

The Road to Discovery

Experimental Search Methods

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Goals

- For theorists (and experimenters):
 - What is it these experimenters *do*?
 - (Except for using fancy equipment to build fancy detectors...)
 - How come it takes them forever to release a result?
 - Why can't they just give us the 4-vectors as they record the data?
 - What can('t) they look for?
- For experimenters:
 - Why does it take so many layers of review to get a result out? My analysis seems pretty simple at first sight....

HEP in 2011

CKM elements:

3 Generations of Fermions

$\frac{2}{3}$ u ~ 5	$\frac{2}{3}$ c ~ 1350	$\frac{2}{3}$ t 175000
$-\frac{1}{3}$ d ~ 9	$-\frac{1}{3}$ s ~ 175	$-\frac{1}{3}$ b ~ 4500
ν_1 ν	ν_2 ν	ν_3 ν
e 0.511	μ 105.66	τ 1777.2

Masses are in MeV

Force Carriers

g 0
γ 0
Z^0 0 91187
W^\pm ± 1 81400

Observable	Central $\pm 1 \sigma$
$ V_{ud} $	0.97430 [+0.00019 -0.00019]
$ V_{us} $	0.22521 [+0.00082 -0.00082]
$ V_{ub} $	0.00350 [+0.00015 -0.00014]
$ V_{cb} $	0.04117 [+0.00038 -0.00115]
$ V_{ud} $ (meas. not in the fit)	0.97444 [+0.00028 -0.00028]
$ V_{us} $ (meas. not in the fit)	0.2257 [+0.0011 -0.0011]
$ V_{ub} $ (meas. not in the fit)	0.00350 [+0.00015 -0.00016]
$ V_{cb} $ (meas. not in the fit)	0.04399 [+0.00069 -0.00397]
$ V_{cd} $	0.22508 [+0.00082 -0.00082]
$ V_{cs} $	0.97347 [+0.00019 -0.00019]
$ V_{td} $	0.00859 [+0.00027 -0.00029]
$ V_{ts} $	0.04041 [+0.00038 -0.00115]
$ V_{tb} $	0.999146 [+0.000047 -0.000016]

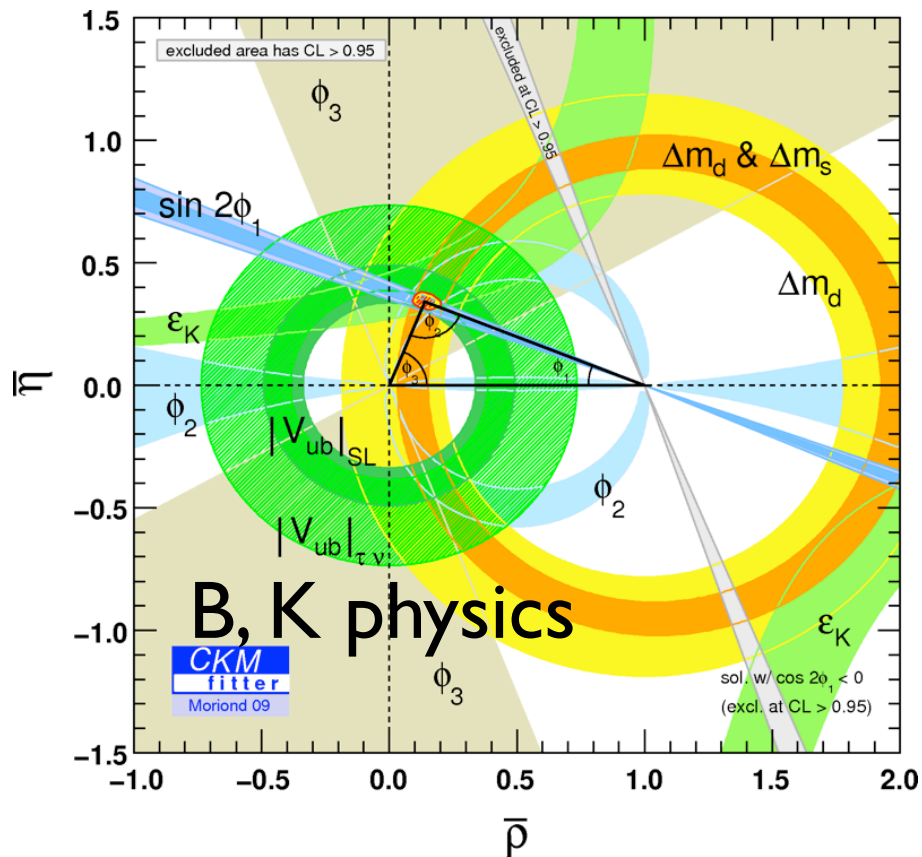
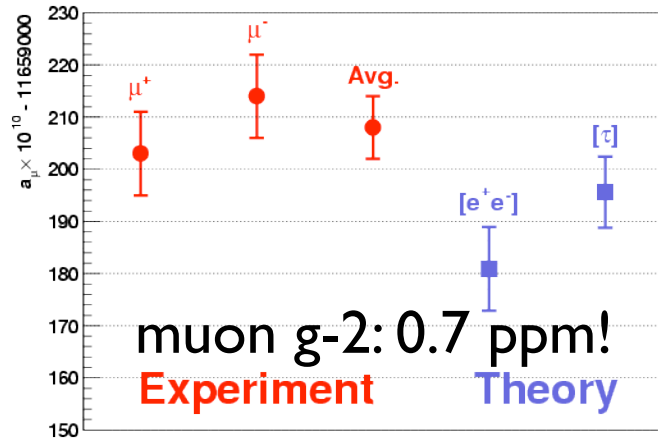
In Words

- Matter is built of spin $1/2$ particles that interact by exchanging 3 different kinds of spin 1 particles corresponding to 3 different (gauge) interactions
- There appear to be 3 generations of matter particles
- The 4 different matter particles in each generation carry different combinations of quantized charges characterizing their couplings to the interaction bosons
- The matter fermions and the weak bosons have “mass”
- Gravitation is presumably mediated by spin 2 gravitons
- Gravitation is extremely weak for typical particle masses
- There appear to be 3 macroscopic dimensions

About the Standard Model

- It's a theory of interactions:
 - Properties of fermions are inputs
 - Properties of interaction bosons in terms of couplings, propagations, masses are linked:
 - Measuring a few allows us to predict the rest, then measure and compare with expectation
- It's remarkably successful:
 - Predictions verified to be correct at sometimes incredible levels of precision
 - After ~30 years, still no serious cracks

Precision Results



	Measurement	Fit	$\frac{O^{\text{meas}} - O^{\text{fit}}}{\sigma^{\text{meas}}}$
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	0.02758 ± 0.00035	0.02768	0.1
m_Z [GeV]	91.1875 ± 0.0021	91.1875	0.0
Γ_Z [GeV]	2.4952 ± 0.0023	2.4957	0.1
σ_{had}^0 [nb]	41.540 ± 0.037	41.477	1.7
R_1	20.767 ± 0.025	20.744	0.9
$A_{\text{fb}}^{0,l}$	0.01714 ± 0.00095	0.01645	0.7
$A_1(P_{\tau})$	0.1465 ± 0.0032	0.1481	0.5
R_b	0.21629 ± 0.00066	0.21586	0.7
R_c	0.1721 ± 0.0030	0.1722	0.0
$A_{\text{fb}}^{0,b}$	0.0992 ± 0.0016	0.1038	2.9
$A_{\text{fb}}^{0,c}$	0.0707 ± 0.0035	0.0742	1.0
A_b	0.923 ± 0.020	0.935	0.6
A_c	0.670 ± 0.027	0.668	0.0
$A_1(\text{SLD})$	0.1513 ± 0.0021	0.1481	1.5
$\sin^2\theta_{\text{eff}}^{\text{lept}}(Q_{\text{fb}})$	0.2324 ± 0.0012	0.2314	0.8
m_W [GeV]	80.398 ± 0.025	80.374	0.9
Γ_W [GeV]	2.140 ± 0.060	2.091	1.0
m_t [GeV]	170.9 ± 1.8	171.3	0.2

LEP, SLD & Tevatron

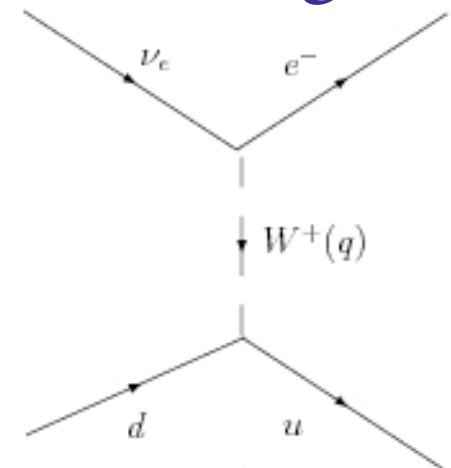
Many Fundamental Questions

- What exactly *is* spin? Or color? Or electric charge?
Why are they quantified?
- Are there only 3 generations? If so, why?
- Why are there e.g. no neutral, colored fermions?
- What is mass? Why are particles so light?
- Is there a link between particle and nucleon masses?
- How does all of this reconcile with gravitation?
How many space-time dimensions are there really?
- ...

The Plot

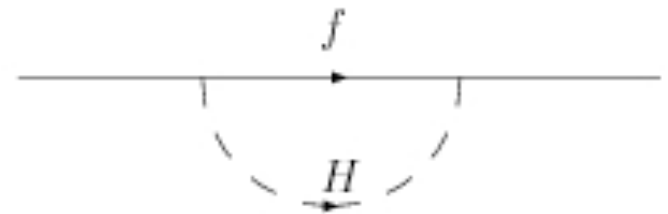
Vector Boson Scattering

- There is in fact one known problem with the Standard Model:
 - If we collide W's and/or Z's (not so easy...), the scattering cross-section grows with the center of mass energy, and gets out of control at about 1.7 TeV
- This is similar to “low” energy neutrino scattering:
 - If $q^2 \ll (M_W)^2$, looks like a “contact interaction”, and cross-section grows with center of mass energy
 - But when $q^2 \approx (M_W)^2$, W-boson propagation becomes visible, and “cures” this problem



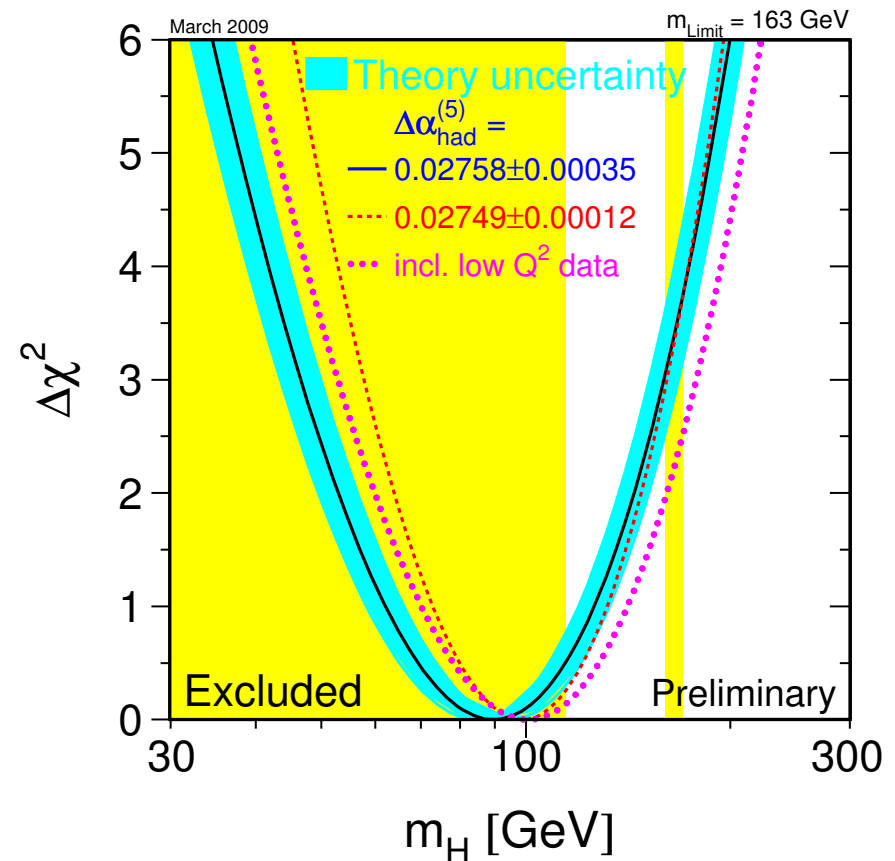
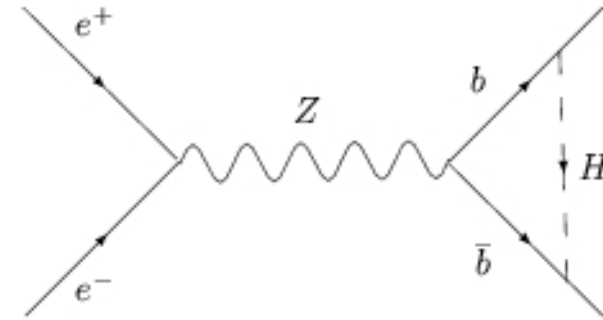
The Higgs Boson

- One way to solve this, is to introduce a massive, spinless particle (of mass $< \sim 1$ TeV)
- Couplings to W and Z are fixed, quantum numbers are known...
- to be those of the vacuum
- Its mass is unknown, and its couplings to the fermions are unknown.... well, maybe
- Fermions can acquire mass by coupling to this Higgs boson, so their couplings could be proportional to their masses. This is called the “Standard Model Higgs”



Precision Measurements

- In fact, we can say something about the standard model Higgs mass
- If the fermions get their masses from the Higgs, we know all couplings and can infer the Higgs mass from precision measurements
- Result is very sensitive to measured top quark, W boson masses
- Really wants a “light” Higgs boson

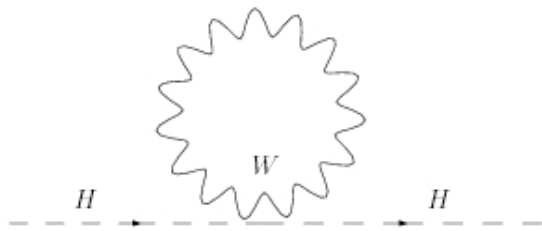


Higgs Drawbacks

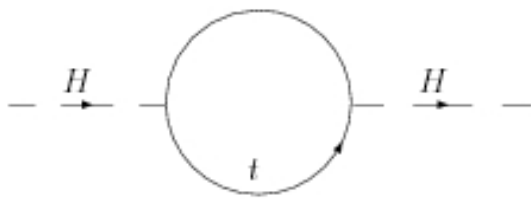
- In principle, with the addition of a Higgs boson around 150 GeV particle physics could be “complete”
 - Like Mendeleev’s table for chemistry
- But by itself, the Higgs is very unsatisfactory:
 - Why are the couplings to the fermions what they are?
 - Dumb luck (aka landscape)?
 - What is the link to gravity?
 - Why does the Higgs break the symmetry?
 - Why are there 3....?

The Plot Thickens

Higgs Mass



$$\longrightarrow \frac{1}{16\pi^2} g^2 E^2$$



$$\longrightarrow \frac{3}{16\pi^2} y_t^2 E^2$$



$$\longrightarrow \frac{1}{16\pi^2} \lambda E^2$$

- Higgs, in fact, also acquires mass from coupling to W's, fermions, and itself!
- These “mass terms” are quadratically divergent
- Drive mass to limit of validity of the theory
- So we expect the Higgs mass to be close to the scale where new physics comes in....

Unravelling the Mystery

Hunting for Answers

- Get more information
 - Measure particles and their interactions in detail
 - Precision measurements
 - Observe new particles or interactions
 - Search in new areas in “phase space”
- Find the underlying pattern(s)
 - Hypothesize, build models
 - Internally consistent? Consistent with data?
 - Suggestions on where to look

Experiment

Theory

Where to Start?

- BSM physics **must** couple to SM (weakly?), but is it
 - “SM-like”?
 - Does it have new massive particles decaying to electrons, muons, quarks,...?
 - Quasi “SM-like”?
 - Same but includes some new long-lived particles in the decay chain... (dark matter candidate)
 - No new “particles” in reach
 - Hidden or too heavy or.... don’t exist
 - Are there new interactions?

Dark Matter??

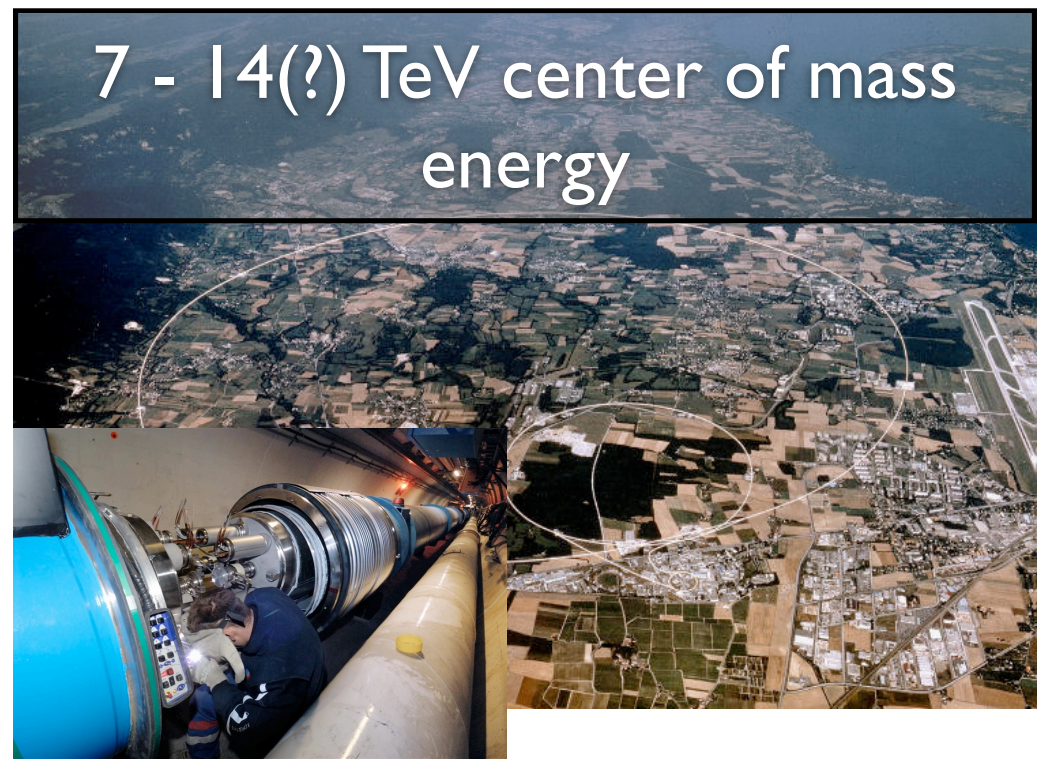
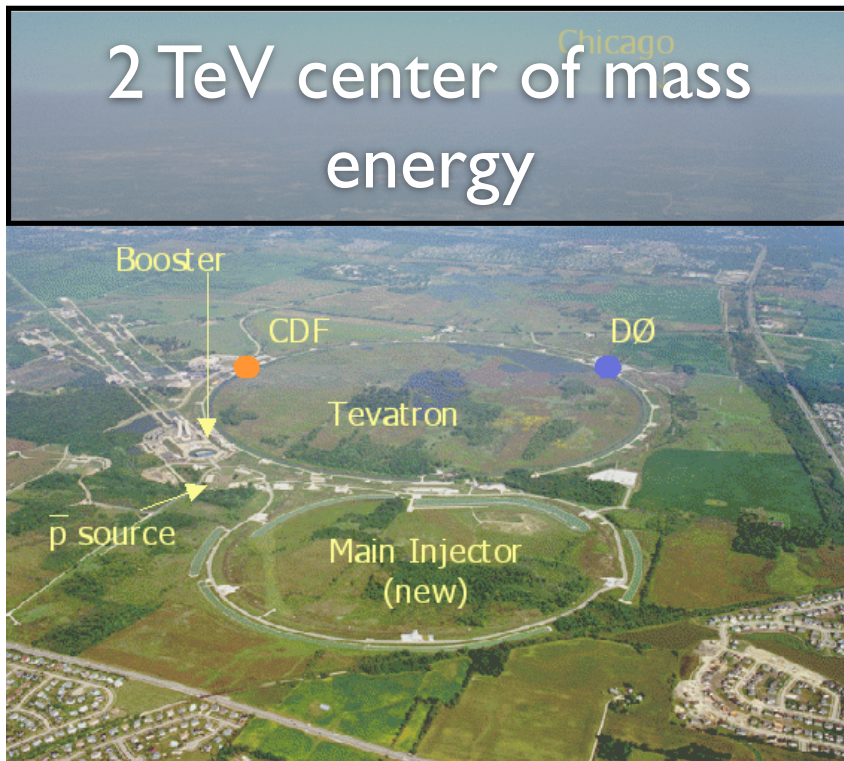
So....

- Go look where the SM breaks down (high energy)...
- ... or for subtle anomalies
- Assume new physics manifestations lead to anomalous production of SM particles
 - Resonant or not (and maybe in loops only)
 - Short-lived or less so
- Rely on guidance from models to some extent
 - What are implications of known constraints? What signatures are “allowed”?
 - Some scenarios do require new approach

The Tools

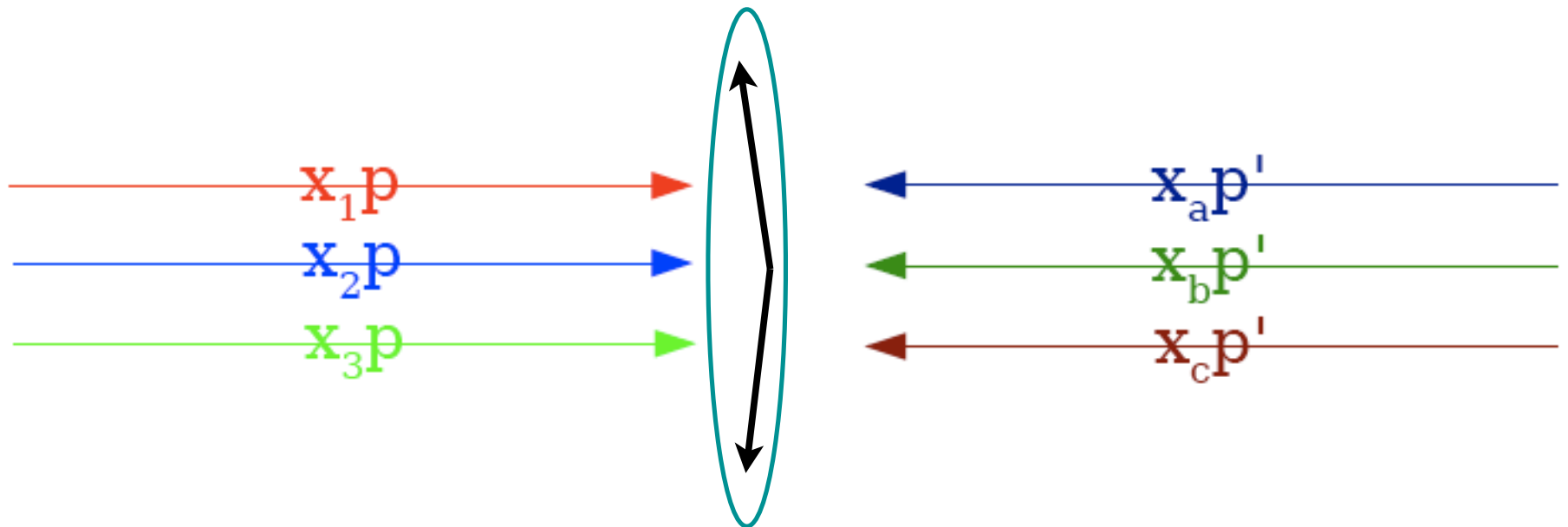
Colliders

- Currently, hadron colliders:
 - High energy implies probing of short distances, and production of other, massive particles



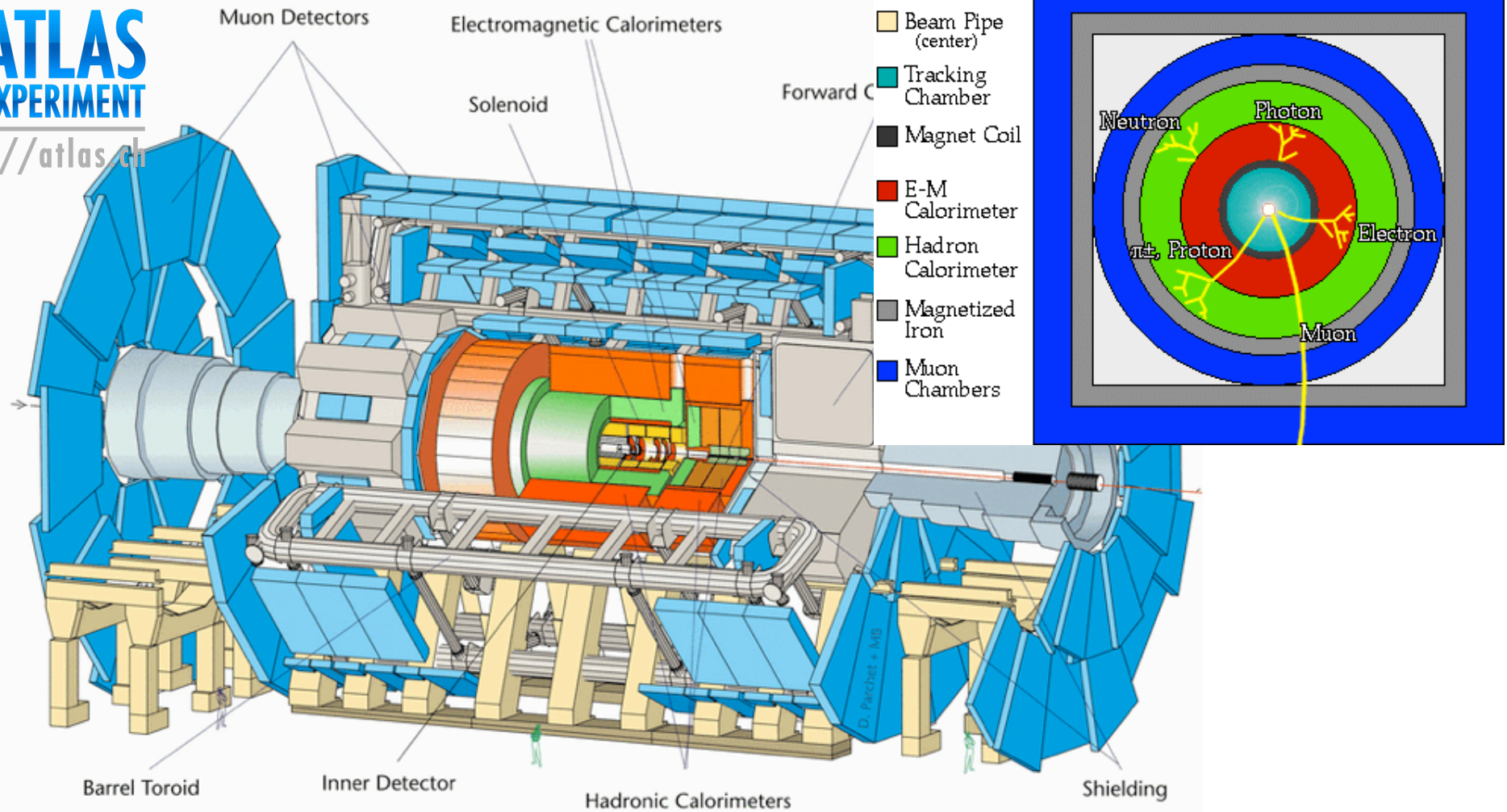
Hadron Colliders

- Incoming longitudinal momentum not known:
 - “Hard interaction” is between one of the quarks and/or gluons from each proton, other quarks/gluons are “spectators”
- Longitudinal boost “flattens” event to a pancake
- We usually work in the plane transverse to the beam



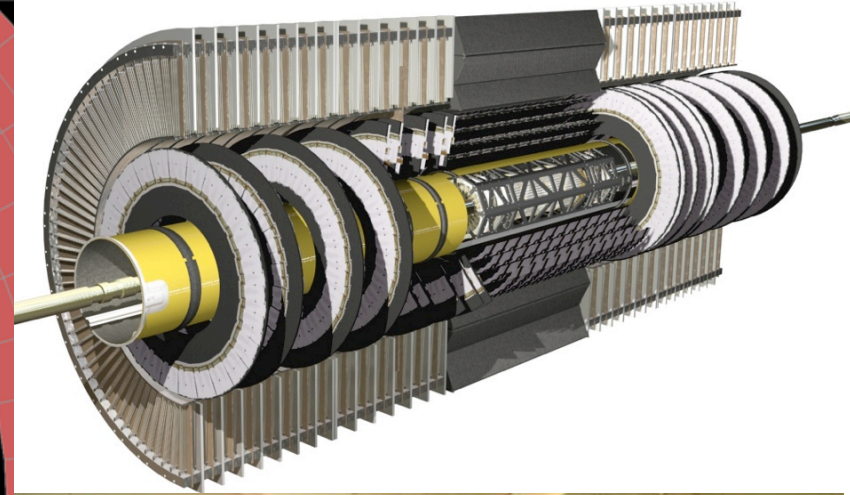
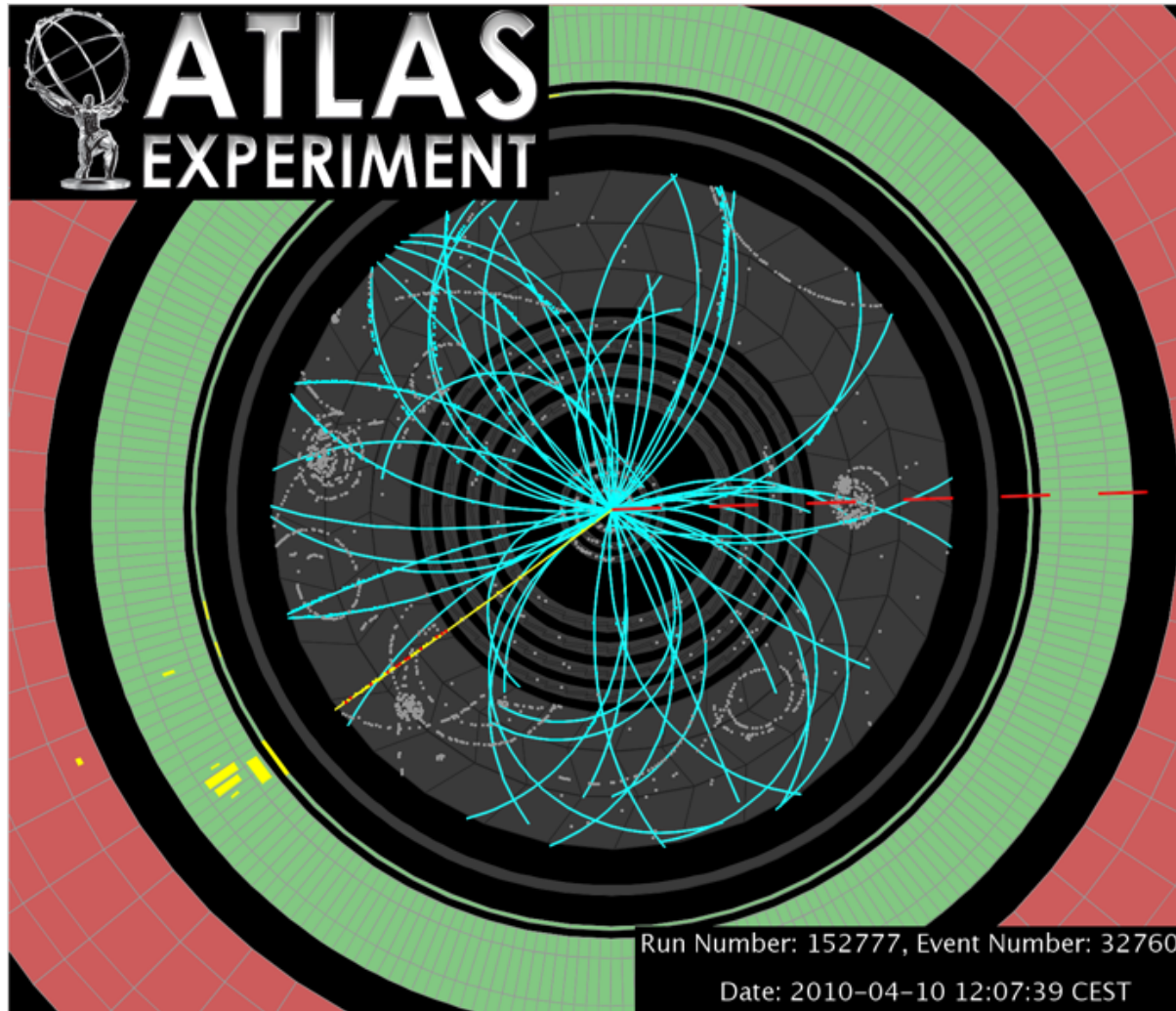
ATLAS

- Make best possible measurement of all particles coming out of collisions



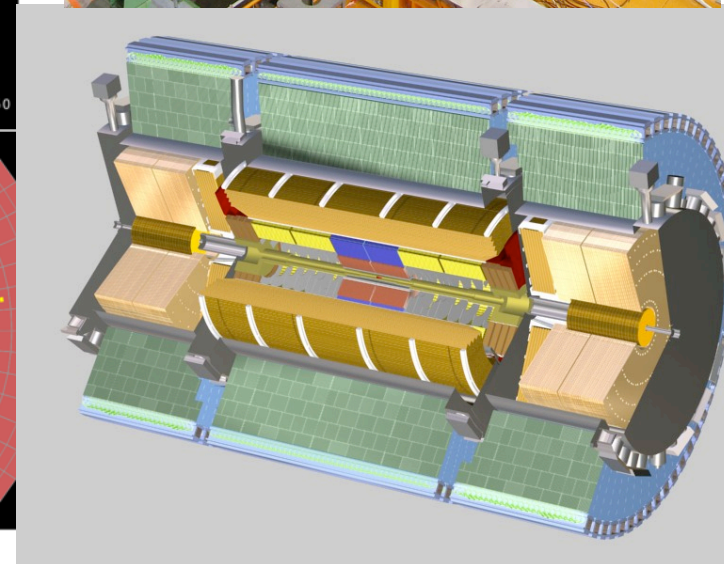
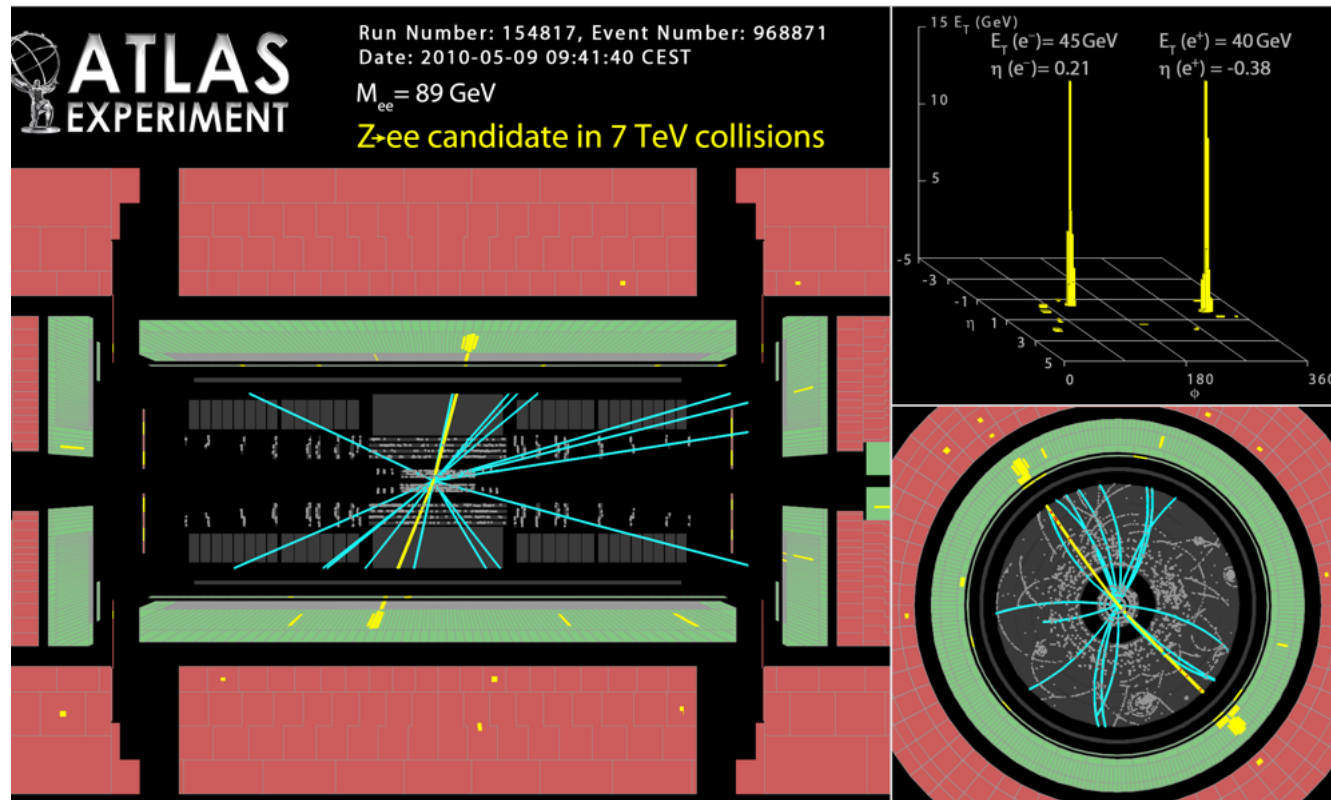
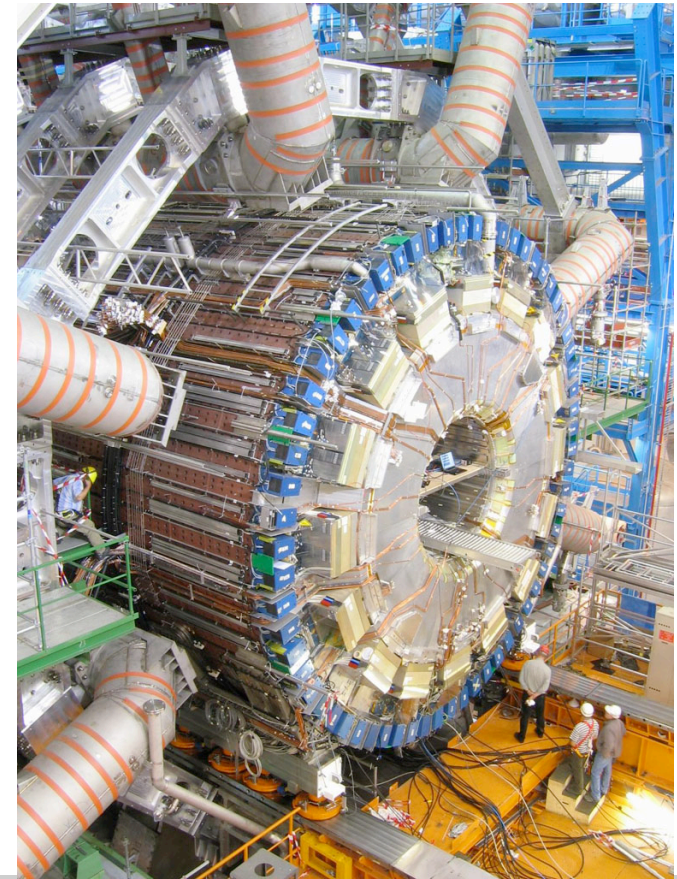
Tracking

- Combination of pixels, silicon strips (“SCT”) and straw tube transition radiation tracker (TRT)



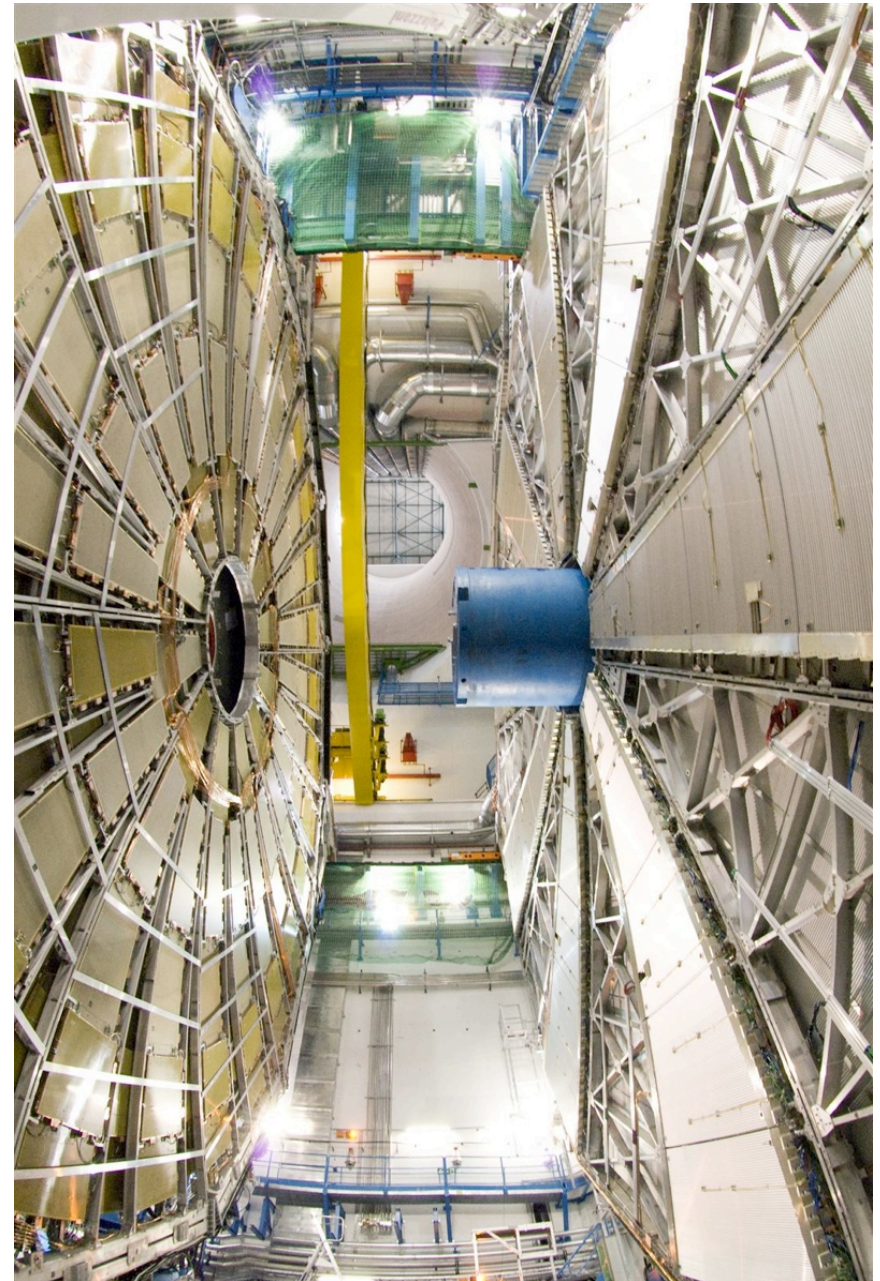
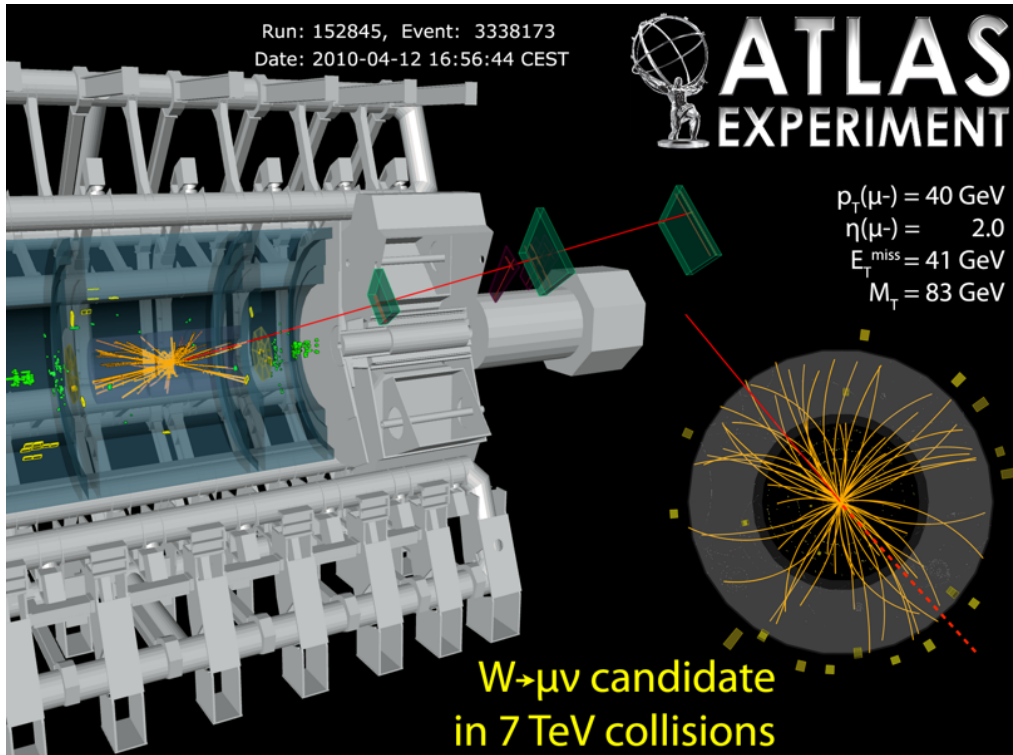
Calorimetry

- Liquid Argon & Pb accordion (EM & forward)
- Scintillator & Pb (hadronic)



Muons

- Air-core toroids
- Makes ATLAS big



Neutrinos*

*(100% acceptance)

Detecting Particles

		3 Generations of Fermions			Force Carriers	
Q u a r k s	2/3	u ✓ ~5	c ✓ ~1350	t ✓ 175000	g ✓ 0	Strong Interactions
	-1/3	d ✓ ~9	s ✓ ~175	b ✓ ~4500	γ ✓ 0	Electro-magnetism
L e p t o n s	0?	ν _e ✓ 0?	ν _μ ✓ 0?	ν _τ ✓ 0?	Z ⁰ ✓ 91187	Weak Interactions
		e ✓ 0.511	μ ✓ 105.66	τ ✓ 1777.2	W [±] ✓ 81400	

Masses are in MeV

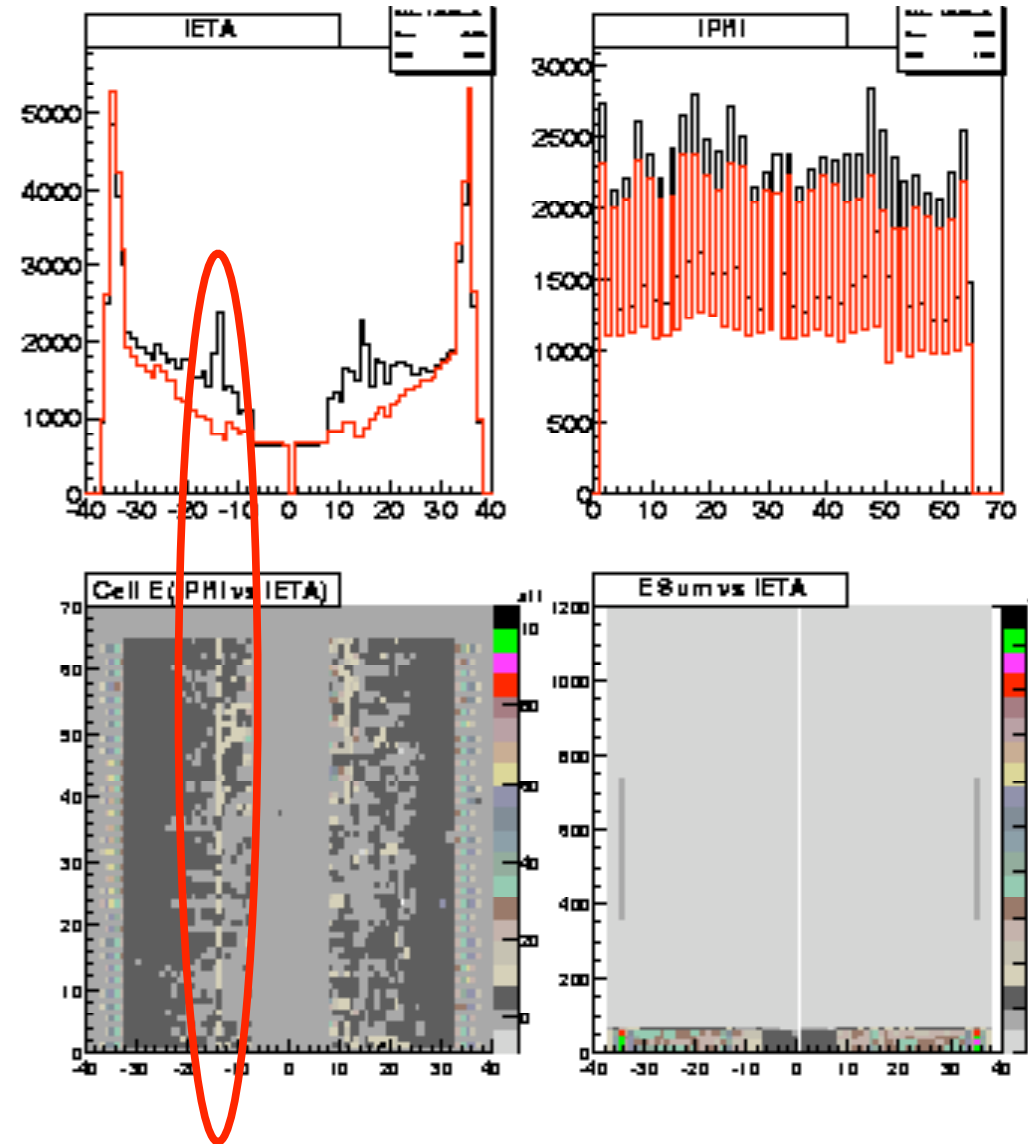
✓ : Detect with high efficiency

✓ : Detect by missing transverse energy

✓ : Detect through decays: $t \rightarrow Wb, W/Z \rightarrow$ leptons

Anecdotes From the Field (I)

- DØ's “ring of fire”
 - Noise in a few “eta rings”
 - Occurred on rare occasions
- Originally thought to be a ground fault in HV distribution
- Found to be concurrent with welding in building
 - Finally traced to liquid Argon purity & temperature monitoring



The Work

Steps in a Physics Search

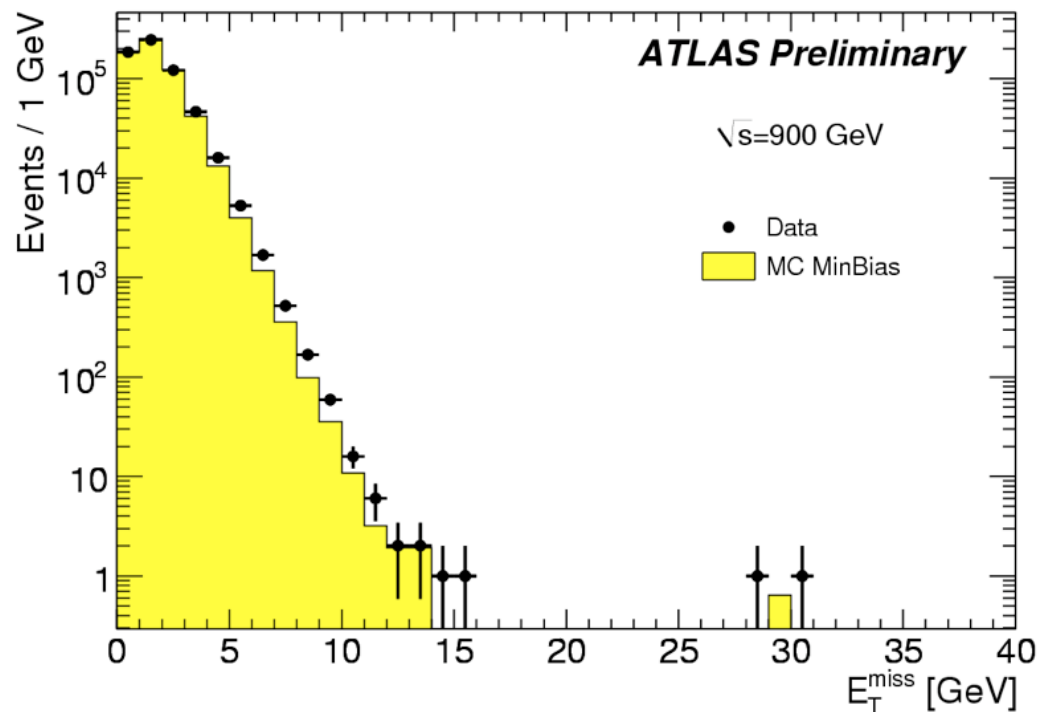
- What is the final state? \Rightarrow “Preselection”
 - Sufficiently loose to be signal-poor
 - Prove you understand the detector response, physics processes contributing
 - But sufficiently tight to have a manageable data volume
 - ATLAS/CMS write 200-400 Hz \times 1+ MB/event = 200-400+ MB/s
 - “4-vectors” is not enough, need some amount of detector info
 - In practice, often have preselected sample for frequent analysis, + looser sample for e.g. multijet background with rare passes
- Note that data volume \sim running time, not $\int \mathcal{L}$

Steps (II)

- Determine preselected sample's composition
 - MC and data to understand each contribution
 - QCD multijet background to leptons often extracted from data: rejection factor $\sim 10^{-4}$, difficult for simulation to be that accurate
 - MC for most other processes, with corrections from data, since generators are LO or NLO
 - Also need to correct MC for real-life data conditions
 - Different alignment, small fraction of dead channels etc.
 - As statistics increase, more difficult, since mis-modelings not hidden by large statistical uncertainties anymore

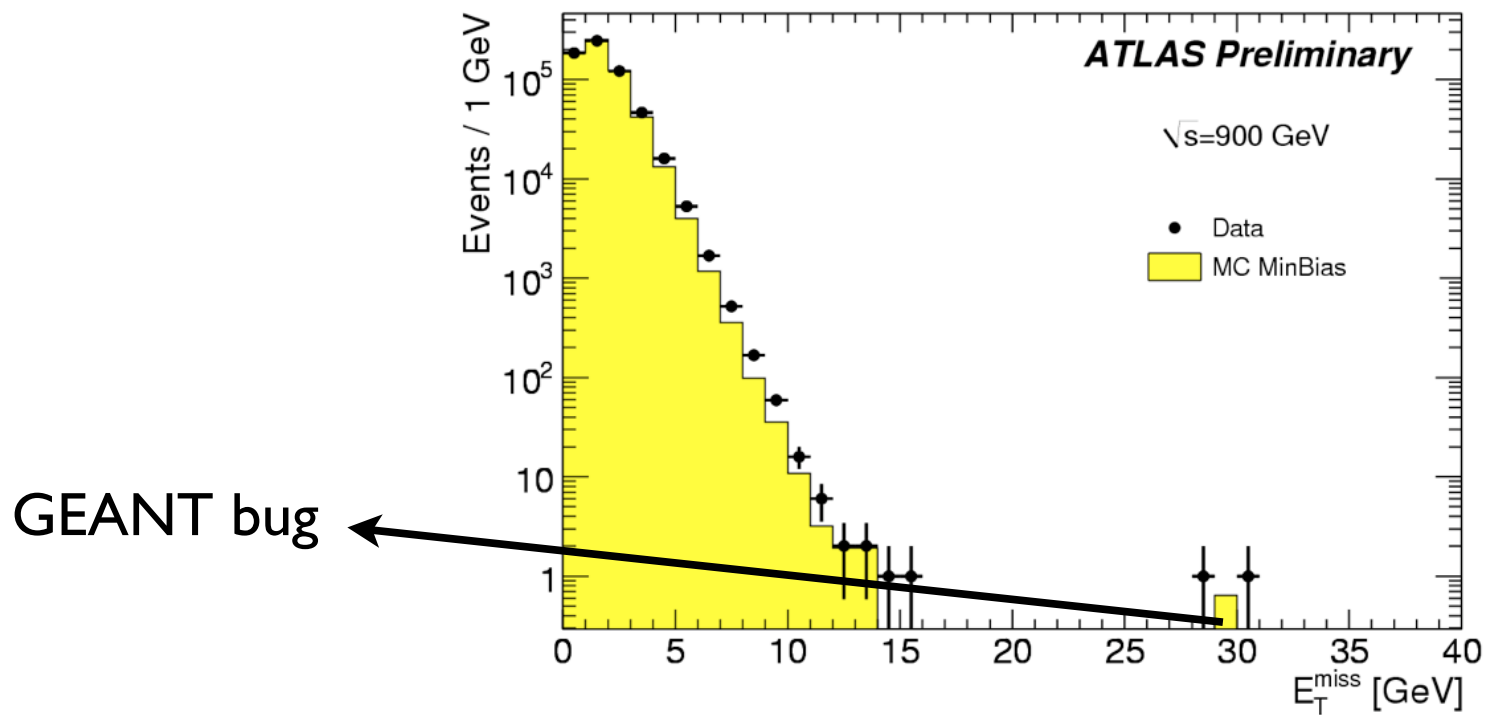
Anecdotes From the Field (II)

- Everybody wants experimenters to produce results fast
- Lots of pressure in the early days of LHC...
 - Only jets, background composition “easy”



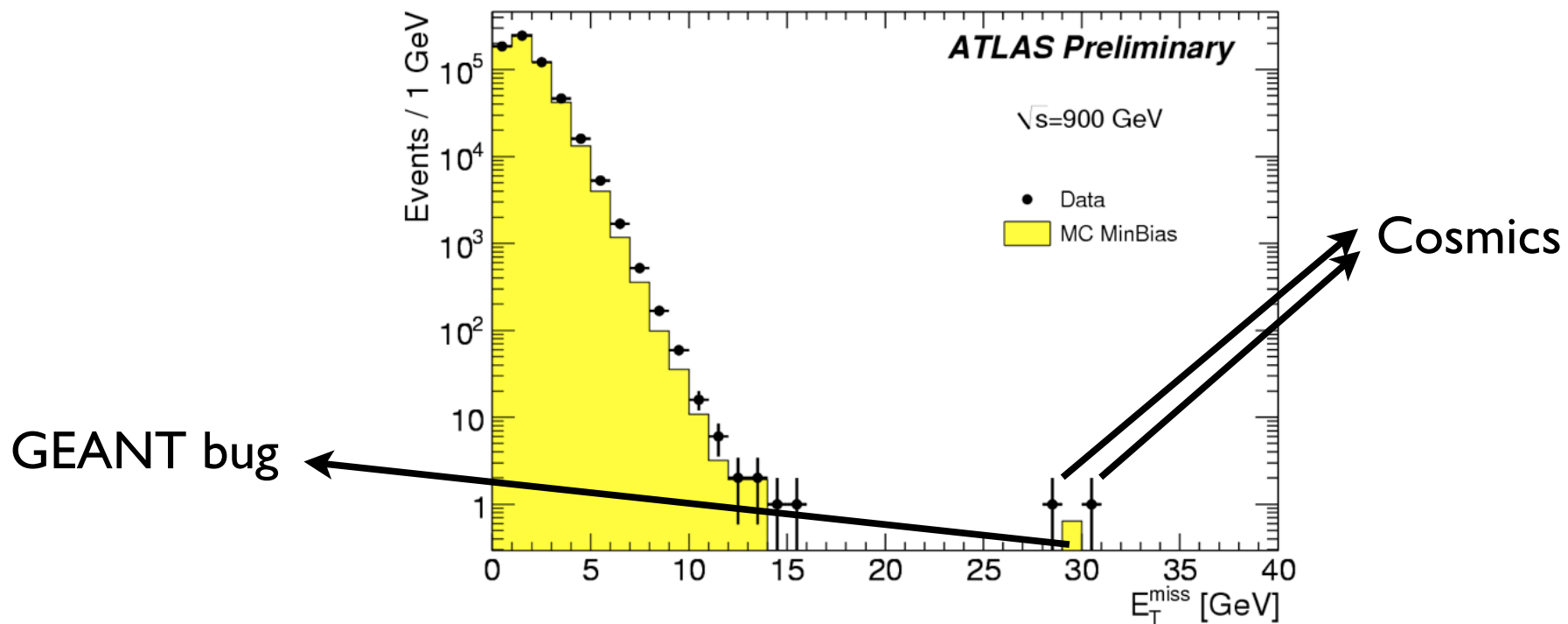
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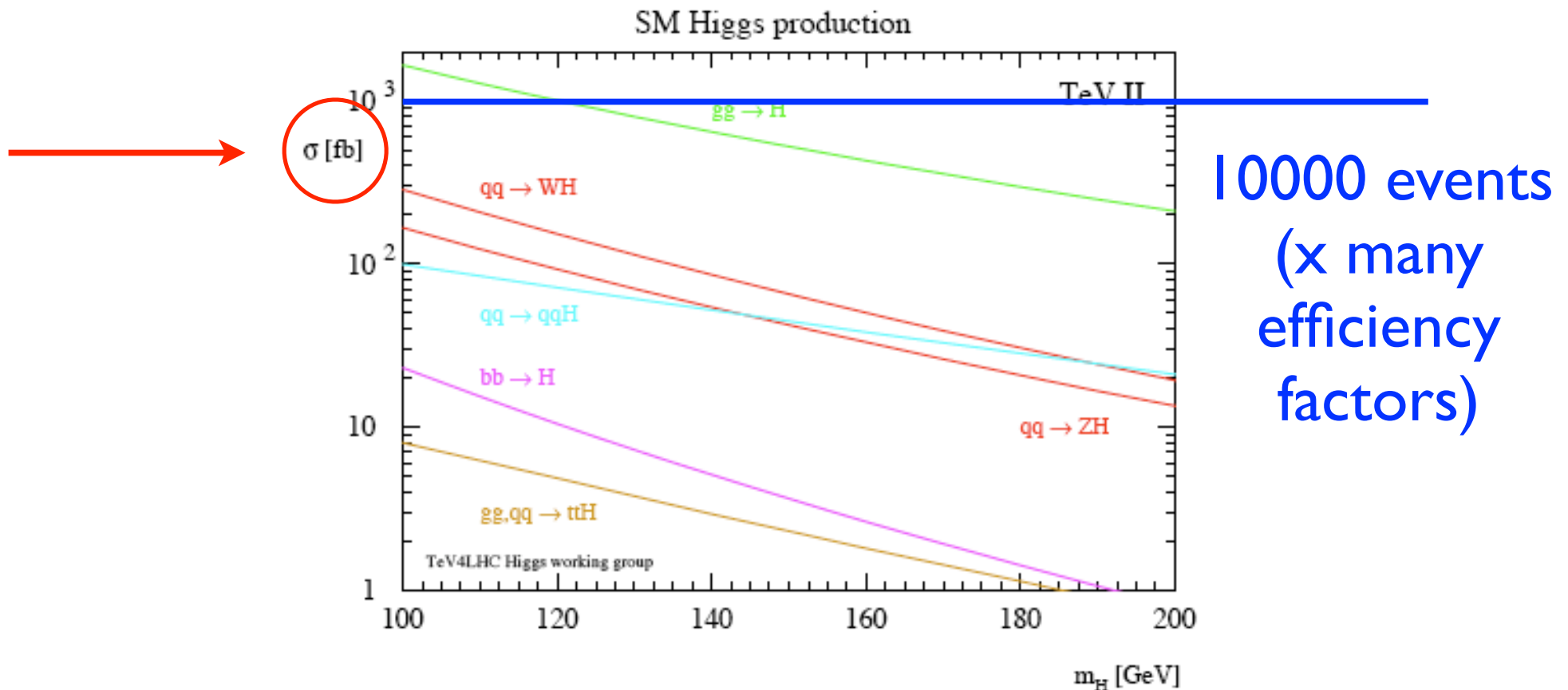


- Sometimes, it's better to take the appropriate time to investigate

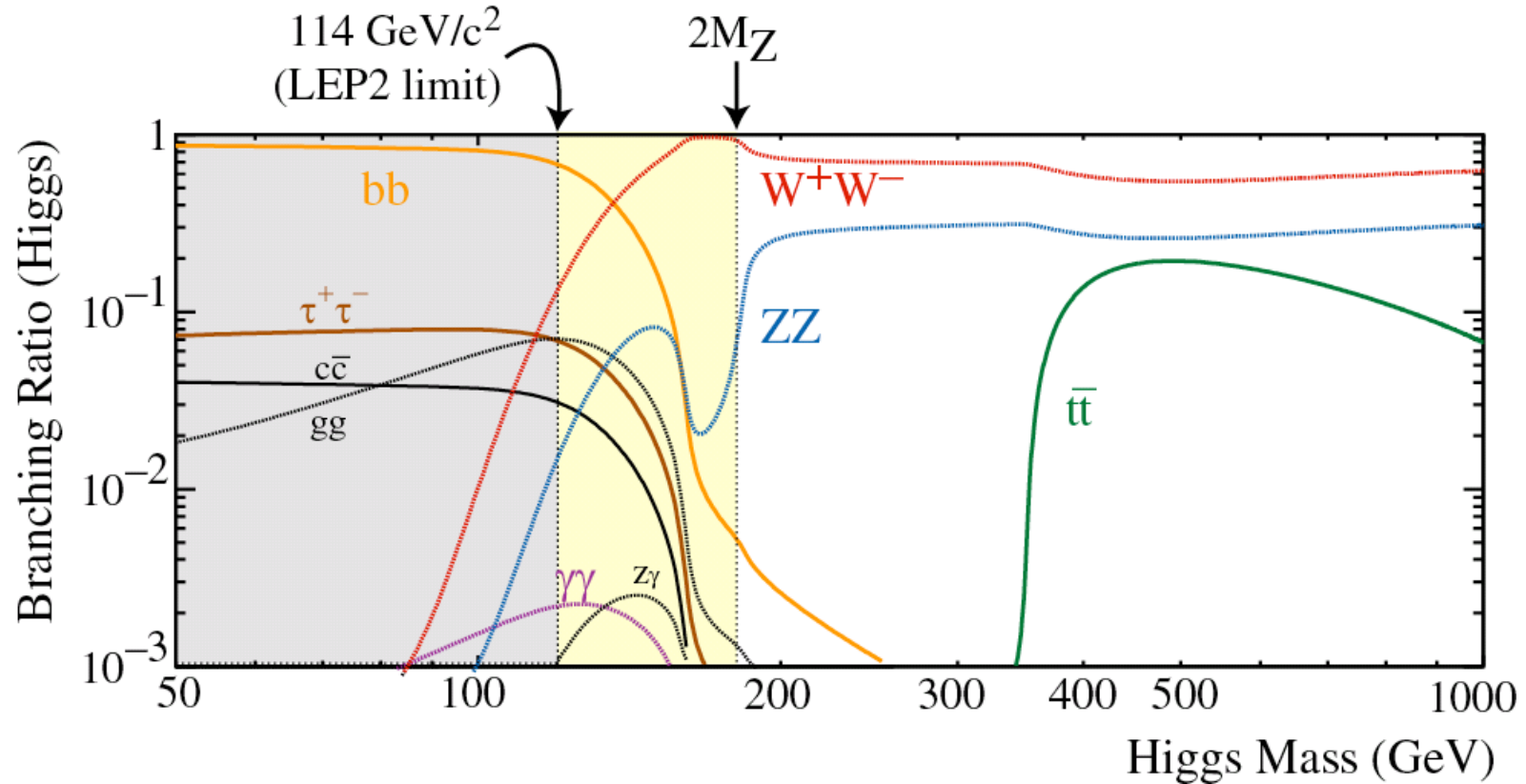
A Challenging Search: The Higgs Hunt at the Tevatron

Producing Higgses

- Tevatron experiments currently have $\sim 10 \text{ fb}^{-1}$ of data on tape
- (Data taking efficiency is $\sim 90\%$)



Higgs Decay

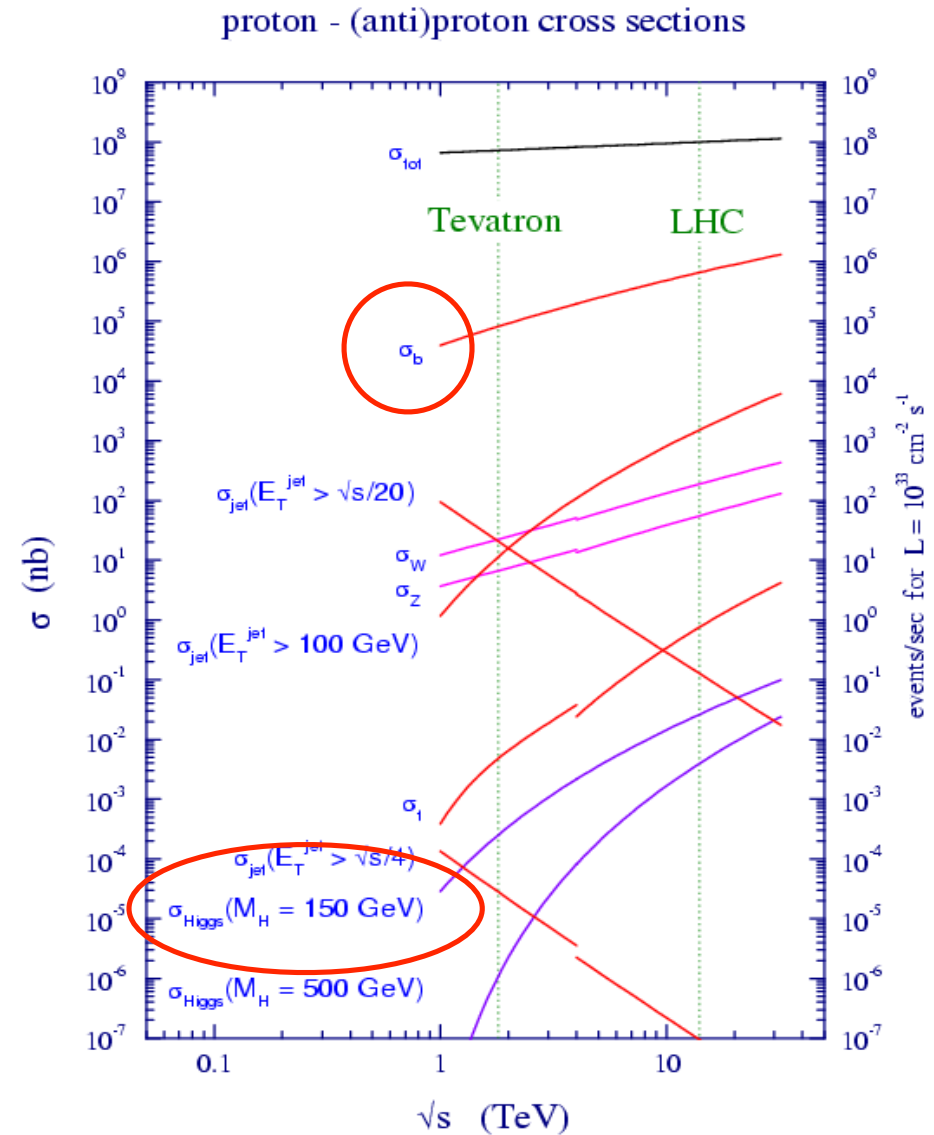


Low Mass
 $H \rightarrow b\bar{b}$

High Mass
 $H \rightarrow WW$

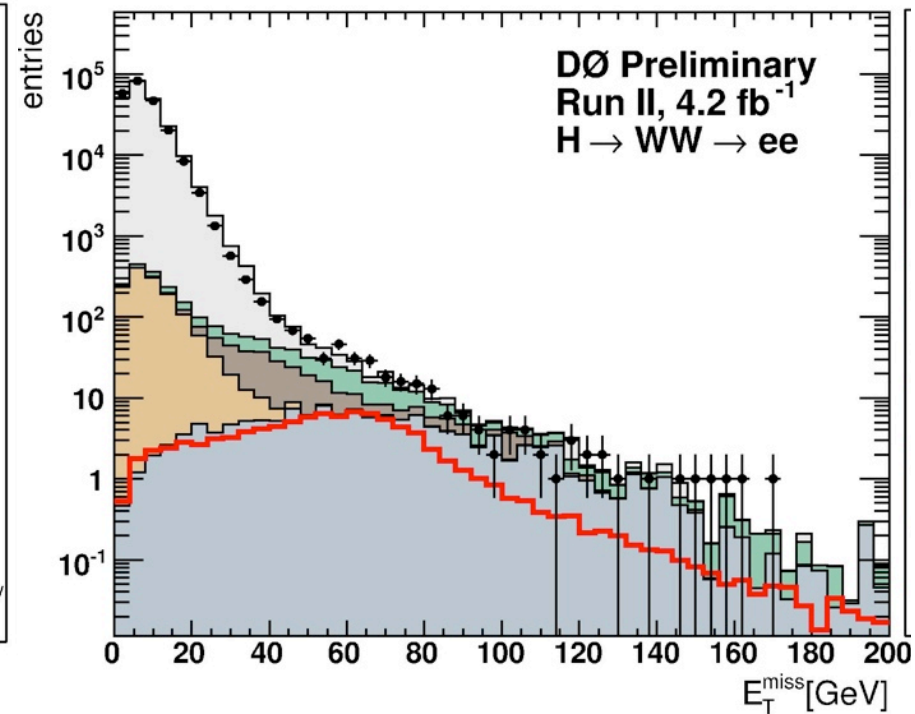
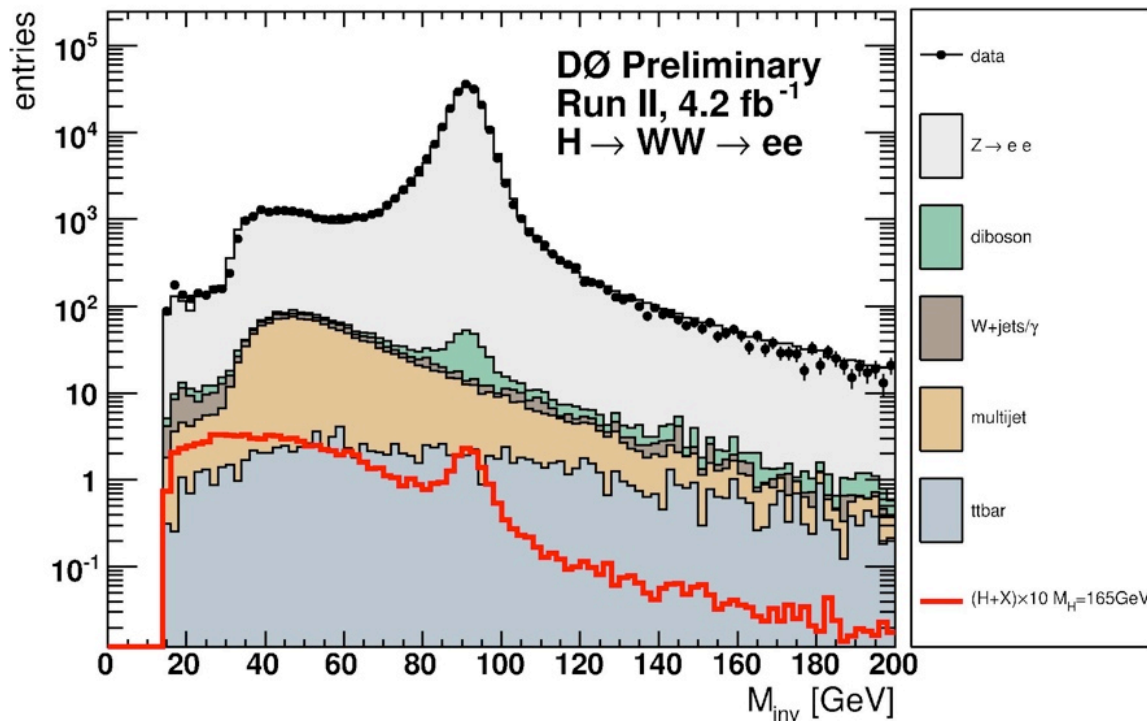
Search Channels

- Hadron colliders
 - $b\bar{b}$ production ~ 9 orders of magnitude larger than H
 - $gg \rightarrow H \rightarrow b\bar{b}$ swamped
- ➔ At low mass look for $pp \rightarrow WH$ or $ZH \rightarrow W/Z b\bar{b}$
- With leptonic W, Z decay, so # of events $\sim 50!$
- At high mass, $gg \rightarrow H \rightarrow WW$ accessible if at least one W decays leptonically



Dilepton + MET

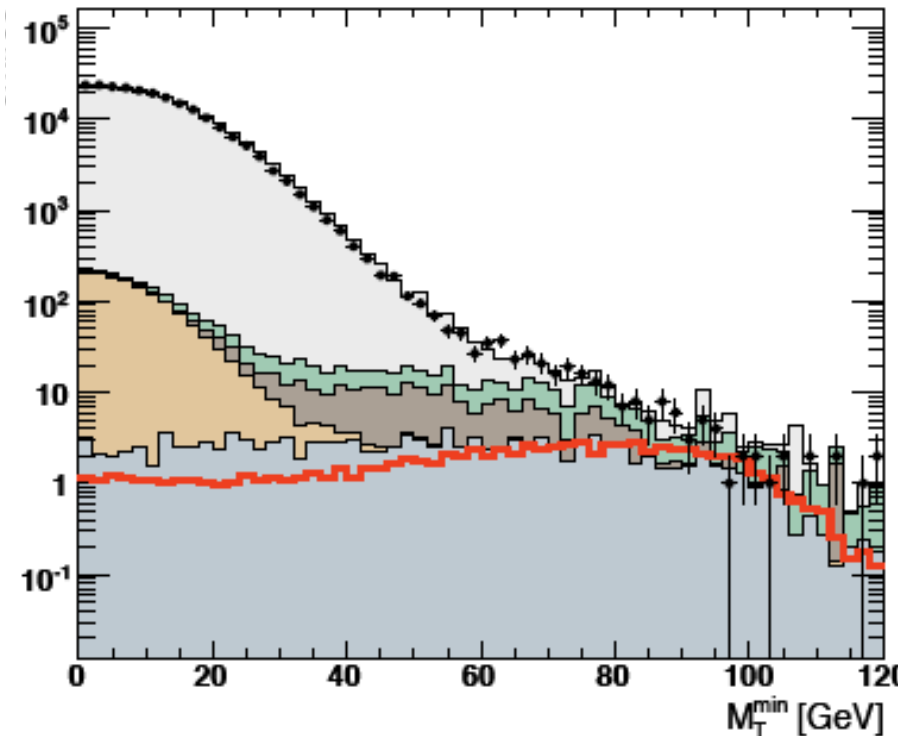
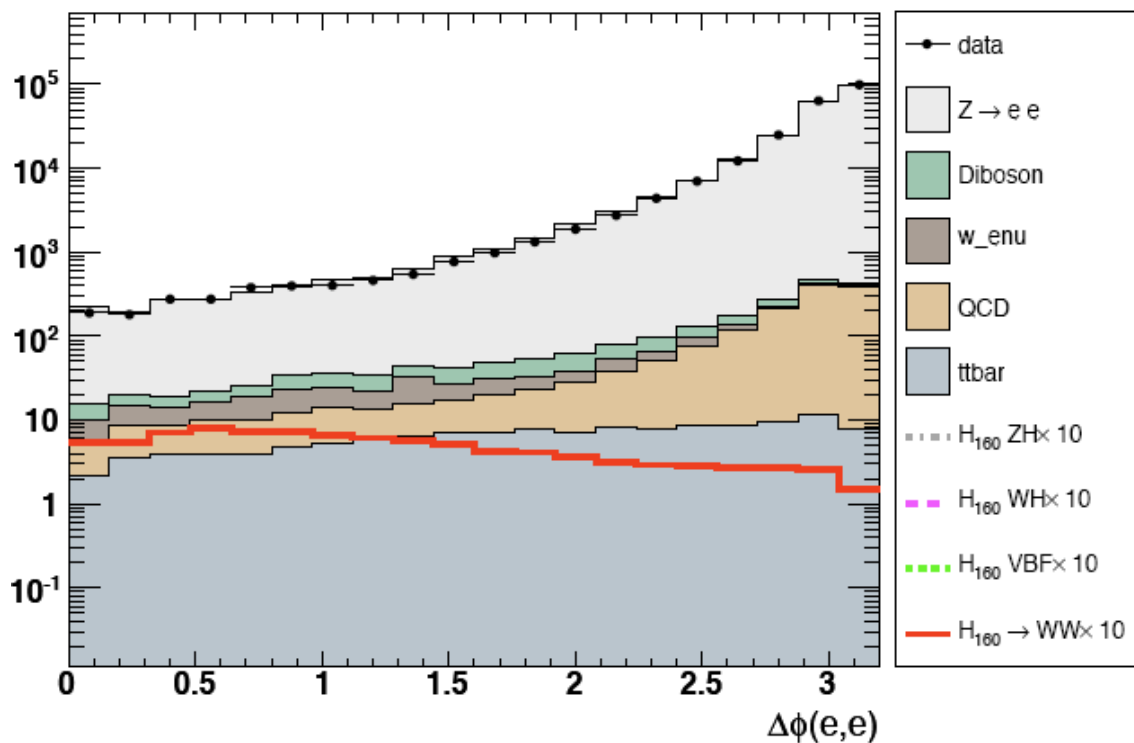
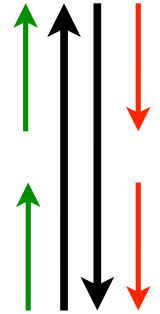
- “Golden” channel:
 - Main background $Z \rightarrow \ell\ell$ also a great reference signal
 - “Easy” to suppress using MET, angle between leptons, ...



Angles

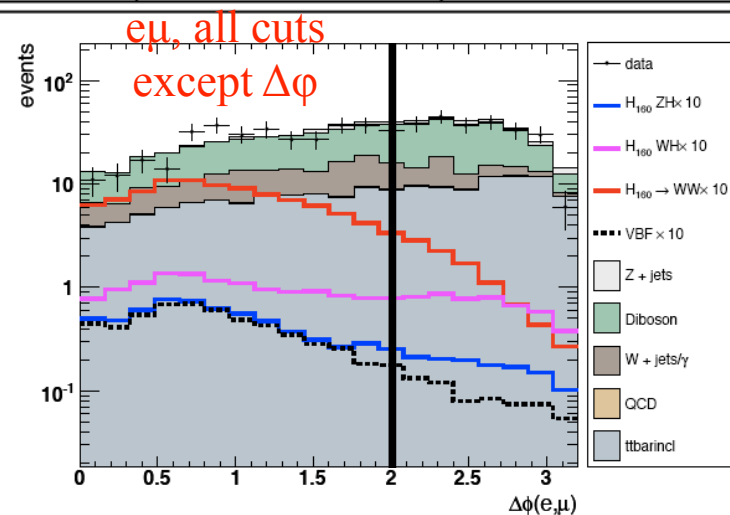
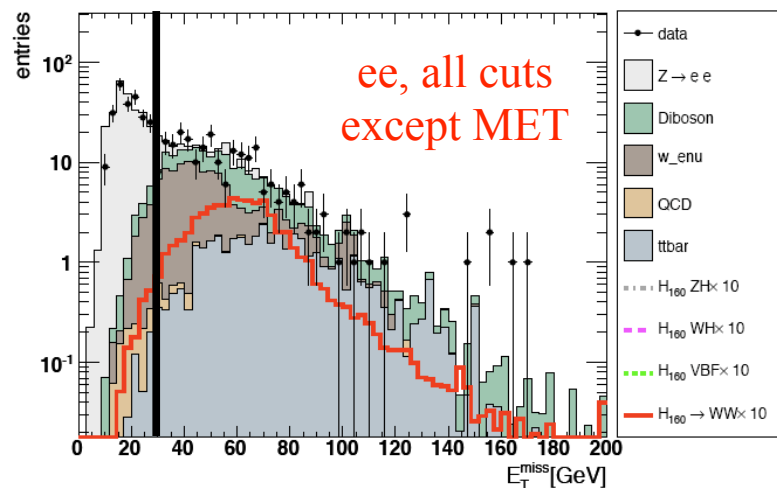
- In $Z \rightarrow ll$ (and dijets faking leptons), leptons preferentially emitted back-to-back
- In Higgs decays, W^+W^- spins back-to back, so charged leptons in similar direction! (One LH, other RH)
- In Z , smallest transverse mass tends to be small

Spins



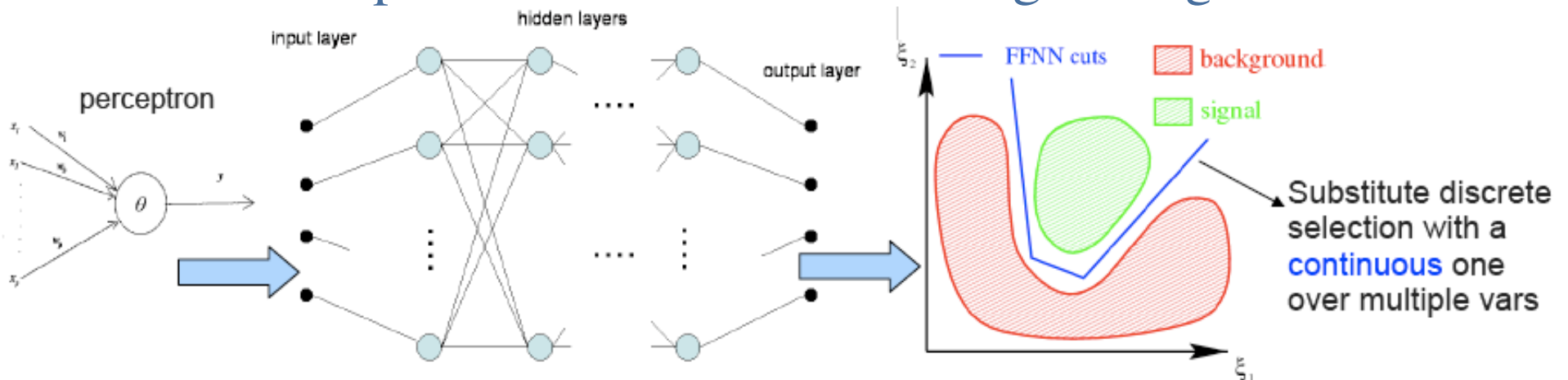
Preselection

Final state	$e\mu$	ee	$\mu\mu$
Cut 0 Pre-selection	lepton ID, leptons with opposite charge and $p_T^\mu > 10$ GeV and $p_T^e > 15$ GeV invariant mass $M_{\ell\ell} > 15$ GeV $\mu\mu$: $n_{\text{jet}} < 2$ for $p_T^{\text{jet}} > 15$ GeV, $\Delta\mathcal{R}(\mu, \text{jet}) > 0.1$ and $p_T^\mu > 15$ GeV for the leading μ		
Cut 1 Missing Transverse Energy \cancel{E}_T (GeV)	> 20	> 20	
Cut 2 $\cancel{E}_T^{\text{Scaled}}$	> 6	> 6	
Cut 3 $M_T^{\text{min}}(\ell, \cancel{E}_T)$ (GeV)	> 20	> 30	
Cut 4 $p_T^{\mu\mu}$ (GeV) for $n_{\text{jet}} = 0$ \cancel{E}_T (GeV) for $n_{\text{jet}} = 1$			> 20 > 20
Cut 5 $\Delta\phi(\ell, \ell)$	< 2.0	< 2.0	< 2.5



Multivariate Tools

- After preselection, S/B not good ($\sim 1/30$, $1/50$, $1/1000$ in $e\mu$, ee and $\mu\mu$ final states)
- Use multivariate tools to exploit correlations between observables for $S \leftrightarrow B$ discrimination
- In the dilepton + MET ($H \rightarrow WW \rightarrow \ell\nu\ell\nu$), use neural nets
- MC samples divided in 2 for training/testing



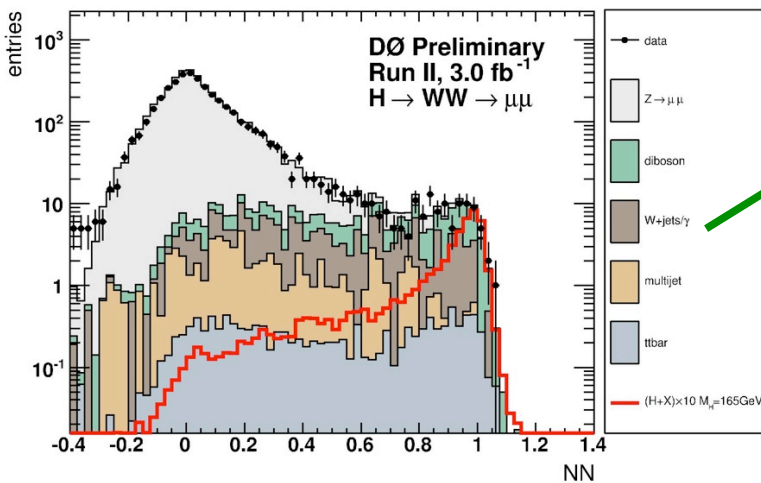
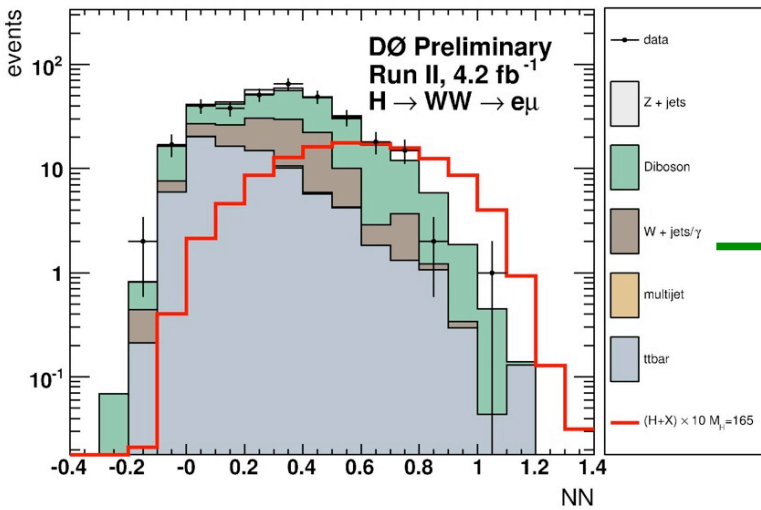
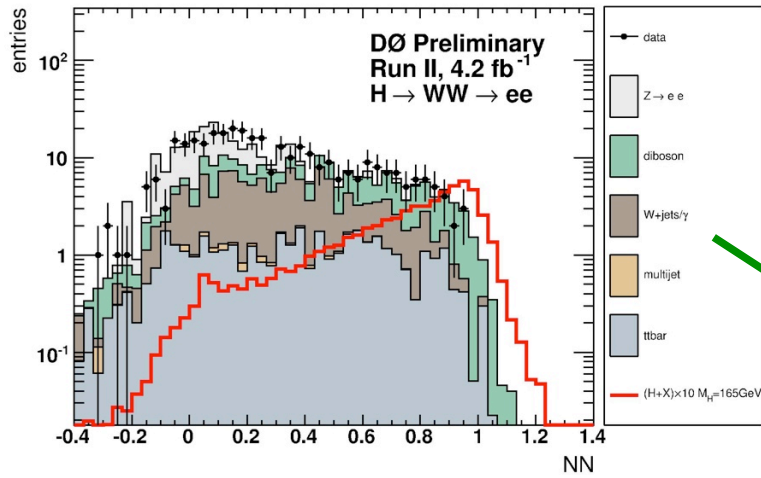
Variables

Only accept variables that are well-modeled!

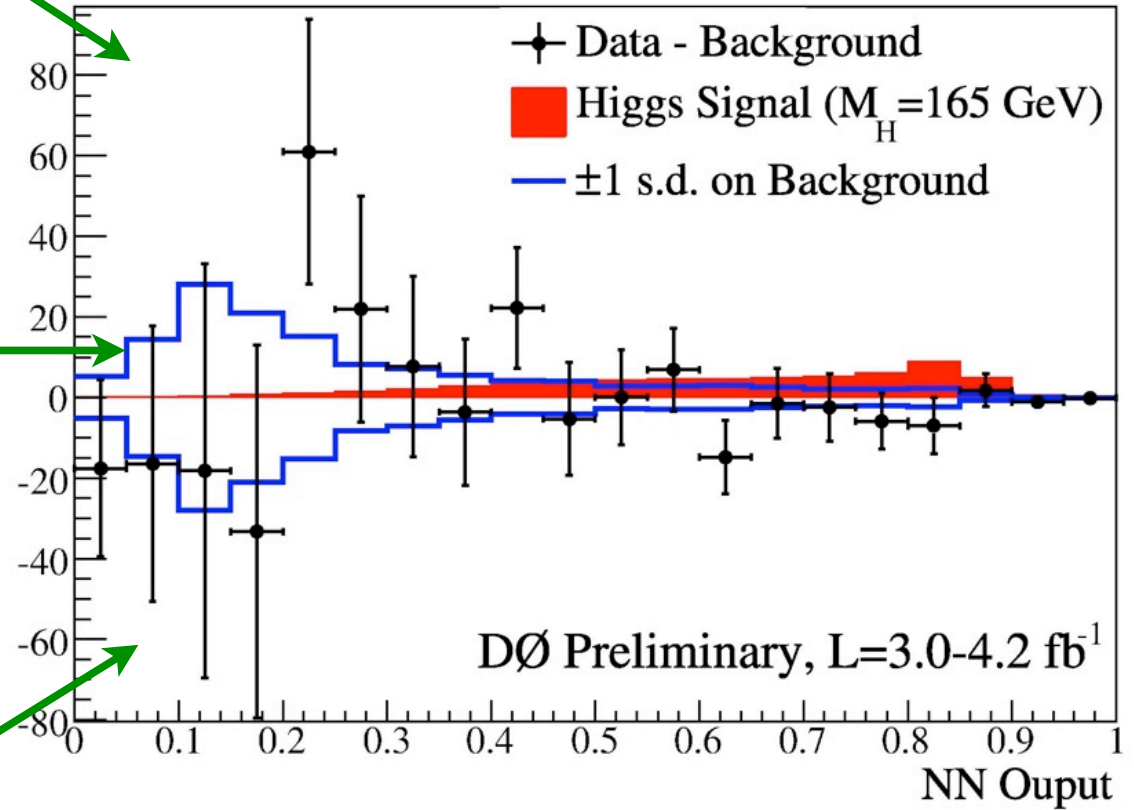
NN Analysis Variables

p_T of leading lepton	$p_T(\ell_1)$
p_T of trailing lepton	$p_T(\ell_2)$
Minimum of both lepton qualities	$\min(q_{\ell_1}, q_{\ell_2})$
Vector sum of the transverse momenta of the leptons:	$p_T(\ell_1) + p_T(\ell_2)$
Scalar sum of the transverse momenta of the jets:	$H_T = \sum_i p_T(\text{jet}_i) $
Invariant mass of both leptons	$M_{\text{inv}}(\ell_1, \ell_2)$
Minimal transverse mass of one lepton and \cancel{E}_T	M_T^{min}
Missing transverse energy	\cancel{E}_T
Scalar transverse energy	E_T^{scalar}
Azimuthal angle between selected leptons	$\Delta\phi(\ell_1, \ell_2)$
Solid angle between selected leptons ($e\mu$ only)	$\Delta\Theta(\ell_1, \ell_2)$
ΔR between selected leptons ($e\mu$ only)	$\Delta R(\ell_1, \ell_2)$
Azimuthal angle between leading lepton and \cancel{E}_T	$\Delta\phi(\cancel{E}_T, \ell_1)$
Azimuthal angle between trailing lepton and \cancel{E}_T	$\Delta\phi(\cancel{E}_T, \ell_2)$

NN Outputs



Events / 0.05

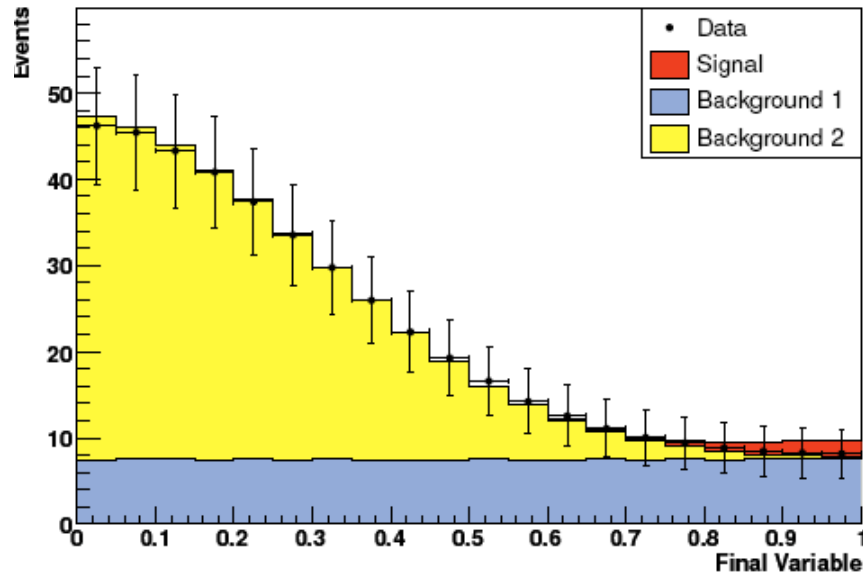


Systematics Profiling

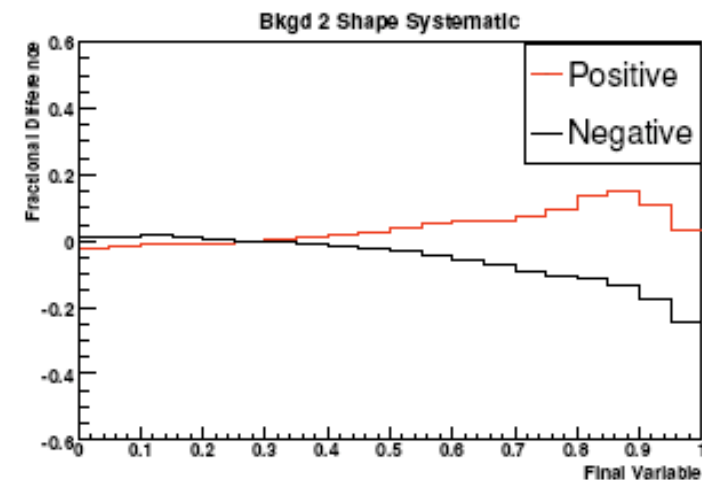
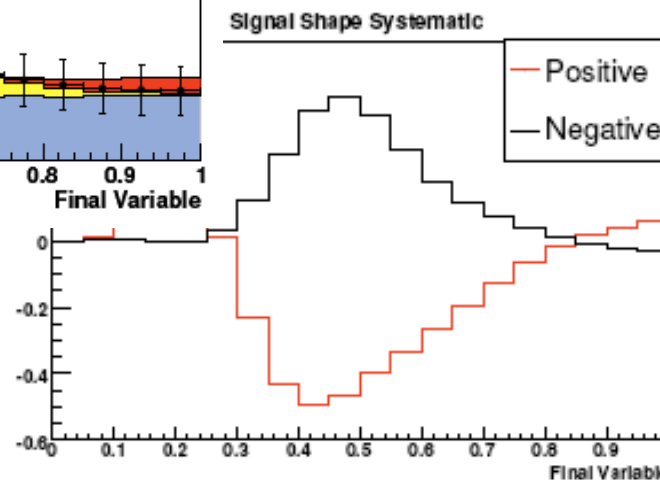
- Systematic uncertainties are propagated through the full analysis chain to the NN output distribution
- E.g. we repeat the analysis with jet energy scale shifted up & down by 1σ
- Some systematic uncertainties affect shape (jet reconstruction efficiency, energy scale and resolution, boson p^T distributions), others only normalization (lepton reconstruction efficiencies and momentum calibration, modeling of multijet background, theoretical cross-sections and luminosity)
- Systematic uncertainties are treated as nuisance parameters

Systematics Profiling

- Nuisance parameters tend to be correlated, but not 100%, among backgrounds
- Can affect rates, shapes, or both (in any distribution), and often asymmetric and non-gaussian

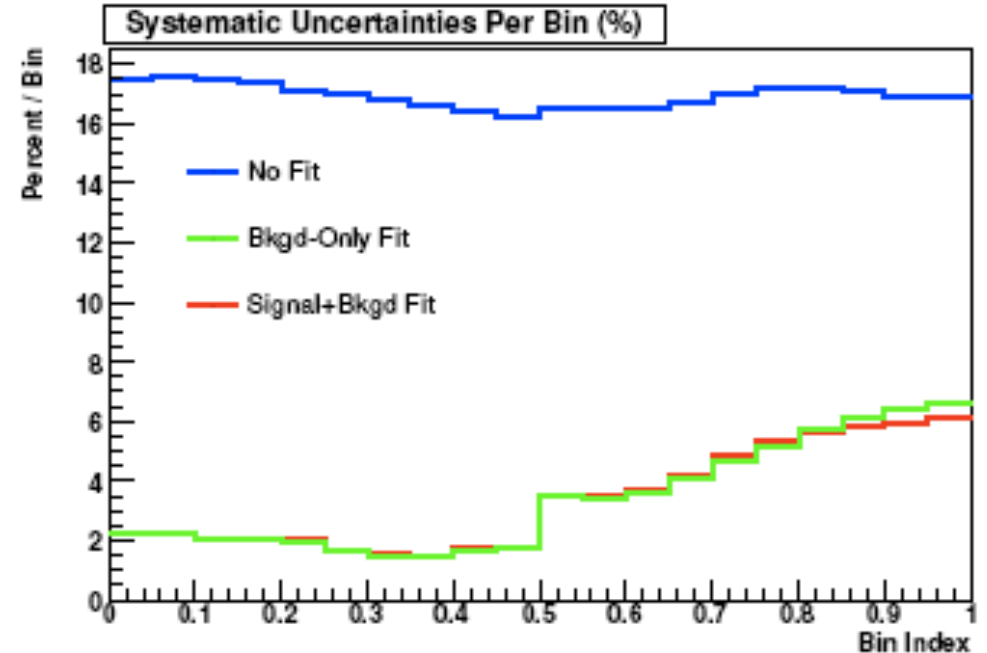
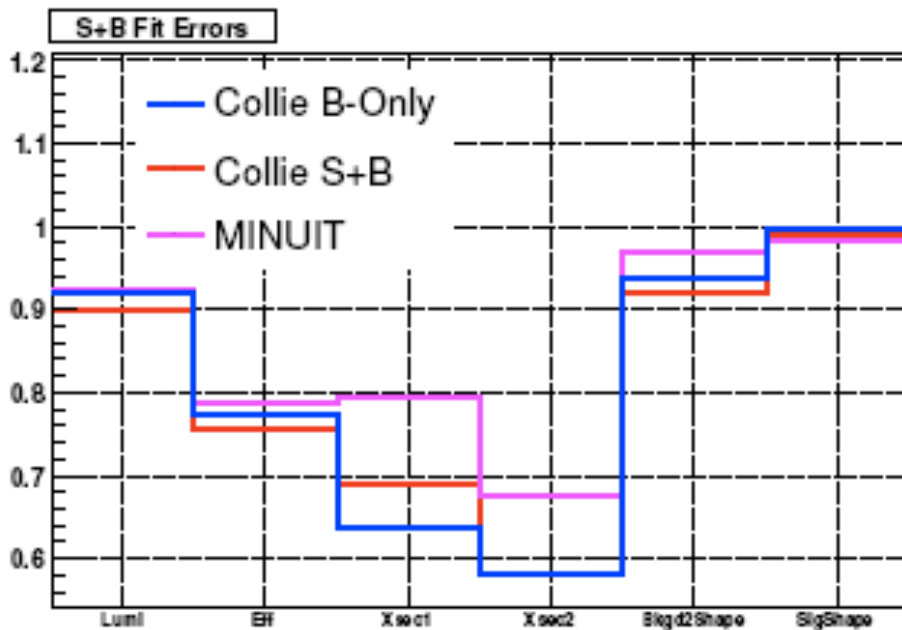


Toy Example



- Can generate pseudo-experiments (events in bins according to poisson), then for each experiment vary nuisance parameters
 - Variations in background (& S+B) prediction
 - Compare results to data using log-likelihood ratio
- So we can maximize likelihood ratio as a function of nuisance parameters → constraint them
 - I.e. use full shape of distribution(s) to see which background uncertainties are over/underestimated
 - Of course limited to size of statistical fluctuations
 - Can remove bins with large S/B if needed
 - Mostly important if uncertainties lead to similar shape distortions

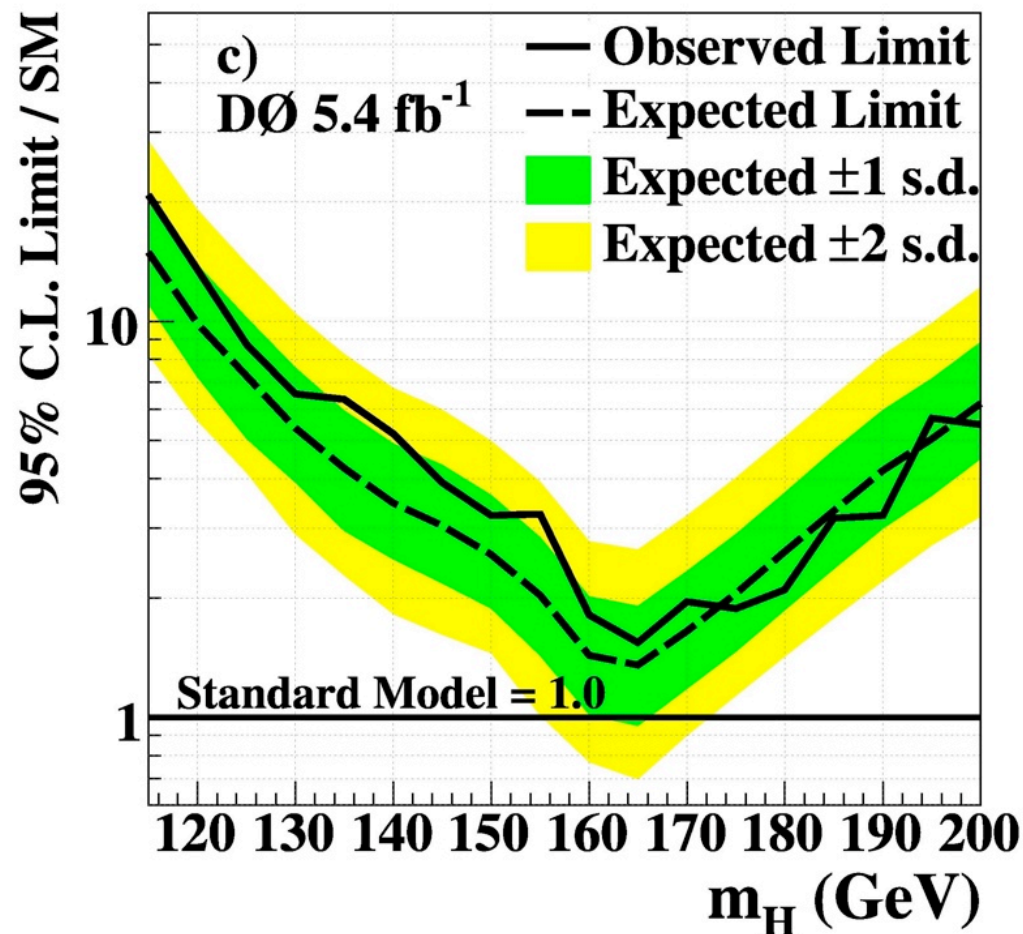
- Test example:
 - Data constructed to disagree with background-only hypothesis (wrong estimates for background uncertainties)
 - But to agree with background-only better than signal+background
 - Improvement quite spectacular (but by construction)



Dilepton + MET Result

- Present result as a 95% C.L. limit in units of the SM Higgs production x-section

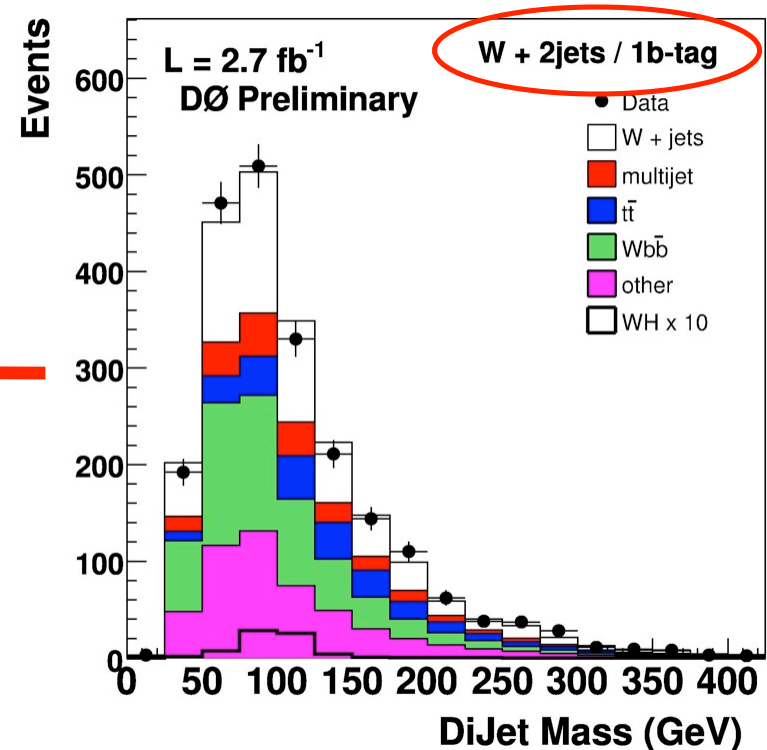
[Phys. Rev. Lett. 104, 061804 \(2010\)](#)




Wjj and the Higgs

- The final state consisting of $W + 2$ jets is critical
 - At low mass ($WH, H \rightarrow bb$), they're b-jets with $m_{bb} = m_H$
 - At high mass ($H \rightarrow WW$), $m_{jj} = m_W, m_{WW} = m_H$
- But dijet mass resolution is so-so:

And lots more background!



Sample Composition

- After preselection, low S/B allows to verify shapes of dominant backgrounds
 - For WH, first before b -tagging, then with 1 tag
 - Determining the sample's composition
 - I.e. which processes contribute, and how
- 
- Diboson from MC simulation (usually small, + “trust” MC)
 - Top from simulation (relatively small @ Tevatron)
 - Z+jets from data & MC (“easy” to get a clean sample, correct MC)
 - QCD multijet from data (no choice)
 - W + jets from MC, but

Generators Used

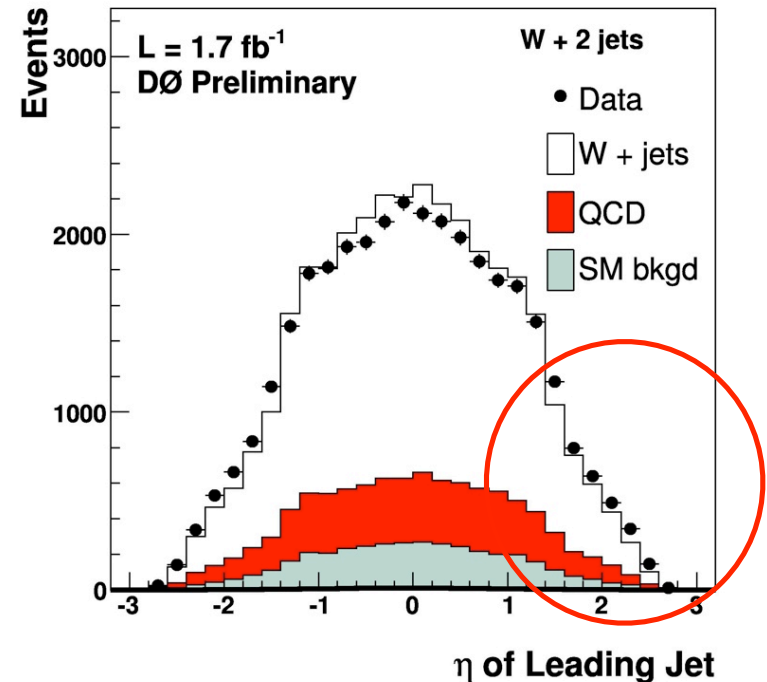
- We use four kinds of Monte Carlo generators
 - “Calculators” (often NNLO) do not actually generate events, they just calculate some (limited) distributions, like $W p^T$
 - Traditional $2 \rightarrow 2$ generators: LO, e.g. $q\bar{q} \rightarrow WZ$
 - Include parton shower, i.e. QCD radiation, and hadronization to jets
 - “Matrix Element” $2 \rightarrow n$ ($n < 9$): LO, e.g. $q\bar{q} \rightarrow evjjjj$
 - Necessary to generate events with multiple hard jets
 - Require matching to parton shower to avoid double counting
 - NLOwPS $2 \rightarrow 2$ generators: include NLO corrections
 - I.e. in a sense they are $2 \rightarrow 2$ & 3 with virtual corrections

Correction Factors

- Of course, the ME's are LO, so “K-factors” needed
- Different ones for heavy flavor etc..... convention to avoid confusion....
- **K-factor is purely theoretical, and denotes a (N)NLO/LO ratio of cross sections;**
- **K'-factor is also theoretical, and denotes a (N)NLO/LL ratio of cross sections.**
According to Steve, ALPGEN cross sections are **Leading Log**;
- **S-factor is empirical, and comes on top of K or K'** to bring MC in agreement with data. MC should be initially normalized to luminosity, and all correction (a.k.a. scale) factors should be applied (trigger, ID...);
- **HF-factor is, in principle, theoretical, but in practice only theory inspired.**
It tells you by how much heavy flavor production should be increased, on top of K or K', and possibly S;
- **S_HF-factor is empirical, and comes on top of K or K', S, and HF, to bring MC in agreement with data, after b-tagging.**

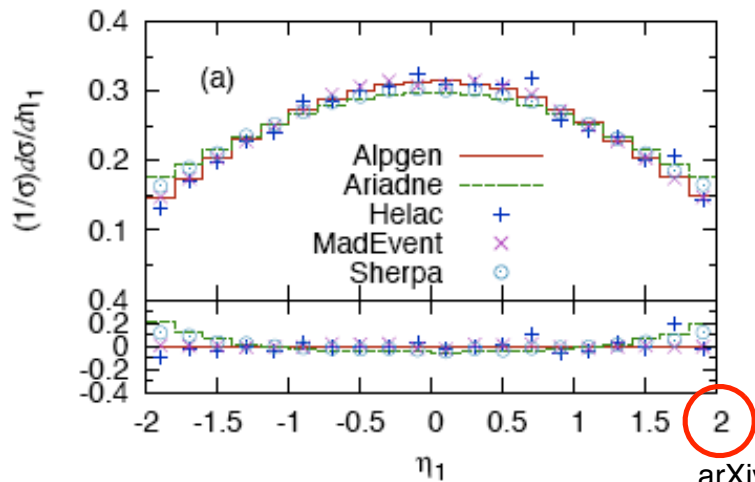
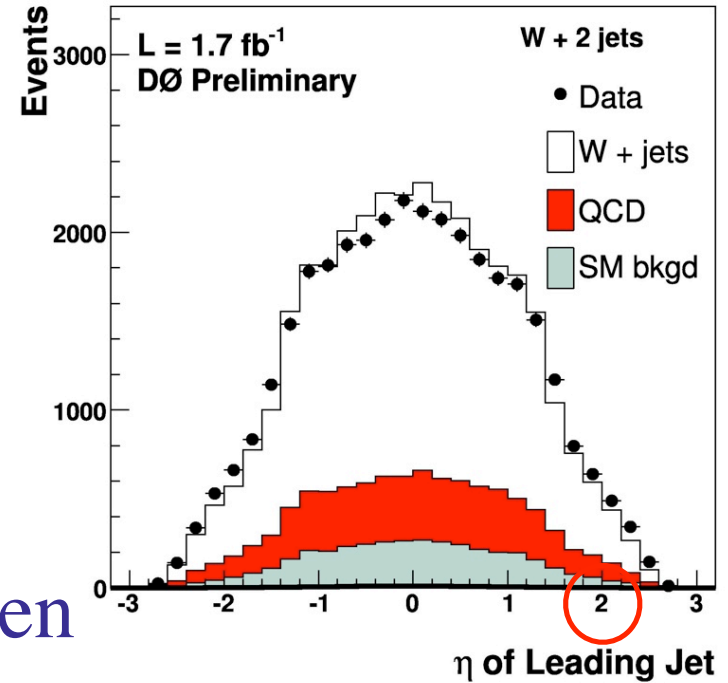
Anecdotes From the Field (III)

- Pile-up events (“minimum bias”) do produce jets
 - At high \mathcal{L} , require that tracks pointing to jets originate from same vertex as lepton
 - High η excess disappeared!
 - Eta-dependence of jet-vertex match turns out to have shape very similar to excess
 - After correcting for this, excess is back....

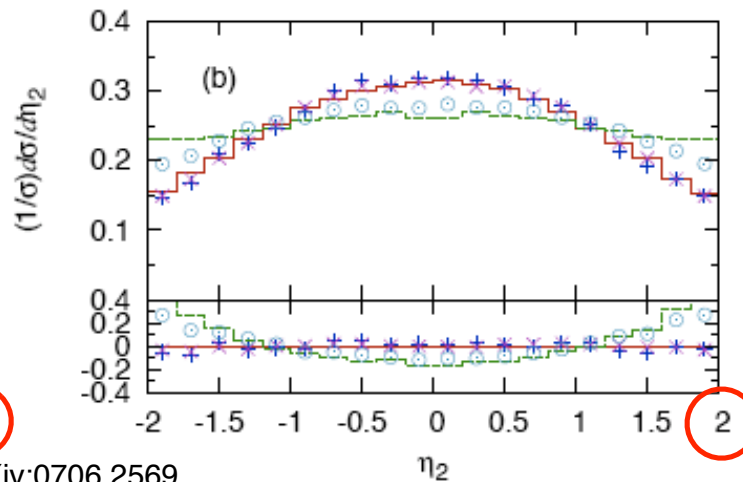


So...

- After all K/K'/S/HF-factors and boson p_T reweighting:
- Similar angular differences between generators: reweigh alpgen to sherpa



arXiv:0706.2569



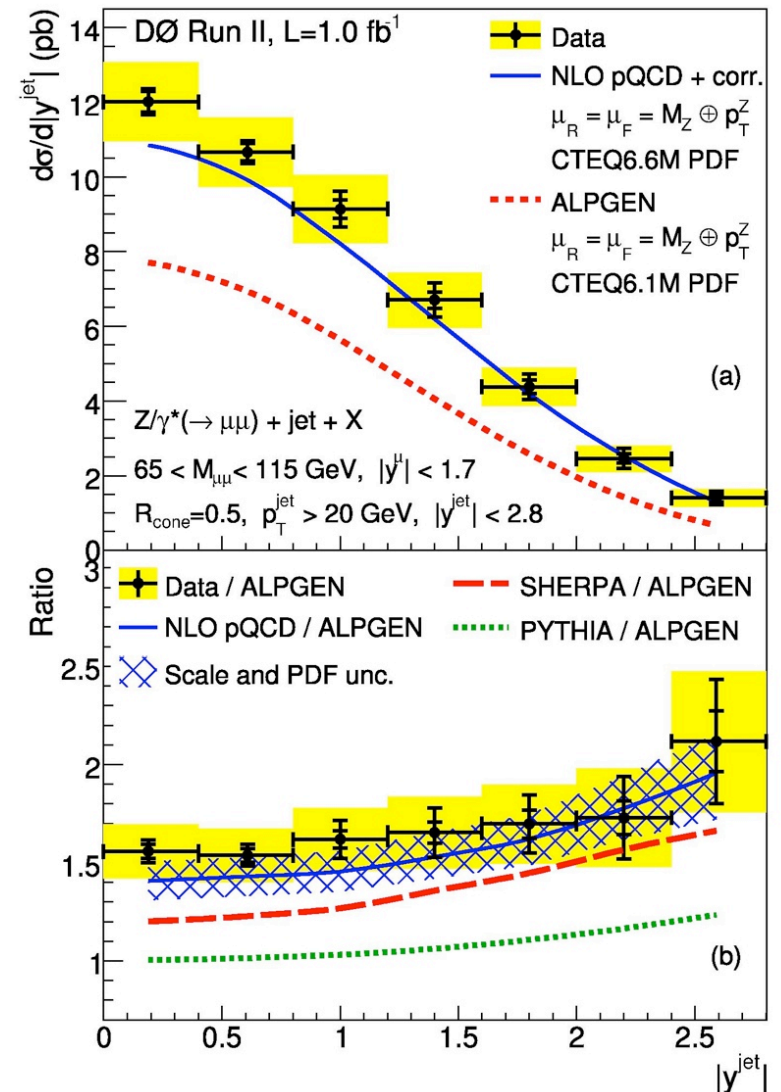
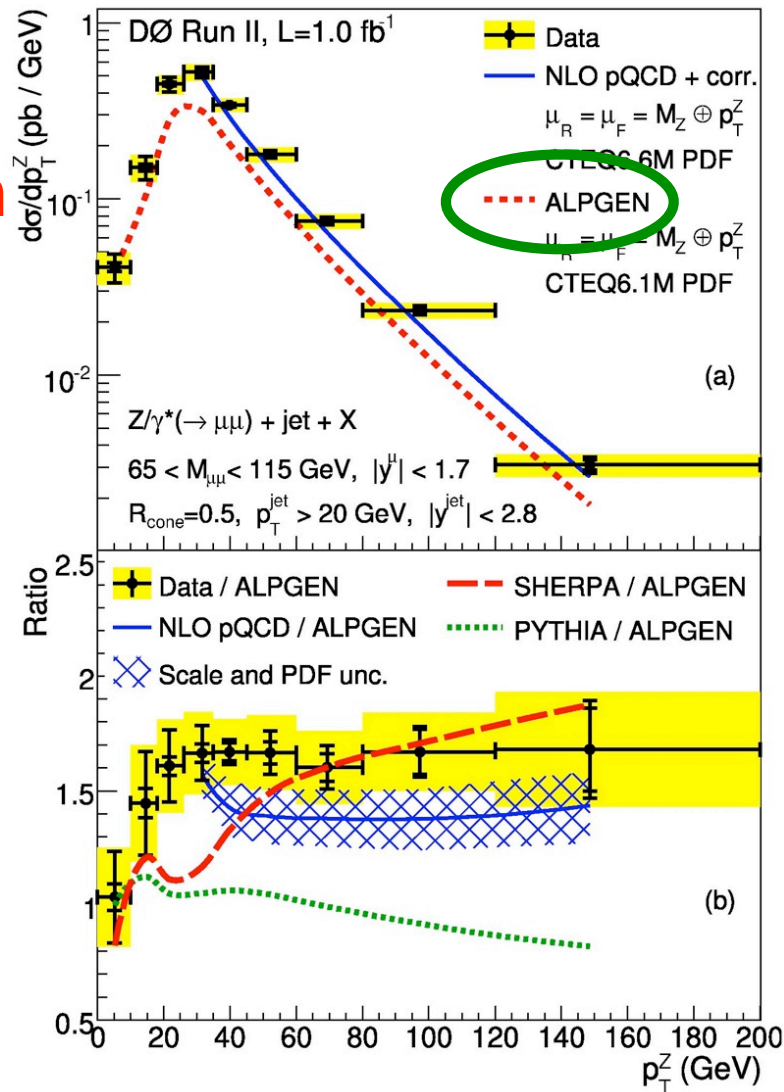
Alpgen, MadEvent,
Helac with MLM,
Sherpa and Ariadne
with CKKW
matching

Z ($\rightarrow ll$) + jets

- Can get a clean sample, check if our simulation reproduces the data

\Rightarrow yes, with
~expected
deviations

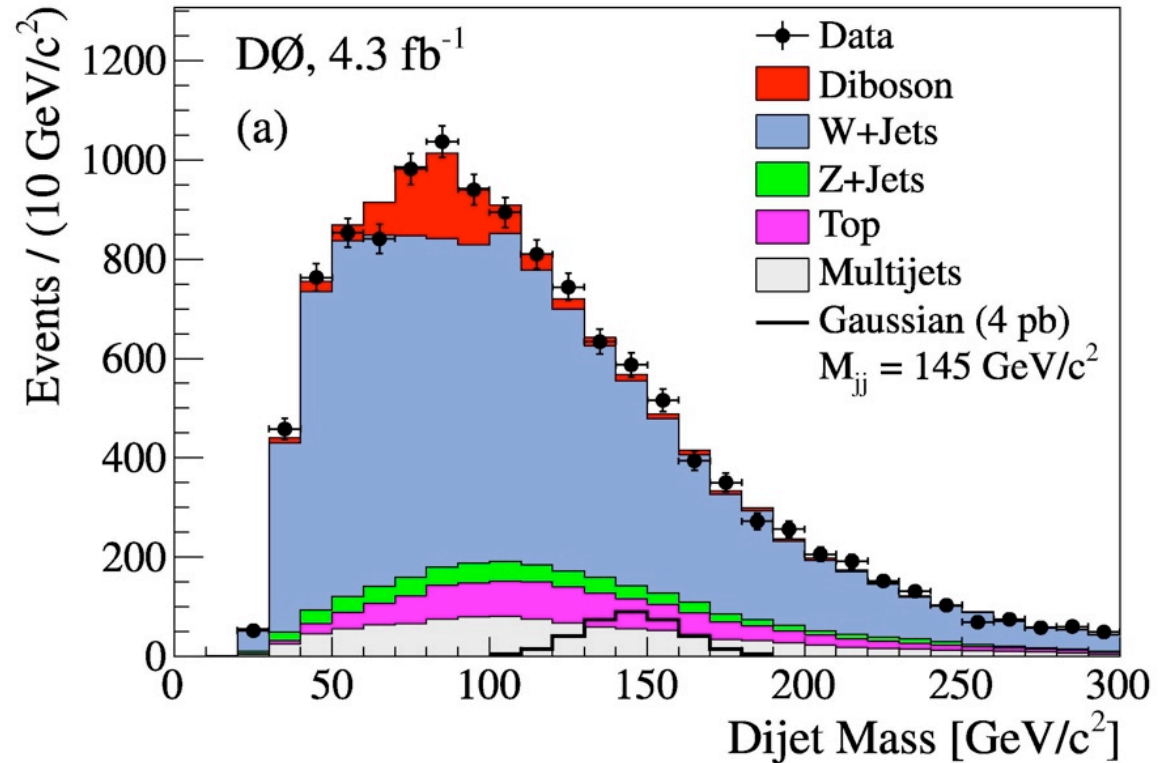
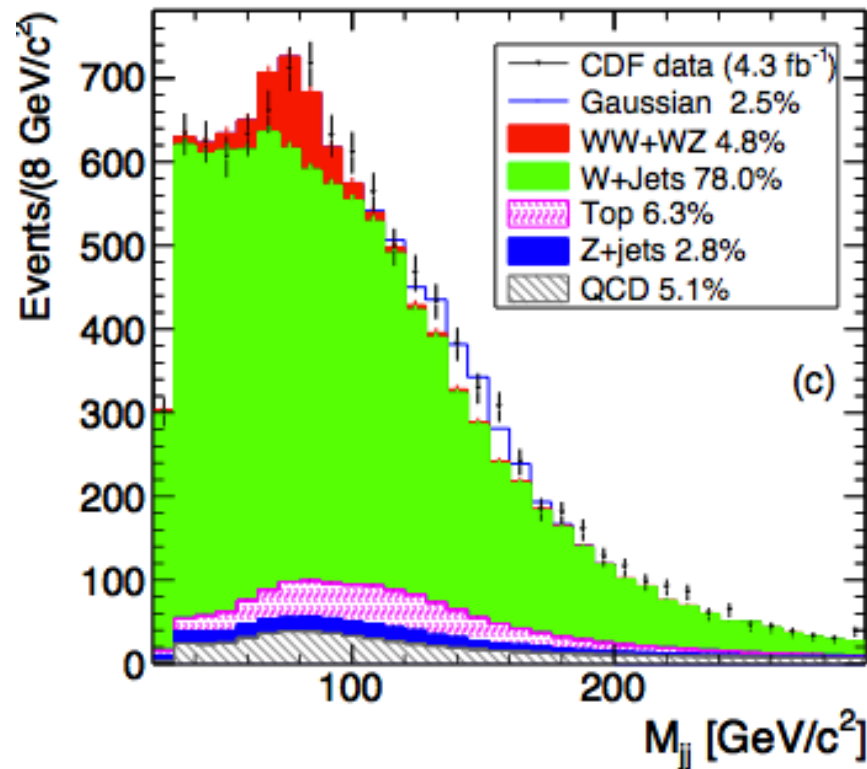
Need
reweighing
of MC



Anecdotes From the Field (III)

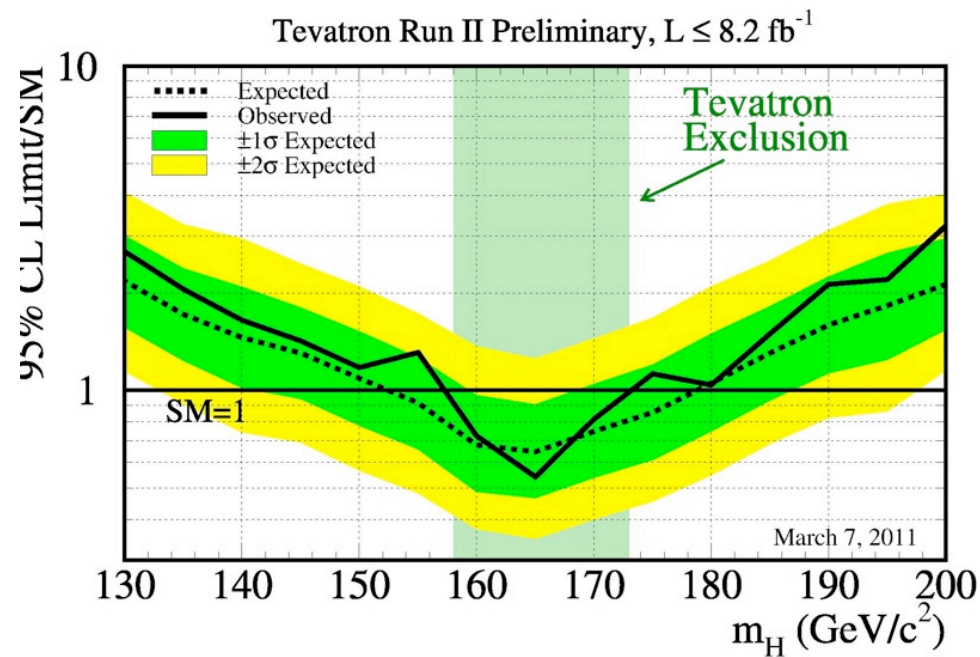
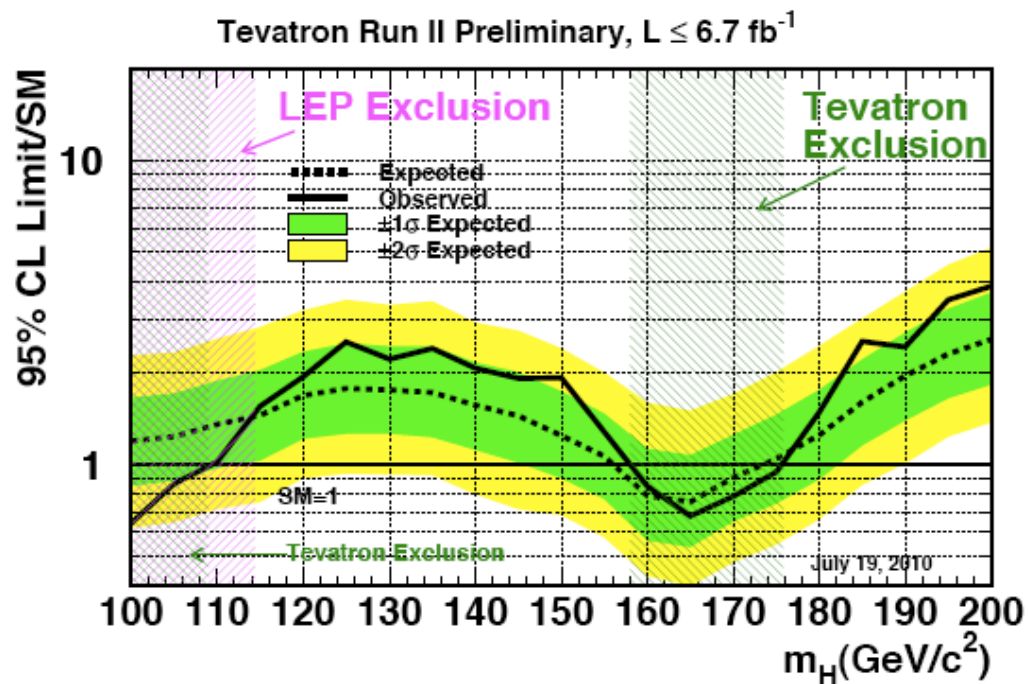
- CDF searched for WW/WZ in $\ell\nu jj$

[Phys.Rev.Lett.106:171801](https://arxiv.org/abs/hep-ex/0605011)



- The background here is not SM, it is uncorrected alpgen!!
- But this is not the issue.....

All Channels, CDF + DØ



- Low mass a struggle!

129 Channels!!

So, Physic Analysis

- Start from:
 - “*How well* do we understand data *and* the SM?”
 - How confident are we in corrections we apply?
- Given that:
 - Which measurements can we make?
 - What do we need to do to improve our understanding?
Complementary measurements!
- Balance the work!
 - Early, low background searches
 - Detailed understanding/verification of SM predictions

