

# Measurement of the ${}^7\text{Be}(d,p){}^8\text{Be}^*$ reaction at 5 MeV/A

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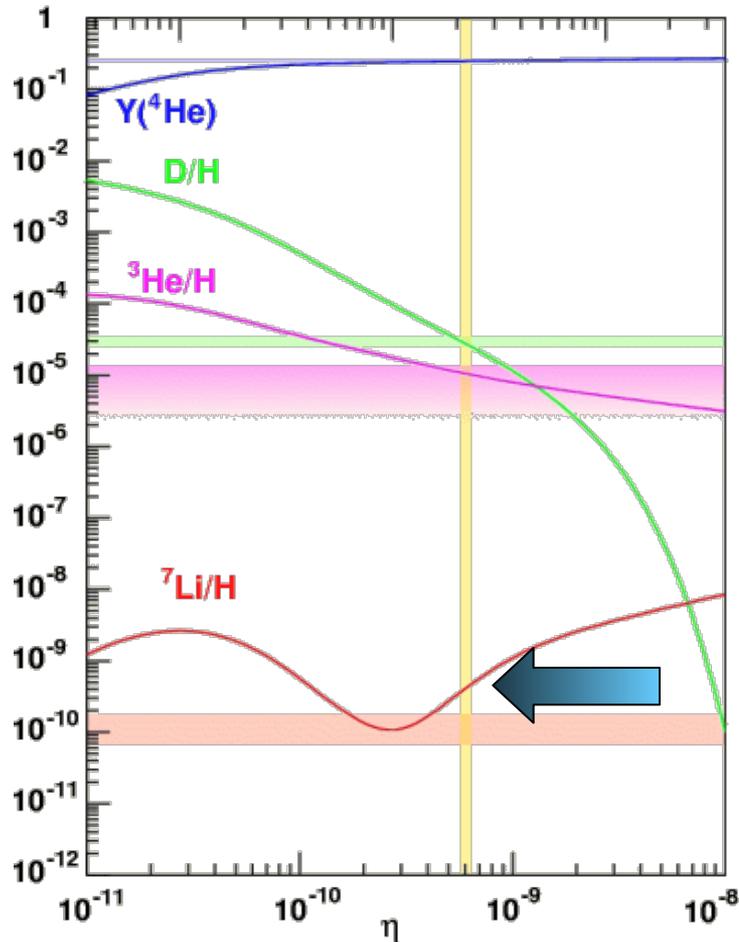


The standard **Big Bang** model of the Primordial Universe is very successful in accounting for the observed relative abundance of the light elements.

The only astrophysical input to the Big Bang Nucleosynthesis (BBN) calculation is the **baryon density** of the Universe, which is now known precisely.

However, BBN theory fails to predict correctly the **observed abundance of  ${}^7\text{Li}$** .

# The Cosmological $^7\text{Li}$ problem



*Observed values represented by bands,  
predicted values represented by lines*

$$\eta = n_B/n_\gamma = 6.104 \pm 0.058 \times 10^{-10}$$

baryon-to-photon ratio

BBN theory over predicts the abundance of  $^7\text{Li}$  by about a factor  $\sim 3$  and up to five sigma deviation from observation. The theory uses the **baryon-to-photon** ratio  $\eta$  from measurements of **cosmic microwave background**.

BBN theory using  $\eta$ :  $\frac{^7\text{Li}}{\text{H}} = 5.12_{-0.62}^{+0.71} \times 10^{-10}$

Observationally extracted:  $\frac{^7\text{Li}}{\text{H}} = 1.58_{-0.20}^{+0.35} \times 10^{-10}$

*Serious discrepancy*

Good agreement of BBN predicted abundances with observations for  $^2\text{H}$ ,  $^3,^4\text{He}$ .

For decades, one of the  
**important unresolved problems**

# Nuclear physics aspects of the primordial lithium problem

## Astrophysical solutions

Improvements in the observationally inferred primordial lithium abundance. Lithium may be destroyed in metal-poor stars through diffusion and turbulent mixing. *Korn, Nature (2006); Ryan (1999)*

## Physics beyond standard BBN

Destruction of mass-7 nuclides through interaction with WIMP particles, unstable particles in the early universe that could have affected BBN. Existence of  $^8\text{Be}$  as a bound nuclide during BBN. Interpretations assumed nuclear reaction rates known accurately *Goudelis (2016), Coc (2012), Fields (2011), Cyburt (2006)*

## Nuclear physics

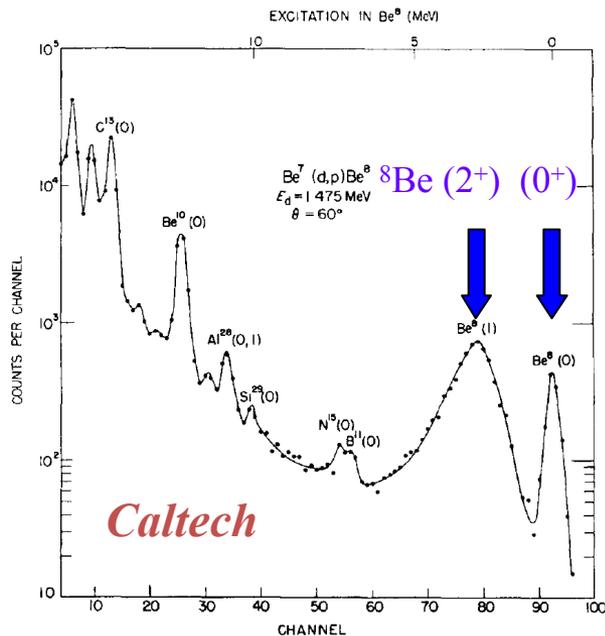
In the condition of BBN,  $^7\text{Li}$  is effectively destroyed through  $^7\text{Li}(p,\alpha)^4\text{He}$ , so that 95% of the primordial  $^7\text{Li}$  is the by-product of the electron capture  $\beta$ -decay of the primordial  $^7\text{Be}$  after the cessation of nucleosynthesis.

Nuclear aspects of the  $^7\text{Li}$  problem would involve the reaction rates of  $^7\text{Be}$  production, mainly  $^3\text{He}(\alpha,\gamma)^7\text{Be}$  and its destruction through  $^7\text{Be}(n,p)^7\text{Li}$ ,  $^7\text{Be}(n,\alpha)^4\text{He}$  and  $^7\text{Be}(d,p)2\alpha$ .

# Incomplete nuclear physics input for BBN calculations: Can resonant enhancement alleviate this discrepancy?

It has been argued that the  ${}^7\text{Li}$  discrepancy could be resolved, if the  ${}^7\text{Be}(d,p)$  reaction rate is substantially larger than previously considered.

*R. W. Kavanagh*  
Nuclear Physics 18 (1960) 492



Experimental data at cm energies of 0.6 – 1.3 MeV. The reaction rate relied on an extrapolation to lower energies. Protons corresponding to the  ${}^8\text{Be}$   $0^+$  g.s and  $1^{\text{st}}$  excited state (3.03 MeV,  $2^+$ ) were detected, up to excitation energies of 11 MeV.

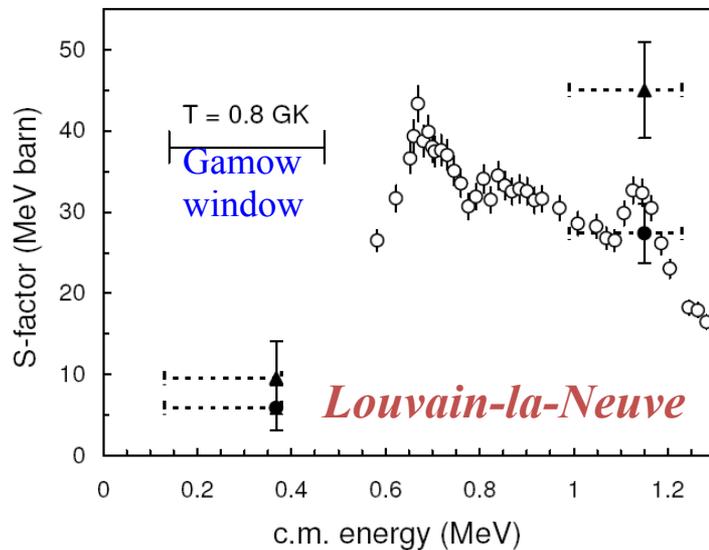
Lacking complete angular distributions, these data were converted to total cross section by multiplying by  $4\pi$  and a factor of  $\sim 3$  to take in to account contributions from **higher excited states in  ${}^8\text{Be}$** . A constant S-factor  $\sim 100$  MeV-barn was adopted. *Parker (1972)*

An experiment performed at lower energy found a significantly reduced cross-section in the BBN Gamow window compared to Parker's estimate.



Angulo et al

Astrophys. Jour. 630 (2005) L105



Cross section was measured at  $E = 5.55, 1.71$  MeV, up to excitation energies in  ${}^8\text{Be}$  of 13.8 MeV. In addition to feeding of the g.s and 1<sup>st</sup> ex states of  ${}^8\text{Be}$ , able to observe **higher energy levels** mainly through the broad 11.4 MeV ( $4^+$ ) state.

Higher energy states not observed by *Kavanagh* contribute about 35% of the total S-factor. Reaction rate is smaller by a factor of  $\sim 2$  at 1.0-1.23 MeV and by  $\sim 10$  at energies relevant to BBN.

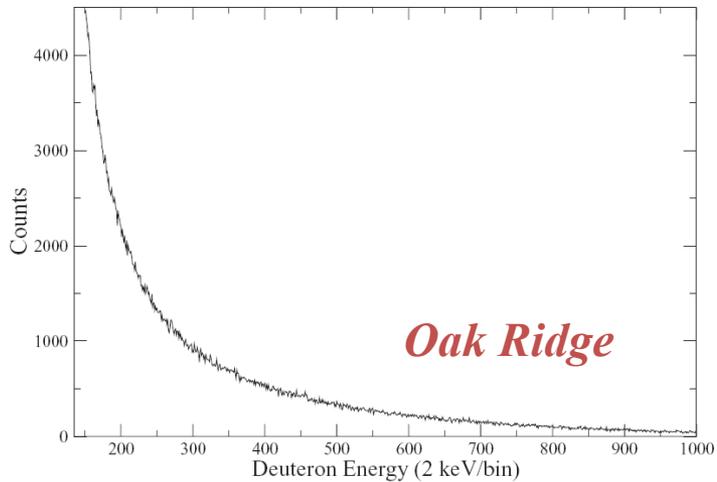
- Kavanagh (1960)
- Angulo (2005), data includes contribution from the g.s + 1<sup>st</sup> excited state of  ${}^8\text{Be}$  only
- ▲ Angulo (2005) Total S-factor

The **S-factor** at BBN energies was not underestimated by Parker, but on the contrary, **overestimated**.

Other works suggested resonant enhancement through a 16.7 MeV ( $5/2^+$ ) resonance state in  $^9\text{B}$  *Cybert (2005), Chakravorty (2011)*

*O'Malley et al*

*Phys. Rev. C 84, 042801( R) (2011)*



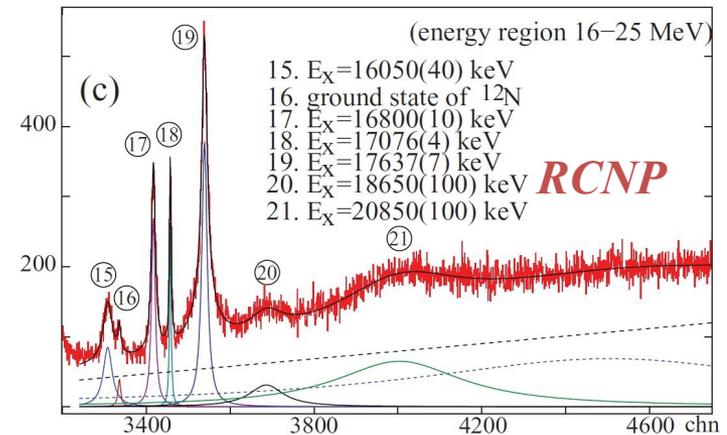
$^2\text{H}(^7\text{Be},d)^7\text{Be}$  ( $E_{^7\text{Be}} = 10 \text{ MeV}$ )

**No evidence for a resonance observed**

*Scholl et al Phys. Rev. C 84, 014308 (2011)*

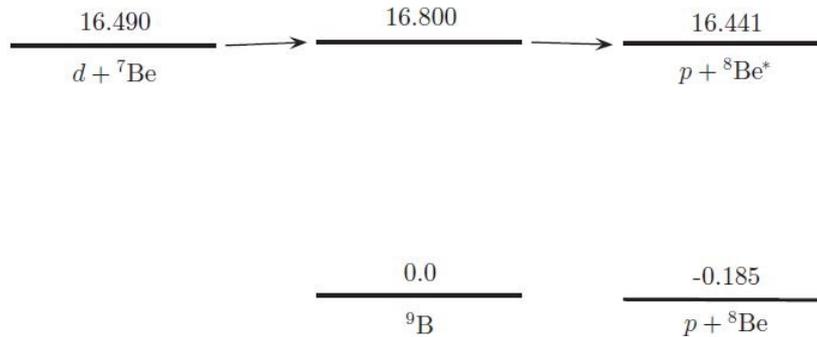
High resolution study of  $^9\text{Be}(^3\text{He},t)^9\text{B}$ ,  $E = 140 \text{ MeV/A}$ , the state is strongly excited.

Energy: **16.800(10) MeV**, width: 81(5) keV



Without experimental knowledge on its decay properties, conclusion about resonant enhancement to the  $d + ^7\text{Be}$  reaction remain uncertain.

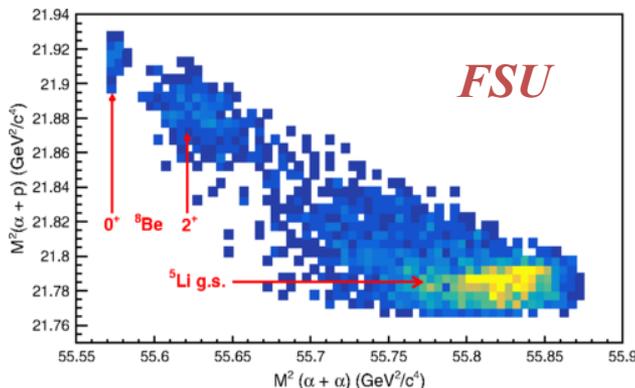
## Proposed ${}^7\text{Be}$ destruction mechanism, $d + {}^7\text{Be} \rightarrow {}^9\text{B}^* \rightarrow p + {}^8\text{Be}^*$



The 16.8 MeV state in  ${}^9\text{B}$  formed by fusion of  ${}^7\text{Be} + d$  and decays by proton emission to a **highly excited state in  ${}^8\text{Be}$ , 16.626 MeV** above the ground state, which subsequently breaks up into two  $\alpha$  particles.

*O.S.Kirsebom et al., Phys. Rev. C 84, 058801 (2011)*

However, recent work (2019) shows,  $d + {}^7\text{Be} \rightarrow 2\alpha + p$  may proceed through intermediate state in  ${}^8\text{Be}$  by  ${}^7\text{Be}(d,p){}^8\text{Be}(\alpha){}^4\text{He}$  or  ${}^5\text{Li}$  by  ${}^7\text{Be}(d,\alpha){}^5\text{Li}(p){}^4\text{He}$  sequence, or in a “democratic” three-particle decay of the  ${}^9\text{B}$  compound system.



*Rijal et al Phys. Rev. Lett. 122 (2019) 182701*

${}^7\text{Be} + d$  measured at  $E_{\text{cm}} \approx 0.2 - 1.5$  MeV, measured cross sections dominated by the  $(d,\alpha)$  channel towards which prior experiments mostly insensitive.

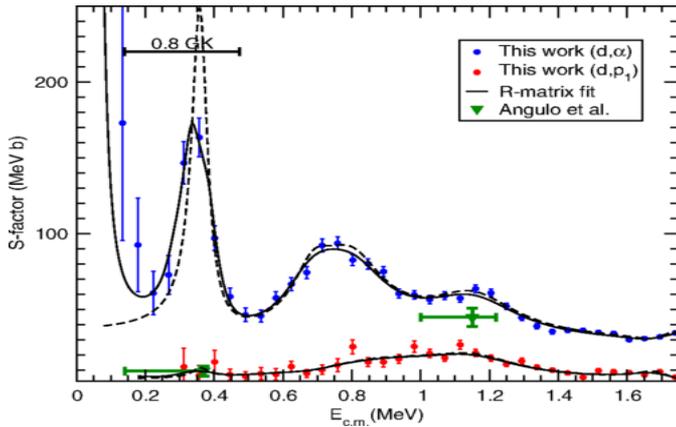
Rijal et al

Phys. Rev. Lett. 122 (2019) 182701

A new resonance at 0.36(5) MeV observed claims to **reduce the predicted abundance** of primordial  ${}^7\text{Li}$ .

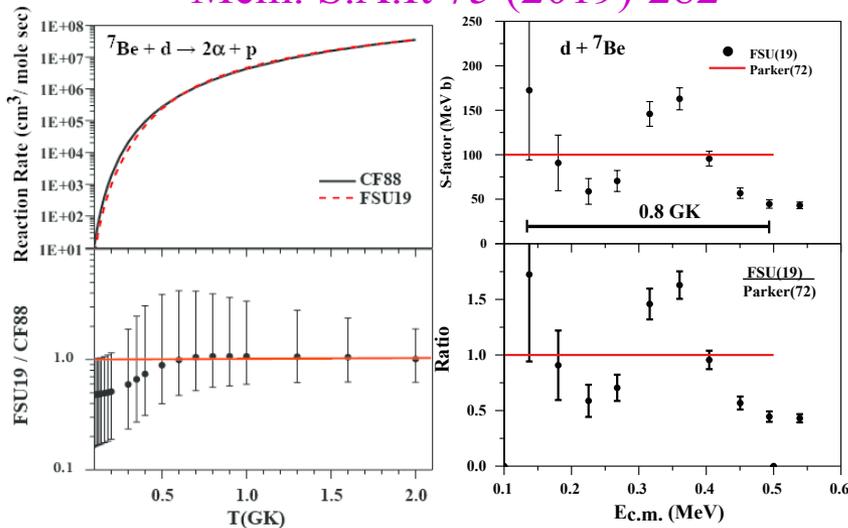
R-matrix analysis: **16.849(5) MeV, 5/2<sup>+</sup>** state in  ${}^9\text{B}$ ?  
Speculation:  ${}^9\text{B}$  resonance at **16.80 MeV?** *Scholl*

Additional experiments with improved statistics needed to reduce the **uncertainty in the resonance energy**.



Moshe Gai

Mem. S.A.It 75 (2019) 282



FSU rate uncertain by a factor of 10, due to uncertainty of the resonance energy. Old BBN  $d + {}^7\text{Be}$  rate (CF88) and Rijal (FSU) rates are **hardly different** if not identical over the entire range.

Since no state is known in  ${}^9\text{B}$  at the proposed "new resonance" energy of 16.85 MeV, resolving such a major systematical uncertainty is required.

**No reduction in  ${}^7\text{Li}$  abundance.**

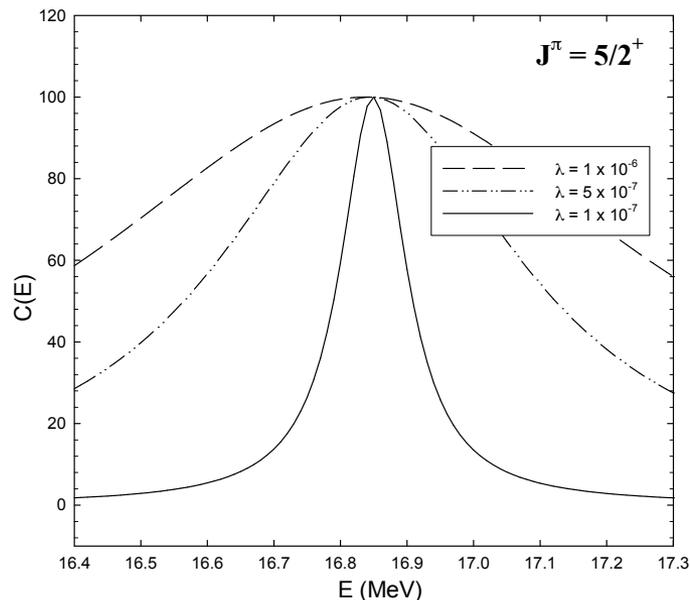
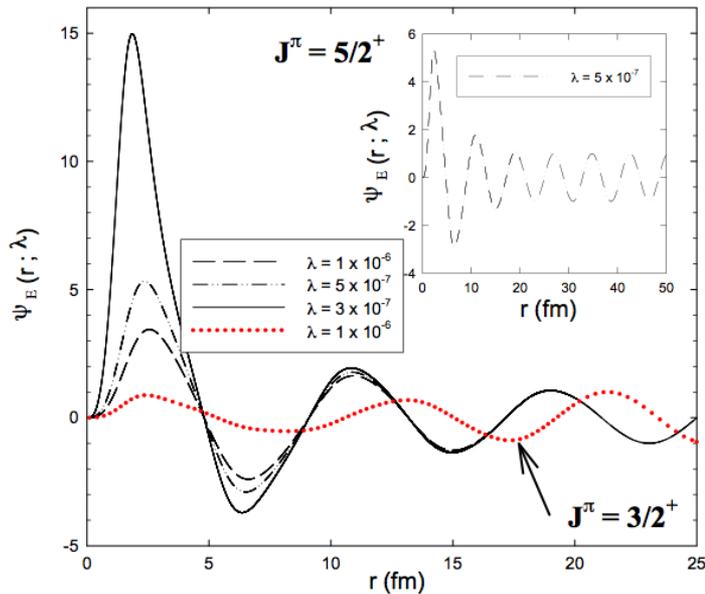
# Supersymmetric quantum mechanics to study the ${}^9\text{B}$ resonance

*S. K. Dutta, D. Gupta, S.K. Saha*

arXiv:2004.09105 [nucl-th] (2020)

Phys. Lett. B 776, 464 (2018)

J. Phys. G: Nucl. Part. Phys. 41, 095104 (2014)



Unstable/unbound systems, with very shallow potentials, pose serious numerical challenges in detecting **resonance states**. We could successfully circumvent this problem by using supersymmetric quantum mechanics.

This transforms the shallow well to a deep well-barrier isospectral potential, generating **resonance state wave-function**. The resonance state energies obtained were found to be in excellent agreement with the experimental values.

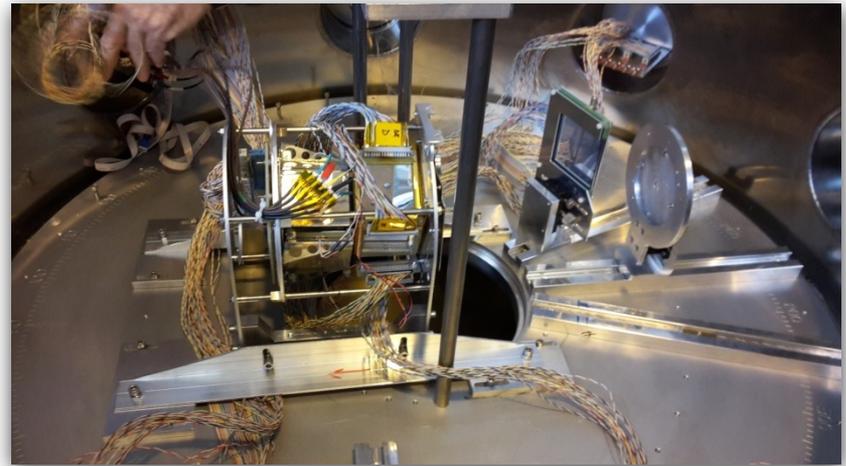
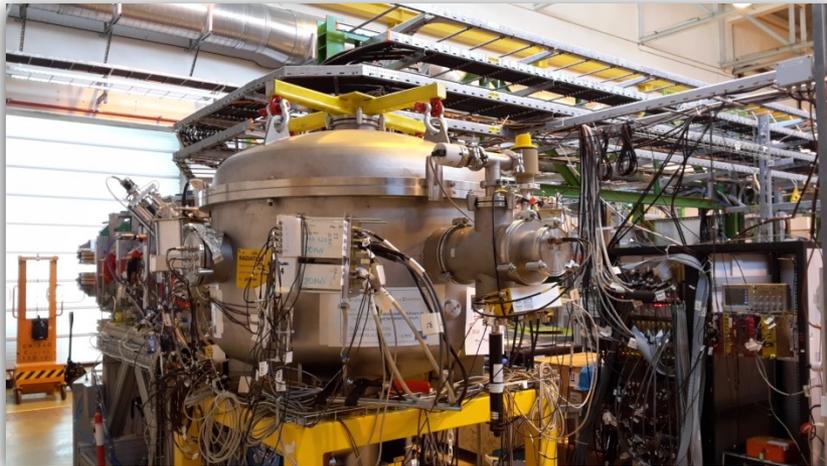
Resonance energy  $E_R = 16.84 \text{ MeV } (5/2^+)$

Width  $\Gamma = 69 \text{ keV}$

# Experiment IS 554 @



**5 MeV/u  $^7\text{Be}$**  on  $\text{CD}_2$  (15  $\mu\text{m}$ ),  $\text{CH}_2$  (15  $\mu\text{m}$ ) and  $^{208}\text{Pb}$  (1 mg/cm<sup>2</sup>) targets, beam intensity  $I \sim 5 \times 10^5$  pps



## Charge particle detector setup

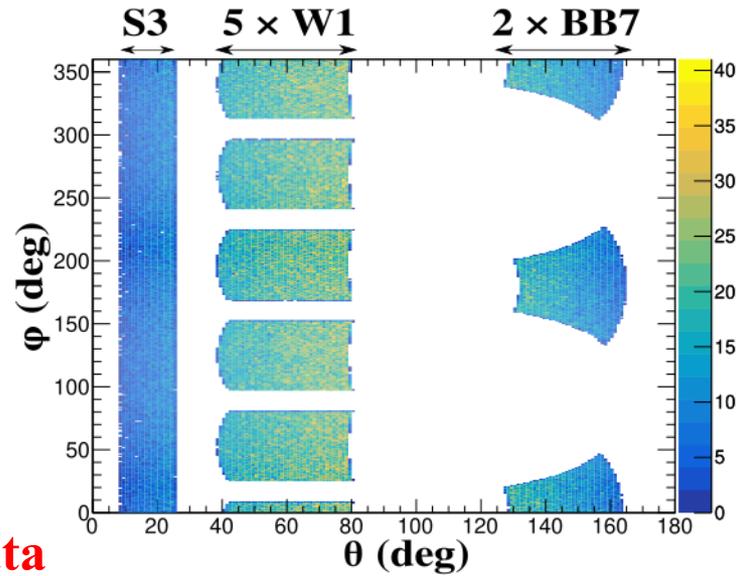
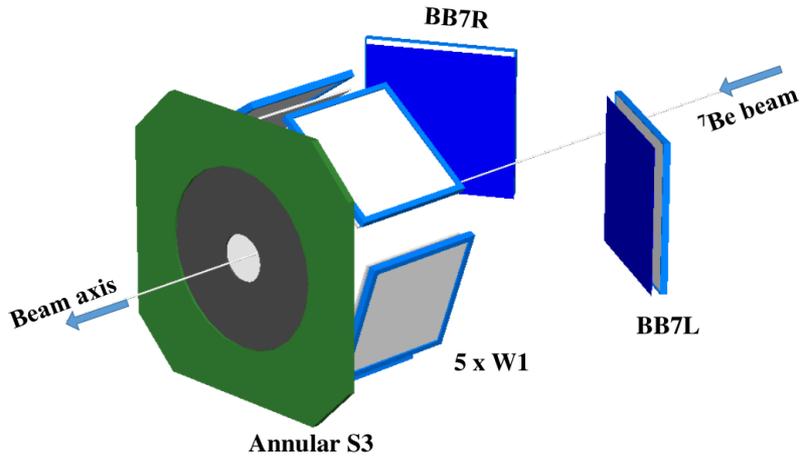
1 x S3 annular DSSD (24 x 32 strips, 1000  $\mu\text{m}$ ) covering front angles **8° – 25°**

5 x W1 DSSD (16 x 16 strips, 60  $\mu\text{m}$ ) in pentagon geometry covering angles **40° – 80°**

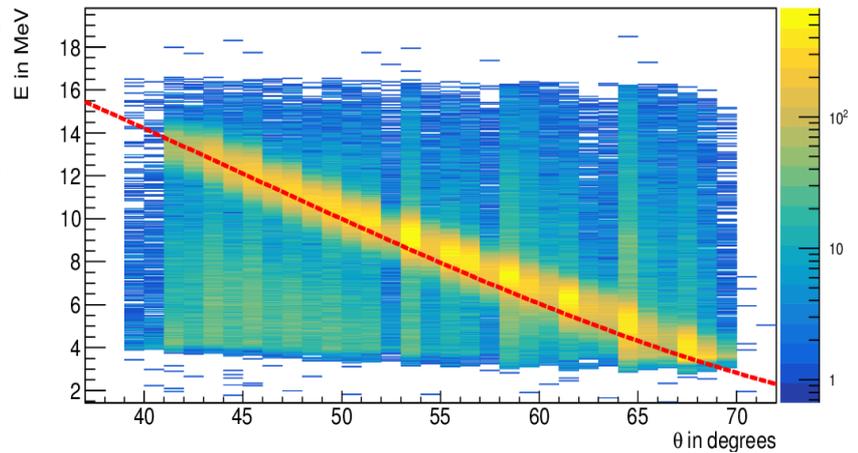
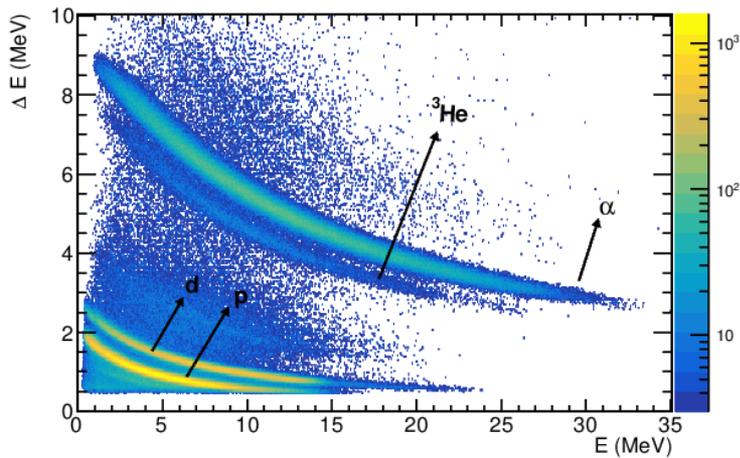
2 x BB7 DSSD (32 x 32 strips, 60  $\mu\text{m}$  and 140  $\mu\text{m}$ ) at backward angles **130° – 170°**

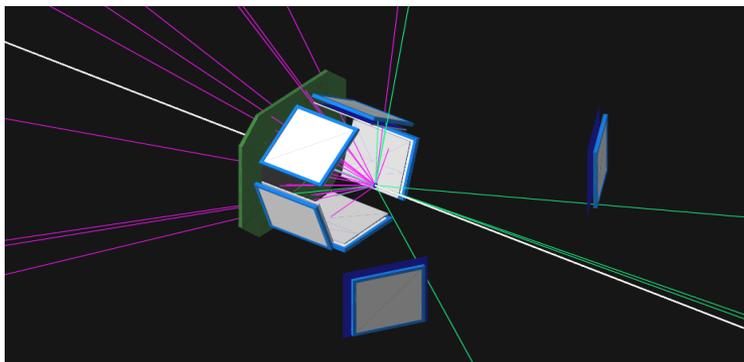
The W1 and BB7 DSSDs are backed by 1500  $\mu\text{m}$  thick unsegmented pads

# Experiment IS 554 @



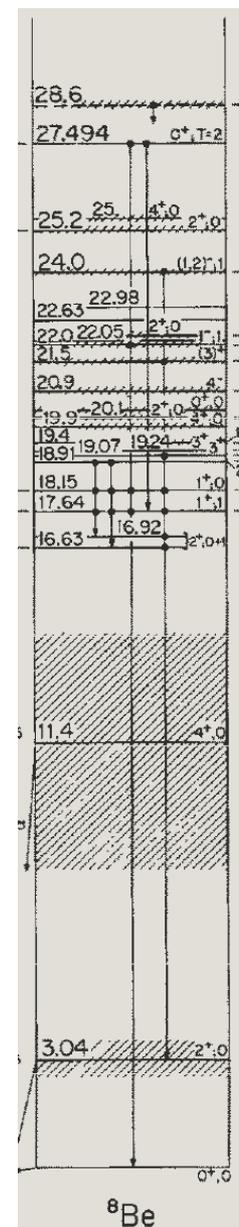
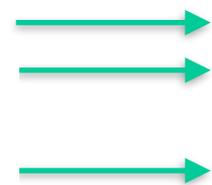
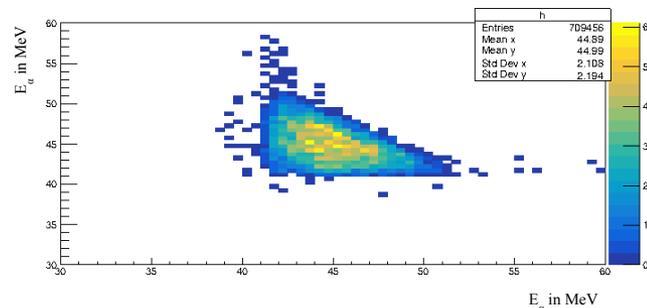
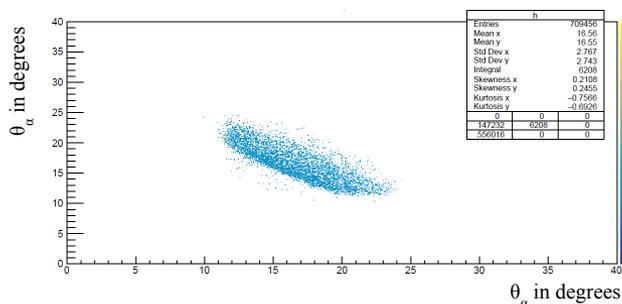
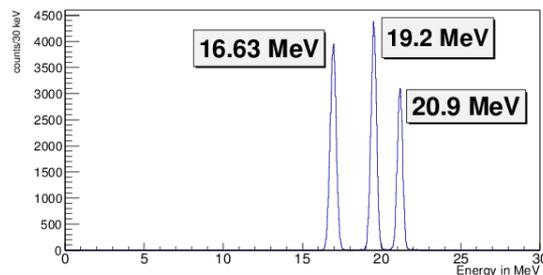
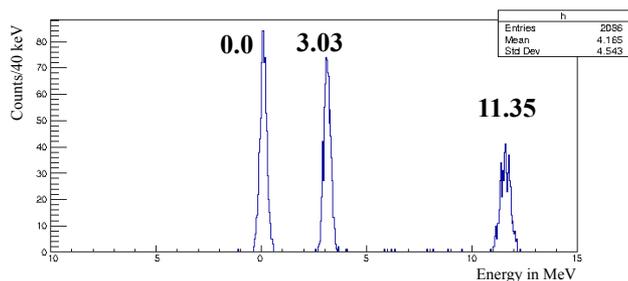
IS554 data





# ${}^7\text{Be}(d,p){}^8\text{Be}^*(2\alpha)$

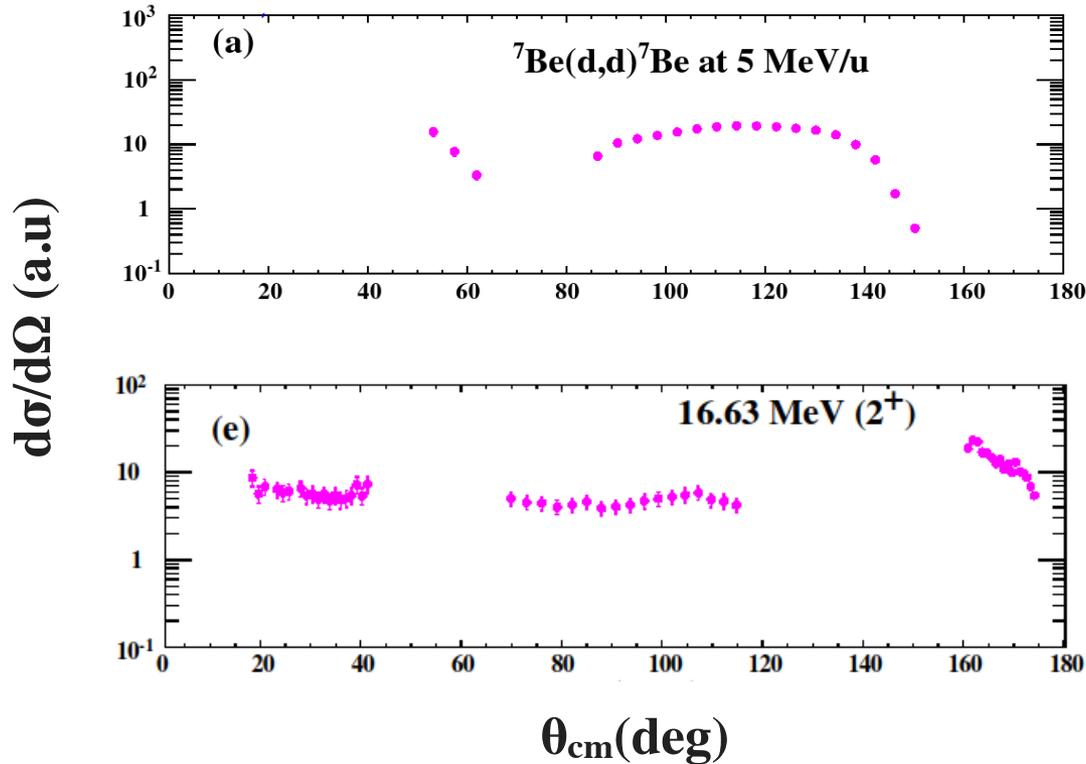
## Simulations



**Energy and angular correlations** of coincident alphas detected by the pentagon DSSDs. Simulations correspond to the correlation of the alphas emitted from the **16.63 MeV state of  ${}^8\text{Be}$** .



# Experiment IS 554 @



Angular distributions from the present work  
(a)  ${}^7\text{Be}(d,d){}^7\text{Be}$  elastic scattering

(b) the  ${}^7\text{Be}(d,p){}^8\text{Be}^*$  to the 16.63 MeV ( $2^+$ ) state of  ${}^8\text{Be}$ .

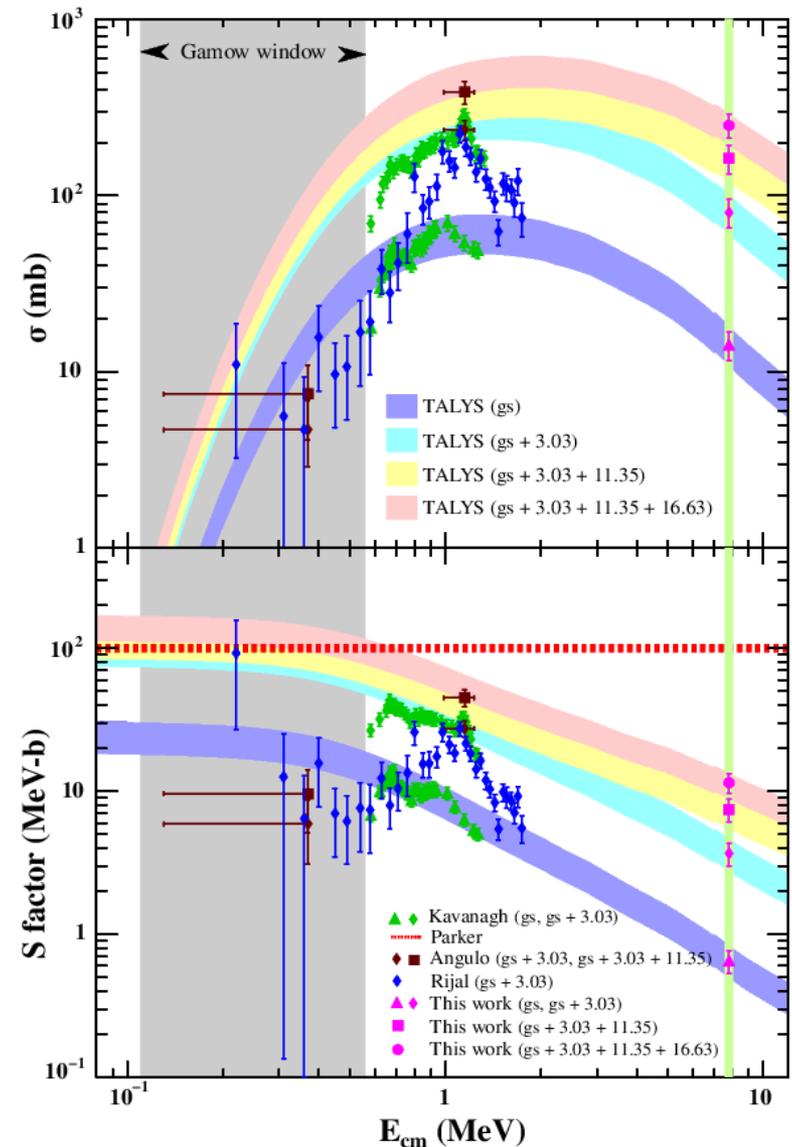
The present work studies the effects of all the resonances in the  ${}^7\text{Be}(d,p){}^8\text{Be}^*$  channel up to **16.63 MeV** and their contributions to the total cross section.

## ${}^7\text{Be}(d,p){}^8\text{Be}^*$ Excitation function

The bands are TALYS calculations normalized to present data at higher energy, giving an estimate of the contributions of the resonance excitations.

The existing data within Gamow window (0.11-0.56 MeV) has **large error bars in energy as well as cross sections**. Good agreement with data outside Gamow window.

The S factor due to contribution of gs, 3.03 and 11.35 MeV state agrees with Parker's estimate of 100 MeV-b. Addition of the **16.63 MeV** resonance leads to a maximum value of 167 MeV-b but is **not sufficient** to solve the Lithium abundance anomaly.



***Contribution of higher excited states in  ${}^7\text{Be}(d,p){}^8\text{Be}^*$  do not solve the Cosmological lithium problem***

# Outlook

## *Search for standard nuclear physics solution to the Cosmological Lithium problem*

Destruction of  ${}^7\text{Be}$  involving neutrons  ${}^7\text{Be}(n,p){}^7\text{Li}$ ,  ${}^7\text{Be}(n,\alpha){}^4\text{He}$  yield a decrease of the lithium abundance but insufficient to solve the anomaly. *Damone (2018), Barbagallo (2016)*  
The destruction channel  ${}^7\text{Be}(d,\alpha)\alpha p$  leads to speculation of a new resonance at **0.36 MeV** corresponding to the **16.8 MeV state of  ${}^9\text{B}$**  *Rijal (2019)*. No reduction of the abundance of  ${}^7\text{Li}$  can be deduced from the data *Gai (2020)*.

*IS 554* data includes the **(d,p)** channel for **higher excitations of  ${}^8\text{Be}$** . Including the **contribution of the 16.63 MeV state**, the maximum total S-factor inside the Gamow window is 167 MeV-b as compared to Parker's earlier estimate of 100 MeV-b. However, this value is also insufficient to alleviate the lithium discrepancy.

## **The cosmological lithium problem persists!**

We conclude from this work that our present understanding of nuclear physics is inadequate to solve this problem. All alternative physics and astronomical scenarios to solve the anomaly is still open.

*It would be interesting in future to see if the lithium problem truly points to new fundamental physics.*

# IS 554 collaboration



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