

INVISIBLE TRACES OF CONFORMAL SYMMETRY BREAKING

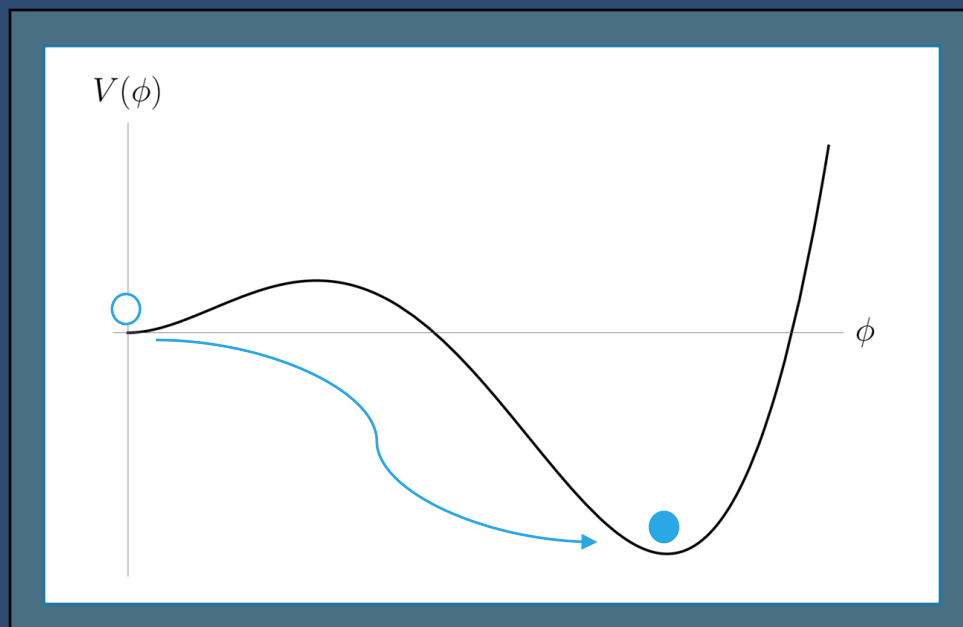
Maciej Kierkla

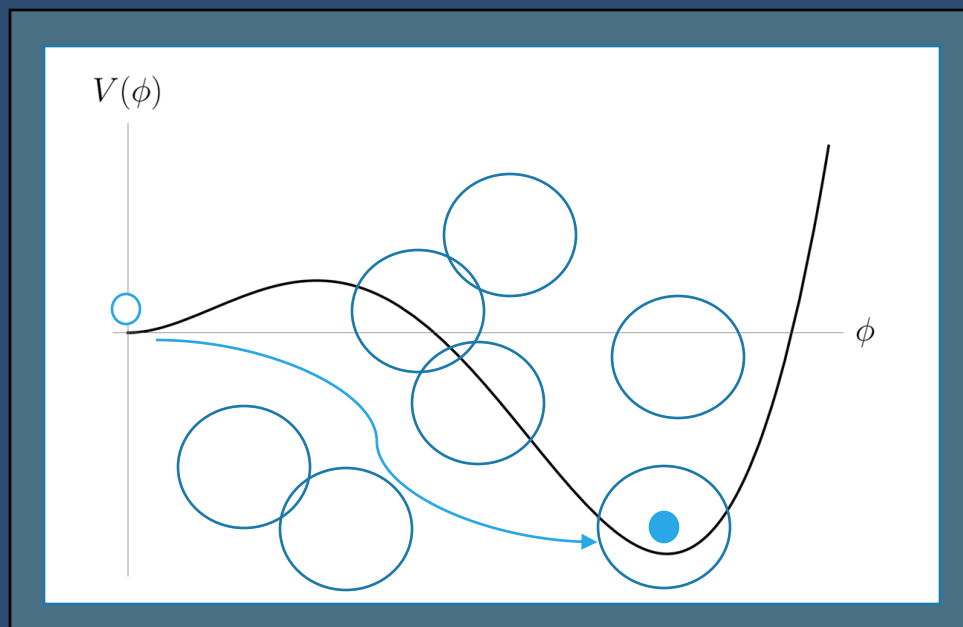
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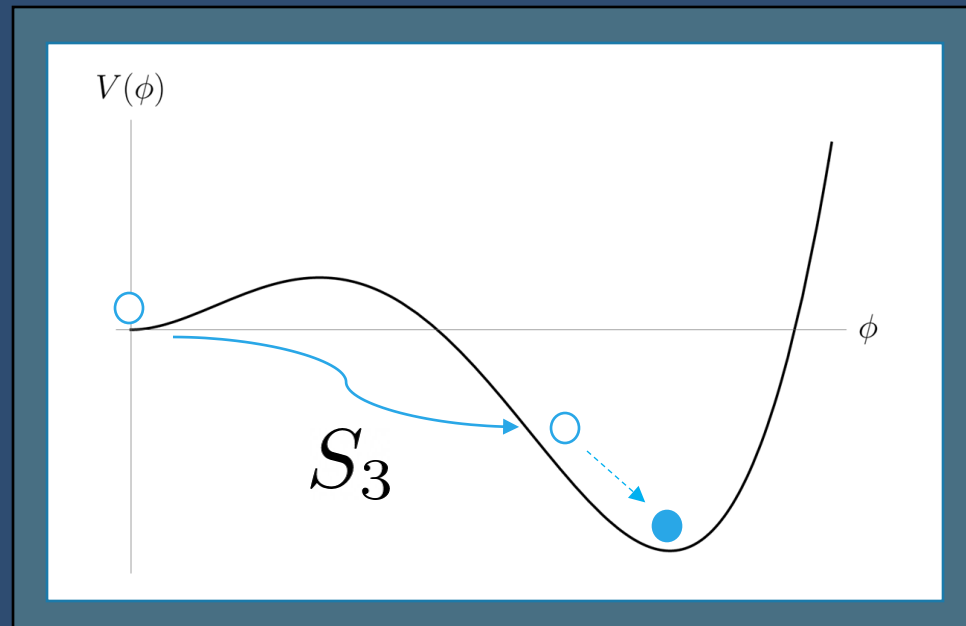






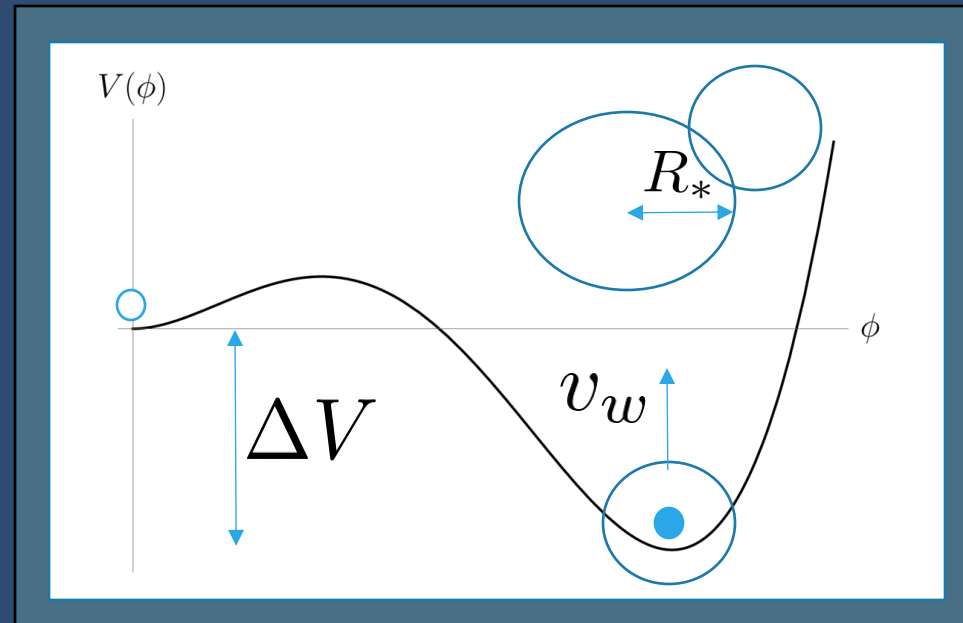
Thermal tunnelling

- Decay rate is given as: $\Gamma(T) \simeq T^4 \left(\frac{S_3}{2\pi T} \right)^{3/2} e^{-S_3/T}$
- The action of the field configuration $S_3 = 4\pi \int r^2 dr \left[\frac{1}{2} \left(\frac{d\varphi}{dr} \right)^2 + V_{\text{eff}}(\varphi, T) \right]$



Cosmological phase transitions

- Strength of the transition $\alpha \sim \frac{\Delta V}{\rho_{\text{rad}}(T_p)}$
- Average bubble radius R_* or inversed time scale $\beta \sim R_*^{-1}$
- Bubble wall velocity v_w



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graph LR; Model[Model] --> Properties[Phase Transition Properties]; Properties --> Spectrum["GW power spectrum  
ΩGWh2"]; style Model fill:#0096d6,color:#fff; style Properties fill:#0096d6,color:#fff; style Spectrum fill:#0096d6,color:#fff;
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Model

Phase Transition
Properties

GW power
spectrum
 $\Omega_{\text{GW}}h^2$

$SU(2)_cSM$



$$V_{\text{tree}} = \frac{1}{4} (\lambda_1 h^4 + \lambda_2 h^2 \phi^2 + \lambda_3 \phi^4)$$

Classical
conformal symmetry

All masses generated
via Coleman-
Weinberg mechanism

Only two free
parameters



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$SU(2)_cSM$



Perturbative and stable
up to the Planck scale

Vector DM candidate,
gauge boson X

Exhibits supercooling

$$V_{\text{tree}} = \frac{1}{4} (\lambda_1 h^4 + \lambda_2 h^2 \phi^2 + \lambda_3 \phi^4)$$

Dark Matter

New gauge boson X , with mass $m_X^2 = \frac{1}{4}g_X^2\phi^2$

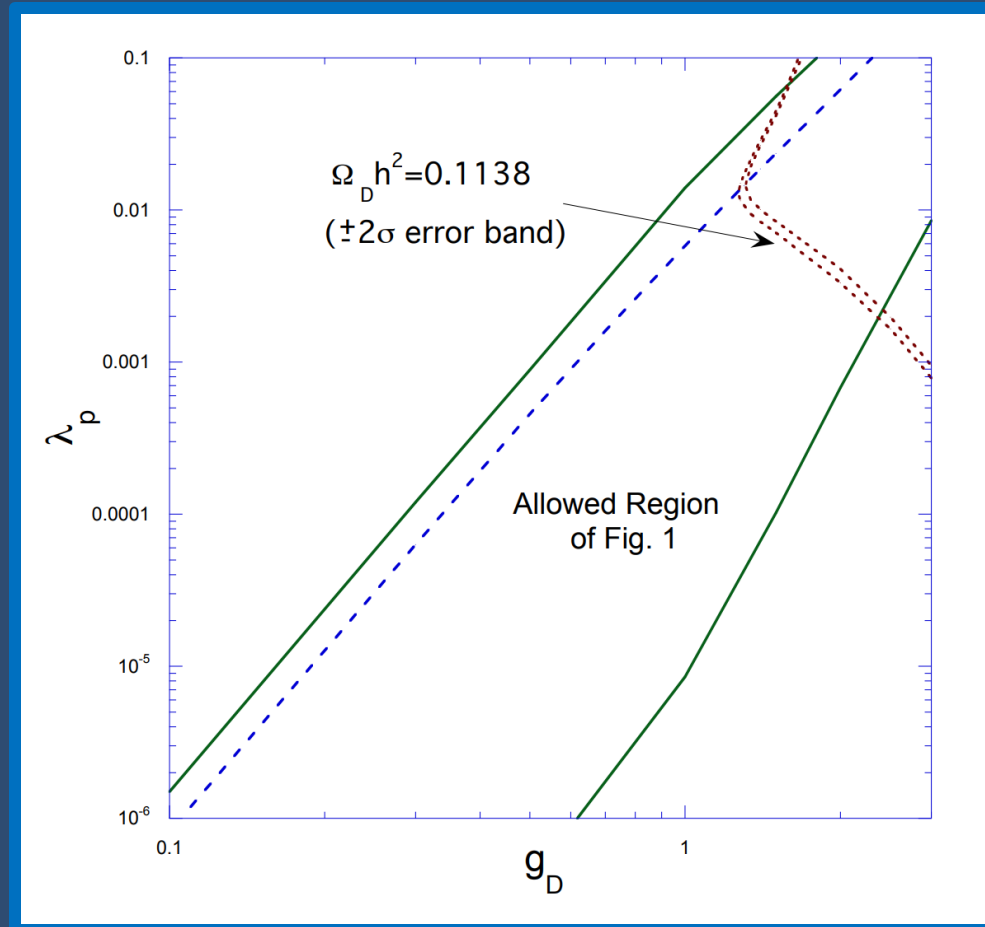
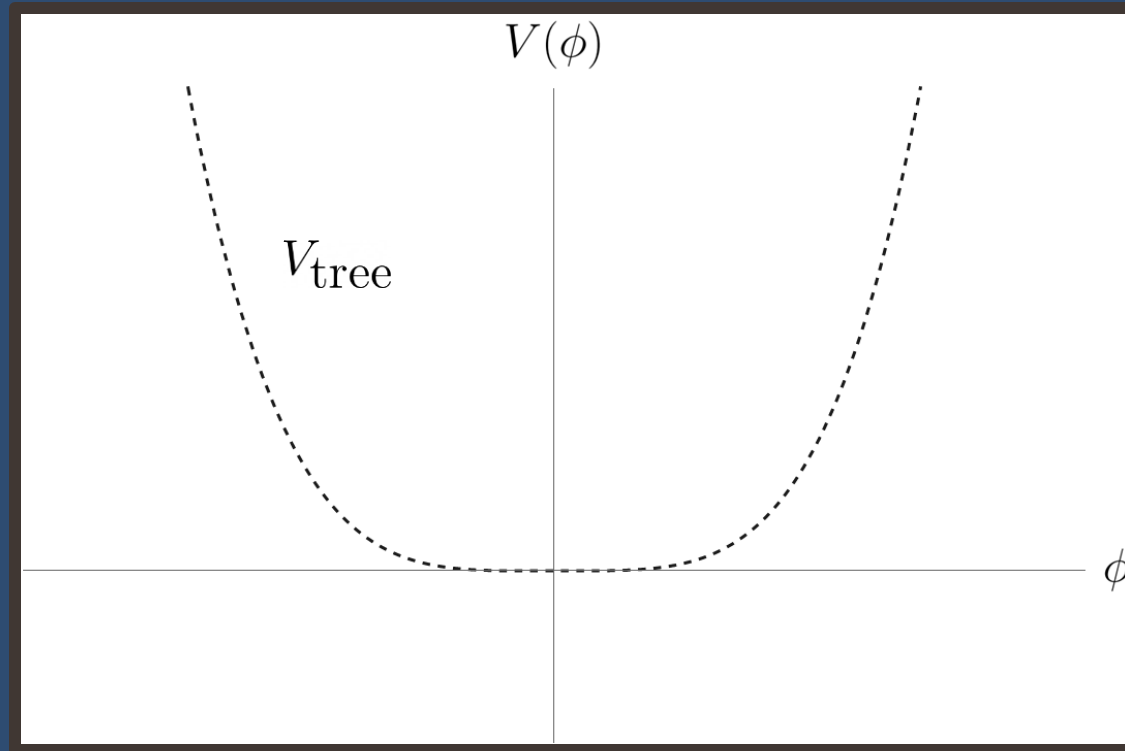


Figure from C. D. Carone, R. Ramos , arXiv:1307.8428

Coleman-Weinberg Mechanism

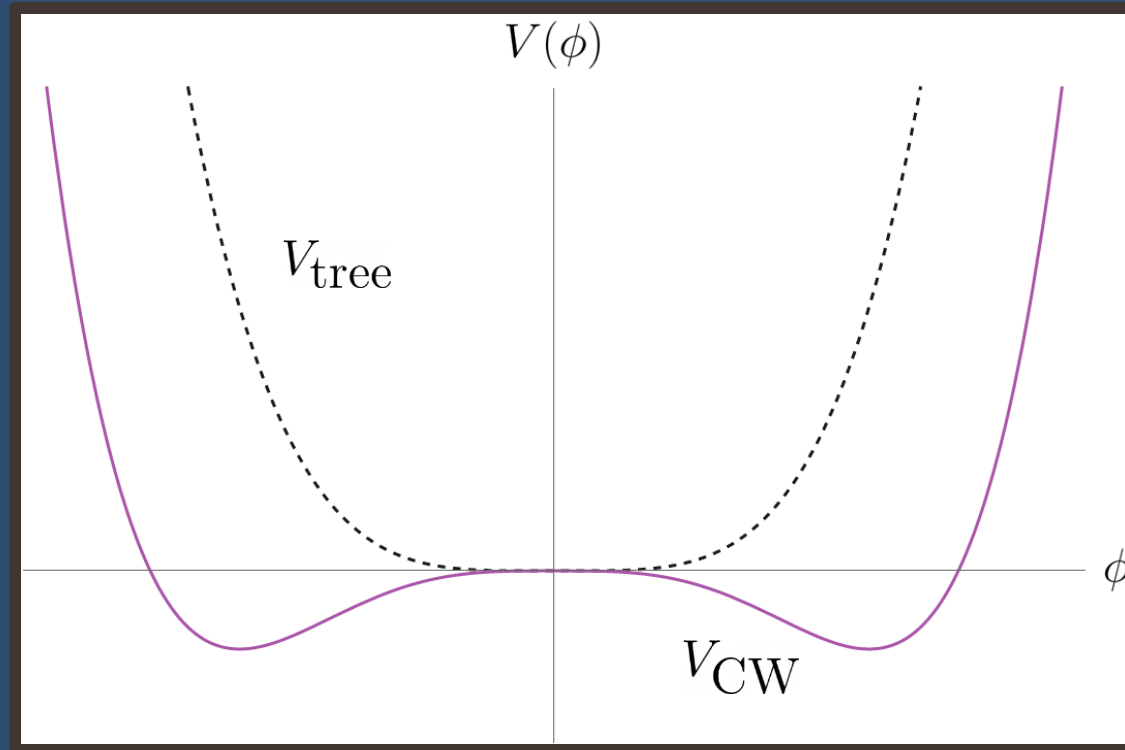
Example: massless
scalar electrodynamics.



$$V_{CW} = \underbrace{\frac{\lambda}{4!} \phi^4}_{V_{\text{tree}}}$$

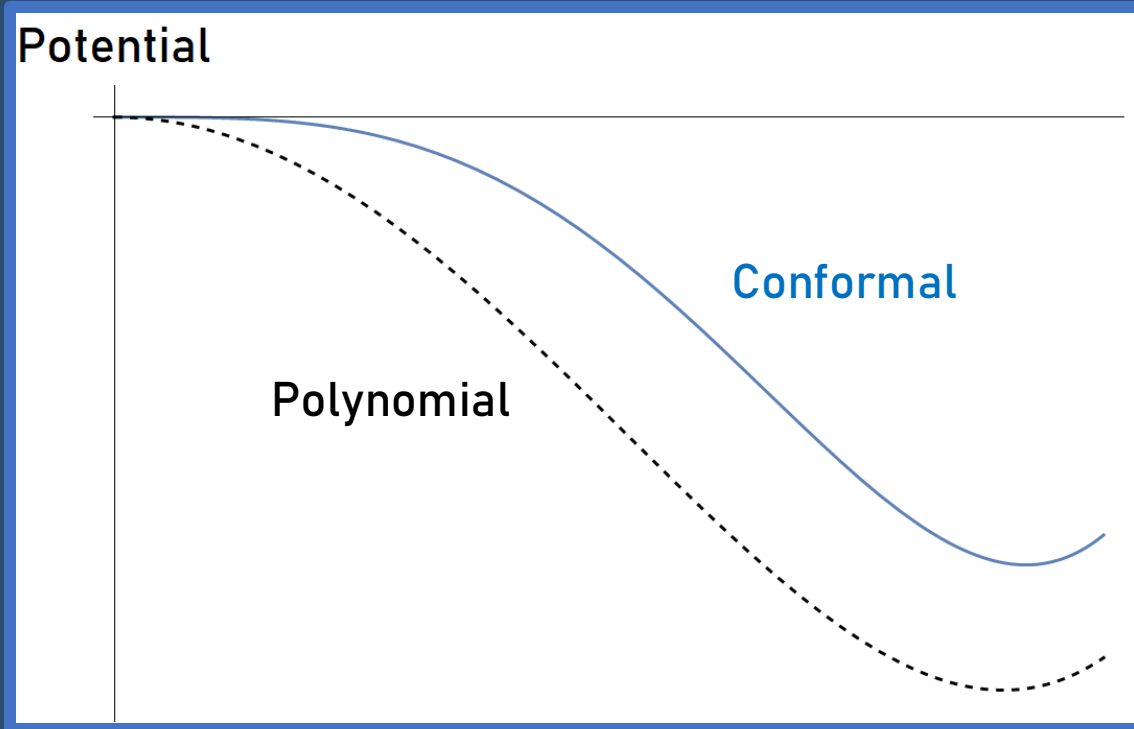
Coleman-Weinberg Mechanism

Example: massless scalar electrodynamics.



$$V_{CW} = \underbrace{\frac{\lambda}{4!} \phi^4}_{V_{\text{tree}}} + \underbrace{\frac{\lambda^4}{64\pi^2} \phi_c^4 \left(\ln \frac{\phi_c^2}{\mu^2} - \frac{3}{2} \right)}_{\text{"scalar" correction}} + \underbrace{\frac{3e^4}{64\pi^2} \phi_c^4 \left(\ln \frac{\phi_c^2}{\mu^2} - \frac{5}{6} \right)}_{\text{"boson" correction}}$$

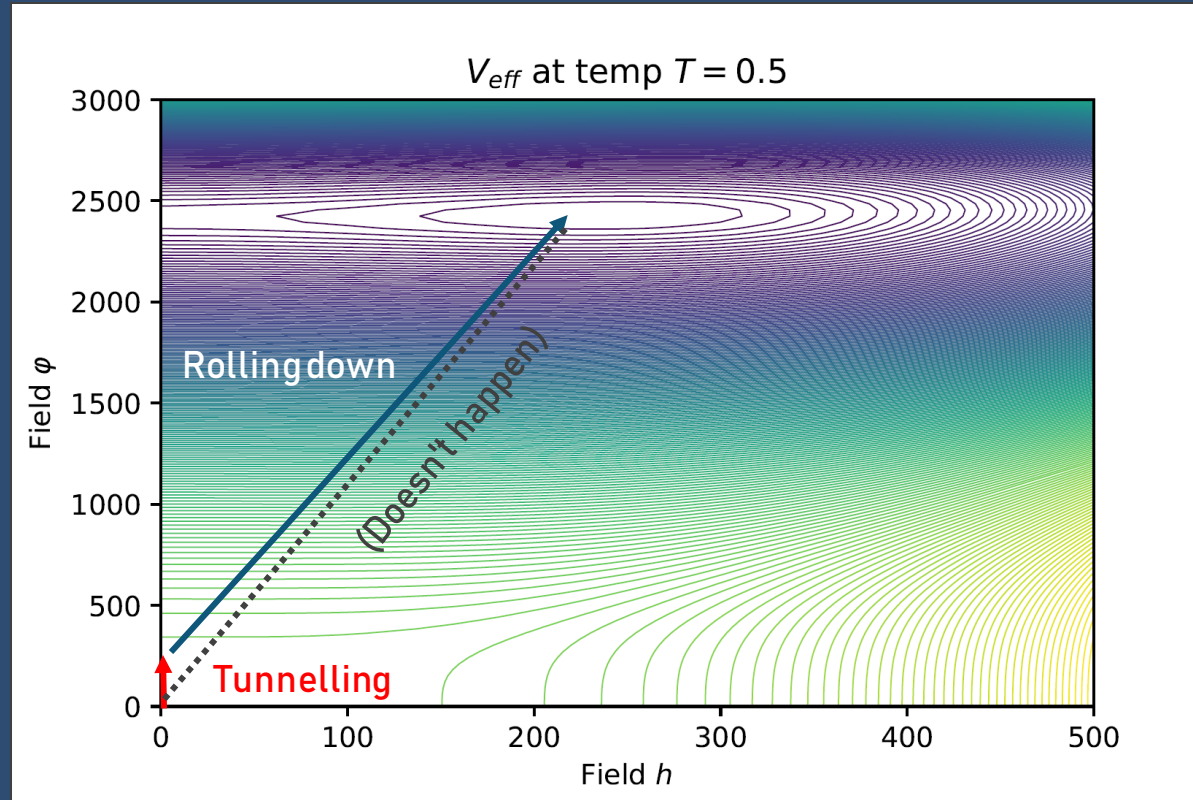
Introducing: supercooling



Features:

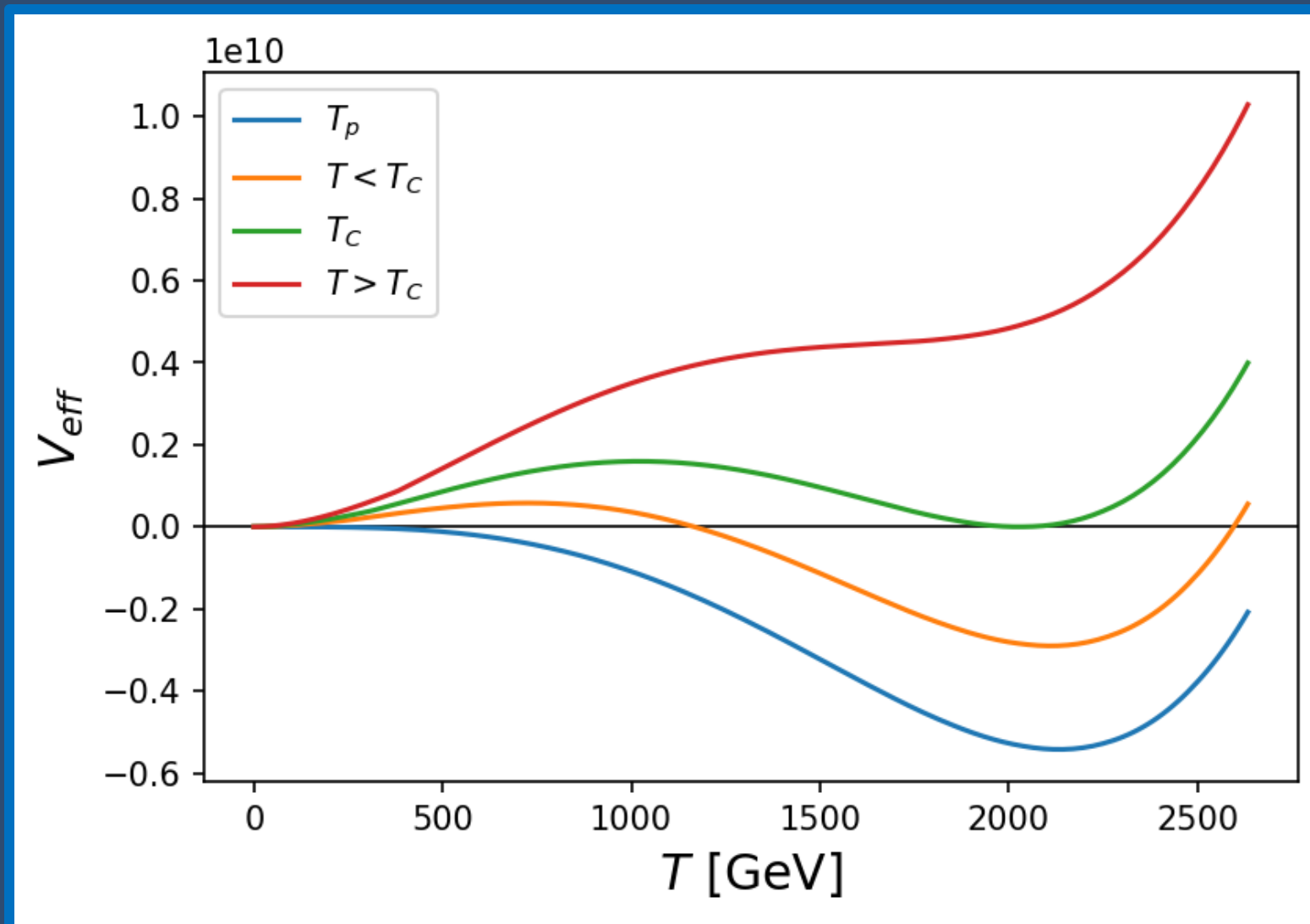
- phase transition happens at temperatures significantly below EW scale,
- thermally produced barrier lasts till $T=0$,
- Induces strong Gravitational Wave signal.

Tunneling scenario in SU(2)cSM



Tunnelling occurs only in the new scalar direction!

Thermal evolution of the effective potential



Thermal inflation

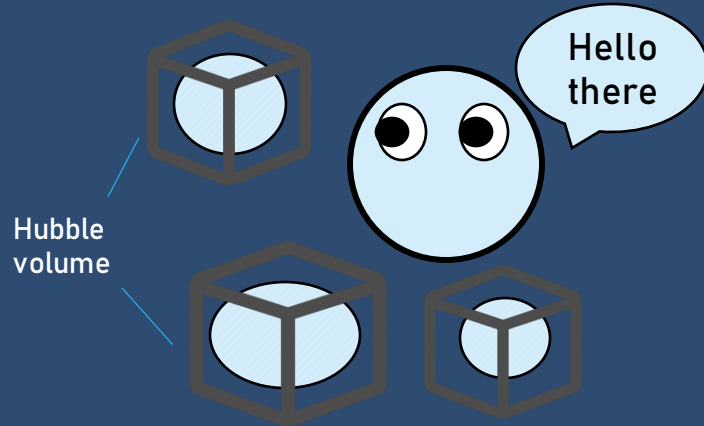
Hubble parameter is given as:

$$H^2 = \frac{1}{3M_{\text{pl}}^2} (\rho_{\text{R}} + \rho_{\text{V}}) = \frac{1}{3M_{\text{pl}}^2} \left(\frac{T^4}{\xi_g^2} + \Delta V \right)$$

We enter into the vacuum domination at the temperature $T_V = \sqrt[4]{\Delta V \xi_g^2}$

Due to non-zero vacuum energy there is a period of *thermal inflation*.
This affects the calculation of thermodynamical parameters.

- Nucleation temperature



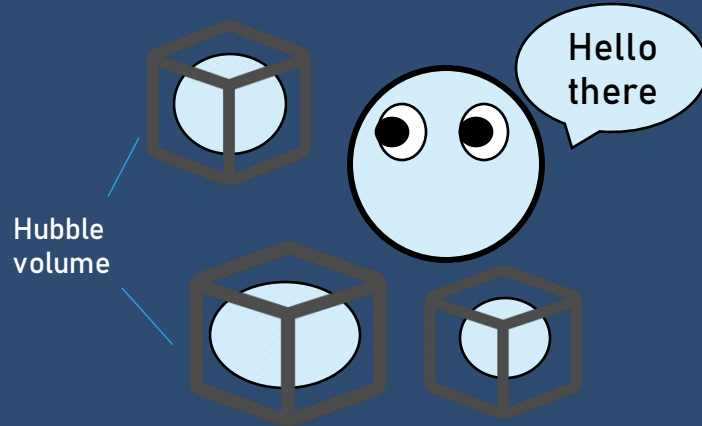
We calculate:
$$N(T_n) = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

One can also use an approximation:

But not this one
$$\frac{S_3}{T} \simeq 140$$

$$\frac{\Gamma(T_n)}{H(T_n)^4} \simeq 1 \Rightarrow \frac{S_3}{T_n} = 4 \log \left(\frac{T_n}{H(T_n)} \right)$$

- Nucleation temperature

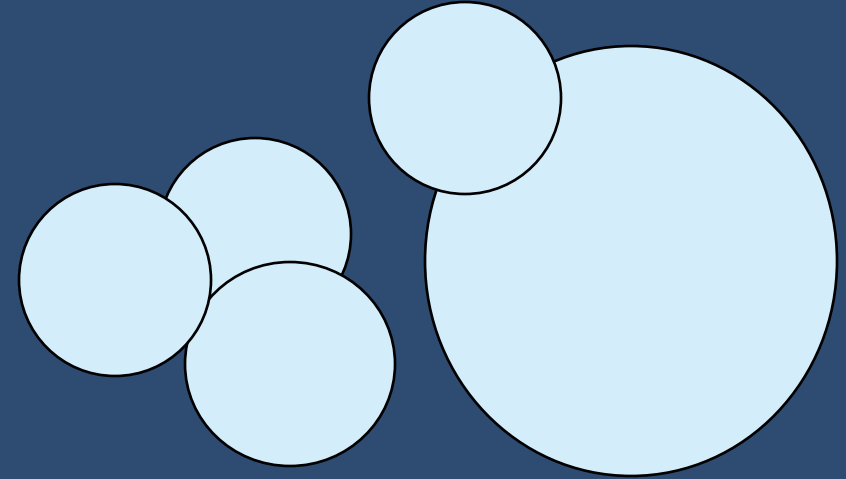


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- Percolation temperature



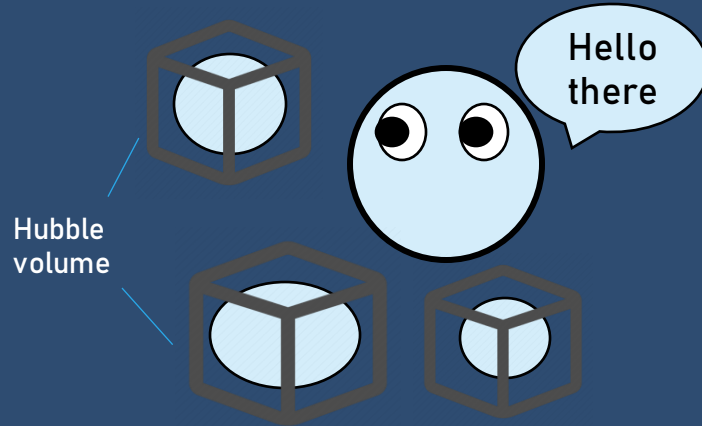
Probability of point still in false vacuum is $P = e^{-I(T)}$, where

$I(T)$ is the volume converted into true vacuum

Then we solve for condition:

$$I(T_p) \simeq 0.34$$

- Nucleation temperature



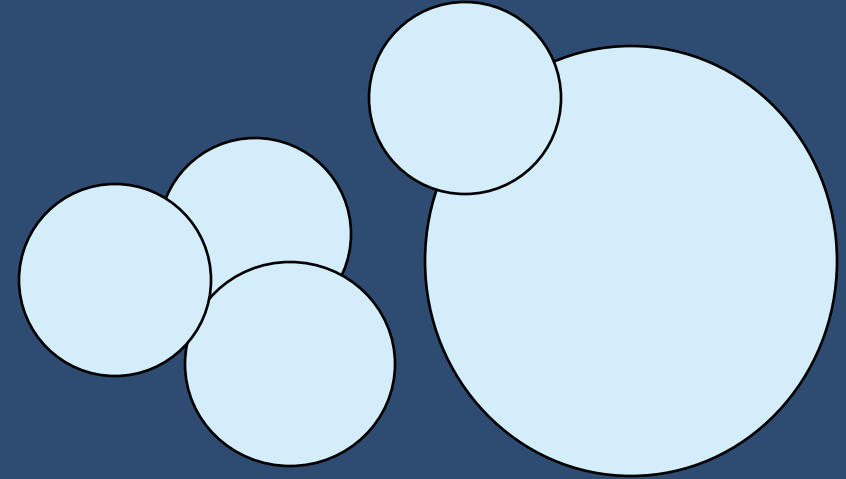
Are not equal
in models with
supercooling

We calculate:
$$N(T_n) = \int_{T_n}^{T_c} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

One can also use an approximation:

$$\frac{\Gamma(T_n)}{H(T_n)^4} \simeq 1 \Rightarrow \frac{S_3}{T_n} = 4 \log \left(\frac{T_n}{H(T_n)} \right)$$

- Percolation temperature



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Reheating temperature...

$$\Gamma_{\varphi} > H_*$$

but if....

$$\Gamma_{\varphi} < H_*$$

- Reheating is instantaneous
- Released energy transforms into radiation
- Universe reheats up to the temperature T_V

- Energy will be stored in the scalar field oscillating about the true vacuum
- Matter domination until temperature at which decay rate is equal to Hubble parameter
- This matter domination period changes the shape of GW spectrum


$$\Omega_{\text{GW}} = \Omega_{\text{collisions}} + \Omega_{\text{sound waves}} + \Omega_{\text{turbulence}}$$

How do we know which source dominates?

Efficiency factors:

$$\kappa_{col} = \frac{E_{wall}}{E_V}$$

$$\kappa_{sw} \sim 1 - \kappa_{col}$$

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And the main GW source is...

Where the energy goes?

There is a lot of friction

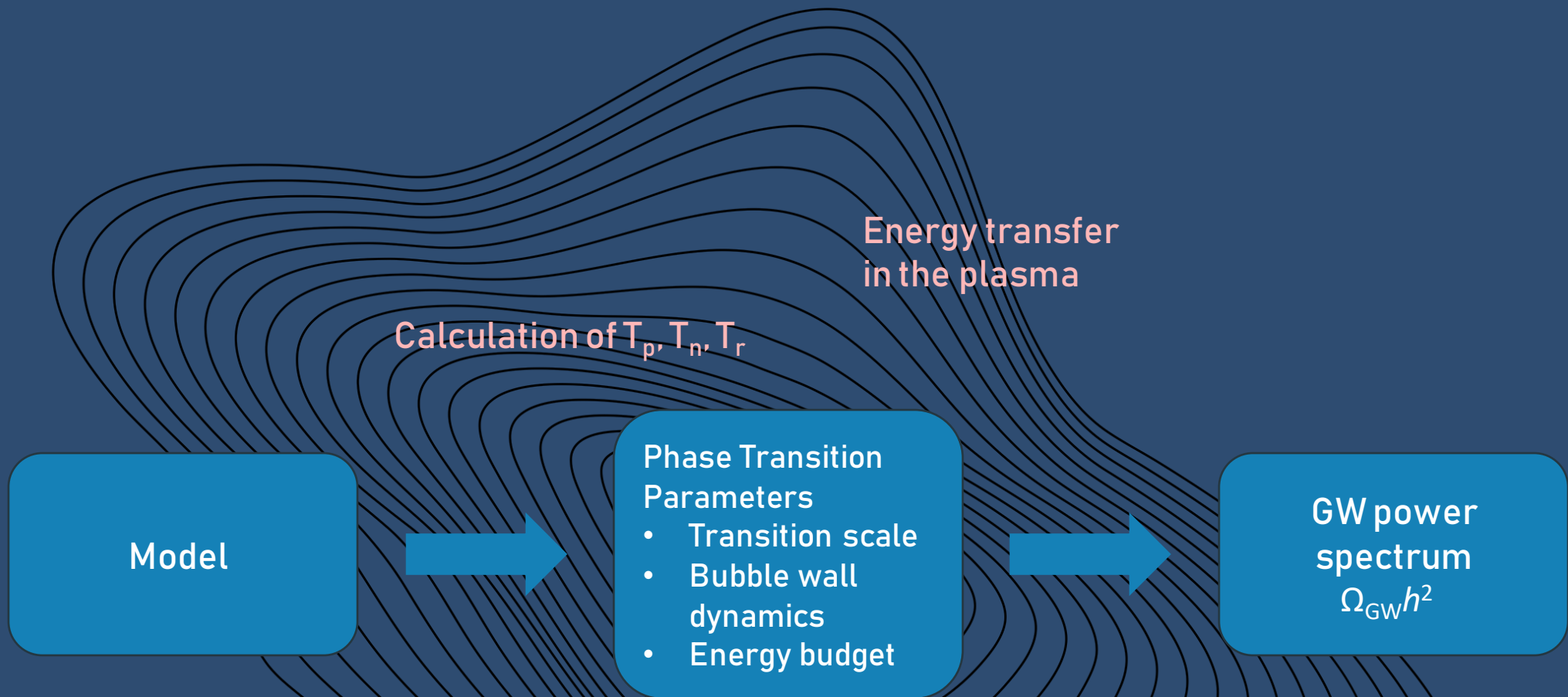
Bubble expansion accelerates

Energy is dissipated in the surrounding plasma

Energy goes to the bubble's wall

Sound waves + Turbulences

Bubble collisions



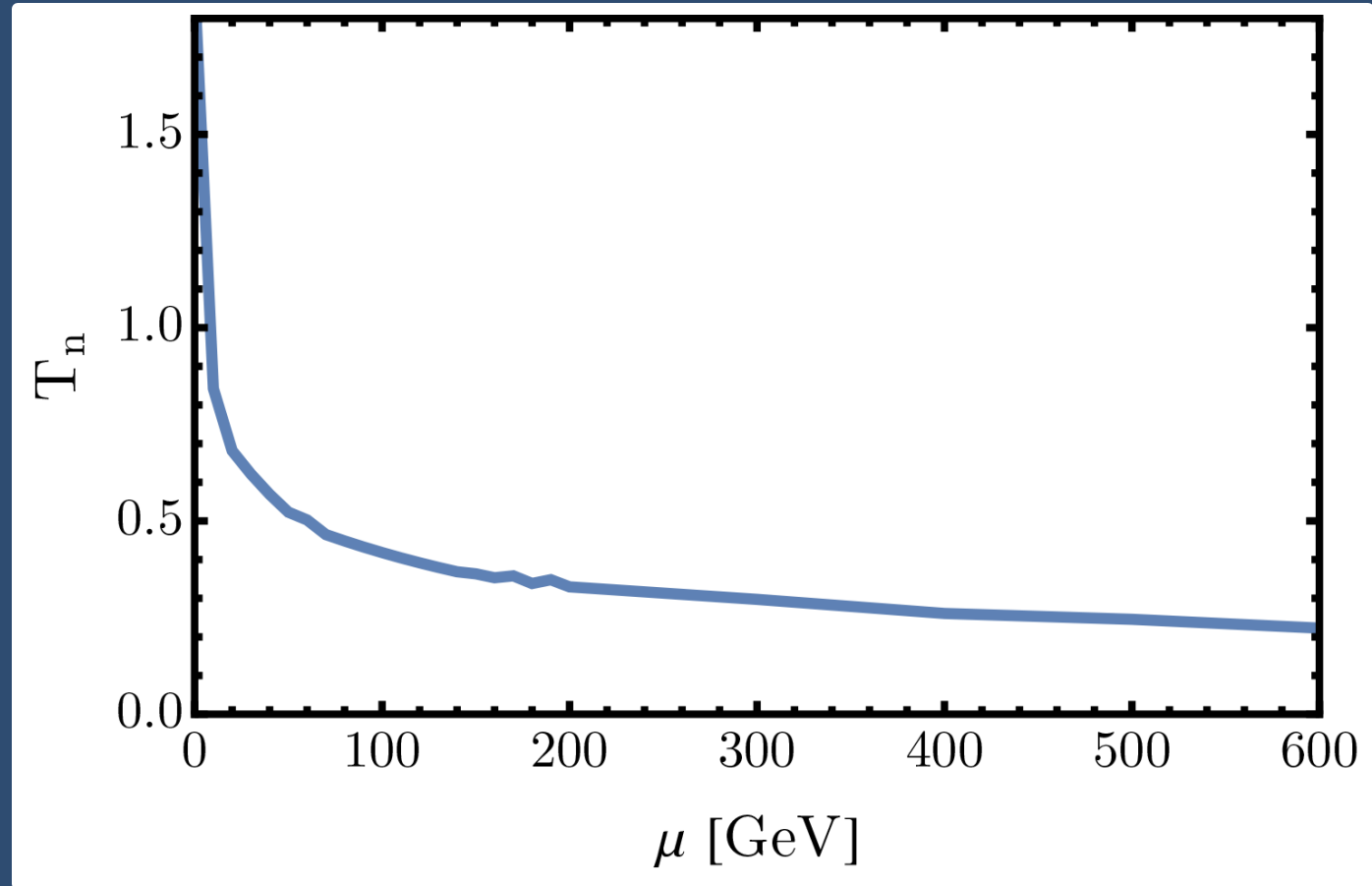
Calculation of T_p, T_n, T_r

Energy transfer in the plasma

RG scale dependence

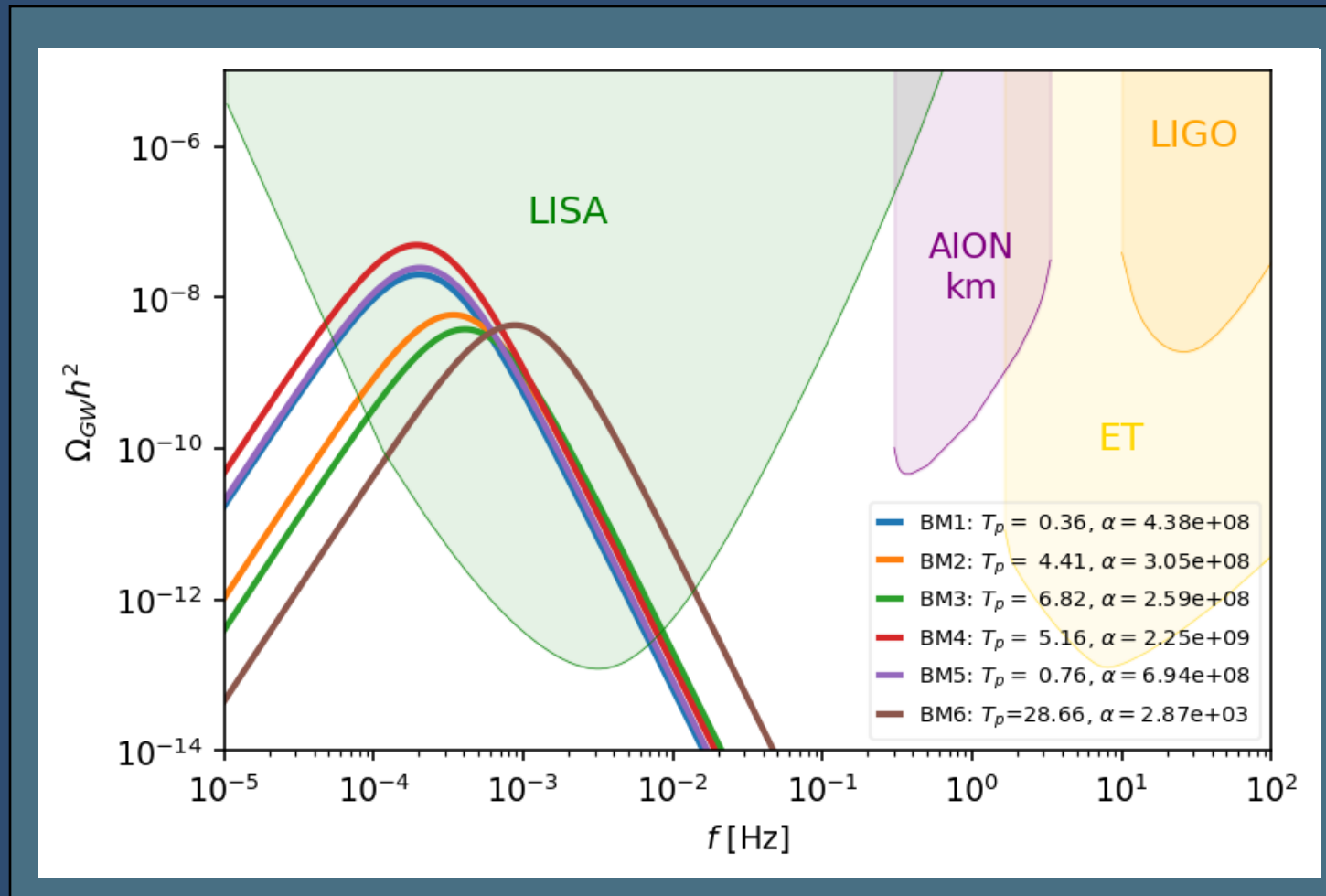
Efficiency factors i.e inclusion of all possible sources

- RG scale dependence



This affects other parameters and therefore the resulting spectrum.

Gravitational Waves spectra in SU(2)cSM



Goal: provide accurate predictions for LISA.

Thank you



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