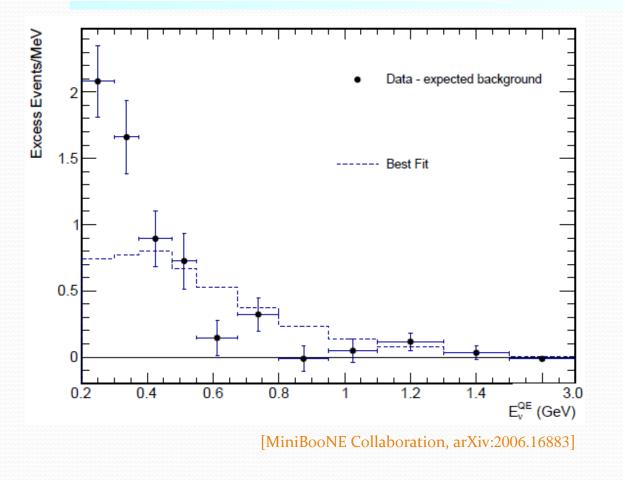
# A DM Interpretation of the MiniBooNE Excess and Its Implications in CEvNS Experiments



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In collaboration with Bhaskar Dutta, Adrian Thompson, Remington Thornton, Richard Van de Water, to appear soon

### **MiniBooNE Low Energy Excess**



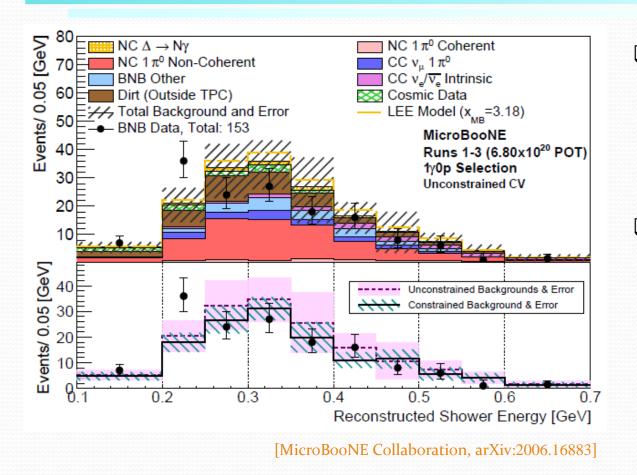
- $\Box$  An observational motivation of new physics (4.8 $\sigma$ )
- Numerous explanations, mostly involving neutrino-sector
   new physics [Karagiorgi, Djurcic, Conrad, Shaevitz, Sorel (2009); Collin,
   Arguelles, Conrad, Shaevitz (2016); Giunti, Lavender (2011); Gariazzo, Giunti,
   Lavender, Li (2017); Kopp, Maltoni, Schwetz (2011); Doring, Pas, Sicking, Weiler
   (2018); Dutta, Ghosh, Li (2020), and many more]
- Interpretations with dark-sector new physics (mostly coming from neutral mesons) less favored because off-target mode measurements [MiniBooNE DM Collaboration, arXiv:1807.06137] show null signal.

We propose an idea of reinstating the dark-sector scenarios

for the MB excess (see also Adrian Thompson's talk).

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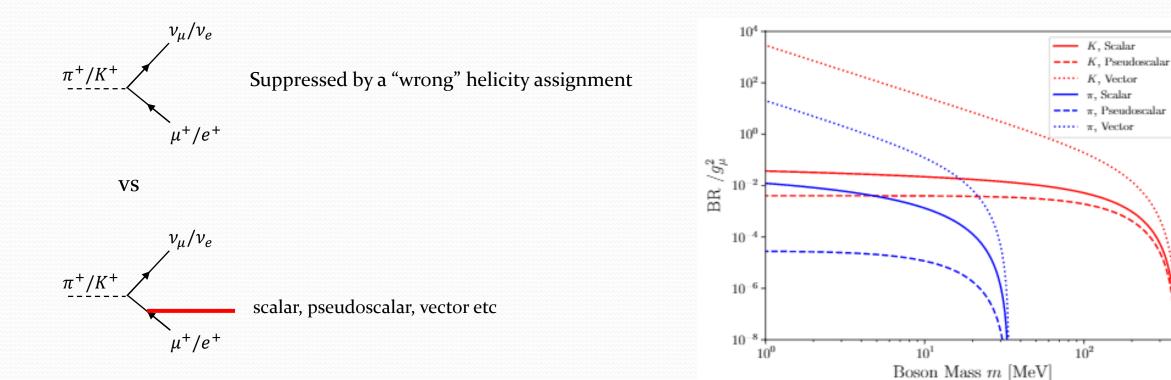
#### **Recent MicroBooNE Result and MiniBooNE Excess**



□ The recent MicroBooNE result constrains NC  $\Delta \rightarrow N\gamma$ event rates more stringently, supporting that the MB excess requires a new physics interpretation!

 The MicroBooNE data may not be sensitive enough to the MB excess, yet, because of (~3 times) smaller
 POTs, (~8 times) smaller detector volume, (~5 – 6 times) smaller efficiency.

#### **Dark-Sector Particles Sourced by Charged Mesons**

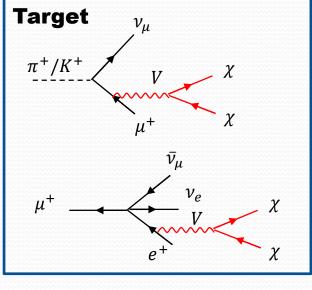


By adding the third particle, the helicity suppression can be evaded, i.e., 3-body decays can be hugely enhanced. The decay to a massive vector is even more enhanced due to the longitudinal polarization. [e.g., Carlson, Rislow, arXiv:1206.3587]

## A Dark Matter Interpretation for the MiniBooNE Excess

 $-\mathcal{L}_{V,\text{int}} \supset e(\epsilon_1 V_{1,\mu} + \epsilon_2 V_{2,\mu}) J_{\text{EM}}^{\mu} + (g_1 V_{1,\mu} + g_2 V_{2,\mu}) J_D^{\mu} + (g_1' V_{1,\mu} + g_2' V_{2,\mu}) J_D^{\prime \mu}$ 

[See Adrian Thompson's talk for (pseudo)scalar scenarios]



Scattering may happen through an exchange of a different mediator.

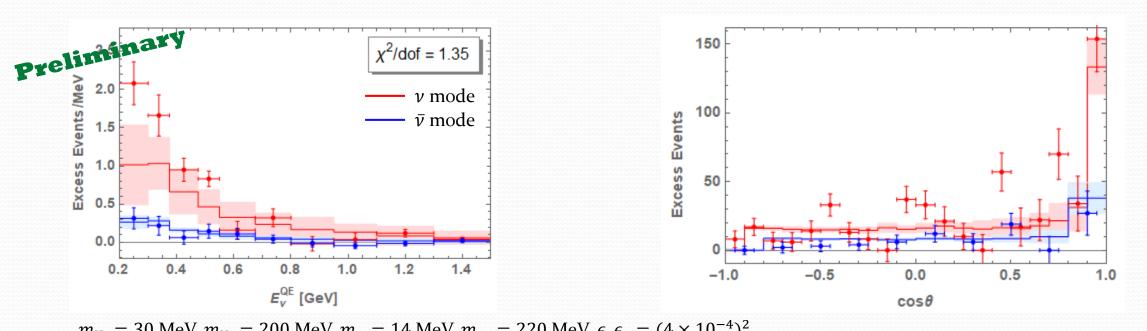
Detector

 $BR(V \rightarrow 2\chi):BR(V \rightarrow 2e) = 50\%: 50\%$  for illustration

□ Neutral meson contributions are small as they are **not focused** and their decays involve **no BR enhancement**.

**Collimated rings** 

# **Example Fit**

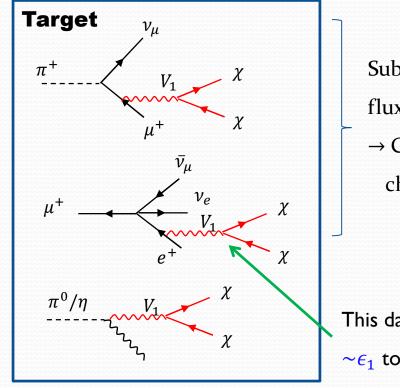


 $m_{V_1} = 30$  MeV,  $m_{V_2} = 200$  MeV,  $m_{\chi} = 14$  MeV,  $m_{\chi'} = 220$  MeV,  $\epsilon_1 \epsilon_2 = (4 \times 10^{-4})^2$ 

- □ The best-fit parameters are consistent with various limits including limits for (in)visibly decaying dark photons, limits for the exotic decays of charged  $\pi^+/K^+$ , and the MiniBooNE off-target mode (~30 times smaller POTs) measurement
- □ We have found equally good fits in the single-mediator scenario.

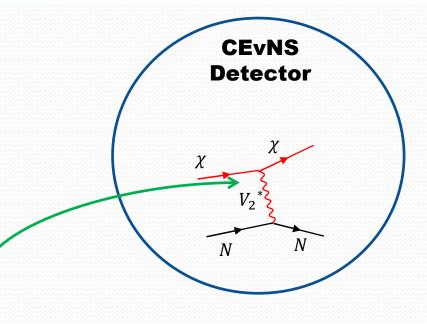
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# **Predictions for CEvNS Experiments**



Subdominant because the resulting χ
flux is not focused, but isotropic.
→ Conventional neutral meson
channels are important.

This dark-sector coupling is as small as  $\sim \epsilon_1$  to have 50%:50% BRs.



This dark-sector coupling in the double-mediator scenario can be large enough for CEvNS experiments to observe MB signal events.

□ Moreover, more energetic beam based experiments, e.g., JSNS<sup>2</sup>, can be sensitive to signals coming from the decay of kaons.

## Conclusions

 The MiniBooNE excess can be explained by darksector scenarios, using three-body decays of charged mesons that are focused by the horn system.
 CEvNS experiments can test new physics scenarios of explaining the MiniBooNE excess.