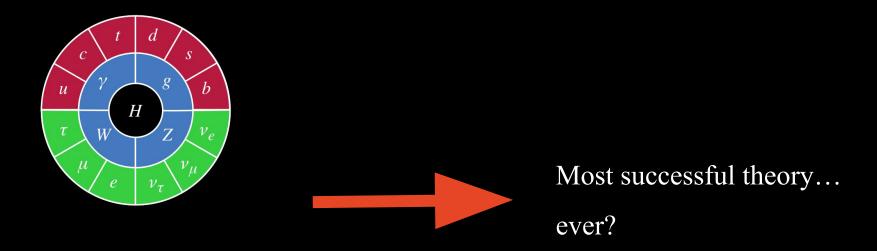
Probing the universe at the intersection of gravity and particle physics

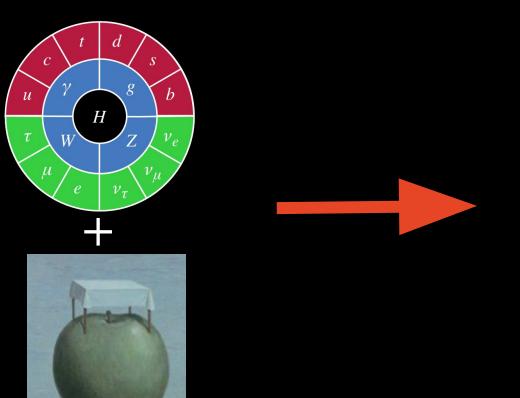
Zachary Picker
CPPC Seminar, Oct 2021



Gravity and particle physics are exciting partners

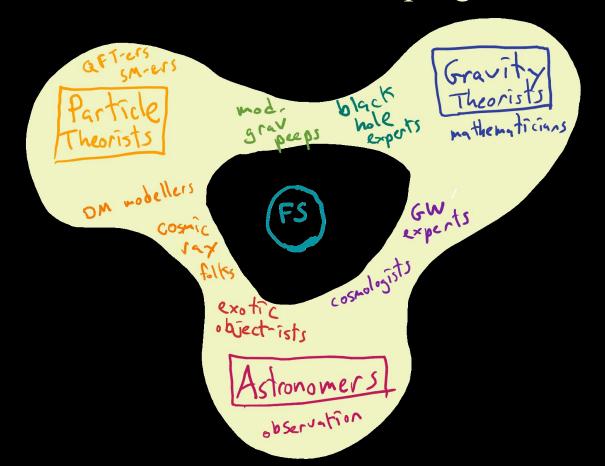


Gravity and particle physics are exciting partners

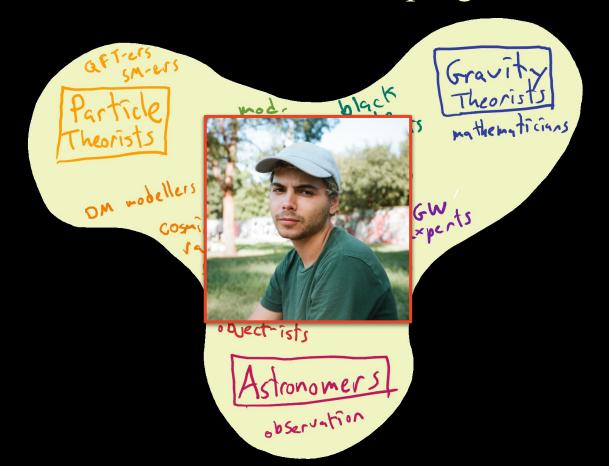


- Quantum gravity?
- Black holes?
- Cosmological constant problem?
 - (least successful prediction... ever?)
- Dark matter?

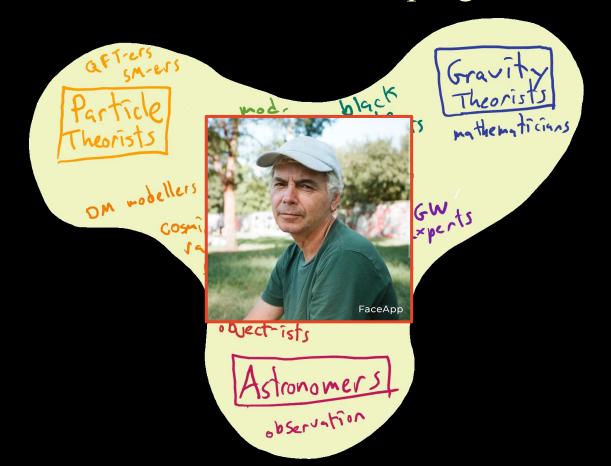
You have to know a lot to make progress...



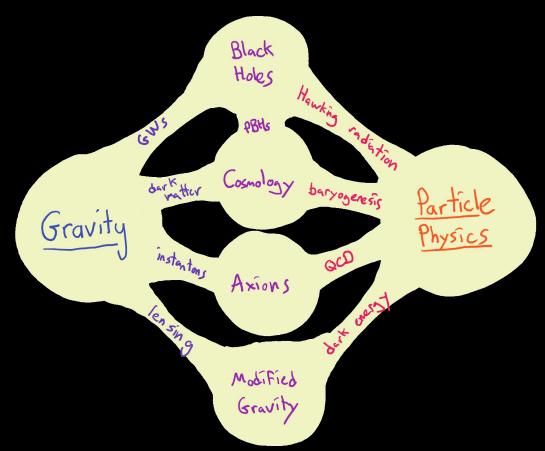
You have to know a lot to make progress...

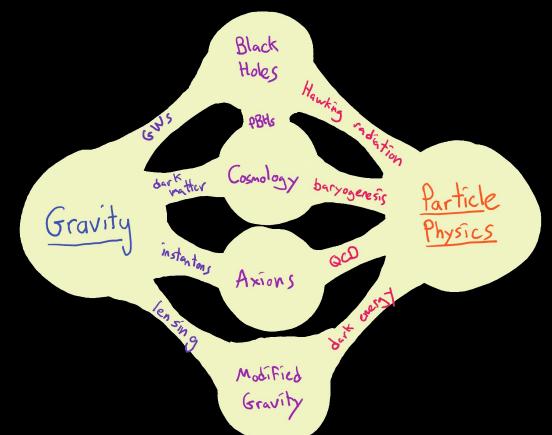


You have to know a lot to make progress...

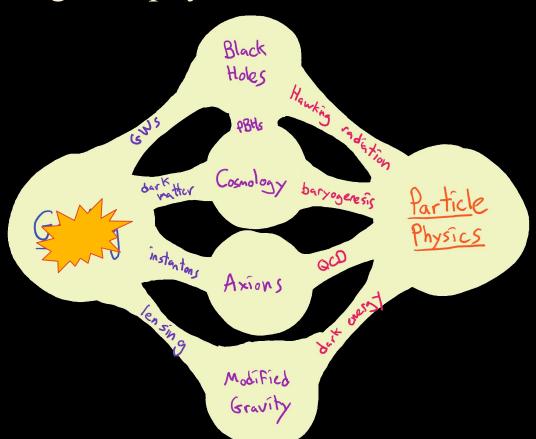


Gravity + PP is a useful lens to study phenomenology

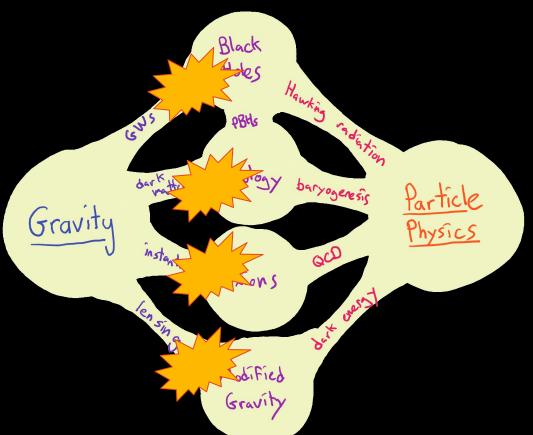




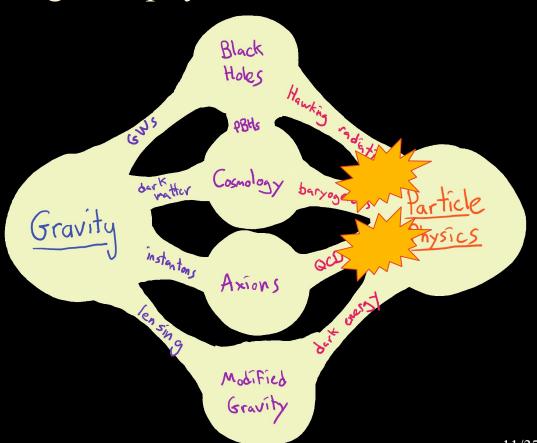
1. What happens when we modify gravity?



- 1. What happens when we modify gravity?
- 2. What happens when we find new solutions within GR?



- 1. What happens when we modify gravity?
- 2. What happens when we find new solutions within GR?
- 3. What happens when we add gravity to particle physics?



1. What happens when we change gravity?

Using observation to constrain theory...GWs

'Quadratic gravity':

- Extra massive scalar and spin-2 gravitons
- BH Binaries deplete energy faster
 - \Rightarrow tightests constraints on theory

$$m_{\phi,\pi} \gtrsim 7.0 \cdot 10^{-12} \text{ eV}$$



arXiv:1906.12034, Kim, Kobakhidze, Picker

Using theory to reexamine constraints...microlensing

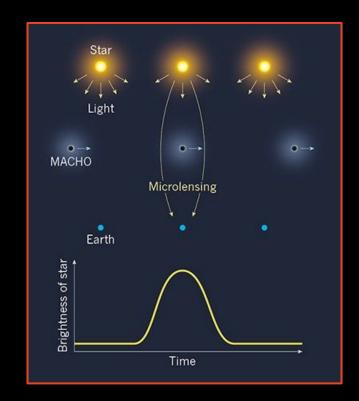
Quadratic gravity

Too constrained (by us!)

'Bimetric gravity'

- Coupled to massive graviton
- Screened at close range
 - \Rightarrow need much farther sources,

or more massive lenses...



arXiv:1906.12034 **Probing Quadratic** Black Gravity with Holes Hawking **Binary Inspirals** BRG Cosmology baryogenesis darkter Particle Gravity Physics instantons QCD Axions Modified Gravity

'Null' result...

2. What happens when we find new solutions?

Primordial black holes can't be Schwarzschild BHs...



Primordial black holes can't be Schwarzschild BHs...

THE HYPOTHESIS OF CORES RETARDED DURING EXPANSION AND THE HOT COSMOLOGICAL MODEL Ya. B. Zel'dovich and I. D. Novikov

Translated from Astronomicheskii Zhurnal, Vol. 43, No. 4, pp. 758-760, July-August, 1966 Original article submitted March 14, 1966

The existence of bodies with dimensions less than $R_g=2GM/c^2$ at the early spansion of the cosmological model leads to a strong accretion of radiation by If further calculations confirm that accretion is catastrophically high, the hyperores retarded during expansion [3, 4] will conflict with observational data.

BLACK HOLES IN THE EARLY UNIVERSE

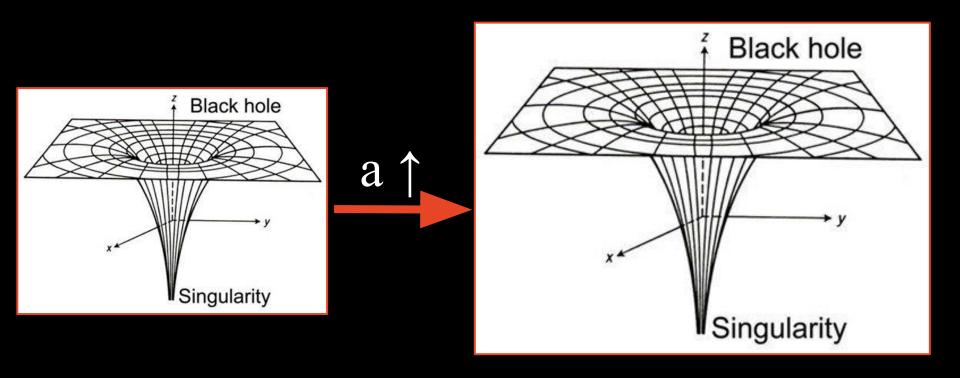
B. J. Carr and S. W. Hawking

(Received 1974 February 25)

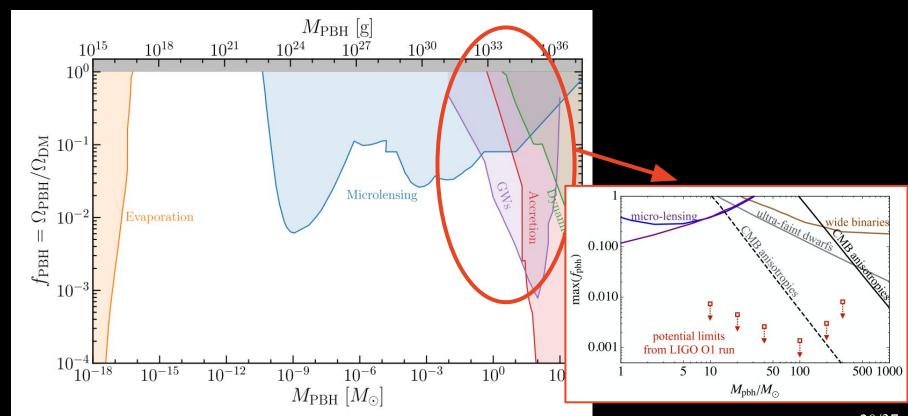
SUMMARY

The existence of galaxies today implies that the early Universe must have been inhomogeneous. Some regions might have got so compressed that they underwent gravitational collapse to produce black holes. Once formed, black holes in the early Universe would grow by accreting nearby matter. A first estimate suggests that they might grow at the same rate as the Universe during the radiation era and be of the order of 10¹⁵ to 10¹⁷ solar masses now. The observational evidence however is against the existence of such giant black holes. This motivates a more detailed study of the rate of accretion which shows that black holes will not in fact substantially increase their original mass by accretion. There could thus be primordial black holes around now with masses from 10⁻⁵ g upwards.

The Thakurta metric is a simple cosmological black hole



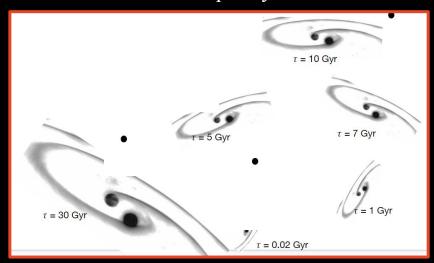
Could LIGO-sized PBHs be dark matter?



Thakurta black holes do not form binaries

Schwarzschild PBHs:

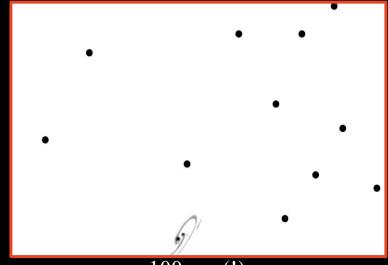
At matter-radiation equality:



Many of these coalesce ~today

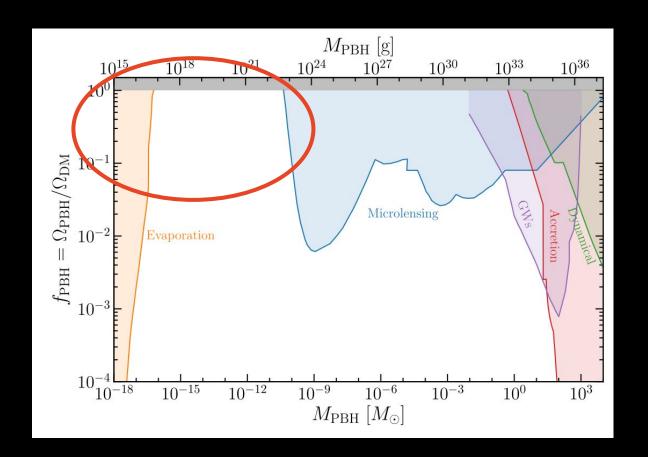
Thakurta PBHs:

At matter-radiation equality:

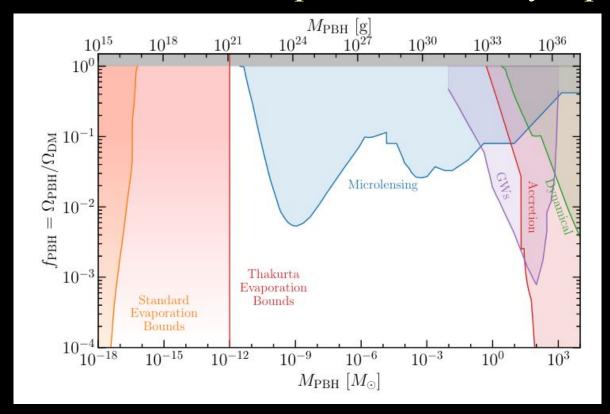


 $\tau_{\rm max} \sim 100 \ {\rm sec} \ (!)$

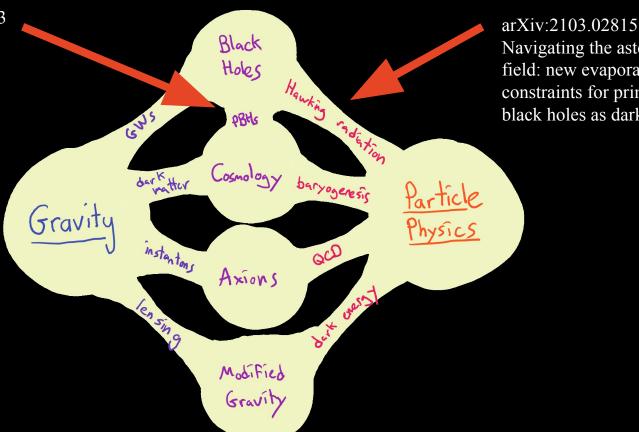
Asteroid-mass PBHs are not constrained



Thakurta black holes evaporate extremely rapidly

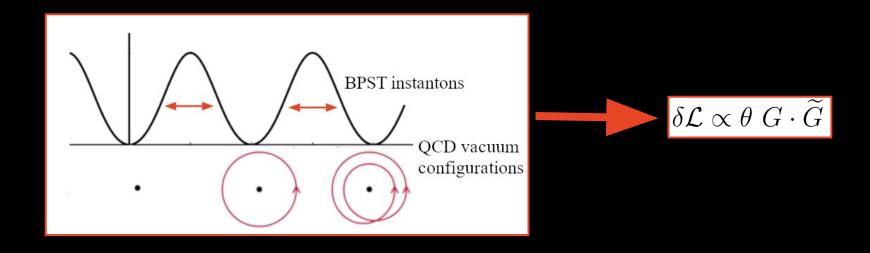


arXiv:2008.10743 Eliminating the LIGO bounds on primordial black hole dark matter

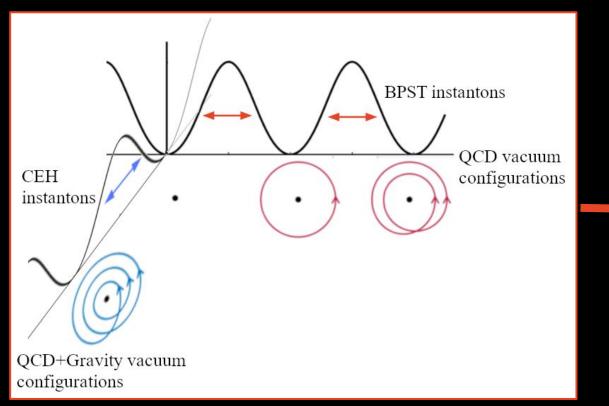


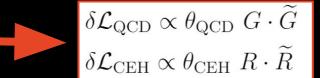
Navigating the asteroid field: new evaporation constraints for primordial black holes as dark matter 3. What happens when we add gravity to PP?

The 'axion' solves the strong-CP problem

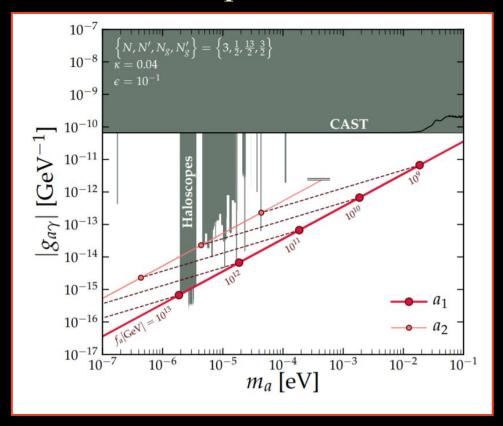


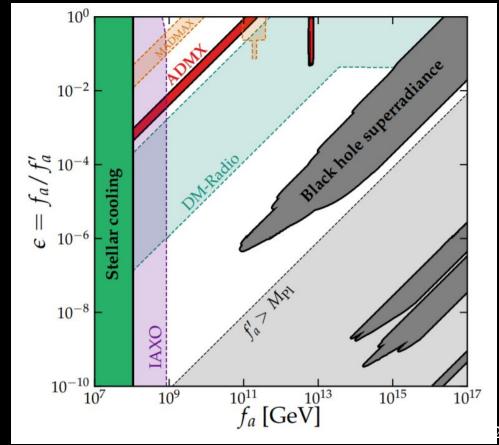
When we add gravity, we need a second axion



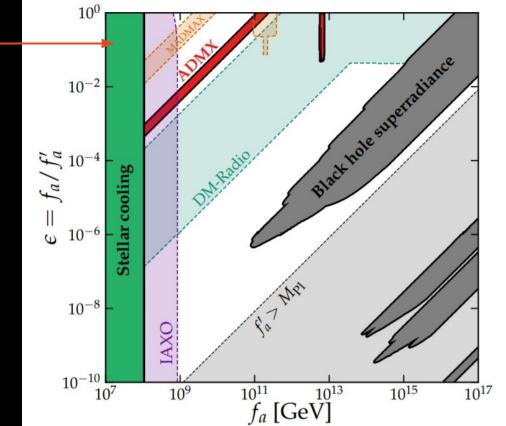


The second axion is coupled to the 'usual' one

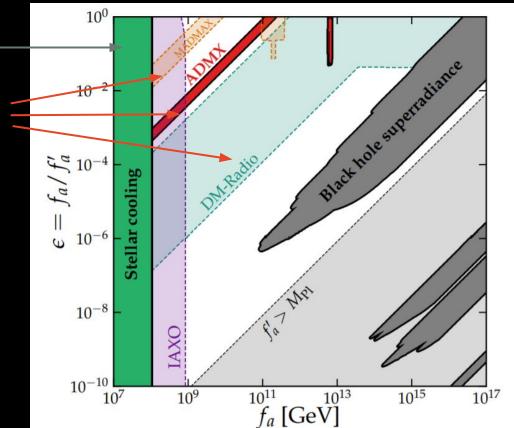




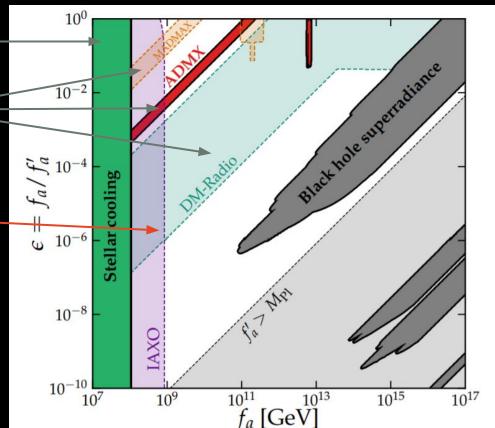
Axion production cools stars



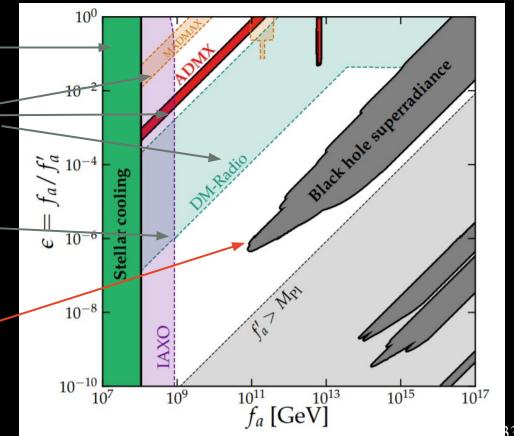
- Axion production cools stars
- Haloscopes: detect axions in dark matter halo using
 resonant cavity



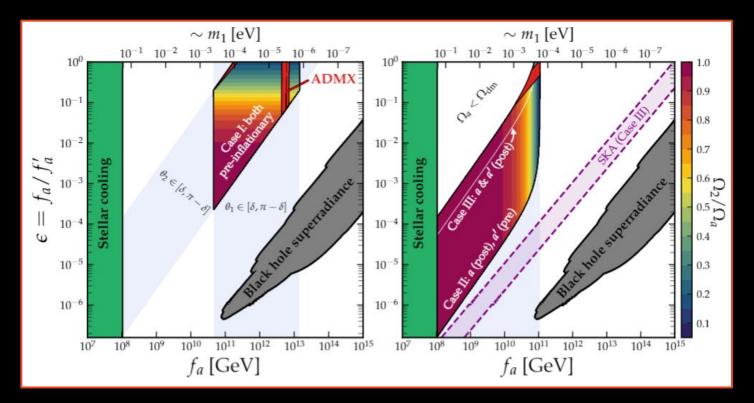
- Axion production cools stars
- Haloscopes: detect axions in dark matter halo using resonant cavity
- Helioscopes: detect stellar
 axions by converting back to
 photons

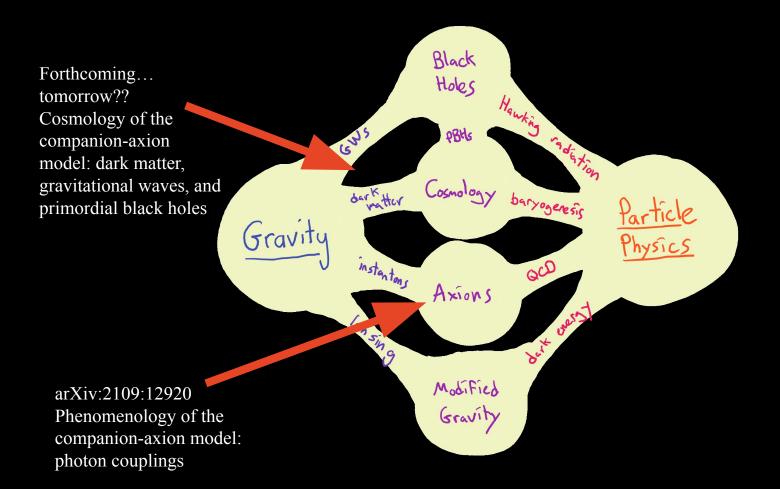


- Axion production cools stars
- Haloscopes: detect axions in dark matter halo using resonant cavity
- Helioscopes: detect stellar
 axions by converting back to
 photons
- Spin down black holes



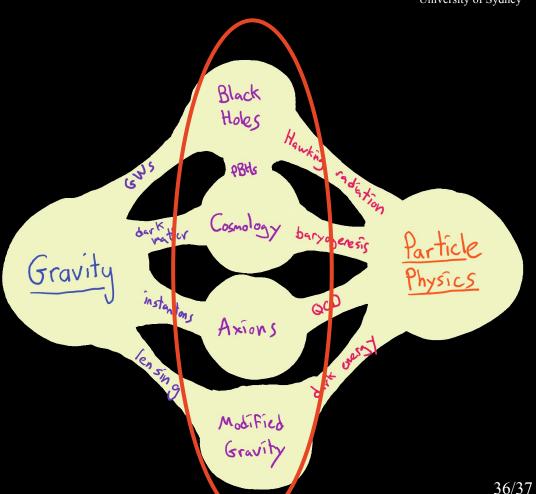
Both axions can be DM, and produce GWs... (and solve the domain wall problem)





Morals and takeaways

- Frontier of physicsbetween gravity and PP
- Lots of unansweredquestions...
 - ⇒ lots of surprising phenom
- Tweak any bit and seewhat happens!



Thanks!

Bonus: quadratic gravity

$$S = \int d^4x \sqrt{-g} \left[\frac{R}{2\kappa} + \beta R^2 + \gamma R^{\mu\nu} R_{\mu\nu} \right]$$



$$S = \int d^4x \sqrt{-\tilde{g}} \left[\frac{\tilde{R}}{2\kappa} - \frac{1}{2} \left(\partial_{\mu}\phi \partial^{\mu}\phi + m_{\phi}^2 \phi^2 \right) - \frac{1}{2} \left(\partial_{\mu}\pi^{\alpha\beta}\partial^{\mu}\pi_{\alpha\beta} + m_{\pi}^2\pi^{\alpha\beta}\pi_{\alpha\beta} \right) \right]$$

$$0 \le \gamma \lesssim 5.9 \cdot 10^{76},$$

$$-\frac{\gamma}{4} \le \beta \lesssim 9.8 \cdot 10^{75} - \frac{\gamma}{4},$$

Bonus: thakurta metric details

$$\mathrm{d}s^2 = a^2 \mathrm{d}s_{schw}^2.$$

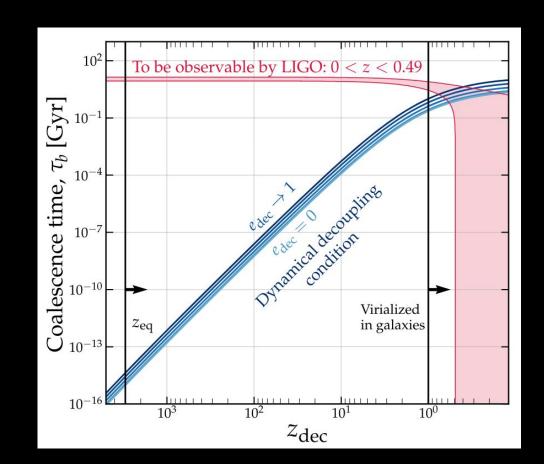
$$f(R) = 1 - 2Gma(t)/R$$

$$ds^{2} = f(R) \left(1 - \frac{H^{2}R^{2}}{f^{2}(R)} \right) dt^{2} + \frac{2HR}{f(R)} dt dR - \frac{dR^{2}}{f(R)} - R^{2} \left(d\theta^{2} + \sin^{2}\theta d\phi^{2} \right)$$

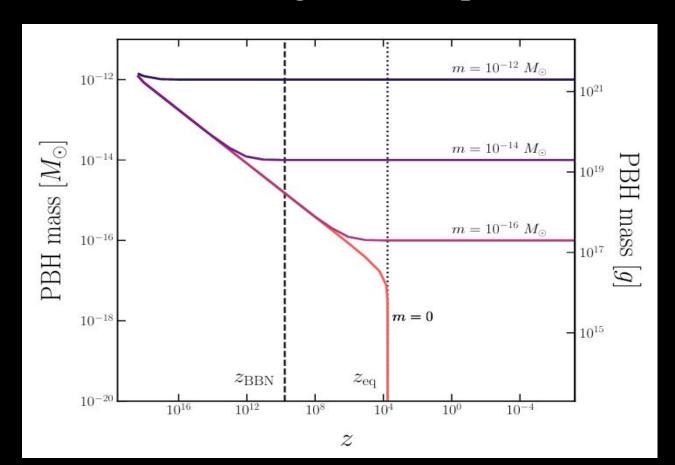
$$m_{\rm MS} = ma(t) + \frac{H^2 R^3}{2Gf(R)}$$

$$R_{\rm b} = \frac{1}{2H} \left(1 - \sqrt{1 - 8HGma(t)} \right) \approx 2ma$$

Bonus: thakurta binary formation plot



Bonus: thakurta hawking radiation plot



Bonus: companion axion details

$$V(a, a') = -2K\cos\left(N\frac{a}{f_a} + N'\frac{a'}{f_a'} + \theta\right) - 2\kappa K\cos\left(N_g\frac{a}{f_a} + N'_g\frac{a'}{f_a'} + \theta_g\right)$$

$$m_1 \propto 1/f_a$$

 $m_2 \approx \epsilon \sqrt{\kappa} m_1$
 $\epsilon \equiv f_a/f_a'$

$$\frac{\Omega_{a_2}}{\Omega_{a_1}} \sim \frac{\theta_2^2}{\theta_1^2} \kappa^{0.41} \epsilon^{-1.19}$$

For gravity, $\sim 0.04 - 0.6$

$$\mathcal{L}_{a\gamma} = \frac{1}{4} \left(a g_{a\gamma} + a' g'_{a\gamma} \right) F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$ds^{2} = \frac{1}{1 - \frac{a^{4}}{r^{4}}} dr^{2} + \frac{r^{2}}{4} \left[d\theta^{2} + \sin^{2}\theta d\phi^{2} + \left(1 - \frac{a^{4}}{r^{4}} \right) \left(d\psi + \cos\theta d\phi \right)^{2} \right],$$