

Spinodal instability with varying criticality in holography

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arXiv:2012.15687 (JHEP)

HI Group - TH Retreat 2021

before MSCA at CERN TH:

PhD Technical University of Vienna



View from TU Freihaus to National Librational

before MSCA at CERN TH:

IGFAE

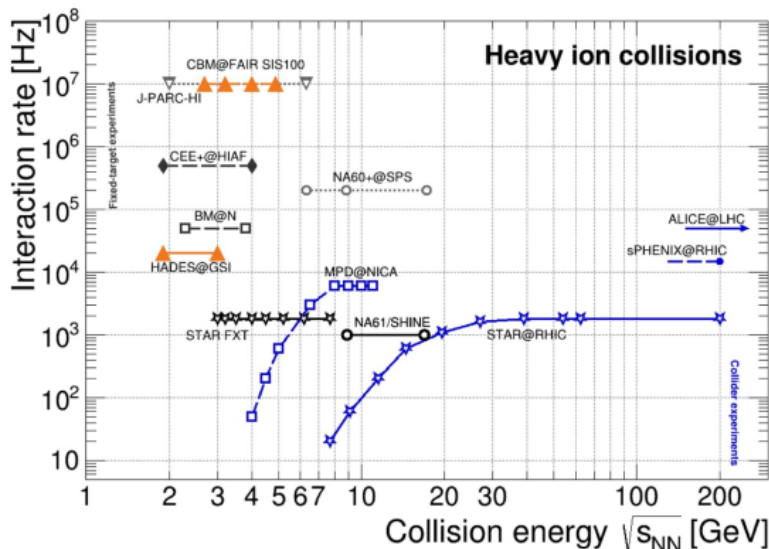


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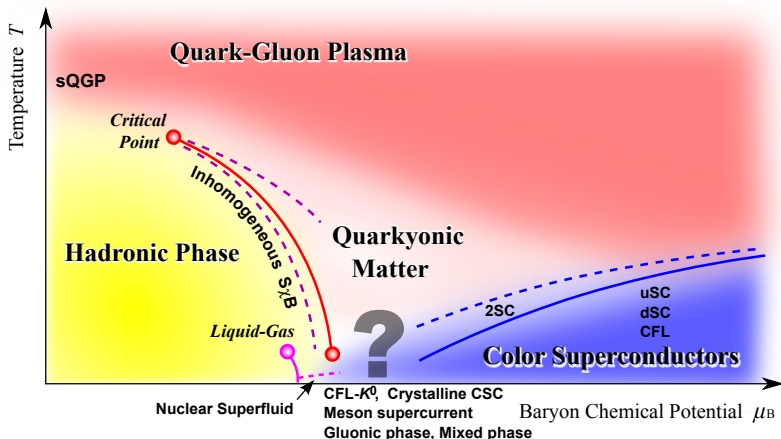




CBM Collaboration, EPJA 53 3 (2017) 60
 T. Galatyuk, NPA982 (2019), update (2021)

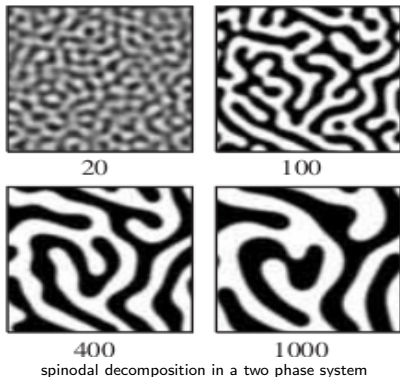
Motivation - many unknowns in the QCD phase diagram

The critical point (CP) is the endpoint of the 1st order phase transition between the Quark-Gluon Plasma and the hadrons:



Assumed QCD phase diagram [Fukushima, Hatsudo 2010]

spinodal instability:



[Onuki 1987]

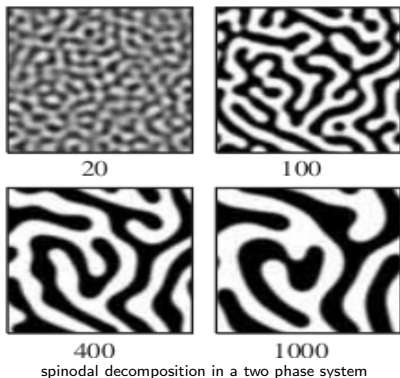
Gregory-Laflamme instability:



thin black ring pinching off
[Figueras, Kunesch, Tunyasuvunakool
2015]

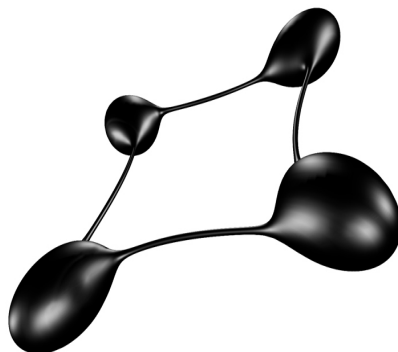
gauge/gravity correspondence:
bridge between physical phenomena in gauge theories and gravity.

spinodal instability:



[Onuki 1987]

Gregory-Laflamme instability:

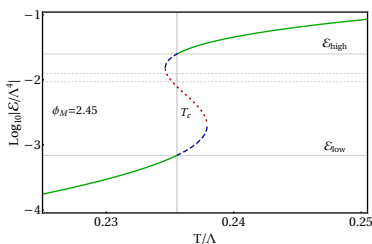
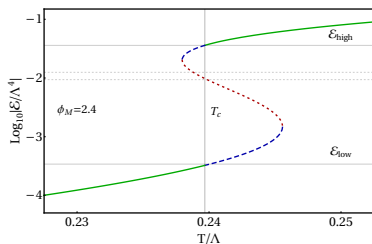
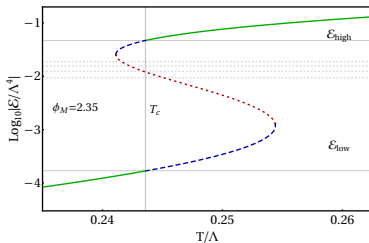
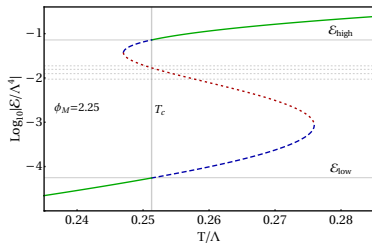


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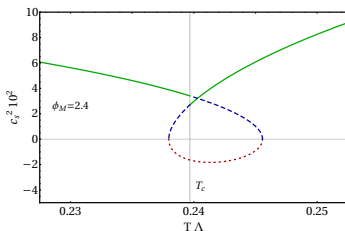
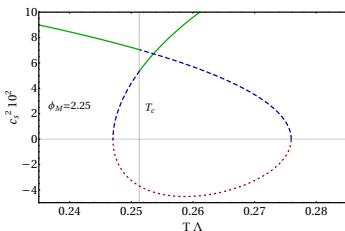
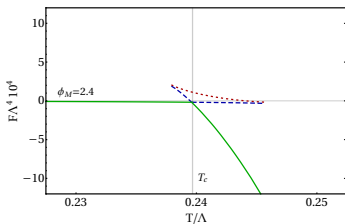
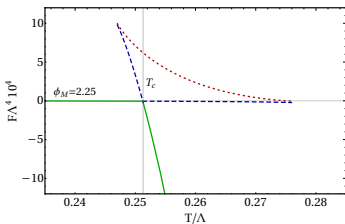
bridge between physical phenomena in gauge theories and gravity.

Bottom-up model: tuning criticality



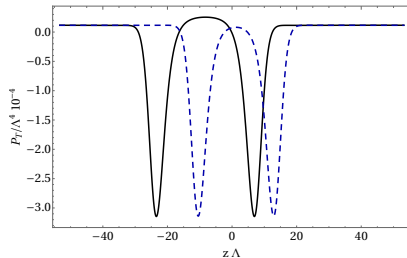
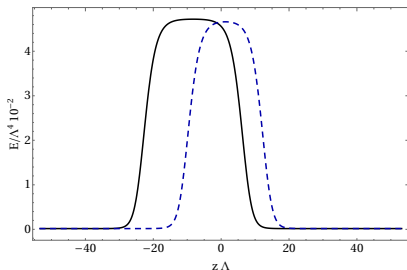
Energy density over T for theories with first order phase transition (blue) metastable (red) spinodal region (green) stable phases

Bottom-up model: tuning criticality II



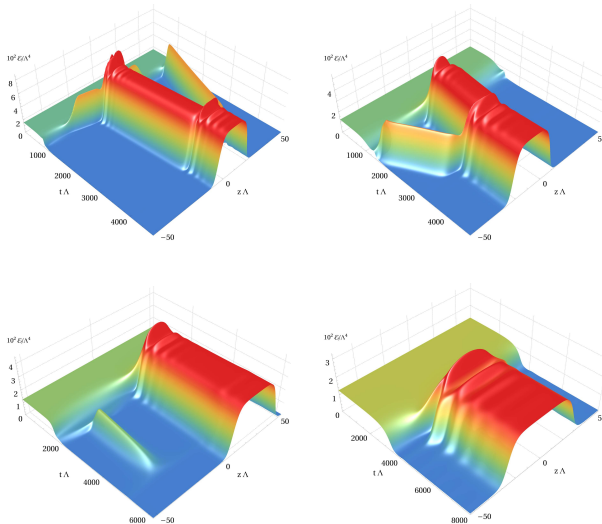
Thermodynamical properties of hard/soft transition:
(top row) Free energy over temperature
(bottom row) Speed of sound squared over temperature

New criterium plateau (full black) where $P_T \geq 0.8P_c$ versus peak (dashed blue) where $P_T < 0.8P_c$



Longitudinal profiles of (left) local energy density (right) transverse pressure P_T

Spinodal instability evolution



Phase separated final solution seen in the evolution of the local energy density of harder/softer phase transitions

Summary: spinodal instability approaching a critical point

Demonstration of the full dynamics of strongly coupled spinodal instability first order phase transitions.

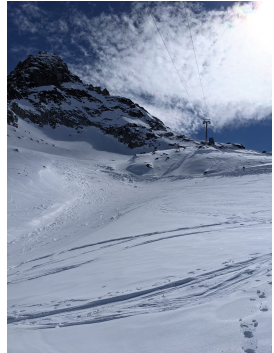
Near a critical point the interface between cold and hot stable phases, given by its width and surface tension, features

- a **wider phase separation**,
- a **smaller surface tension** and
- a **longer formation time**.

Hence the spinodal instability is easier to detect far from a critical point along the first order phase transition.

New discovery of an atypical setup with dissipation of a peak into a plateau.

Private outdoor activities



The formation time is defined as the time from the start of the evolution to when the system first reaches one of the two stable phases.

Formation time of the spinodal instability with varying criticality:

ϕ_M	2.25	2.35	2.4	2.45
$t_{\text{formation}}\Lambda$	880	1380	1725	2660

The formation time correlates with the strength of the first order phase transition: the strongest phase transition with the largest spinodal region shows the fastest formation time.

- 1) “Holographic approach of the spinodal instability to criticality”
By Maximilian Attems arXiv:2012.15687 [hep-th].
- 2) “Dynamics near a first order phase transition” By Loredana Bellantuono, Romuald A. Janik, Jakub Jankowski, Hesam Soltanpanahi arXiv:1906.00061 [hep-th], JHEP 10 (2019) 146.
- 3) “Dynamics of Phase Separation from Holography” By Maximilian Attems, Yago Bea, Jorge Casalderrey-Solana, David Mateos, Miguel Zilhao, arXiv:1905.12544 [hep-th], JHEP 01 (2020) 106.

Backup: setup bottom-up model

Dual field theory: 'mimics' a deformation of N=4 SYM with a dimension 3 operator O and source Λ as 'mass'

$$S_{\text{GaugeTheory}} = S_{\text{conformal}} + \int d^4x \Lambda O$$

Einstein-Hilbert action coupled to a scalar with non-trivial potential in five-dimensional bottom-up model:

$$S = \frac{2}{\kappa_5^2} \int d^5x \sqrt{-g} \left[\frac{1}{4} \mathcal{R} - \frac{1}{2} (\nabla\phi)^2 - V(\phi) \right]$$

We derive the potential from the superpotential [Bianchi, Freedman, Skenderis 2002]: $V(\phi) = -\frac{4}{3} W(\phi)^2 + \frac{1}{2} W'(\phi)^2$

$$\ell^2 W(\phi) = -\frac{3}{2} - \frac{\phi^2}{2} - \frac{\phi^4}{4\phi_M^2}$$

using a single parameter with critical value $\phi_M = 2.521$

$$\ell^2 V(\phi) = -3 - \frac{3\phi^2}{2} - \frac{\phi^4}{3} - \frac{\phi^6}{3\phi_M^2} + \frac{\phi^6}{2\phi_M^4} - \frac{\phi^8}{12\phi_M^4}$$

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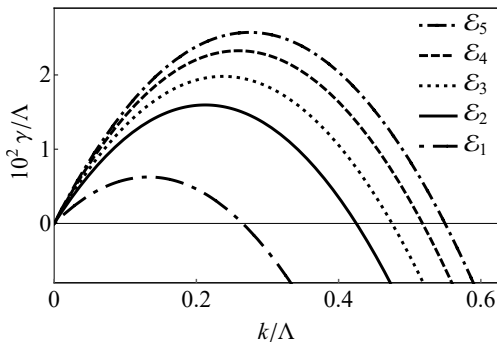
Growth rate

$$\gamma(k) = |c_s| k - \Gamma k^2$$

sound attenuation
constant

$$\Gamma = \frac{1}{2T} \left(\frac{4\eta}{3s} + \frac{\zeta}{s} \right)$$

raise dynamical instability



four generic stages:

- 1 linear stage
- 2 reshaping
- 3 merger
- 4 final:
static + phase-separated

Homogeneous initial state in periodic unstable region, initial instability triggered by $n = 3$

endstate:

phase-separated configuration

conjecture:

all static, non-phase separated configurations are dynamically unstable

$$10^2 \varepsilon / \Lambda^4$$

