Superlso Relic tutorial

# Alexandre Arbey

IP2I & Lyon U. & CERN TH

Development team: AA, F.N. Mahmoudi, M. Palmiotto & G. Robbins

Code references: arXiv:0906.0369 and 1806.11489

https://superiso.in2p3.fr/

**Computational Tools for High Energy Physics and Cosmology** 

IP2I, Lyon, France – November 22nd, 2021

### Main purpose

Computing dark matter observables in standard and non-standard cosmological scenarios

## Main features

- public and open-source C program
- (hopefully) easy to modify and user-friendly
- includes different cosmological scenarios
- computes dark matter relic density
- computes dark matter direct and indirect detection observables
- computes flavour physics observables
- incorporates supersymmetric models
- no external library needed
- interfaced with AlterBBN to constrain cosmological scenarios with BBN

### Webpage

Download from

```
https://superiso.in2p3.fr/
```

A detailed manual is also available.

## Compilation instructions

- to uncompress: tar xjvf superiso\_relic\_v4.1.tar.bz2
- enter directory: cd superiso\_relic\_v4.1/
- to set-up compilation: ./configure --with-cc=gcc --with-fc=gfortran
- ullet to compile the library: make  $\longrightarrow$  This will take a few minutes...
- to compile the program xxx.c: make xxx

Let us focus on the observable linking particle physics and cosmology:

## Dark matter relic density

Radiation domination true at temperatures below  $\sim\,\text{MeV}$ 

However, this may not hold at higher temperatures.

If the calculated relic density is different from the measured dark matter density, it could be due to novel phenomena in the early Universe. In the Standard Model of Cosmology:

• before and at nucleosynthesis time, the expansion is dominated by radiation

$$H^2 = 8\pi G/3 imes 
ho_{\mathsf{rad}}$$

• the evolution of the number density of all NP particles follows the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\rm eff} v \rangle (n^2 - n_{\rm eq}^2)$$

• the time and temperature are related through the adiabaticity condition:

$$rac{ds_{\mathsf{rad}}}{dt} = -3Hs_{\mathsf{rad}}$$

with  $s_{rad} \propto h_{eff}(T) T^3$  ( $h_{eff}$ : radiation entropy degrees of freedom)

 $\langle \sigma_{\rm eff} v \rangle$ : related to the amplitudes of (co-)annihilations of BSM particles into SM particles

 $\langle \sigma_{\rm eff} v \rangle$ : Thermal average of effective cross section ( $\kappa_{1,2}$ : modified Bessel functions):

$$\langle \sigma_{\mathrm{eff}} v 
angle = rac{\int_{0}^{\infty} dp_{\mathrm{eff}} p_{\mathrm{eff}}^{2} W_{\mathrm{eff}} \mathcal{K}_{1}\left(rac{\sqrt{s}}{T}
ight)}{m_{1}^{4} T \left[\sum_{i} rac{g_{i}}{g_{\chi}} rac{m_{i}^{2}}{m_{\chi}^{2}} \mathcal{K}_{2}\left(rac{m_{\chi}}{T}
ight)
ight]^{2}}$$

where: (ij: coannihilating BSM particles / kl: SM outgoing particles)

$$\frac{dW_{\rm eff}}{d\cos\theta} = \sum_{ijkl} \frac{p_{ij}p_{kl}}{32\pi p_{\rm eff}S_{kl}\sqrt{s}} \sum_{\rm helicities} \left|\sum_{\rm diagrams} \mathcal{M}(ij \to kl)\right|^2$$

Dark matter density normalised to radiation entropy density as a function of  $m_{\chi}/T$ .



AA & F. Mahmoudi, Prog.Part.Nucl.Phys. 119 (2021) 103865

The moment at which the dark matter density leaves the equilibrium density is called freeze-out.

The differential equations are solved from an initial temperature  $T_{init}$  down to the present temperature  $T_0 = 2.725$  K

The relic density is then obtained:

$$\Omega_{\chi}h^2(T_0)\equiv 2.755 imes 10^{-8}rac{
ho_{\chi}(T_0)}{s_{rad}(T_0)} \qquad ext{with} \ 
ho_{\chi}=m_{\chi}\,n(T_0)$$

Very precise measurements of cold dark matter density by Planck (+ others):  $\Omega_c h^2({\cal T}_0)=0.120\pm 0.001$ 

The Planck results lead to very strong constraints on BSM parameters.

#### Relic density in alternative cosmological scenarios

For example, the expansion rate can be modified:

$$H^2 = 8\pi G/3 \times (\rho_{rad} + \rho_D)$$

The entropy content of the Universe can also be altered!

$$rac{ds_{\mathsf{rad}}}{dt} = -3Hs_{\mathsf{rad}} + \Sigma_D$$

 $\Rightarrow$  Modified relation between time, expansion rate and temperature!

And relics can be generated non-thermally:

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\rm eff} v \rangle (n^2 - n_{\rm eq}^2) + N_D$$

 $\rho_D$ ,  $\Sigma_D$  and  $N_D$  are model-dependent.

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### Main programs

- amsb.c, cmssm.c, gmsb.c, hcamsb.c, mmamsb.c, nuhm.c: observable calculation in different MSSM scenarios
- cnmssm.c, ngmsb.c, nnuhm.c, mmamsb.c: observable calculation in different NMSSM scenarios
- test\_modeleff.c: effect of QCD eos on relic density
- test\_widthcalc.c: effect of Higgs widths on relic density
- test\_standmod.c: effect of non-standard cosmological models on relic density
- test\_reheating.c: effect of reheating cosmological scenario on relic density
- test\_phi.c: effect of decaying scalar field on relic density
- direct.c: calculation of direct detection observables
- indirect.c: calculation of indirect detection observables
- create\_propagation.c: calculation of indirect detection observables within user-defined propagation and DM profile models
- blackholes.c: indirect detection constraints for PBH spectra from BlackHawk
- flha.c, modelindep\_chi2.c, slha\_chi2.c, sm.c, sm\_chi2.c, thdm.c: flavour observables only

### With which arguments?

Running ./program.x without argument will show you the possible arguments

```
Example: ./test_standmod.x
```

```
This program assumes that in the early Universe, during the dark times (before BBN), a dark entropy and/or
a dark density could modify the expansion and thermal properties.
The program consider the presence of a dark density such as
rho dark(T) = rho 0 T^ndd
and of a dark entropy such as
s_dark(T) = s_0 T^nsd
rho dark(T) = rho 0 T^ndd
and of non-thermal production of relic particles such as
N_nt(T) = nt0 T^nnt
It needs 5 parameters:
          name of the SLHA file
  name
  dd0
          dark energy proportion to photon density at BBN time (1 MeV)
          dark energy decrease exponent (preferentially >4)
  ndd
  sd0
          dark entropy proportion to photon entropy at BBN time (1 MeV)
          dark entropy decrease exponent
  ns d
Auxiliary parameters are:
  ЪТ
          dark energy cut temperature (in GeV)
  Τs
          dark entropy cut temperature (in GeV)
          Non thermal production rate at BBN time
  nt 0
          Non thermal production rate decrease exponent
  nnt
  Τt
          Non thermal production rate cut temperature (in GeV)
```

```
./stand_cosmo.x example.lha 3 6 0 0
SuperIso Relic v4.1 - A. Arbey, F. Mahmoudi & G. Robbins
AlterBBN v2.1 - A. Arbey, J. Auffinger, K. Hickerson & E. Jenssen
For the cosmological standard model:
omega=1.429e+01
For the specified model with dark density/entropy/non thermal relics:
omega=1.263e+04
Model excluded by BBN constraints
```

Starting from an SLHA point with a relic density of 14.29, adding a dark density with 3 times the radiation density at BBN time and decrease exponent 6, the relic density becomes  $1.263 \times 10^4$ .

This model is nevertheless excluded by BBN constraints, as tested automatically by AlterBBN.

#### Three Dark Matter benchmark cases

- CONSERVATIVE: AMS-02 antiprotons with Burkert profile and MED propagation model + local density of 0.2 GeV/cm<sup>3</sup>
- STANDARD: Fermi-LAT gamma rays with NFW profile + local density of 0.4 GeV/cm<sup>3</sup>

 STRINGENT: AMS-02 antiprotons with NFW profile and MAX propagation model + local density of 0.6 GeV/cm<sup>3</sup>

#### ./direct.x example.lha

```
/----WIMP-NUCLEON cross-section(pb) at 0-momentum transfer----/
 /-Spin-Independent-/
Conservative
                      Standard
                                      Stringent
Proton 1.838746e-10 1.838746e-10 1.838746e-10
Neutron 1,907715e-10 1,907715e-10
                                      1.907715e-10
/-Spin-Dependent-/
Conservative
                      Standard
                                      Stringent
Proton 5.242343e-07 5.242343e-07
                                      5 242343e-07
Neutron 5.332762e-07 5.332762e-07
                                      5 332762e-07
/----PANDAX-2 2017 Poisson delta-loglikelihood (point excluded at 2 sigma if <-4.000000 with 1 d.o.f.)----/
Conservative
                St and ar d
                               Stringent
-2.265334e+00 -2.265334e+00 -2.265334e+00
/----XENON1T 2017 Poisson delta-loglikelihood (point excluded at 2 sigma if <-4.000000 with 1 d.o.f.)----/
Conservative
                St and ar d
                               Stringent
-2.076444e+00 -2.076444e+00 -2.076444e+00
/----PICD60 2017 Poisson delta-loglikelihood (point excluded at at 2 sigma if <-4.000000 with 1 d.o.f.)----/
Conservative
                St and ar d
                               Stringent
 -2 261627e-02 -2 261627e-02
                                -2 261627e-02
```

### Example of indirect detection constraints

## ./indirect.x example.lha

channel	<sigmav>[cm^3/</sigmav>	s]	
<pre>clickeric c</pre>	<pre>6 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1</pre>	-	
o1 o1 -> h z	2.05907e-31		
Total annihilation cross section 4.082059e-30 [cm^3/s]			
Fermi-LAT			
delta-loglikelihood : excluded at 2 sigma if <-12.850000 with 6 d.o.f.			
conservative	standard	stringent	
2.905715e-04	2.728380e-04	- -2.857854e-04	
 AMS-02			
delta-loglikelihood : excluded at 2 sigma % CL if <-12.850000 with 6 d.o.f.			
conservative	standard	stringent	
1.655830e-02	2.213955e-02	7.296463e-02	

./slha.x example.lha	
()	
Relic density Oh2	1.429e+01
SI proton xsection SD proton xsection	1.839e-10 5.242e-07
excluded_Xenon1T (standard) excluded_PANDAX (standard) excluded_PICO60 (standard)	0 0 0
Tot annihilation xsection	4.082e-30
excluded_Fermi (standard) excluded_AMS02 (standard)	0 0

#### Relic density: decaying scalar field

Scenario with a pressureless decaying scalar field (e.g. modulus, late inflaton, dilaton, ...) of energy density  $\rho_{\phi}$ :

$$H^2 = 8\pi G/3 \left(\rho_{rad} + \rho_{\phi}\right)$$

We define the scalar field decay width  $\Gamma_{\phi}$ , with a large branching fraction to radiation and a (tiny) branching ratio b to WIMPs:

$$\begin{aligned} \frac{d\rho_{\phi}}{dt} &= -3H\rho_{\phi} - \Gamma_{\phi}\rho_{\phi} \\ \frac{ds_{rad}}{dt} &= -3Hs_{rad} + \frac{\Gamma_{\phi}\rho_{\phi}}{T} \\ \frac{dn}{dt} &= -3Hn - \langle \sigma_{\text{eff}}v \rangle \left(n^2 - n_{eq}^2\right) + \frac{b}{m_{\phi}}\Gamma_{\phi}\rho_{\phi} \end{aligned}$$

Reheating temperature  $T_{RH}$  (at which the scalar field is mostly decayed) defined by:

$$\Gamma_{\phi} = \sqrt{rac{4\pi^3 g_{eff}(T_{RH})}{45}} rac{T_{RH}^2}{M_P}$$

Non-thermal production parameter:

$$\eta = b\left(rac{1 ext{ GeV}}{m_{\phi}}
ight)$$

Initial (relative) scalar field density:

$$\kappa_{\phi} = rac{
ho_{\phi}(T_{init})}{
ho_{\gamma}(T_{init})}$$

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AA, J. Ellis, F. Mahmoudi & G. Robbins, arXiv:1807.00554

Evolution of the scalar field density, WIMP density and entropy injection  $\tilde{\Sigma}^* \equiv \frac{\Gamma_{\phi} \rho_{\phi}}{3HT_{s_{rad}}}$ as a function of  $x = m_{\chi}/T$ , in absence of non-thermal production of WIMPs  $(\eta = 0)$ 



(a)  $T_{RH} = 0.01 \text{ GeV}$ ,  $\kappa_{\phi}^{init} = 100$ ,  $T_{init} = 40 \text{ GeV}$ 

(b)  $T_{RH} = 10 \text{ GeV}$ ,  $\kappa_{\phi}^{init} = 100$ ,  $T_{init} = 40 \text{ GeV}$ 

Complex interplay between expansion rate and entropy injection...

In absence of non-thermal WIMP production, results in a decrease of the relic density

AA, J. Ellis, F. Mahmoudi & G. Robbins, arXiv:1807.00554

Evolution of the WIMP density as a function of  $x = m_{\chi}/T$  in presence of non-thermal production of WIMPs



Standard scenario can be strongly modified by the non-thermal production of WIMPs.

Alexandre Arbey

#### Relic density: decaying scalar field vs. Big-Bang nucleosynthesis

AA, J. Ellis, F. Mahmoudi & G. Robbins, arXiv:1807.00554

Value of the relic density as a function of  $T_{RH}$  and  $\kappa_{\phi}$  for a pMSSM example point with  $\Omega_{\rm standard} h^2 = 1.27$ , in absence of non-thermal production  $(\eta = 0)$ 



The gray region is excluded by Big-Bang nucleosynthesis constraints

The dark strip is compatible with Planck results

AA, J. Ellis, F. Mahmoudi & G. Robbins, arXiv:1807.00554

Minimal value of  $\kappa_{\phi}$  to obtain the observed relic density as a function of original relic density  $\Omega_{\text{standard}} h^2$  for a sample of pMSSM points, with  $T_{RH} = 6$  MeV,  $T_{init} = 40$  GeV and in absence of non-thermal production ( $\eta = 0$ )



This sets constraints on the primordial Universe...

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- Possibility to input tables of dark density or entropy for non-standard cosmological scenarios.
- Interface with MARTY for automatic dark matter observable calculation in any BSM scenario
  - one type of dark matter particles
  - several types of dark matter particles
  - density of each NP particle through time
  - freeze-out and freeze-in scenarios
- Interface with BlackHawk
  - NP particles radiated by PBHs
  - DM as PBHs and particles
  - Indirect detection of PBHs through Hawking radiation
- Automatic calculations of Wilson coefficients, flavour physics observables and  $(g-2)_{\mu}$  in any BSM scenario, using MARTY