



# **Evidence of multistage charge collection in Si irradiated detectors operated as monitors of intensive fragmented proton beams**

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# Outline

- Motivation
- Experiment and experimental results
- Multistage process with avalanche multiplication
- Data treatment and simulations
- Third stage of charge collection
- Electric field profile
- Internal charge gain and hole trapping time constant degradation rate
- Conclusions

# Motivation

Motivation of semiconductor detector application for monitoring of proton beam loss in LHC machine:

- installation of compact detectors in the vicinity of superconductive magnets, increase in sensitivity of monitoring the particle loss of proton beam, **initiated by CERN BE-BI** (Beam Div., Beam Instrumentation)

The goal of the presented study:

- further analysis of the results of the unique *in situ* irradiation test of Si detectors;
- new findings in detector and radiation physics at low T

# Experiment

***In situ* irradiation tests** (the only method of experimental study at 1.9 K)

2012 – *in situ* irradiation test 1 at 1.9K, 6 weeks of operation:

**Si p<sup>+</sup>-n-n<sup>+</sup> detectors**;  $\rho \sim 10\text{-}15 \text{ k}\Omega$ , 300  $\mu\text{m}$ ; processed in Russia

## **Irradiation:**

23 GeV protons beam fragmented into 400 ms spills

Beam intensity  $1.3 \times 10^{11}$  p/cm<sup>2</sup> per 400 ms spill ( $\sim 10^{10}$  p/s on detectors)

- ◆ Fluence to  $1 \times 10^{15}$  -  $1 \times 10^{16}$  p/cm<sup>2</sup>
- ◆ Beam position monitoring (BPM + Si beam telescopes)
- ◆ **Current pulse response, TCT**, LeCroy WavePro, 3 GHz bandwidth, 630 nm laser, width 45 ps

Laser illumination of the n<sup>+</sup> contact



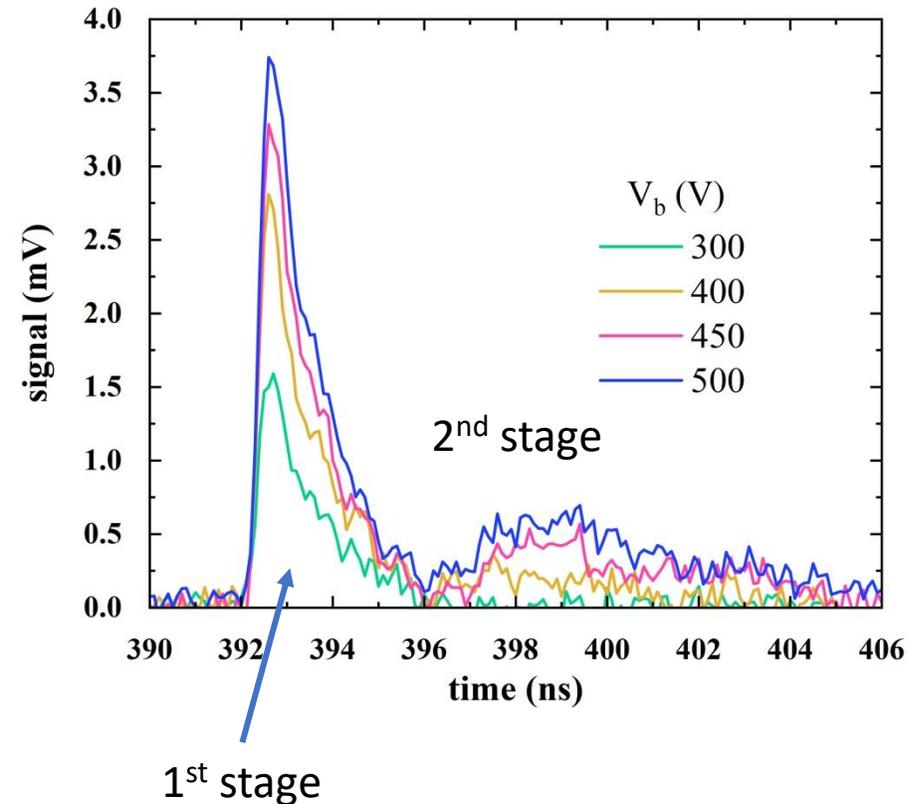
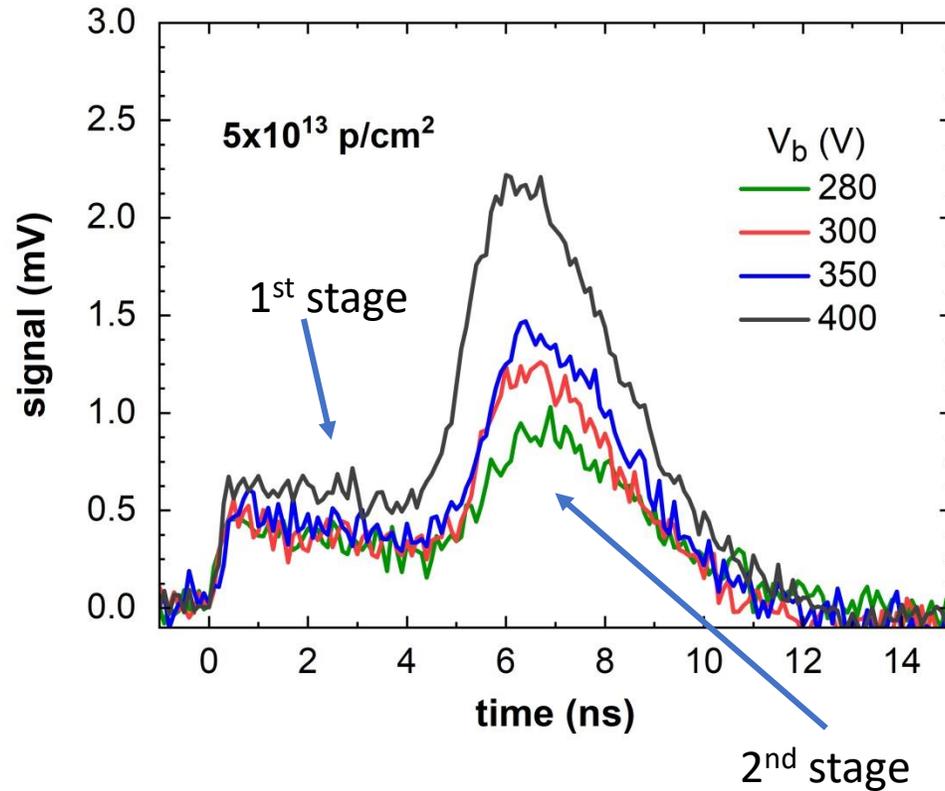
# Experimental results

No reflection as seen at lower F  
At higher F – details of fine structure

irradiation dose

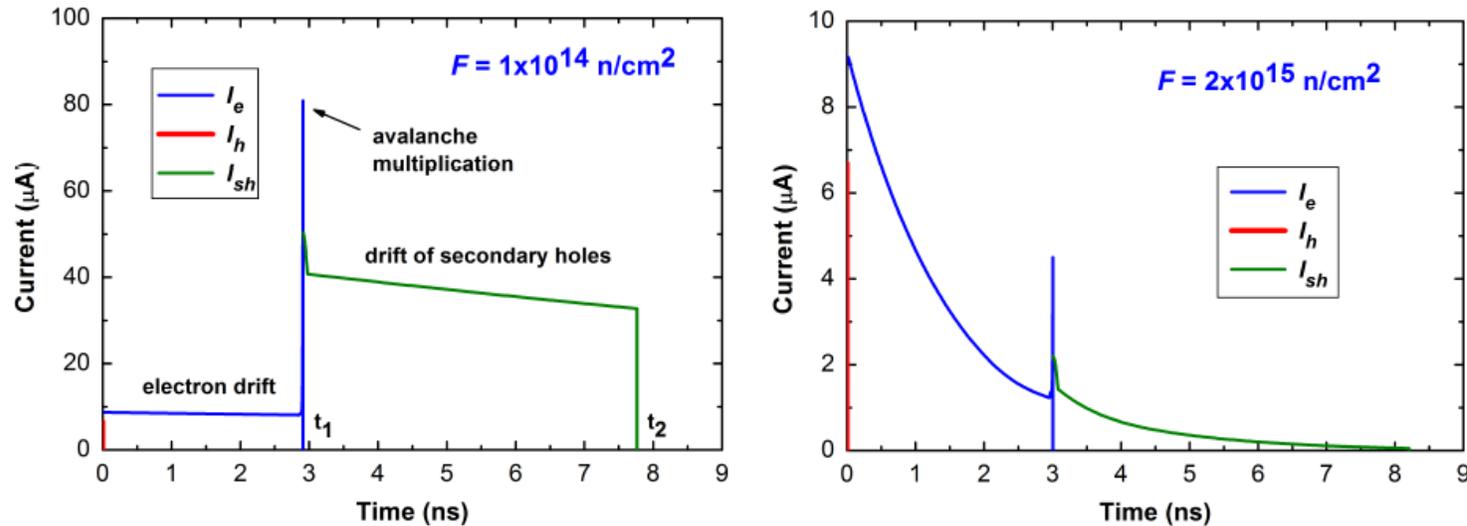
$5 \times 10^{13} \text{ p/cm}^2$

$2.7 \times 10^{14} \text{ p/cm}^2$

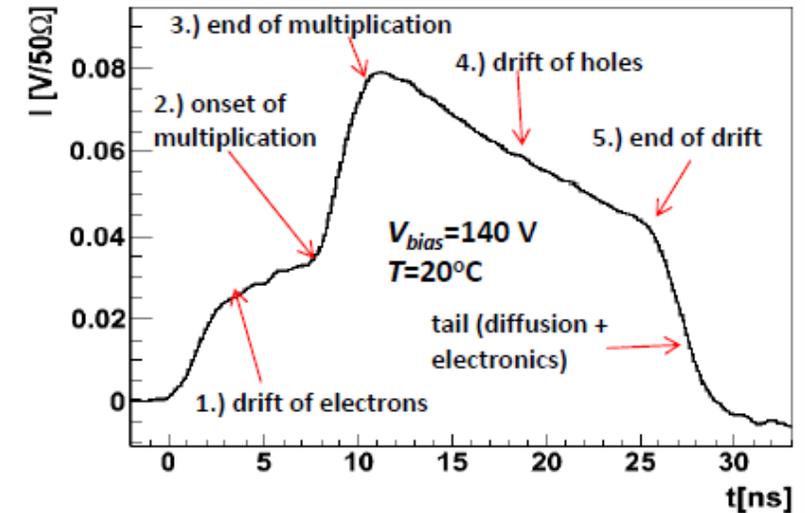


# Motivation to consider multistage process with avalanche multiplication

Simulation: irradiated LGAD



Experiment: nonirradiated LGAD

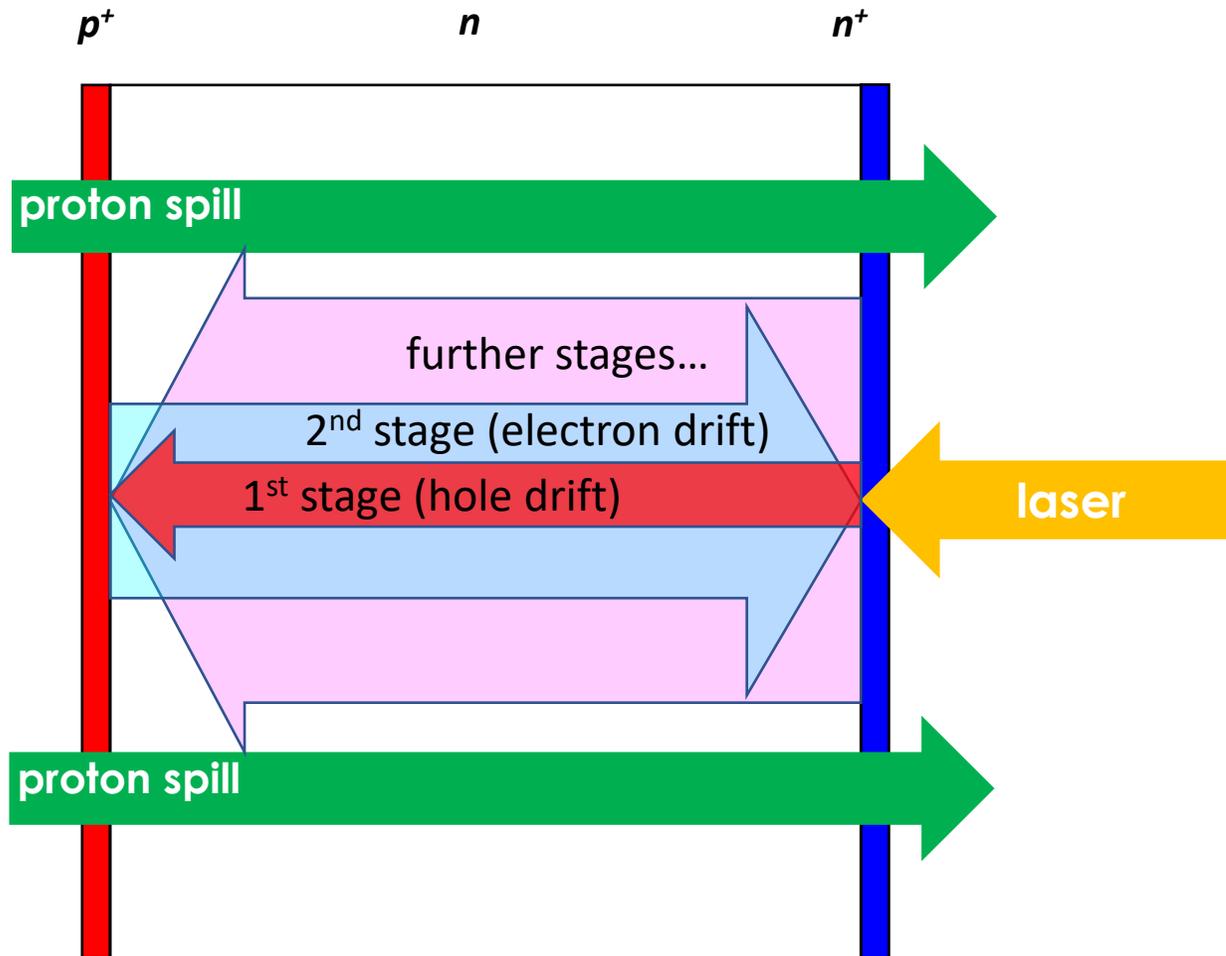


G. Kramerberger, et al., 2015 JINST 10 P07006

E. Verbitskaya, et al.,  
2016 JINST 11 P12012

Charge multiplication in structure **without** built-in layer

# Multistage process of charge collection



## Two stages of drift:

### 1<sup>st</sup> stage

- initial hole drift from n<sup>+</sup> to p<sup>+</sup>
- impact ionization and charge multiplication initiated by holes at p<sup>+</sup>

### 2<sup>nd</sup> stage

- electron drift from p<sup>+</sup> to n<sup>+</sup>

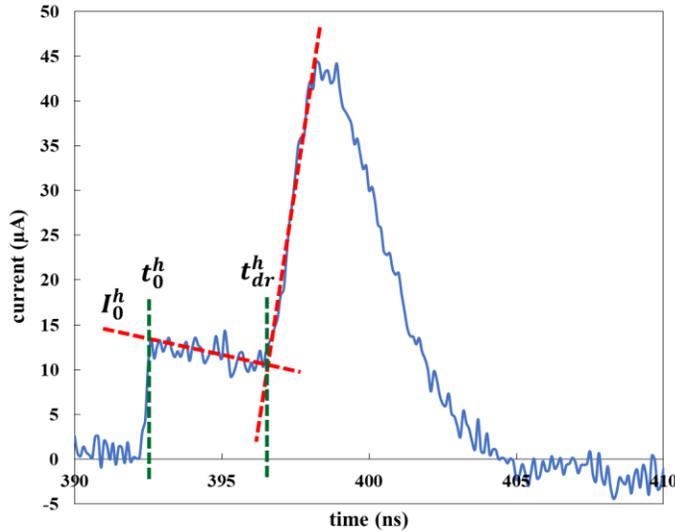
If electric field near n<sup>+</sup> contact is **high**: electrons initiate multiplication and **3<sup>rd</sup> stage** of charge collection (hole drift again)

# Data treatment and simulations

LOWER DOSE:  $5 \times 10^{13}$  p/cm<sup>2</sup>

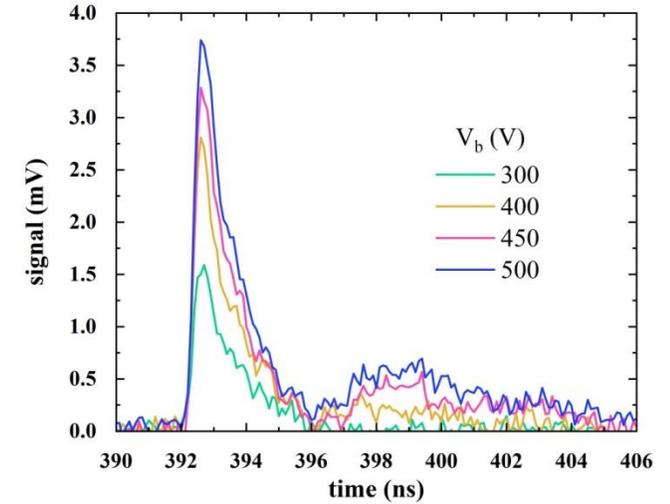
HIGHER DOSE:  $2.7 \times 10^{14}$  p/cm<sup>2</sup>

1<sup>st</sup> stage approximation



Signal construction

$$s(t) = \begin{cases} I_0^h \exp\left(-\frac{t}{\tau_{tr}^h}\right), & 0 < t < t_{dr}^h \\ \dots & \dots \\ \dots & \dots \end{cases}$$



$$Q_c^h = \int_{t_0^h}^{t_{dr}^h} Q_0 \frac{v_s}{d} \exp\left(-\frac{t}{\tau_{tr}^h}\right) dt = I_0^h \tau_{tr}^h \left[1 - \exp\left(-\frac{t_{dr}^h}{\tau_{tr}^h}\right)\right]$$

difficult to perform exp. fit

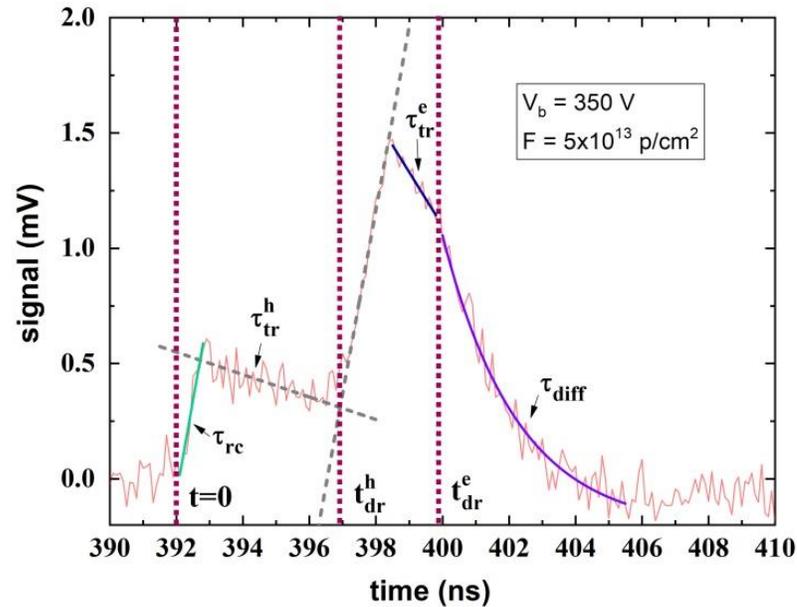
ONLY 1<sup>st</sup> stage approximation

+ accurate exp. fit  
(direct  $\tau_{tr}^h$  extraction)

# Data treatment and simulations

2<sup>nd</sup> stage and pulse tail

**complicated** because of carrier cloud size



Beer-Lambert law

$$p(x, 0) = p_0 e^{-\alpha x}$$

Hole distribution during drift

$$p(x, t) = \frac{1}{2} p_0 \exp\left(\frac{\alpha^2 \sigma^2}{2} - \alpha x\right) \operatorname{erfc}\left(\frac{\alpha \sigma^2 - 2x}{2\sigma}\right)$$

$$\sigma = \sqrt{2Dt}$$

Signal construction

$$s(t) = \begin{cases} I_0^h e^{-\frac{t}{\tau_{tr}^h}}, & 0 < t \leq t_{dr}^h \\ s(t = t_{dr}^h) [1 + G_Q f(t)] e^{-\frac{t - t_{dr}^h}{\tau_{tr}^e}}, & t_{dr}^h < t \leq t_{dr}^e \\ s(t = t_{dr}^e) e^{-\frac{t - t_{dr}^e}{\tau_{diff}}}, & t > t_{dr}^e \end{cases}$$

$f(t)$  - the function defined by the change of charge in the detector determined by integration of  $p(x, t)$

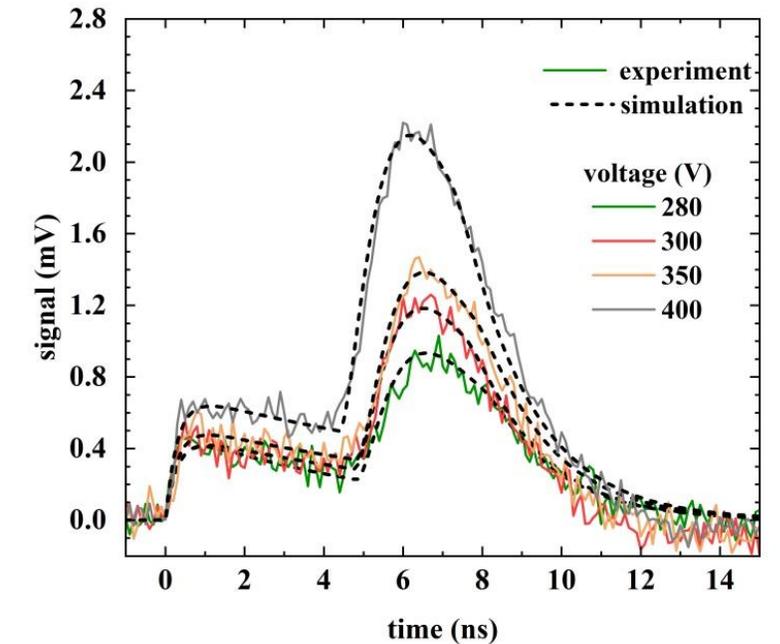
+ RC distortion

$$y(t) = \int_0^t s(t') h(t - t') dt'$$

taking into account

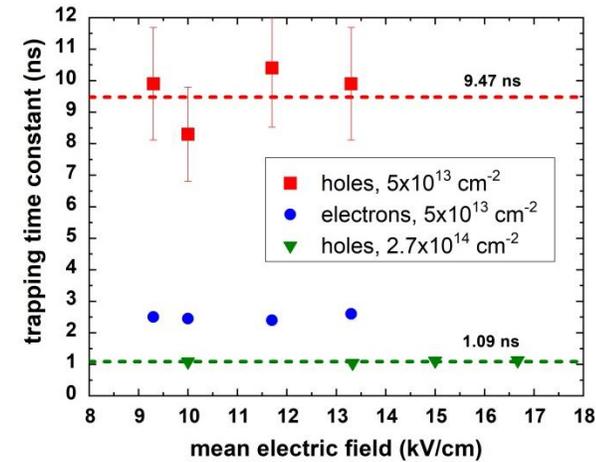
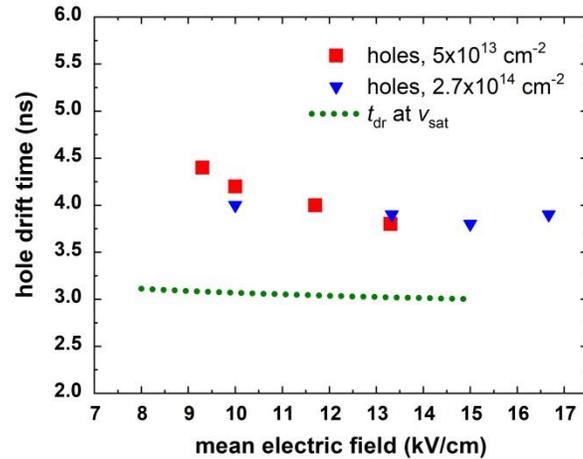
- initial nonuniformity
- diffusion spreading

dash lines: simulated current responses



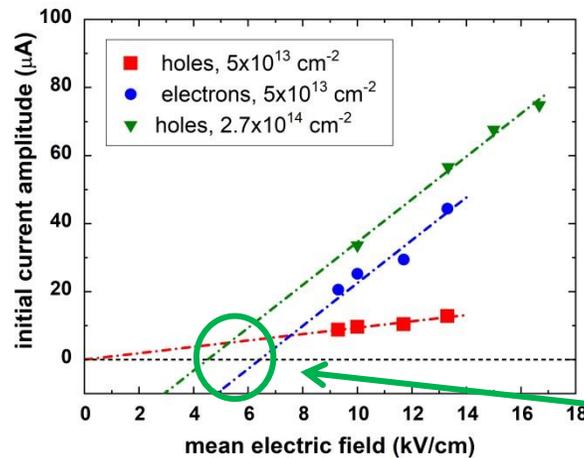
# Carrier transport characteristics

bias voltage  $\rightarrow$  mean electric field  $\langle E \rangle = \frac{V_b}{d}$

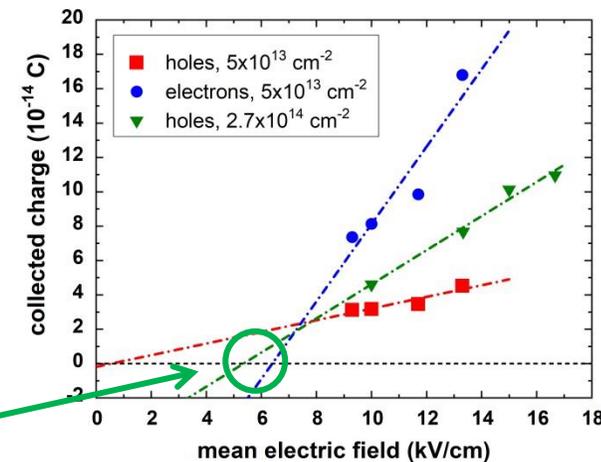


more accurate for higher dose because of exp. fit

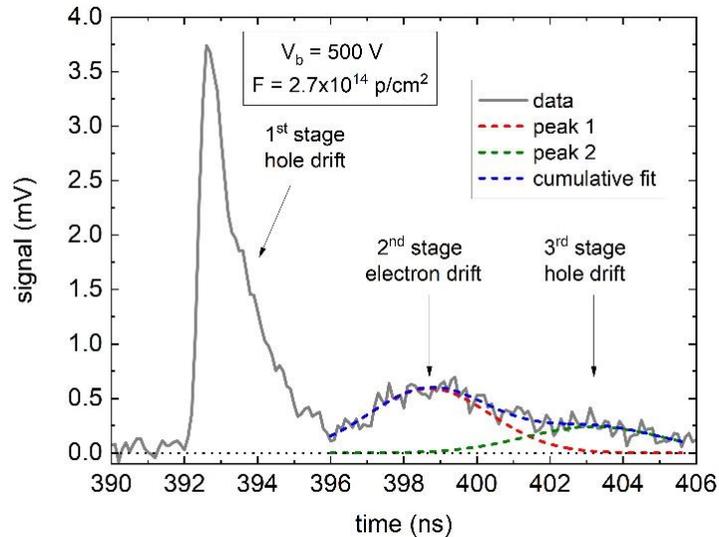
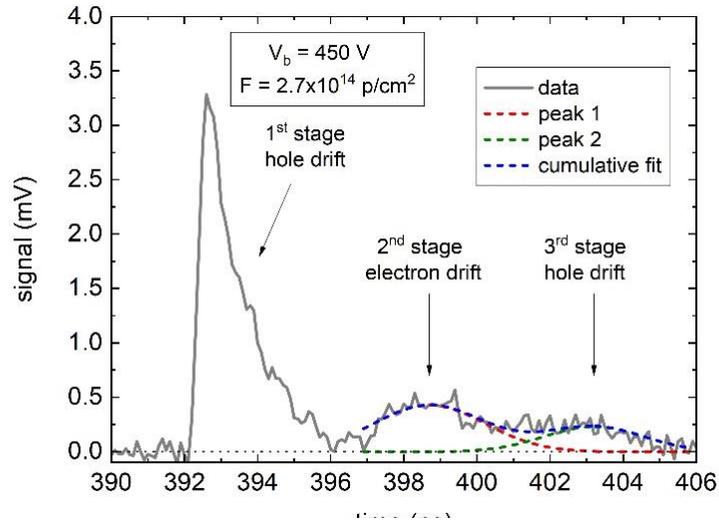
hole parameters at  $F = 5 \times 10^{13} \text{ p/cm}^2$  cross (0, 0)



threshold



# Third stage of charge collection



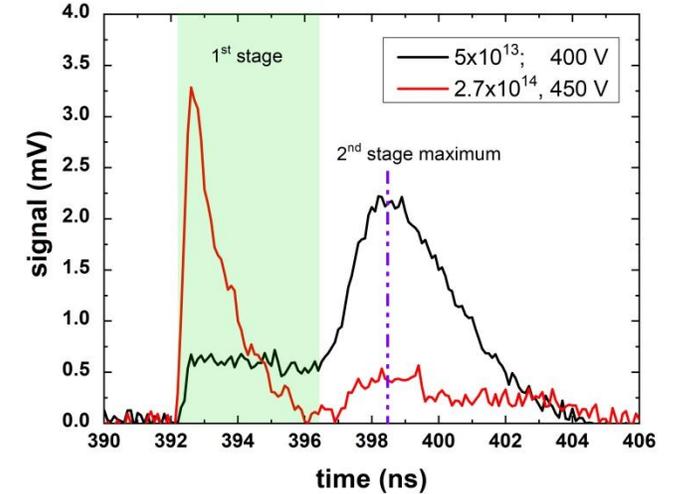
3<sup>rd</sup> stage is the **drift of holes** after impact ionization produced by electrons near n<sup>+</sup> contact

Sum of Gaussian fit

$$y(t) = y_0 + \frac{A}{\sigma\sqrt{2\pi}} \exp\left(-\frac{(t - t_m)^2}{2\sigma^2}\right)$$

the result of the fit

Dose comparison

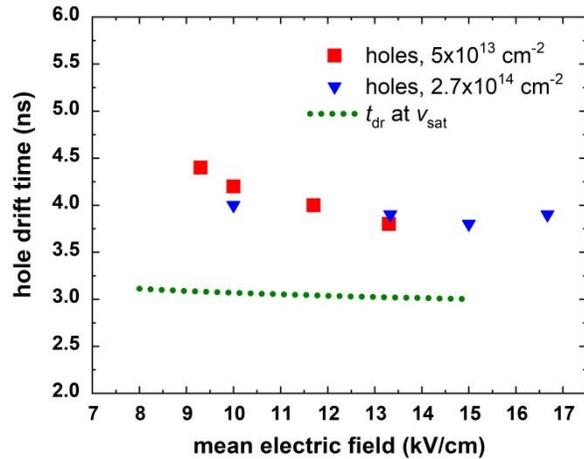


Parameter	450 V		500 V	
	S2	S3	S2	S3
$t_m$ (ns)	398.7	403.2	398.7	403.2
$\sigma$ (ns)	$3.0 \pm 0.2$	$2.9 \pm 0.3$	$3.3 \pm 0.1$	$3.4 \pm 0.4$
$A$ ( $10^{-14}\text{ C}$ )	$3.3 \pm 0.2$	$1.6 \pm 0.1$	$5.0 \pm 0.1$	$2.2 \pm 0.2$

The time between maxima **correlates** with  $t_{dr}$  with drift **velocity below saturated**, but the **current maximum** is determined by **charge** drifting inside the detector

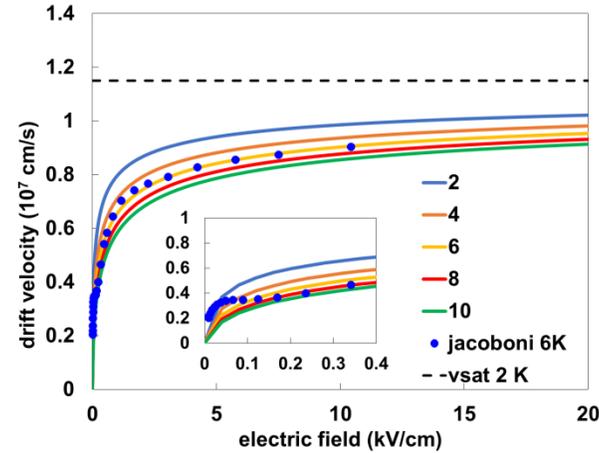
# Electric field profile: low field

result for hole drift time

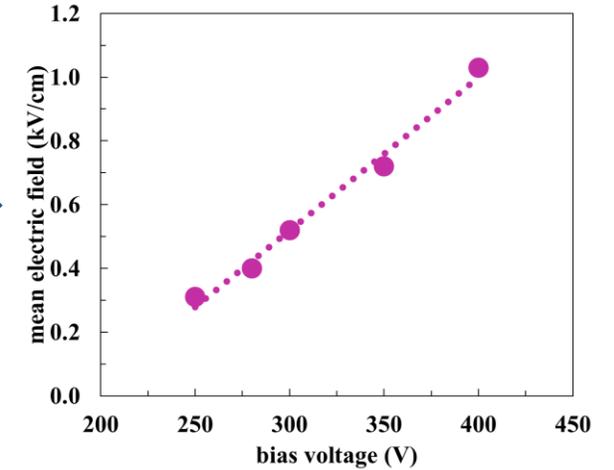


$$\langle v_{dr} \rangle = \frac{d}{t_{dr}}$$

drift velocity vs electric field

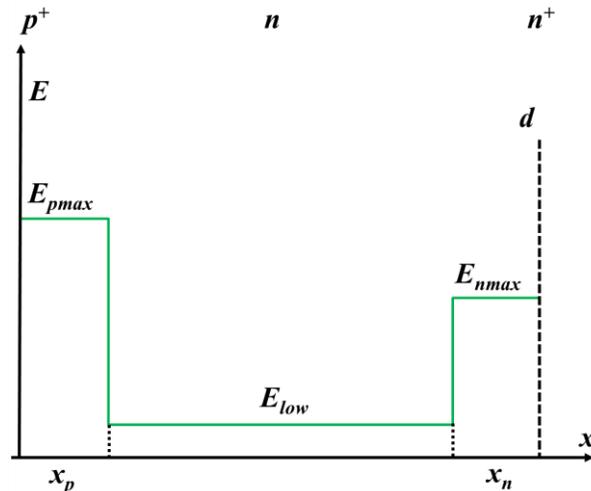


low electric field *in between*



electric field distribution

$x_p$  and  $x_n$  are thin  
( $\leq 10 \mu\text{m}$ )

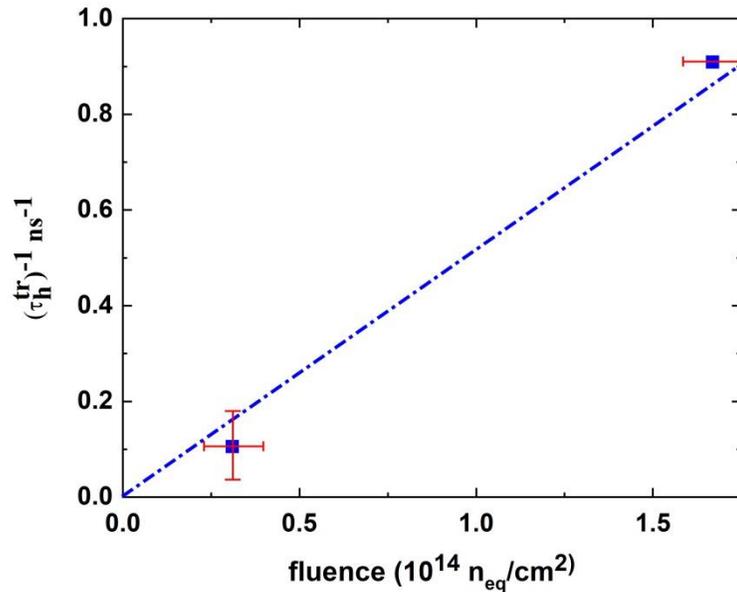


1<sup>st</sup> stage decay  
**controlled by  
hole trapping**

$$E_{pmax} > E_{nmax}$$

# Internal charge gain and $\tau_h$ degradation rate

trapping time constant vs eq. fluence



1.9 K:

$$\beta_h = 3.65 \times 10^{-15} \text{ cm}^2/\text{ns}$$

at -10° C:

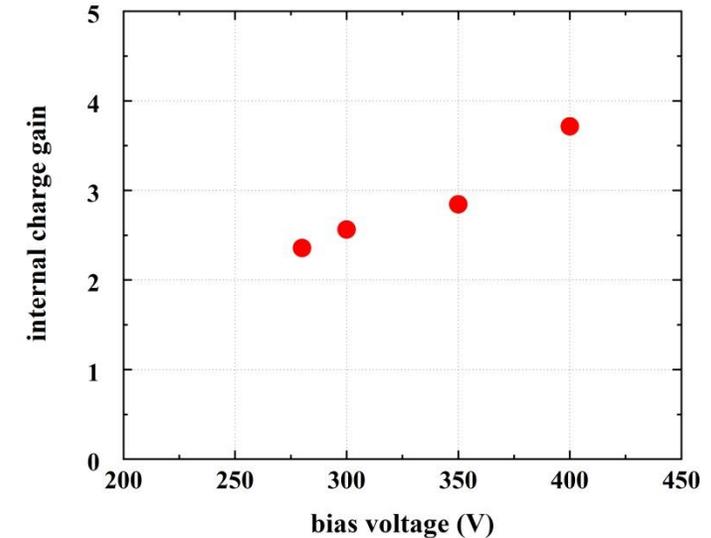
$$\beta_h = 7.7 \times 10^{-16} \text{ cm}^2/\text{ns}$$

(NIM A 481 (2002) 297)

$$\frac{1}{\tau_h} = \beta_h F$$

Internal charge gain

$$G_Q = \frac{Q_{col}^e + Q_{col}^h}{Q_0}$$



negative feedback due to carrier trapping

# Conclusions

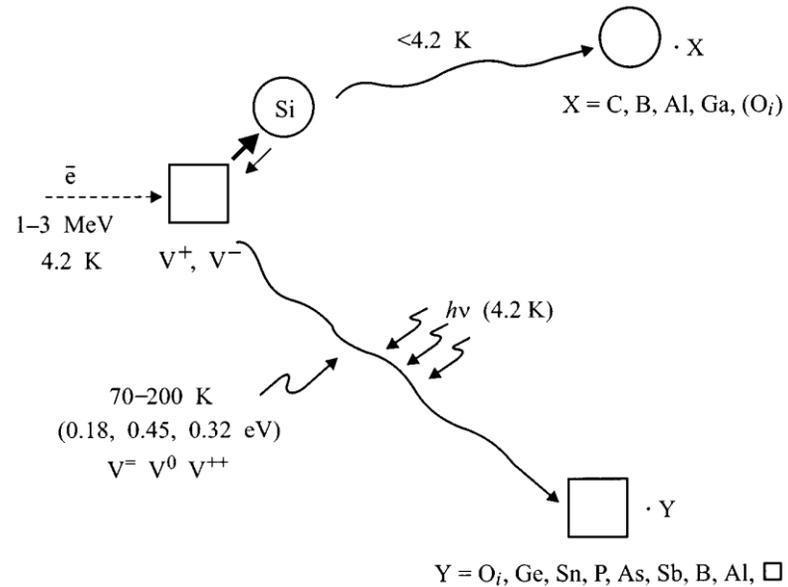
- In the unique experiment at 1.9 K, Si p<sup>+</sup>/n/n<sup>+</sup> detector maintains operation as a regular diode structure with highly doped contacts.
- The effect of sequential stages of charge multiplication was found. The third stage of charge collection was observed.
- The developed procedure of TCT data treatment for irradiated Si detectors with avalanche multiplication allowed extraction of full set of carrier transport parameters and internal charge gain.
- Comparison of the trapping time constant degradation rate at 1.9 K and room temperature showed fivefold increase of  $\beta_h$  at 1.9 K.

## **New questions have arisen:**

- ❖ conservation of the low-field region after polarization,
- ❖ relaxation of current response after spill escape.

**Thank you for your attention!**

# Why $E_{pmax} > E_{nmax}$ ?



At such a low temperature defect formation is not a well-known mechanism. There is J.D. Watkins hypothesis that mobility of initial defects is different and at  $T < 4.2$  K more donors (hole traps) are forming, while at a RT more acceptors (electron traps) are forming.

J.D. Watkins