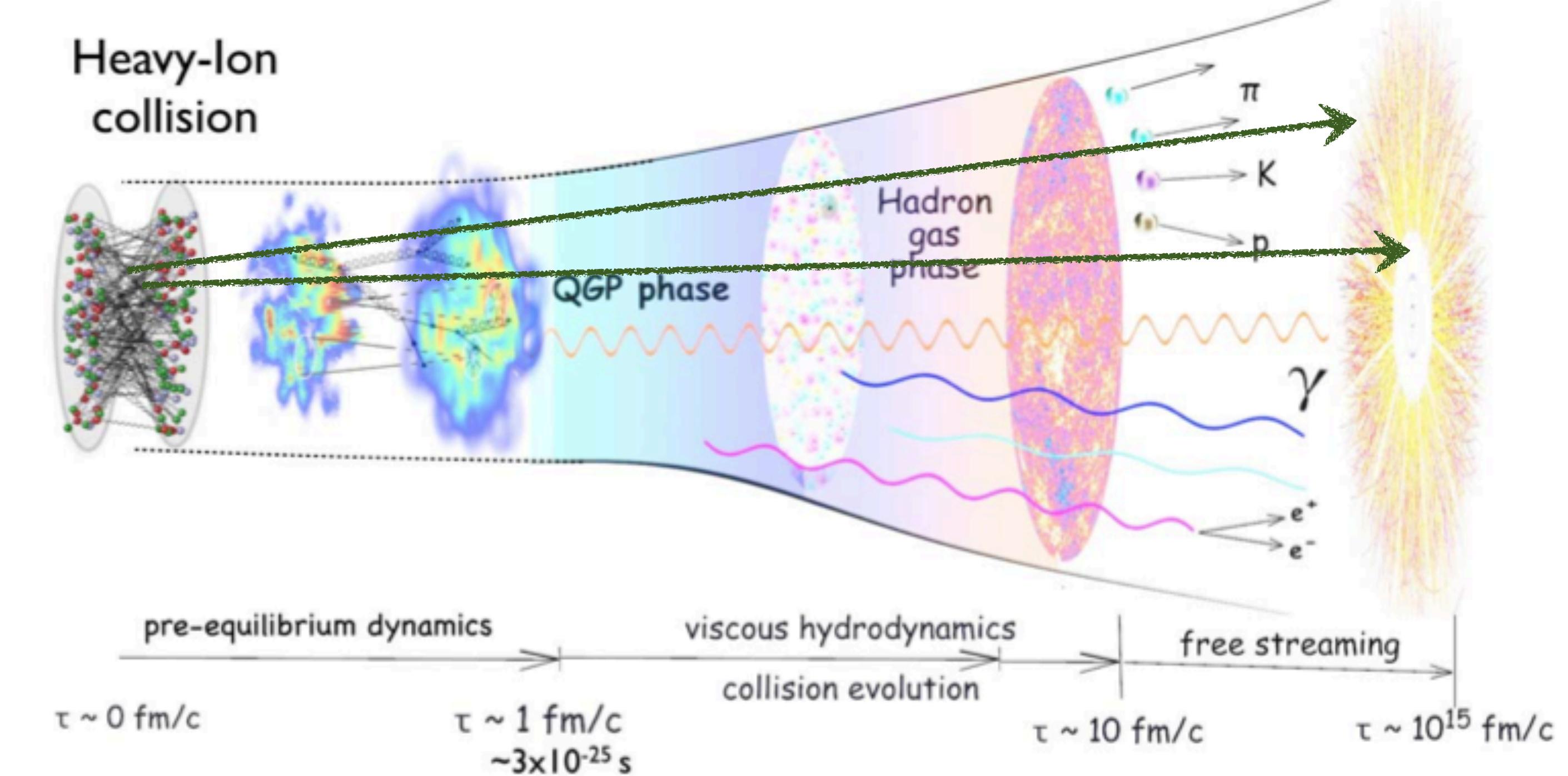
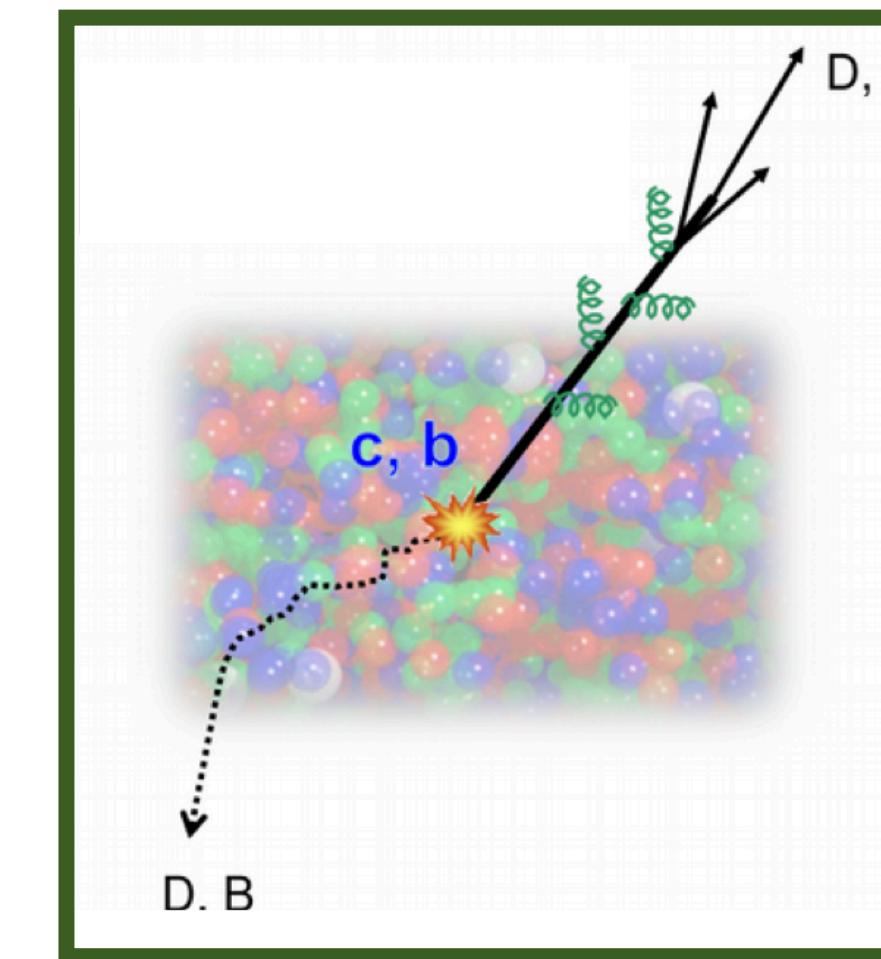
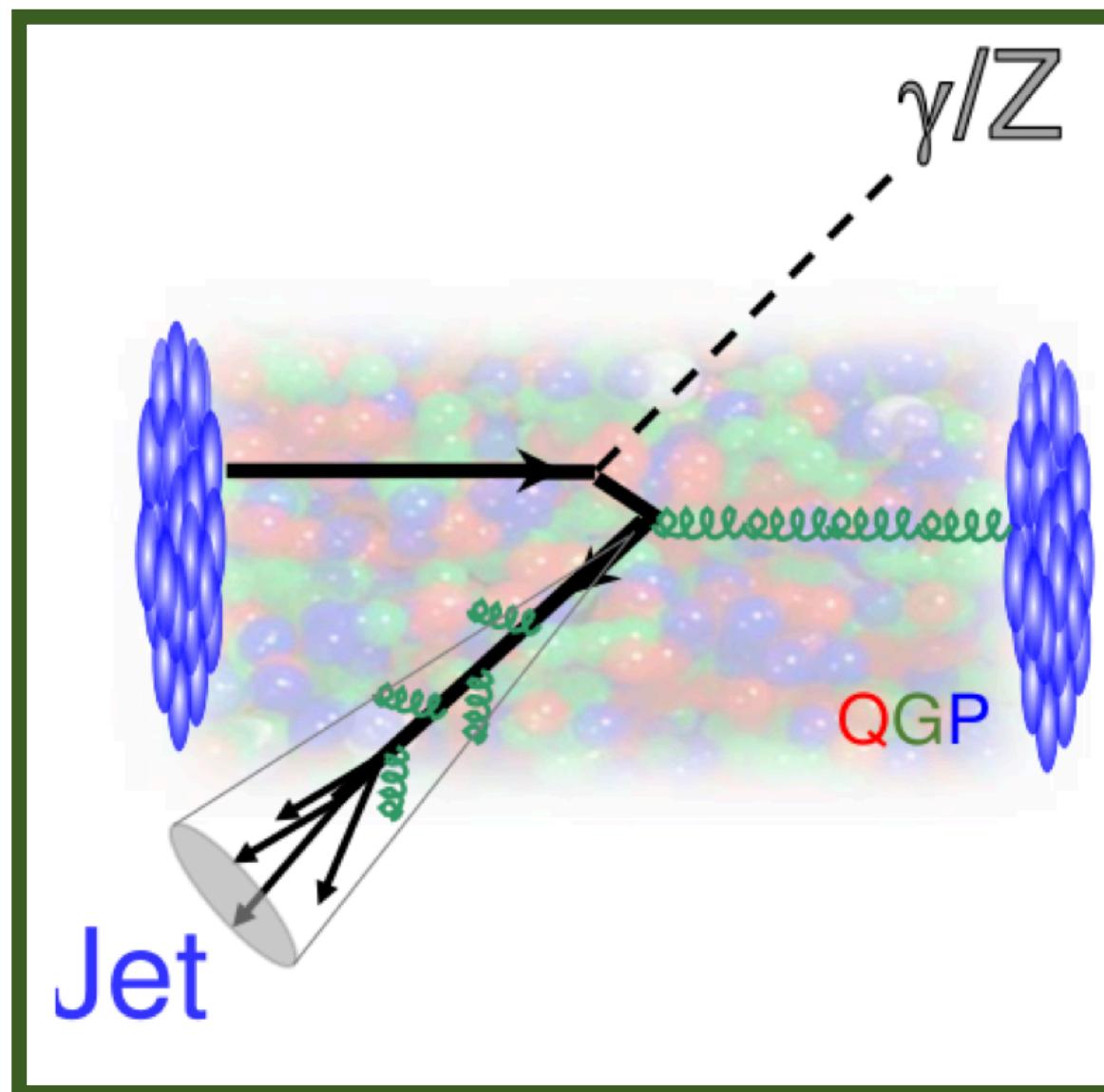


Heavy-Ion Collisions

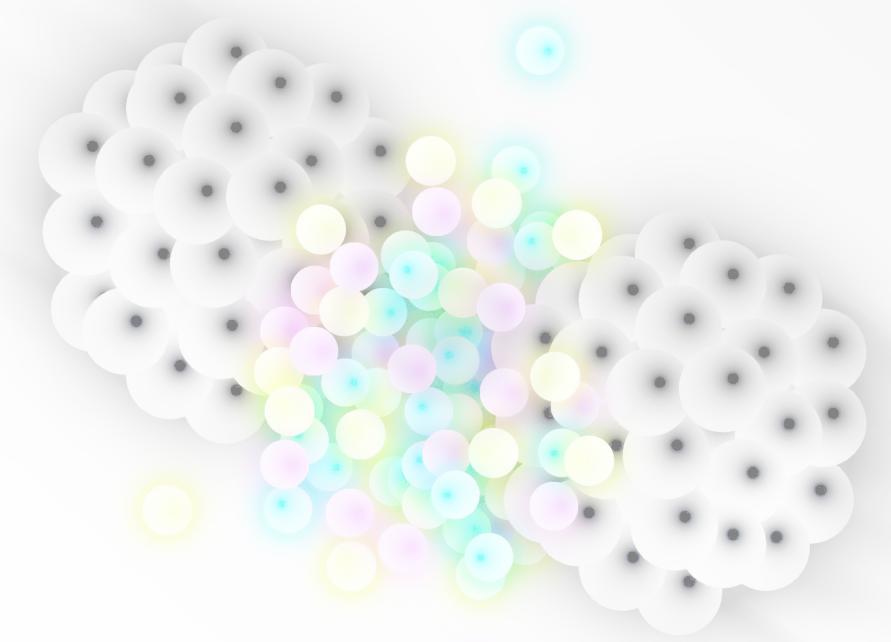


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Focus of this talk



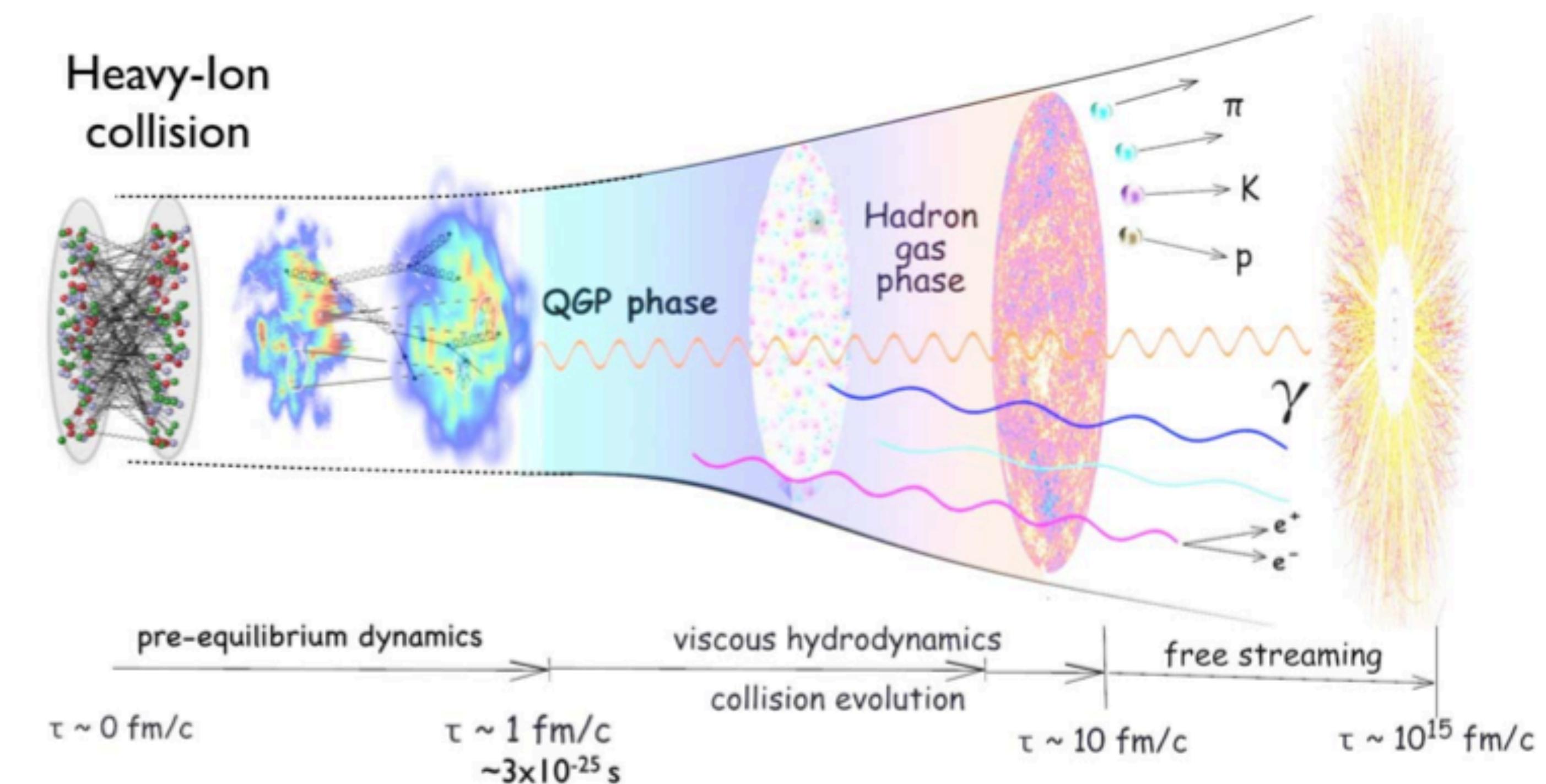
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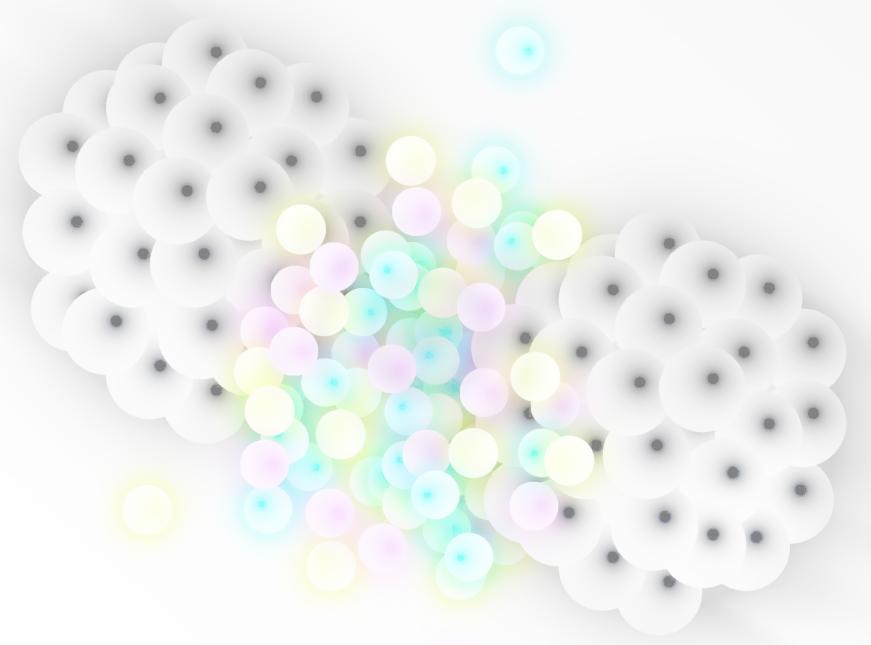
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Common difficulty: QGP is dynamically evolving system

All observables require interpretation in the framework of transport models



Heavy-Ion Collisions



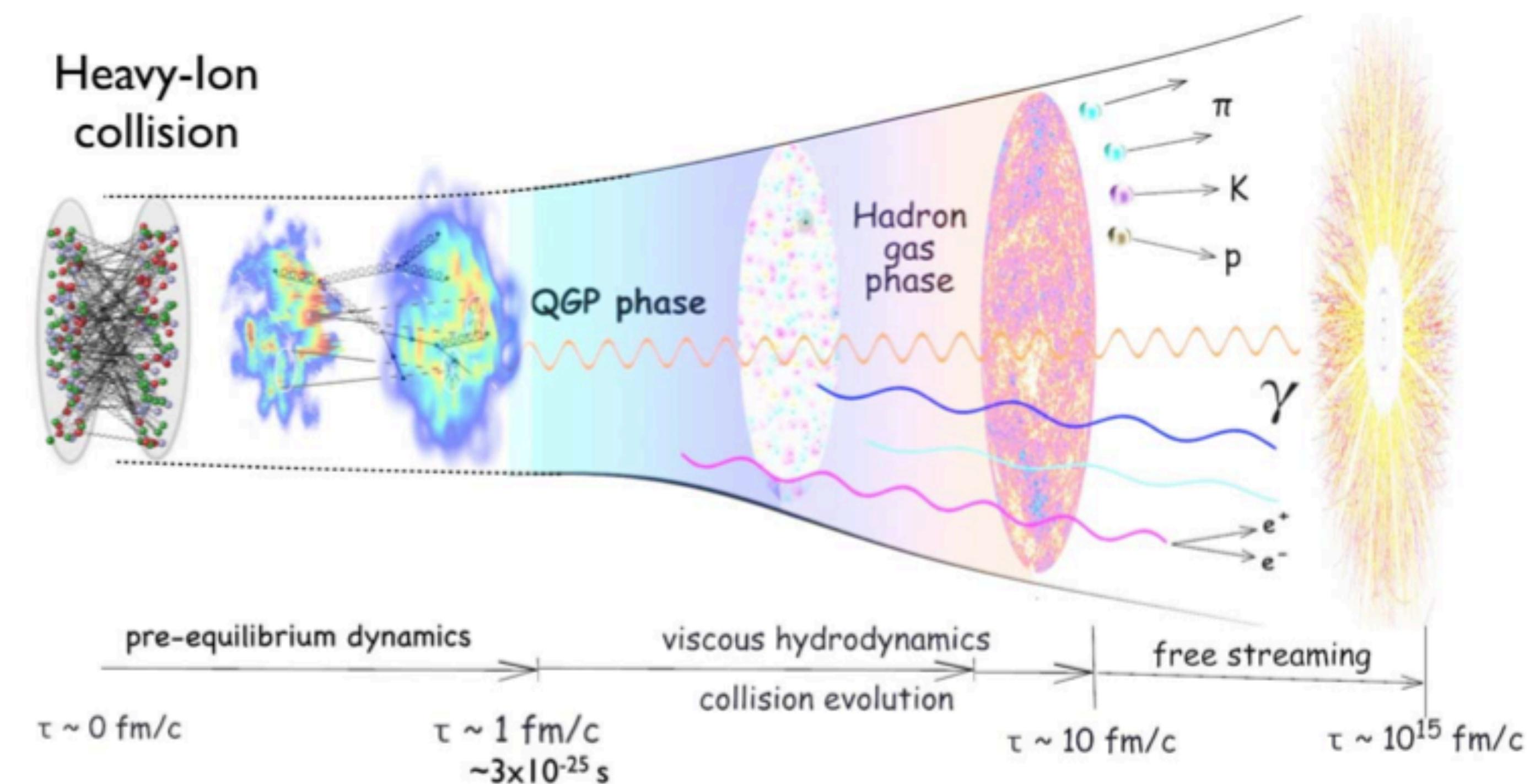
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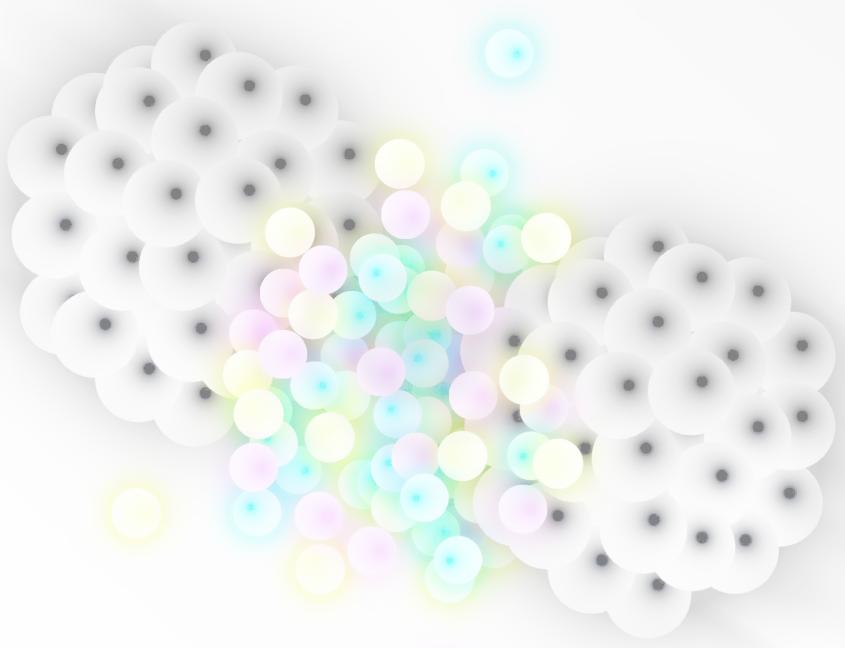
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Heavy-ion collision characterisation:

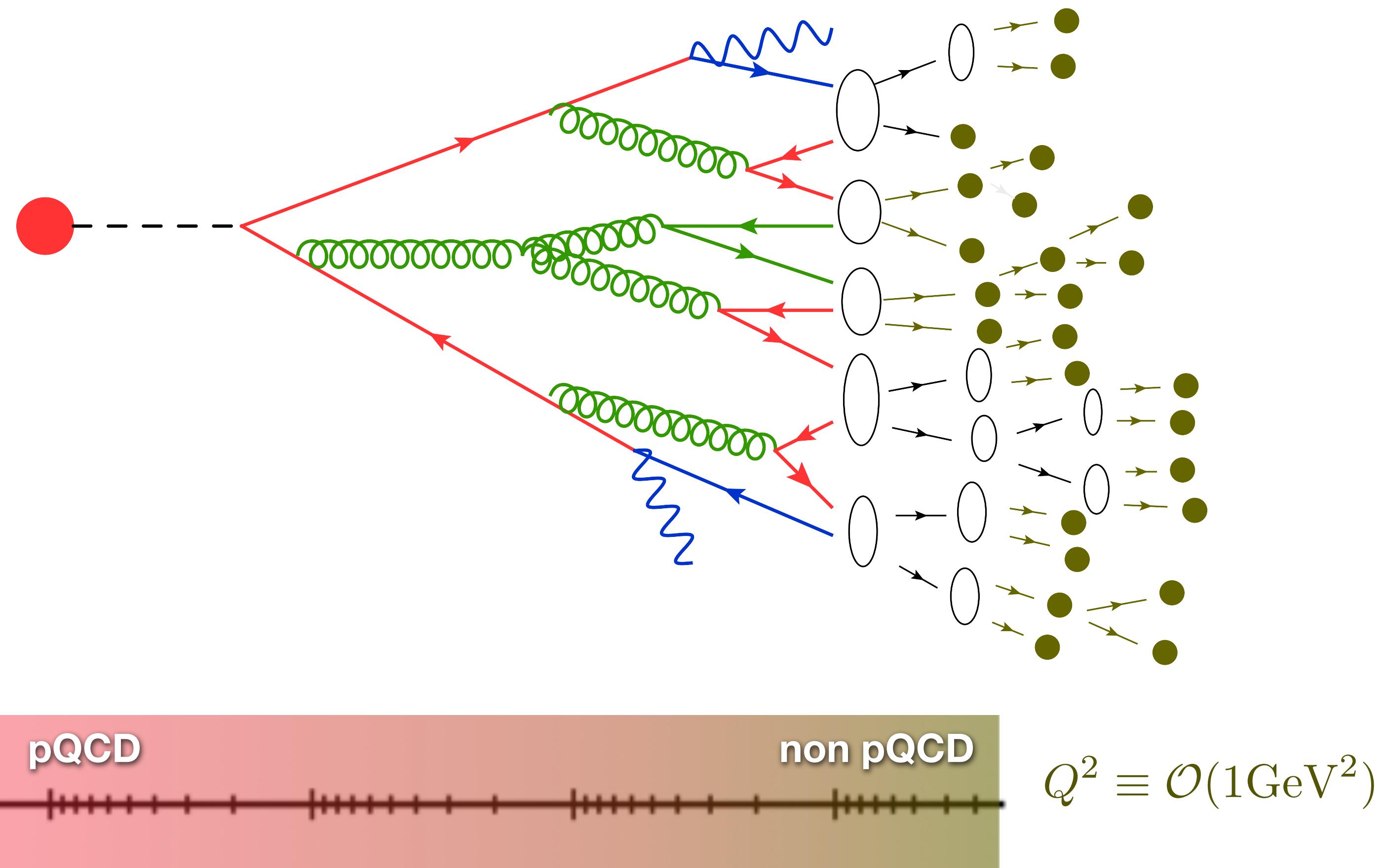
A multi-scale problem!



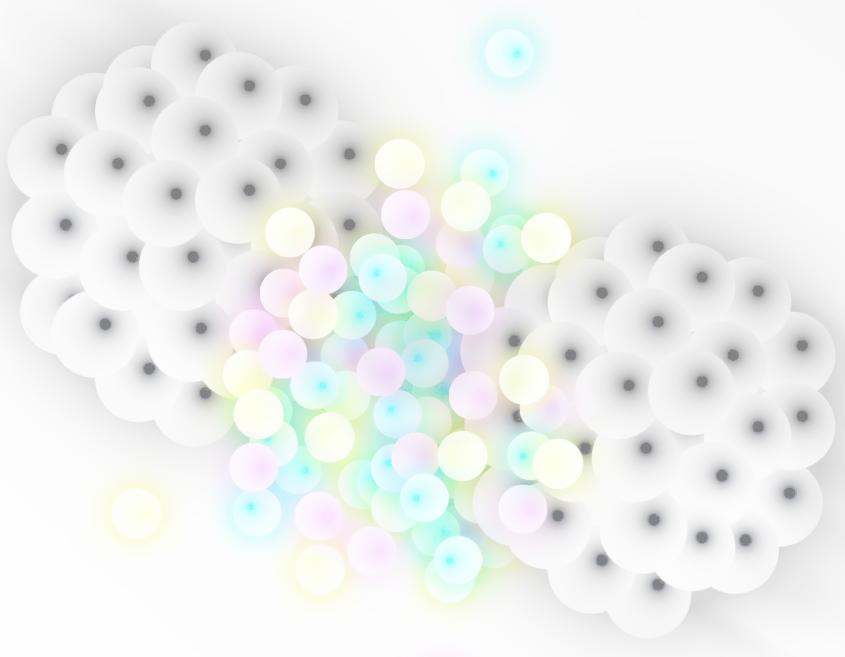
Jets in heavy-ions



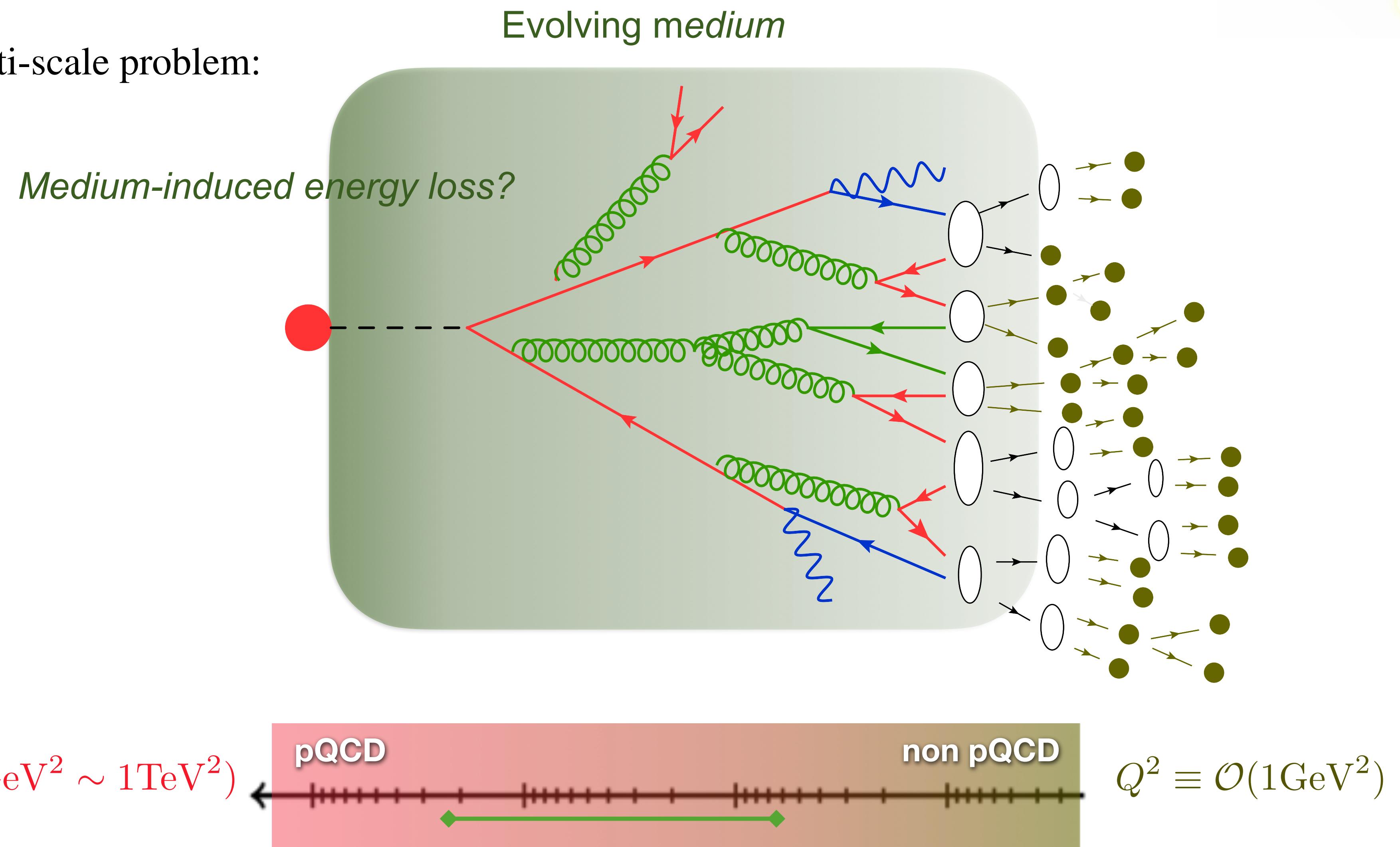
- Also a multi-scale problem:



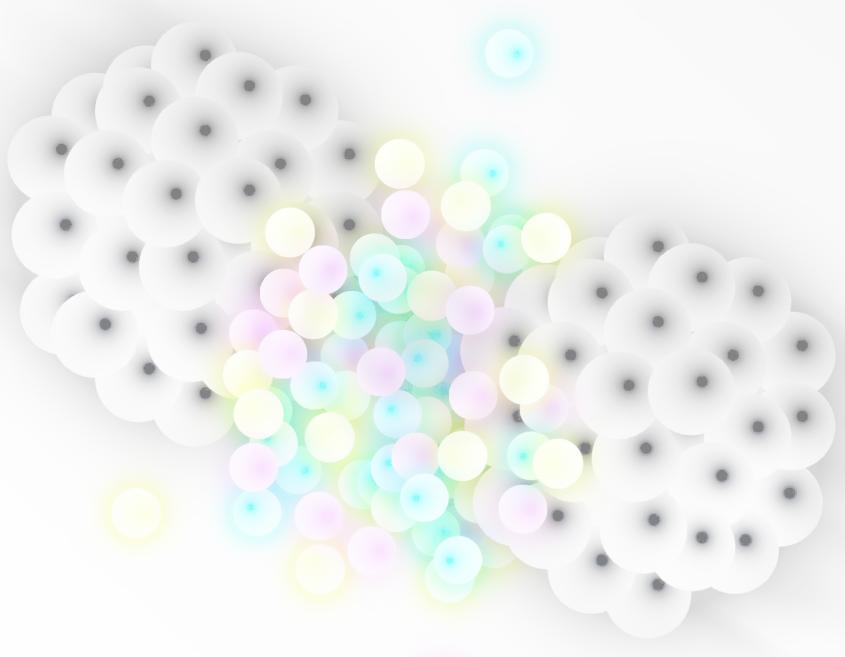
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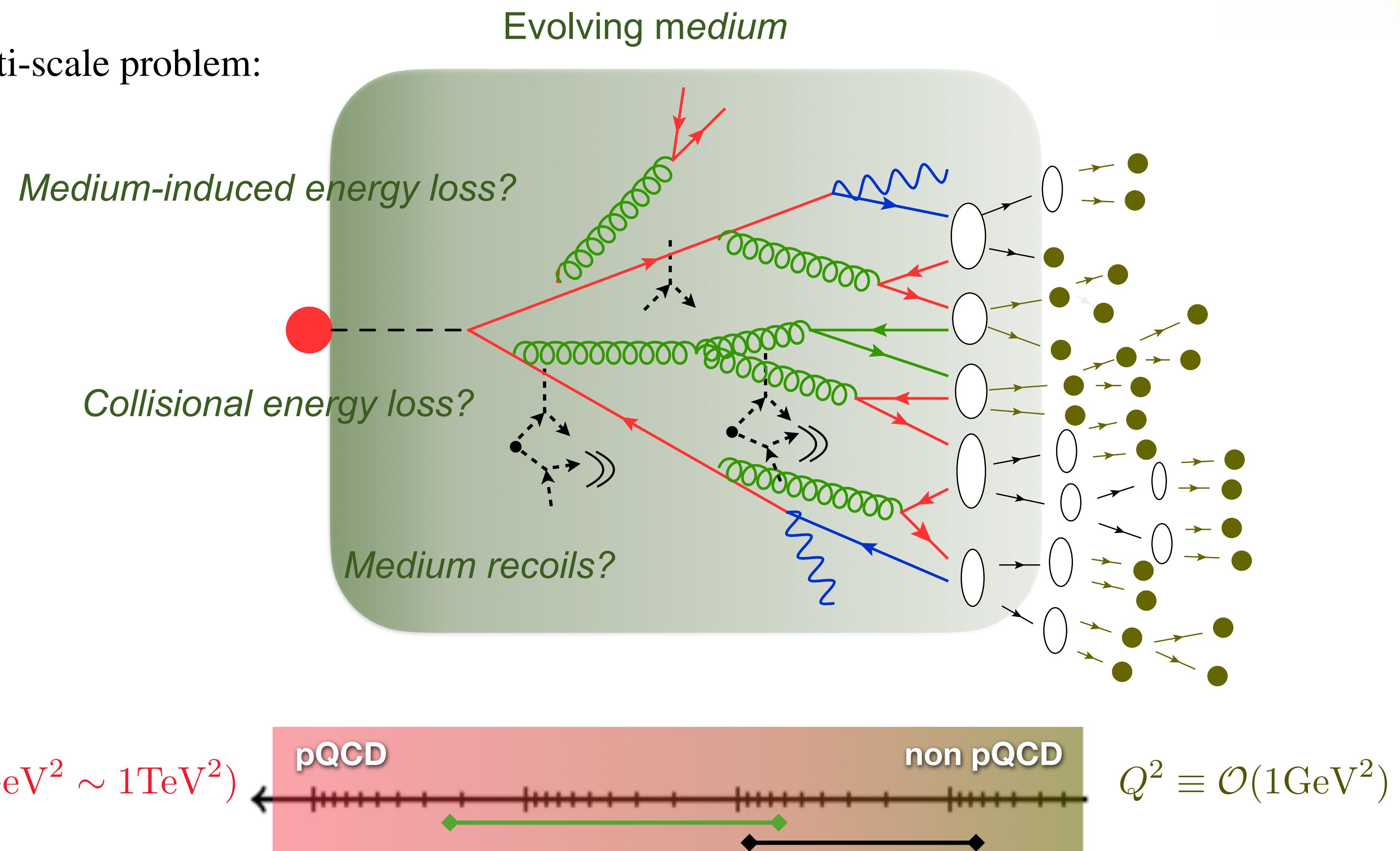
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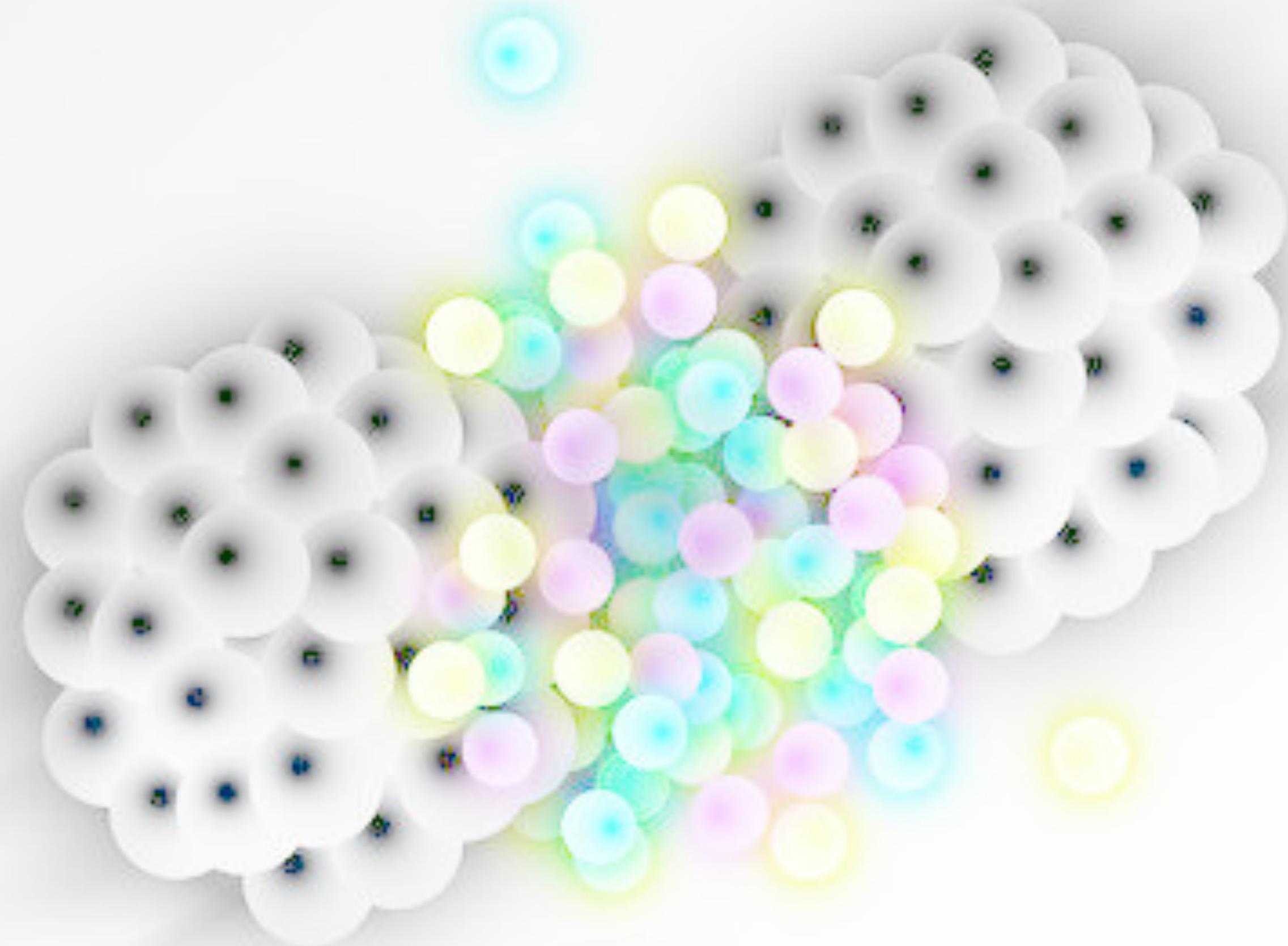
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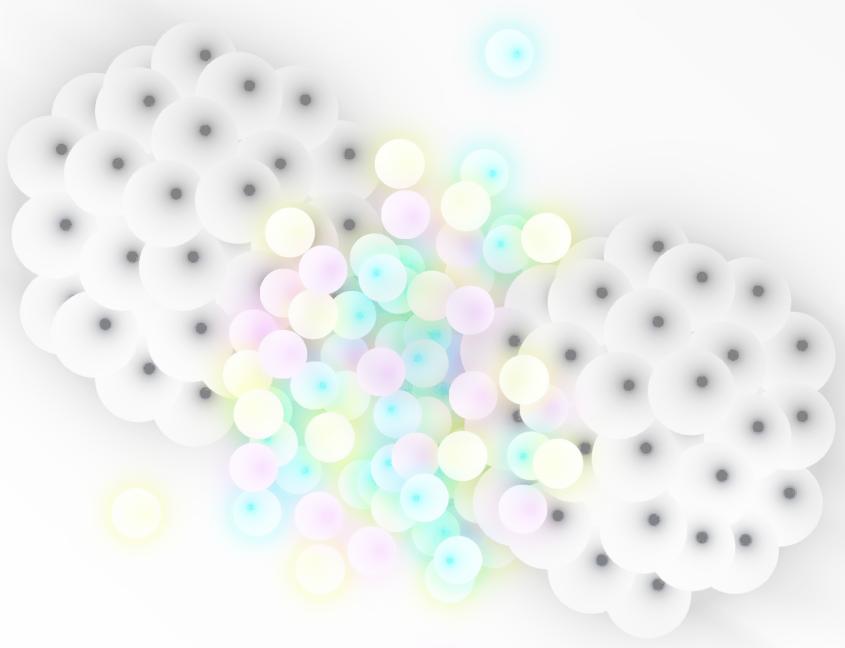
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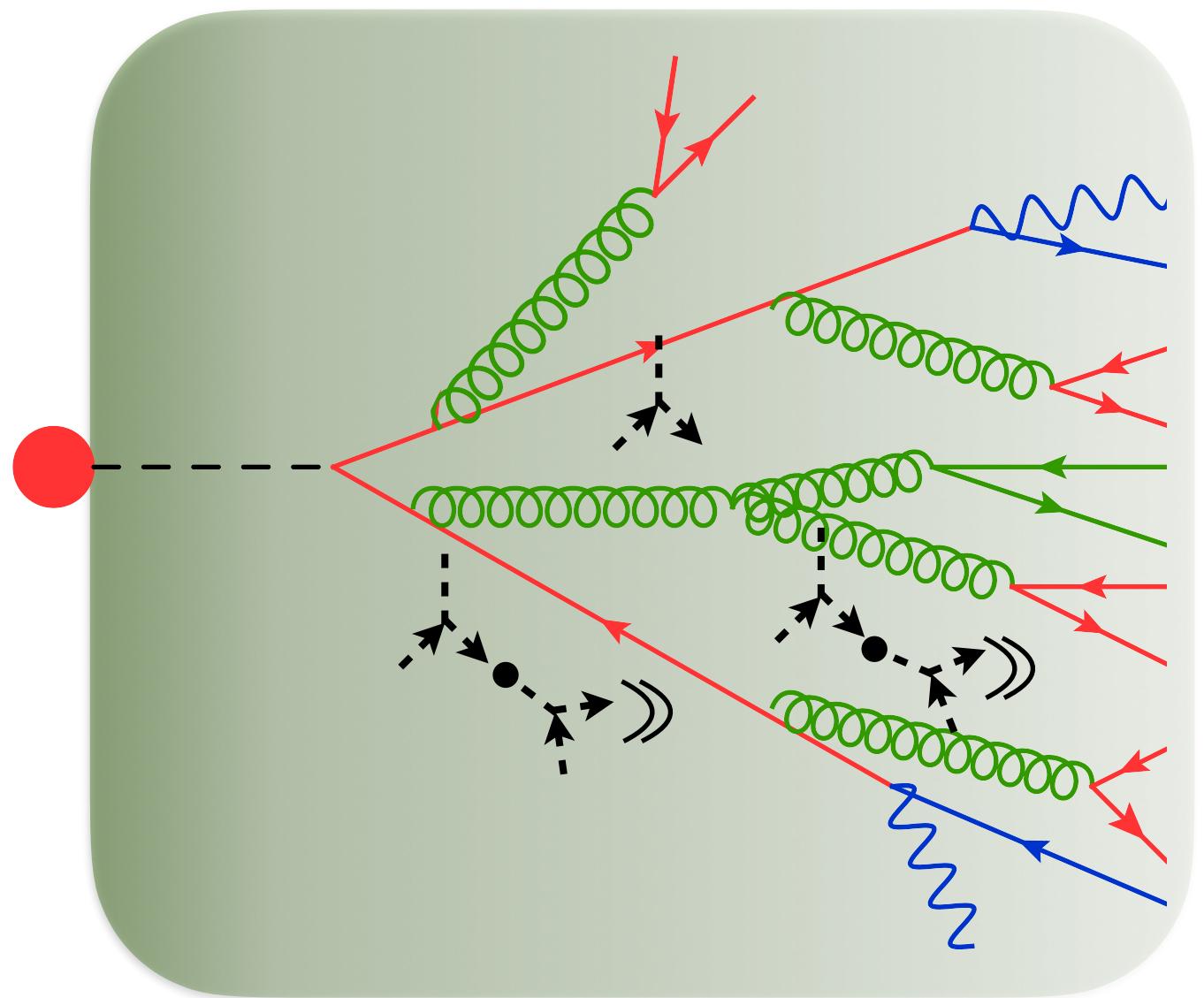
Improving theoretical control



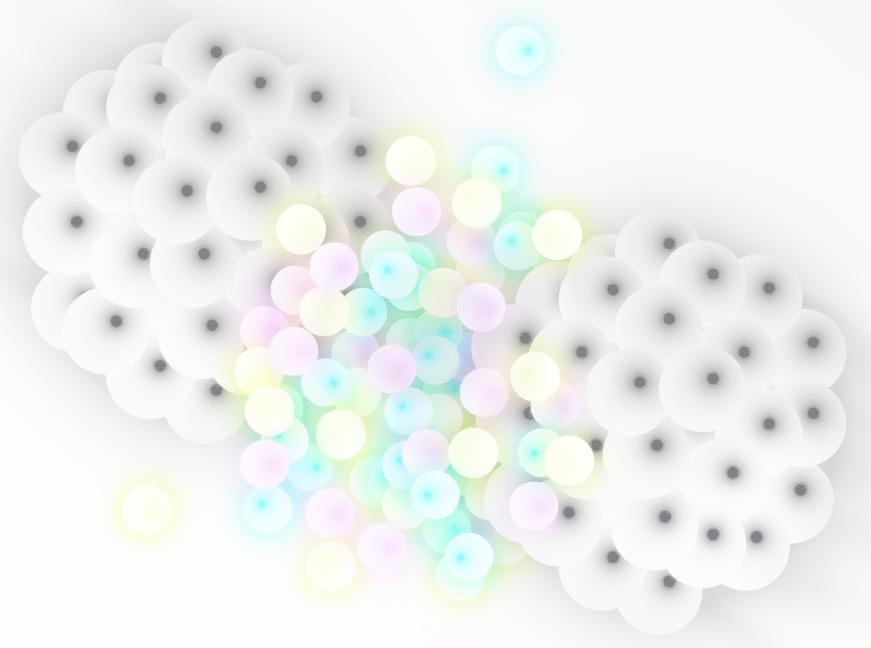
In-medium processes



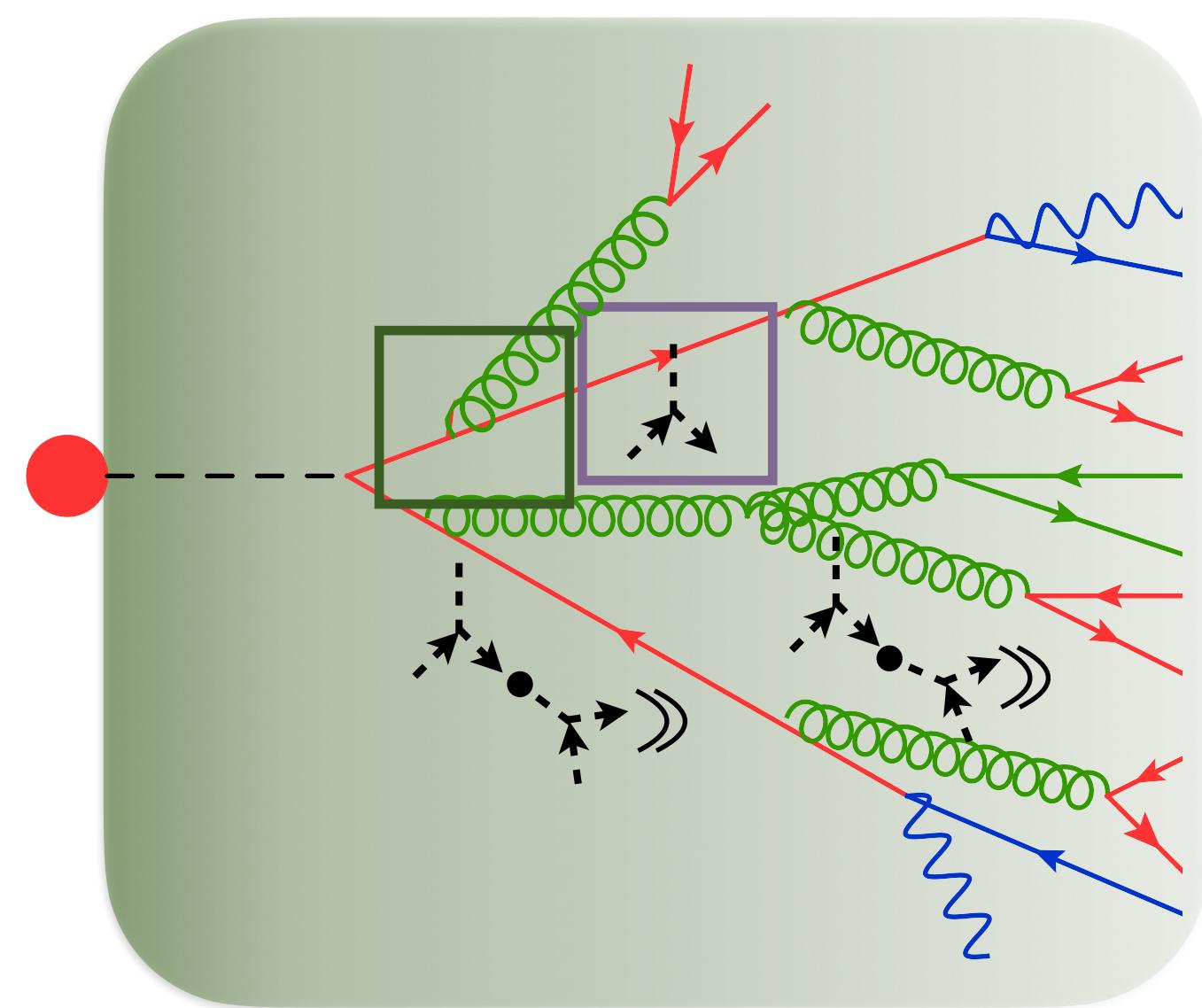
- Amount of energy loss measures transparency to the passage of a high momentum particle:
 - Towards higher accuracy in elementary building blocks of the parton shower



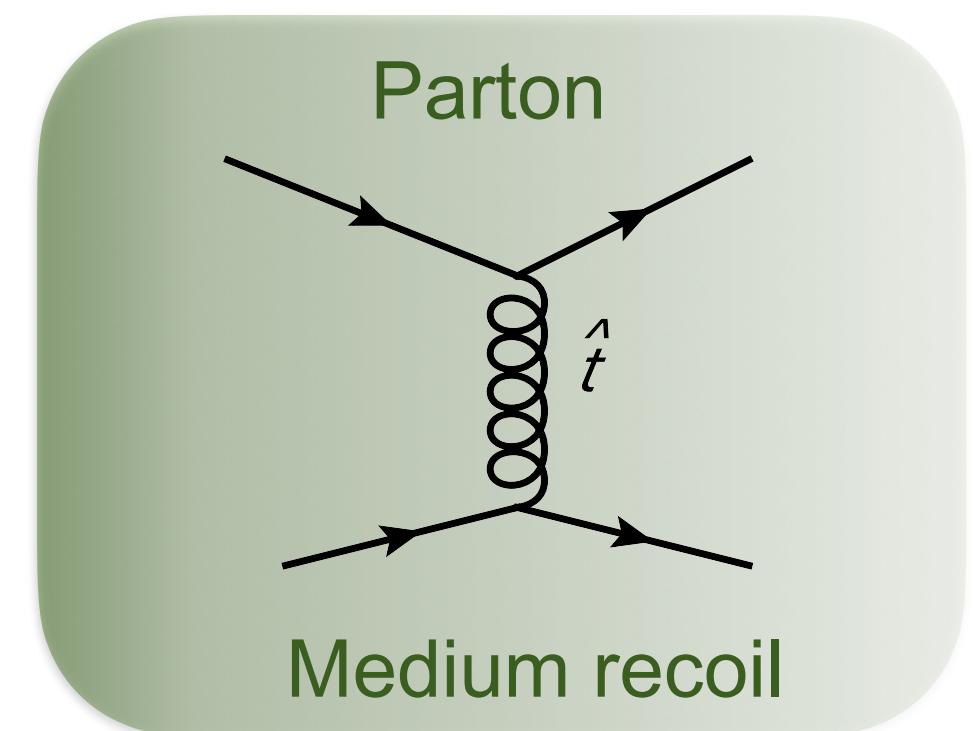
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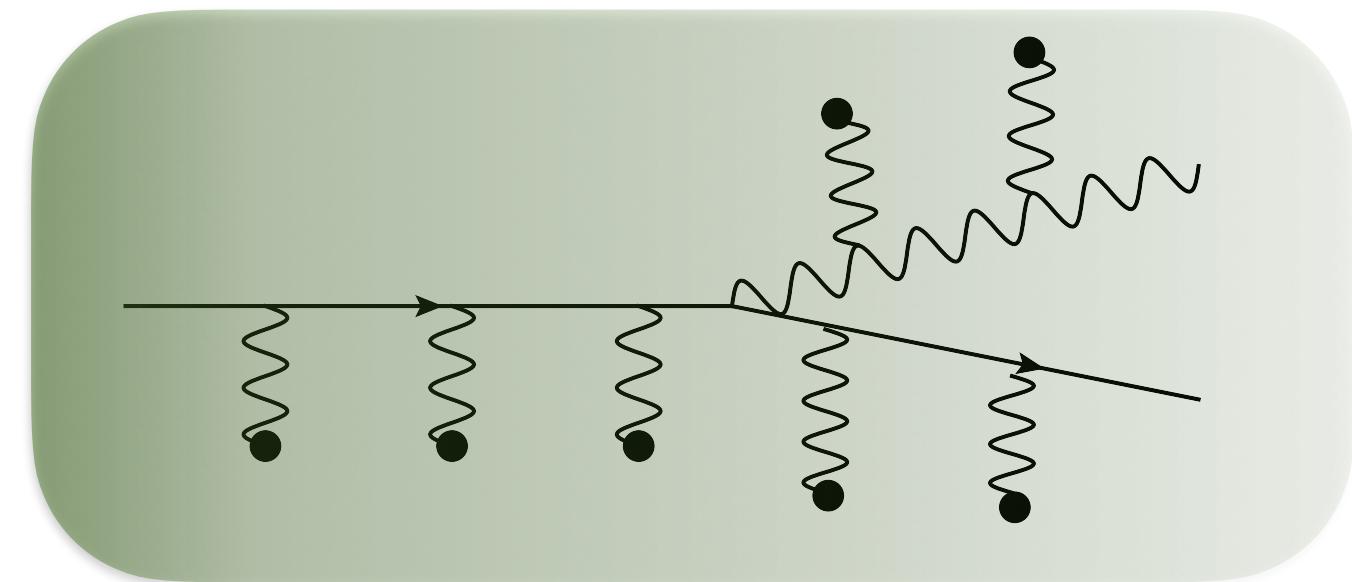
Elastic scattering processes:



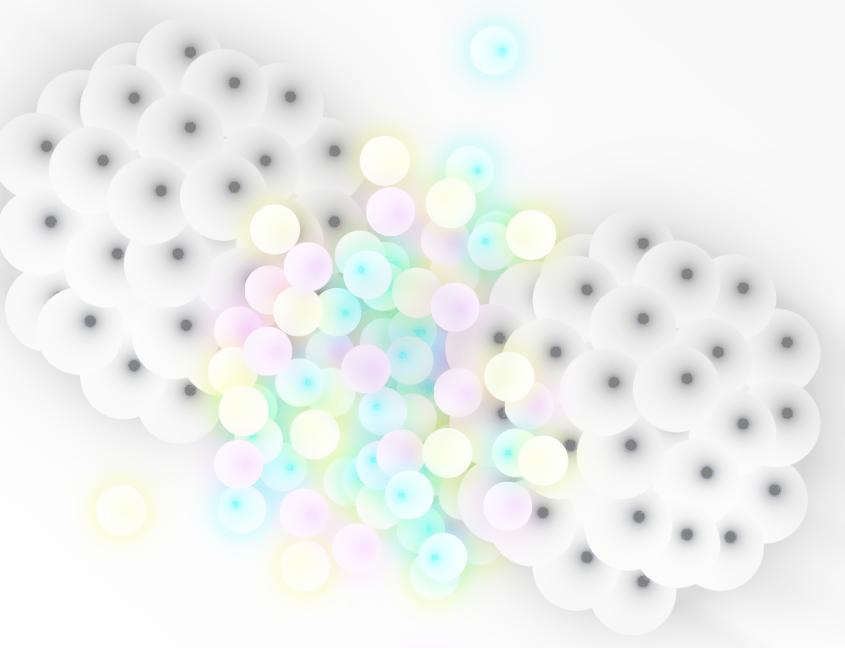
Relevant for heavy (low-energy) partons

Dominant for light (high-energy) partons

Inelastic scattering processes:

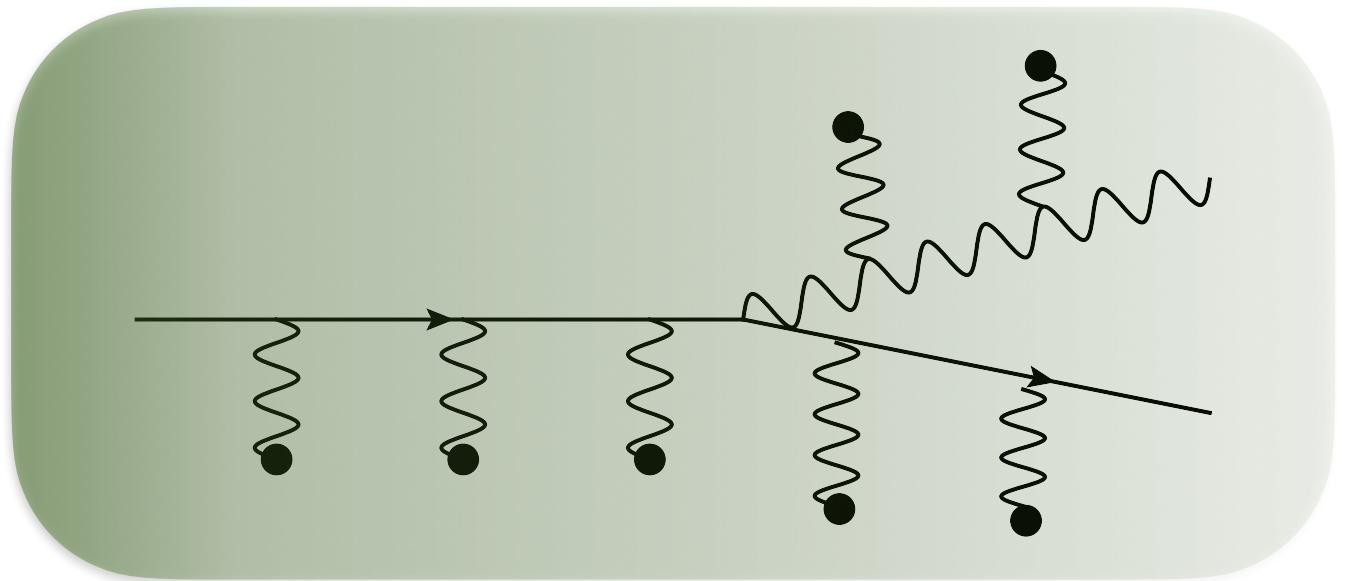


Medium-induced radiation

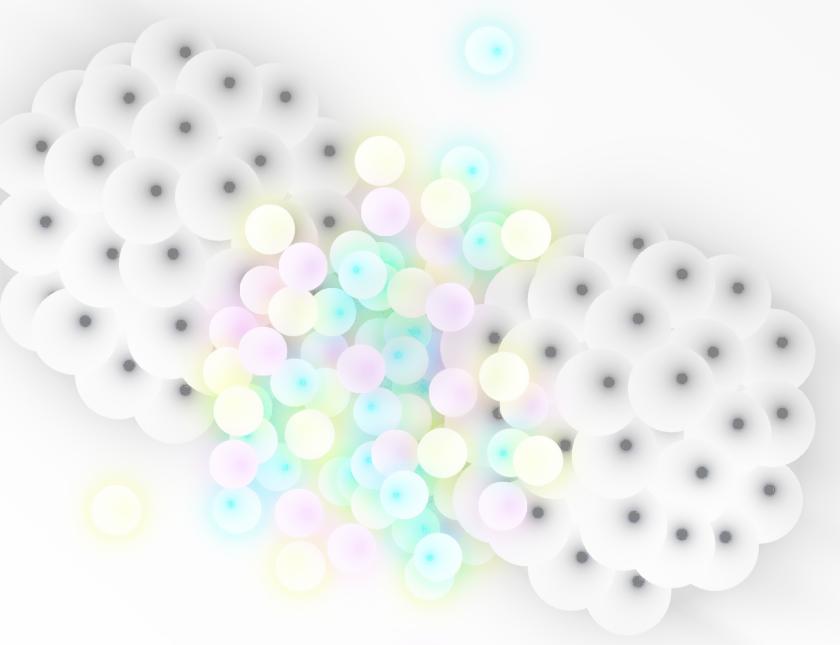


- Accumulation of momenta enhances gluon radiation:
 - Single-gluon emission spectrum:

$$\omega \frac{dI}{d\omega d^2\mathbf{k}} = \frac{2\alpha_s C_R}{(2\pi)^2 \omega^2} \text{Re} \int_0^\infty dt' \int_0^{t'} dt \int_{\mathbf{p}, \mathbf{q}} \mathbf{p} \cdot \mathbf{q} \tilde{\mathcal{K}}(t', \mathbf{q}; t, \mathbf{p}) P(\infty, \mathbf{k}; t', \mathbf{q})$$

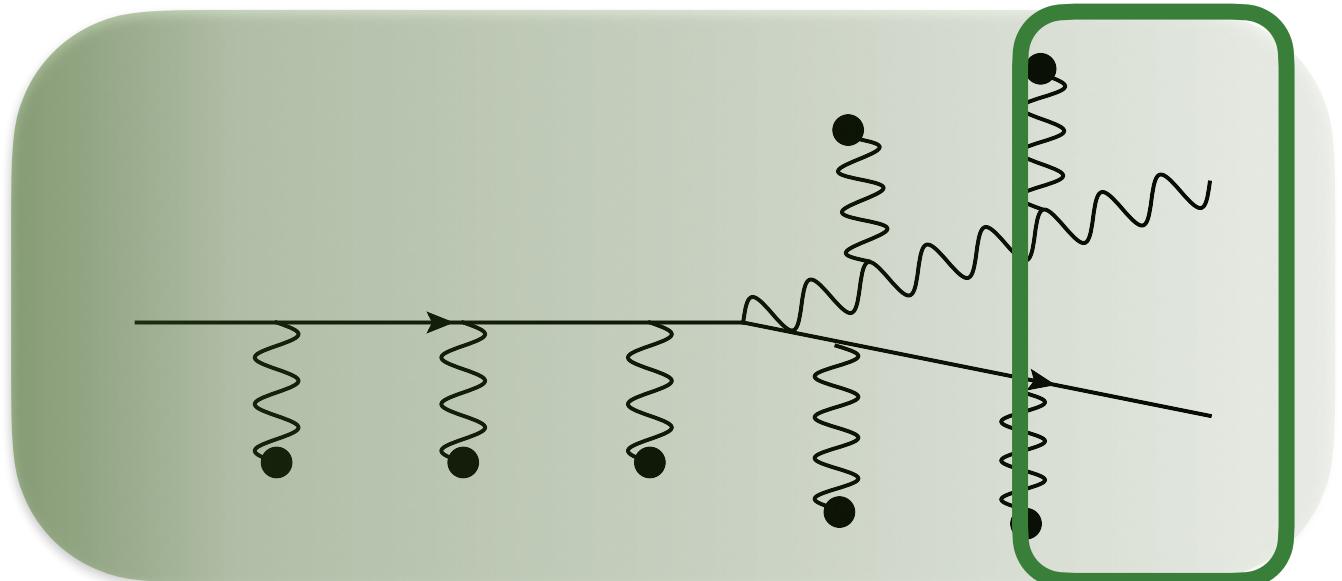


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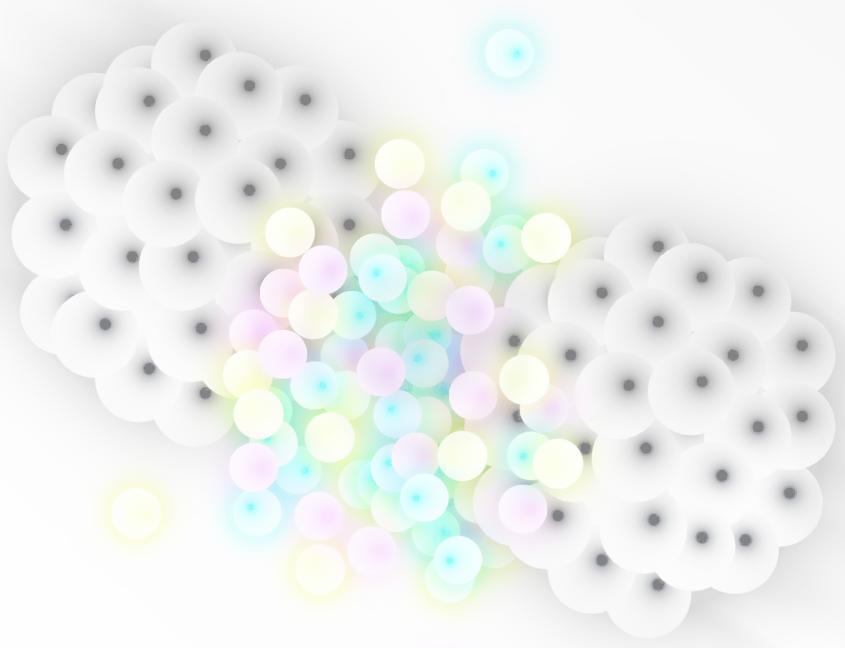
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Momentum Broadening:

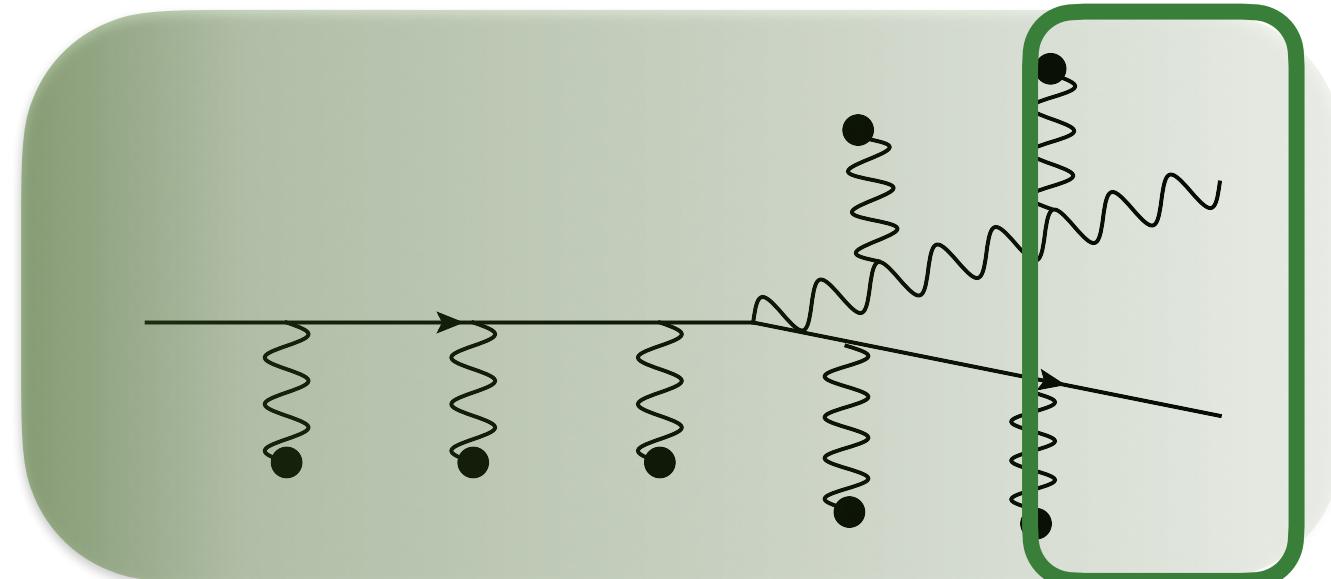
$$\mathcal{P}(t'', \mathbf{k}; t', \mathbf{q}) \equiv \int d^2 z e^{-i(\mathbf{k}-\mathbf{q}) \cdot z} \exp \left\{ -\frac{1}{2} \int_{t'}^{t''} ds n(s) \sigma(z) \right\}$$

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Density of scattering centres:

$$n(x_+) = \int dx_{i+} \delta(x_+ - x_{i+}).$$

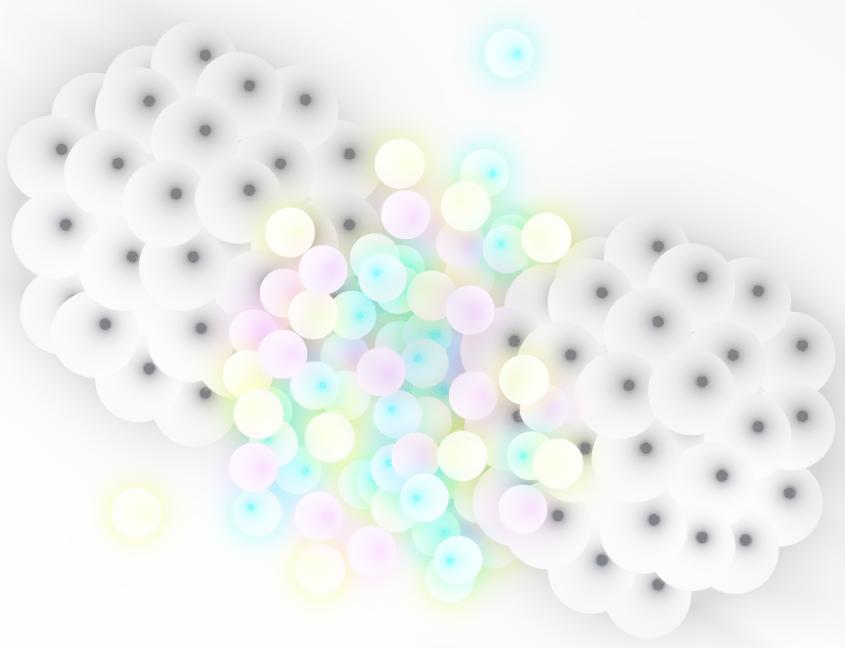
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Dipole cross-section (collision rate):

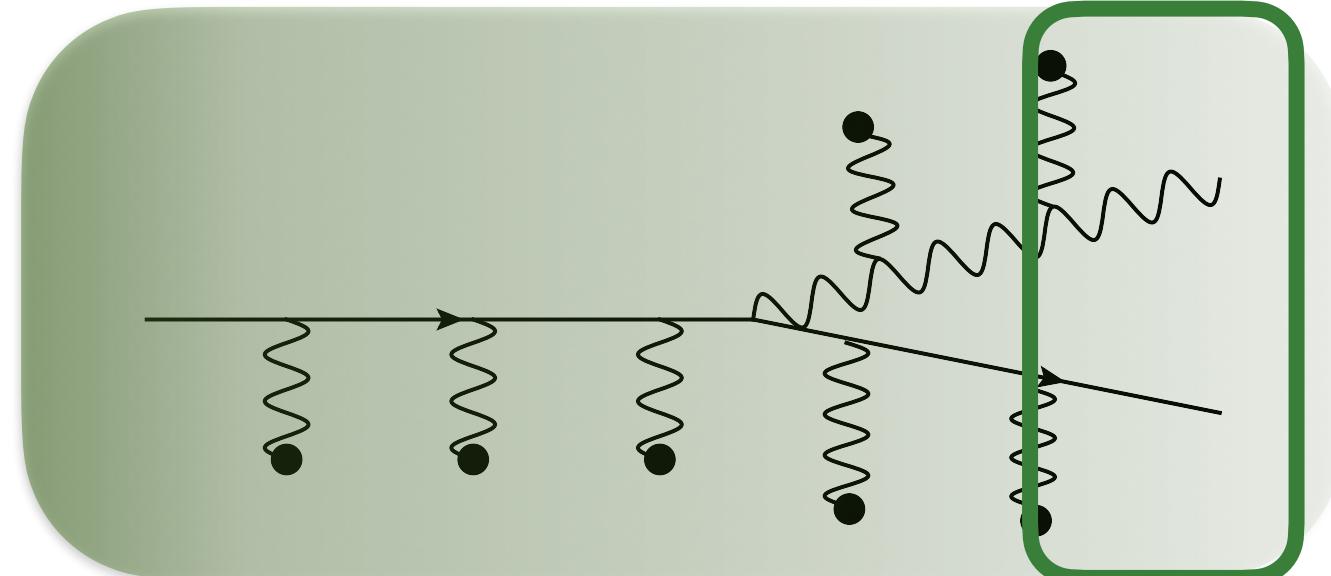
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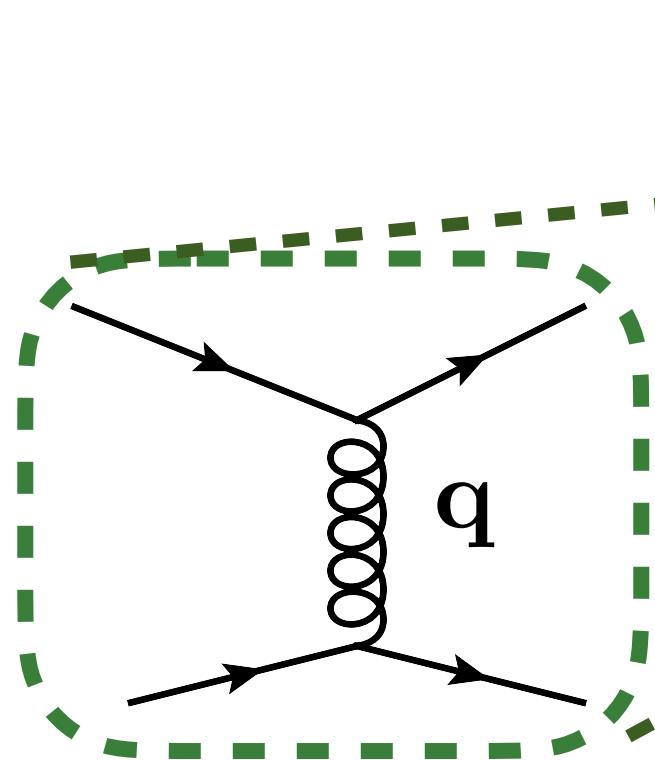


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Parton-medium
interaction

Medium-induced radiation

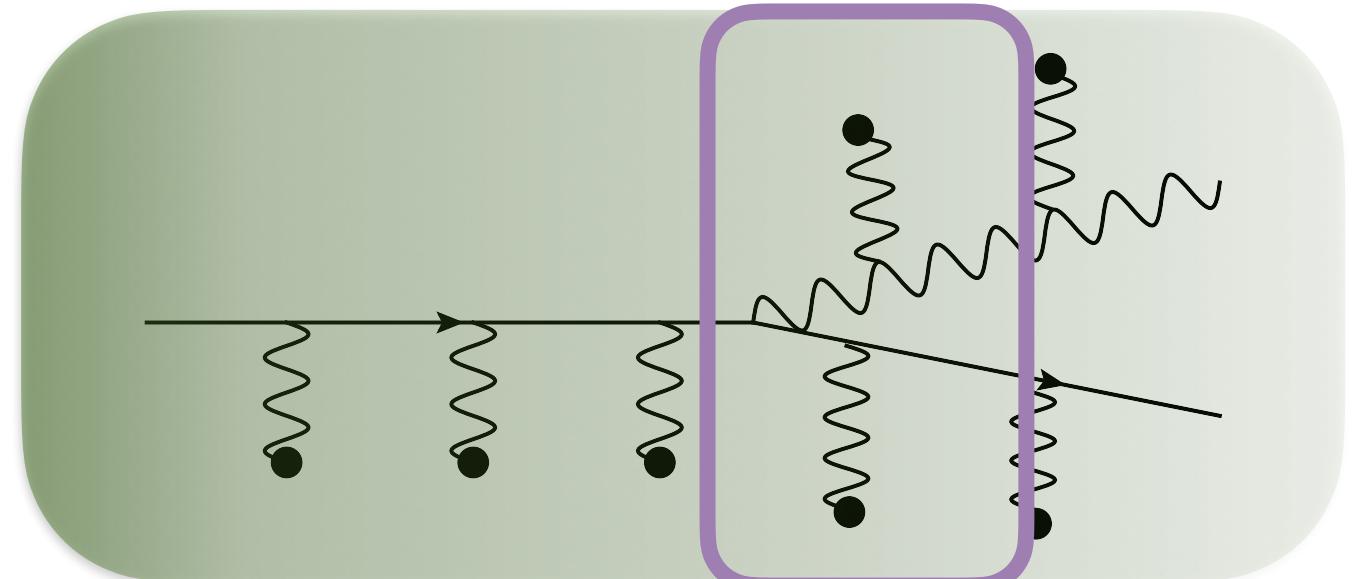


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Emission Kernel:

$$\begin{aligned} \mathcal{K}(t', \mathbf{z}; t, \mathbf{y}) &\equiv \int_{\mathbf{p}, \mathbf{q}} e^{i(\mathbf{q} \cdot \mathbf{z} - \mathbf{p} \cdot \mathbf{y})} \tilde{\mathcal{K}}(t', \mathbf{q}; t, \mathbf{p}) \\ &= \int_{\mathbf{r}(t) = \mathbf{y}}^{\mathbf{r}(t') = \mathbf{z}} d\mathbf{r} \exp \left[\int_t^{t'} ds \left(\frac{i\omega}{2} \dot{\mathbf{r}}^2 - \frac{1}{2} n(s) \sigma(\mathbf{r}) \right) \right] \end{aligned}$$



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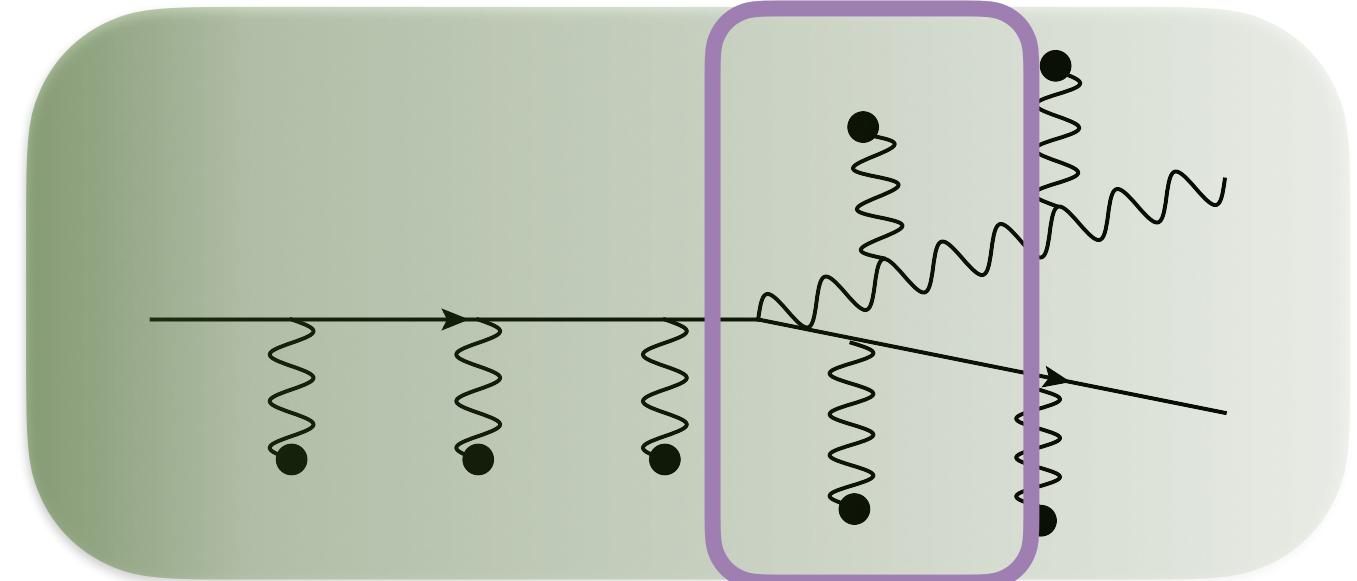


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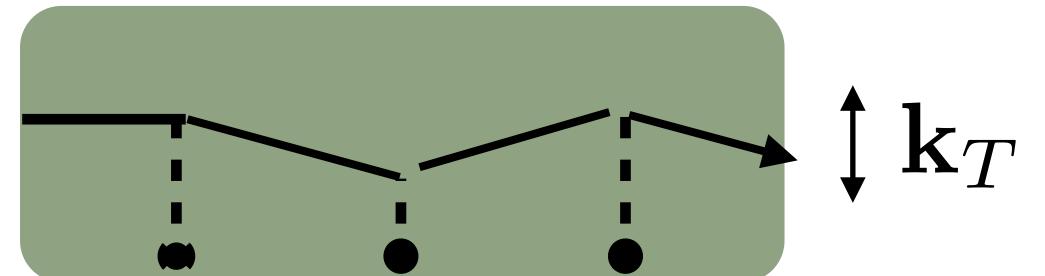
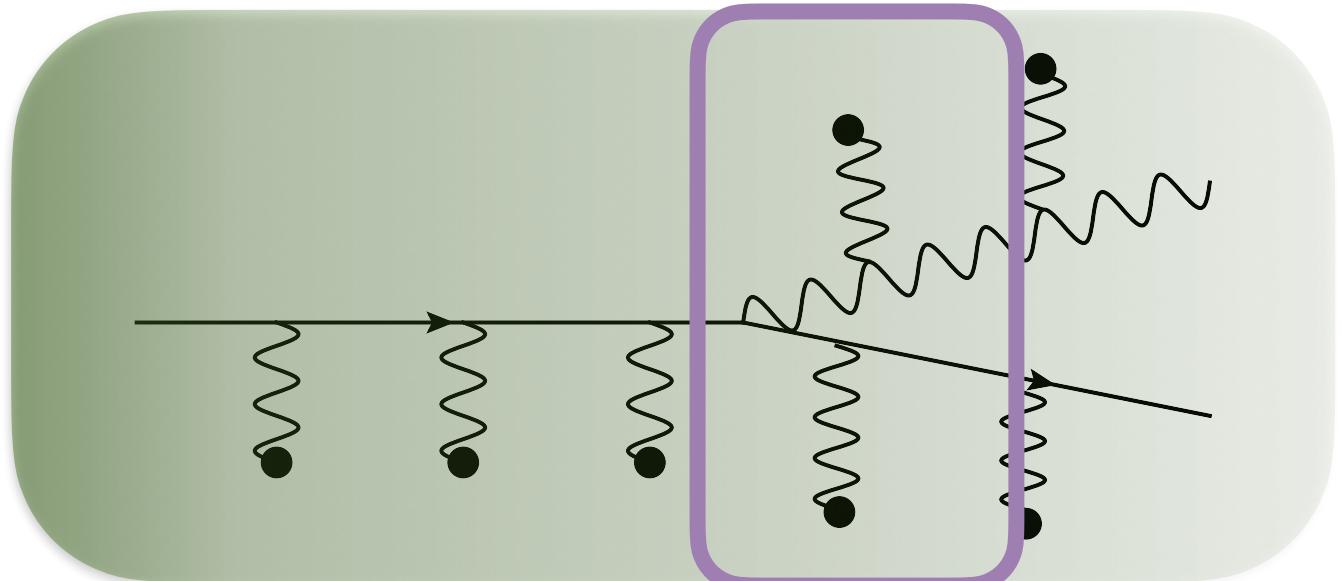
$$\sigma(\mathbf{r}) = \int_{\mathbf{q}} V(\mathbf{q}) (1 - e^{i\mathbf{q}\mathbf{r}})$$

Solution to the path integral (for an arbitrary potential) poses significant technical challenges...

Medium-induced radiation



- Accumulation of momenta enhances gluon radiation:
 - In addition to energy loss, parton also undergoes transverse momentum diffusion **See also Sievert talk (Tue)**
 - Medium-induced transverse momentum broadening



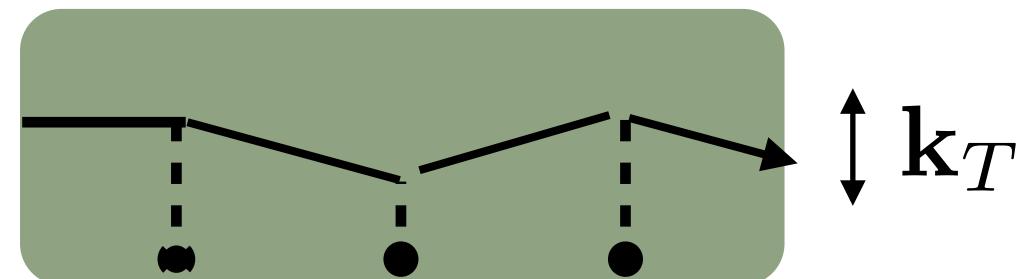
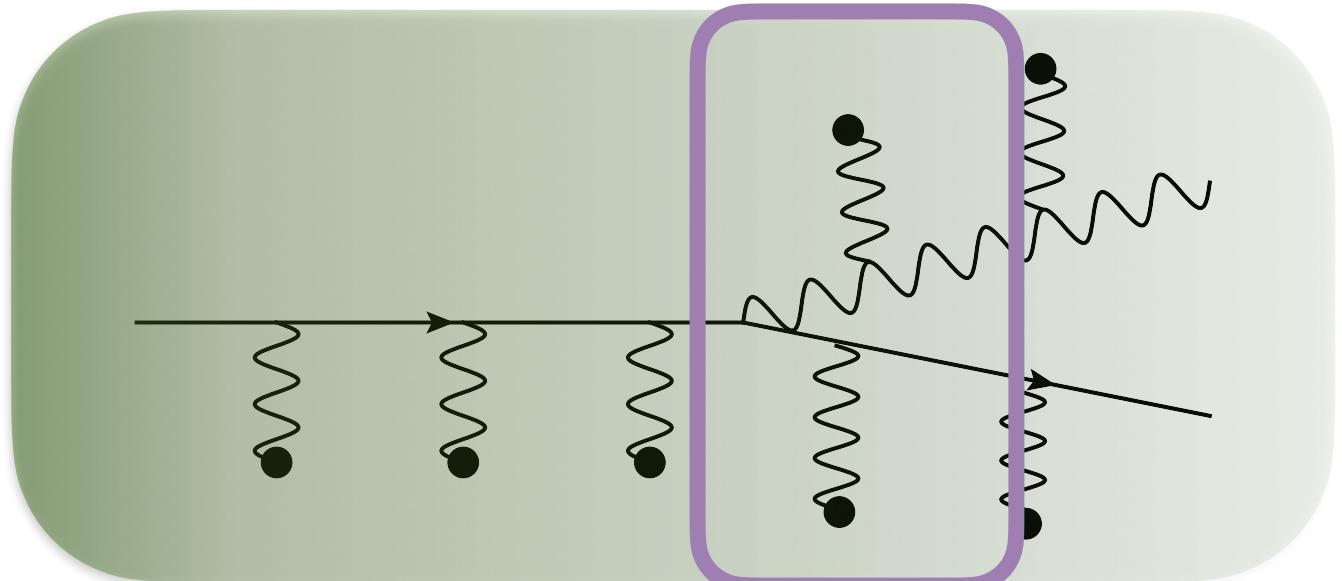
Transport coefficient:

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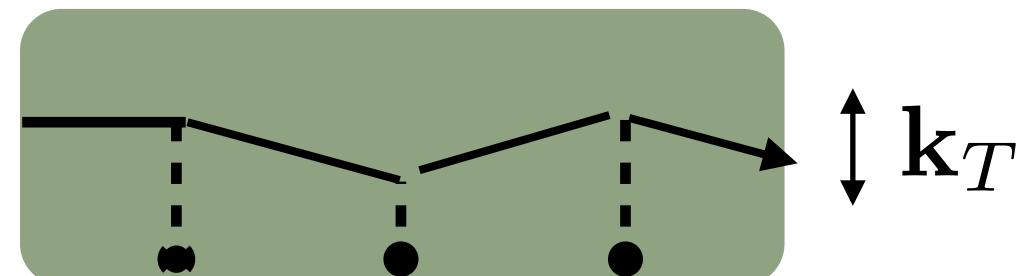
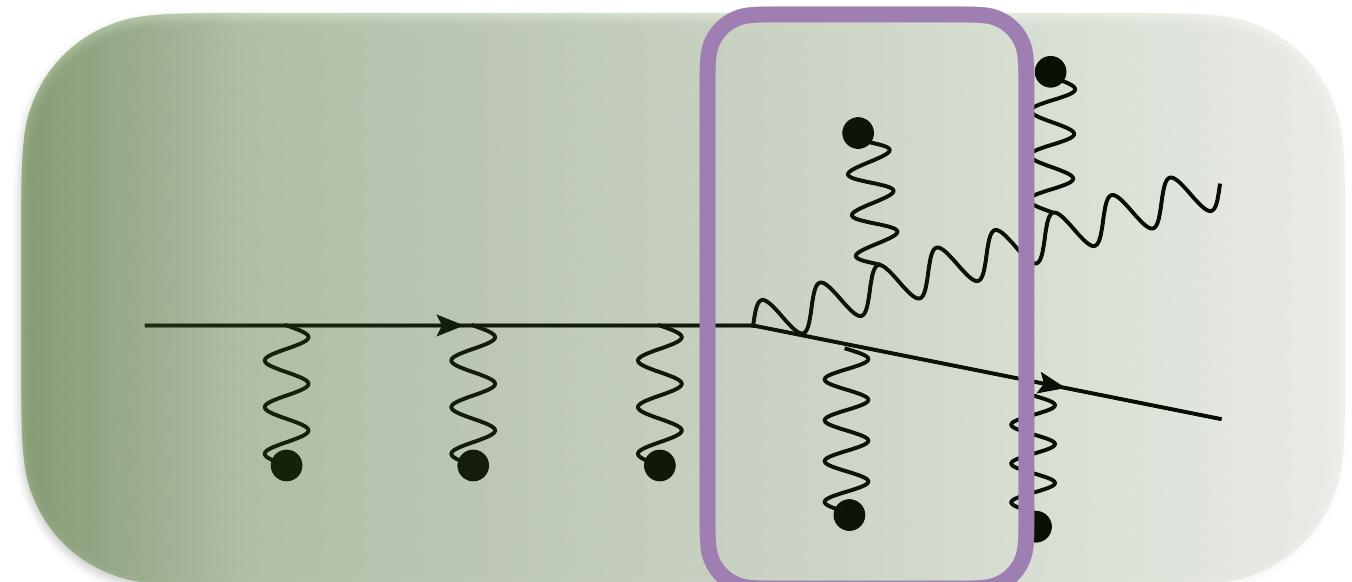
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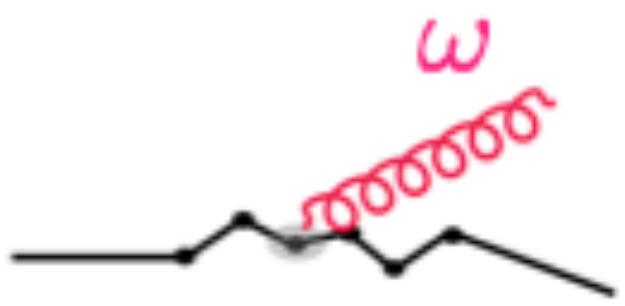
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Medium-interactions per emission?



Multiple-soft scattering



Single-hard emission

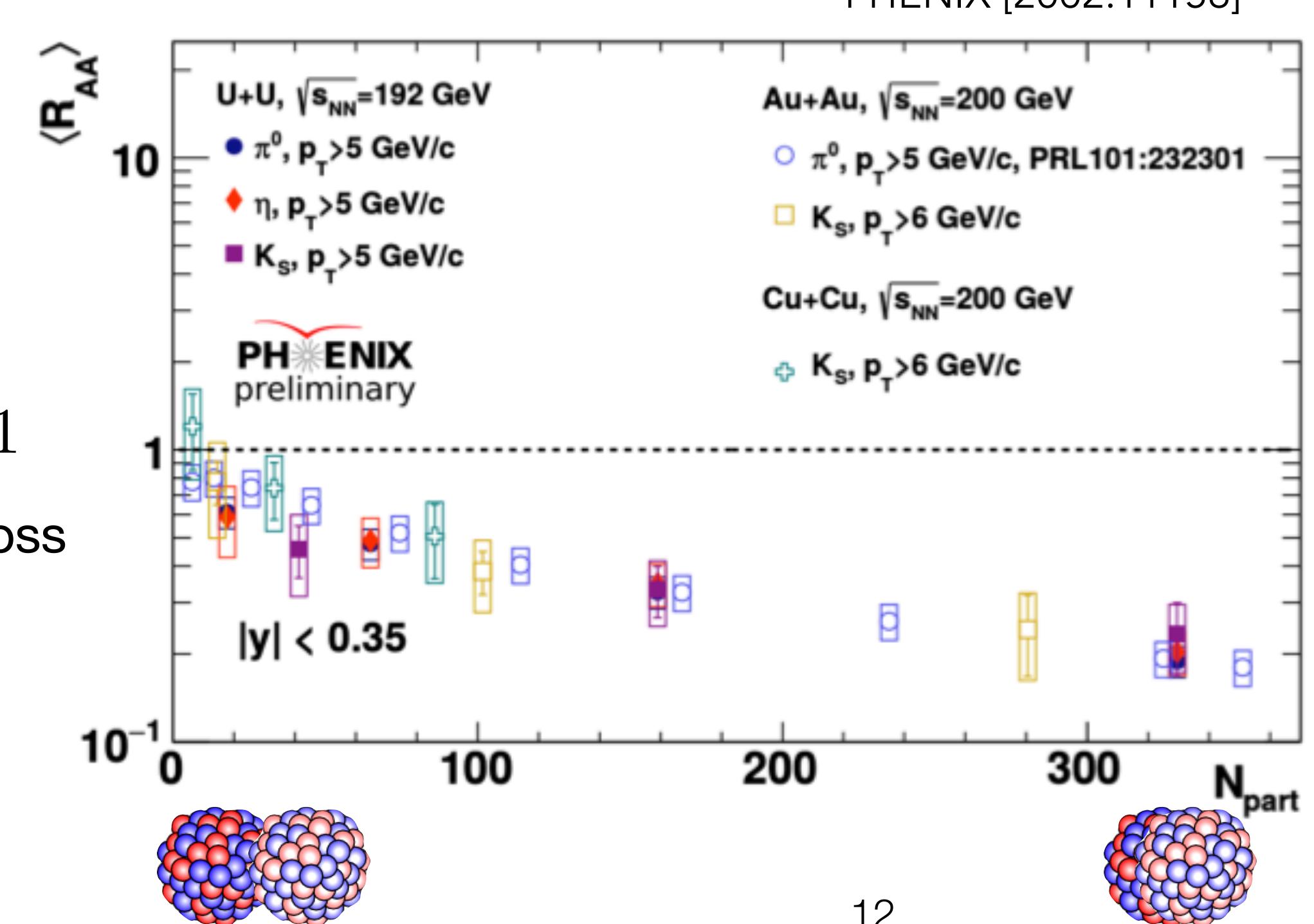
Current landscape

- Medium-induced energy loss and momentum broadening closely connected
 - From single-particle or jet suppression, recover \hat{q}

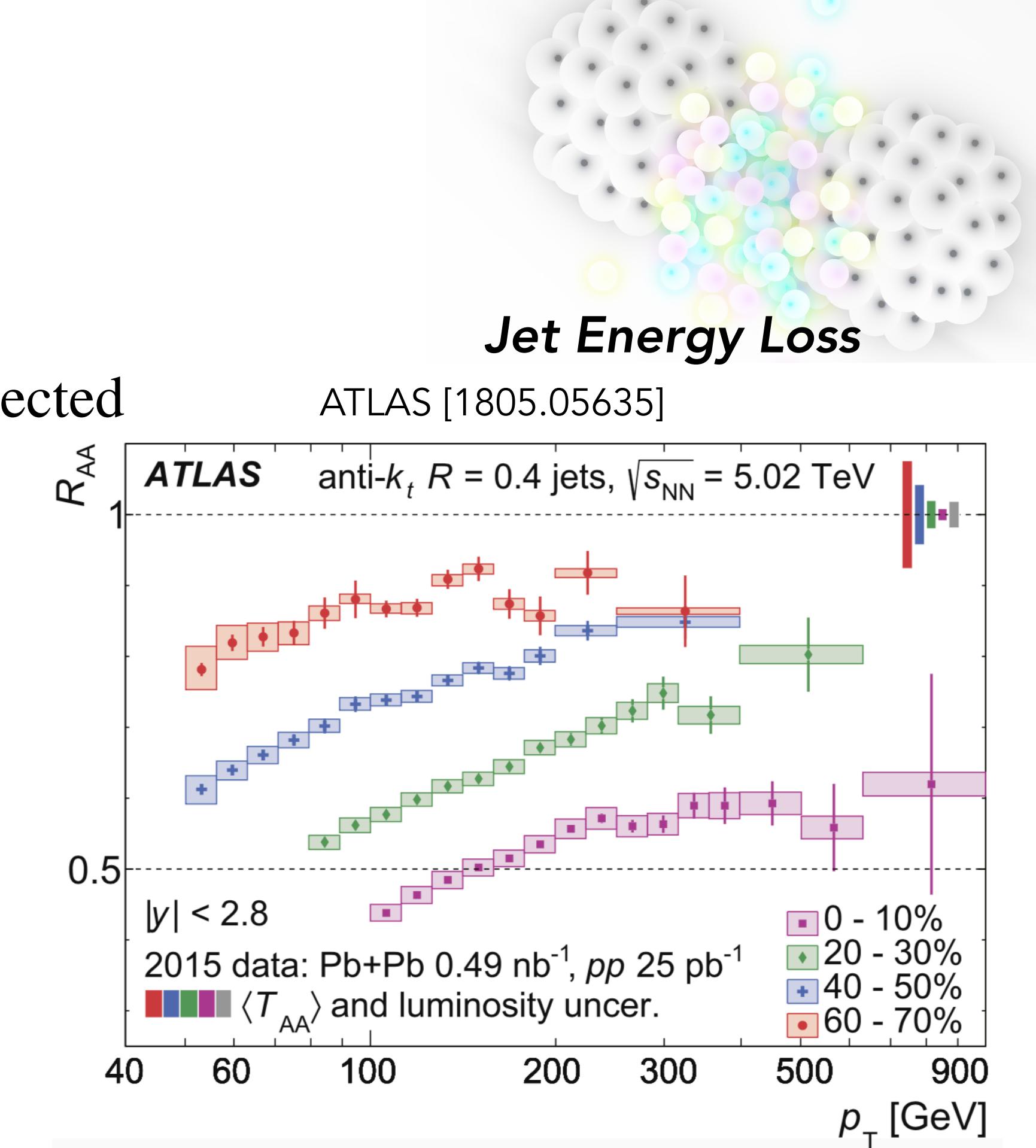
$$R_{AA} = \frac{Y_{AA}^X}{\langle T_{AA} \rangle \cdot \sigma_{pp}^X}$$

$$R_{AA} = 1$$

No energy loss



$R_{AA} < 1$
Energy loss



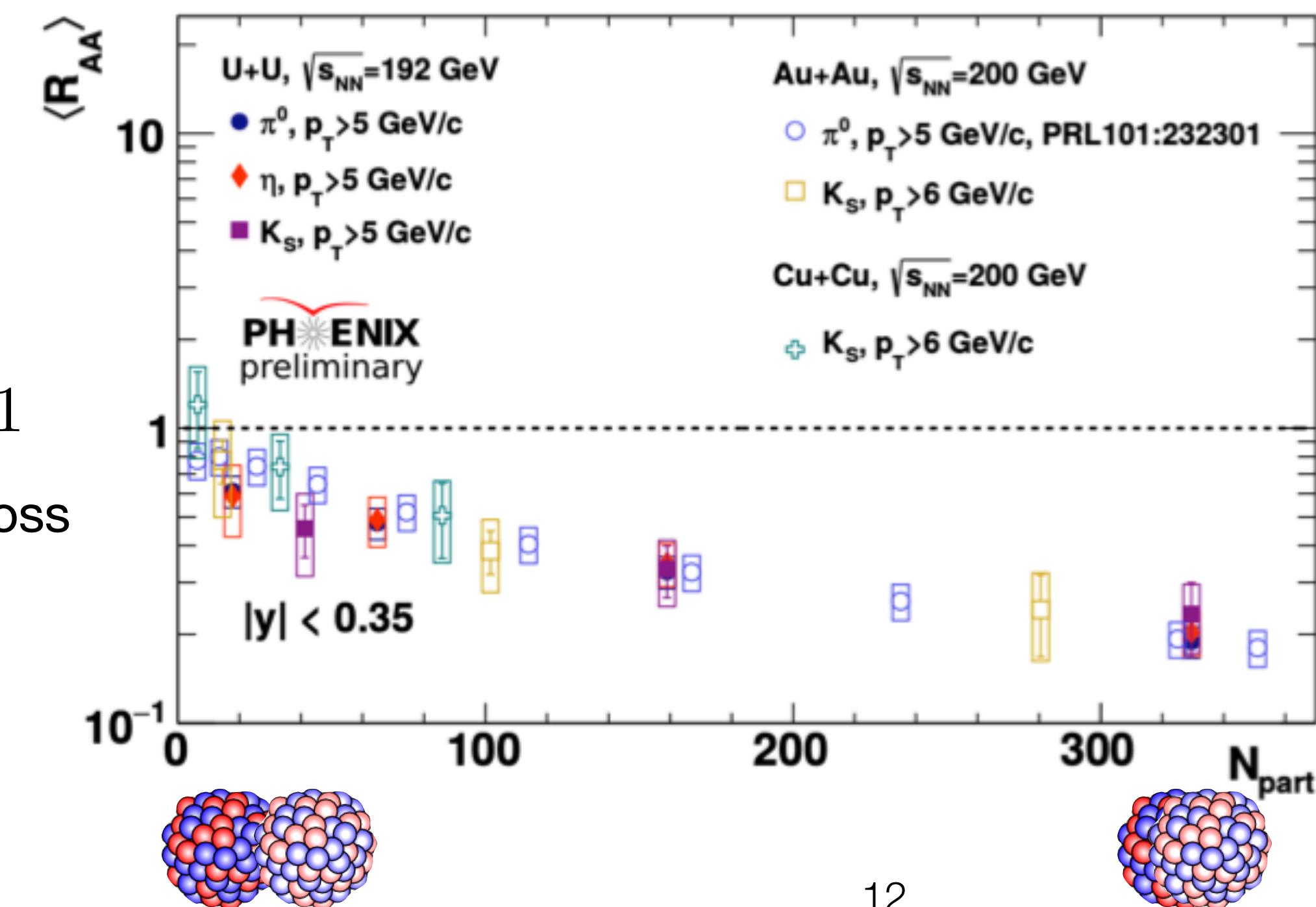
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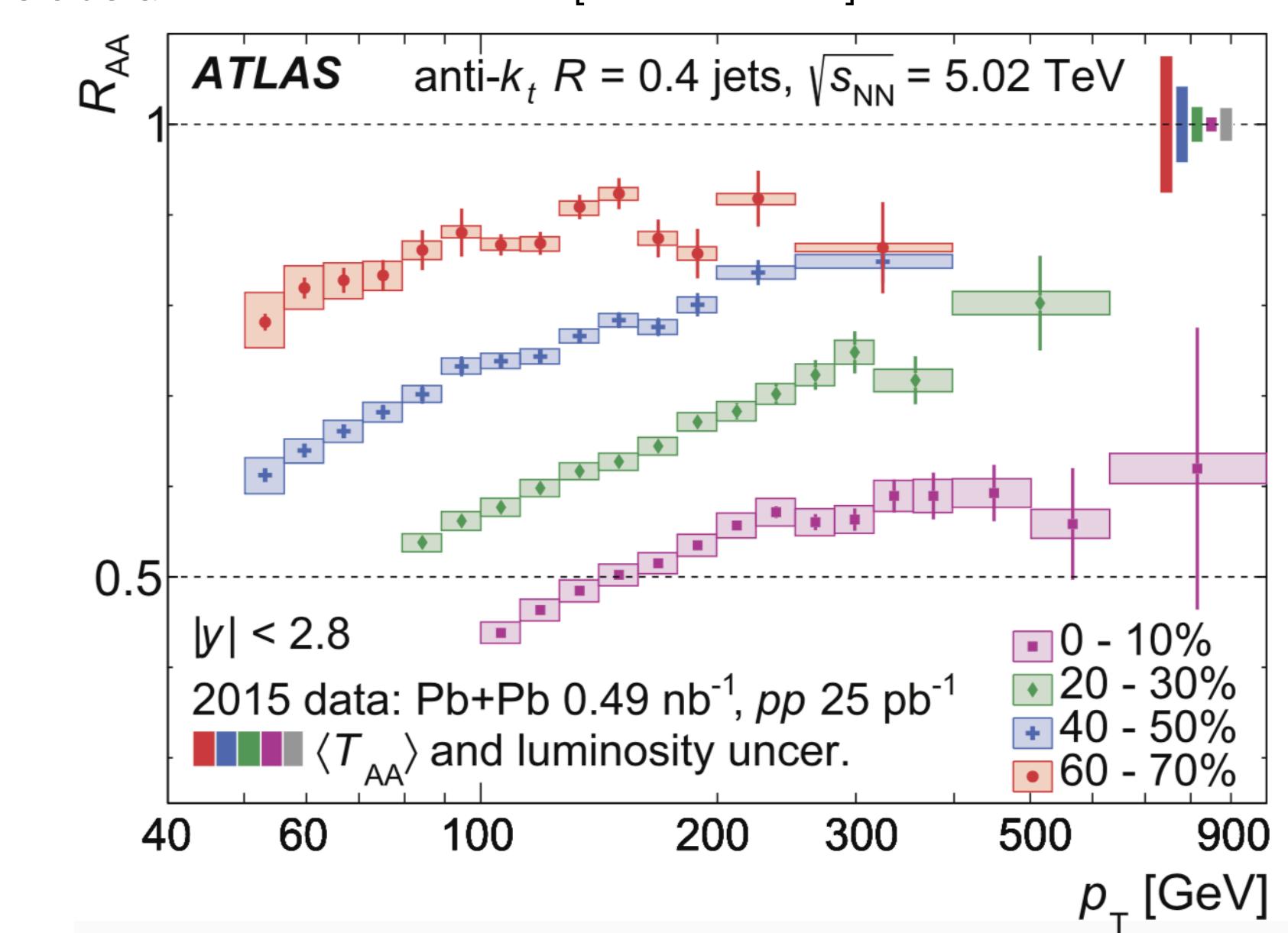
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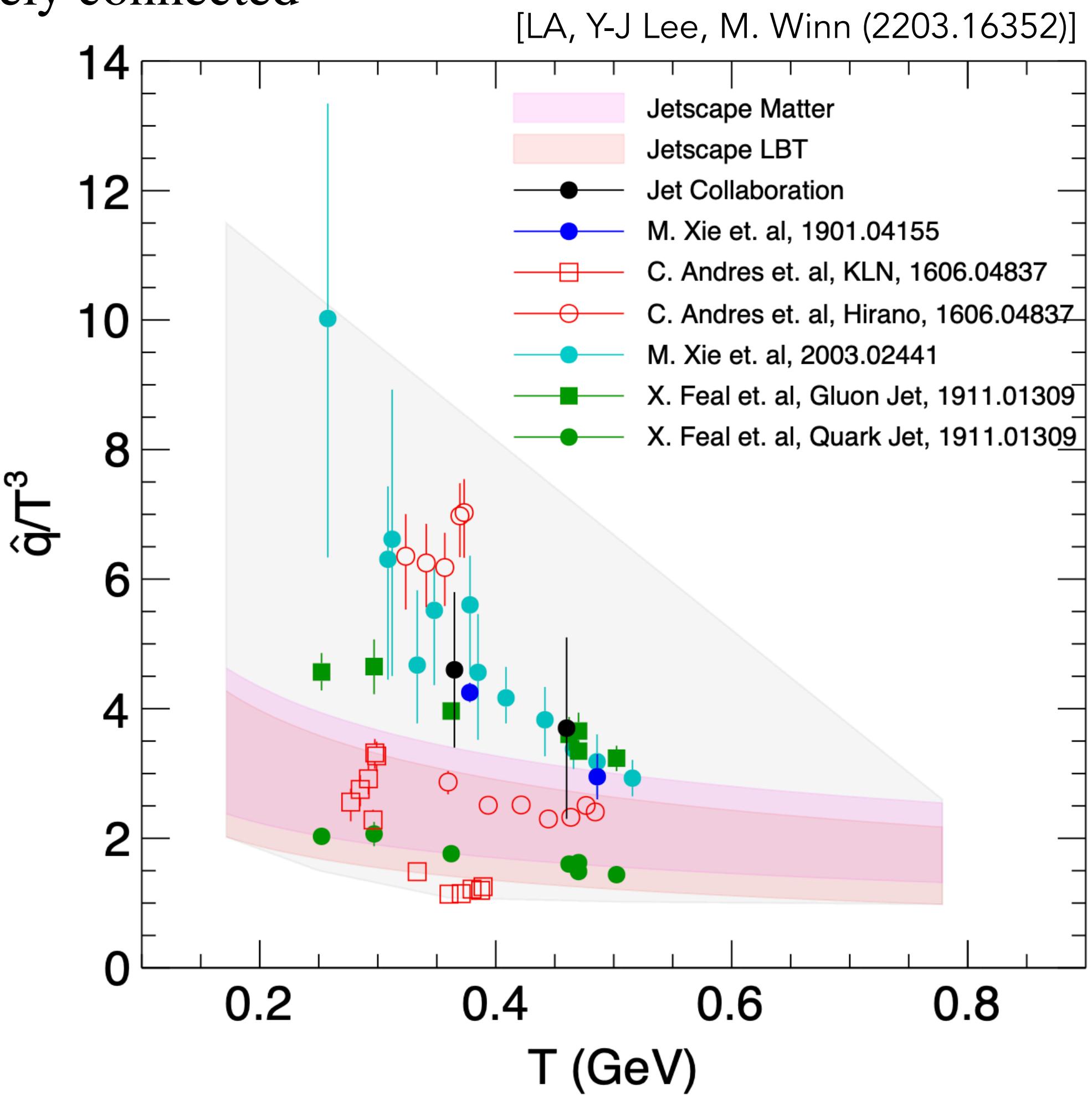
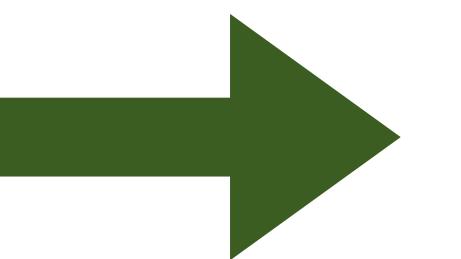
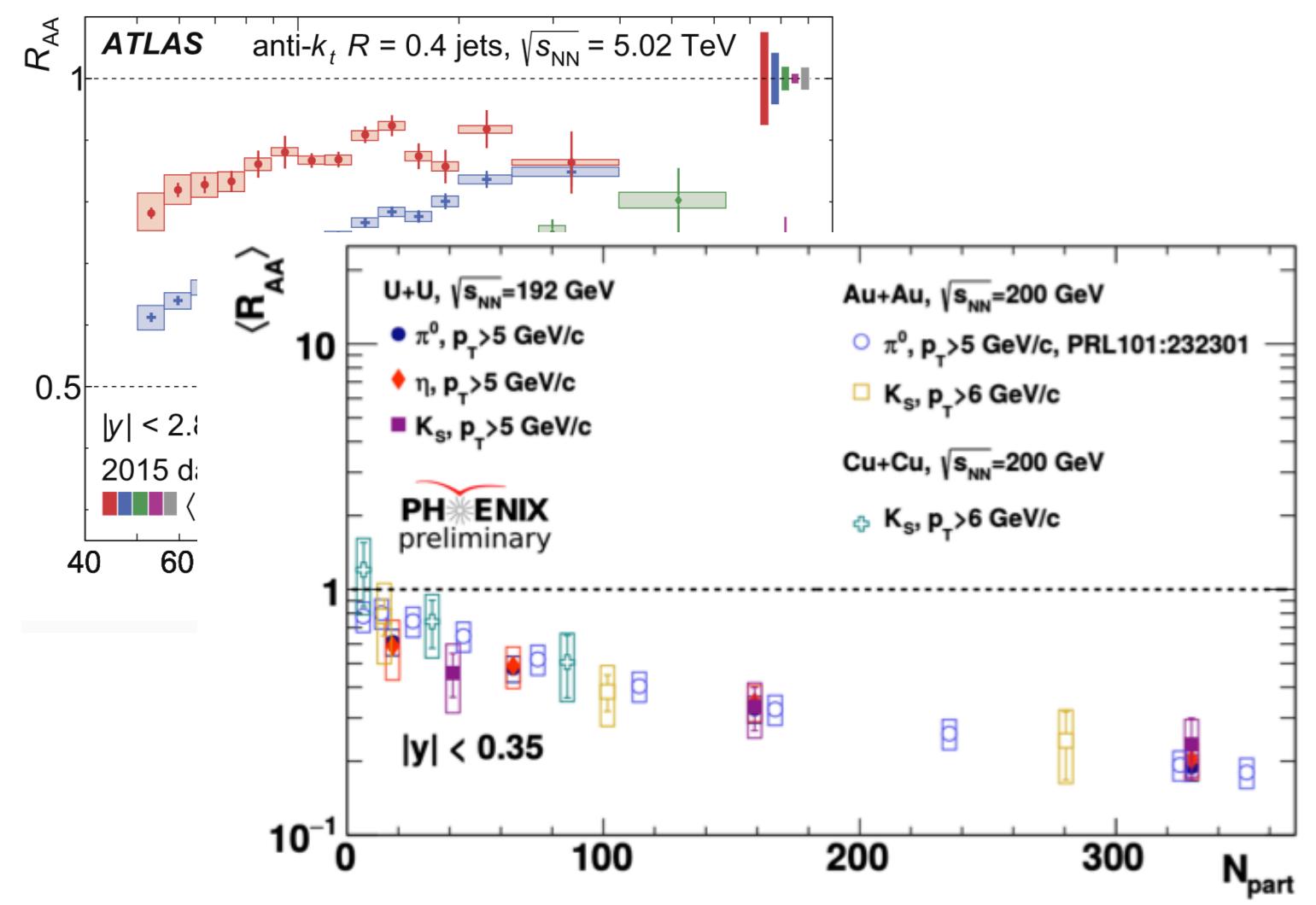


Current landscape

[HQ: Beraudo et al (1803.0382),
Cao et al (1809.07894)]

See also Escobedo (Quarkonia - Th) and Ru (CNM - Th) talks

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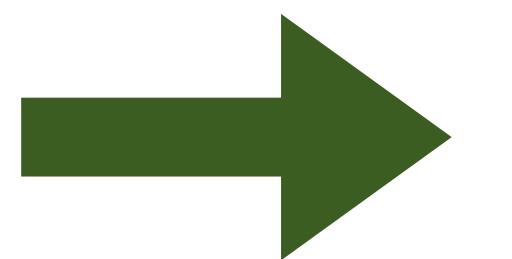
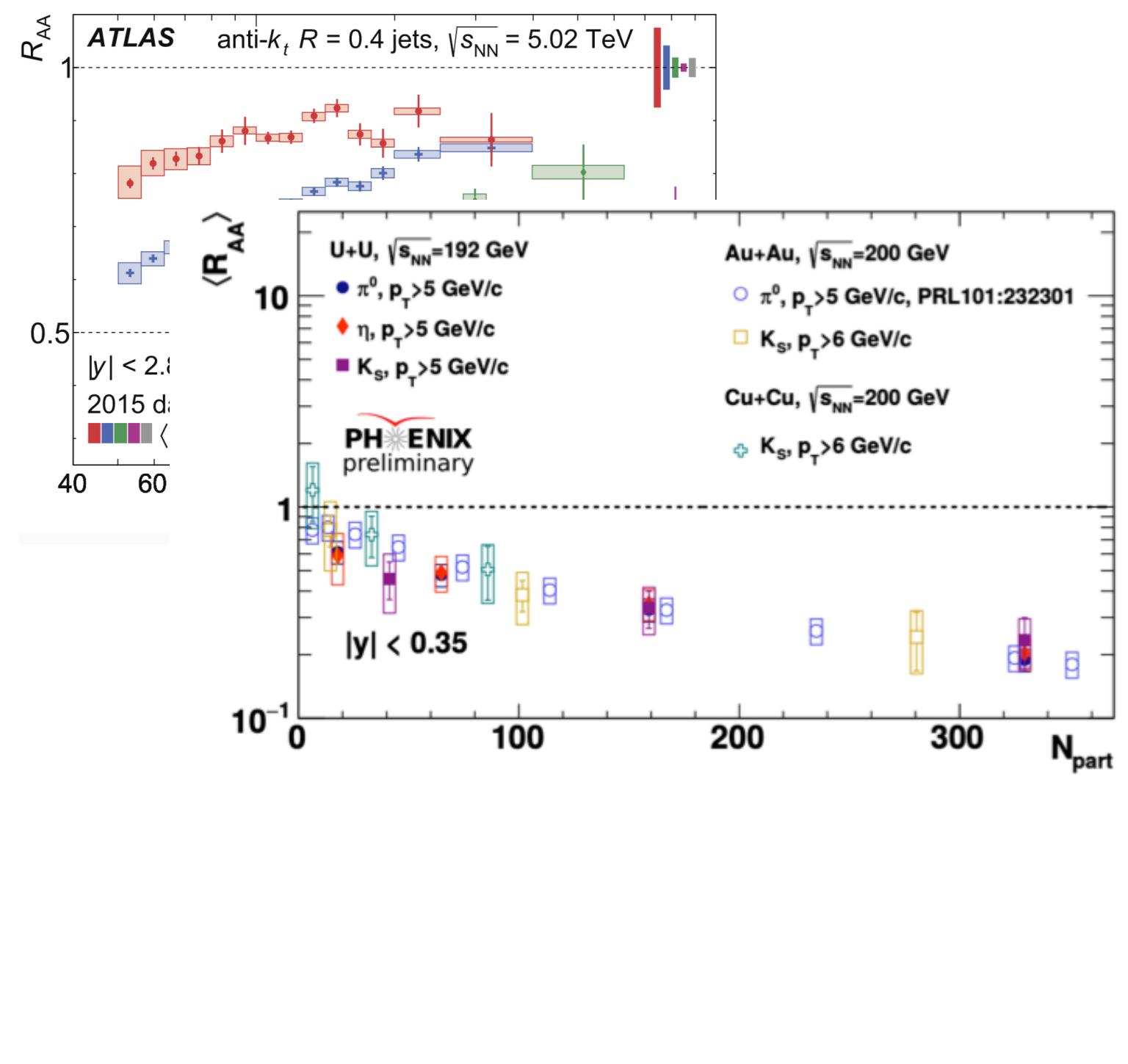
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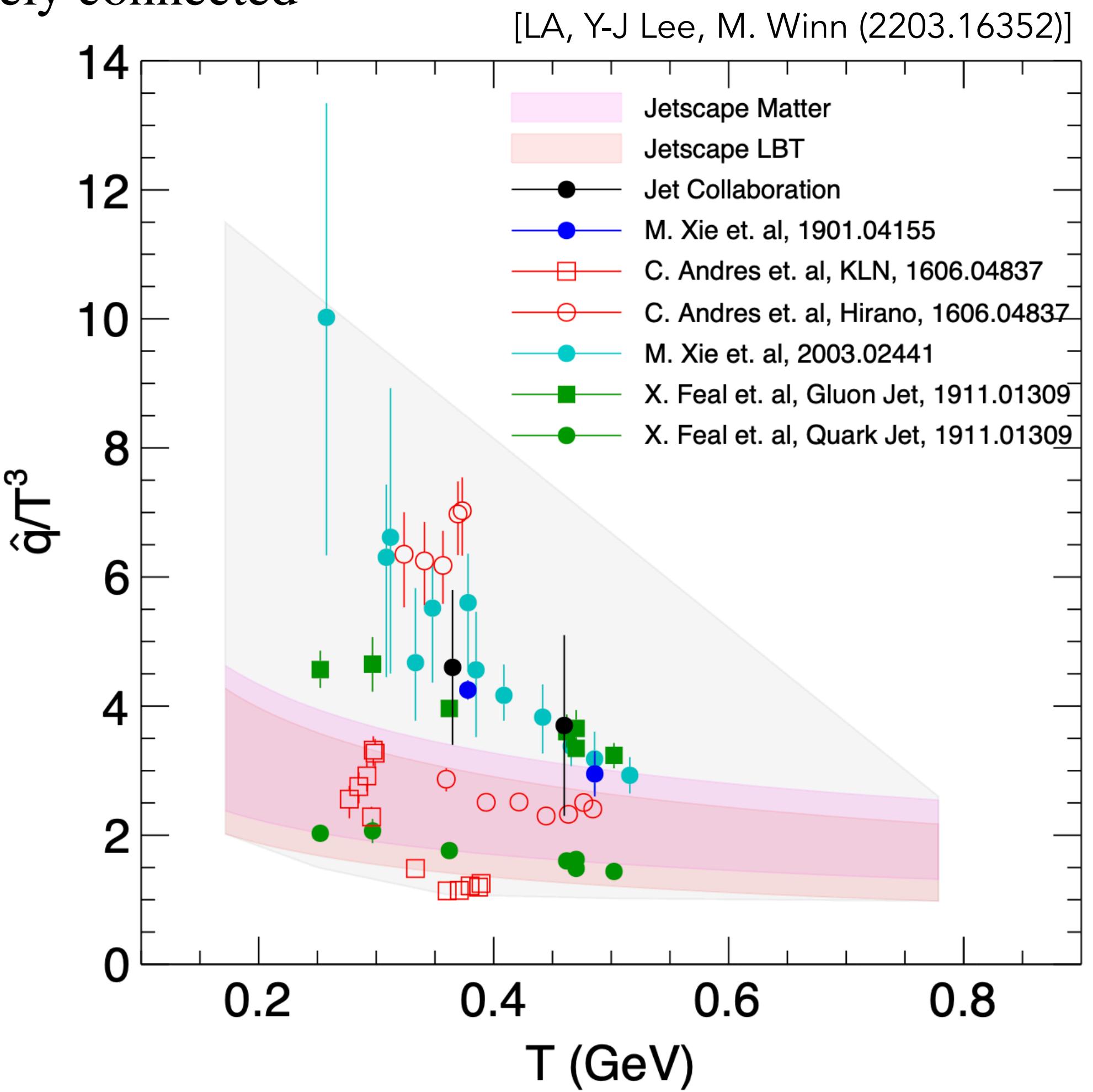
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- Several ansatz:**
- Initial state (factorisation to final-state effects)?
 - Medium temperature and energy-density time-evolution profiles?
 - QGP phase initialisation time?
 - Energy loss during partonic and hadronic phases?
 - QGP EoS and degrees of freedom?
 - ...



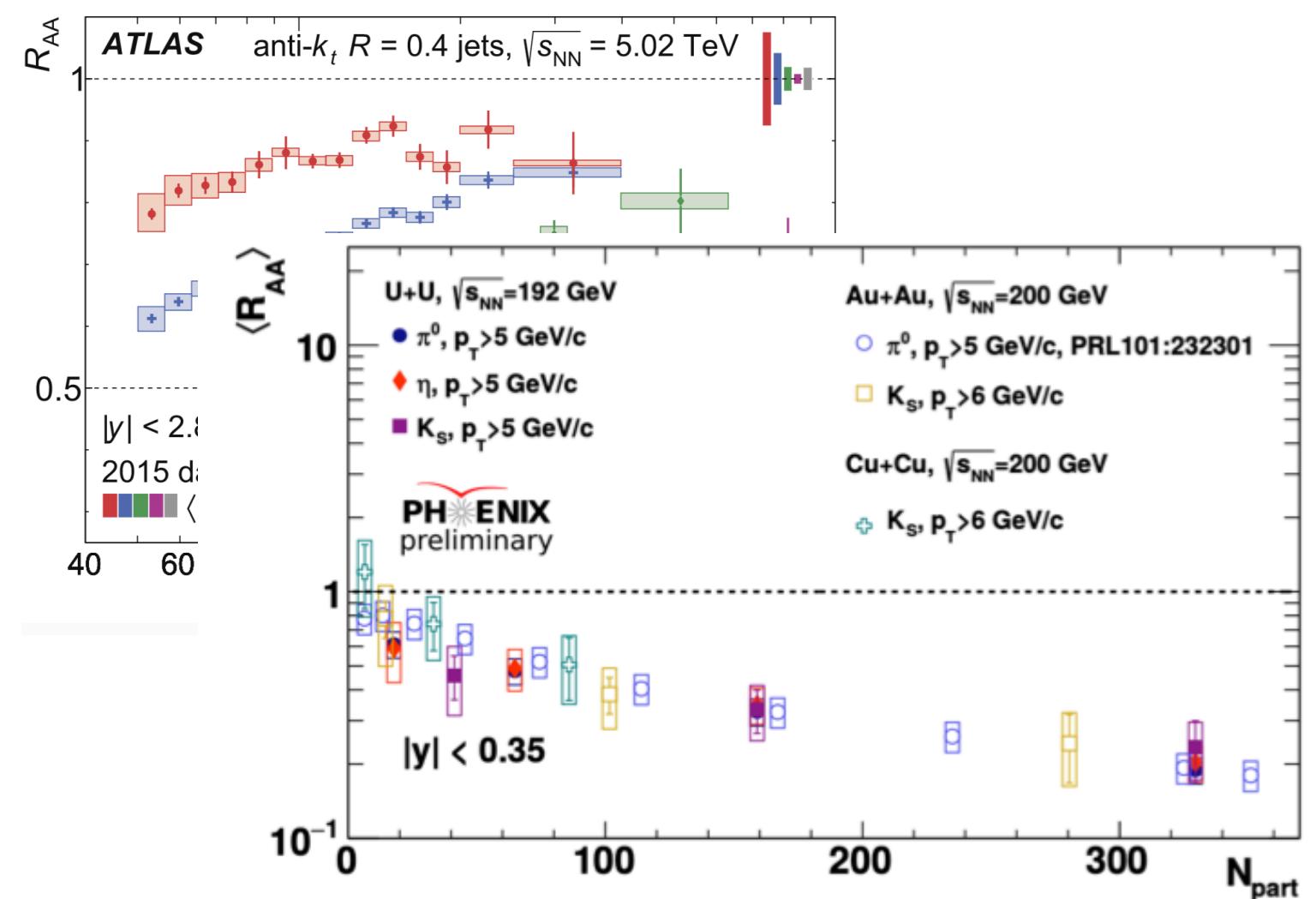
Current landscape

[HQ: Beraudo et al (1803.0382),

Cao et al (1809.07894)]

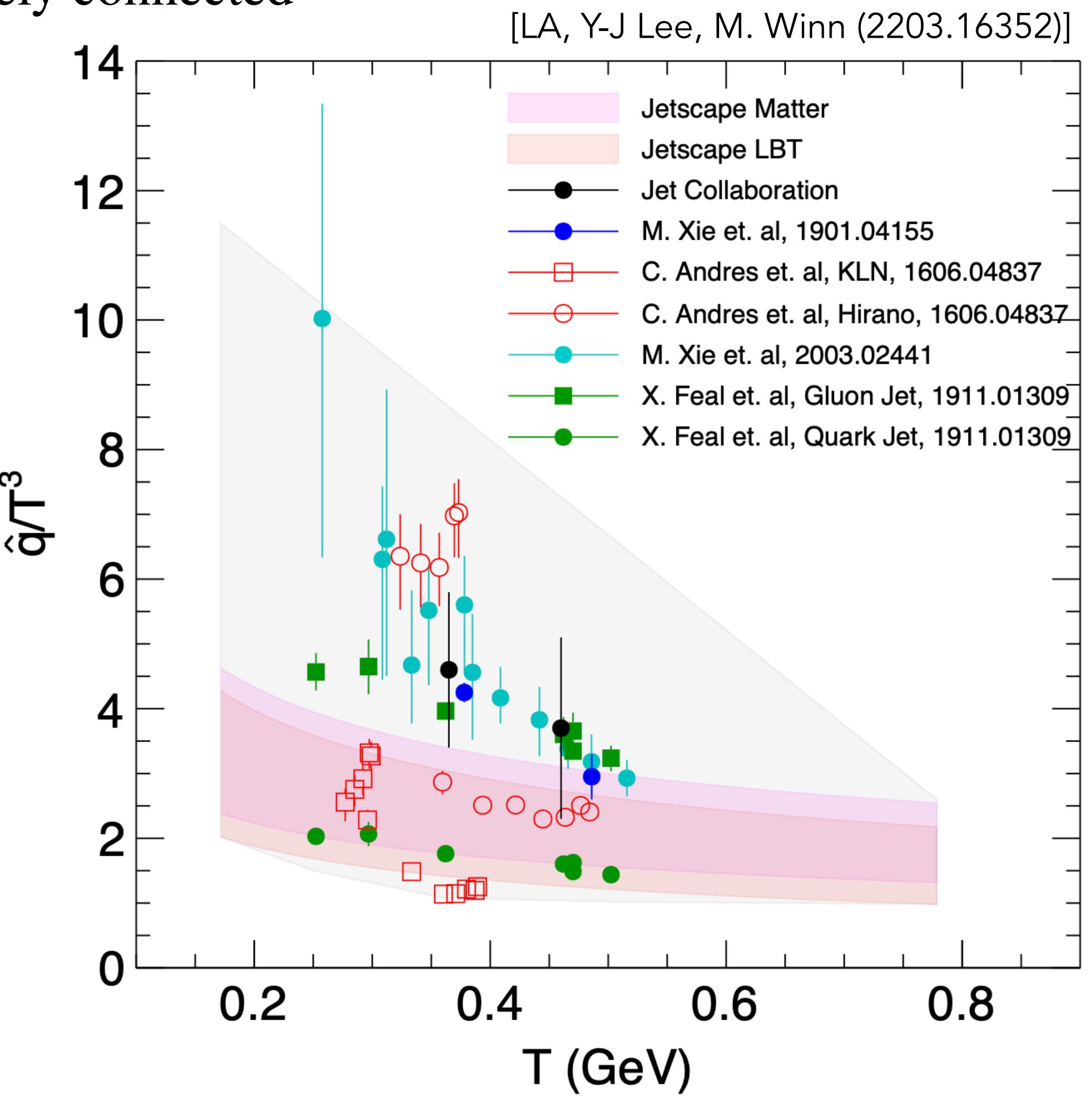
See also Escobedo (Quarkonia - Th) and Ru (CNM - Th) talks

- Medium-induced energy loss and momentum broadening closely connected
 - From single-particle or jet suppression, recover \hat{q}

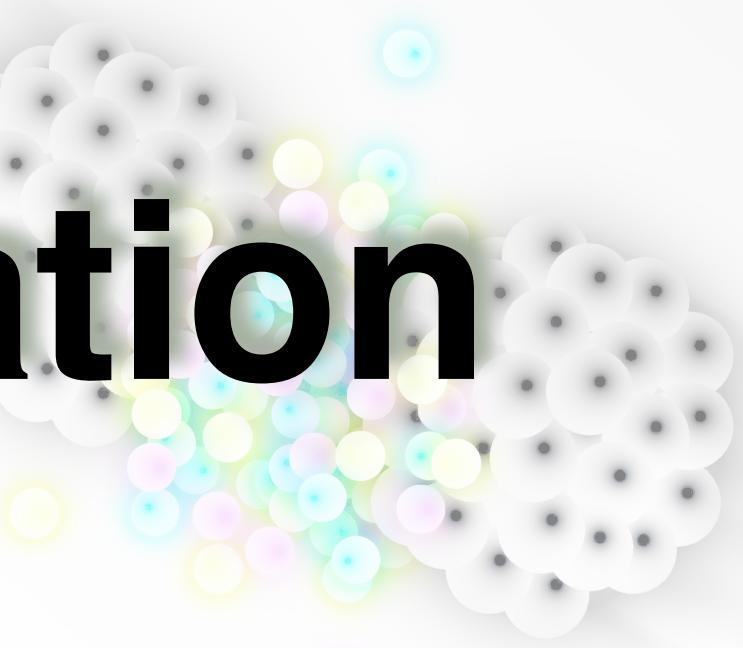


How can we improve it?

- Several ansatz:
- Initial state (factorisation to final-state effects)?
 - Medium temperature and energy-density time-evolution profiles?
 - QGP phase initialisation time?
 - Energy loss during partonic and hadronic phases?
 - QGP EoS and degrees of freedom?
 - ...



Improving medium-induced radiation



- Accuracy of radiation spectrum:



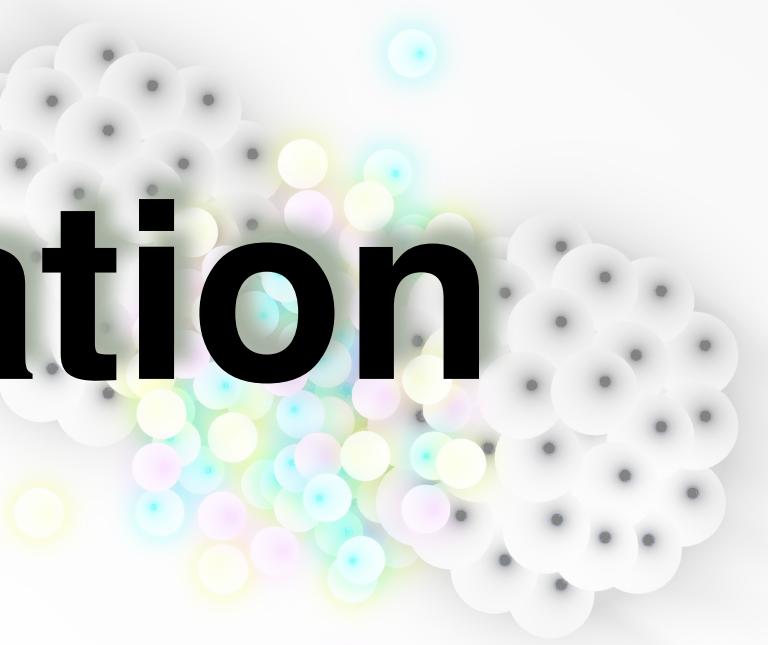
or

- Improved analytic opacity expansion (expand multiple soft interaction)

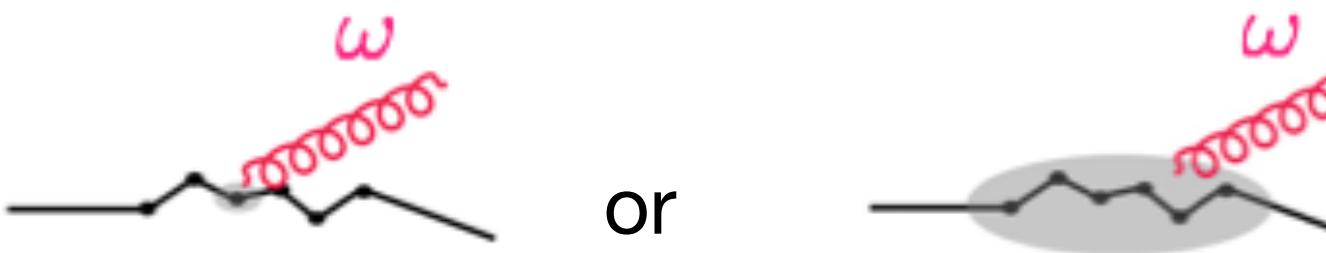
[Barata, Mehtar-Tani, Soto-Ontoso, Tywoniuk
(1910.02032, 2106.07402)]

$$n(s)\sigma(\mathbf{r}) \simeq \frac{1}{2}\hat{q}\mathbf{r}^2 + \mathcal{O}(r^2 \ln r^2) \Rightarrow v(r, s)_{HO} + \delta v(r, s)$$

Improving medium-induced radiation



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- Full numerical solution:

[Andrés, LA, Dominguez, Gonzales
(2002.01517, 2011.06522)]

- Solve the spectrum by using Schwinger-Dyson type equations (in momentum space):

$$\partial_\tau \mathcal{P}(\tau, \mathbf{k}; s, \mathbf{l}) = -\frac{1}{2}n(\tau) \int_{\mathbf{k}'} \sigma(\mathbf{k} - \mathbf{k}') \mathcal{P}(\tau, \mathbf{k}'; s, \mathbf{l})$$

$$\partial_t \tilde{\mathcal{K}}(s, \mathbf{q}; t, \mathbf{p}) = \frac{i\mathbf{p}^2}{2\omega} \tilde{\mathcal{K}}(s, \mathbf{q}; t, \mathbf{p}) + \frac{1}{2}n(t) \int_{\mathbf{k}'} \sigma(\mathbf{k}' - \mathbf{p}) \tilde{\mathcal{K}}(s, \mathbf{q}; t, \mathbf{k}')$$

Set of integro-partial differential equations that can be numerically solved to any (realistic) potential

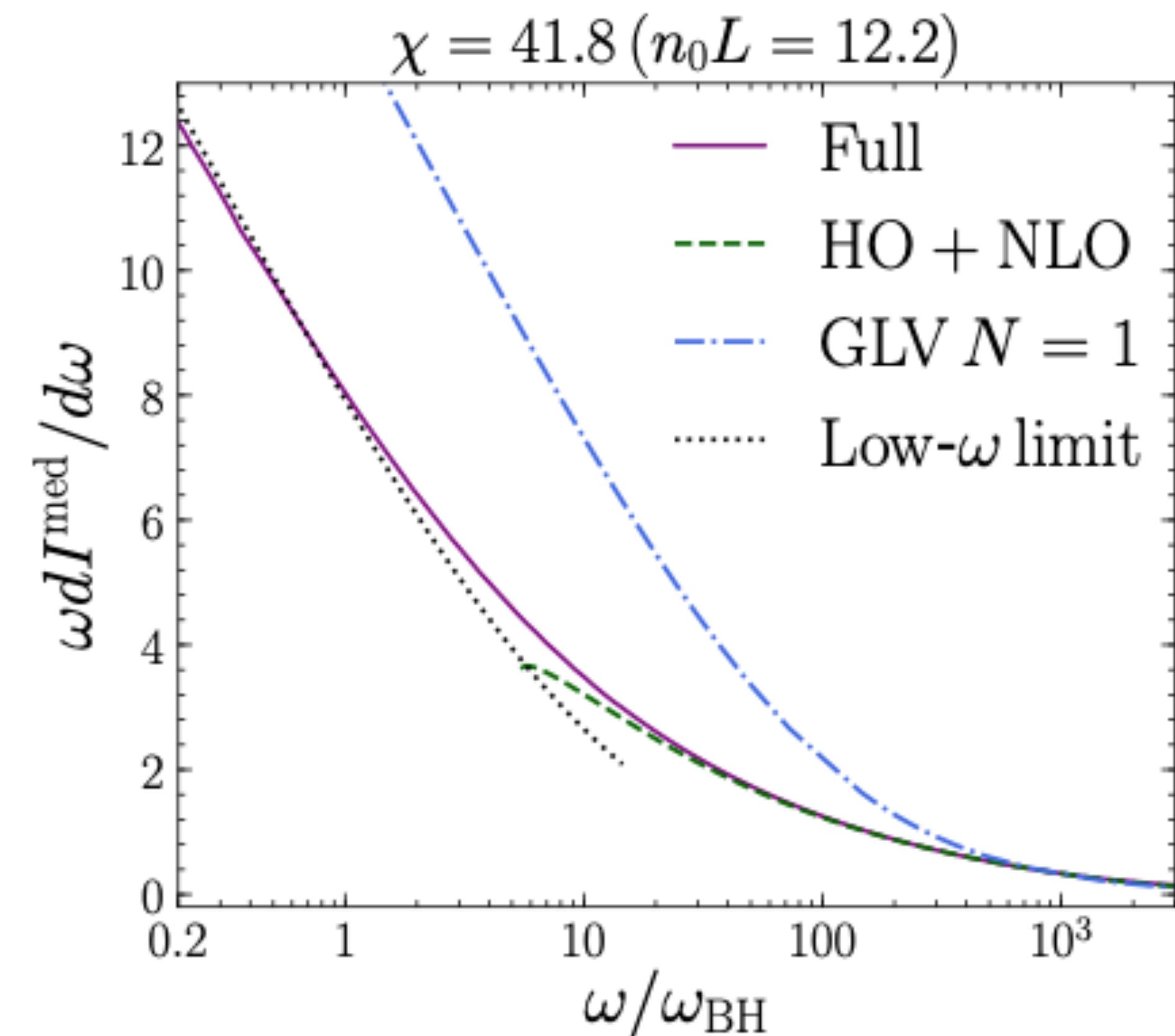
Also: [Feal, Salgado, Vasquez (1911.01309)]

Improving medium-induced radiation



- **Accuracy** of radiation spectrum:
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 - Full numerical solution:
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Improving medium-induced radiation

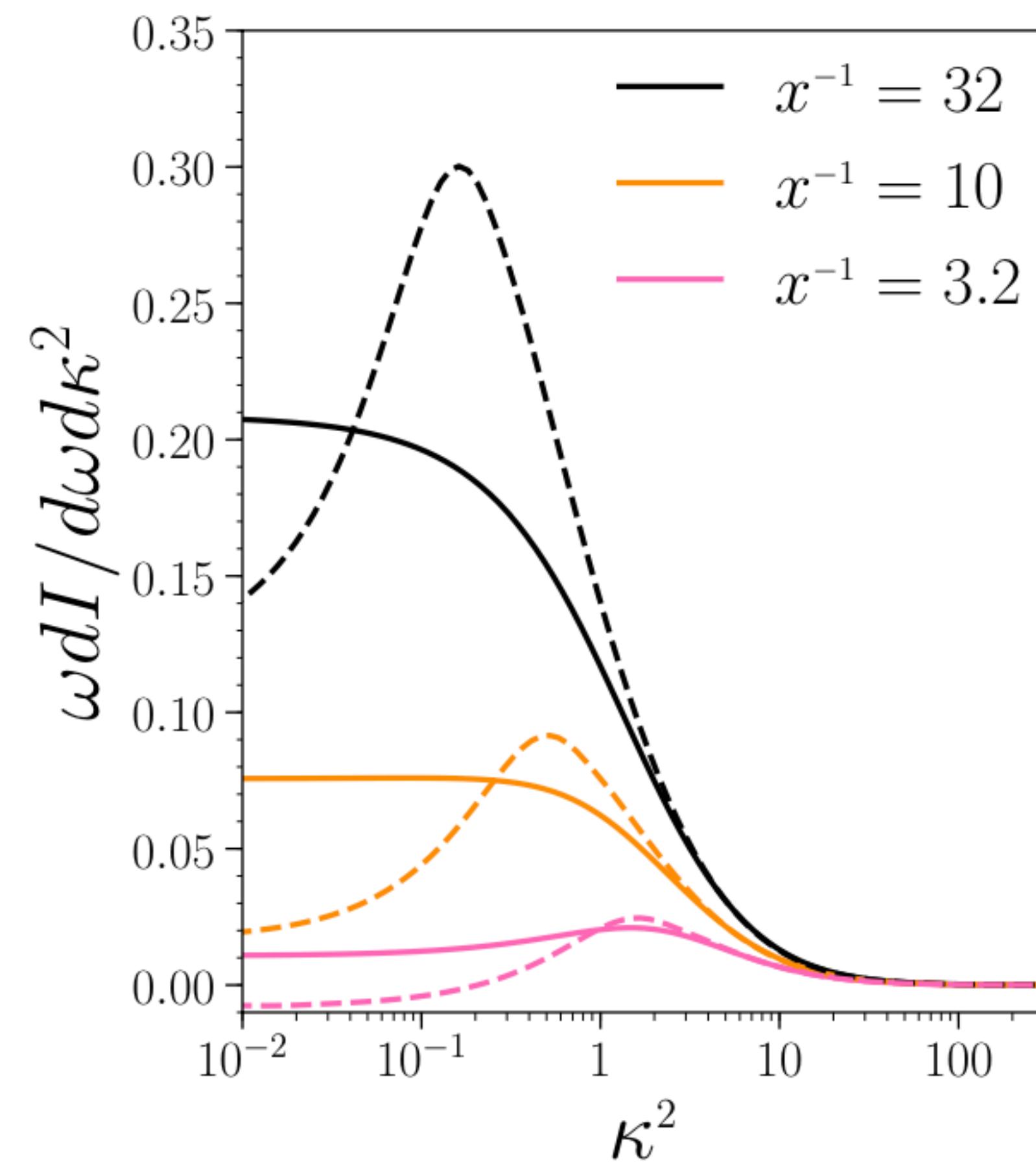
[Andrés, LA, Dominguez, (2002.01517)]

- Accuracy of radiation spectrum:
 - Improved analytic opacity expansion
 - Full numerical solution:
 - Solve the spectrum by using Schwinger-Dyson type equations (in momentum space):

Yukawa potential: $V(\mathbf{q}) = \frac{8\pi\mu^2}{(\mathbf{q}^2 + \mu^2)^2}$

HTL potential: $\frac{1}{2}n V(\mathbf{q}) = \frac{g_s^2 N_c m_D^2 T}{\mathbf{q}^2 (\mathbf{q}^2 + m_D^2)}$

- Full HTL $TL = 0.4$
- Full Yukawa $n_0 L = 1$

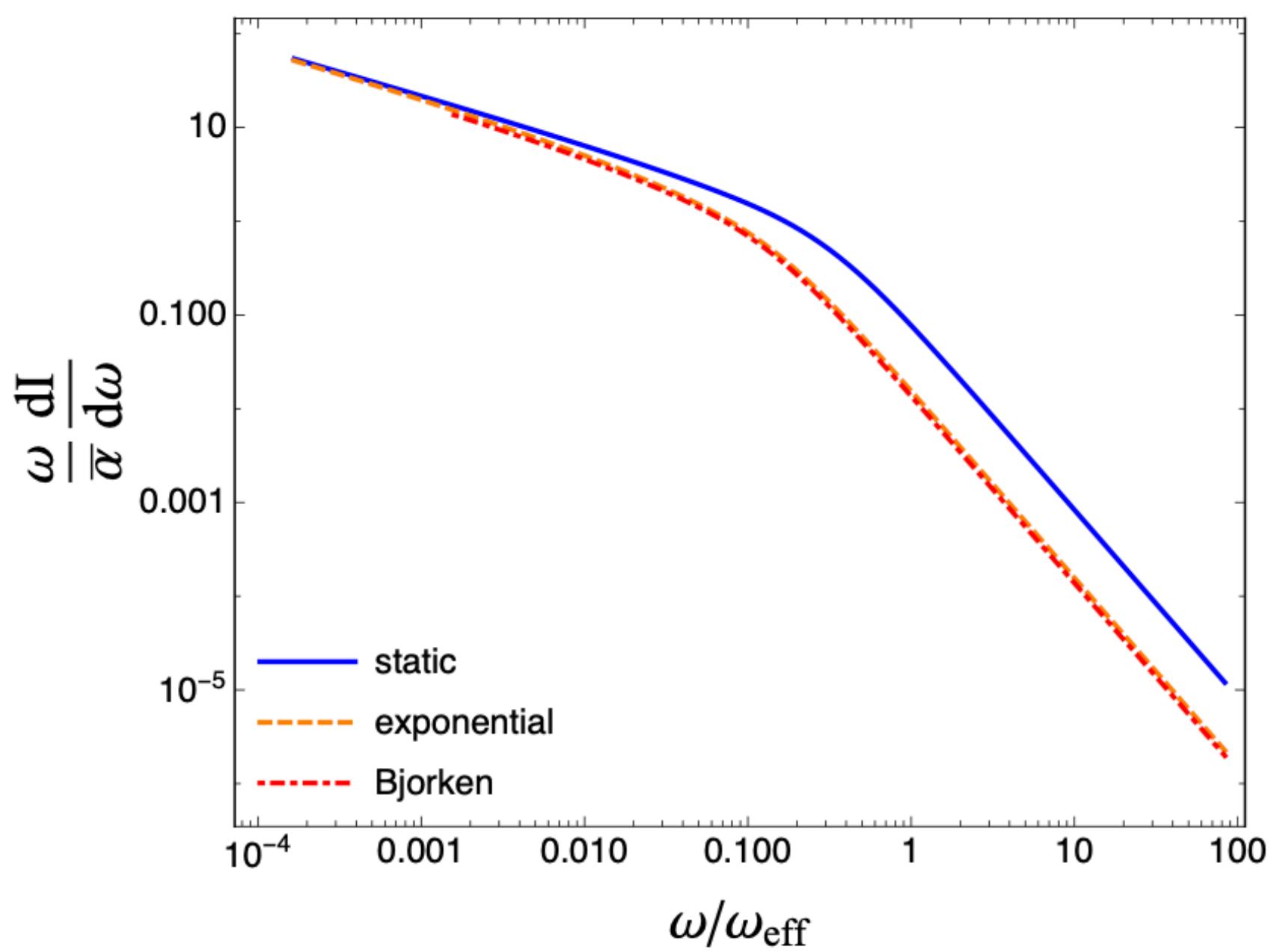
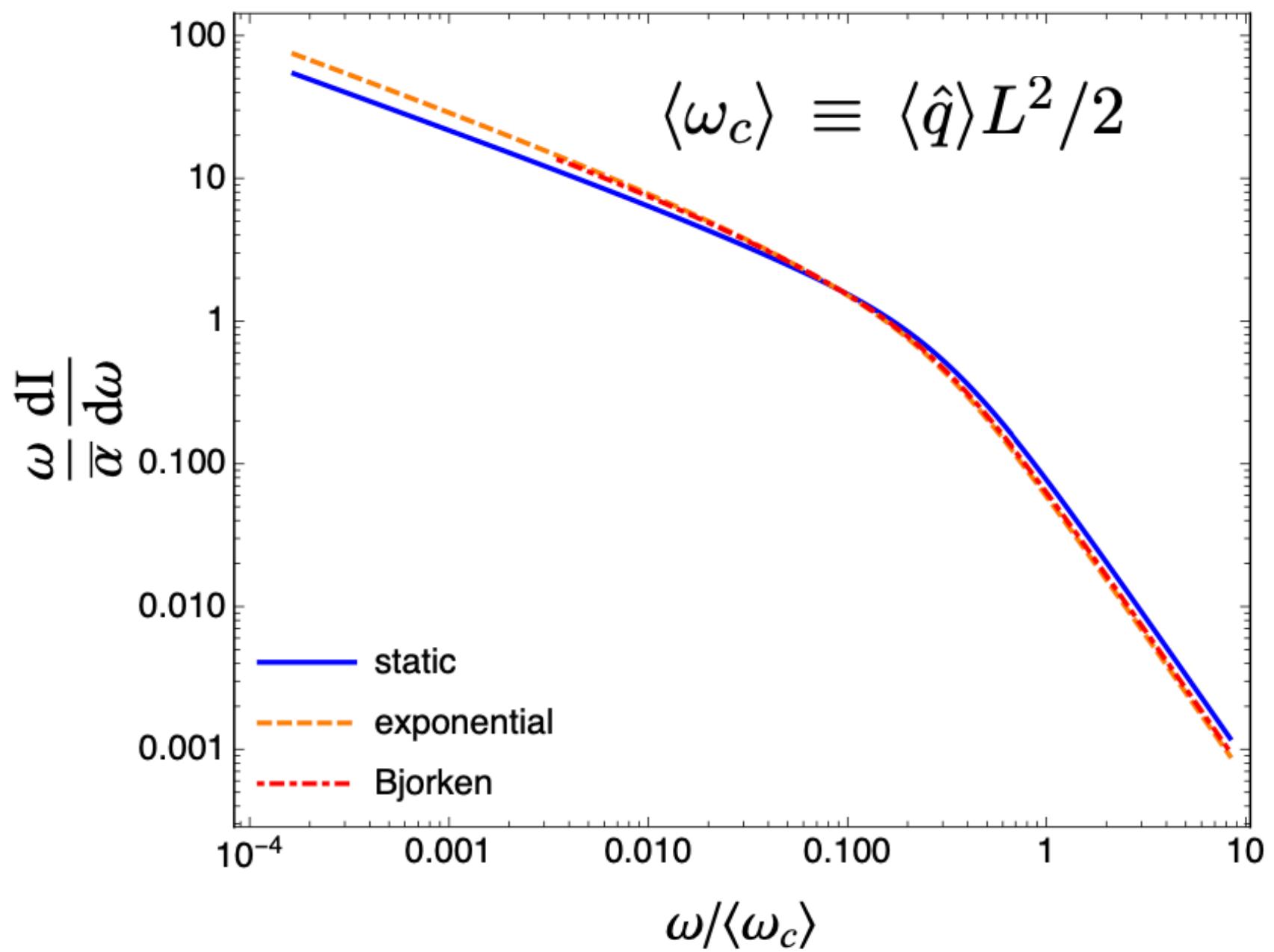


Improving medium-induced radiation

- Effects of **medium expansion** on energy loss (HO): $\hat{q} = \hat{q}(t)$
 - Static equivalent of an expanding medium obtained by scaling laws:

$$\langle \hat{q} \rangle = \frac{2}{L^2} \int_{t_0}^{L+t_0} dt (t - t_0) \hat{q}(t)$$

[Adhya, Salgado, Spousta, Tywoniuk, (1911.12193)]



$$\omega_{\text{eff}} = \begin{cases} \frac{1}{2}\hat{q}_0 L^2 & \text{static medium} \\ 2\hat{q}_0 L^2 & \text{exponentially expansion} \\ 2\hat{q}_0 t_0 L & \text{Bjorken expansion} \end{cases}$$

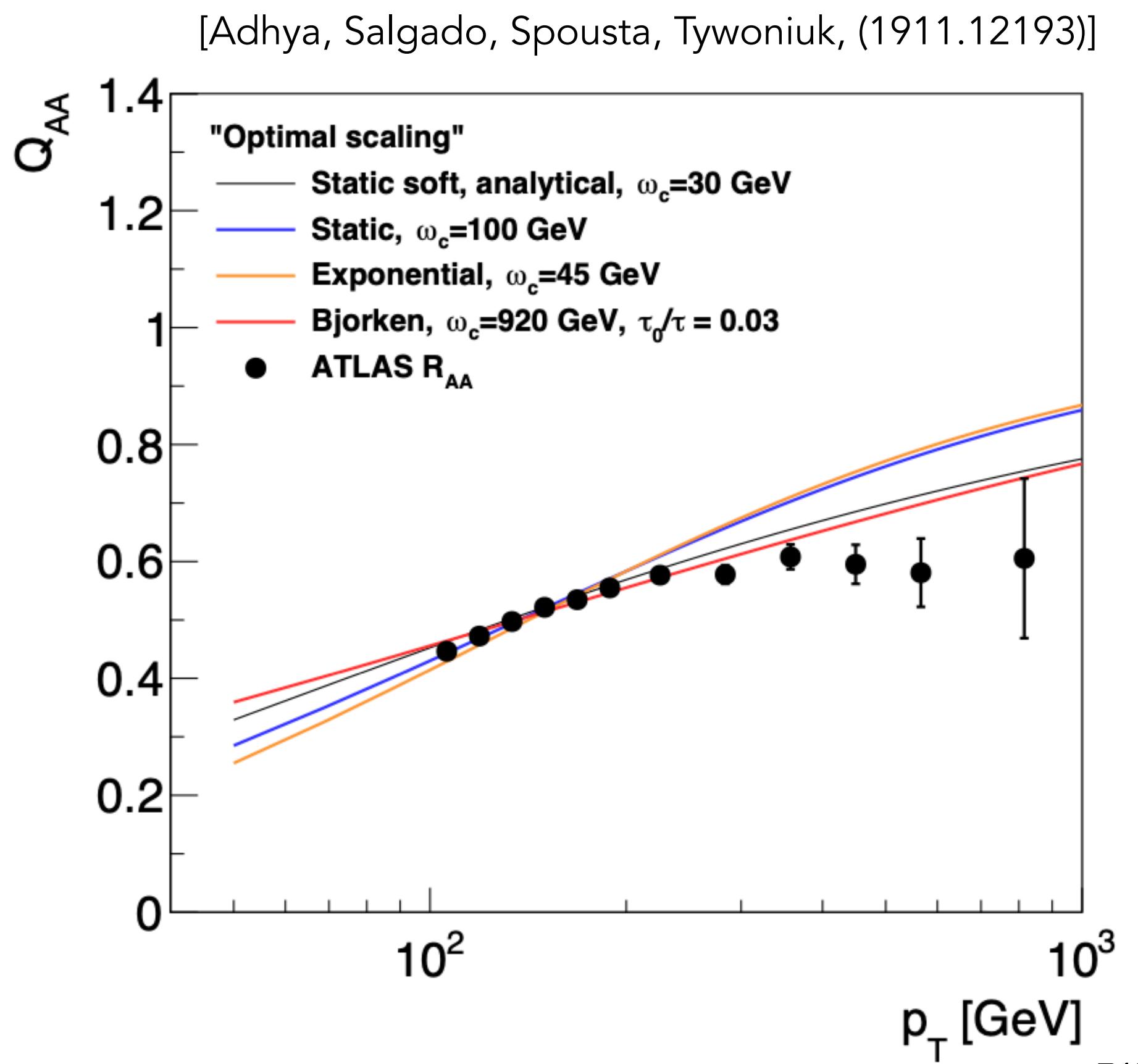
Improving medium-induced radiation

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\hat{q}_0 [GeV ³]	static	exponential	Bjorken
no scaling	0.2	0.2	0.2
soft scaling	0.2	0.05	1.66
optimal scaling	0.2	0.09	1.84
scaling by $\langle \omega_c \rangle$	0.2	0.1	3.33



Improving medium-induced radiation

[Andrés, LA, Dominguez, Gonzalez, Salgado (on-going)]

- Effects of **medium expansion** on energy loss (full solution):

- Static equivalent of an expanding medium obtained by scaling laws:

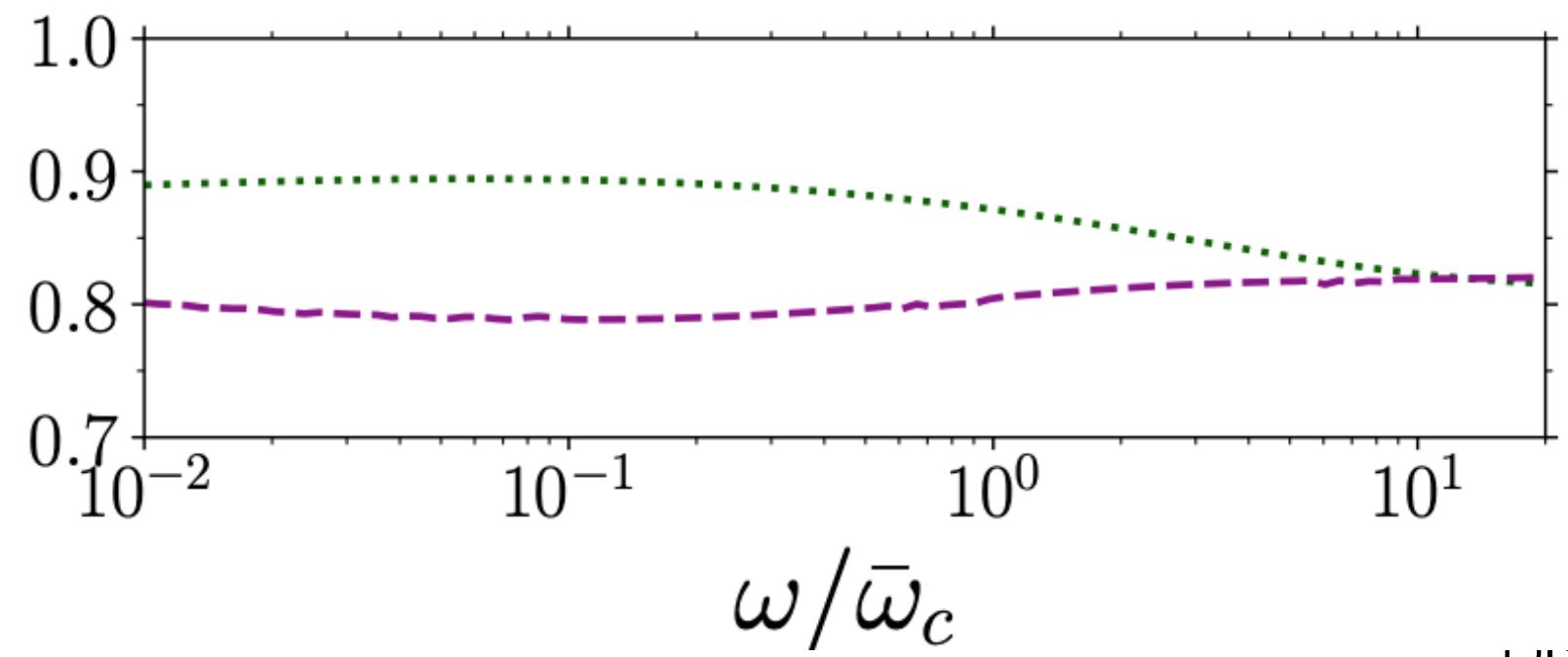
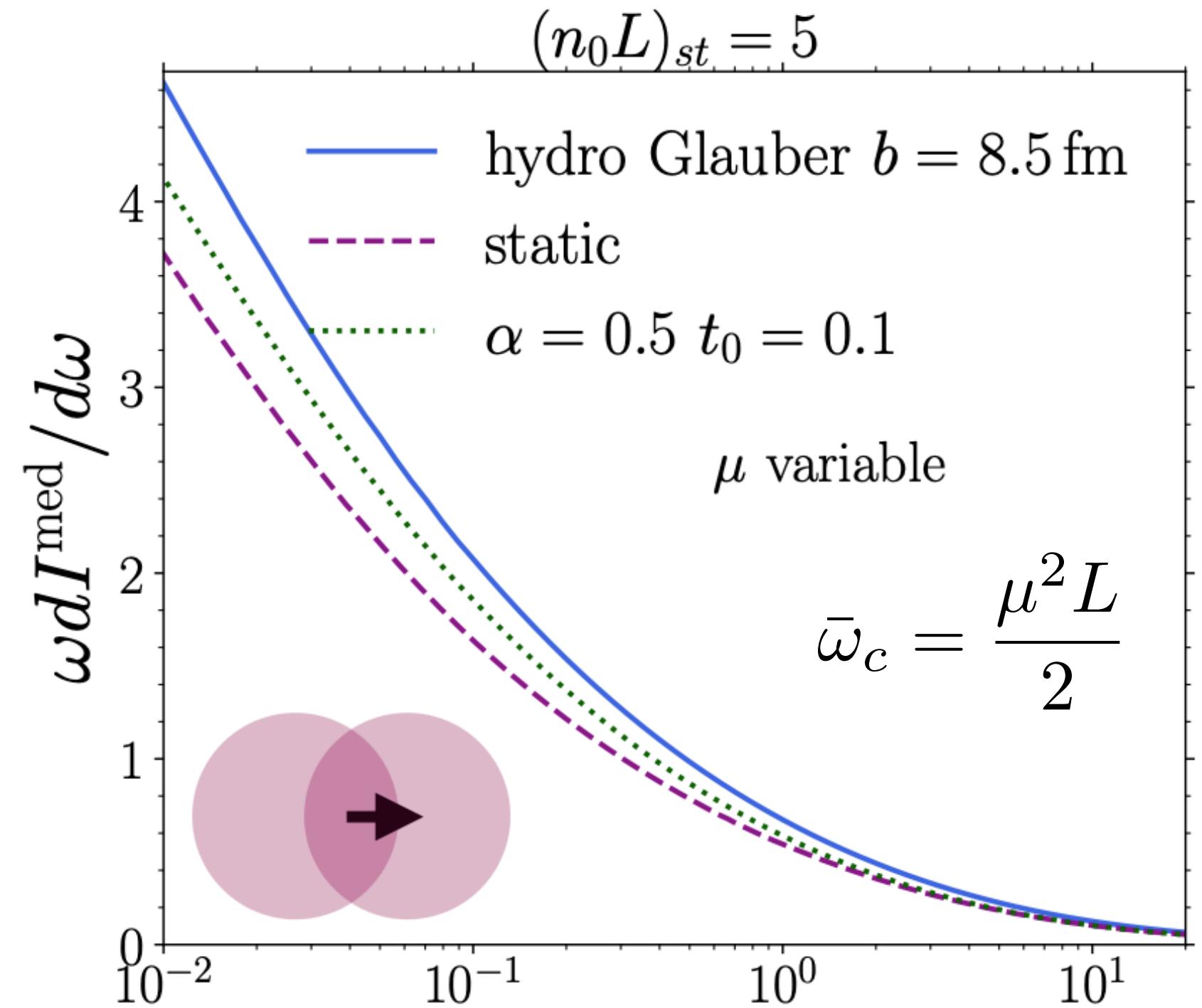
$$n_0 L = \int_0^{L'} dt n(t) \quad \frac{n_0 \mu^2 L^2}{2} = \int_0^{L'} dt t n(t) \mu^2(t)$$

- For a hydrodynamic medium, use instead a power-law equivalent to improve accuracy

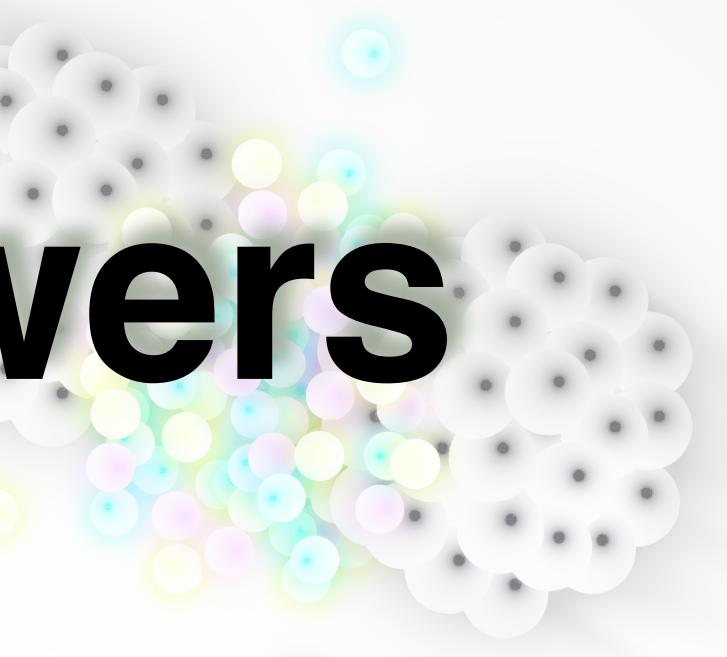
$$n_{hydro}(t) = k_1 T(t)$$

$$\mu_{hydro}^2(t) = k_2 T^2(t)$$

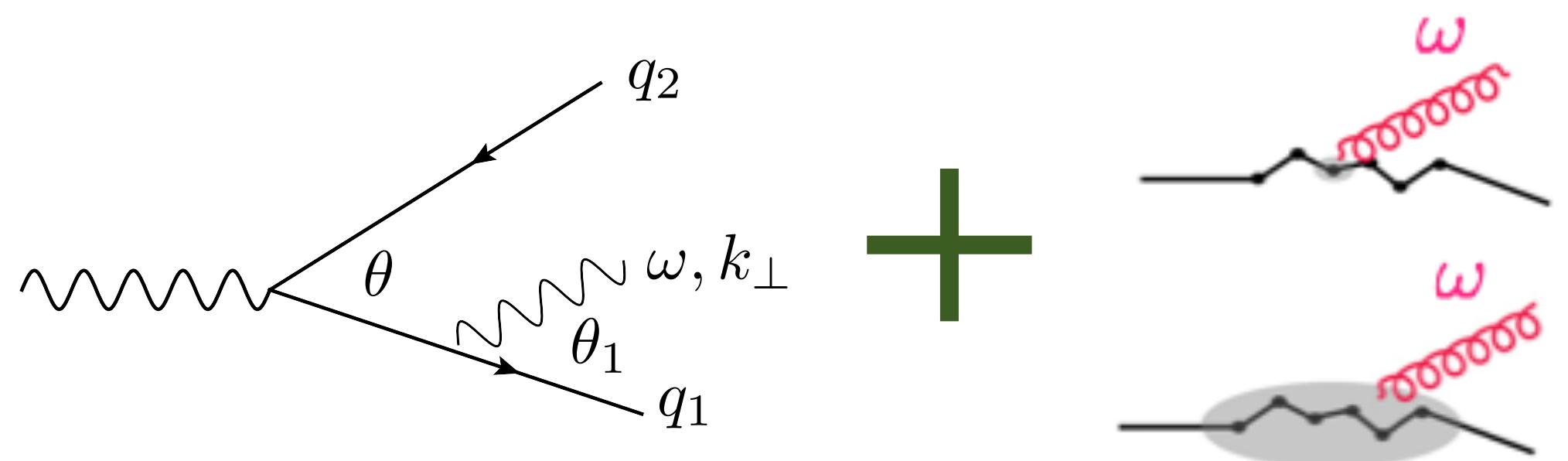
$$n(t) = \frac{n'_0}{(t + t_0)^\alpha} \quad \mu^2(t) = \frac{\mu'^2}{(t + t_0)^{2\alpha}}$$



Improving "medium" parton showers



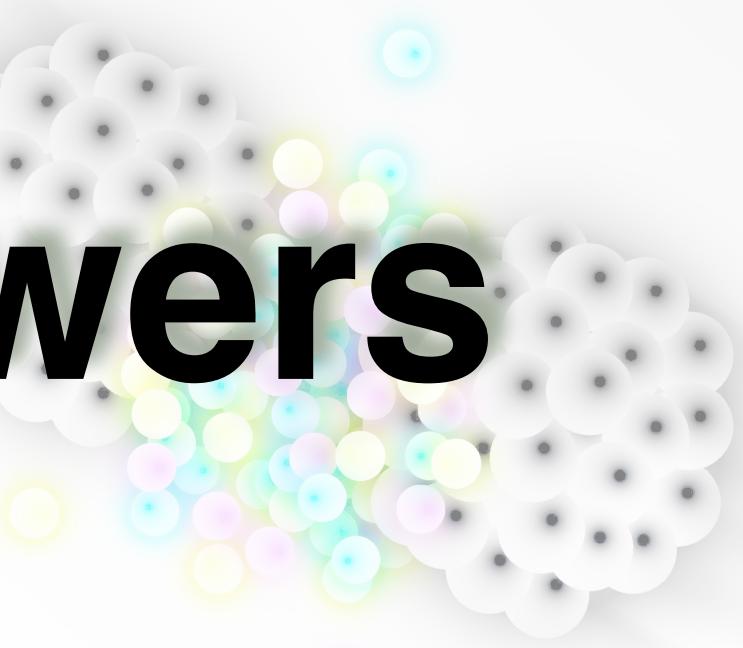
- Multiple emitters:
 - Interference effects suppressed (+ anti-angular ordering)



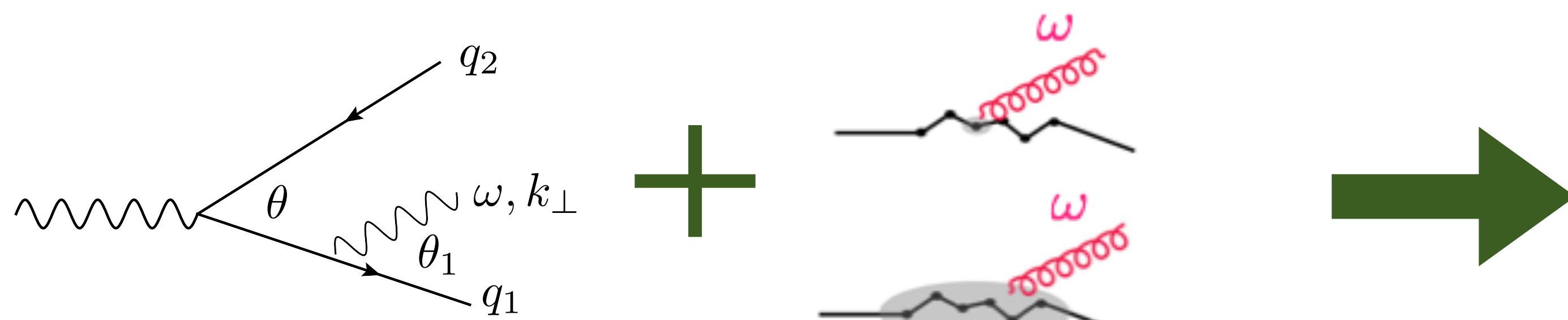
$$dN_q^{\omega \rightarrow 0} \sim \alpha_s C_R \frac{d\omega}{\omega} \frac{\sin \theta d\theta}{1 - \cos \theta} [\Theta(\cos \theta_1 - \cos \theta) + \Delta_{med} \Theta(\cos \theta - \cos \theta_1)]$$

Analytic: [Casalderrey-Solana, Iancu, Mehtar-Tani, Salgado, Tywoniuk (1105.1760, 1210.7765)]

Improving "medium" parton showers

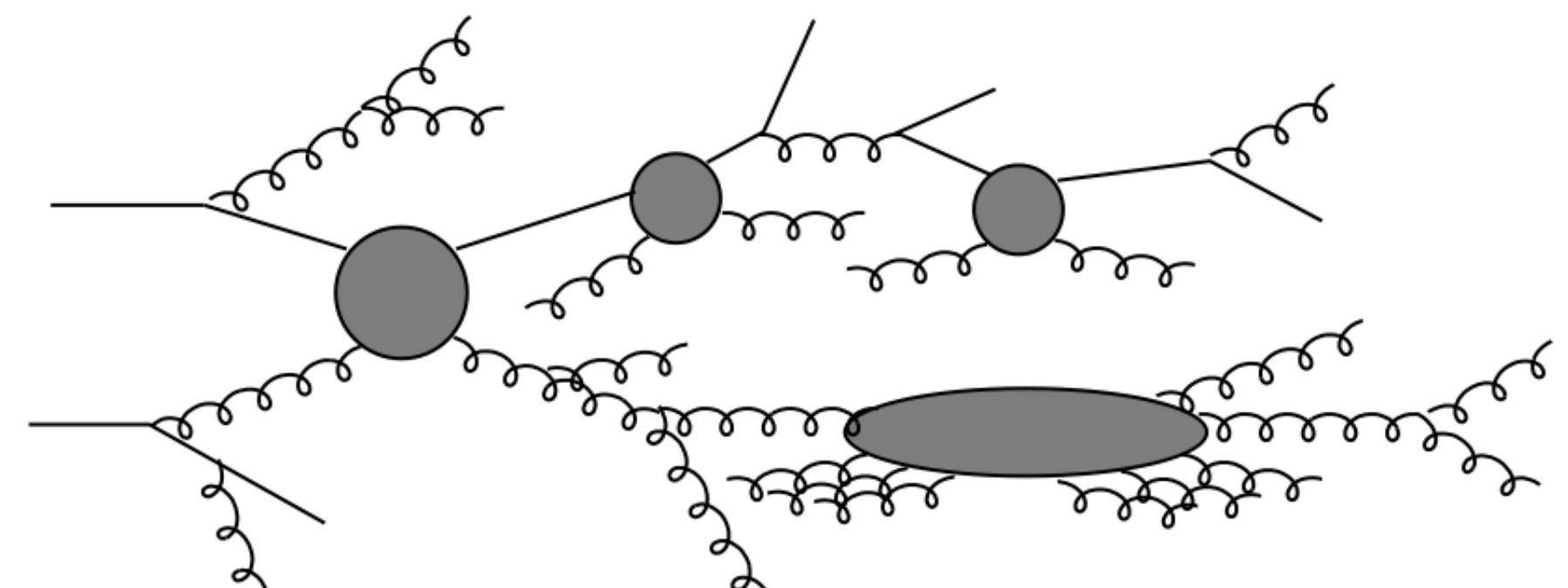


- Multiple emitters:
 - Interference effects suppressed (+ anti-angular ordering)
 - Non-instantaneous emissions will induce modifications to the vacuum parton shower structure:



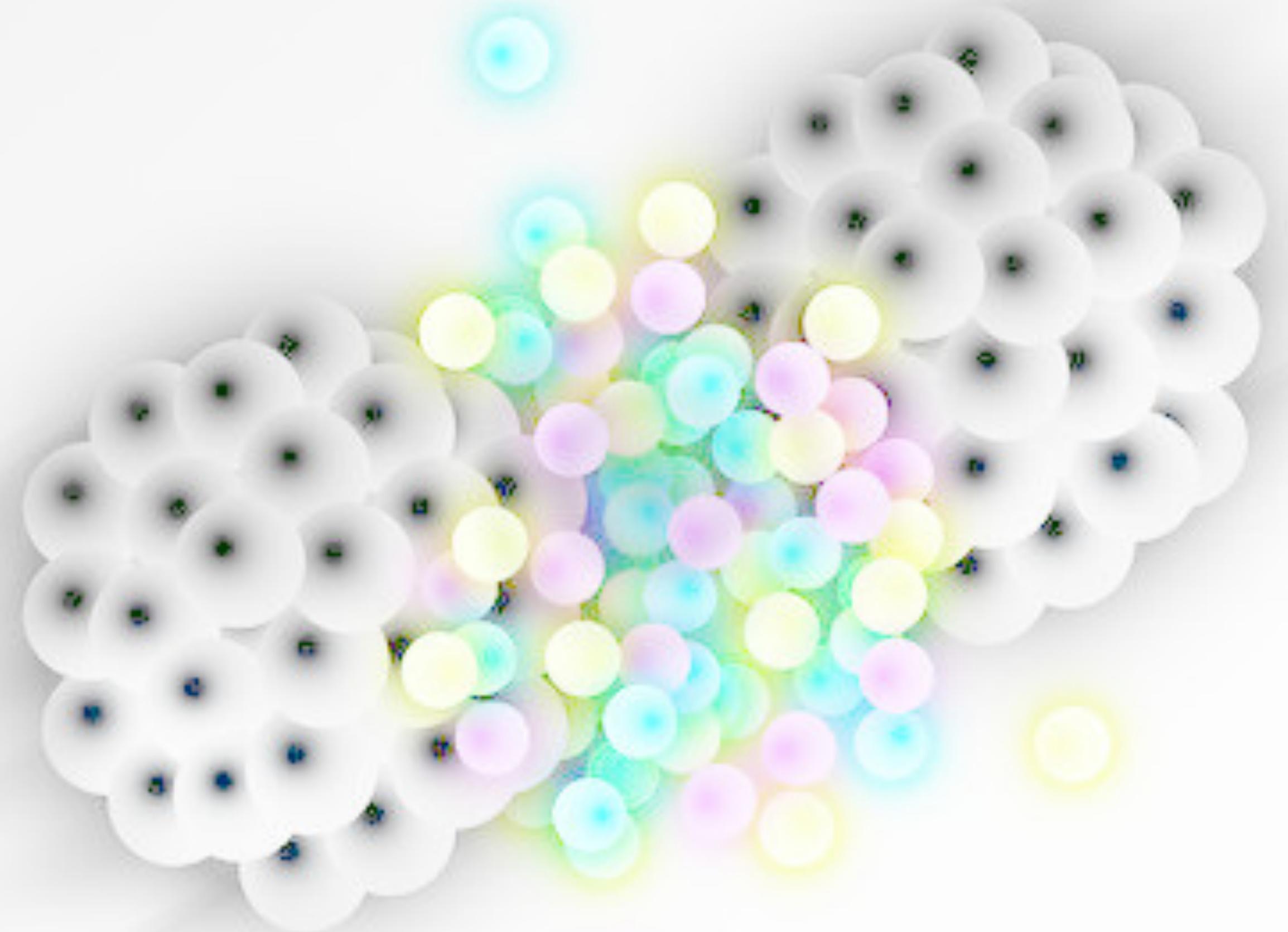
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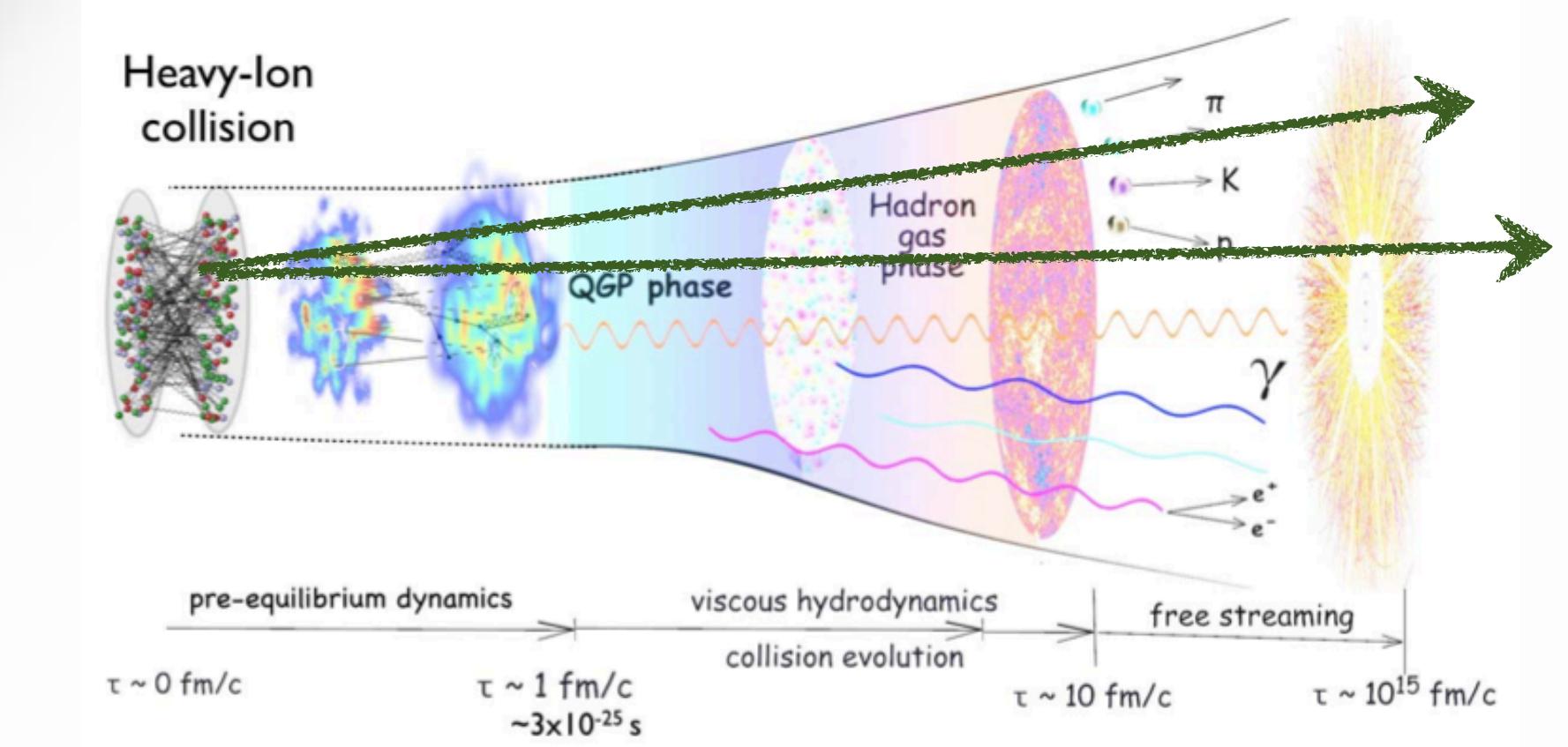
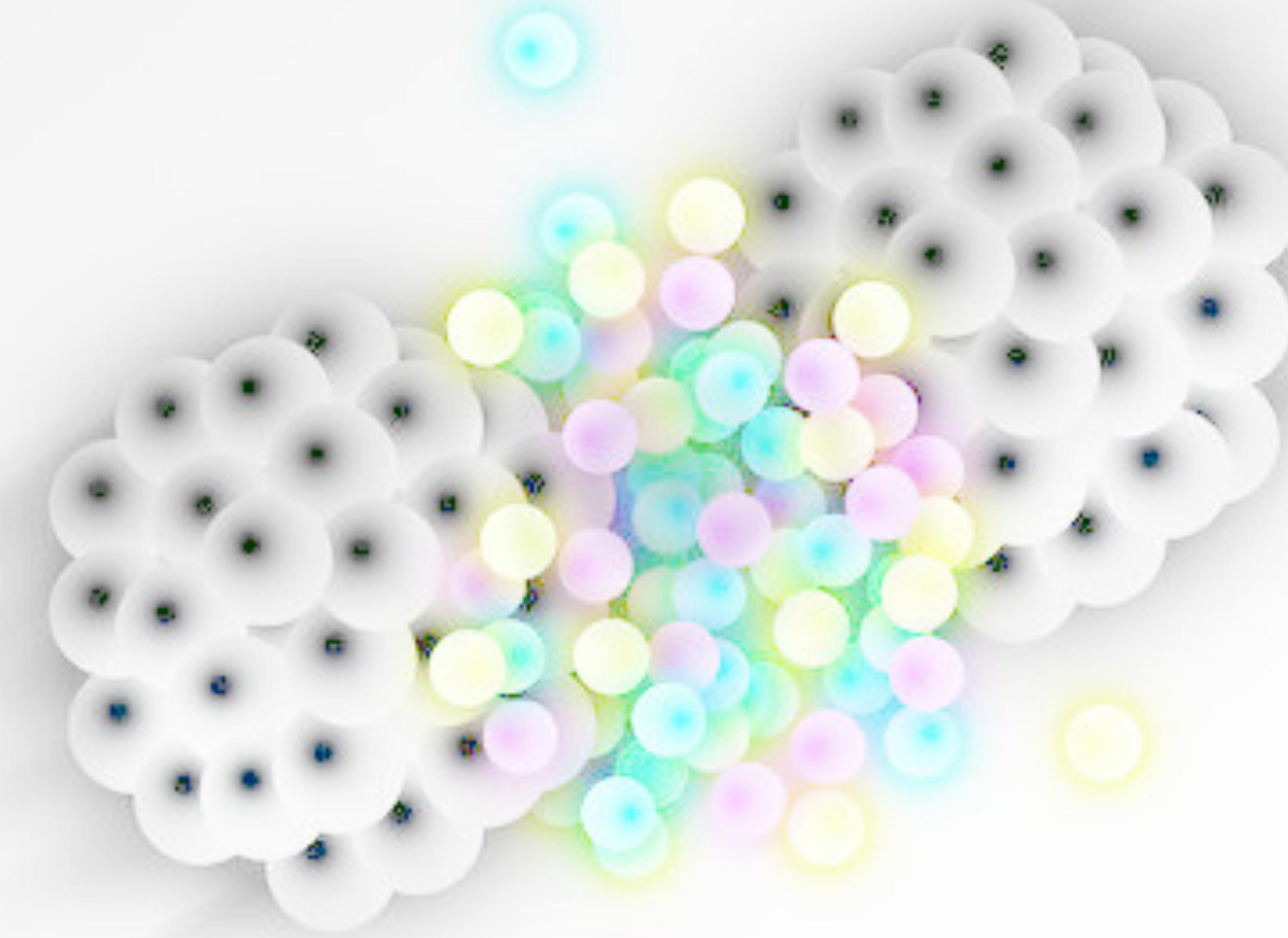


Monte Carlo: [Q-PYTHIA, JEWEL]
[Armesto, Cunqueiro, Salgado (0907.1014),
Zapp (1311.0048)]

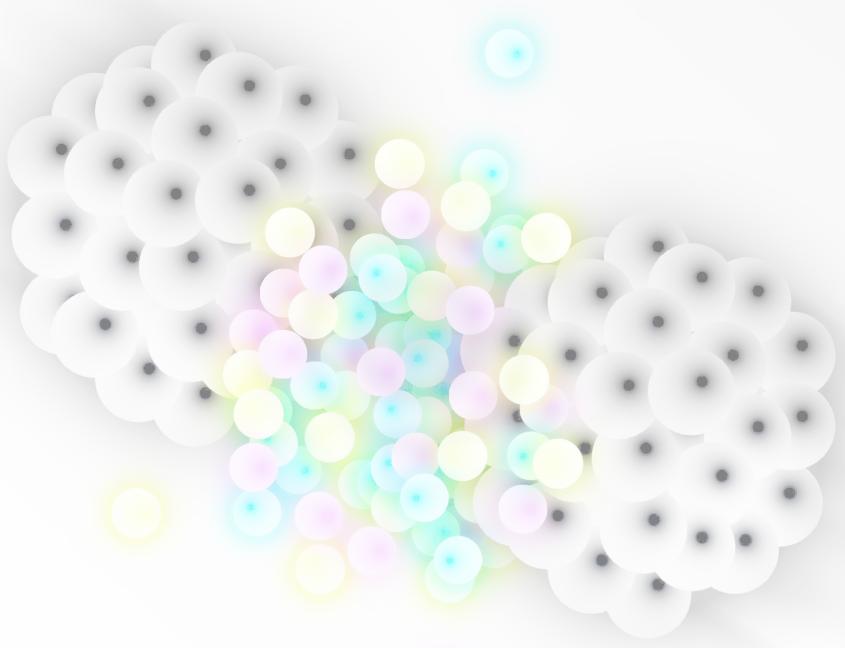
New experimental handles



New experimental handles

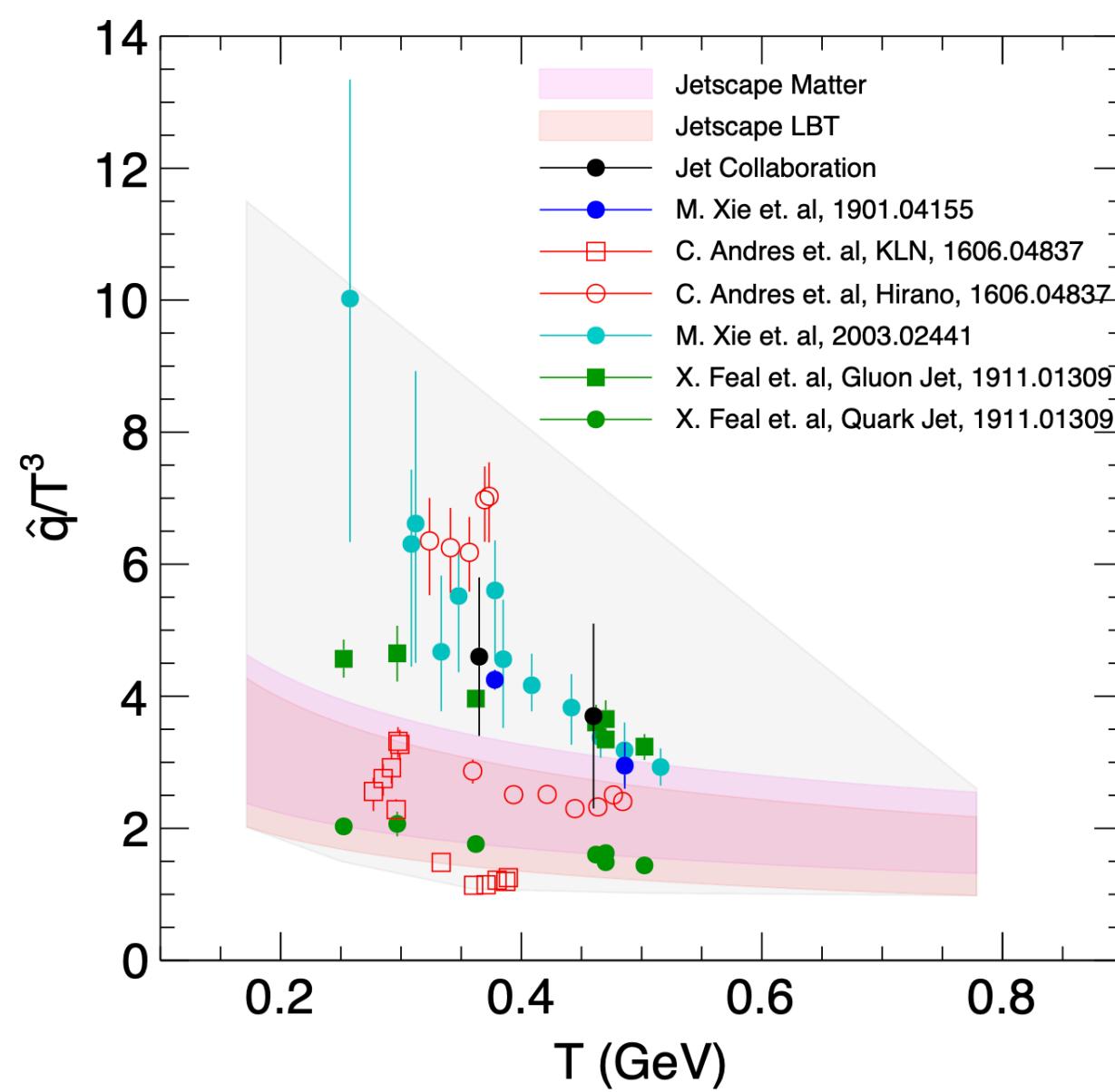
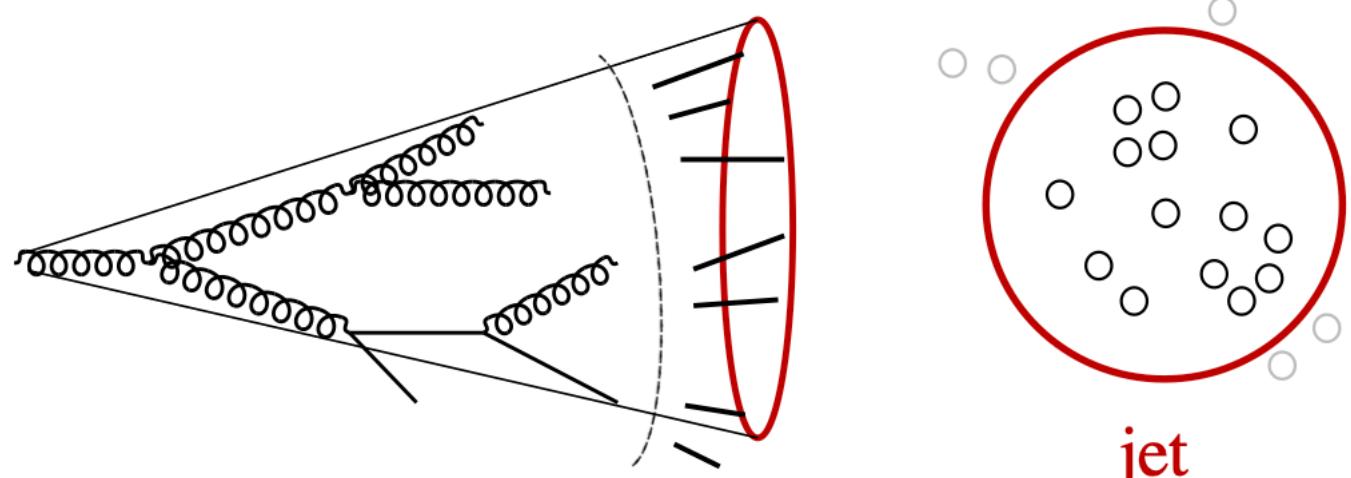
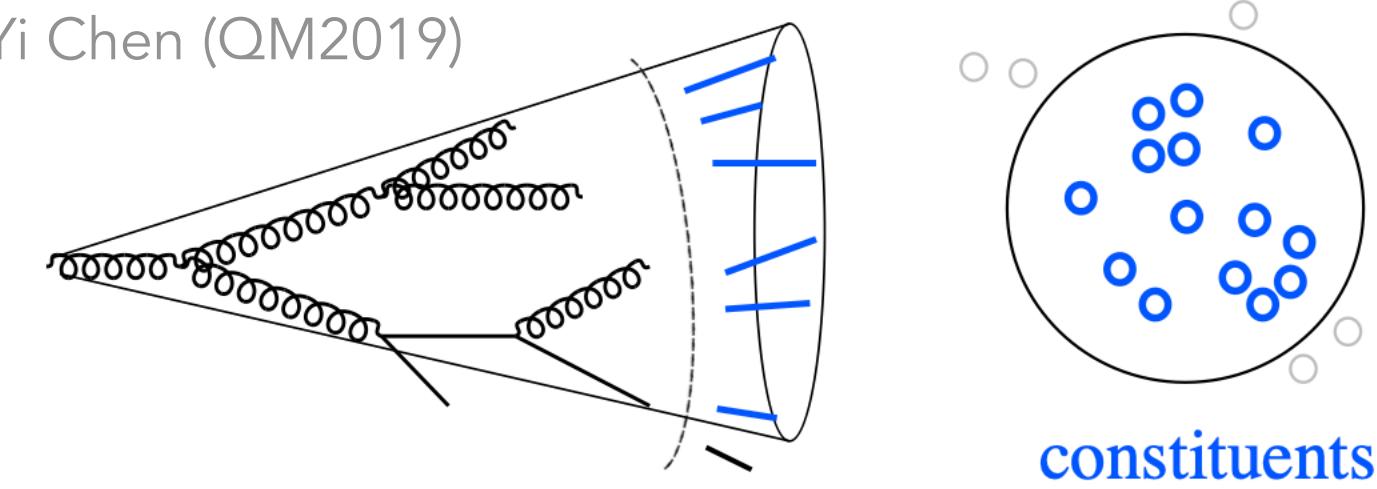


From particles to jets

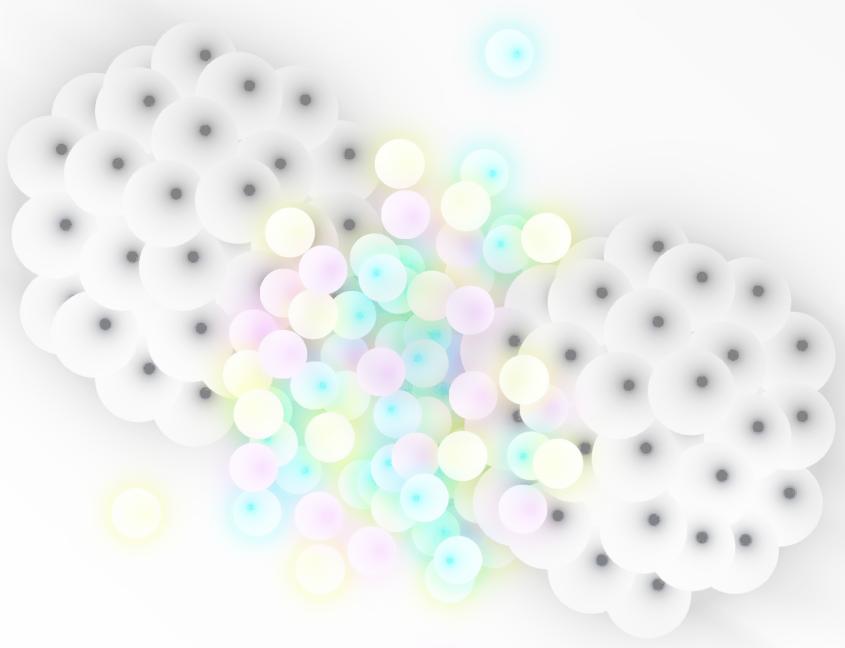


- How can we access QGP-related information?

[Adapted from Yi Chen (QM2019)]

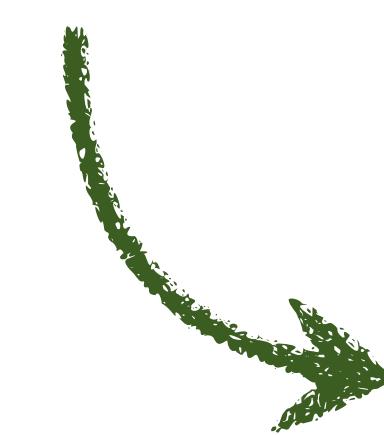
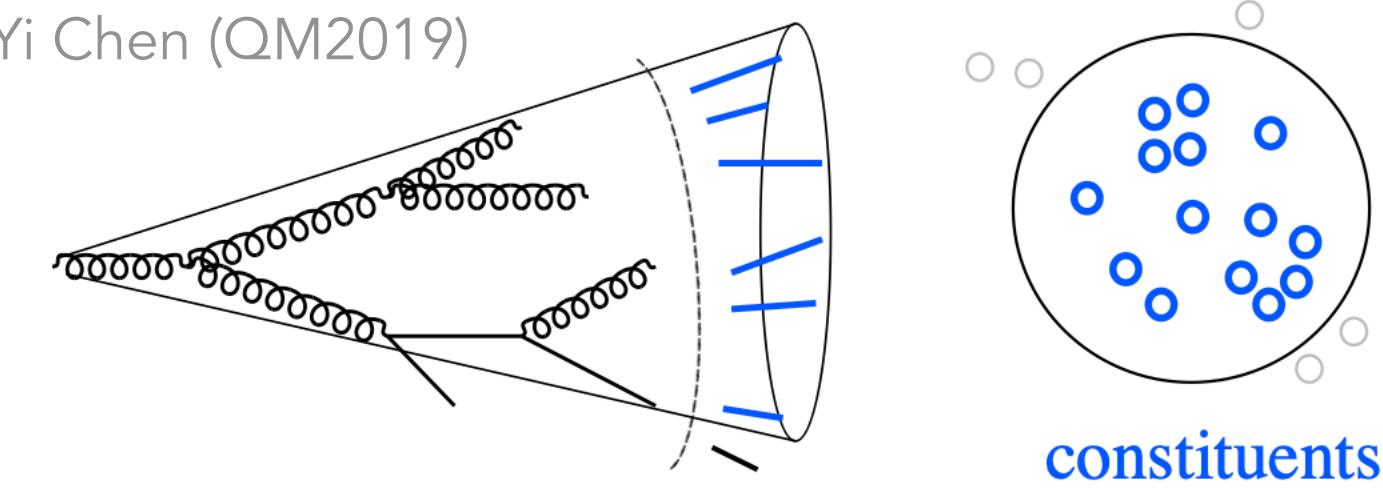


From particles to jets

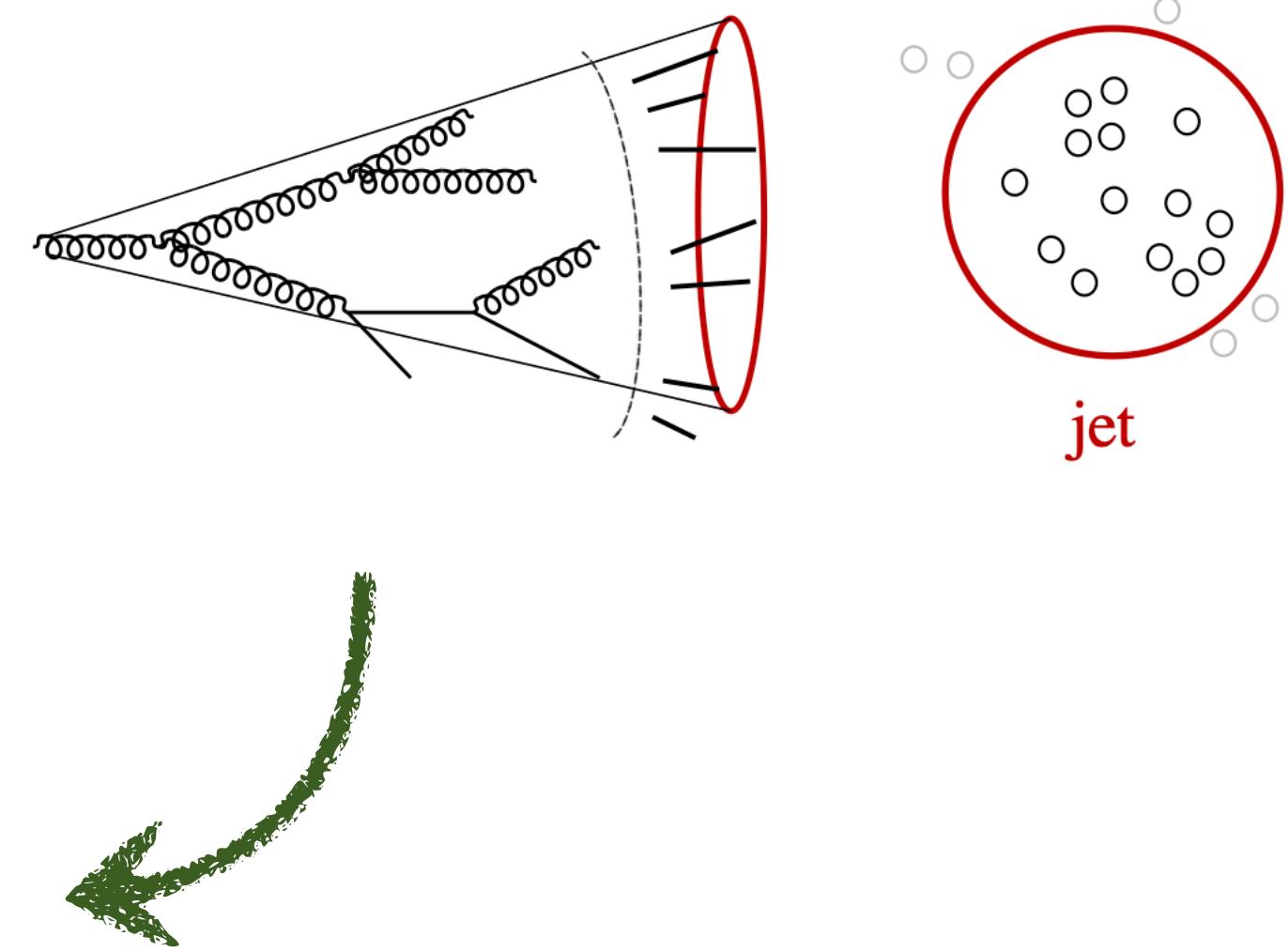
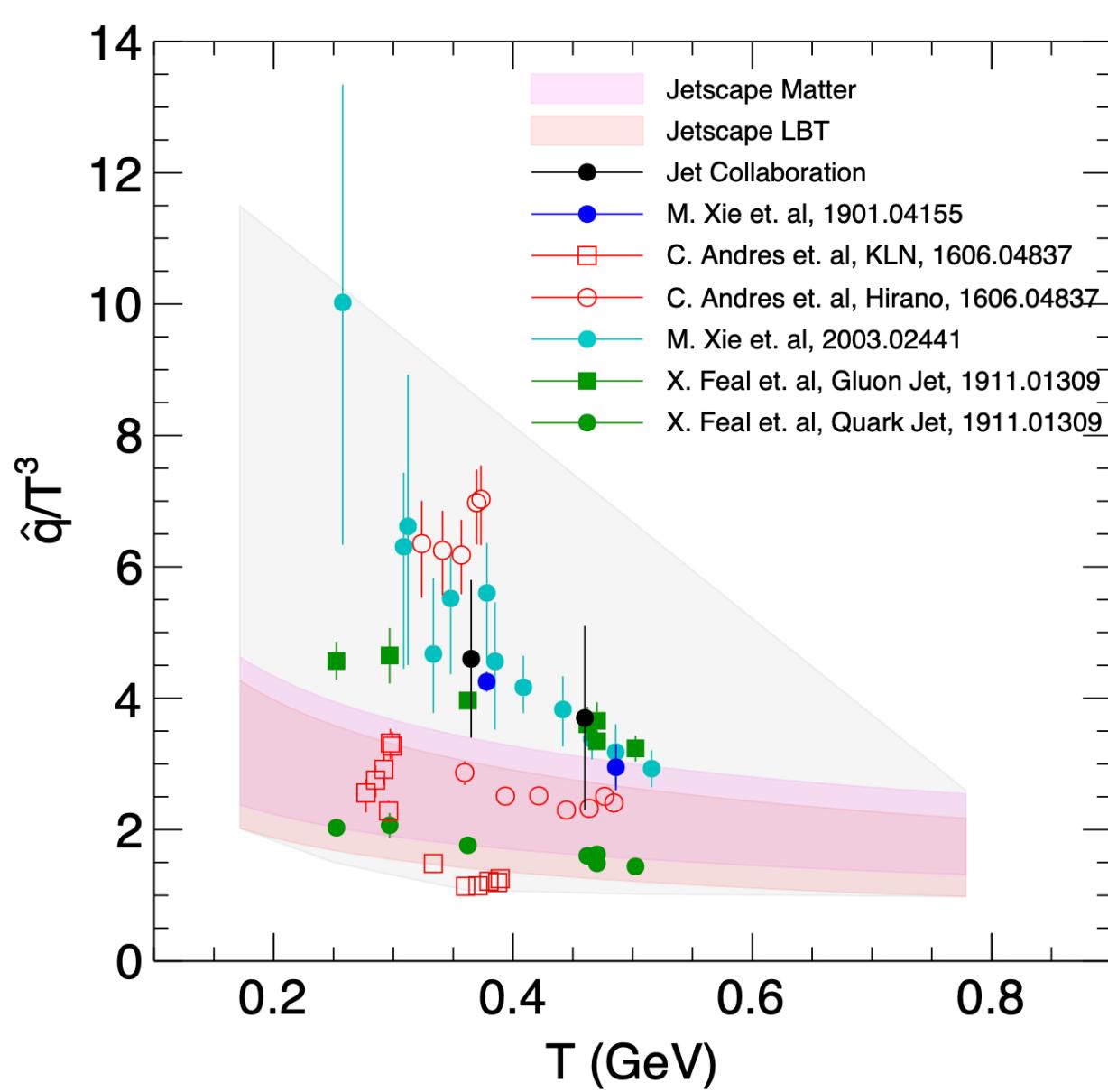


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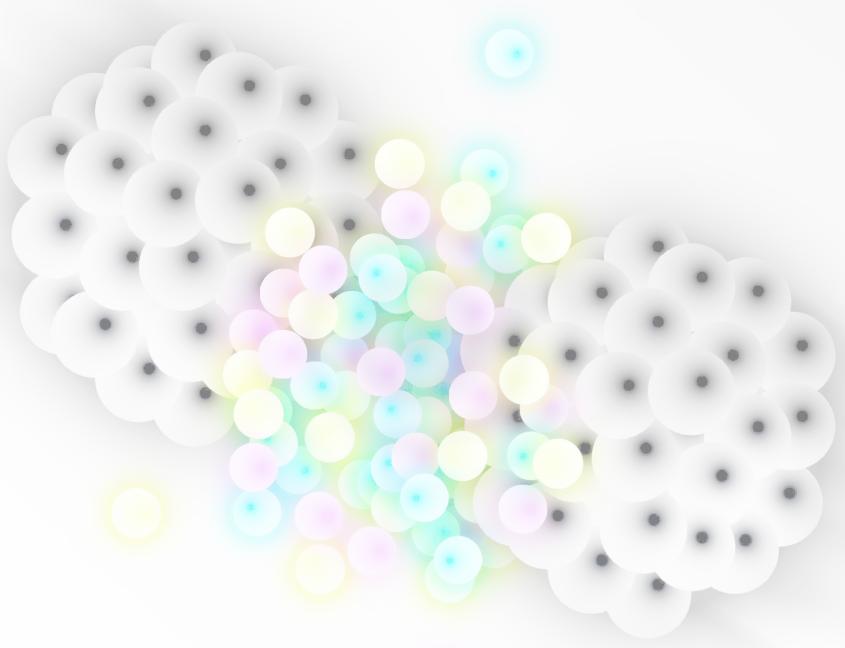


However:
- Susceptible to hadronization effects...



However:
- Sensitive to average quantities...

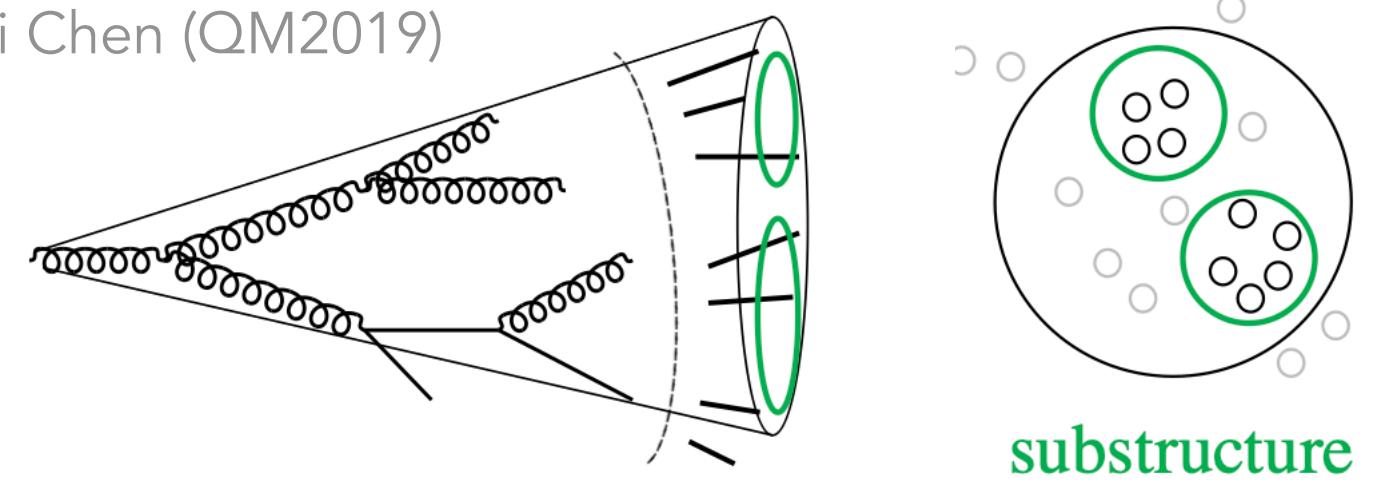
From jets to jet substructure



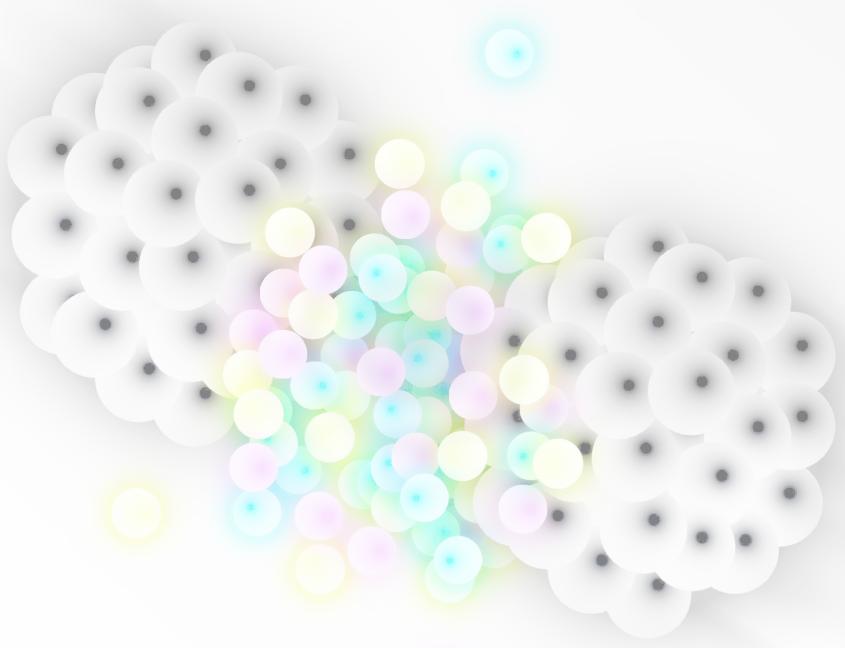
- How can we access QGP-related information?

What more information can they provide?

[Adapted from Yi Chen (QM2019)]

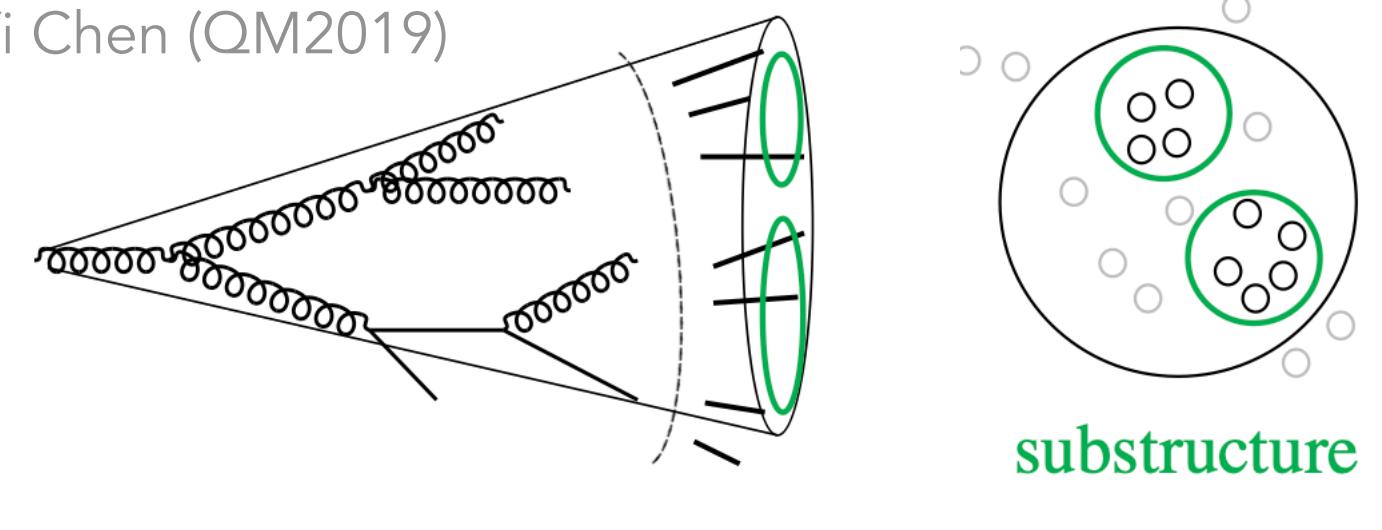


From jets to jet substructure



- How can we access QGP-related information?

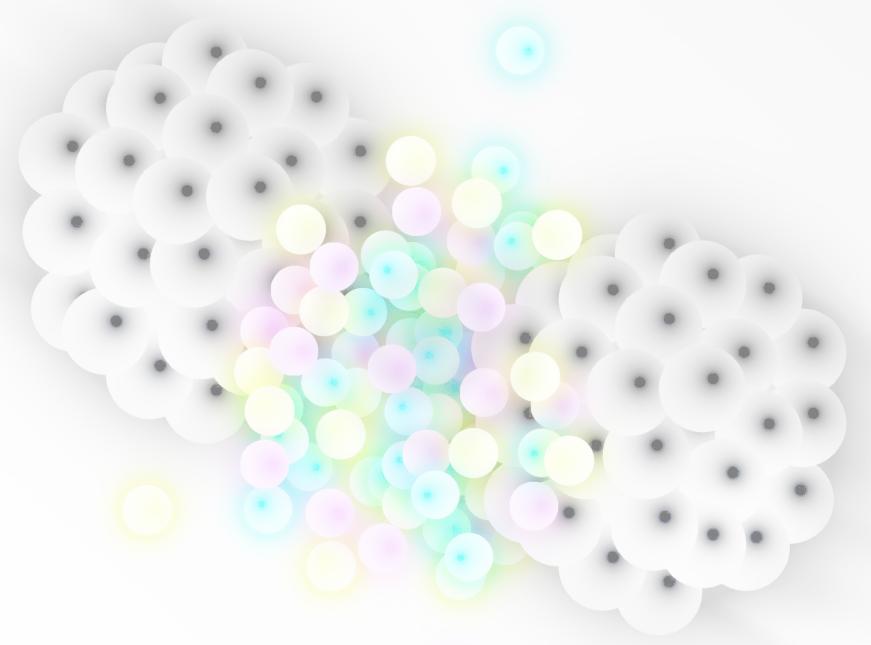
[Adapted from Yi Chen (QM2019)]



What more information can they provide?



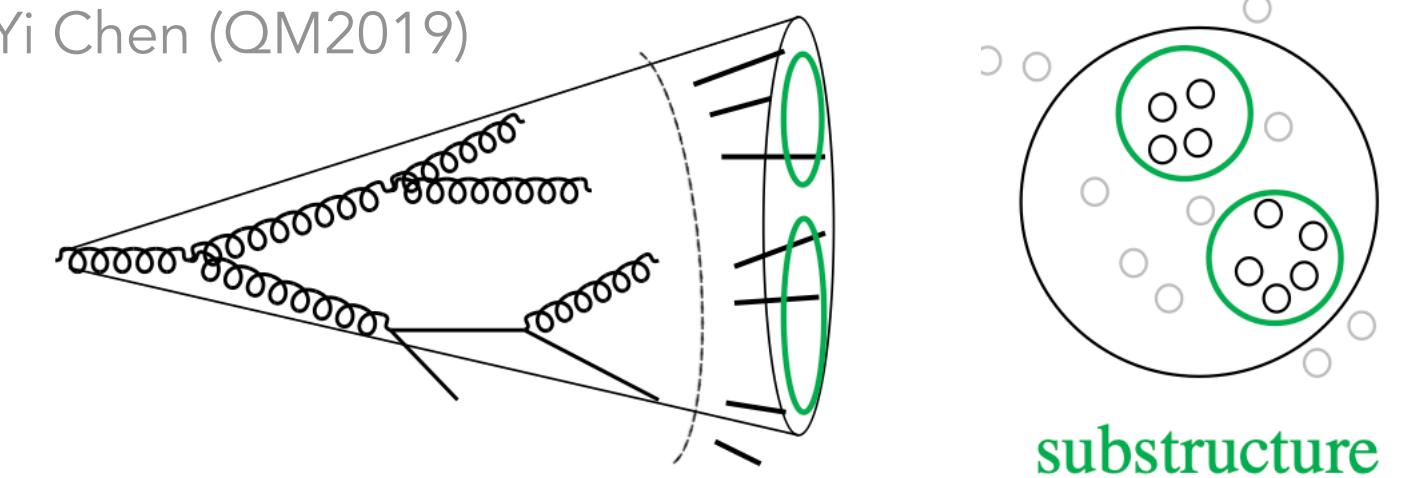
From jets to jet substructure



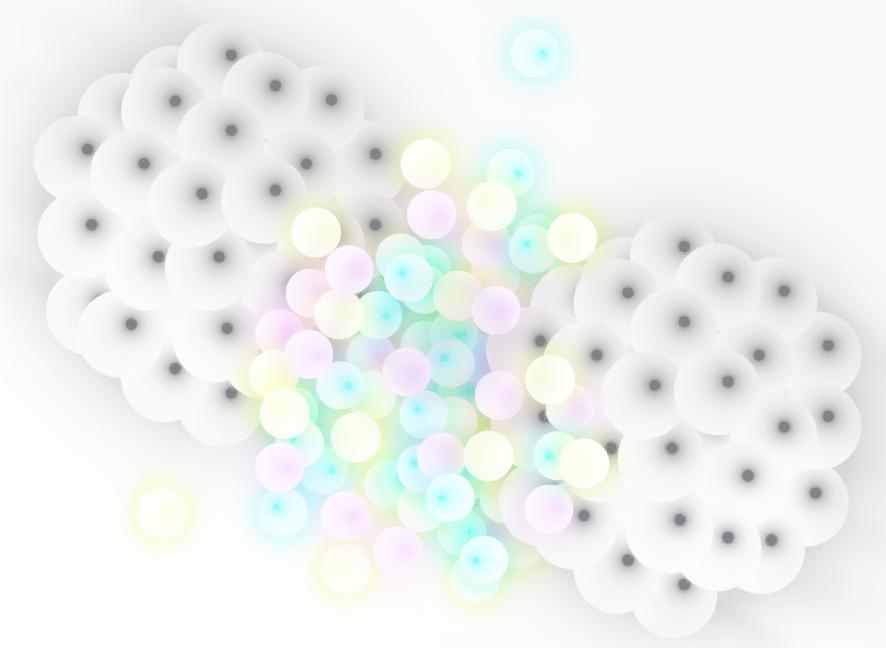
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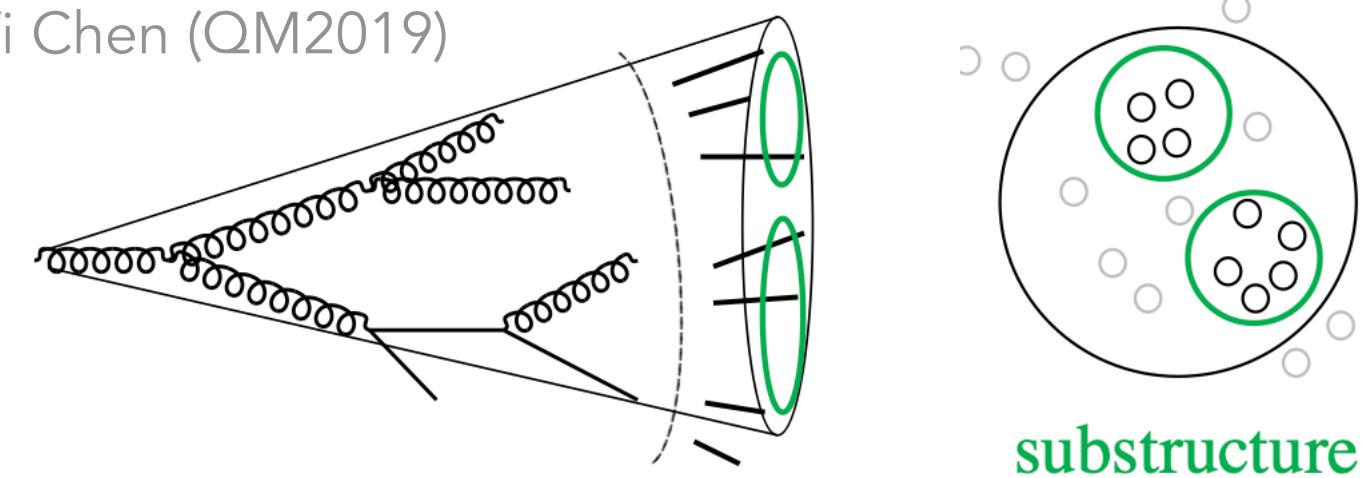


From jets to jet substructure

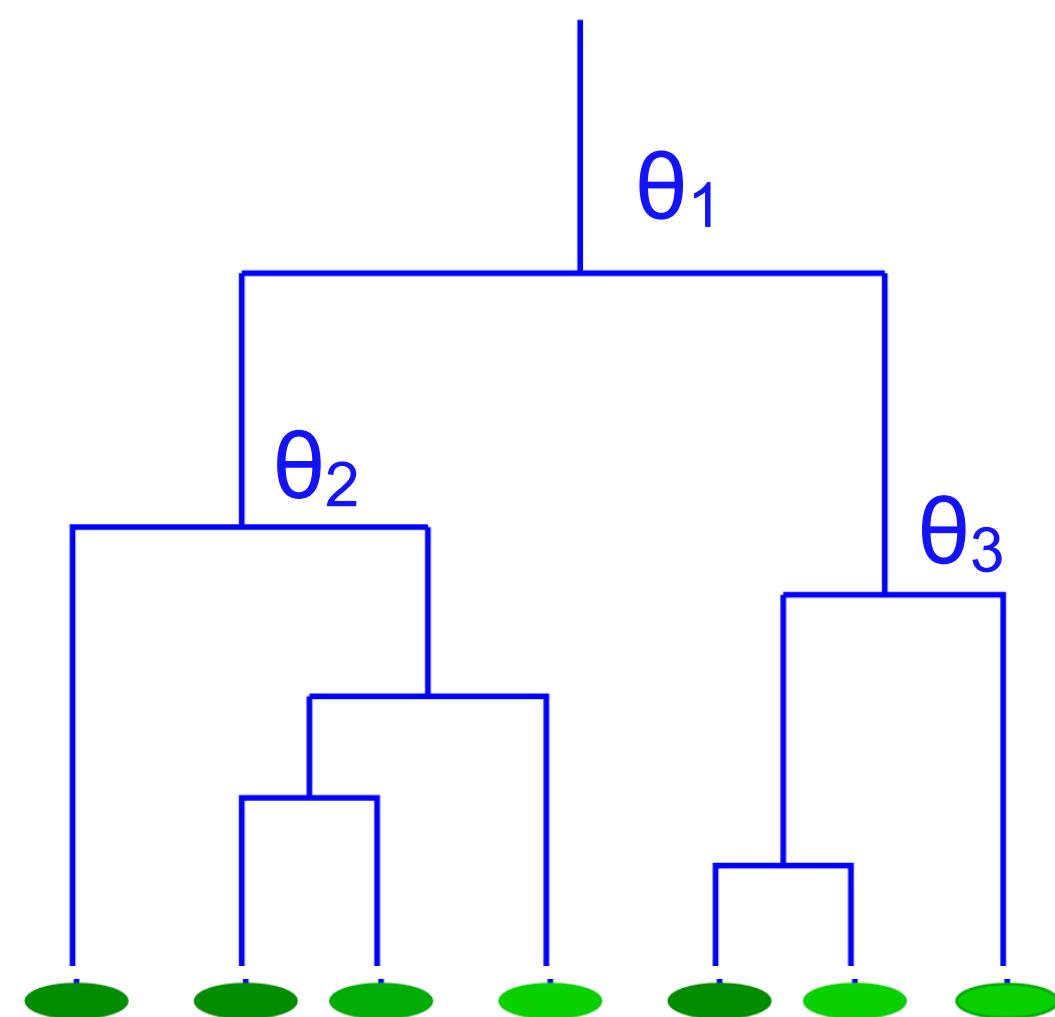


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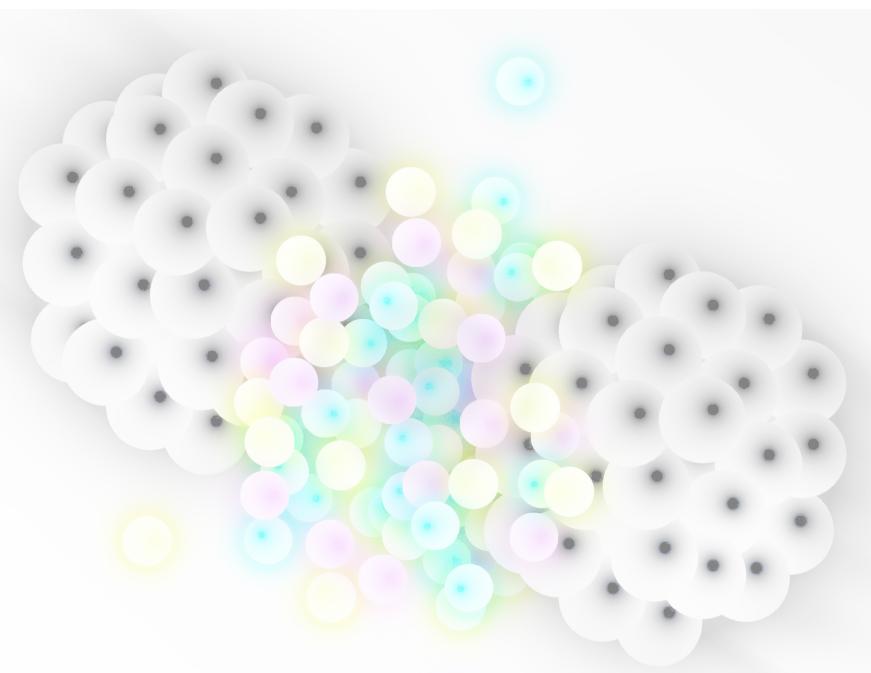


What more information can they provide?



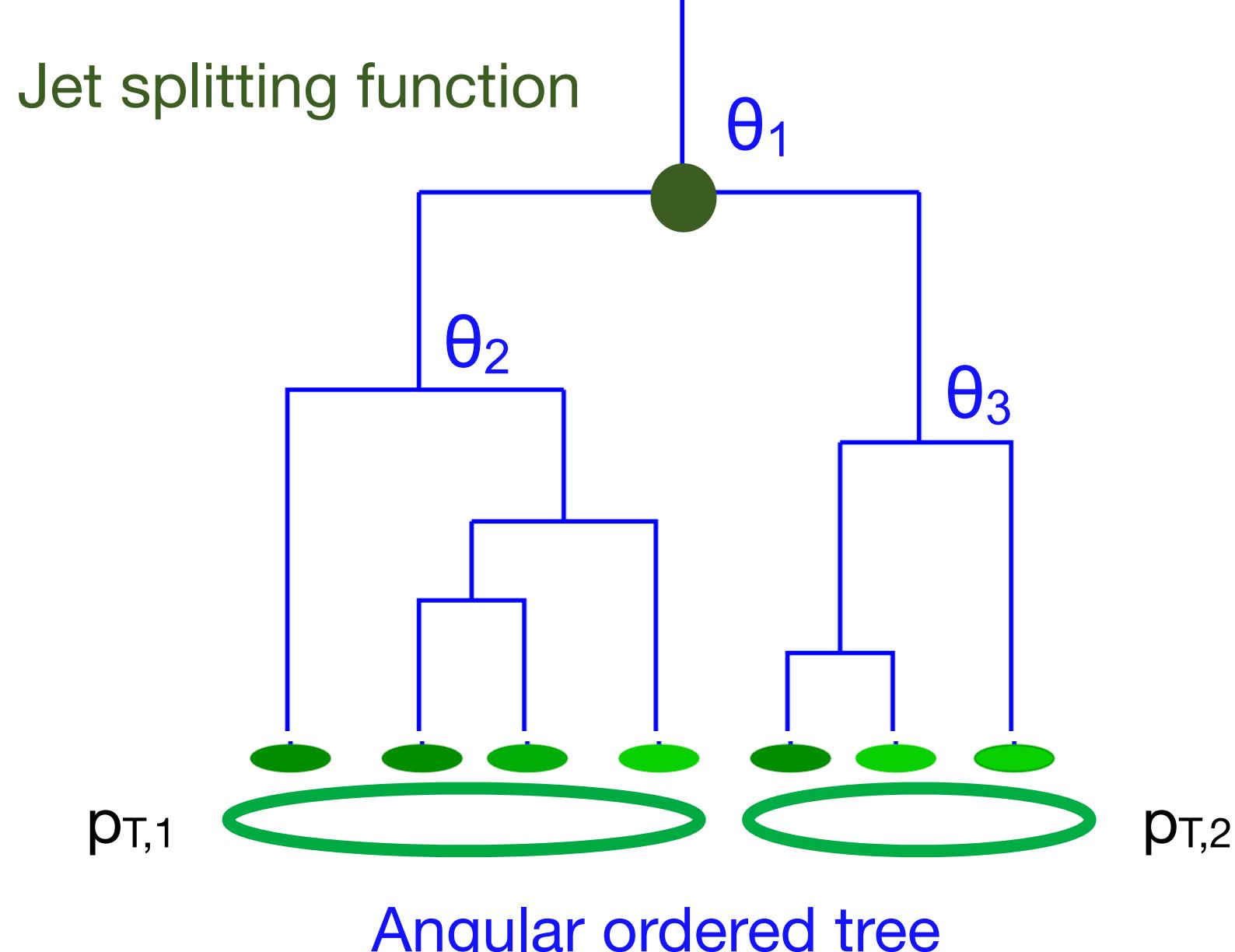
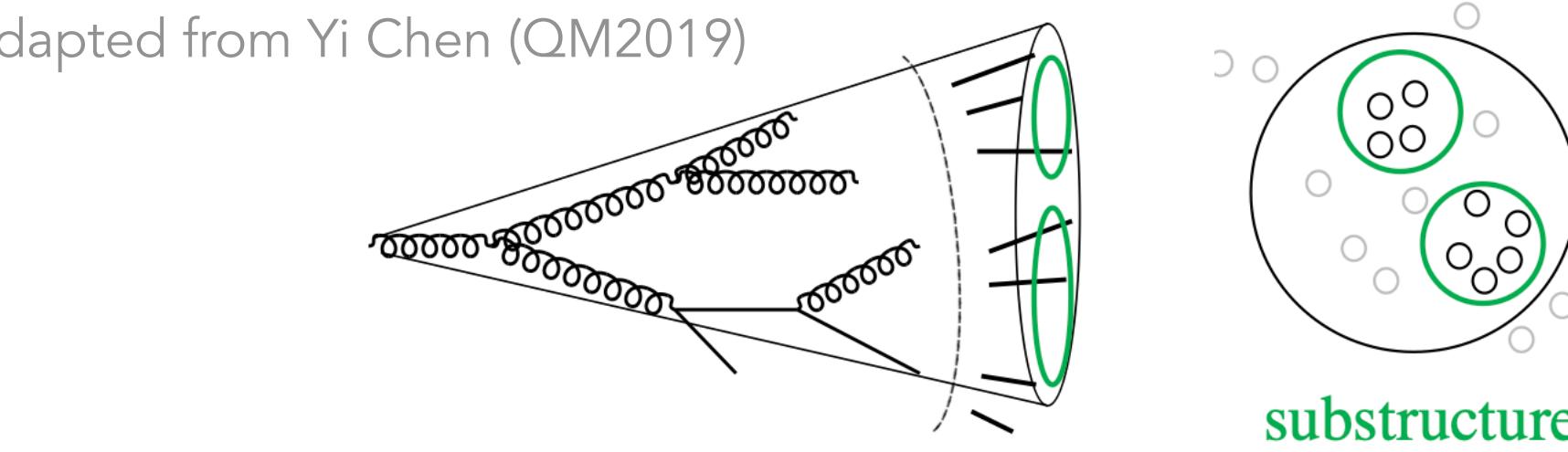
Angular ordered tree

From jets to jet substructure



- How can we access QGP-related information?

[Adapted from Yi Chen (QM2019)]

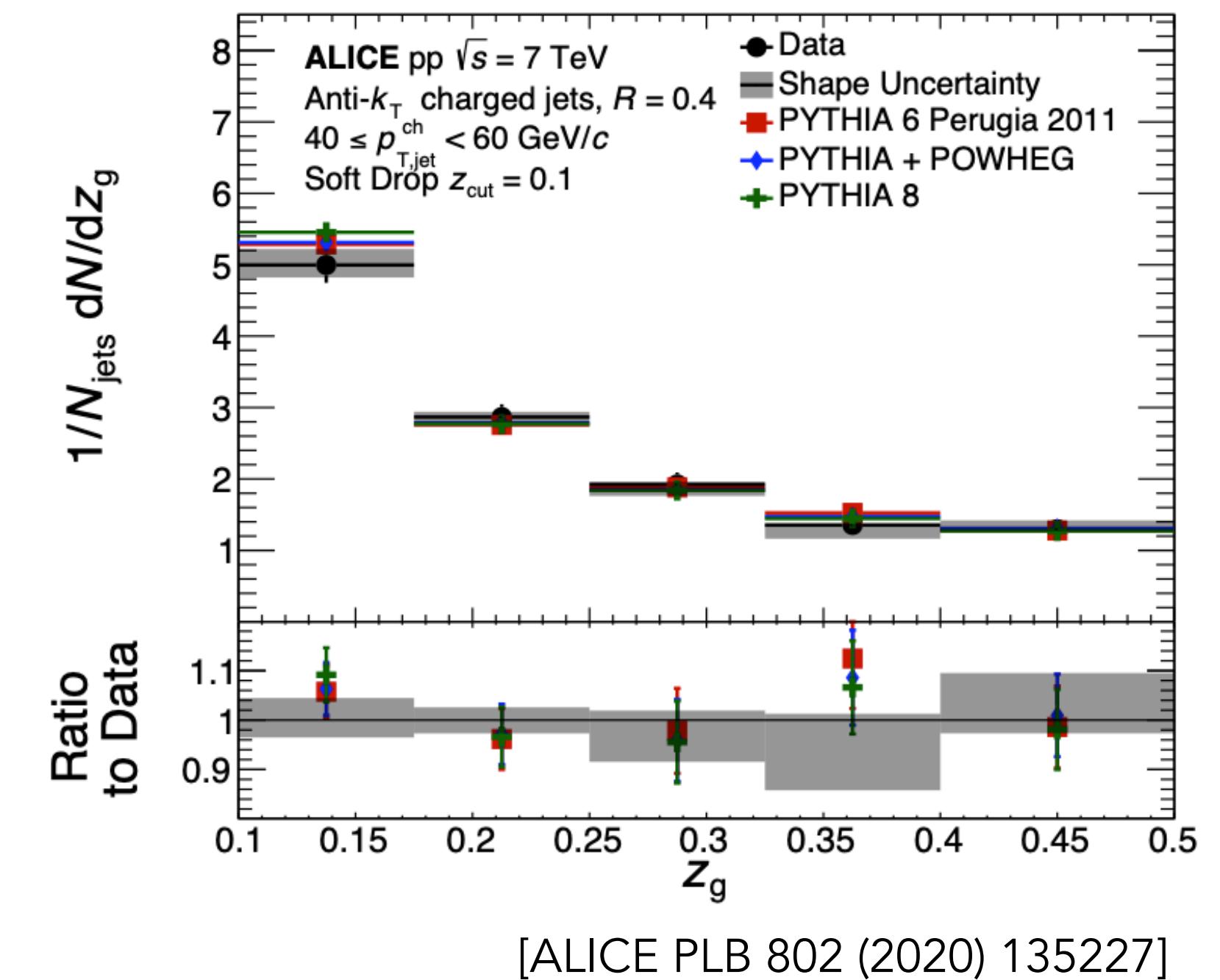


$$\text{when } \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}} > z_{cut} \left(\frac{R_{12}}{R_0} \right)^\beta$$

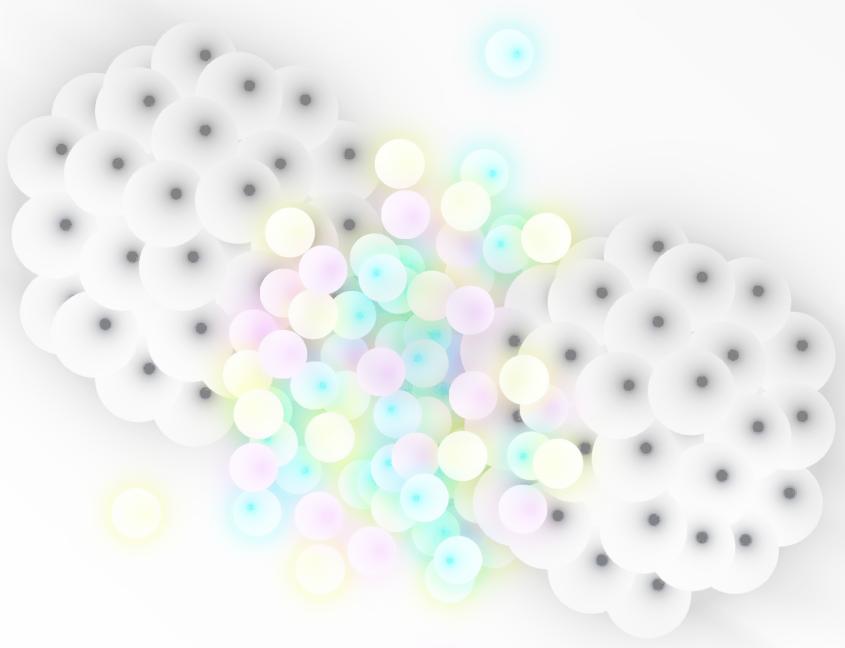
[Larkoski, Marzani, Soyez, Thaler (1402.2657)]
 [Dasgupta, Fregoso, Marzani, Salam (1307.0007)]

$$z_g = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$$

What more information can they provide?

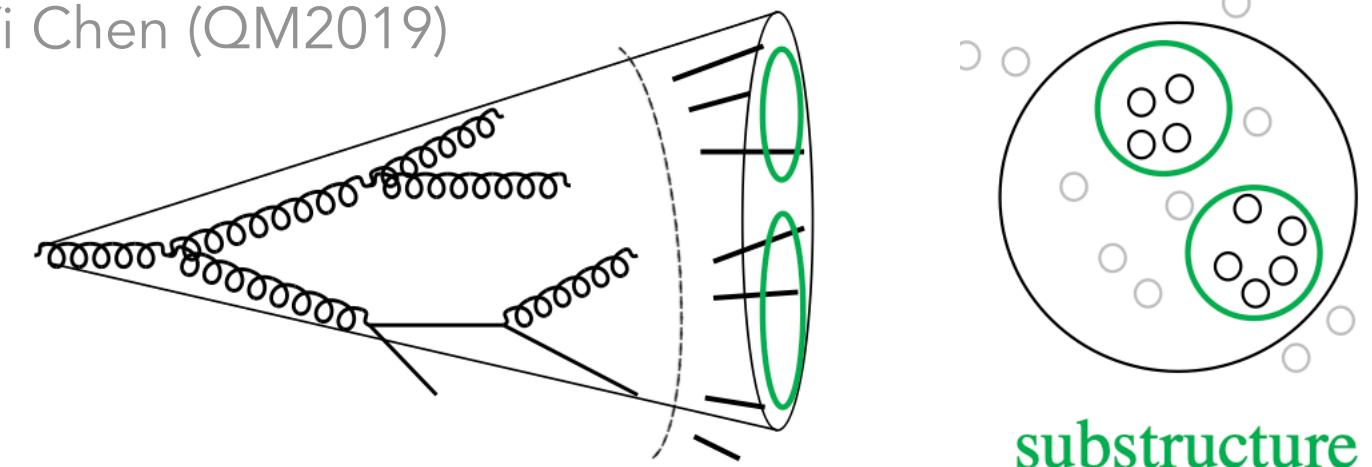


From jets to jet substructure



- *How can we access QGP-related information?*

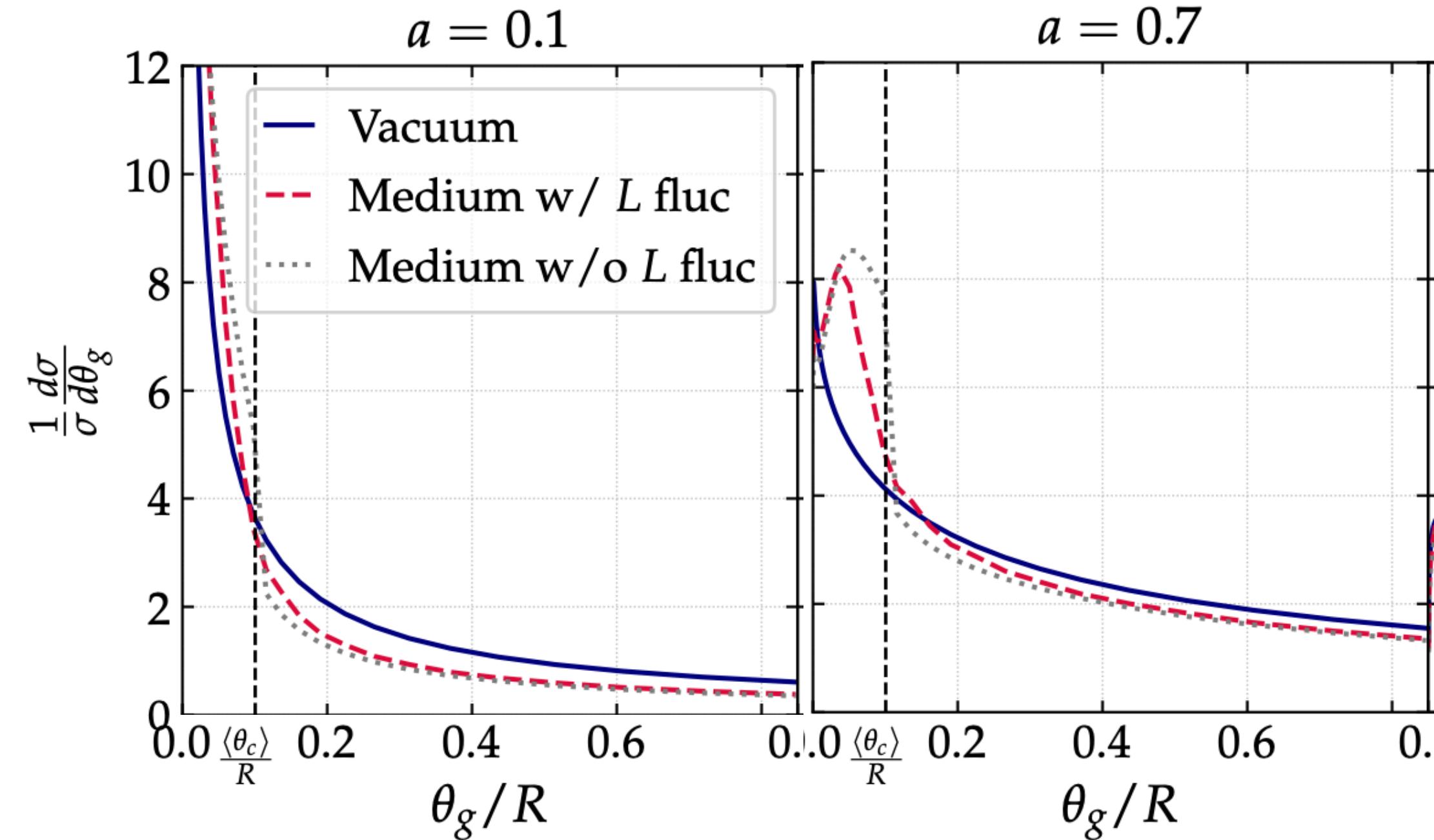
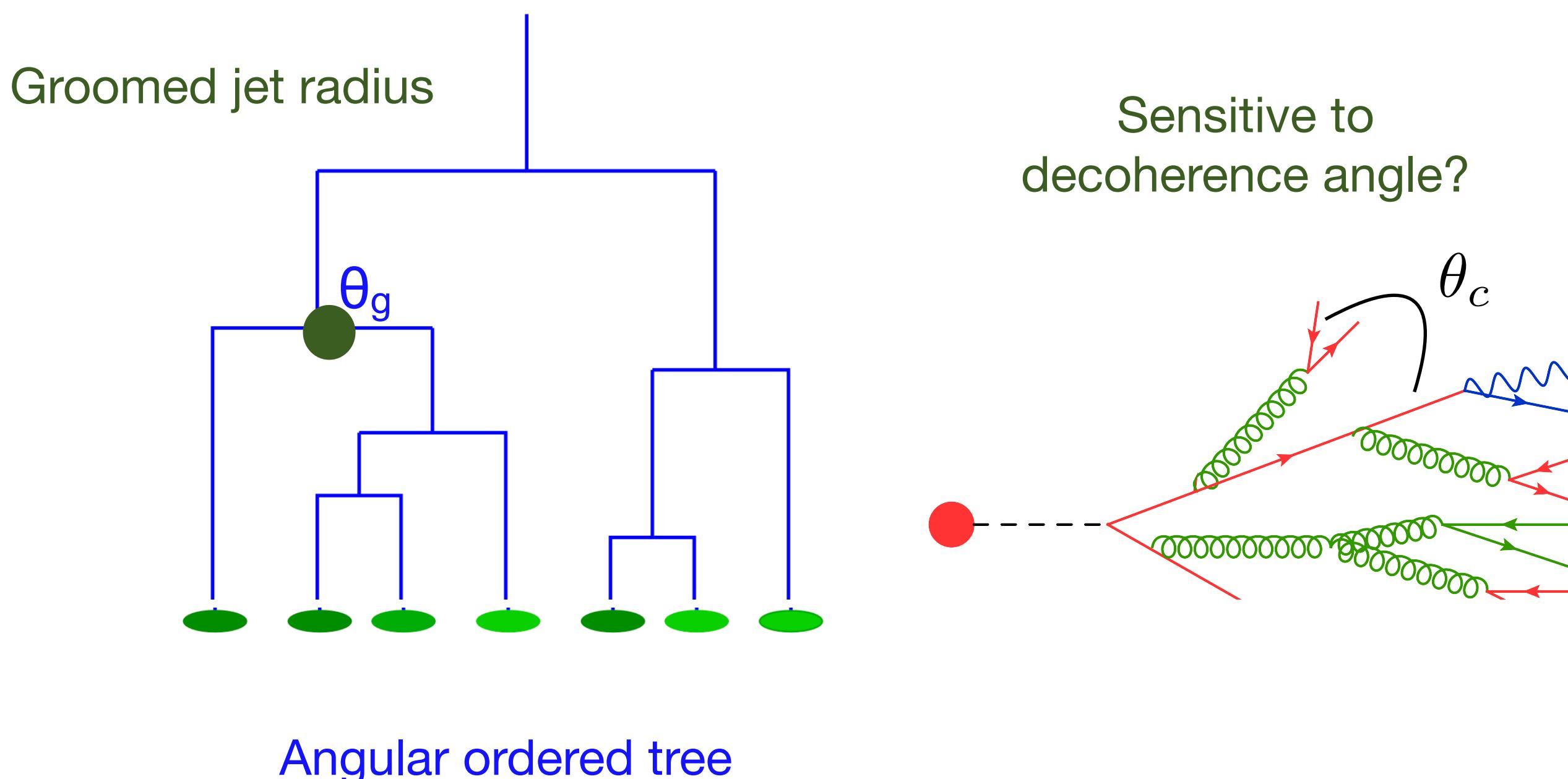
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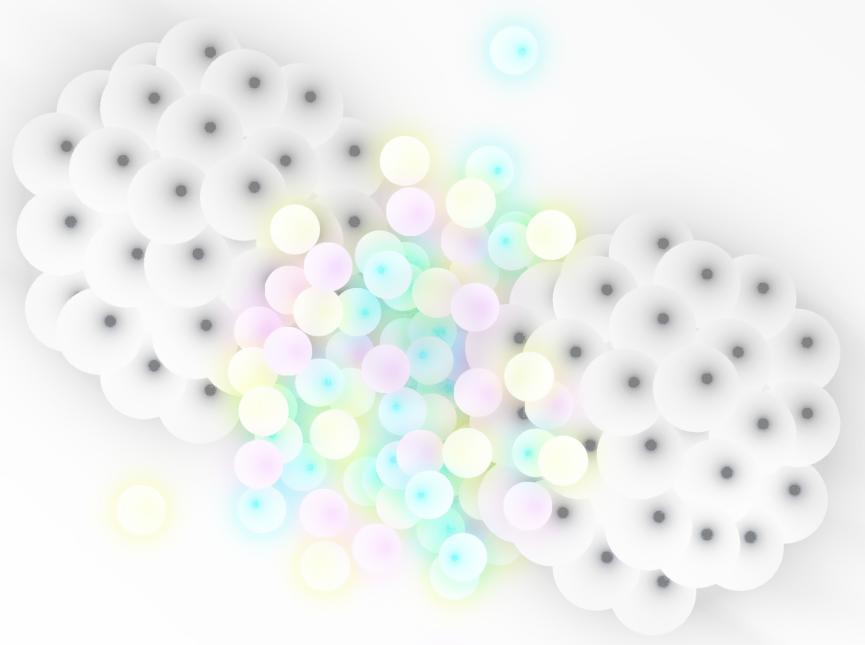
$$\kappa^{(a)} = \frac{1}{p_{t,\text{jet}}} z(1-z)p_t \left(\frac{\theta}{R}\right)^a$$

[Mehtar-Tani, Soto-Ontoso, K. Tywoniuk
(1911.00375)]

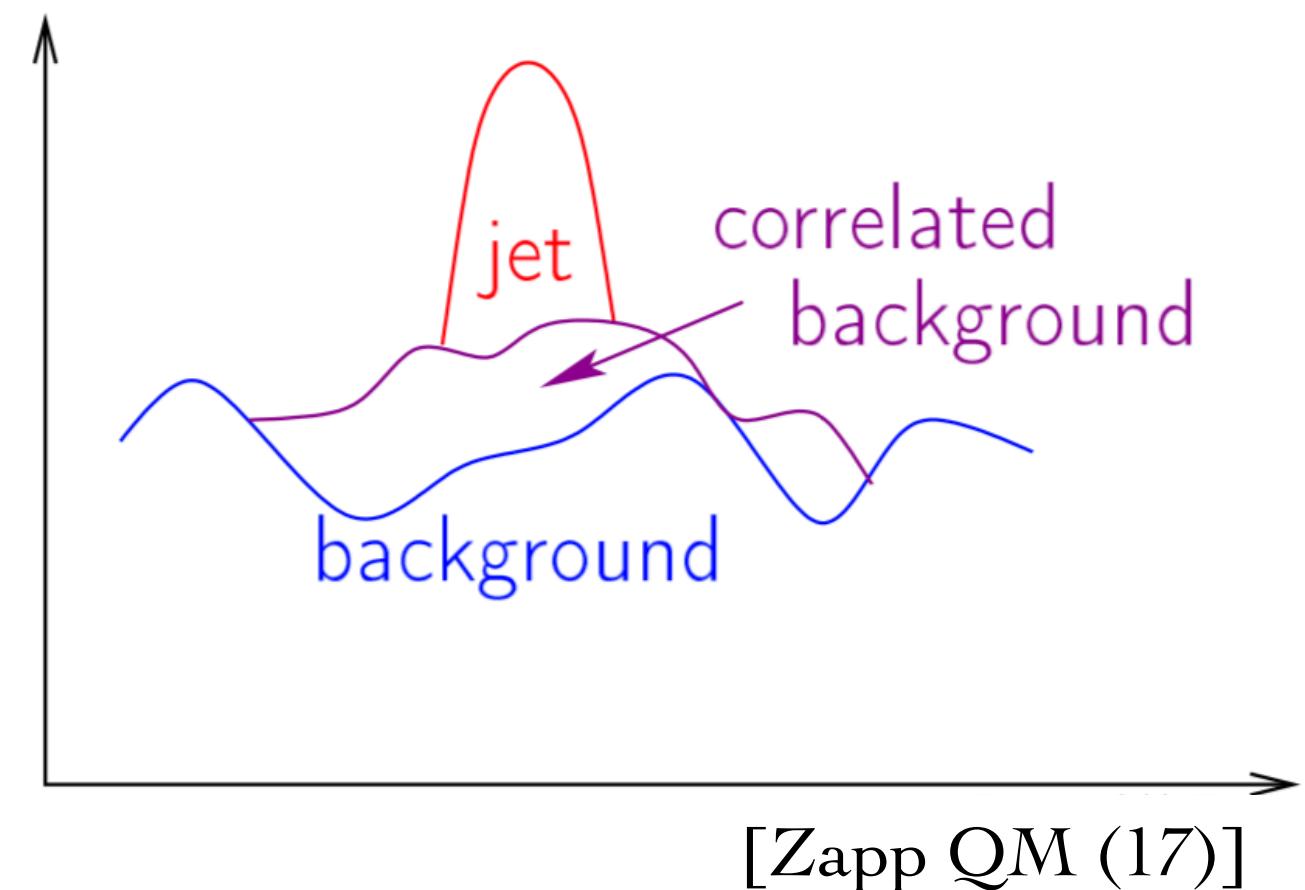
What more information can they provide?



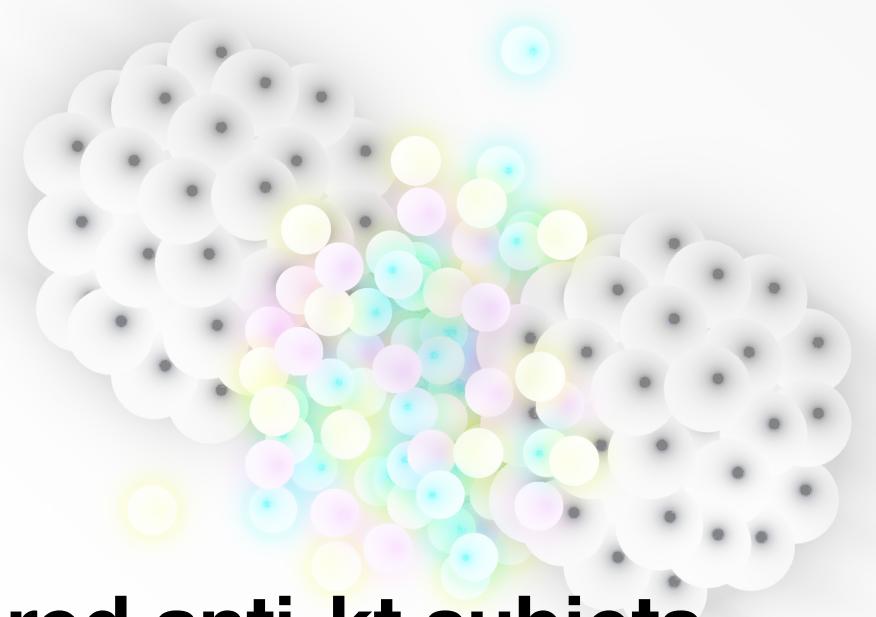
From jets to jet substructure



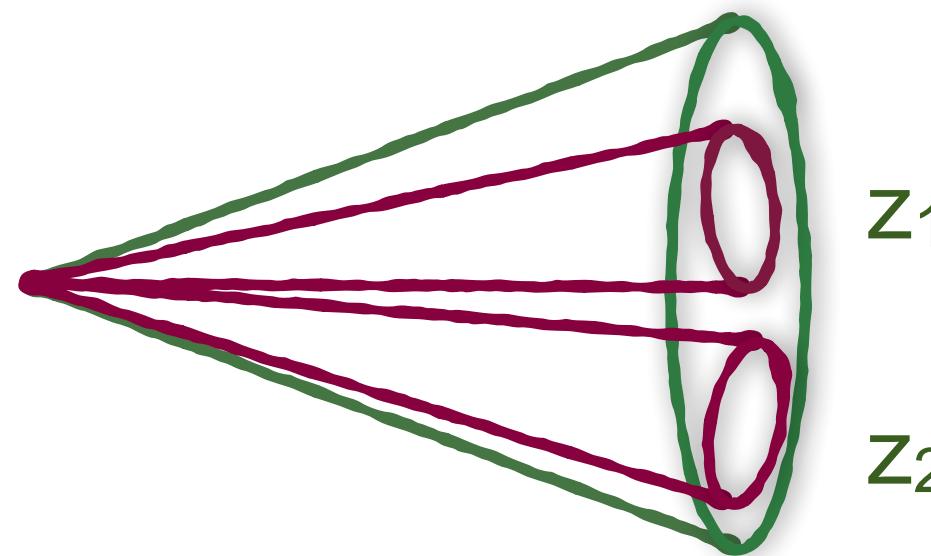
- How can we access QGP-related information?
 - Jets in PbPb \neq Jets in pp + Background



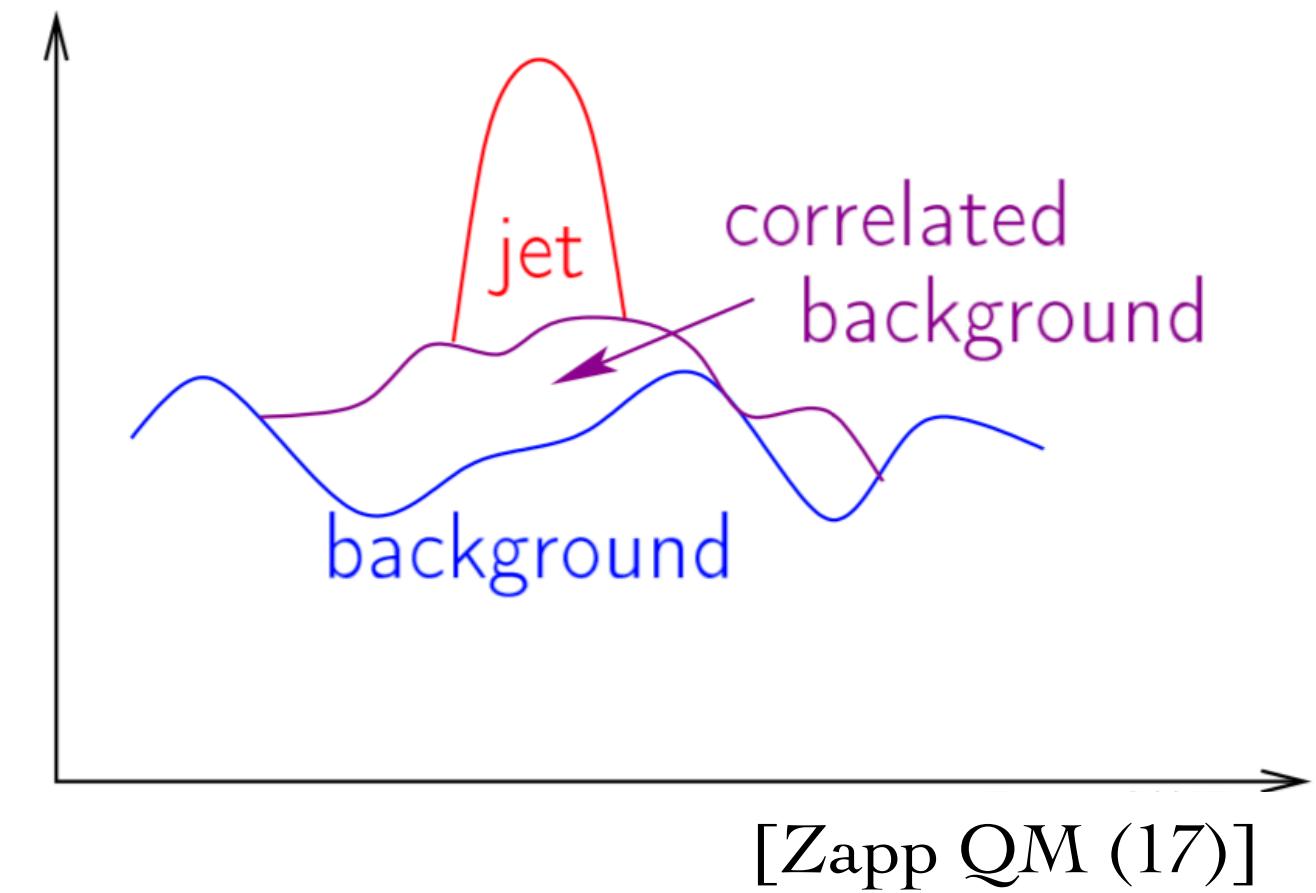
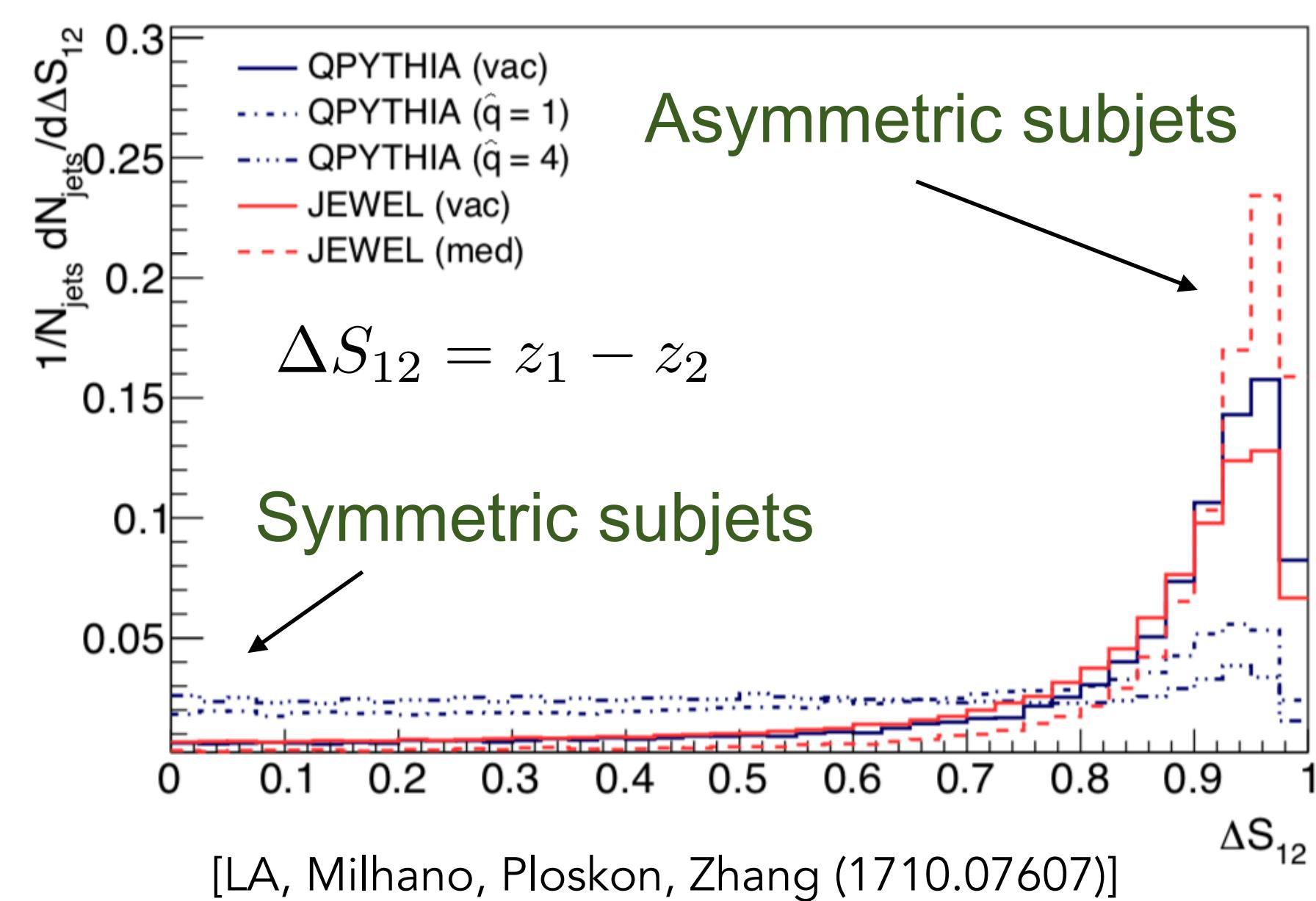
From jets to jet substructure



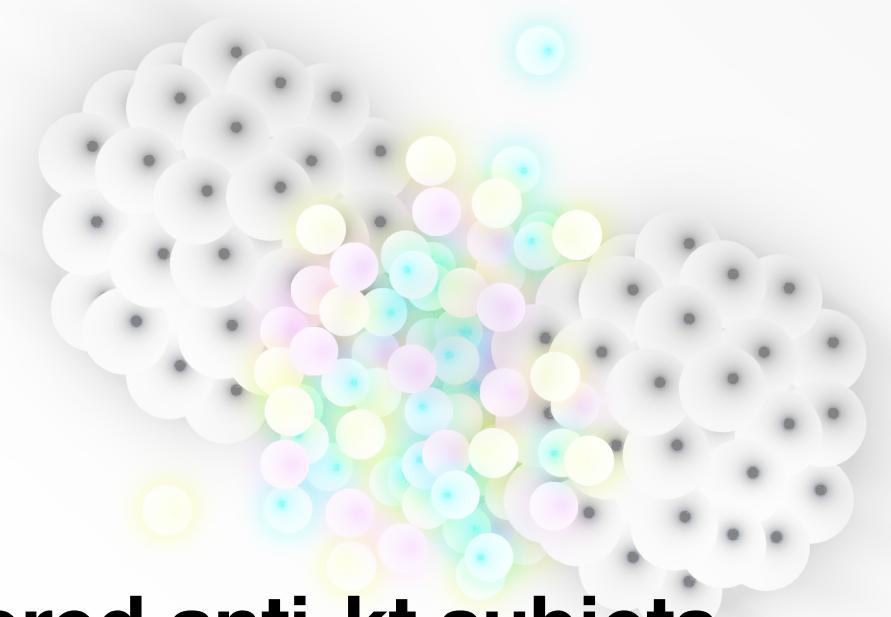
- How can we access QGP-related information?
 - Jets in PbPb \neq Jets in pp + Background
 - Background-resilient to distinguish quenching models



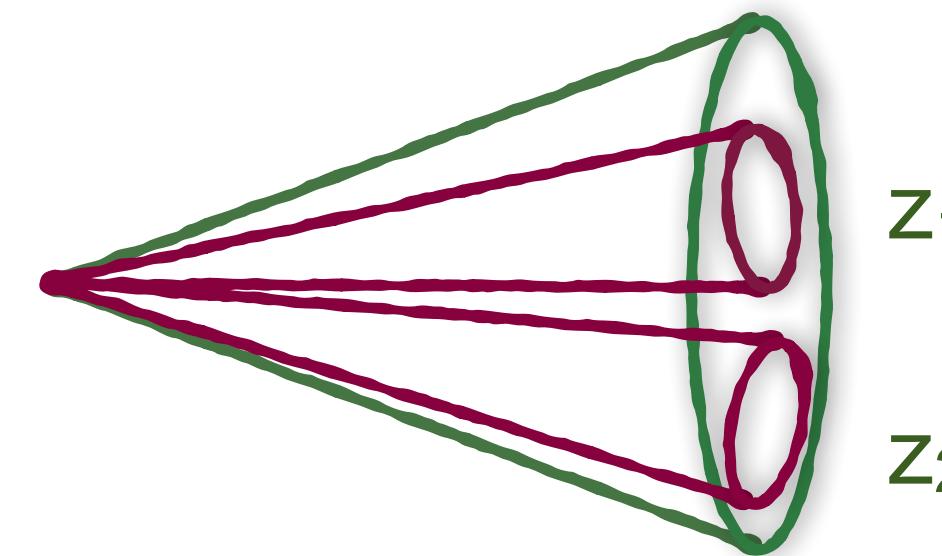
Fully reclustered anti- k_t subjets



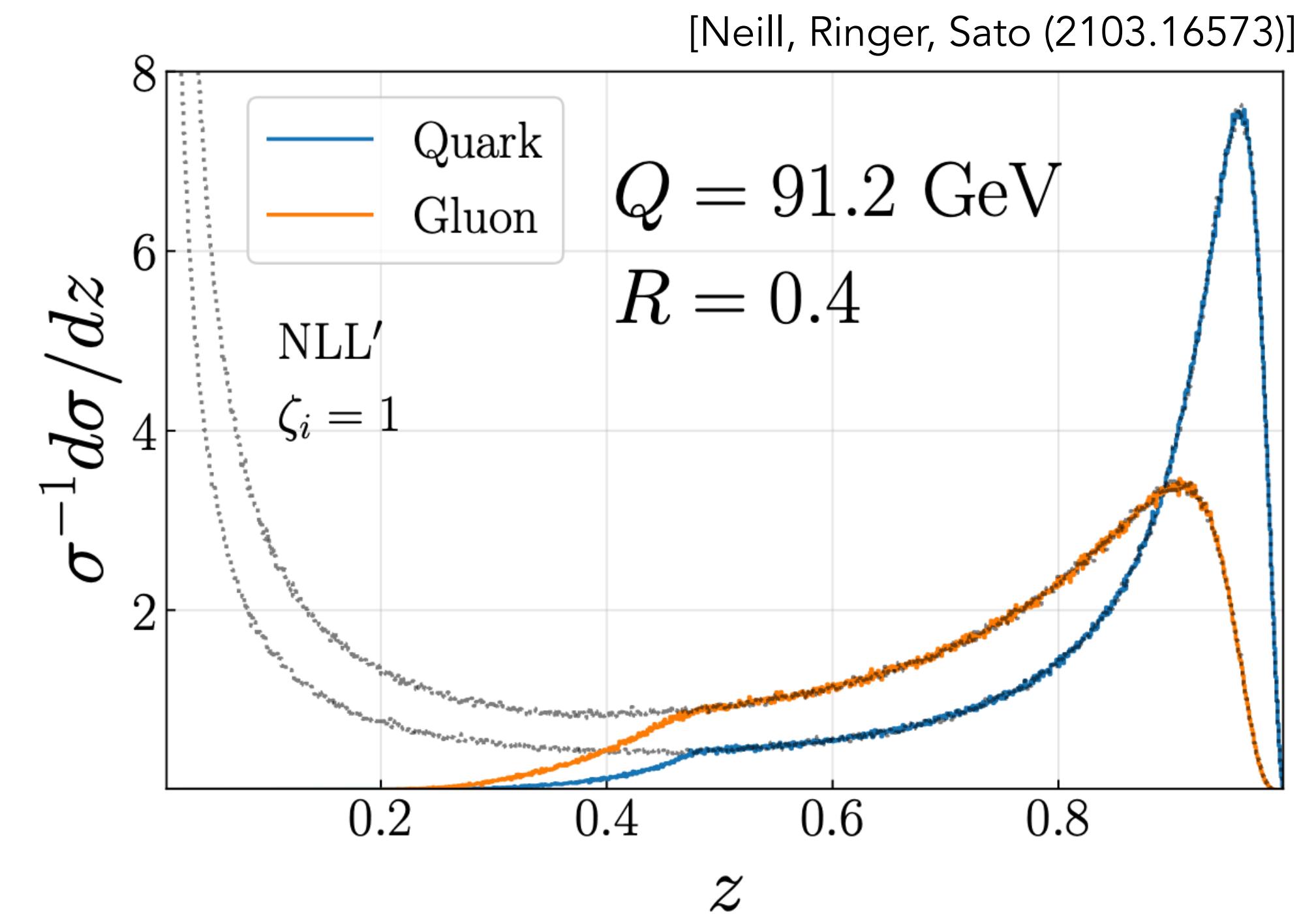
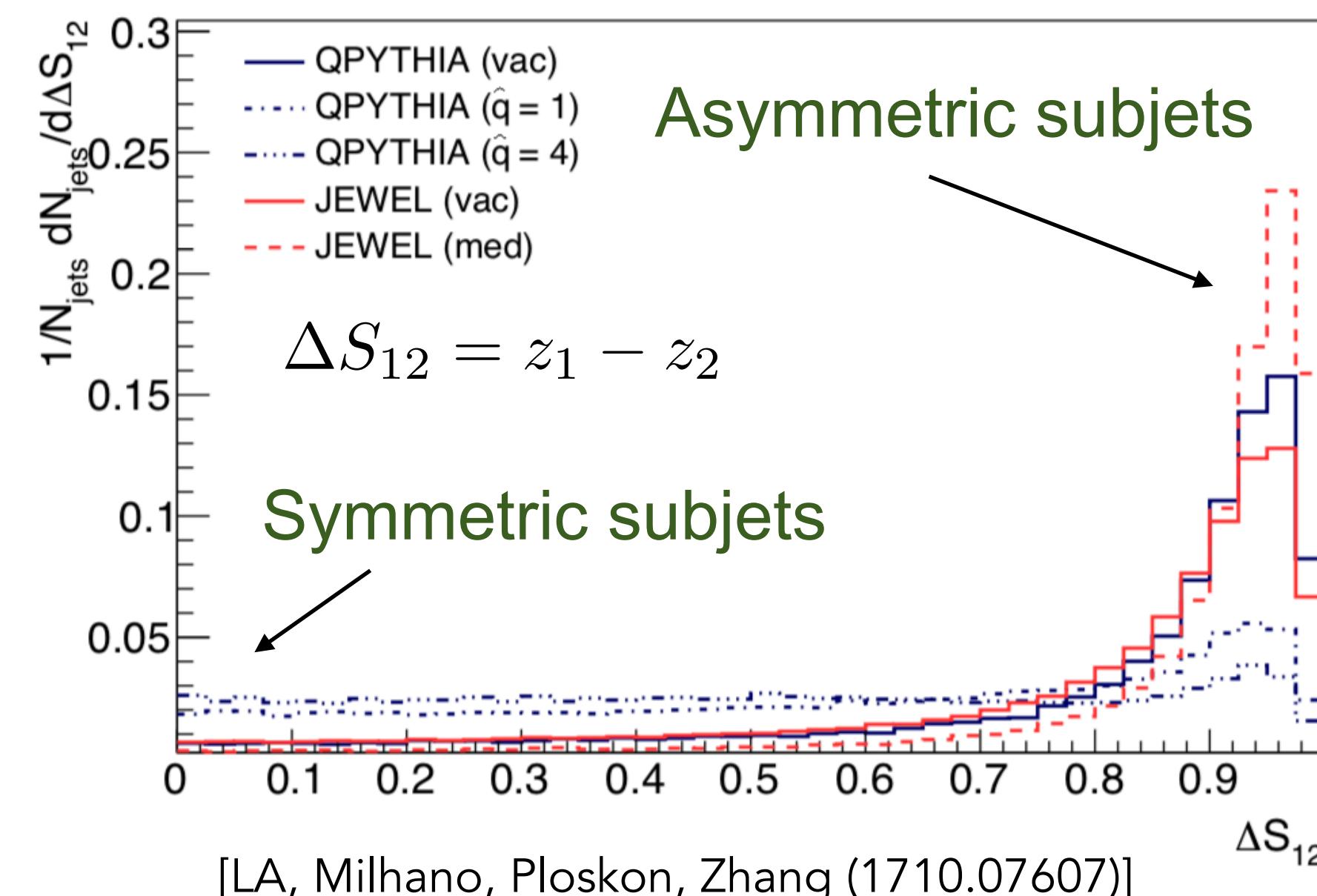
From jets to jet substructure



- How can we access QGP-related information?
 - Jets in PbPb \neq Jets in pp + Background
 - Background-resilient to distinguish quenching models
 - Leading jet: quantifies quark vs gluon in-medium energy loss



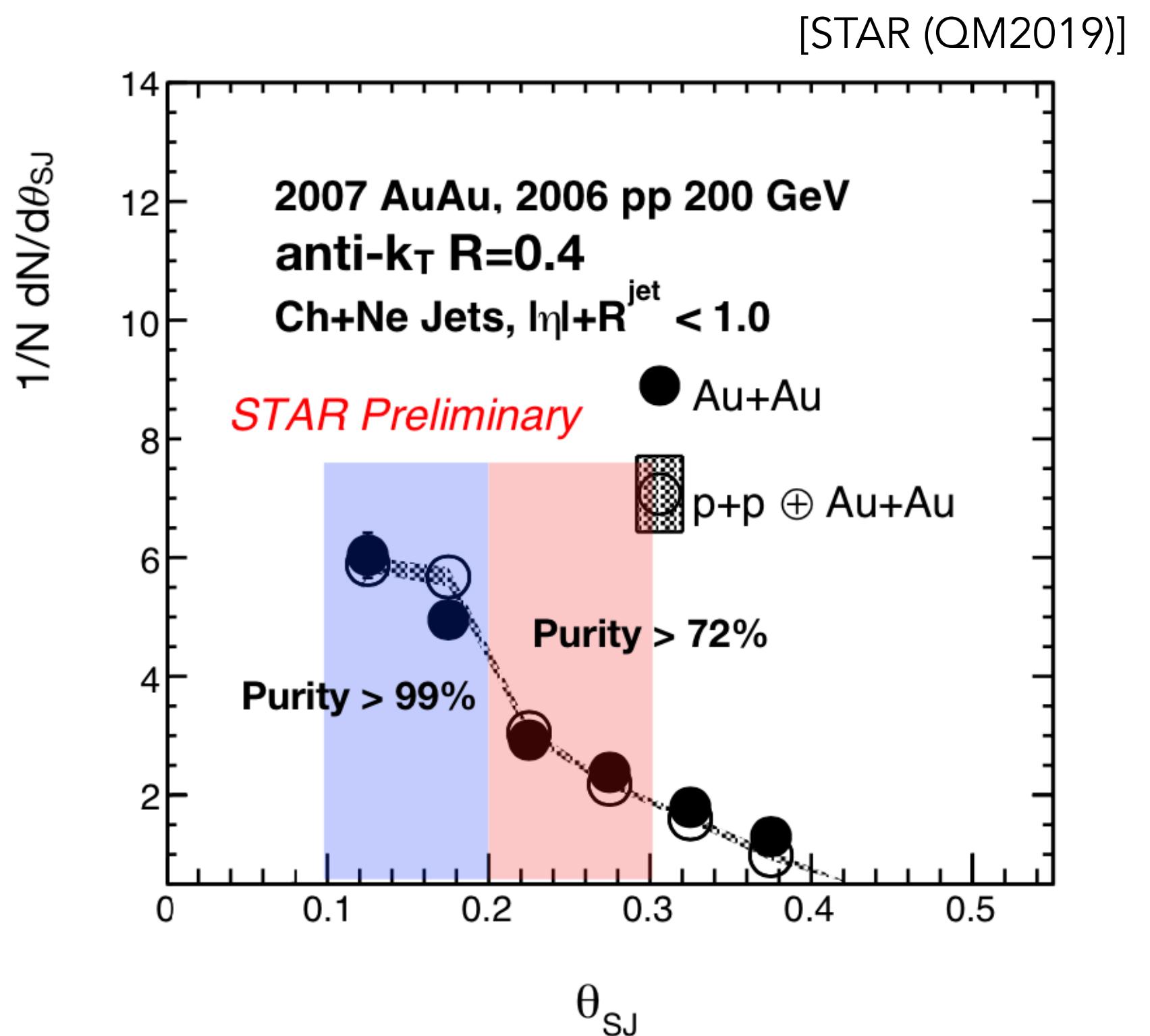
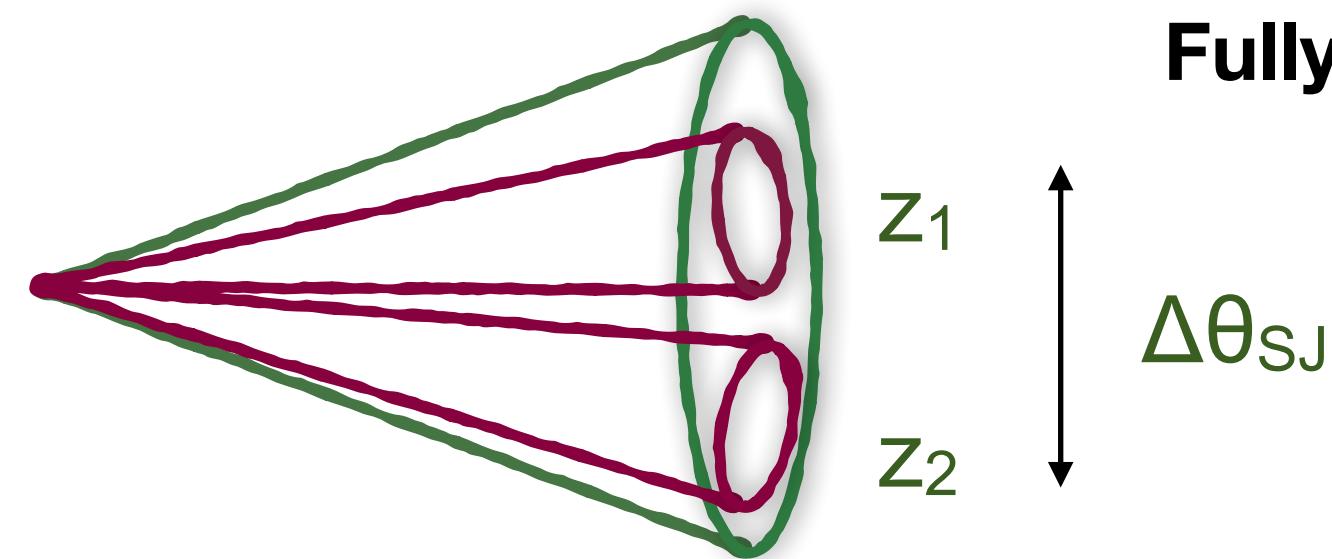
Fully reclustered anti- k_t subjets



From jets to jet substructure



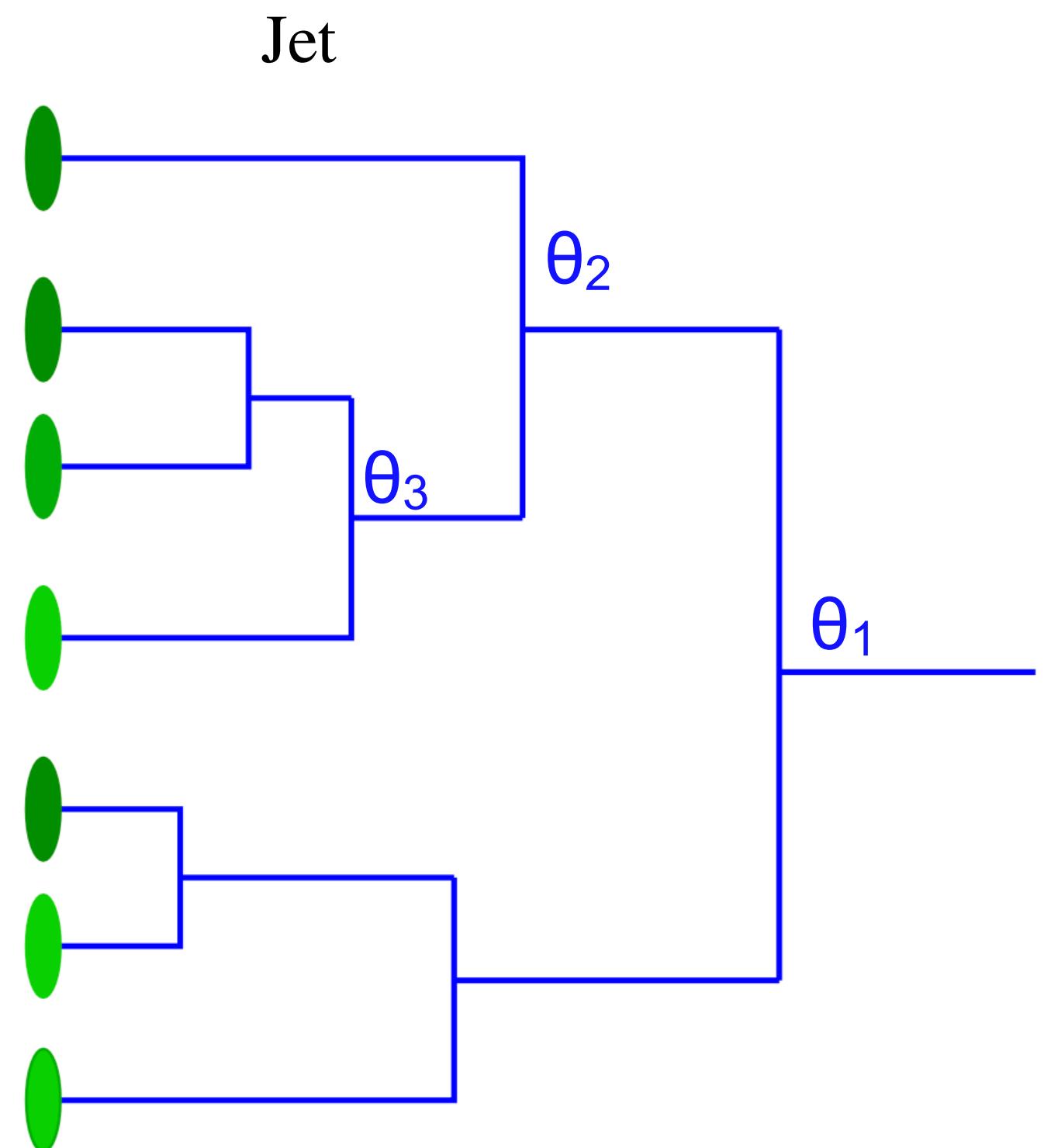
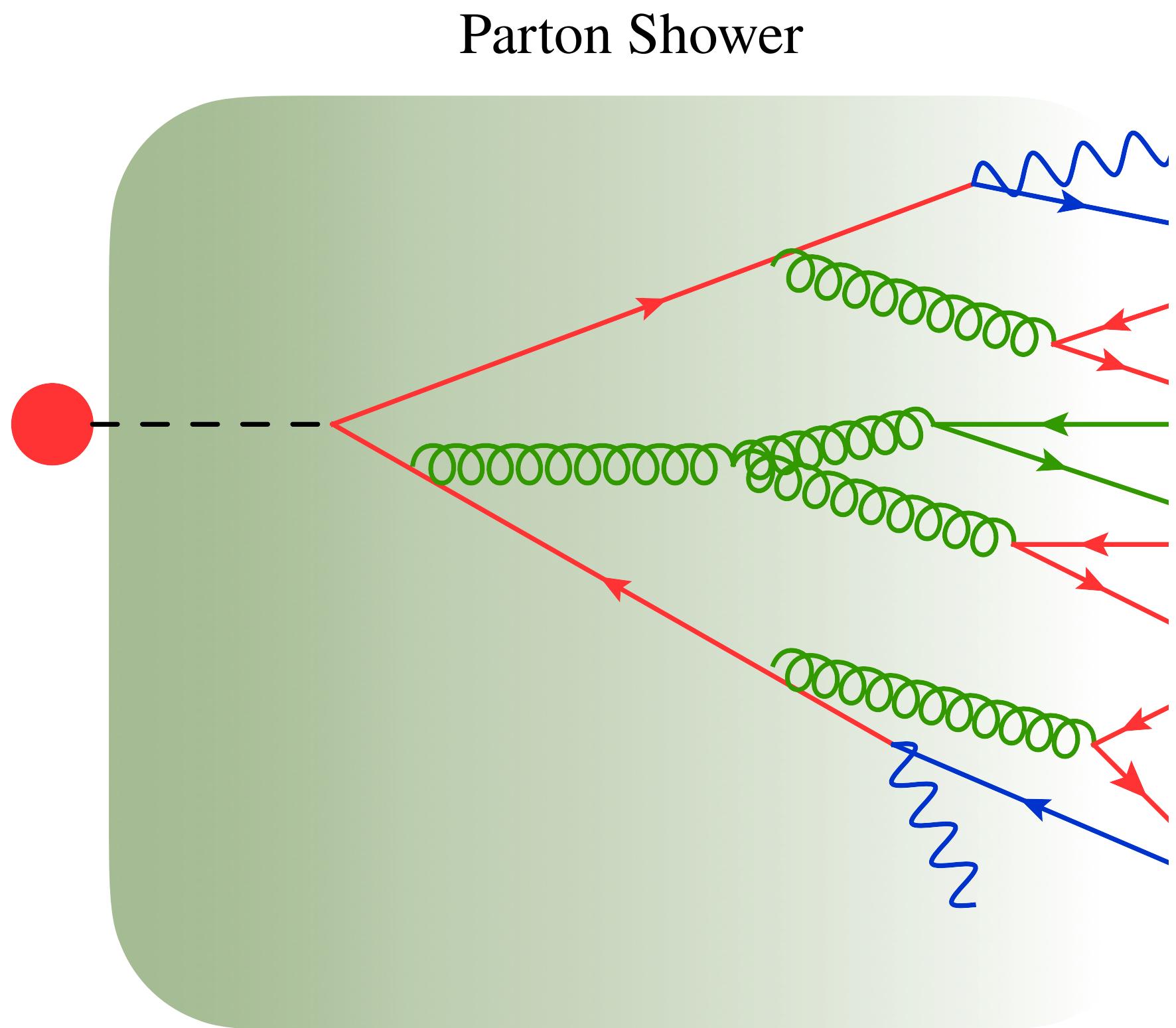
- How can we access QGP-related information?
 - Jets in PbPb \neq Jets in pp + Background
 - Background-resilient to distinguish quenching models
 - Leading jet: quantifies quark vs gluon in-medium energy loss
 - Allows to create samples that are the same in pp and in PbPb





From jets to jet substructure

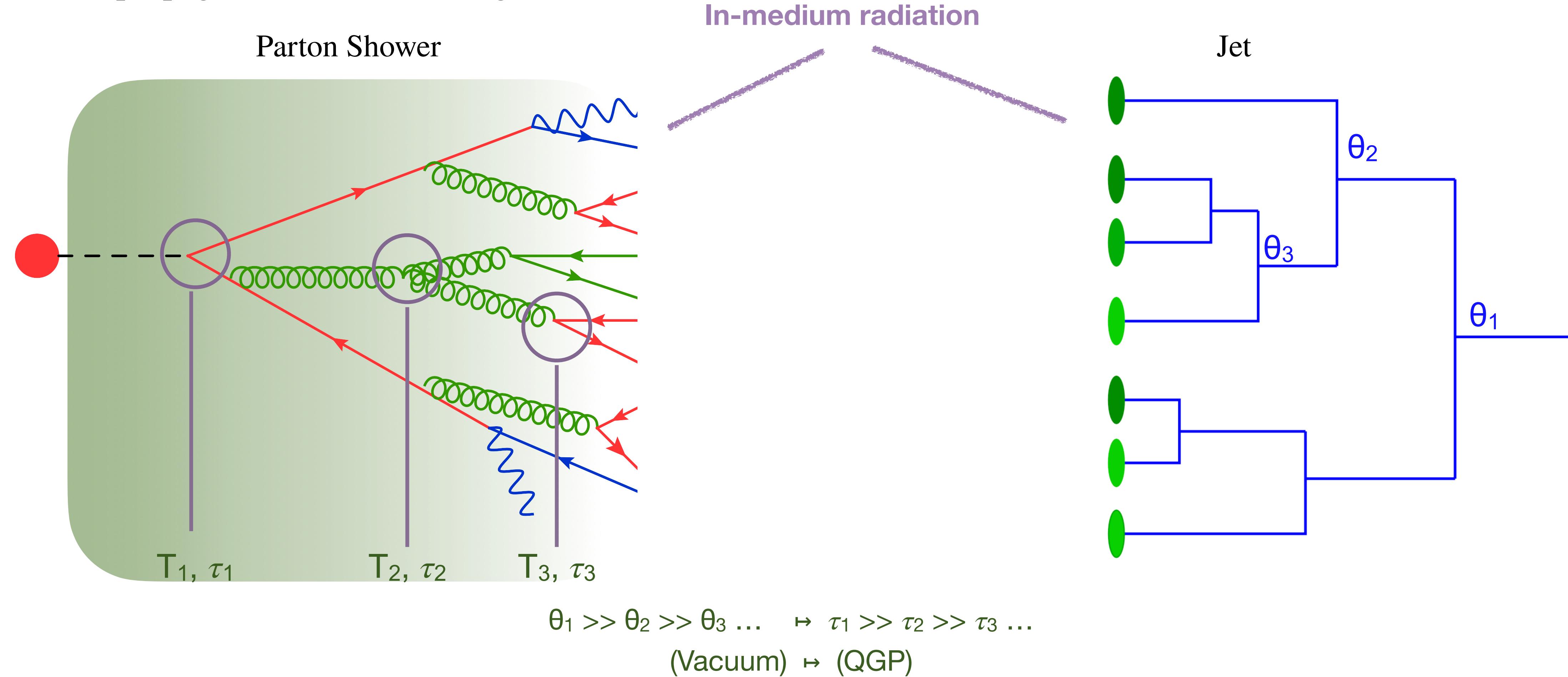
- Jets propagate on a fast evolving medium:





From jets to jet substructure

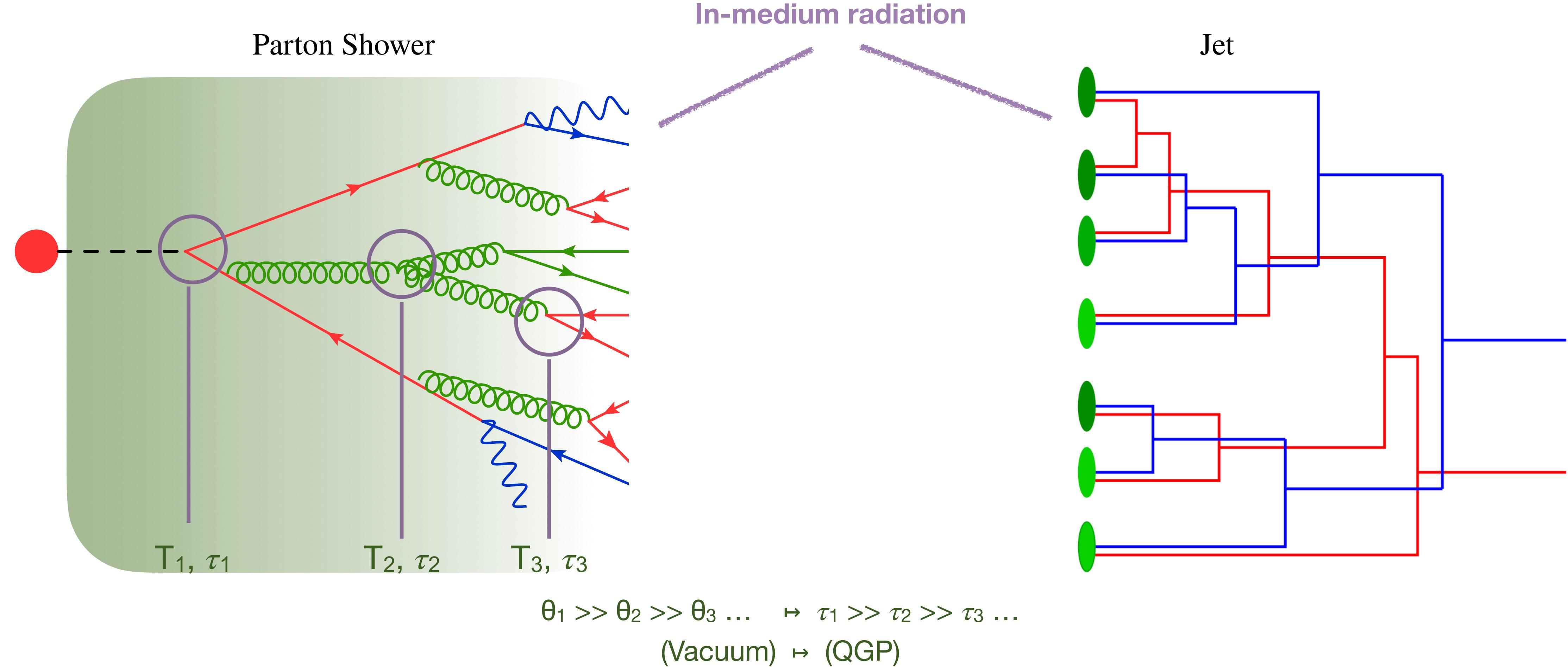
- Jets propagate on a fast evolving medium:





From jets to jet substructure

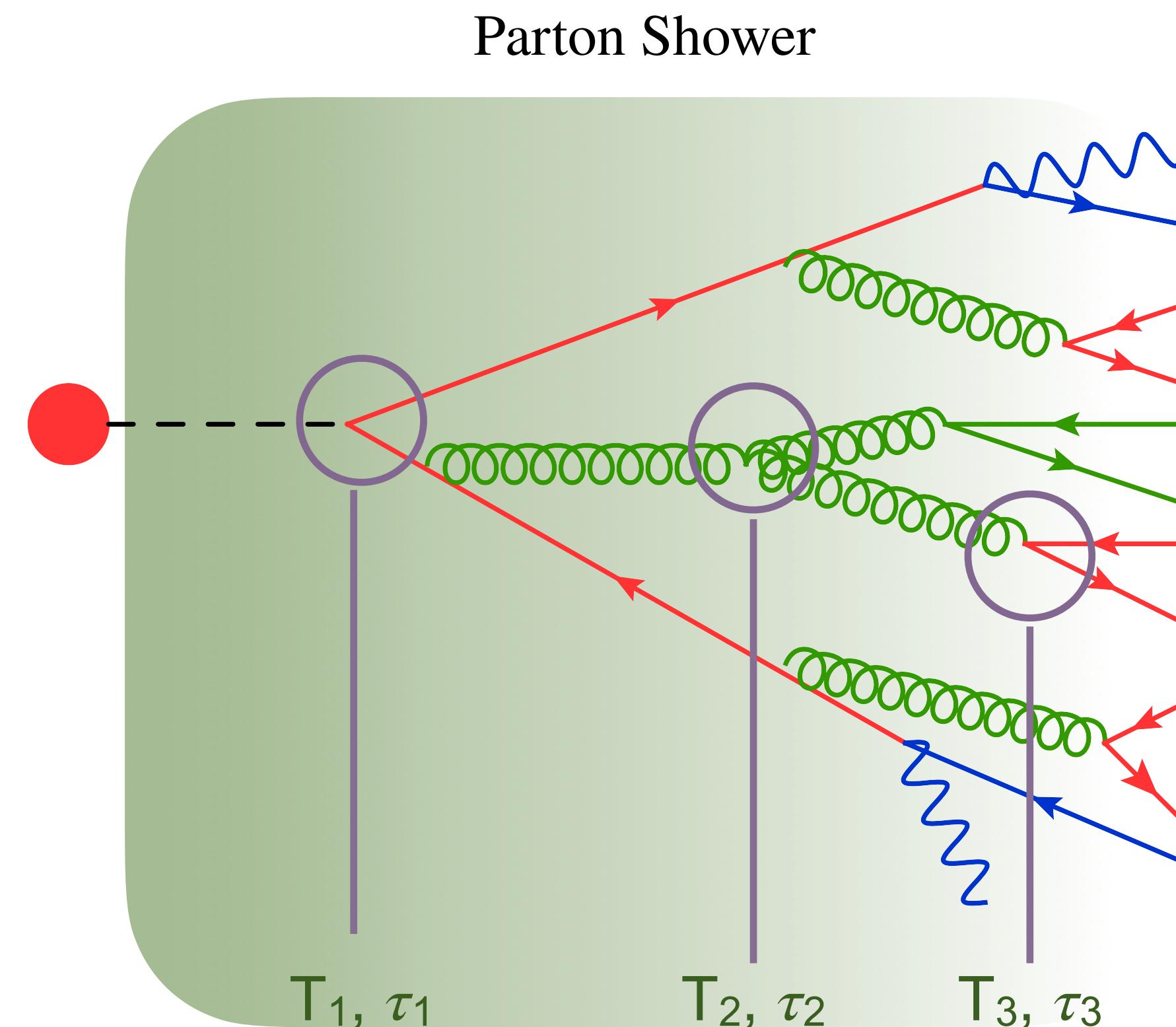
- Jets propagate on a fast evolving medium:





From jets to jet substructure

- Jets propagate on a fast evolving medium:



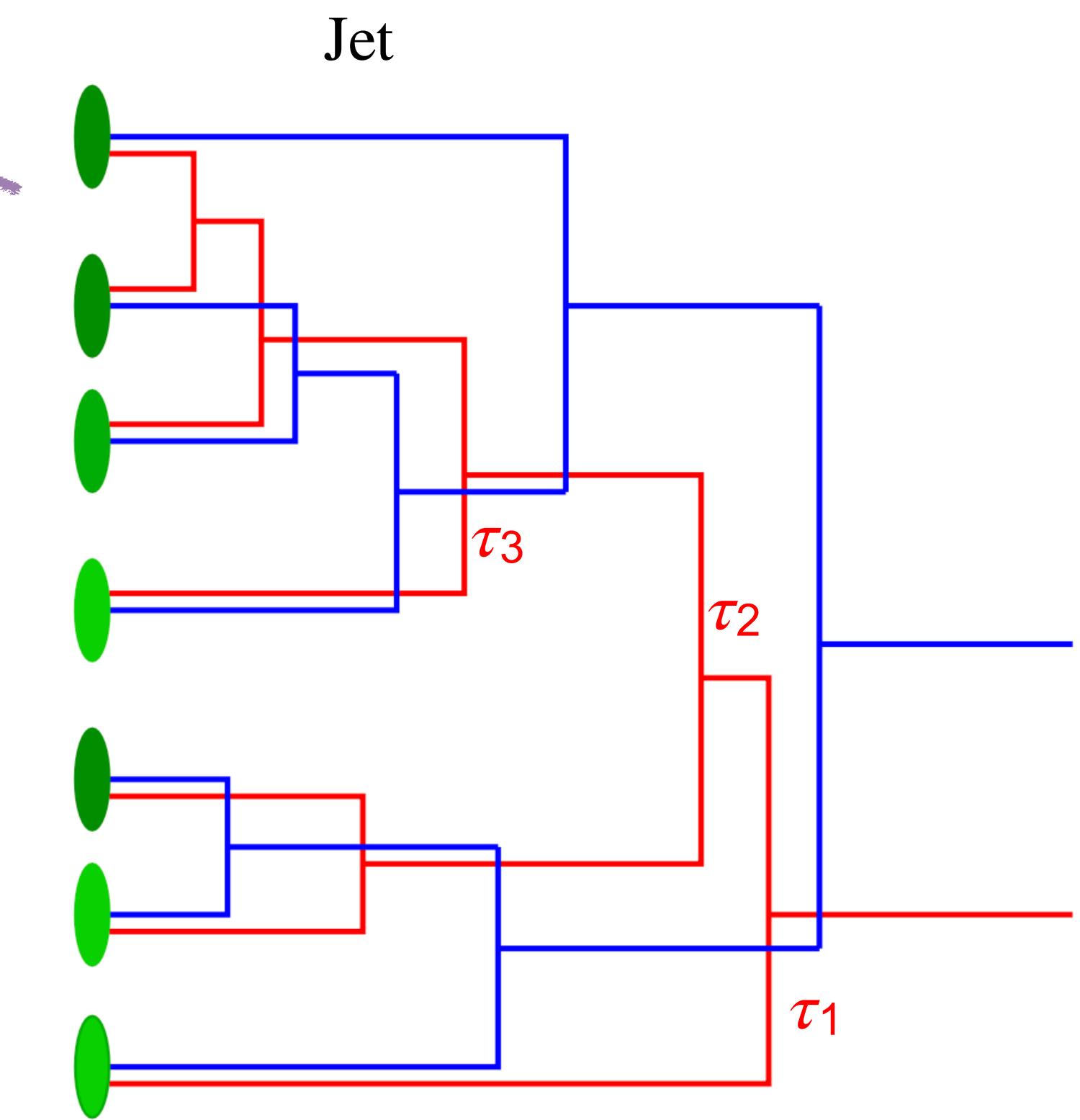
In-medium radiation

$$d_{ij} = \min(p_{t,i}^{2p}, p_{t,j}^{2p}) \frac{\Delta R_{ij}^2}{R^2}$$

$$d_{iB} = p_{t,i}^{2p}$$

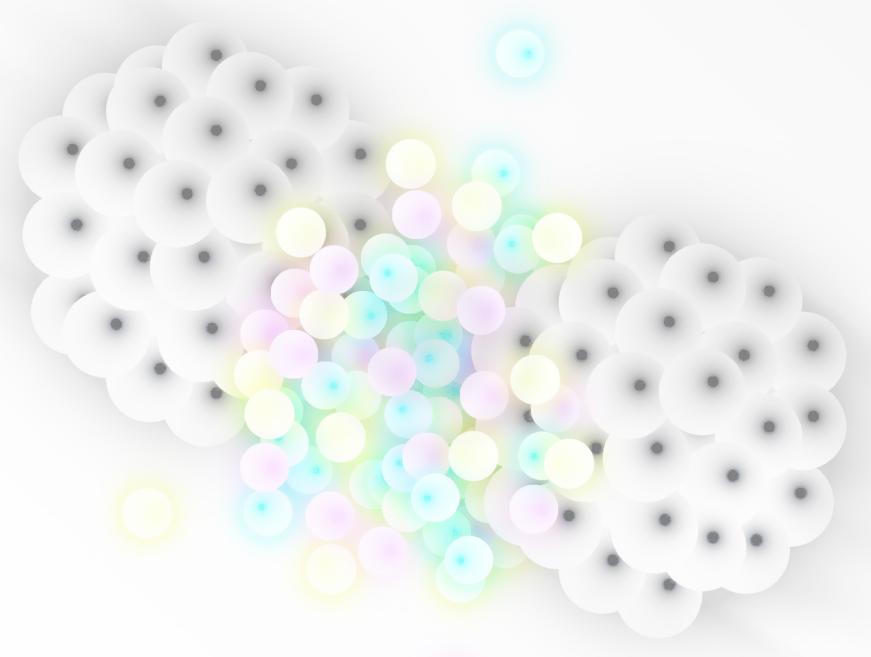
$$p = 0.5$$

$$d_{ij} \sim p_T \theta^2 \sim \frac{1}{\tau_{form}}$$

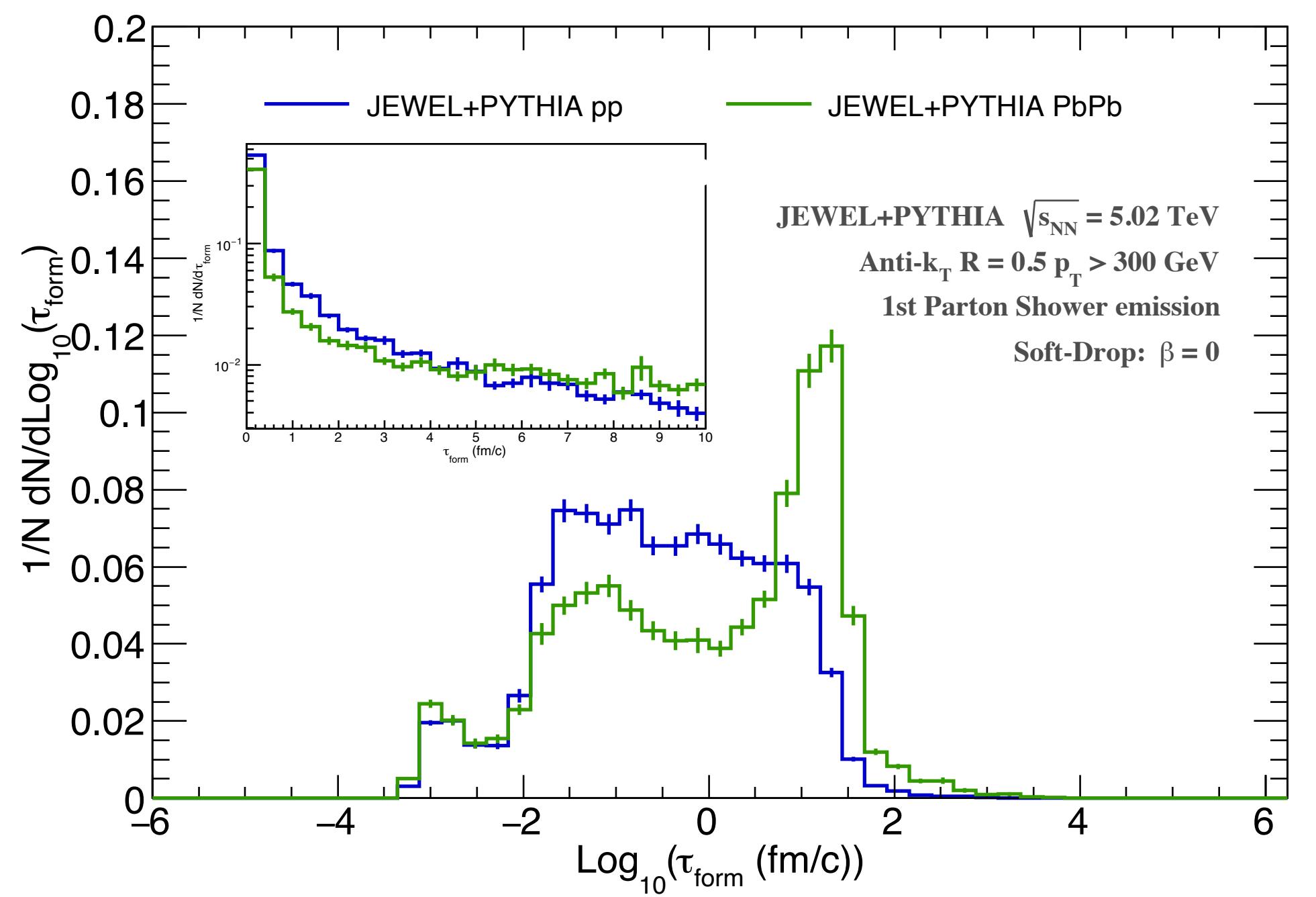


$\theta_1 \gg \theta_2 \gg \theta_3 \dots \rightarrow \tau_1 \gg \tau_2 \gg \tau_3 \dots$
 (Vacuum) \rightarrow (QGP)

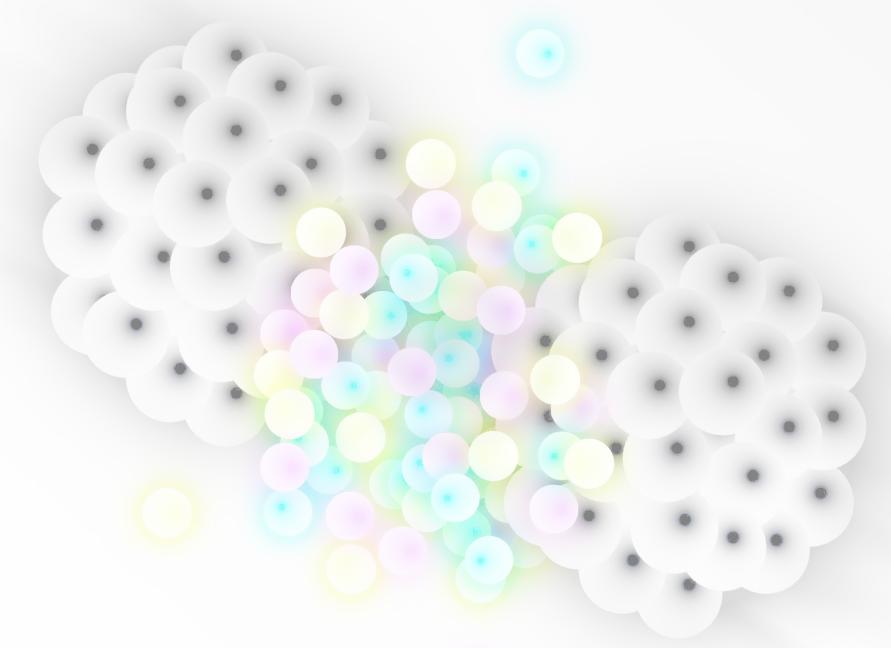
Novel jet reclustering tools



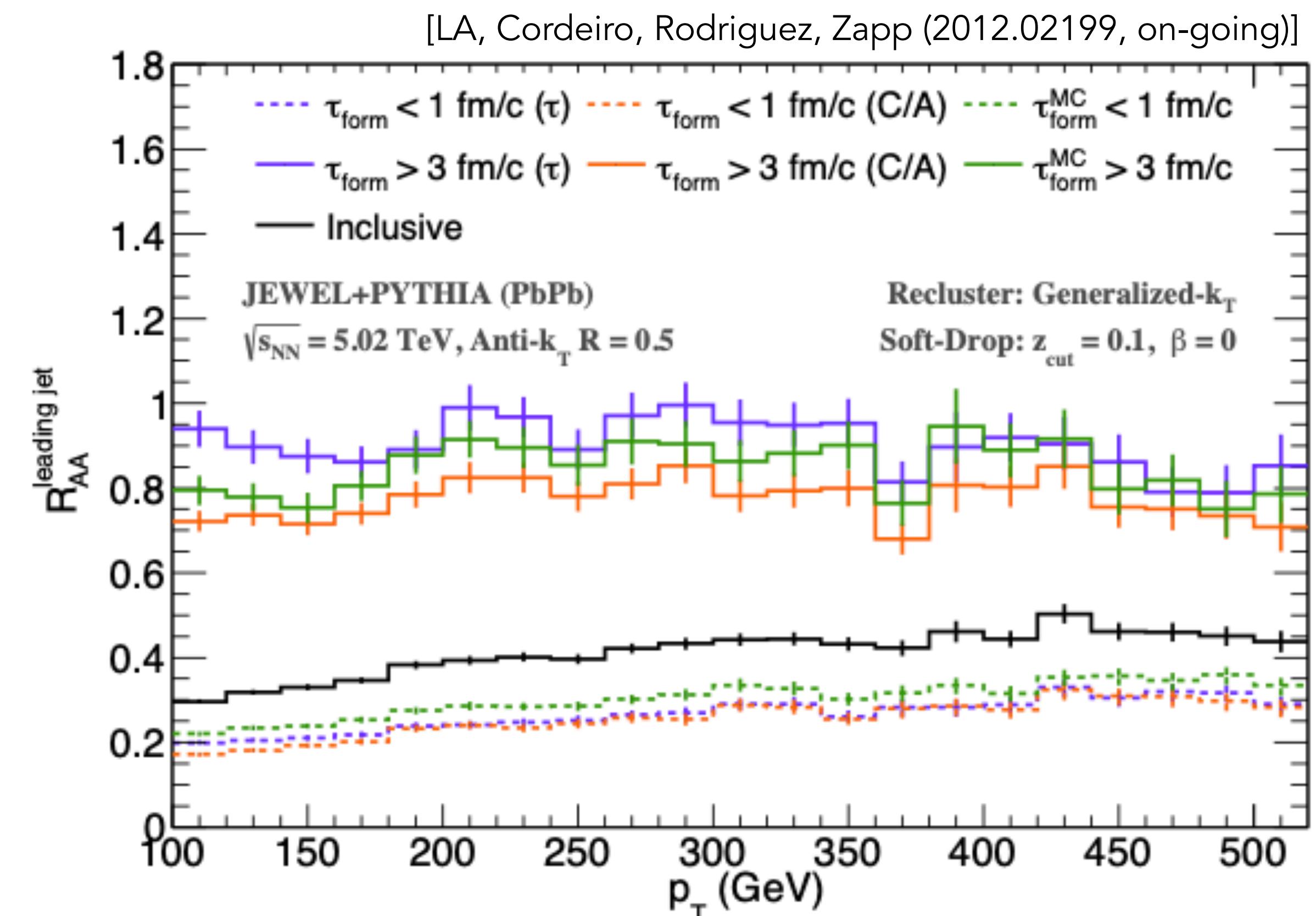
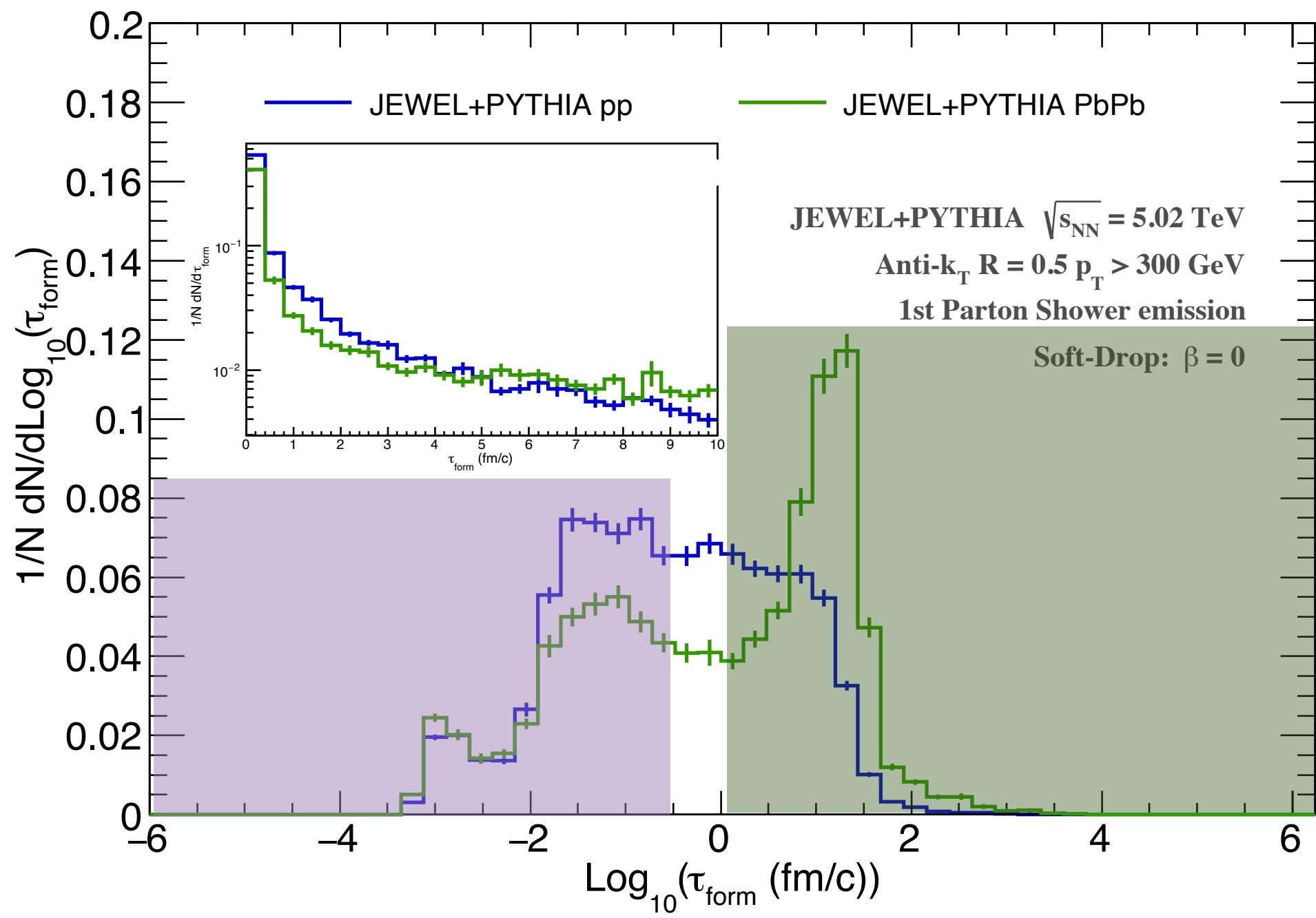
- Easily select two classes of jets:



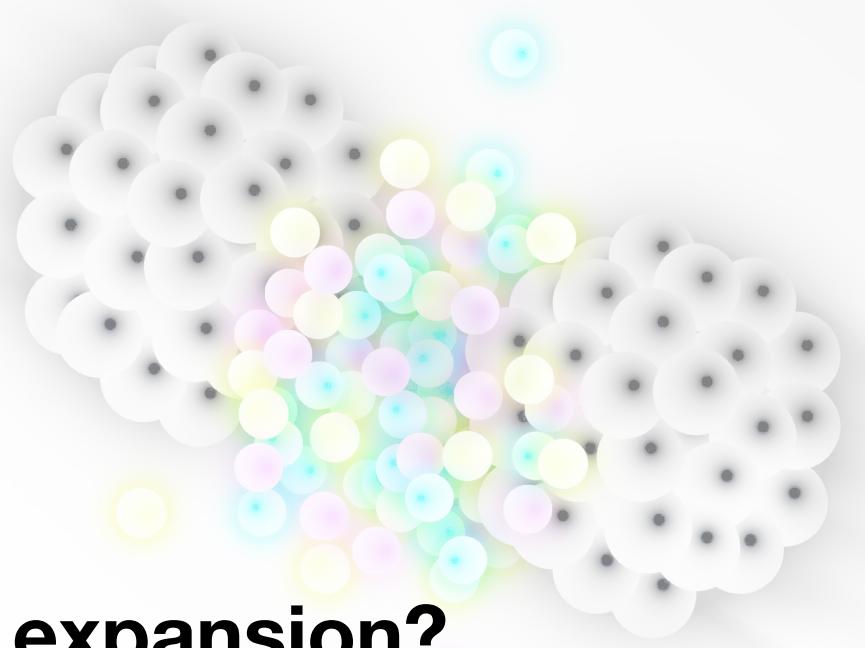
Novel jet reclustering tools



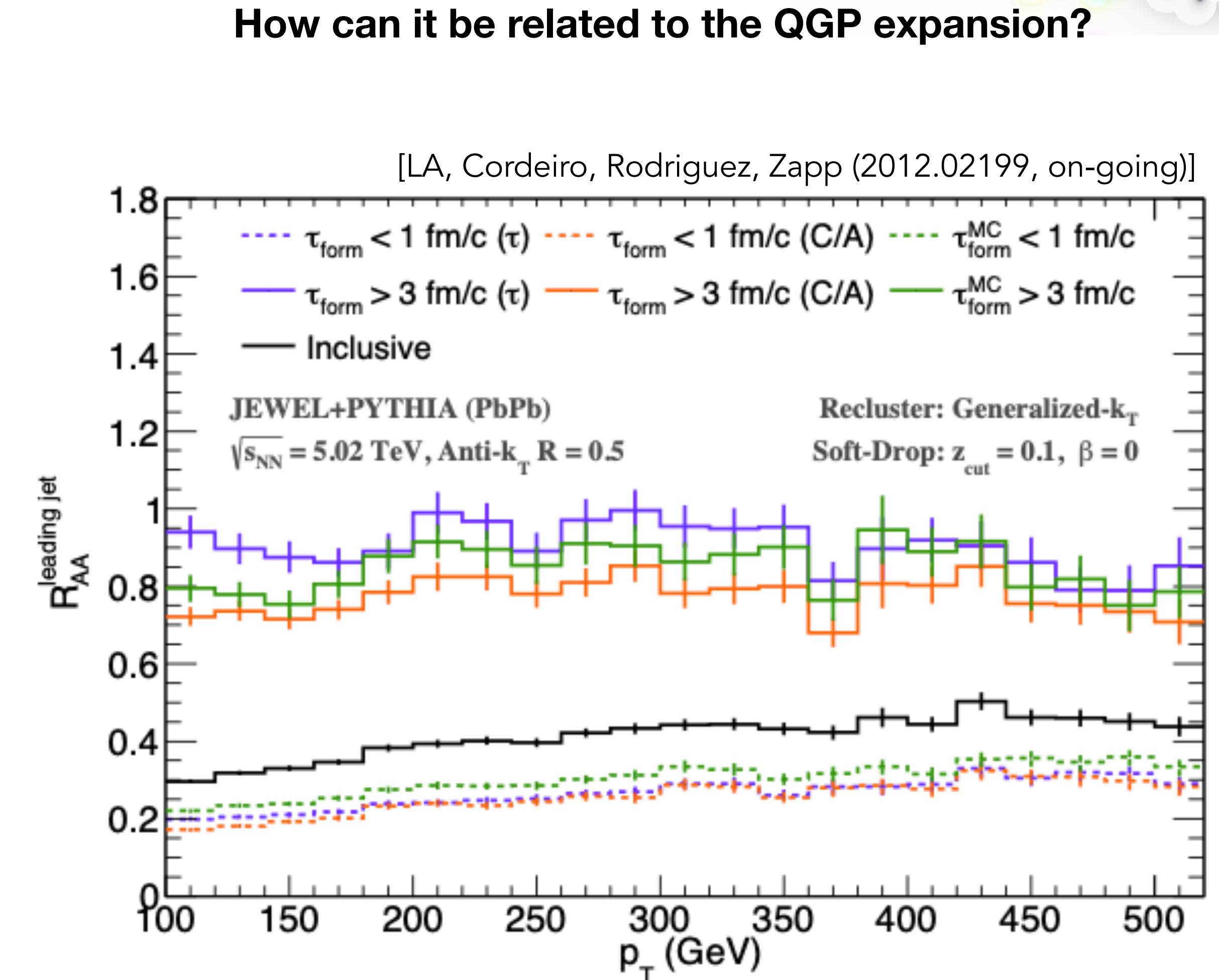
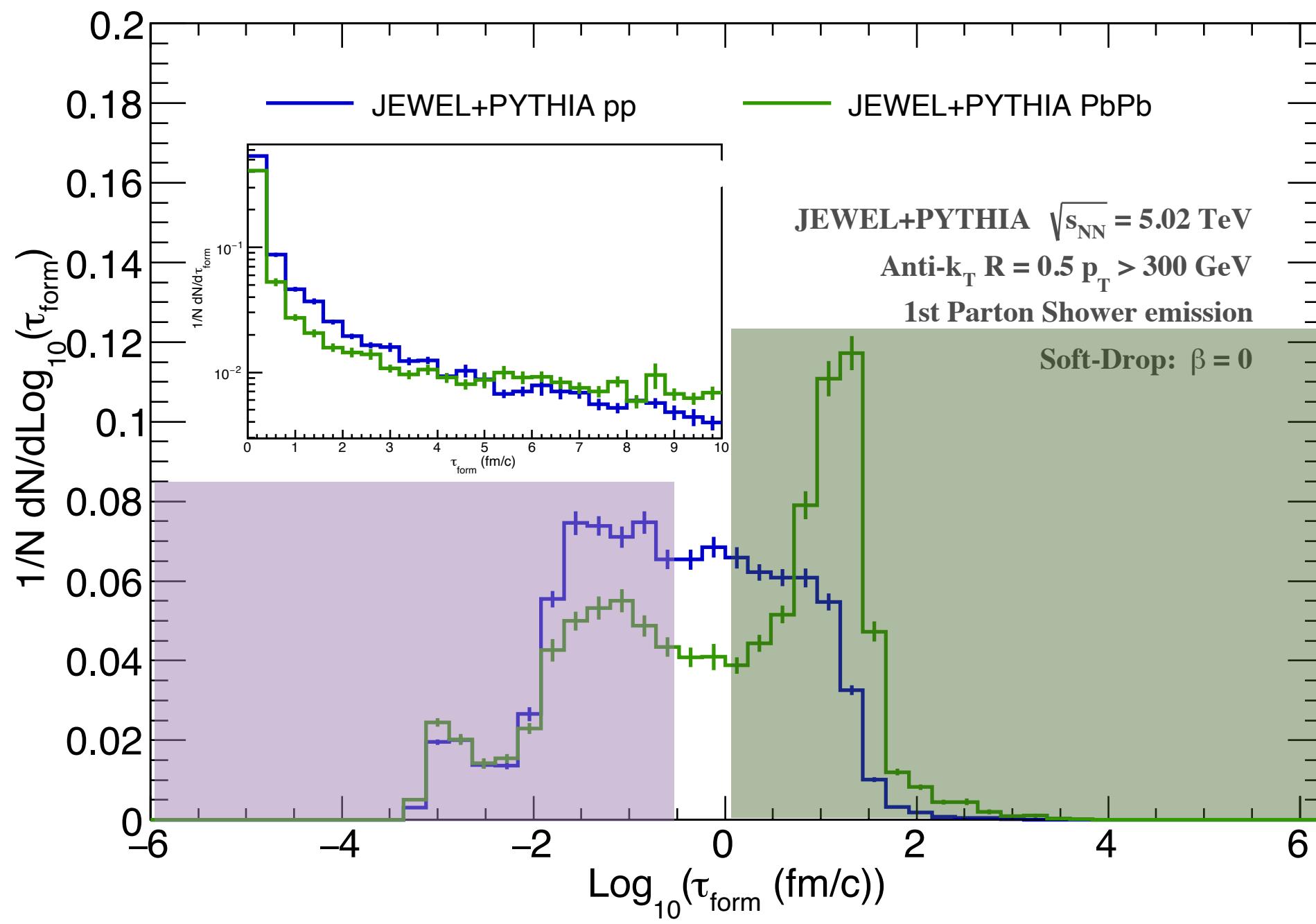
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 - “early” jets: $\tau_1 < 1 \text{ fm/c}$ (strongly modified)
 - “late” jets: $\tau_1 > 3 \text{ fm/c}$ (weakly modified)



Novel jet reclustering tools

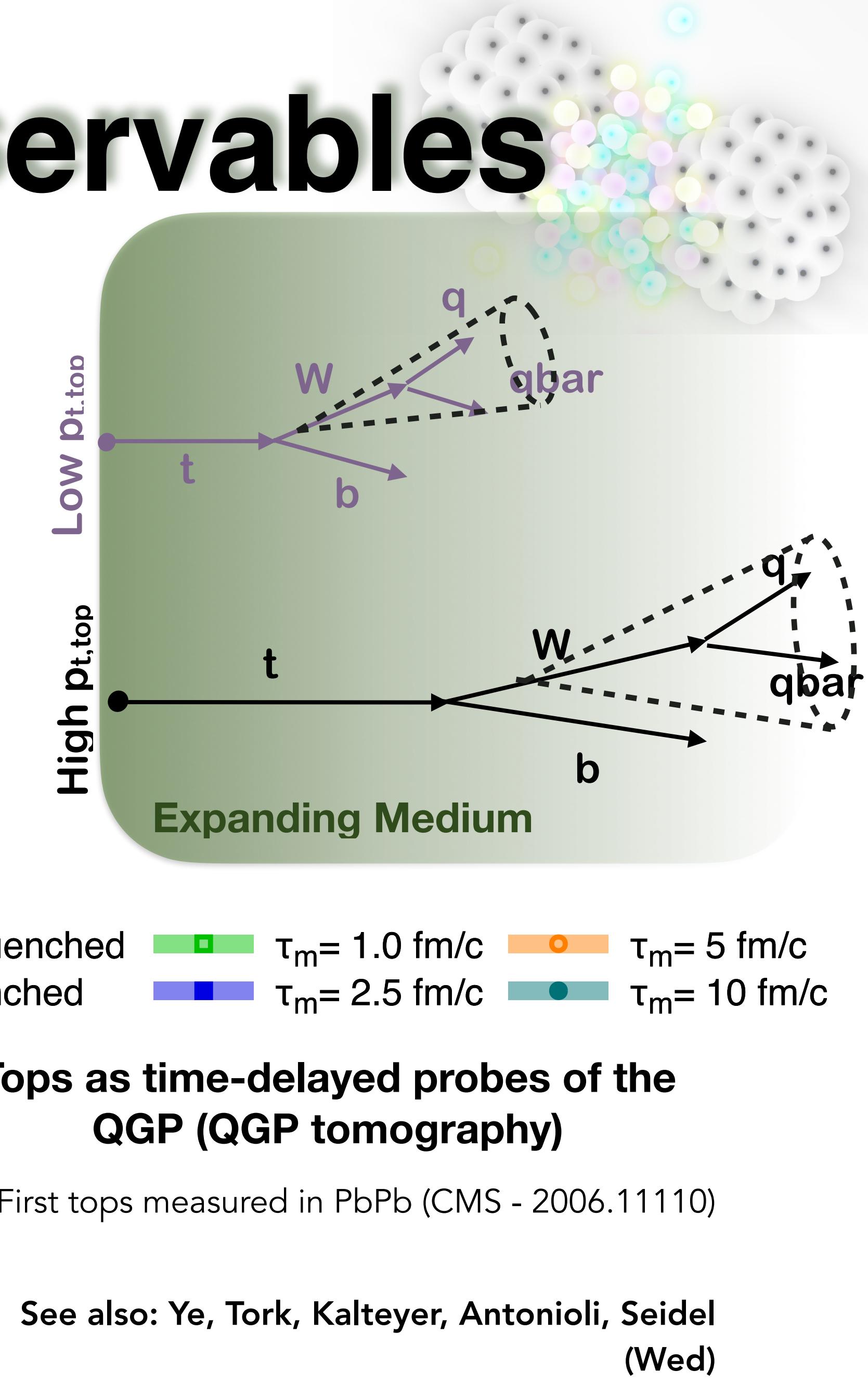
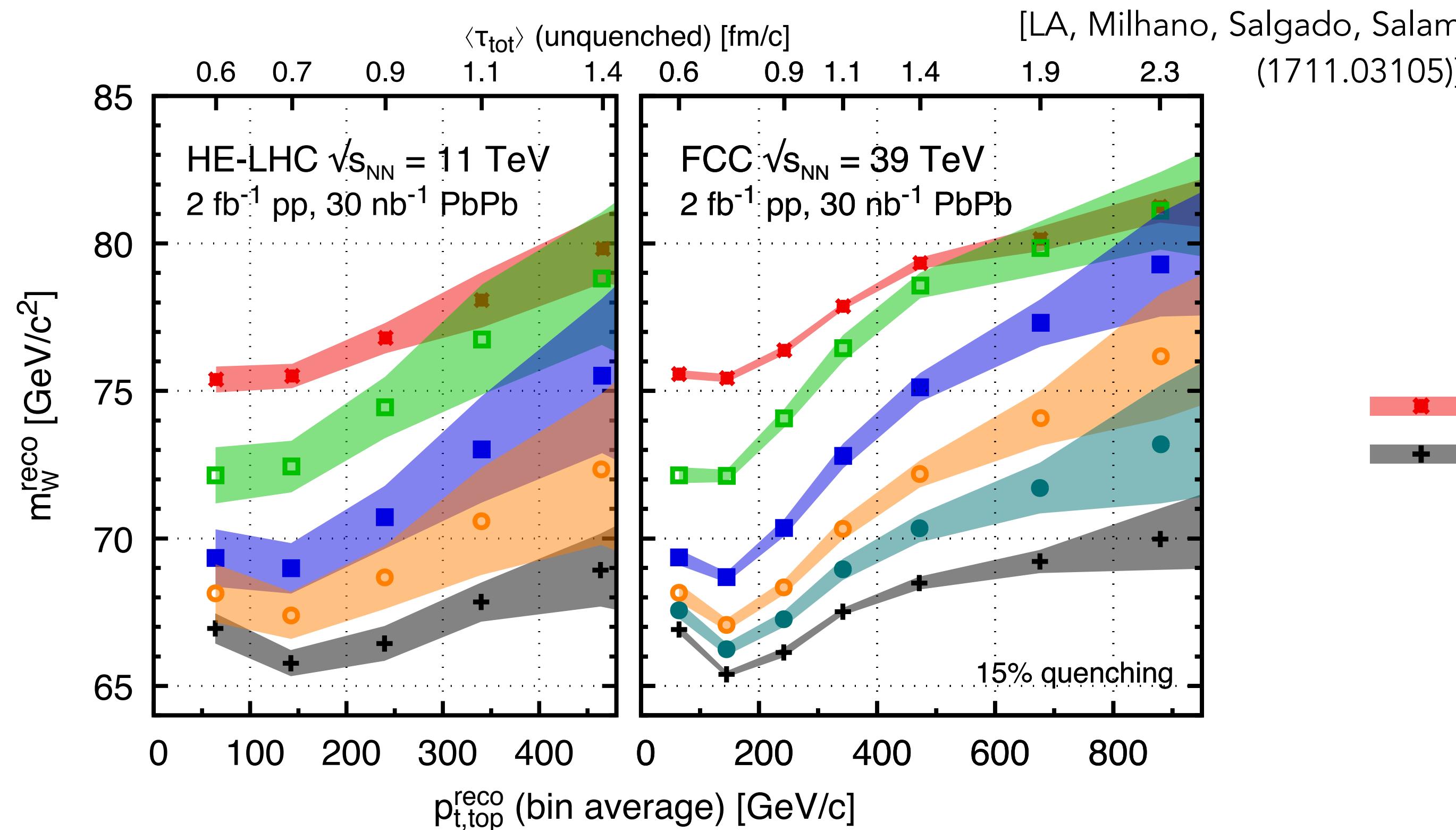


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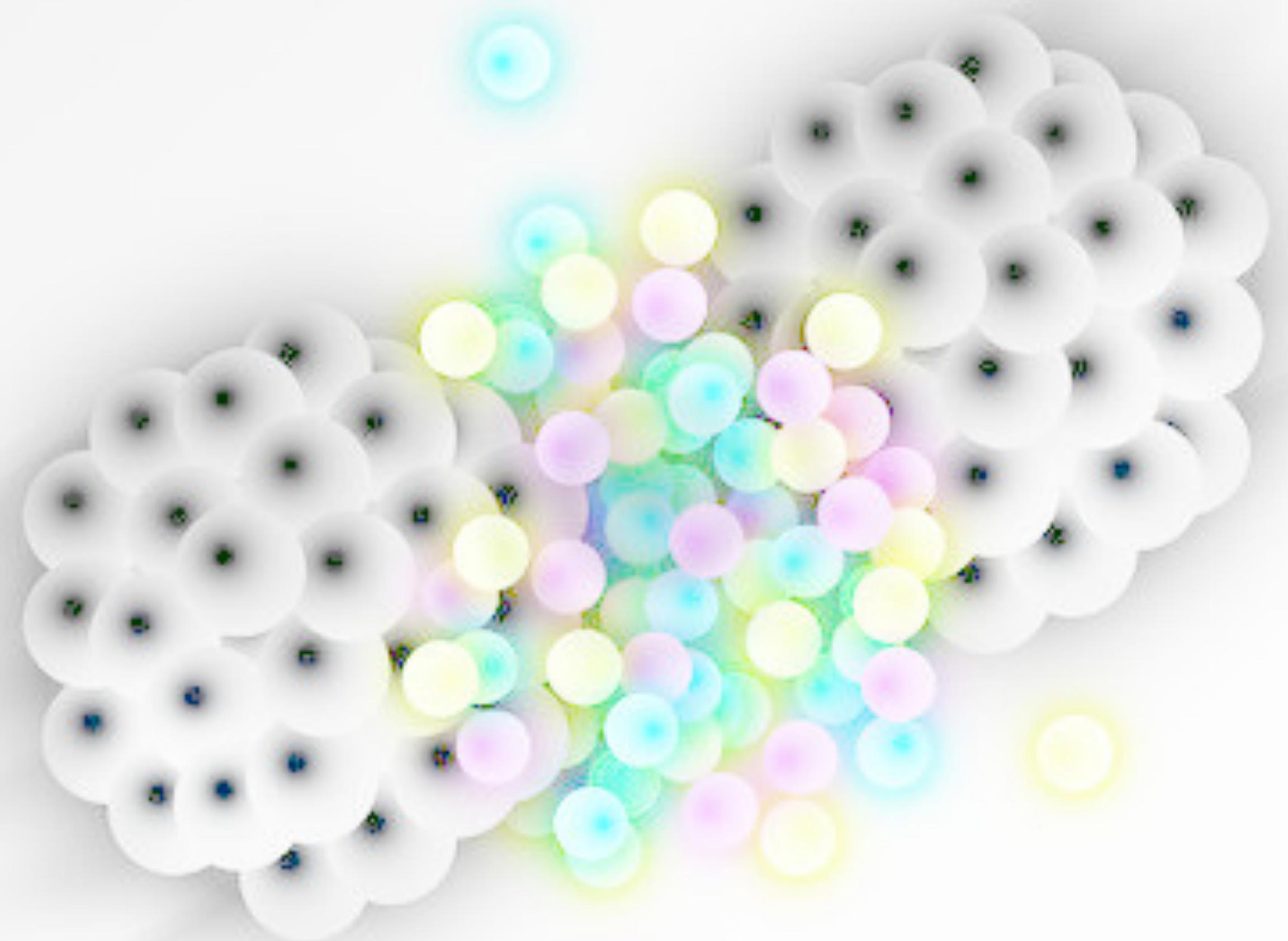


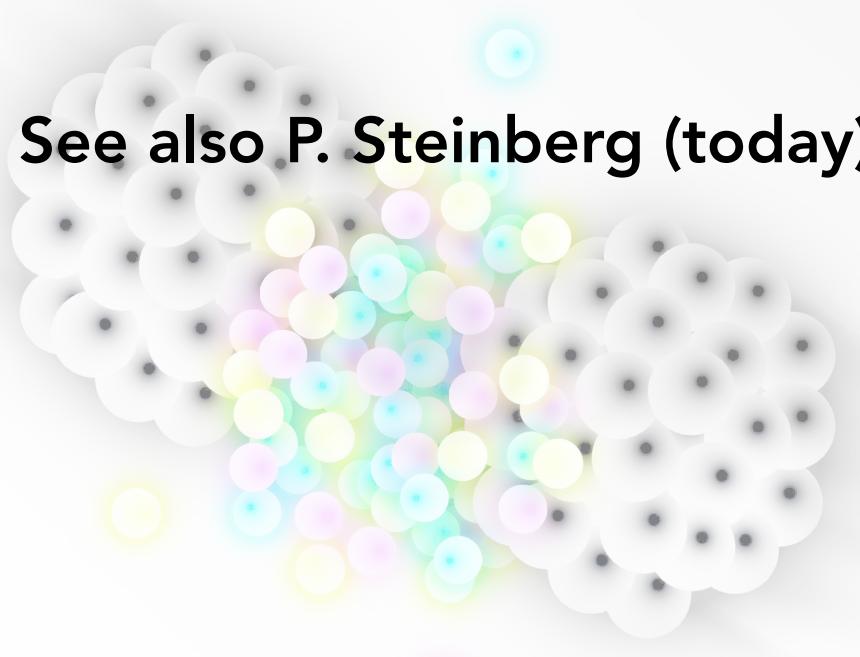
Novel heavy-flavour observables

- Hard probes: result of the full integrated medium evolution
- Time-differential measurements might be possible with tops



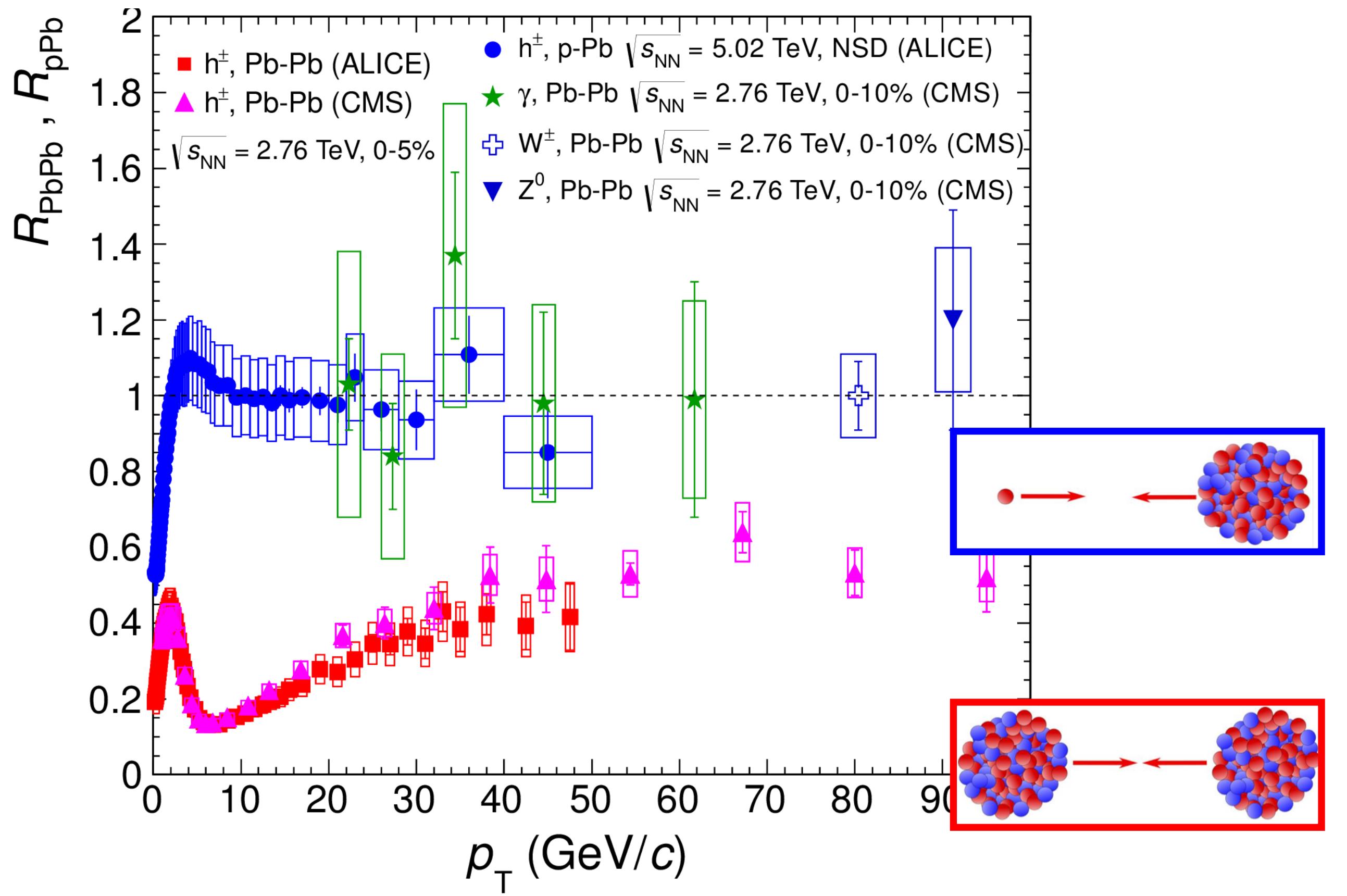
From dense to light





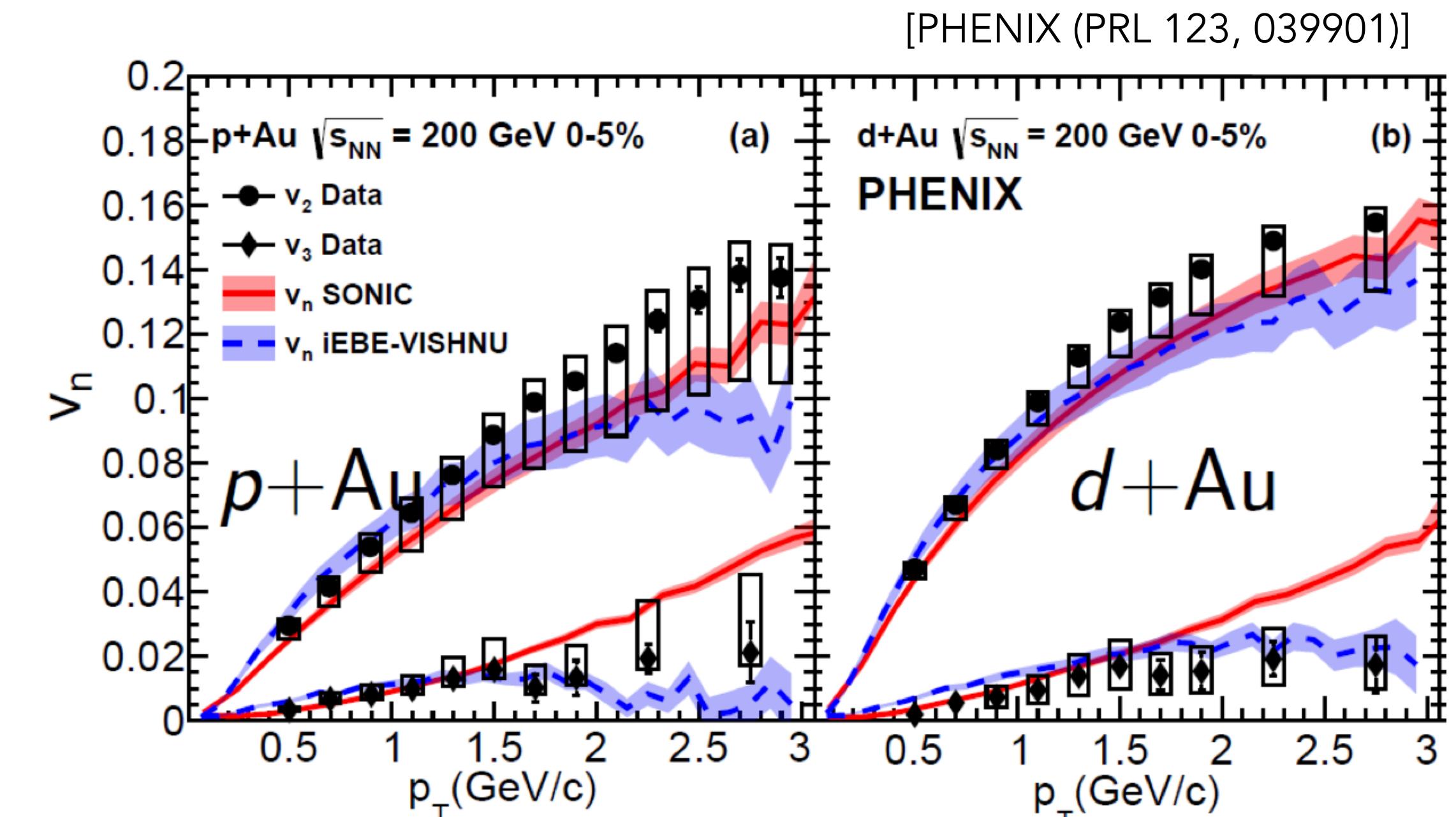
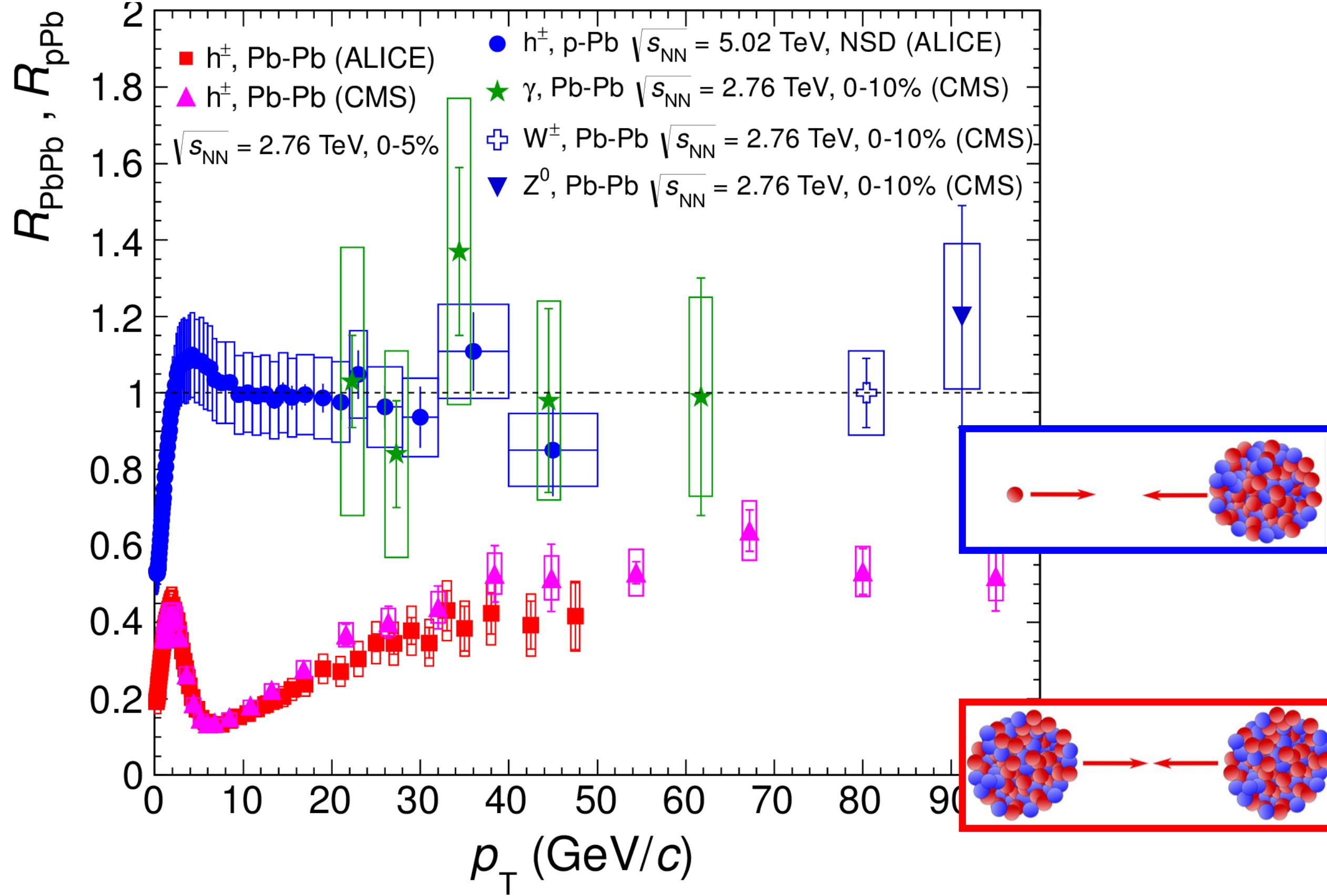
QGP onset

- No energy loss in pA...



QGP onset

- No energy loss in pA... but strong evidence in support of hydrodynamic behavior

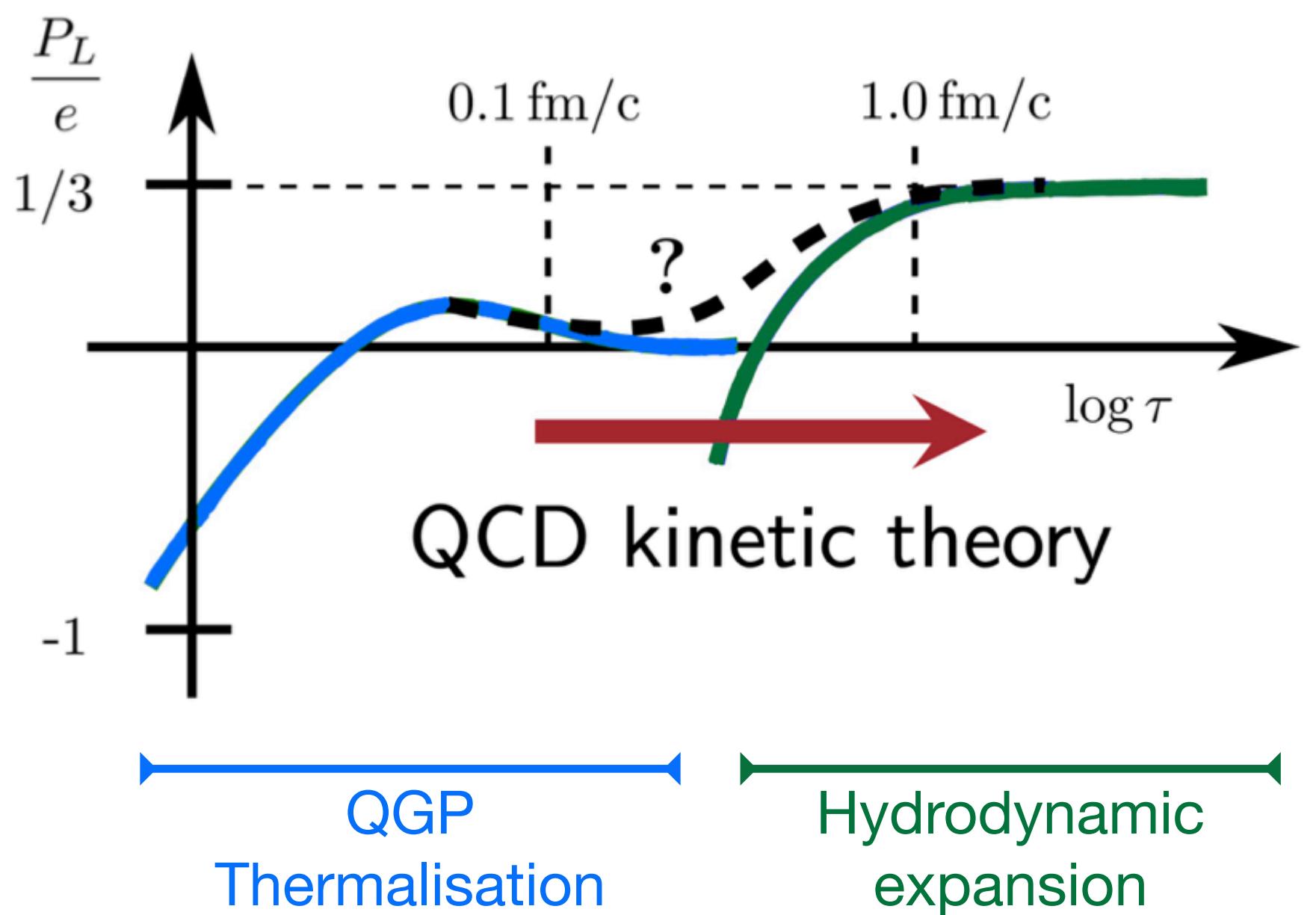


Flow coefficients well reproduced by hydro predictions, but not by initial state effects only



From dense to light systems

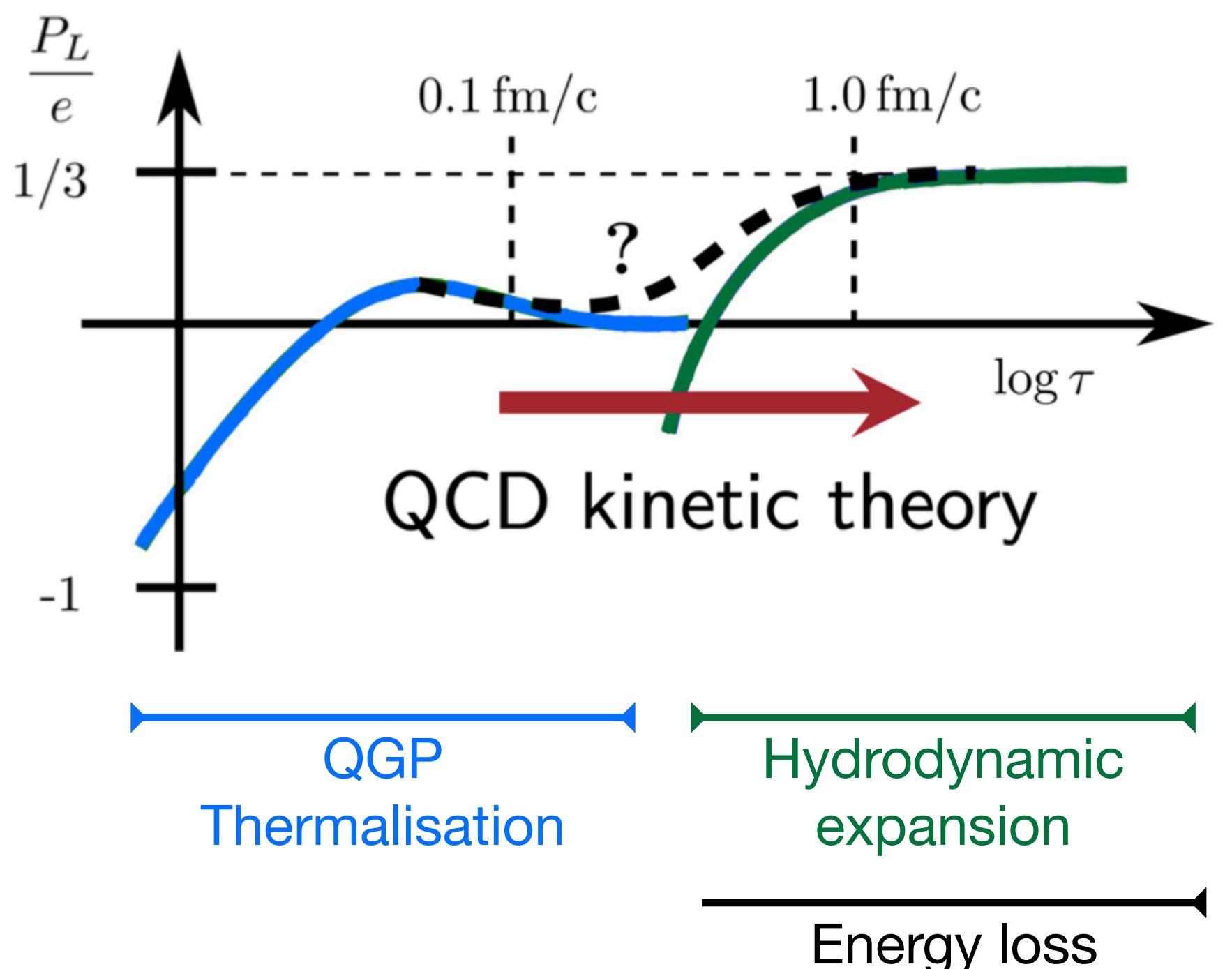
- Extrapolation from dense to light needs further understanding...





From dense to light systems

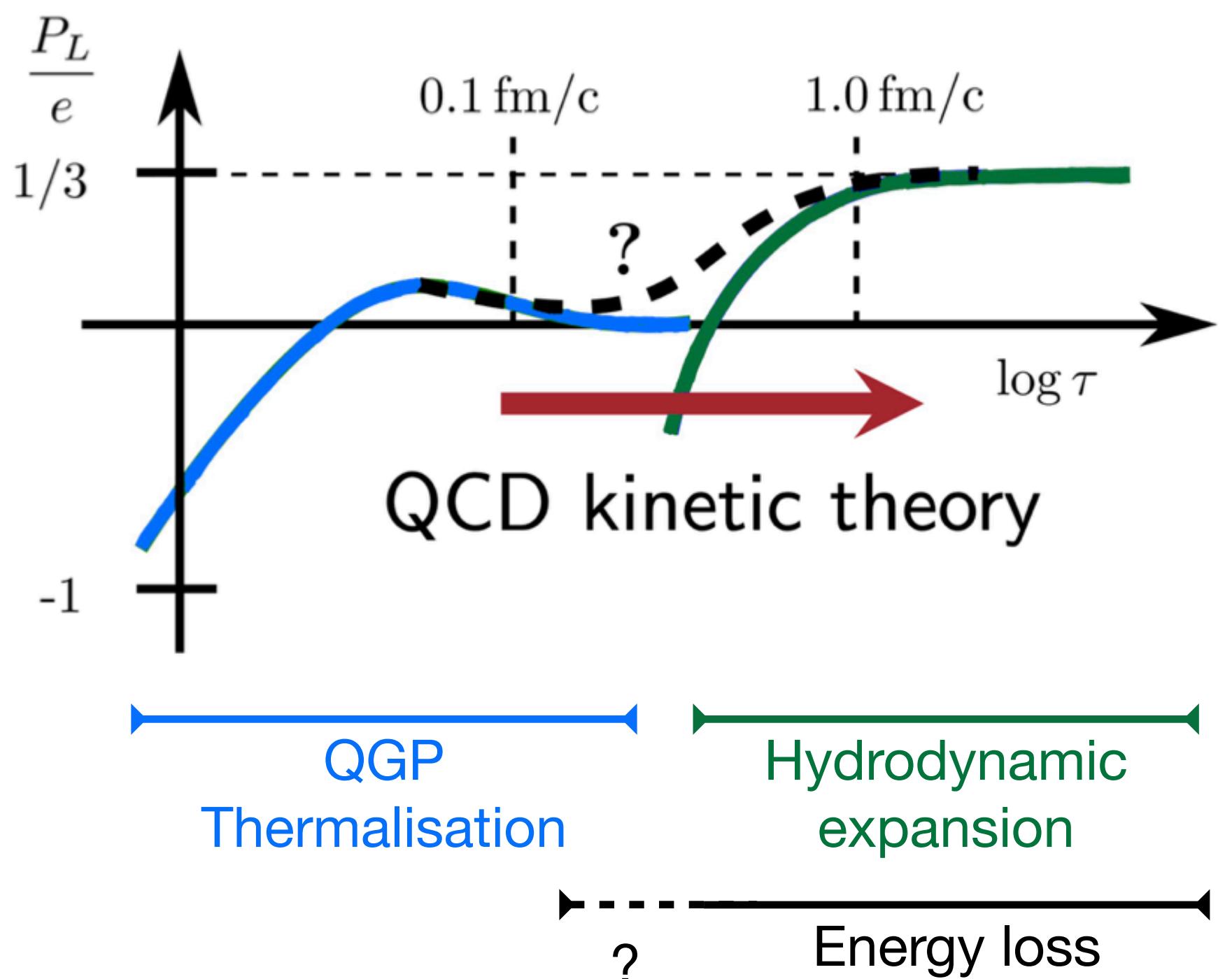
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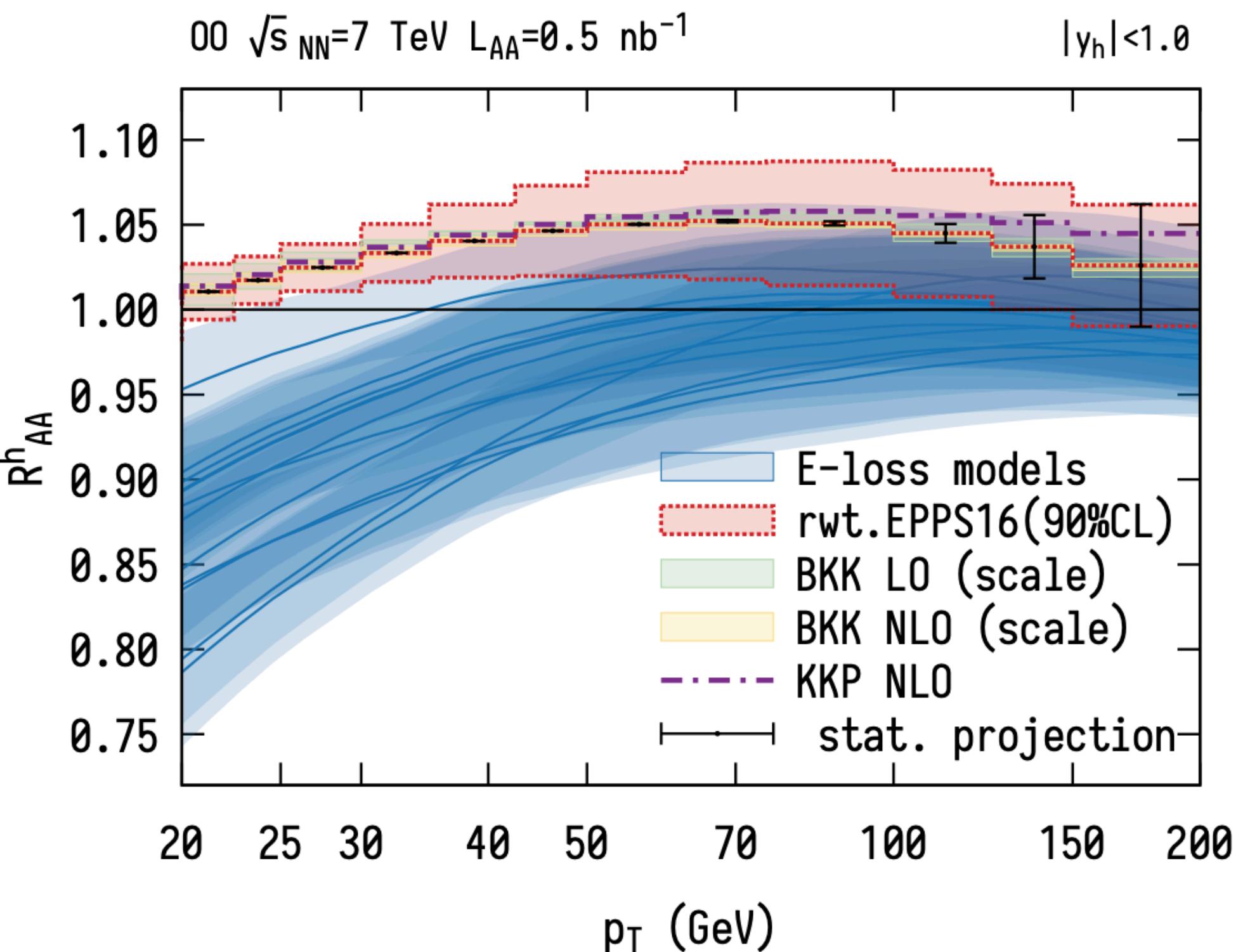
From dense to light systems

- Extrapolation from dense to light needs further understanding...
- Future oxygen runs can help us to determine the smallest amount of energy loss, provided that we control the initial state

[Huss, et al (2007.13754)]

Future OO run similar to PbPb peripheral
 (better suited to system-size dependence)

Future pO run crucial do reduce nPDF
 uncertainties





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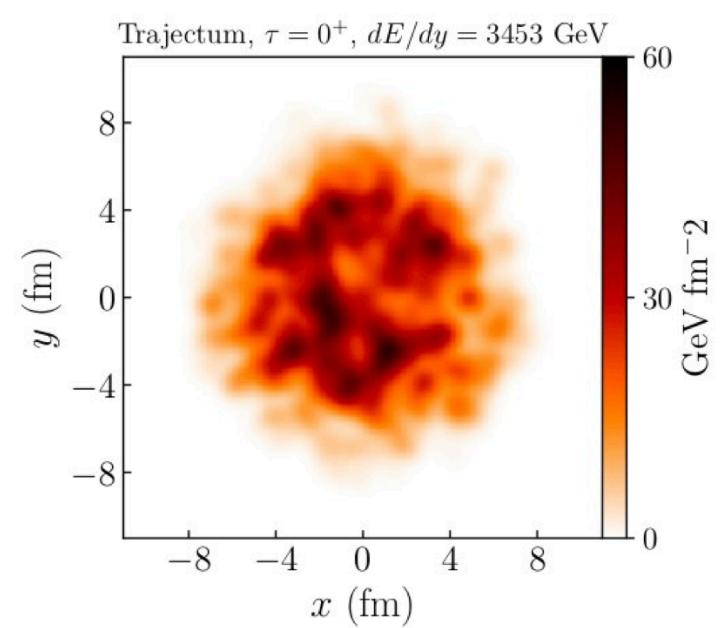
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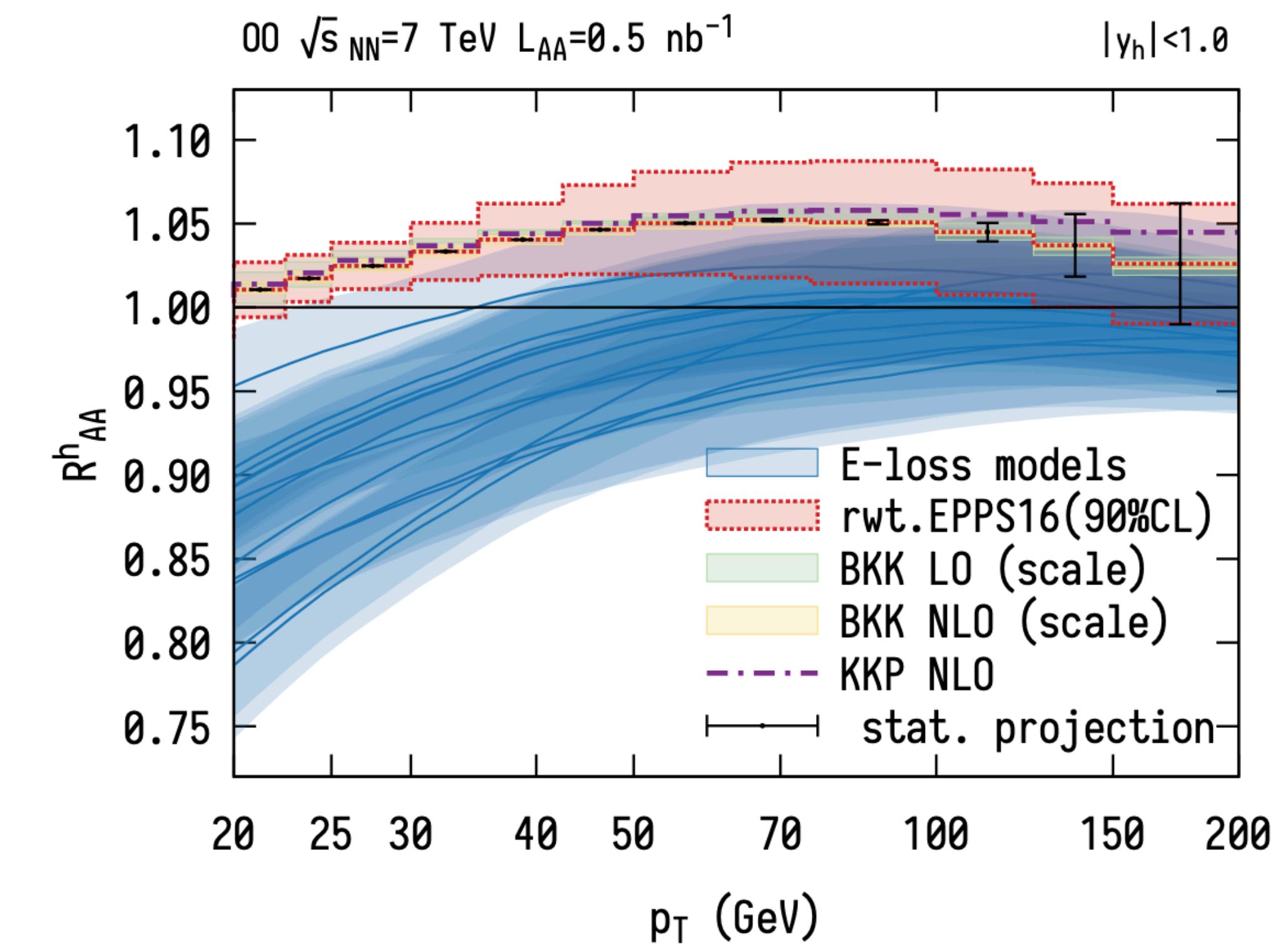
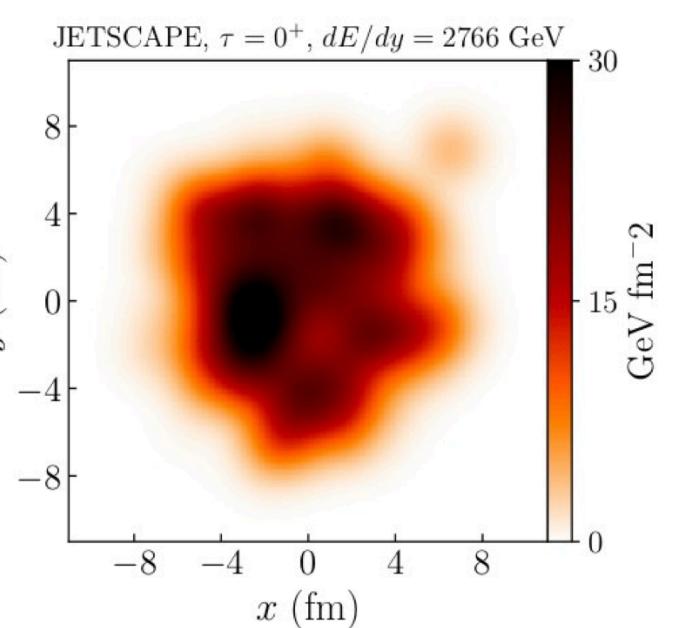
Future pO run crucial do reduce nPDF
 uncertainties

Cold or Hot nuclear matter effects?

Nucleon structure at
 high energy:



or

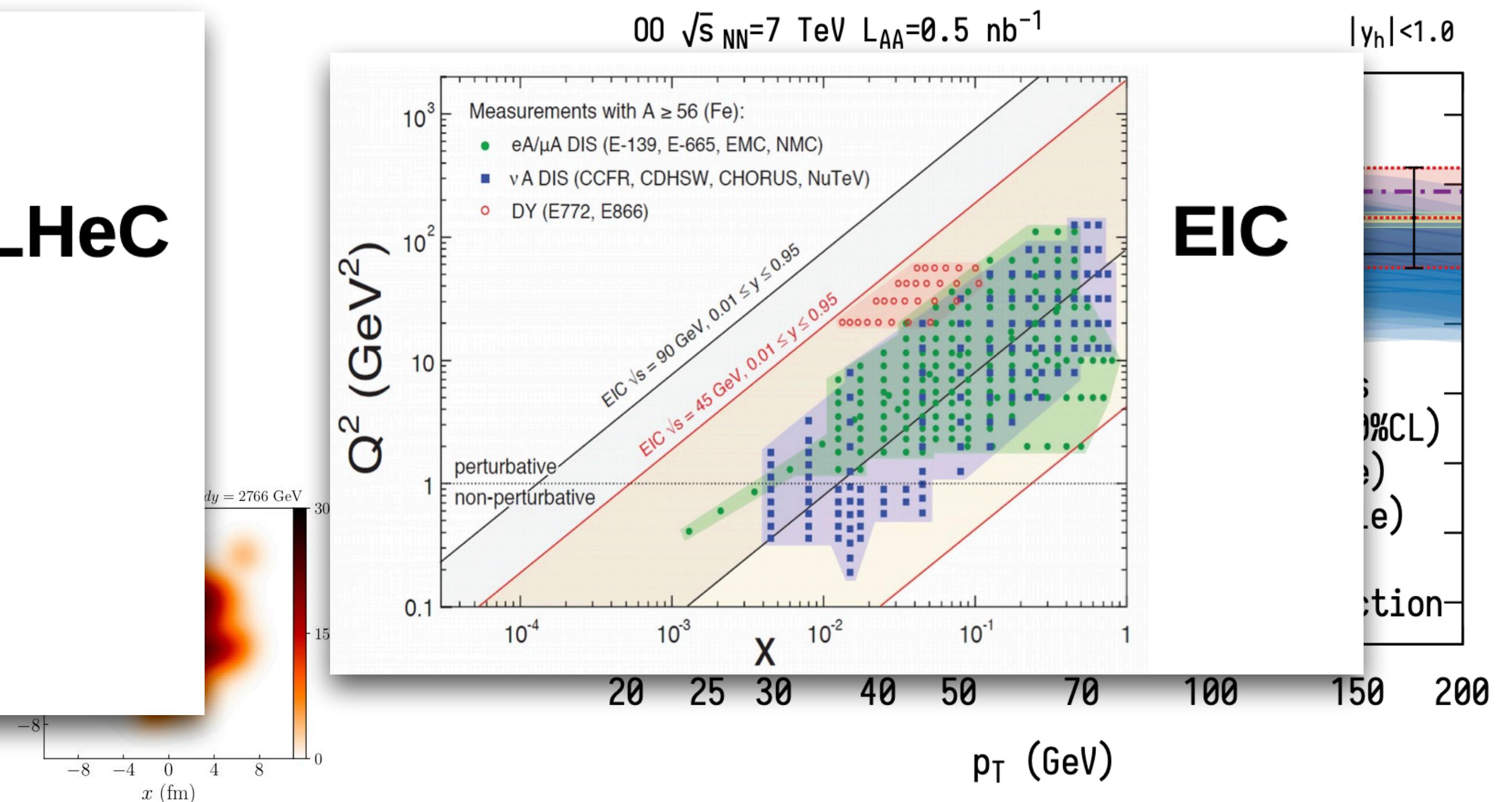
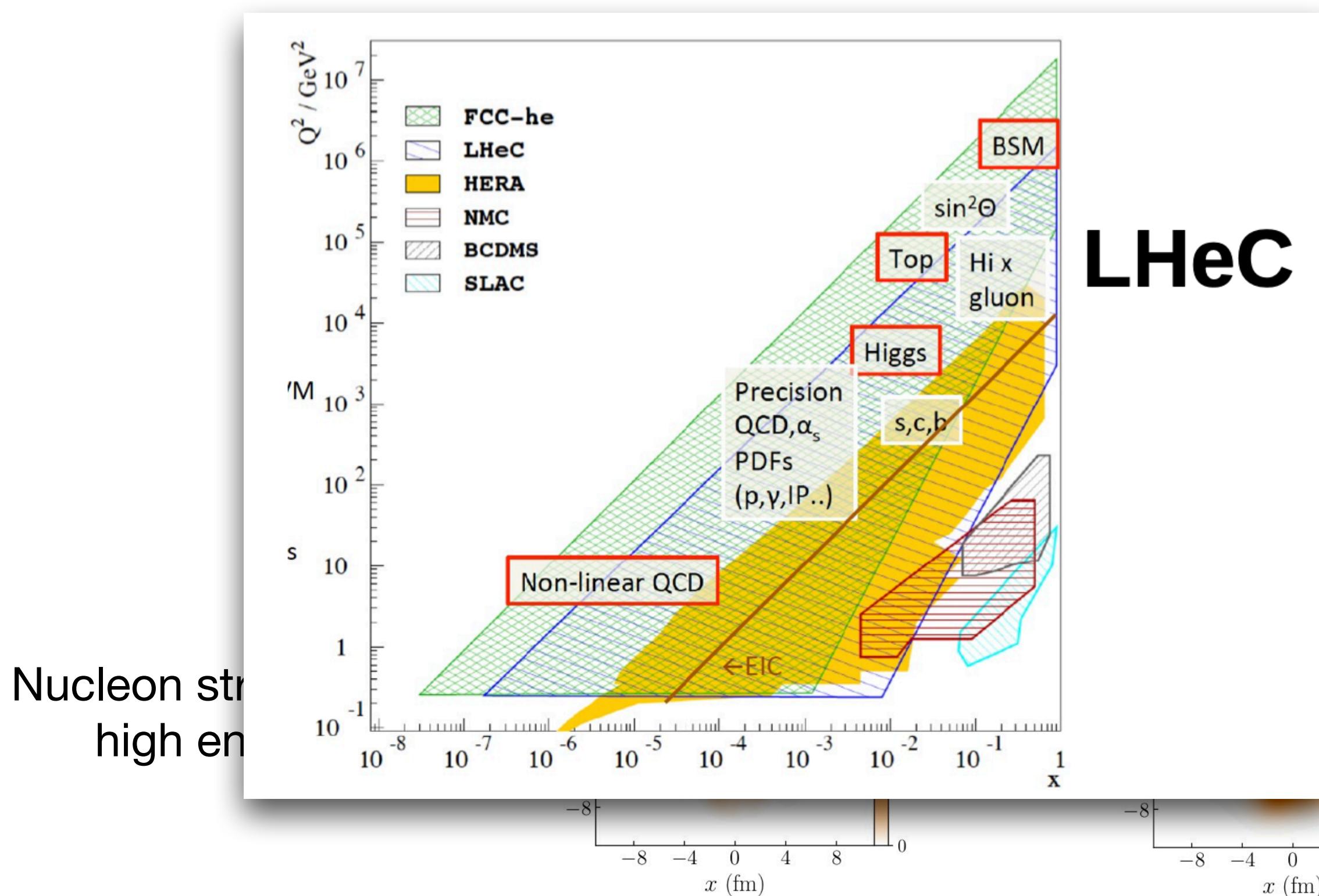


From dense to light systems

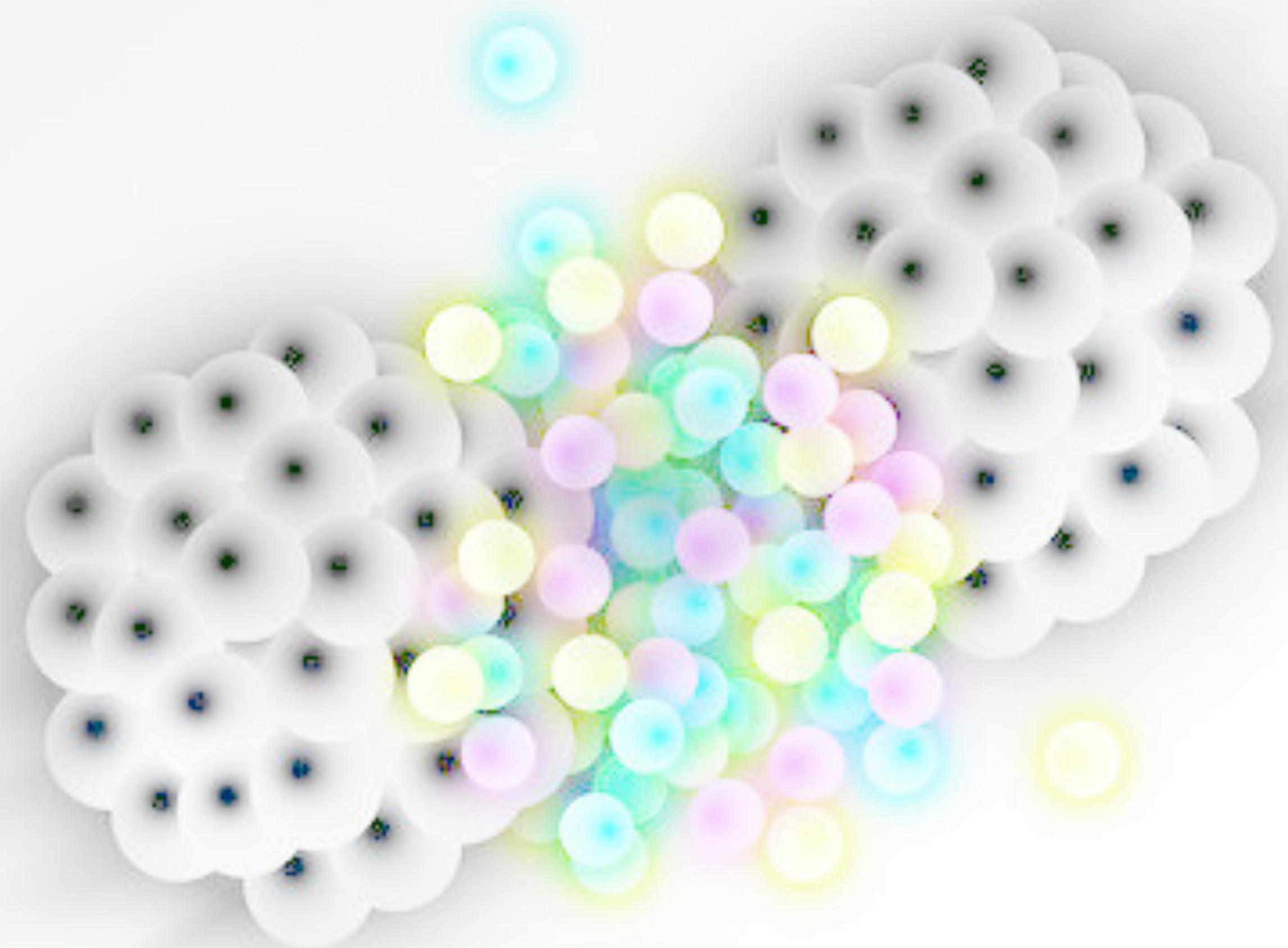
- Extrapolation from dense to light needs further understanding...
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See also Mikuni, Klest (Tue), Lim, Radhakrishnan, Morales, Vitev (Th)

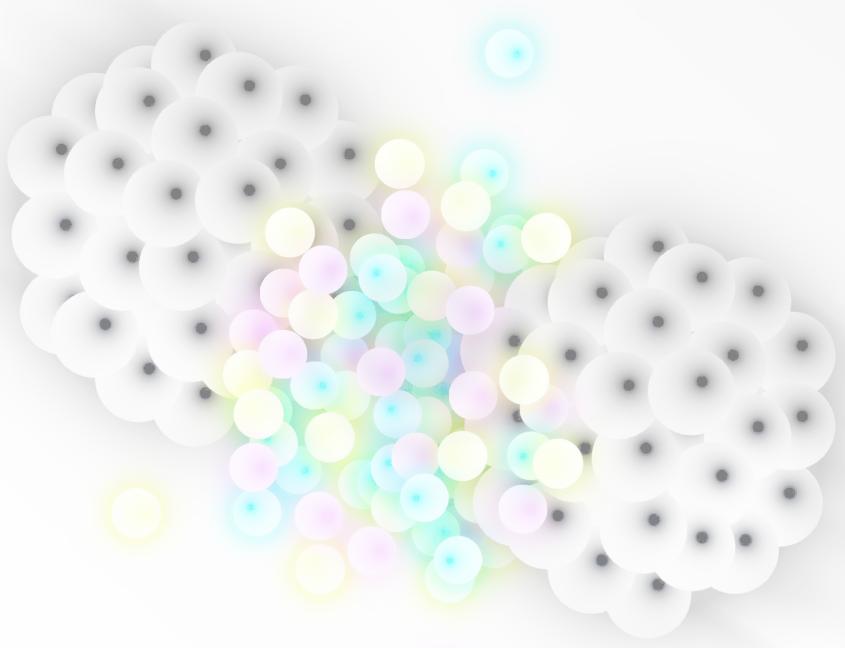
[Huss, et al (2007.13754)]



Wrapping up

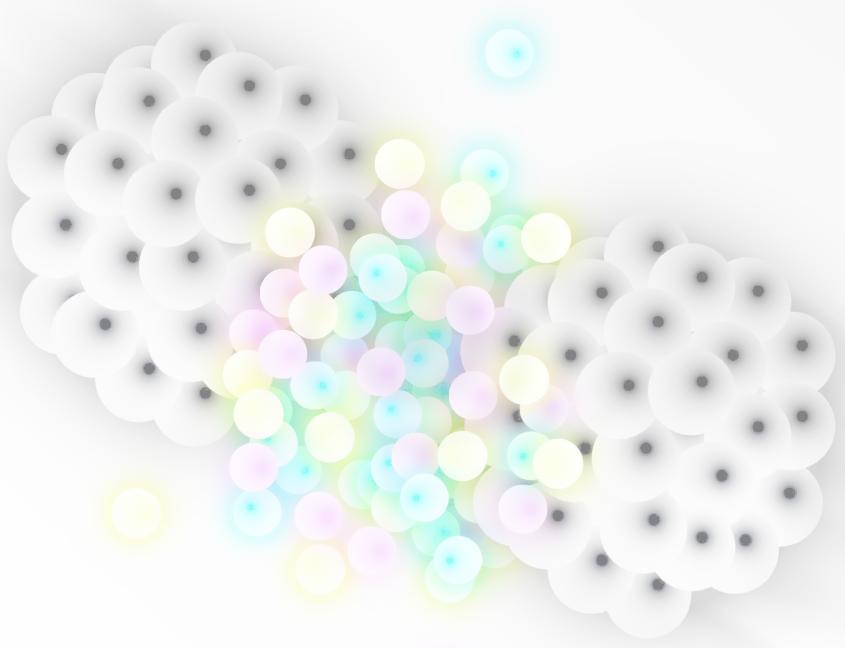


Summary



- **Heavy-ions** are a vibrant field full of **activity**
 - From far-from-equilibrium QCD to a fully thermalised medium
- **Quark-Gluon Plasma** studies have entered **precision** physics era
 - Determination of energy loss, momentum broadening and structure of a medium-modified parton showers
- **Future runs** will provide crucial input to many of our current **unsolved questions**
 - HL-LHC, sPHENIX, LHeC, EIC...

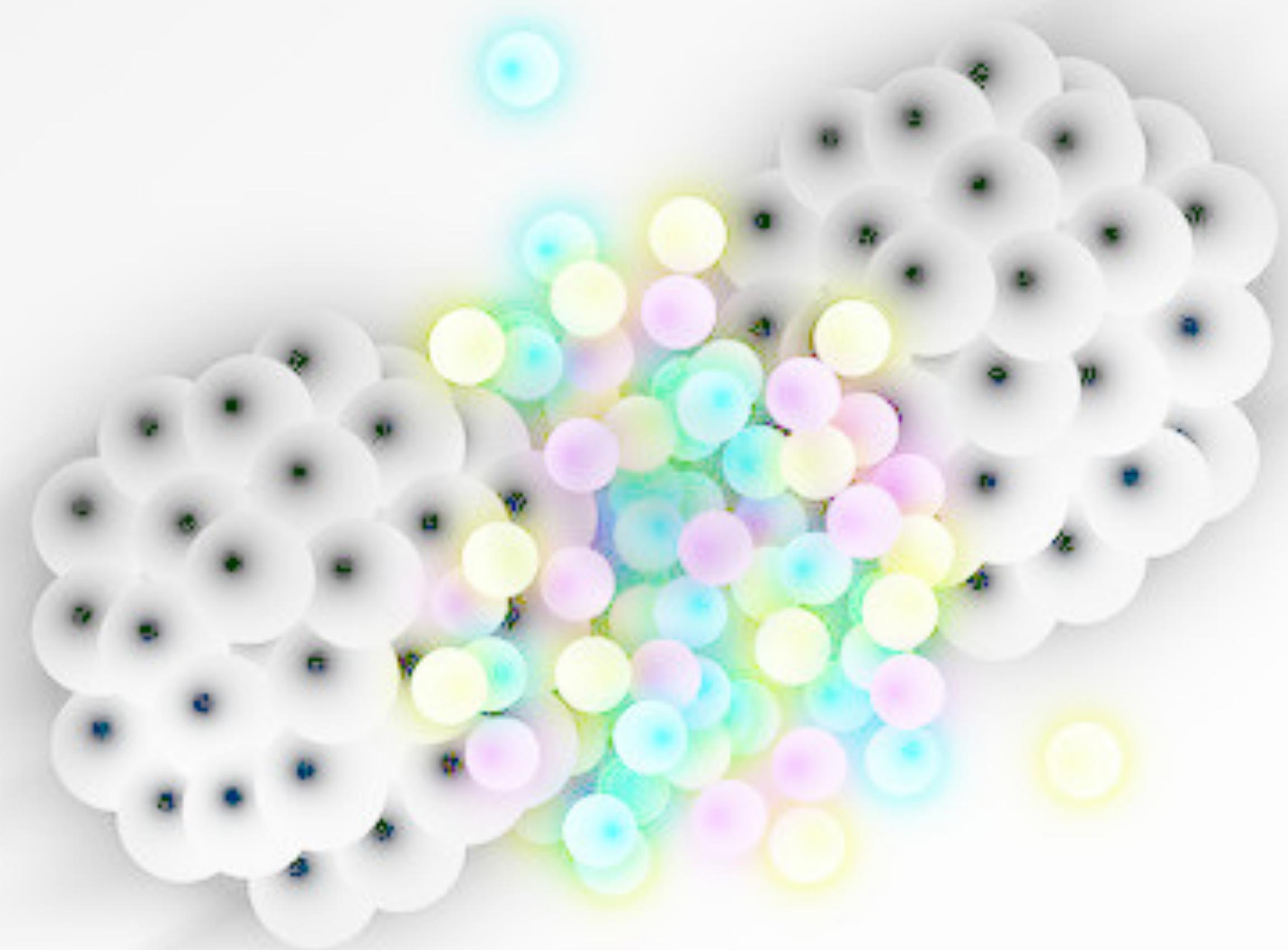
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Thank you!

Backup Slides



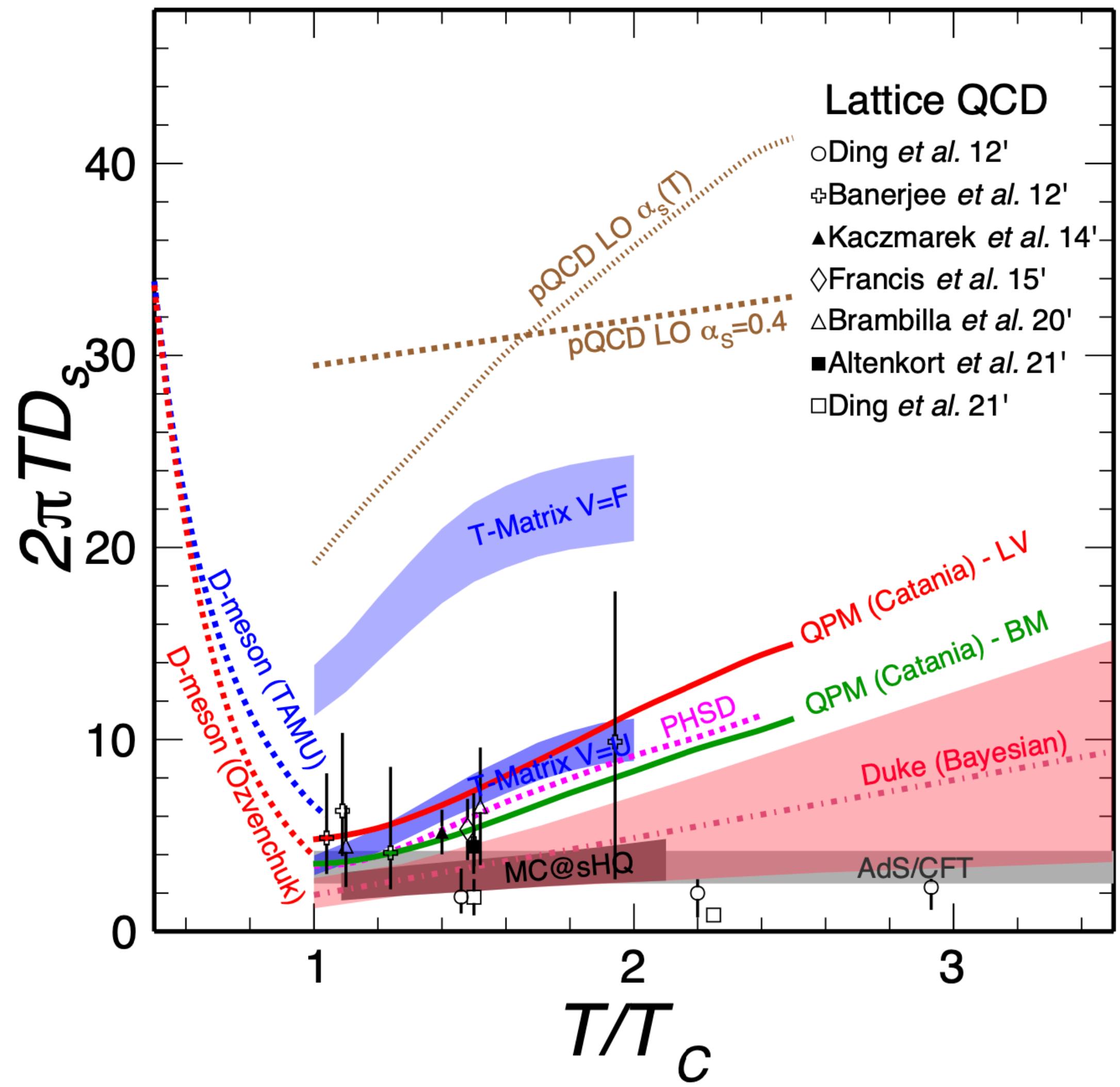
Heavy-Quark transport coefficients

[LA, Y-J Lee, M. Winn (2203.16352)]

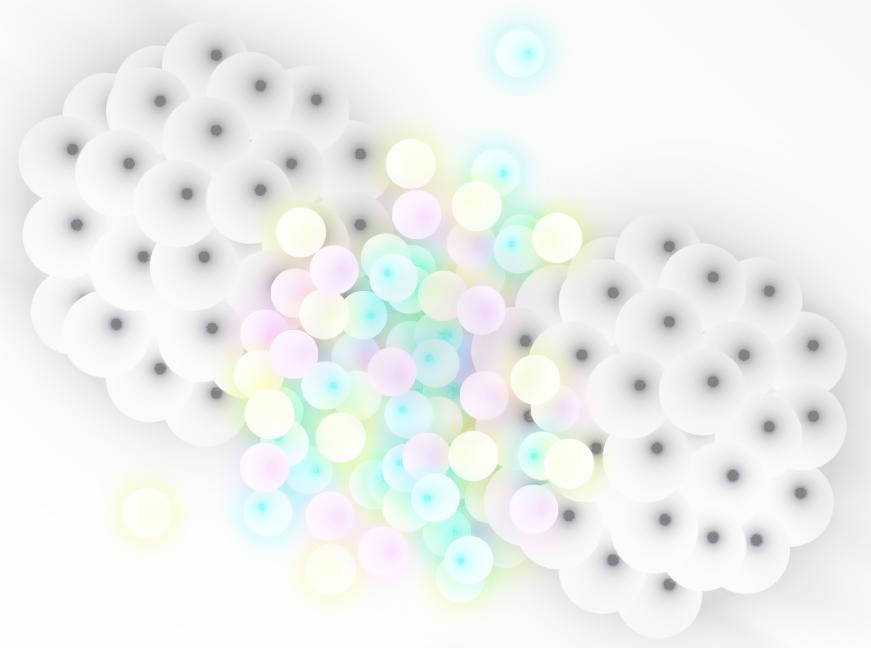
- Heavy-quark transport coefficients

$$D_s = \frac{d(\Delta E)^2}{dt}$$

$$\hat{e} = \frac{dE}{dt}$$



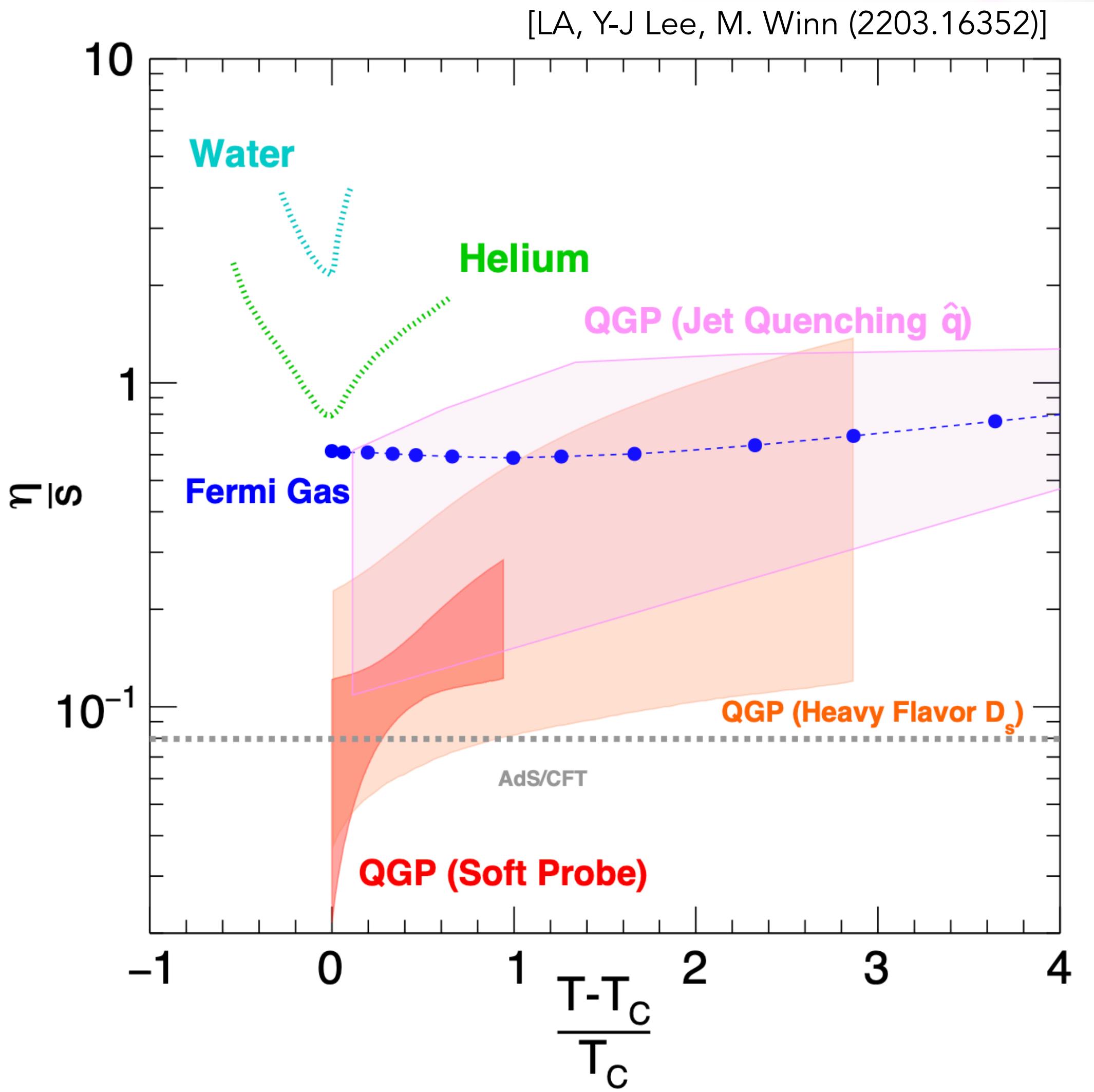
Soft vs Hard



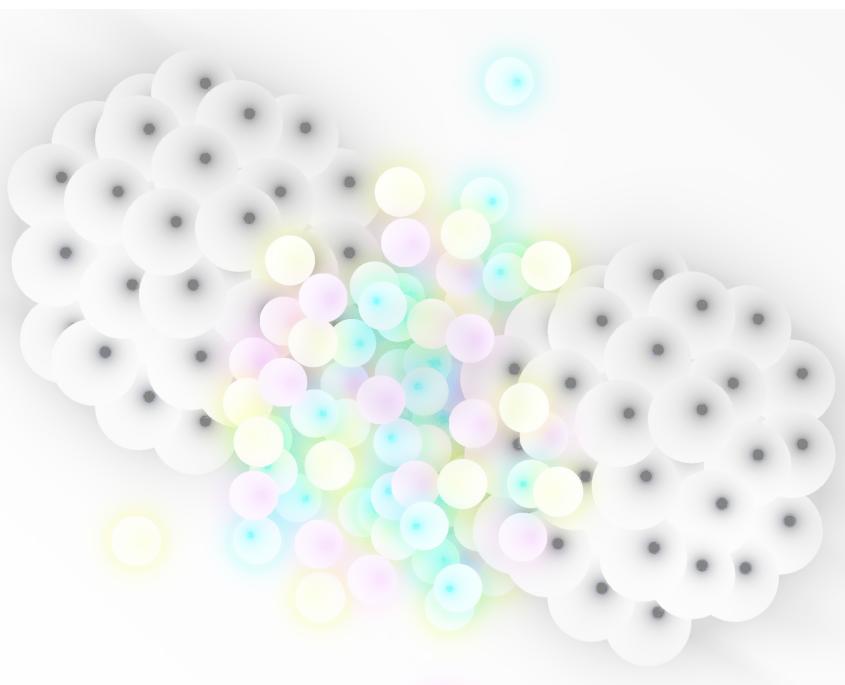
- Compilation of the specific shear viscosity as a function of temperature of the medium.

$$\frac{\eta}{s} = \frac{Ds(2\pi T)}{4\pi k}$$

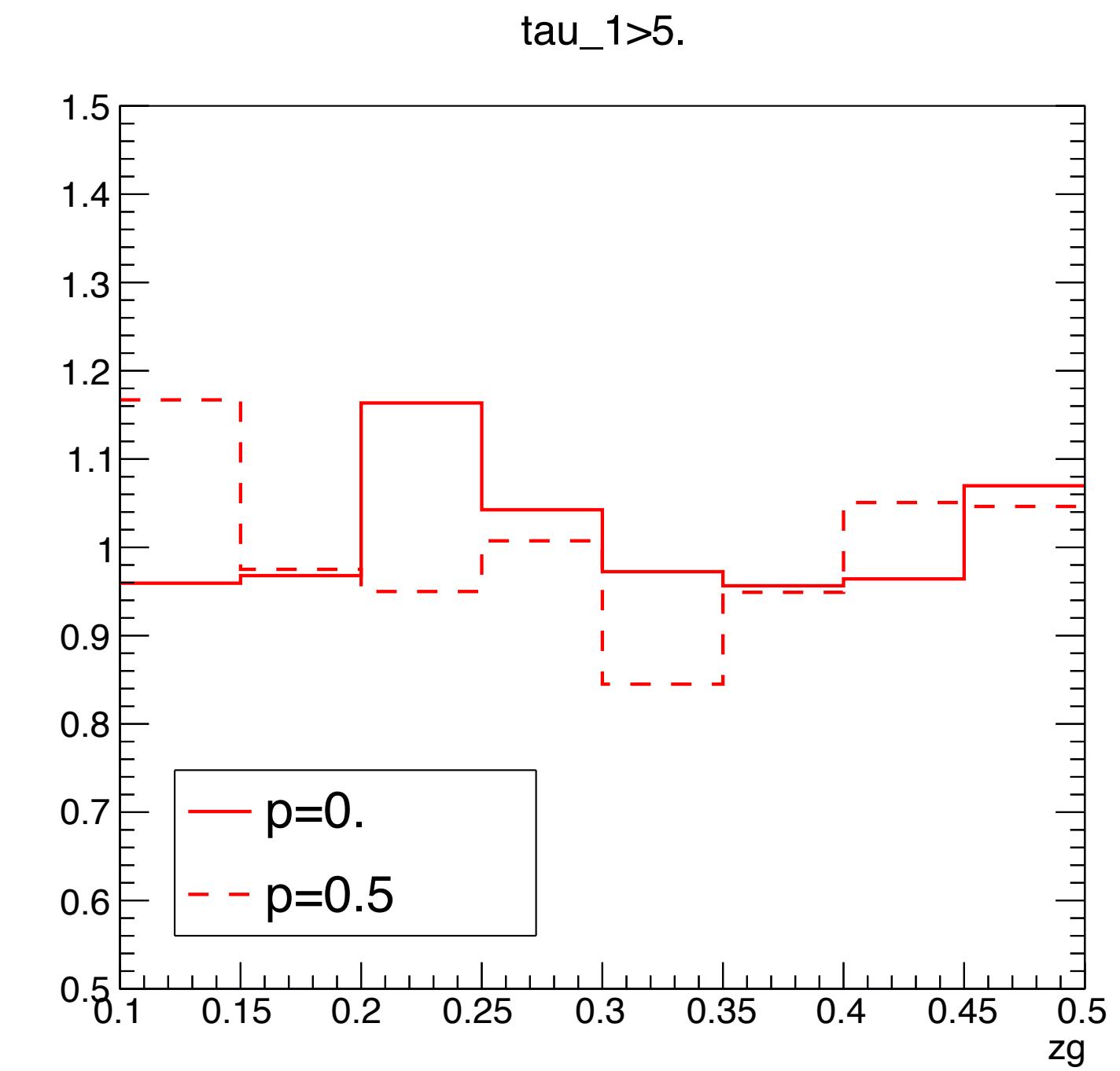
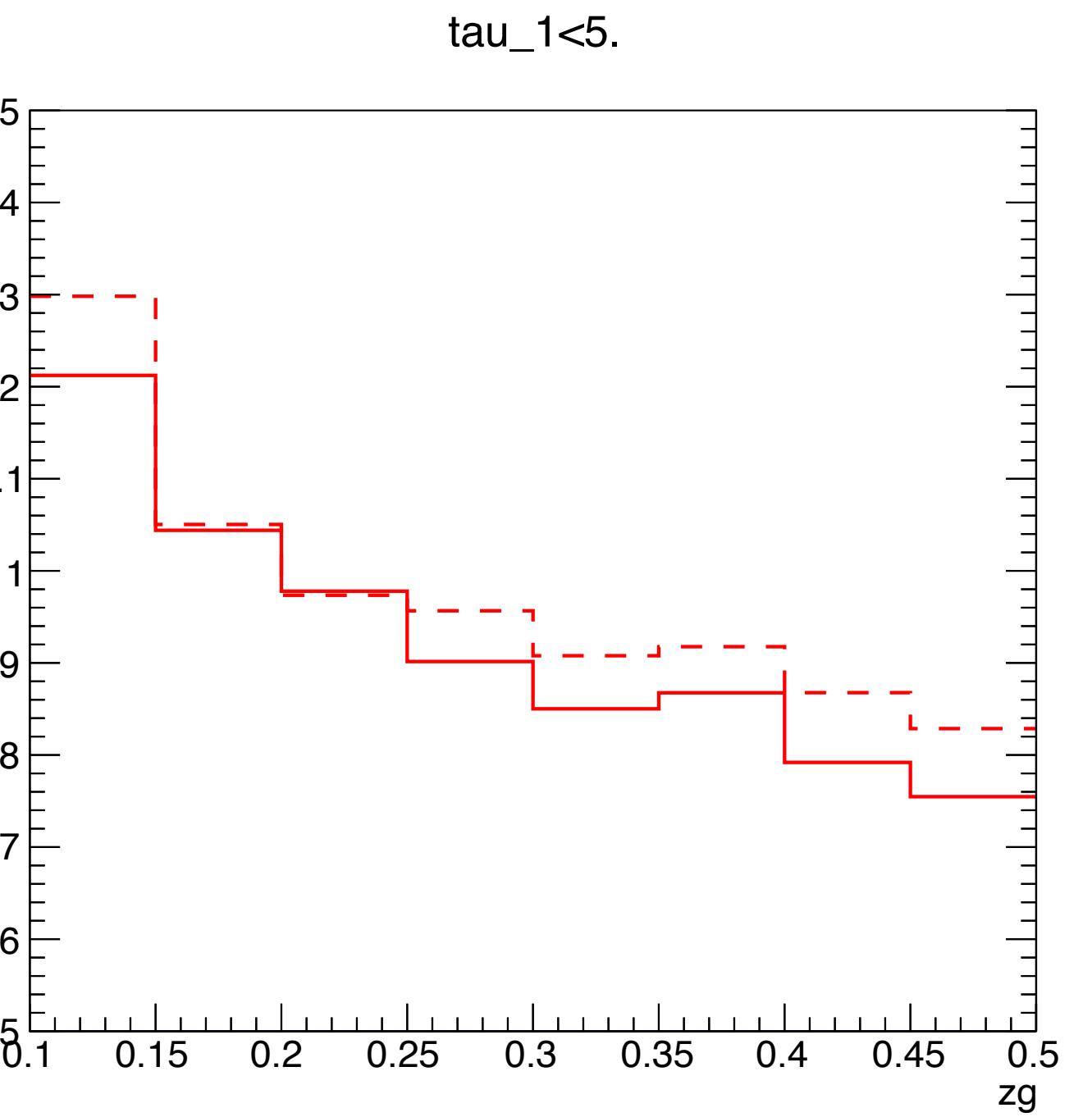
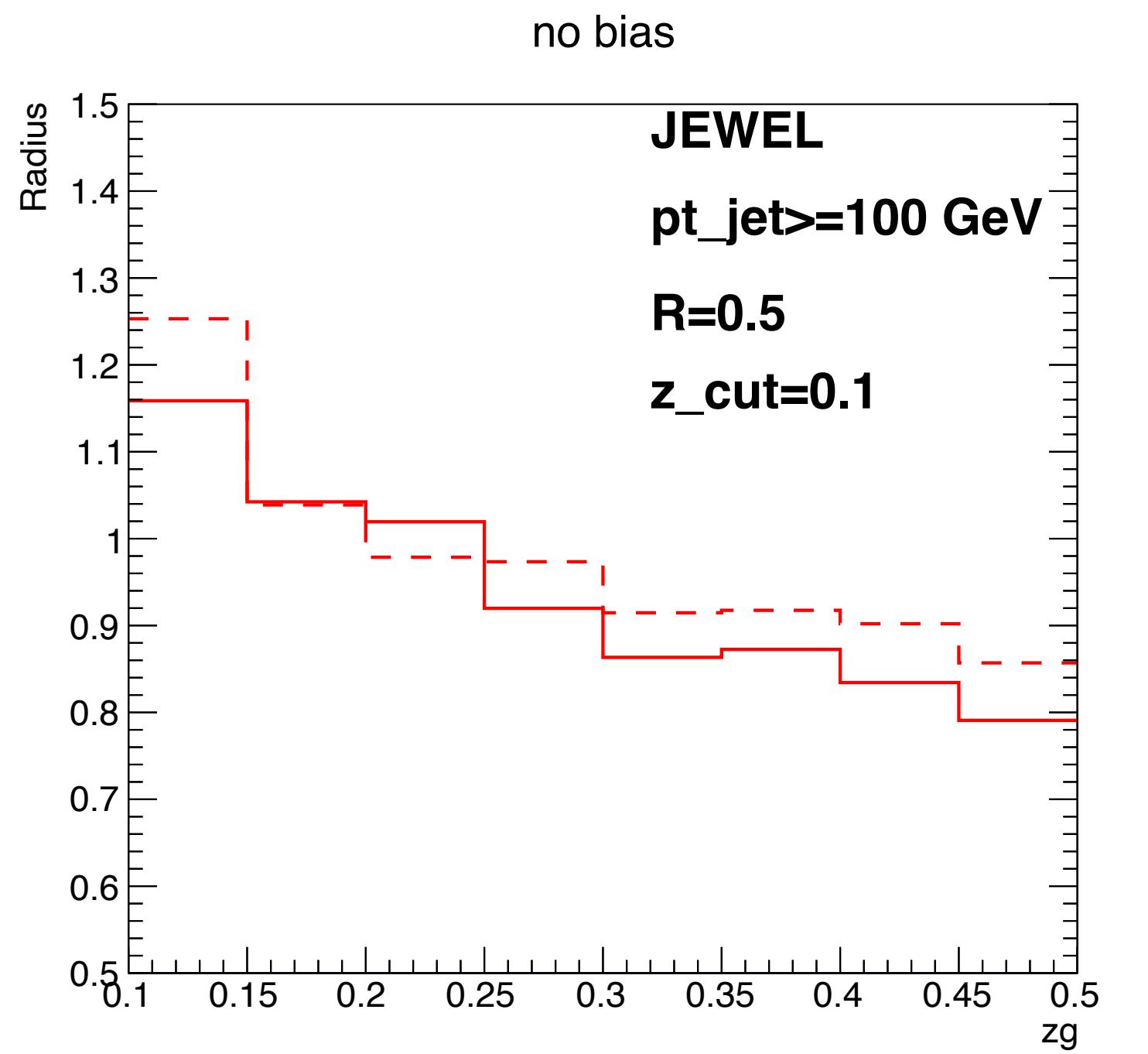
$$\frac{\eta}{s} \approx 1.25 \frac{T^3}{\hat{q}}$$



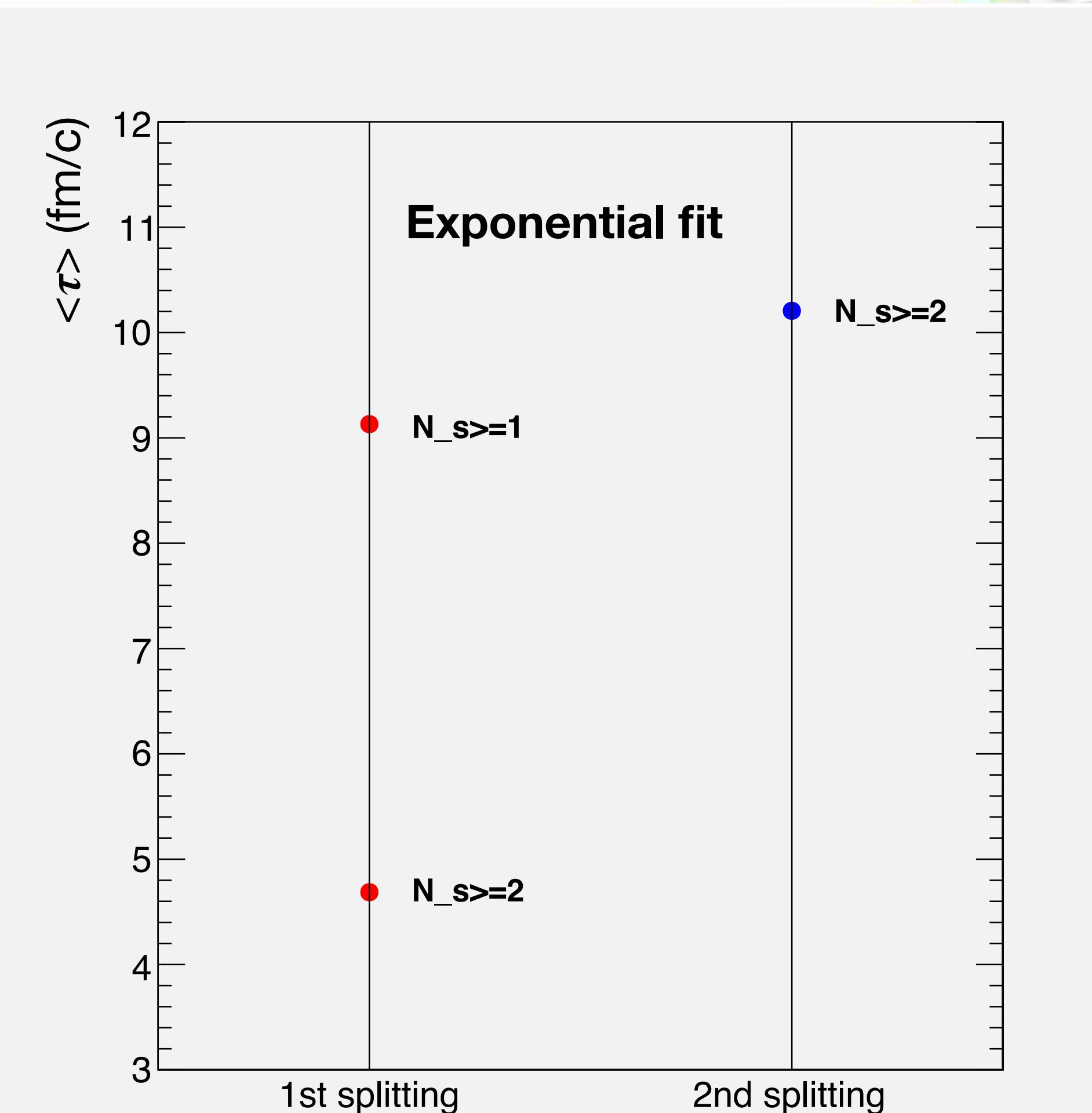
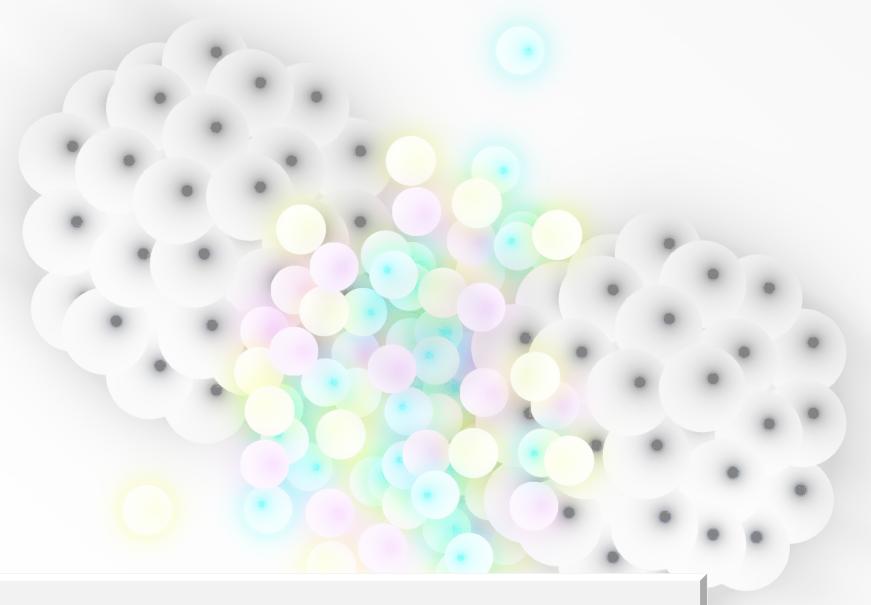
Zg 1st SD



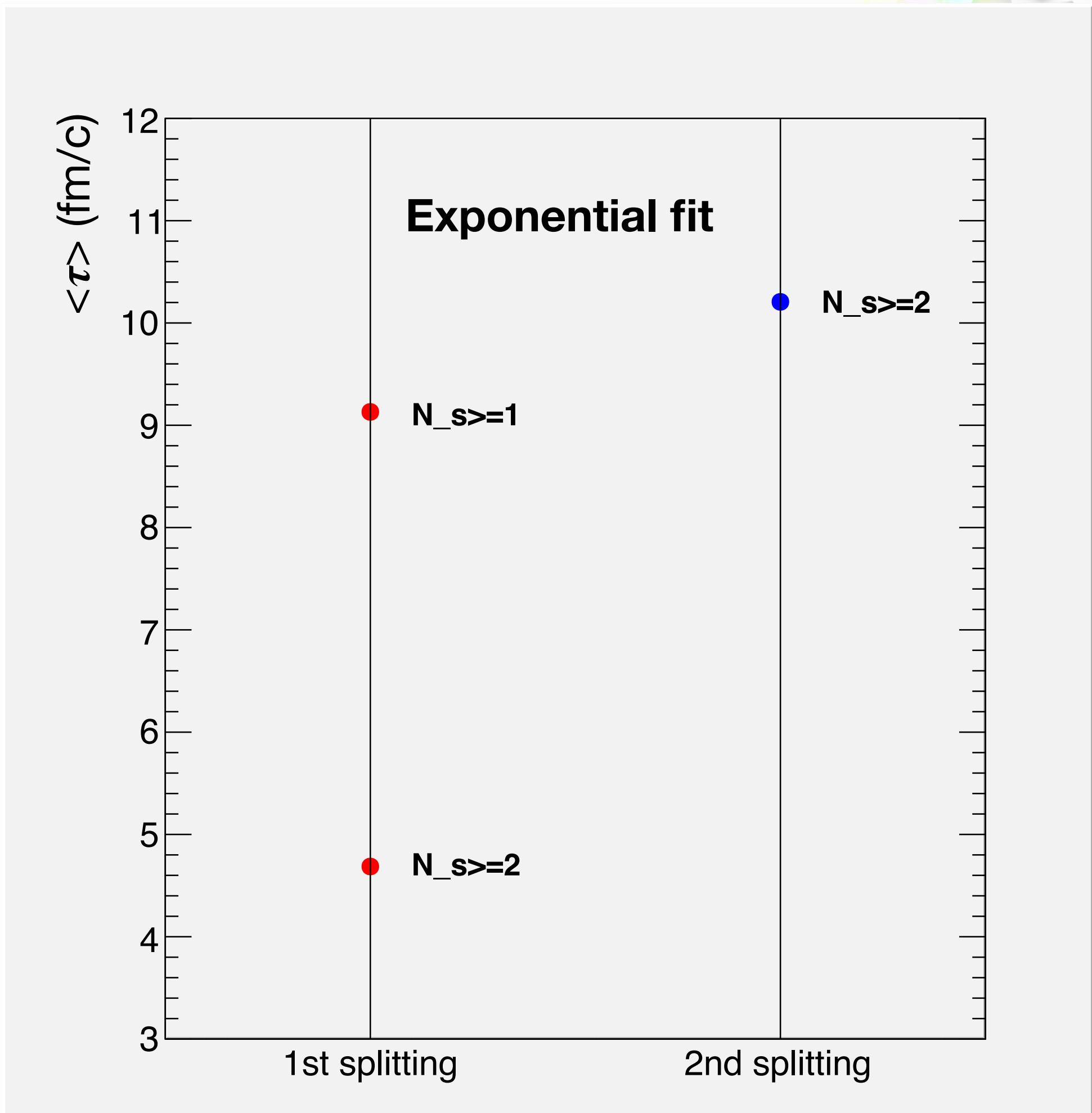
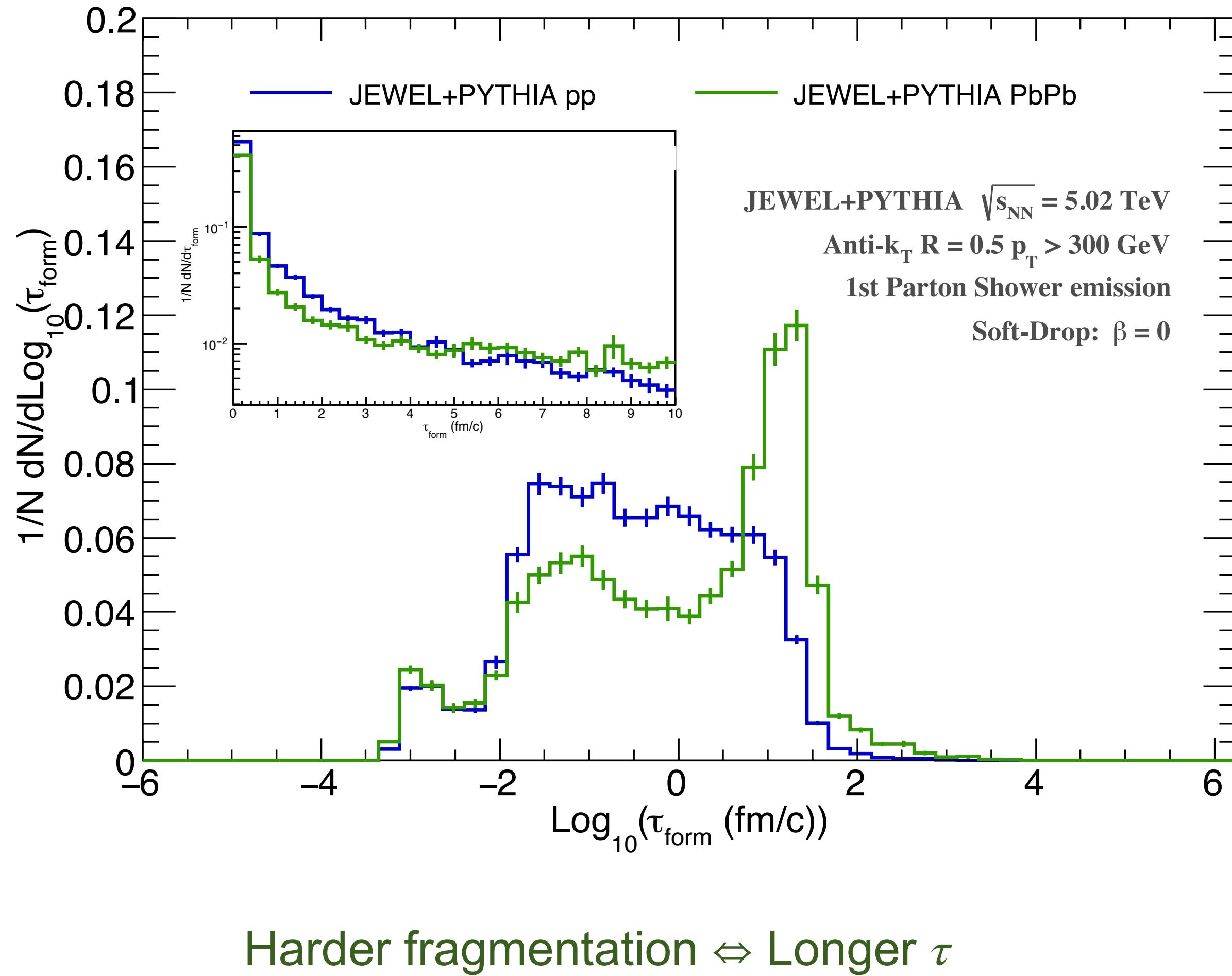
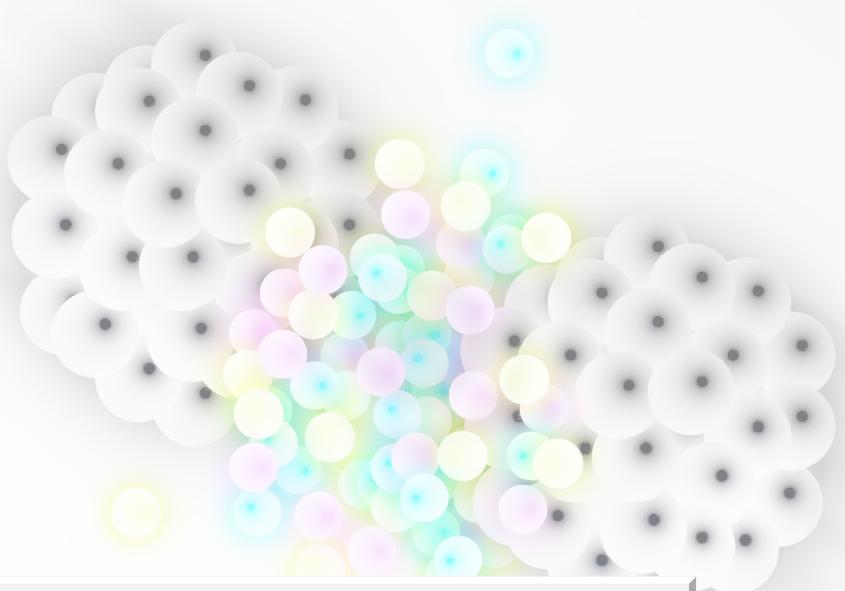
- Ratio of zg JEWEL (PbPb/pp):



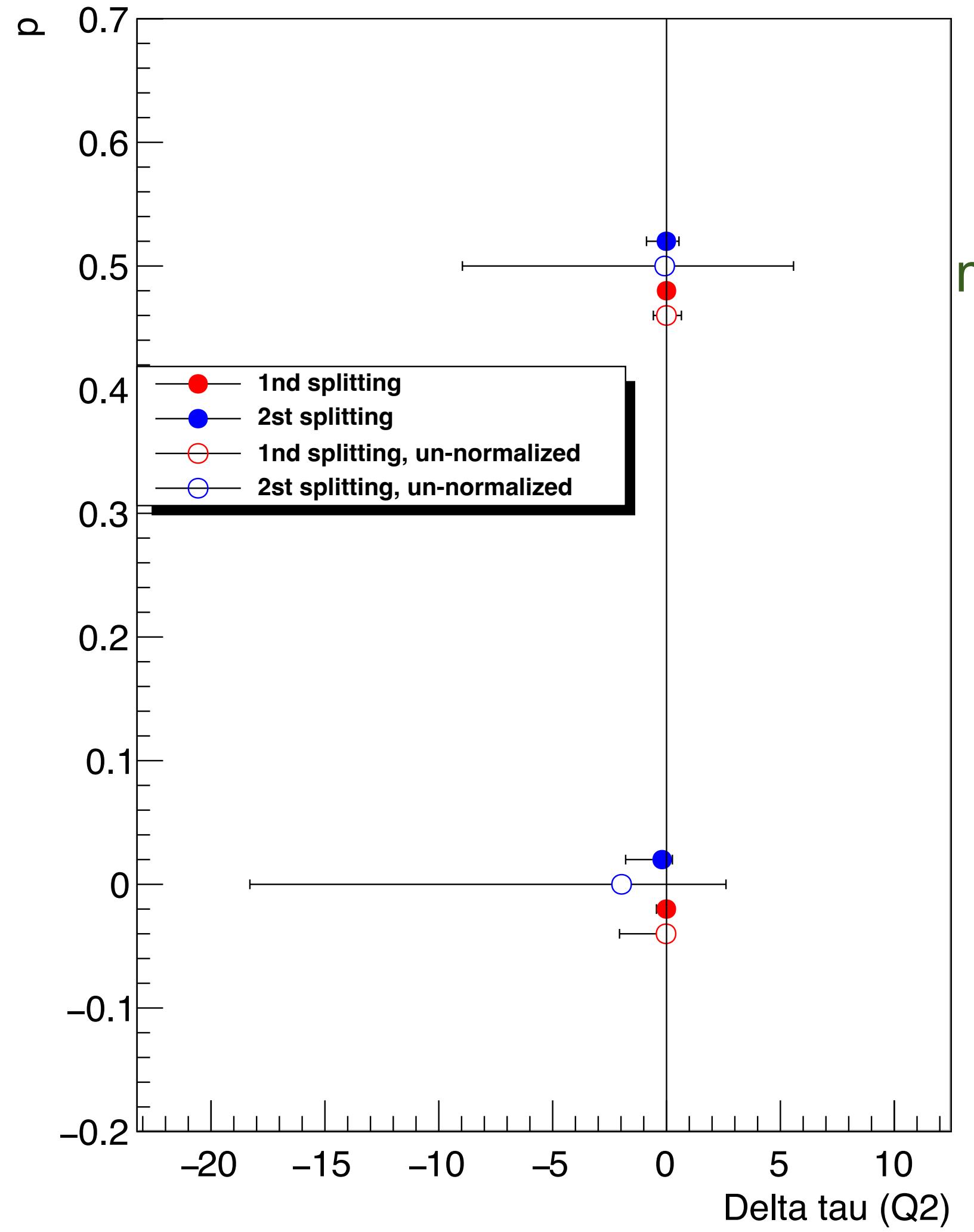
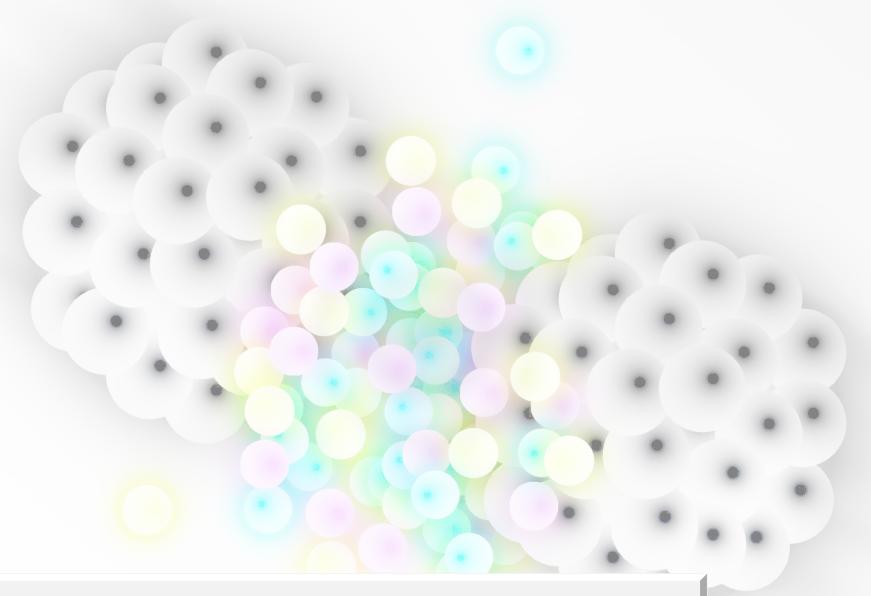
1st and 2nd Emissions



1st and 2nd Emissions

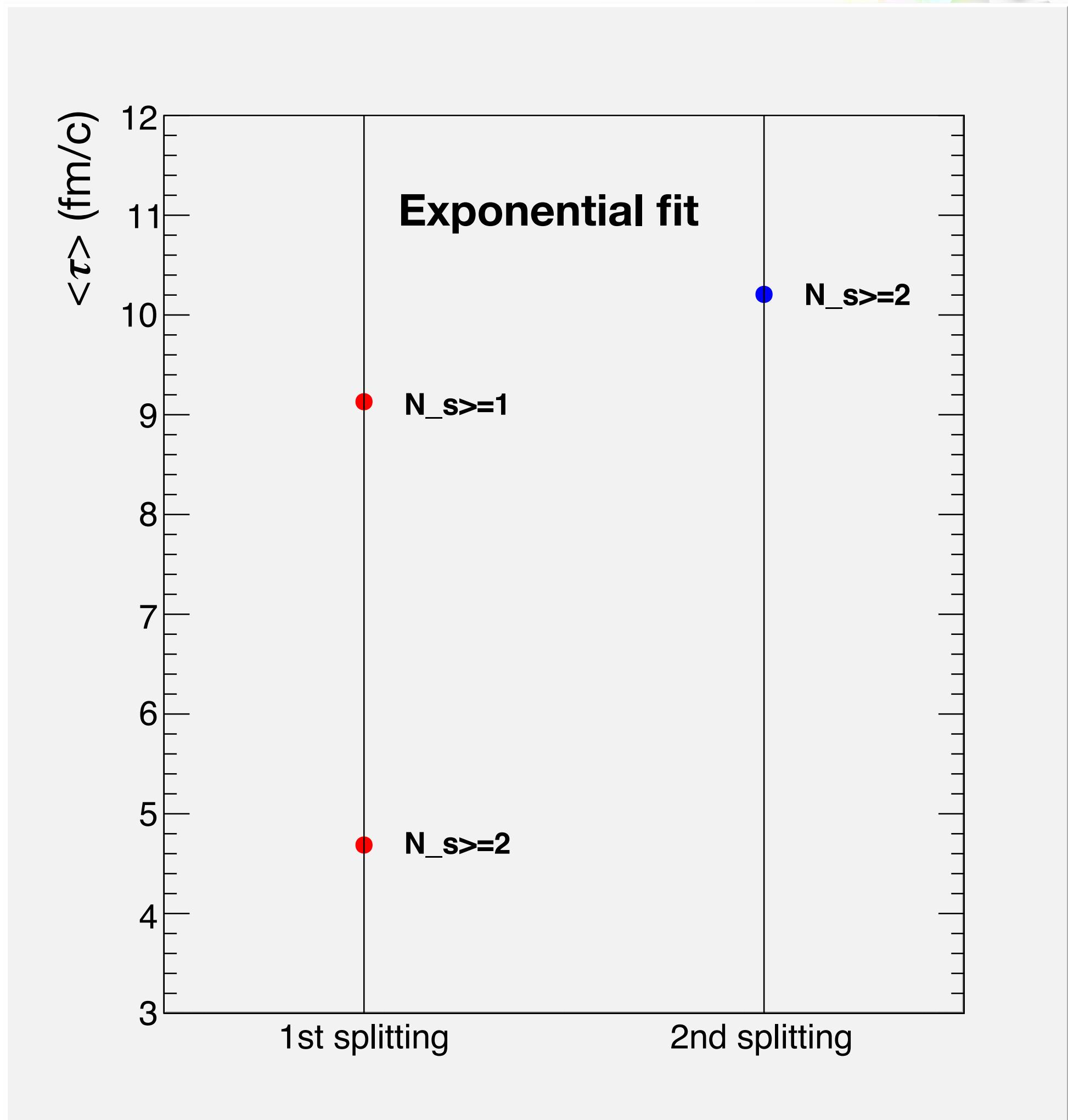


1st and 2nd Emissions



Select jets that fragment
more will increase τ window

Can we use two SD
handles per jet?



Exponential fit

$N_{\text{s}} >= 2$

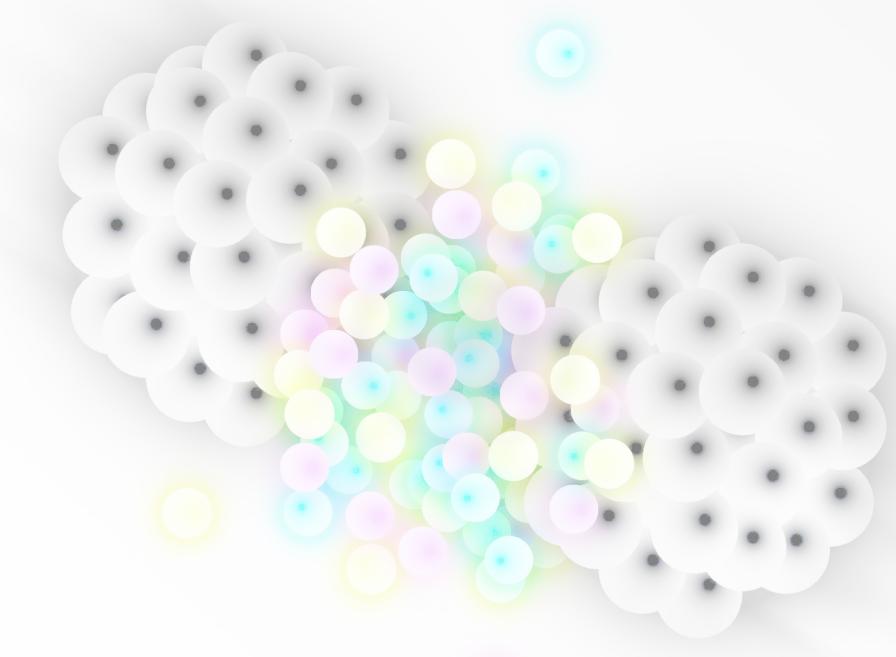
$N_{\text{s}} >= 1$

$N_{\text{s}} >= 2$

1st splitting

2nd splitting

Acknowledgments



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PORTUGUESA

FCT

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MINISTÉRIO DA EDUCAÇÃO E CIÉNCIA



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LISBOA