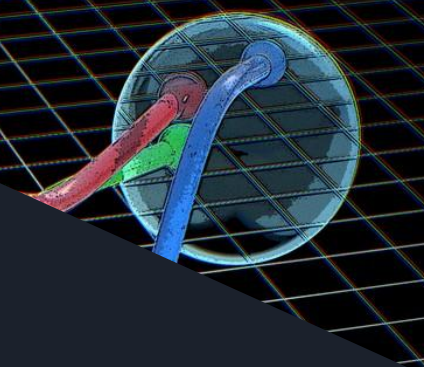


Towards High-Precision Nucleon Parton Distributions via Distillation

JHEP 11 (2021) 148; arXiv: 2107.05199 [hep-lat]



LaMET 2021

December 7th, 2021



Colin Egerer
(For the HadStruc Collaboration)



Matrix Elements of Non-Local Parton Bilinears

Focus of this talk:

A matrix element of a
distinct character

$$M^{\alpha}(p, z) = \langle h(p) | \bar{\psi}(z) \gamma^{\alpha} \Phi_{\hat{z}}^{(f)}(\{z, 0\}) \psi(0) | h(p) \rangle$$



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$$\frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle h(p) | \bar{\psi}(\frac{z}{2}) \gamma^+ \Phi_{\tilde{z}^-}^{(f)}(\{\frac{z}{2}, -\frac{z}{2}\}) \psi(-\frac{z}{2}) | h(p) \rangle$$

$$p^\alpha = \left(p^+, \frac{m_h^2}{2p^+}, \mathbf{0}_\perp \right)$$

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- universality/constrained by global analyses

$$F_i(x, Q^2) = \sum_{a=q, \bar{q}, g} f_{a/h}(x, \mu^2) \otimes H_i^a\left(x, \frac{Q^2}{\mu^2}, \alpha_s(\mu^2)\right) + h.t.$$

J. Collins, D. Soper, G. Sterman, *Adv. Ser. Direct. High Energy Phys.* 5, 1 (1989)

- complement 3D imaging efforts



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A. Radyushkin, Phys.Rev.D 96 (2017) 3, 034025

A frame amenable to Lattice QCD

$$p^\alpha = (E, \mathbf{0}_\perp, p_z)$$

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Ioffe-time Pseudo-distribution (pseudo-ITD)

$$\mathcal{M}(p_z z_3, z_3^2) = \int_{-1}^1 dx e^{i\nu x} \mathcal{P}(x, z_3^2)$$

Generalization of light-cone PDFs onto space-like intervals w/ Lorentz covariant parton momentum fraction



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Generalization of light-cone PDFs onto space-like intervals w/ Lorentz covariant parton momentum fraction

- *non-trivial light-cone limit*
- *short-distance factorization matches ITD to pseudo-ITD*

$$\mathcal{M}(\nu, z^2) = C(z^2 \mu^2, \alpha_s(\mu^2)) \otimes Q(\nu, \mu^2) + \mathcal{O}(z^2 \Lambda_{\text{QCD}}^2)$$



Nucleon Pseudo-ITDs - Numerical Study

Spatial smearing increases operator-state overlaps onto low-lying modes

$$\hat{q}(\vec{x}, T) = \sum_{\vec{y}} S[U](\vec{x}, \vec{y}) q(\vec{y}, T)$$



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Distillation - advantageous spatial smearing kernel

M. Peardon et al., Phys. Rev. D80, 054506 (2009)

- low-rank approximation of a gauge-covariant smearing kernel

$$J_{\sigma, n_\sigma} = e^{\sigma \nabla^2} = \sum_{\lambda} e^{-\sigma \lambda} |\lambda\rangle \langle \lambda|$$
$$\square(\vec{x}, \vec{y}; t)_{ab} = \sum_{k=1}^{R_{\mathcal{D}}} \xi_a^{(k)}(\vec{x}, t) \xi_b^{(k)\dagger}(\vec{y}, t)$$



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Why Distillation ?

- *explicit momentum projections - all times*
- *excited-state control*
- *reusability*
- *efficient realization of variational method*

R. Briceno et al., Phys.Rev.D 97 (2018) 5, 054513

J. Dudek et. al., Phys.Rev.D 88 (2013) 9, 094505

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CE, D. Richards, F. Winter, Phys. Rev. D 99 (2019) 3, 034506

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Goal: Can distillation isolate pseudo-ITDs with better precision than conventional smearing kernels & shed light on any systematics?

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JLab/WM/LANL 2+1 Flavor Isotropic Lattices

ID	a (fm)	m_π (MeV)	$L^3 \times N_t$	N_{cfg}	N_{srcs}	N_{vec}
$a094m358$	0.094(1)	358(3)	$32^3 \times 64$	349	4	64

Parameters/Statistics

t_{sep}/a	$p_z \left(\times \frac{2\pi}{L} \right)$	z/a
4, 6, \dots , 14	0, $\pm 1, \dots, \pm 6$	0, $\pm 1, \dots, \pm 12, \dots$
0.38, \dots , 1.32 fm	0, 0.411, \dots , 2.47 GeV	0, 0.094, \dots , 1.13 fm



Obtaining the Ioffe-time Pseudo-Distribution

Needed correlation functions:

$$C_2(p_z, T) = \langle \mathcal{N}(-p_z, T_f) \overline{\mathcal{N}}(p_z, T_0) \rangle = \sum_n |\mathcal{A}_n|^2 e^{-E_n T}$$

$$\begin{aligned} C_3(p_z, T, \tau; z_3) &= V_3 \langle \mathcal{N}(-p_z, T_f) \hat{\mathcal{O}}_{\text{WL}}^{[\gamma_4]}(z_3, \tau) \overline{\mathcal{N}}(p_z, T_0) \rangle \\ &= V_3 \sum_{n, n'} \langle \mathcal{N} | n' \rangle \langle n | \overline{\mathcal{N}} \rangle \langle n' | \hat{\mathcal{O}}_{\text{WL}}^{[\gamma_4]}(z_3, \tau) | n \rangle e^{-E_{n'}(T-\tau)} e^{-E_n T} \end{aligned}$$

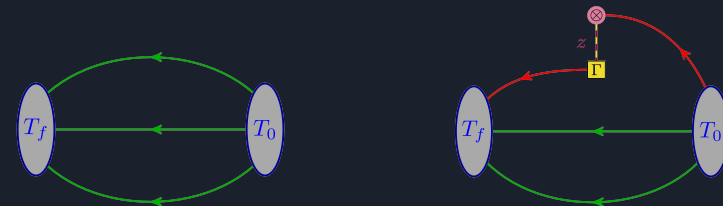


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"Perambulators"

$$\tau_{\alpha\beta}^{kl}(t_f, t_0) = \xi^{(k)\dagger}(t_f) M_{\alpha\beta}^{-1}(t_f, t_0) \xi^{(l)}(t_0)$$

"Elementals"

$$\Phi_{\mu\nu\sigma}^{(i,j,k)}(t) = \epsilon^{abc} (\mathcal{D}_1 \xi^{(i)})^a (\mathcal{D}_2 \xi^{(j)})^b (\mathcal{D}_3 \xi^{(k)})^c(t) S_{\mu\nu\sigma}$$

Irrep.
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Unpolarized PDFs

Space-like Wilson line



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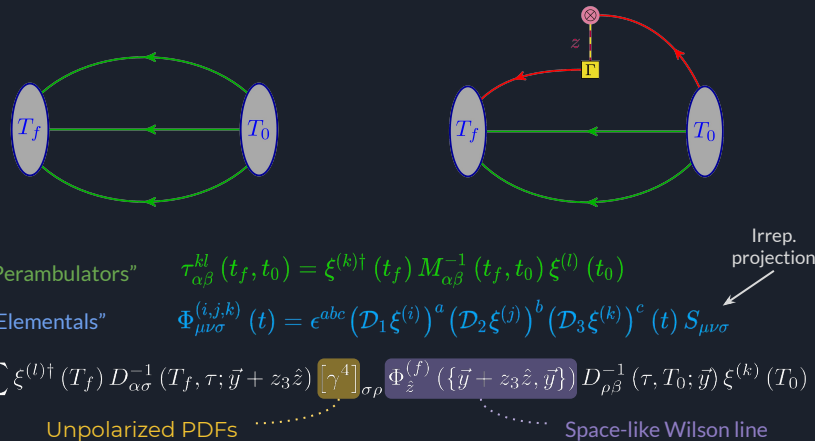
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$$\xi_{\pm}^{(k)}(\vec{z}, t) \equiv e^{i \vec{\zeta}_{\pm} \cdot \vec{z}} \xi_{\pm}^{(k)}(\vec{z}, t)$$

G. S. Bali et al. Phys. Rev. D93, 094515 (2016)
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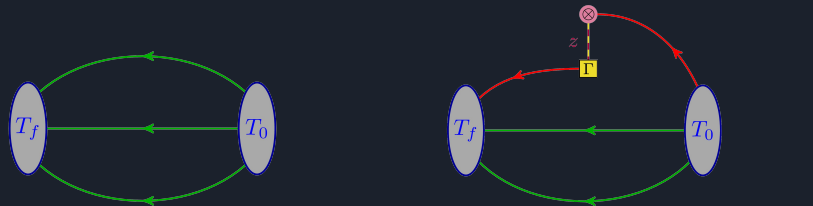
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Summation method

L. Maiani et al., Nucl. Phys. B293 (1987)
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$$R(p_z, z_3; T) = \sum_{\tau/a=1}^{T-1} \frac{C_3(p_z, T, \tau; z_3)}{C_2(p_z, T)}$$

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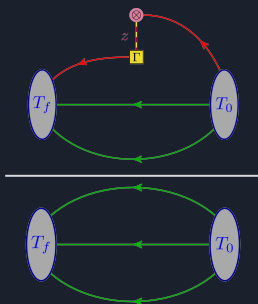
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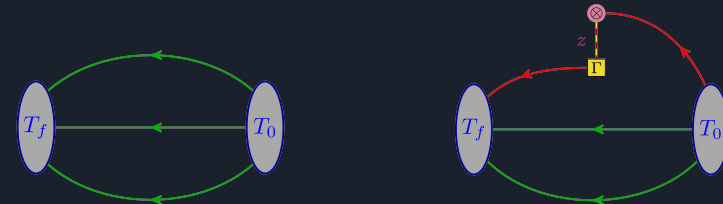
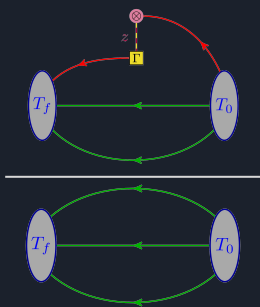
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Reduced Distribution (reduced pseudo-ITD)

➤ cancel Wilson line UV divergences

K. Orginos, et al., Phys. Rev. D96, 094503 (2017)

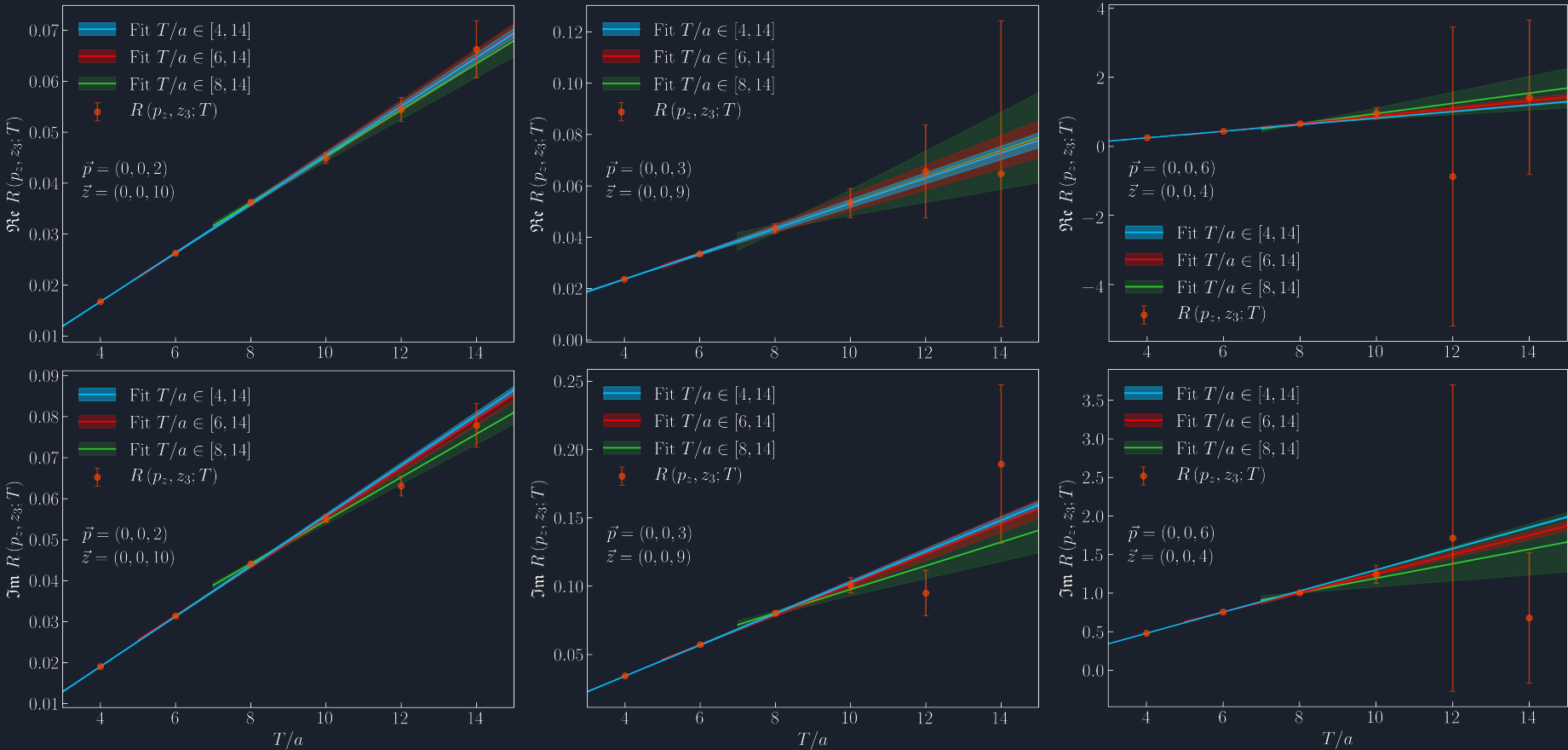
Local vector current in
zero sep. limit
(not conserved)

$$\mathfrak{M}(\nu, z_3^2) \equiv \frac{\mathcal{M}_p(\nu, z_3^2)}{\mathcal{M}_p(0, z_3^2)} \dots$$

$$= \left(\frac{\mathcal{M}_p(\nu, z_3^2)}{\mathcal{M}_p(\nu, 0) |_{z_3=0}} \right) \times \left(\frac{\mathcal{M}_p(0, 0) |_{p=0, z_3=0}}{\mathcal{M}_p(0, z_3^2) |_{p=0}} \right) \Bigg\} \text{RGI!}$$

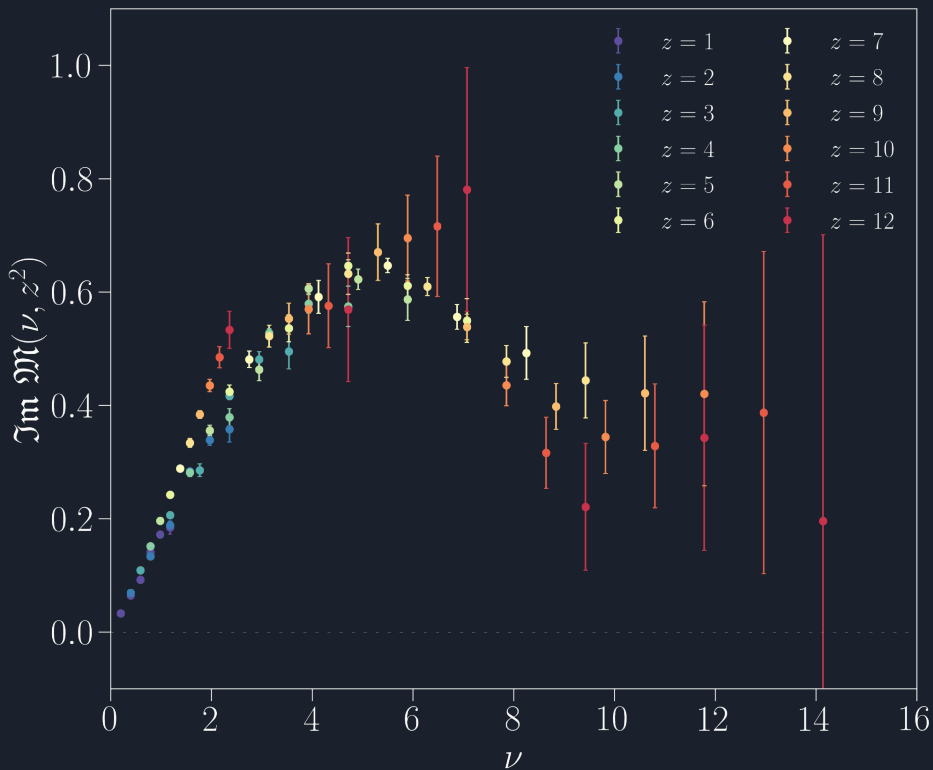
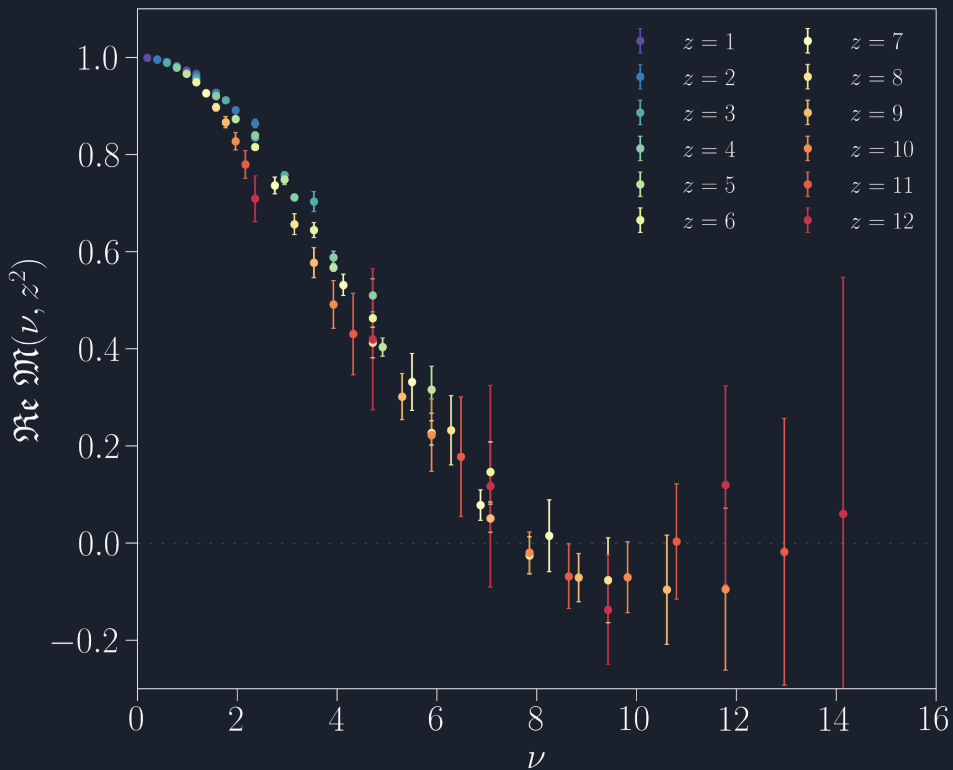


Selected Matrix Element Extractions



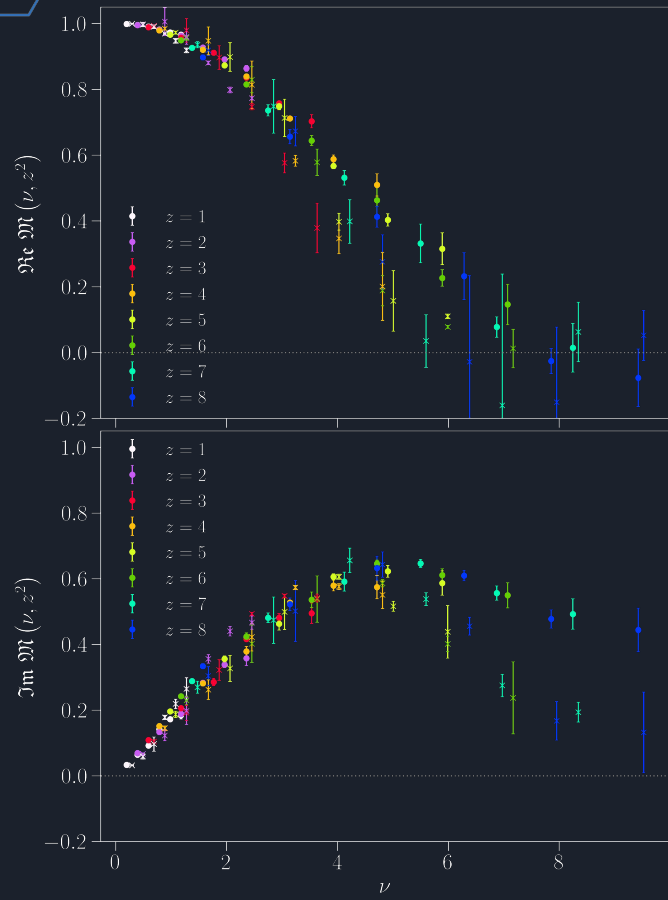


Unpolarized Ioffe-time Pseudo-Distribution

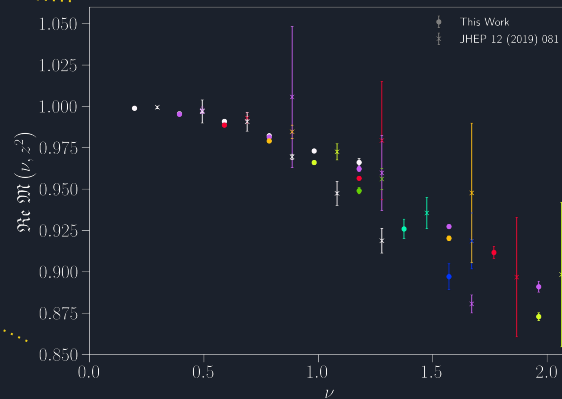
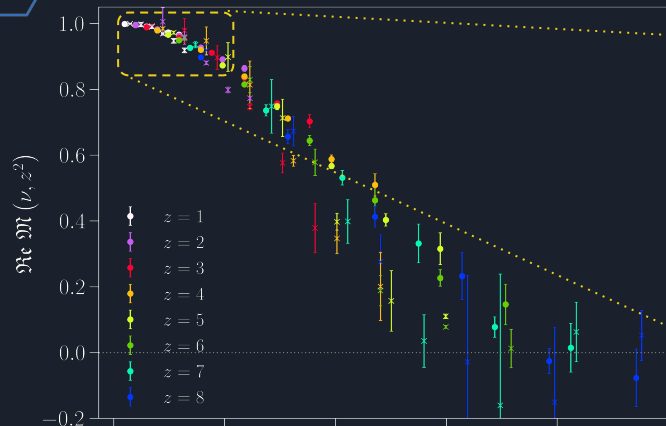




Efficacy of Distillation



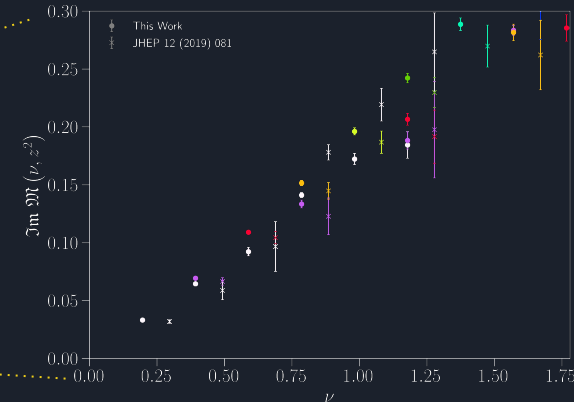
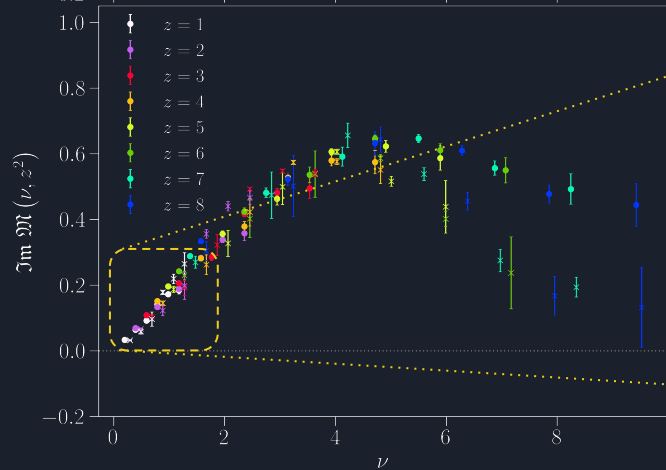
Efficacy of Distillation



B. Joó et al., JHEP 12 (2019) 081
[Gaussian smearing]

$$N_{\text{cfg}} = 417 \quad N_{\text{src}} = 8 \quad N_{\zeta} = 5$$

$$N_{\text{inv}}/\text{cfg} \simeq 8.6\text{k}$$



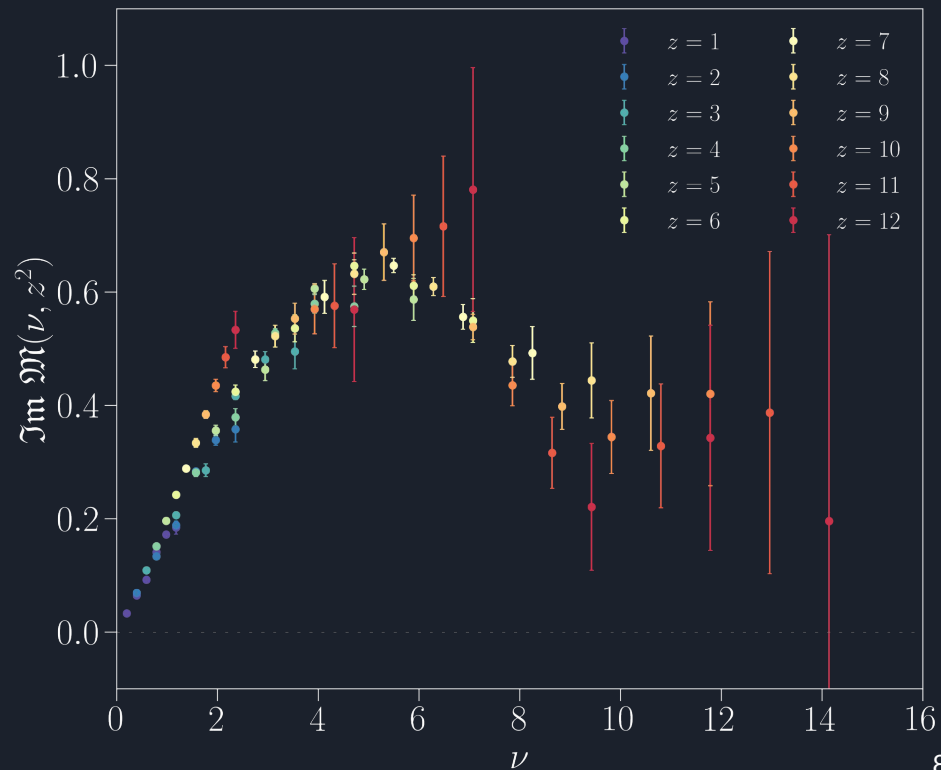
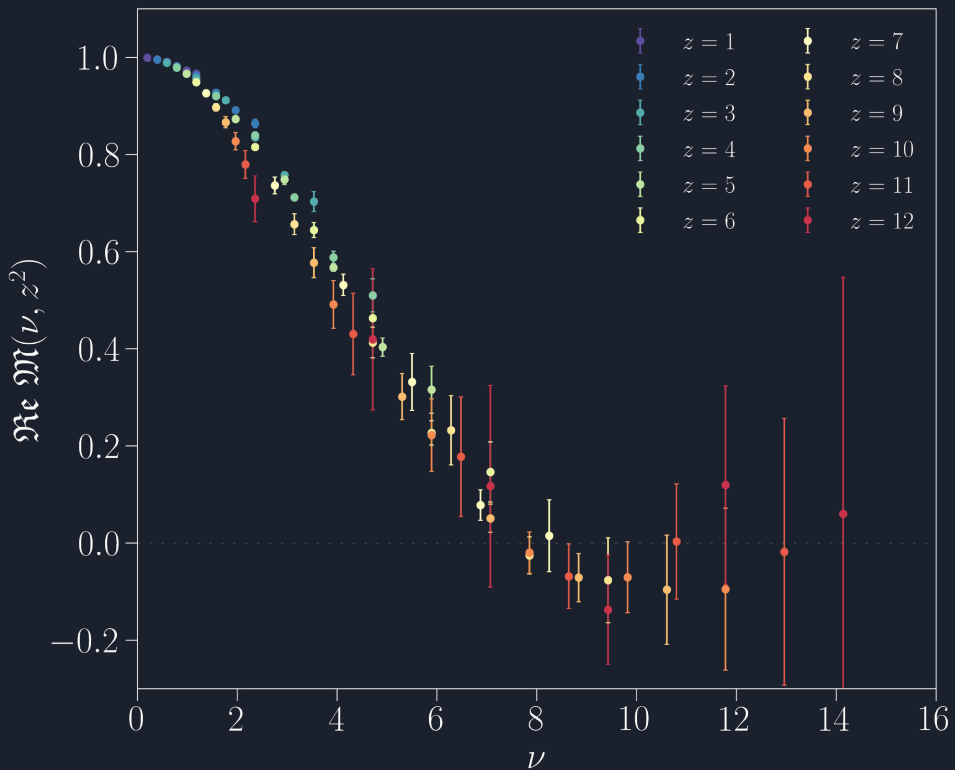
This Work
[Distillation]

$$N_{\text{cfg}} = 349 \quad N_{\text{src}} = 4 \quad N_{\zeta} = 3$$

$$N_{\text{inv}}/\text{cfg} \simeq 16\text{k}$$



Unpolarized Ioffe-time Pseudo-Distribution





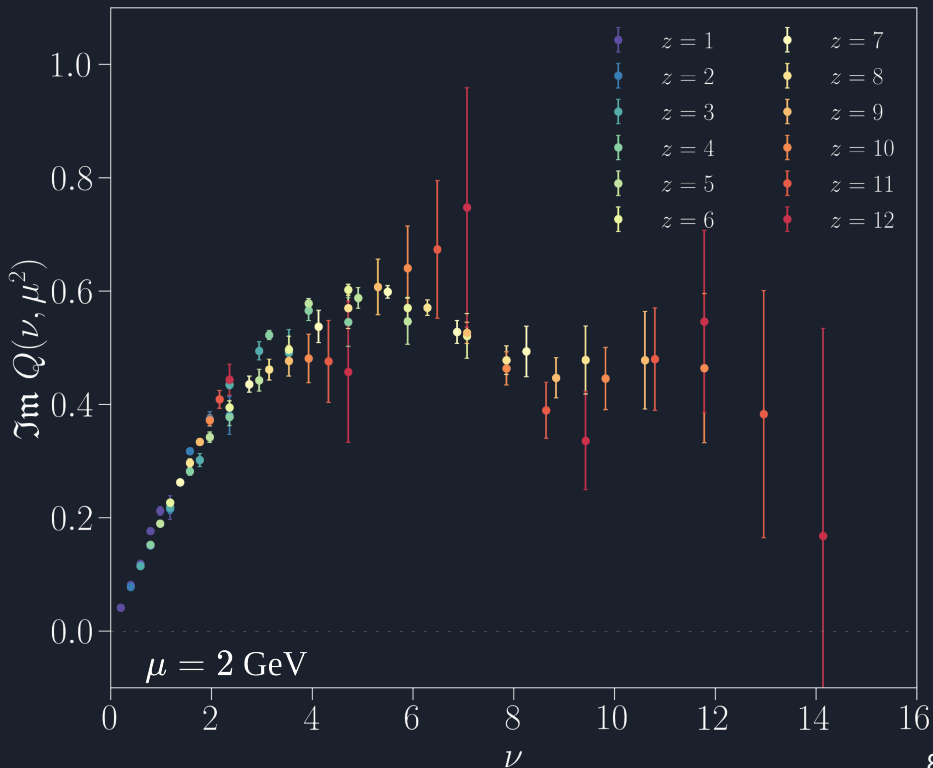
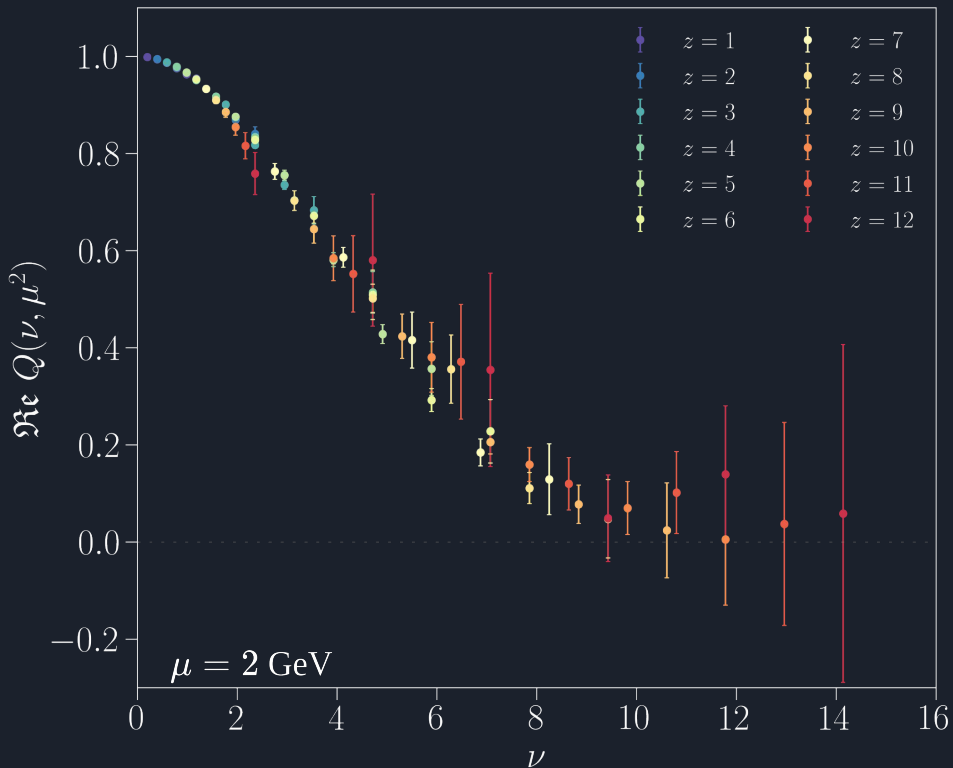
Unpolarized Ioffe-time Distribution

T. Izubuchi, et al., Phys.Rev. D98 (2018) no.5, 056004
 A. Radyushkin, Phys. Rev. D 98 (2018) no.1, 014019
 J.-H. Zhang, et al., Phys.Rev. D97 (2018) no.7, 074508

$$\mathcal{Q}(\nu, \mu^2) = \mathfrak{M}(\nu, z^2) + \frac{\alpha_s C_F}{2\pi} \int_0^1 du \left[\ln \left(\frac{e^{2\gamma_E+1} z^2 \mu^2}{4} \right) B(u) + L(u) \right] \mathfrak{M}(u\nu, z^2)$$

$$B(u) = \left(\frac{1+u^2}{1-u} \right)_+$$

$$L(u) = \left[4 \frac{\ln(1-u)}{1-u} - 2(1-u) \right]_+$$





Expand Pseudo-ITD in Orthogonal Polynomials

Ill-posed (pseudo-)ITD/PDF matching relation:
$$\mathfrak{M}(\nu, z^2) = \int_{-1}^1 dx \, \mathcal{K}(x\nu, z^2\mu^2) f_{q/h}(x, \mu^2) + \sum_{k=1}^{\infty} \mathcal{B}_k(\nu) (z^2)^k$$



Expand Pseudo-ITD in Orthogonal Polynomials

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$$\sigma_n^{(\alpha, \beta)}(\nu, z^2\mu^2) = \int_0^1 dx \mathcal{K}_\nu(x\nu, z^2\mu^2) \boxed{x^\alpha (1-x)^\beta \Omega_n^{(\alpha, \beta)}(x)}$$

$$\eta_n^{(\alpha, \beta)}(\nu, z^2\mu^2) = \int_0^1 dx \mathcal{K}_+(x\nu, z^2\mu^2) \boxed{x^\alpha (1-x)^\beta \Omega_n^{(\alpha, \beta)}(x)}$$



Expand Pseudo-ITD in Orthogonal Polynomials

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$$\Re \mathfrak{M}_{\text{fit}}(\nu, z^2) = \sum_{n=0}^{\infty} \sigma_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{\nu, n}^{lt(\alpha, \beta)}$$

$$\Im \mathfrak{M}_{\text{fit}}(\nu, z^2) = \sum_{n=0}^{\infty} \eta_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{+, n}^{lt(\alpha, \beta)}$$



Expand Pseudo-ITD in Orthogonal Polynomials

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$$\Re \mathfrak{M}_{\text{fit}}(\nu, z^2) = \sum_{n=0}^{\infty} \sigma_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{\nu, n}^{lt(\alpha, \beta)} + \boxed{\frac{a}{|z|}} \sum_{n=1}^{\infty} \sigma_{0, n}^{(\alpha, \beta)}(\nu) C_{\nu, n}^{az(\alpha, \beta)}$$

$$\Im \mathfrak{M}_{\text{fit}}(\nu, z^2) = \sum_{n=0}^{\infty} \eta_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{+, n}^{lt(\alpha, \beta)} + \boxed{\frac{a}{|z|}} \sum_{n=0}^{\infty} \eta_{0, n}^{(\alpha, \beta)}(\nu) C_{+, n}^{az(\alpha, \beta)}$$



Expand Pseudo-ITD in Orthogonal Polynomials

Ill-posed (pseudo-)ITD/PDF matching relation: $\mathfrak{M}(\nu, z^2) = \int_{-1}^1 dx \mathcal{K}(x\nu, z^2\mu^2) \boxed{f_{q/h}(x, \mu^2)} + \sum_{k=1}^{\infty} \mathcal{B}_k(\nu) (z^2)^k$

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Strategy of parametric fits with Jacobi polynomials

1. scan over truncation orders
 - a. search for optimal expansion coefficients for each
2. establish polynomial hierarchy
 - a. preference given to low-order polynomials
 - b. restrict x-space contaminating distributions to be sub-leading to leading-twist PDF
 - c. Bayesian priors (gaussian)
3. separability of non-linear optimization

$$\begin{aligned} \Re \mathfrak{M}_{\text{fit}}(\nu, z^2) &= \sum_{n=0}^{\infty} \sigma_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{\nu, n}^{lt(\alpha, \beta)} + \boxed{\frac{a}{|z|}} \sum_{n=1}^{\infty} \sigma_{0, n}^{(\alpha, \beta)}(\nu) C_{\nu, n}^{az(\alpha, \beta)} \\ &\quad + \boxed{z^2 \Lambda_{\text{QCD}}^2} \sum_{n=1}^{\infty} \sigma_{0, n}^{(\alpha, \beta)}(\nu) C_{\nu, n}^{t4(\alpha, \beta)} + \boxed{z^4 \Lambda_{\text{QCD}}^4} \sum_{n=1}^{\infty} \sigma_{0, n}^{(\alpha, \beta)}(\nu) C_{\nu, n}^{t6(\alpha, \beta)} \\ \Im \mathfrak{M}_{\text{fit}}(\nu, z^2) &= \sum_{n=0}^{\infty} \eta_n^{(\alpha, \beta)}(\nu, z^2\mu^2) C_{+, n}^{lt(\alpha, \beta)} + \boxed{\frac{a}{|z|}} \sum_{n=0}^{\infty} \eta_{0, n}^{(\alpha, \beta)}(\nu) C_{+, n}^{az(\alpha, \beta)} \\ &\quad + \boxed{z^2 \Lambda_{\text{QCD}}^2} \sum_{n=0}^{\infty} \eta_{0, n}^{(\alpha, \beta)}(\nu) C_{+, n}^{t4(\alpha, \beta)} + \boxed{z^4 \Lambda_{\text{QCD}}^4} \sum_{n=0}^{\infty} \eta_{0, n}^{(\alpha, \beta)}(\nu) C_{+, n}^{t6(\alpha, \beta)} \end{aligned}$$

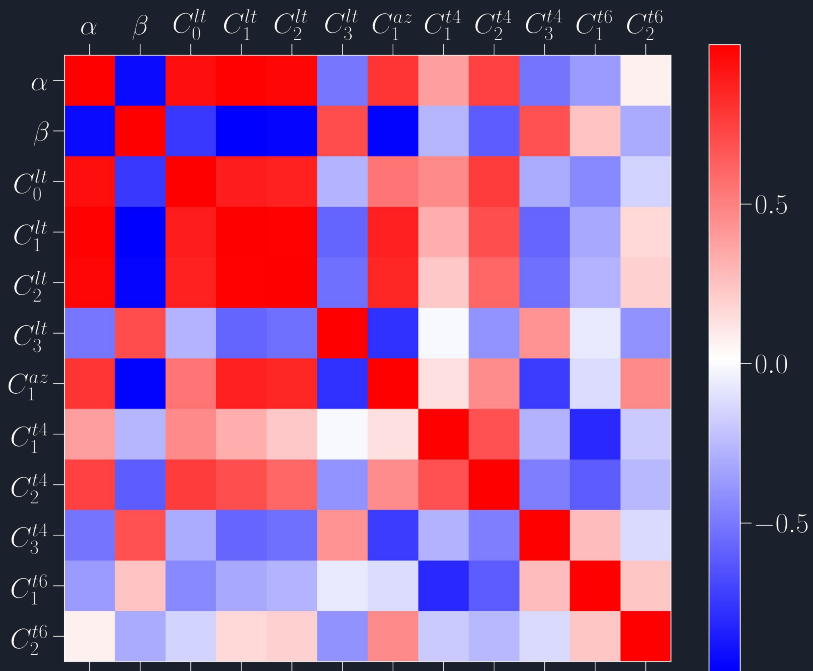
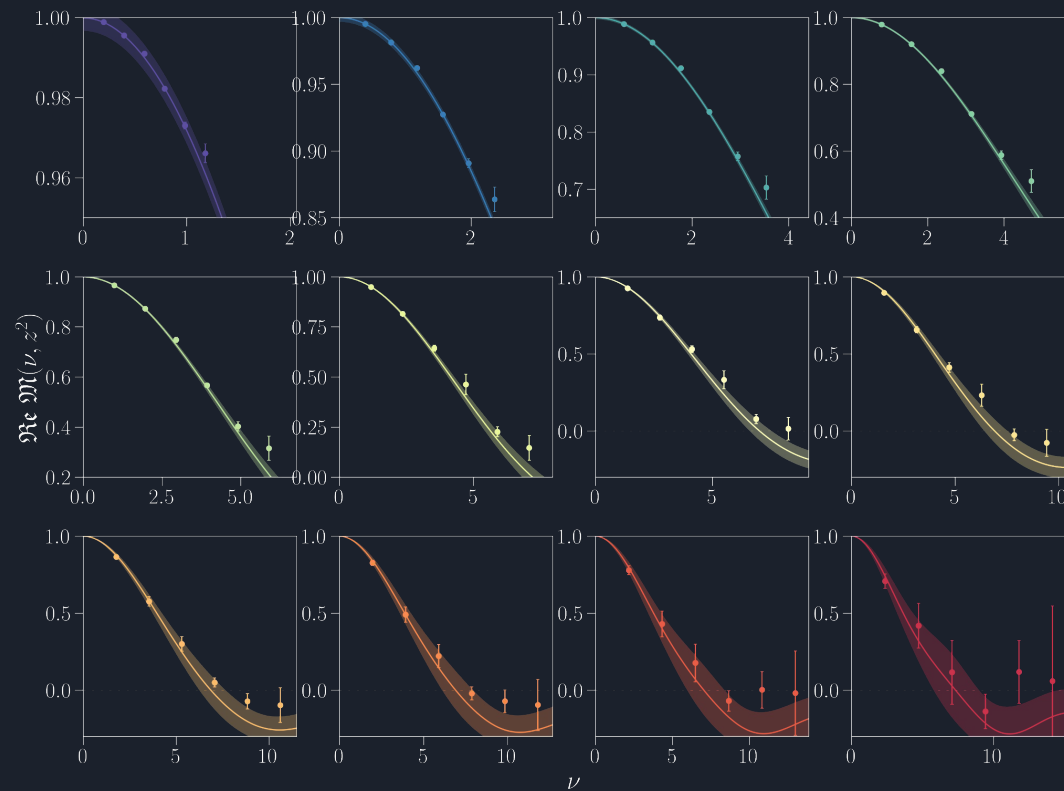


Jacobi polynomial basis are only non-linear terms
Separable non-linear optimization → variable projection

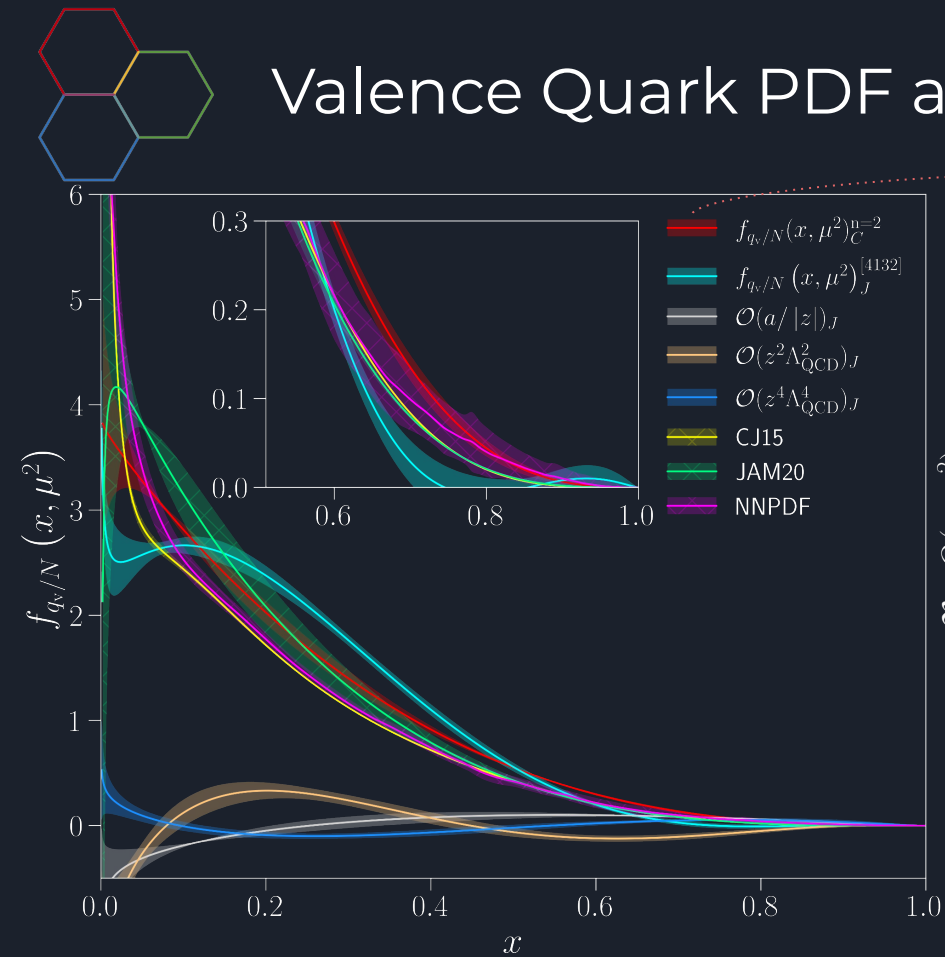
G. Golub and V. Pereyra, SIAM Journal on Numerical Analysis 10, 413 (1973)



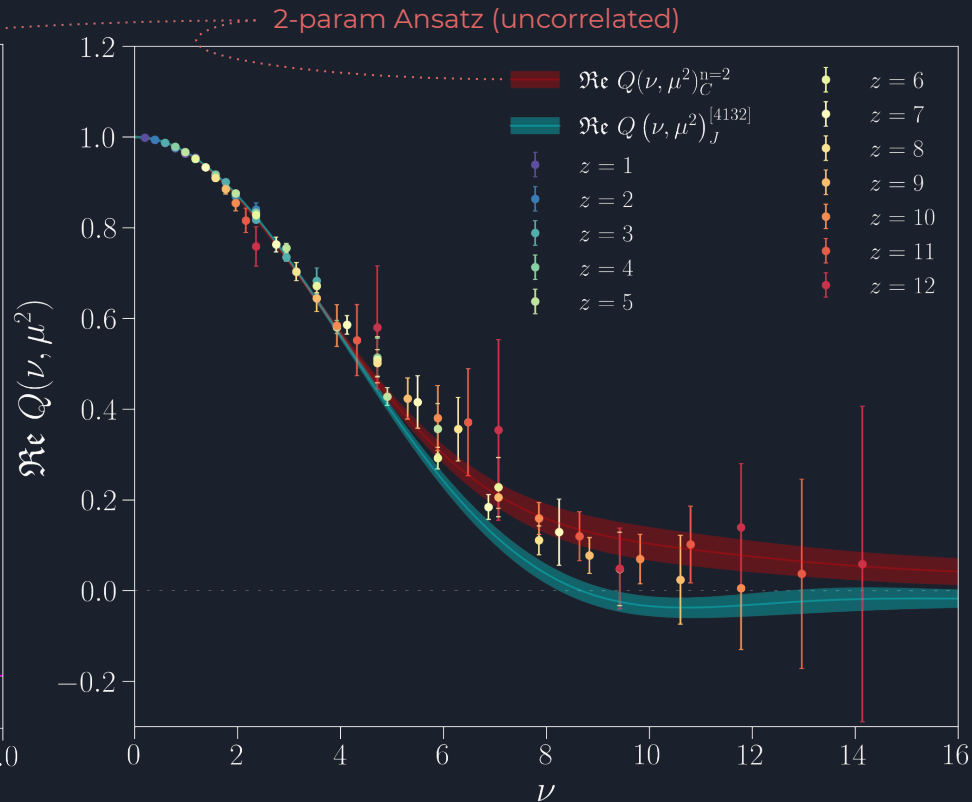
Optimal Fit for Valence PDF



Valence Quark PDF and Leading-Twist ITD



$$f_{q_v/N}(x, \mu^2)_J = x^\alpha (1-x)^\beta \sum_{n=0}^{n_{lt}} C_{v,n}^{(\alpha,\beta)} \Omega_n^{(\alpha,\beta)}(x)$$



$$\mathcal{O}(\text{corr})_J = x^\alpha (1-x)^\beta \sum_{n=1}^{n_{corr}} C_{v,n}^{corr(\alpha,\beta)} \Omega_n^{(\alpha,\beta)}(x)$$



Short-Distance Tension

Dramatic effect of a discretization correction

$\{n_{lt}, n_{az}, n_{t4}, n_{t6}\}_{v/+}$	α	β	$C_{v,0}^{lt}$	$C_{v,1}^{lt}$	$C_{v,2}^{lt}$	$C_{v,3}^{lt}$
$\{4, 1, 3, 2\}_v$	-0.209(147)	1.330(415)	1.606(257)	0.427(752)	-0.880(409)	-0.675(122)
$\{4, 0, 3, 2\}_v$	-0.376(37)	2.032(496)	1.340(165)	0.335(261)	-0.125(100)	-0.651(140)
$C_{v,1}^{az}$	$C_{v,1}^{t4}$	$C_{v,2}^{t4}$	$C_{v,3}^{t4}$	$C_{v,1}^{t6}$	$C_{v,2}^{t6}$	χ_r^2
-0.279(48)	0.052(53)	-0.371(106)	-0.407(122)	-0.045(37)	0.228(52)	2.620(345)
—	-0.090(52)	-0.112(77)	0.274(99)	0.011(39)	0.397(84)	45.68(1.72)



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Dramatic effect of a discretization correction

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$C_{v,1}^{az}$	$C_{v,1}^{t4}$	$C_{v,2}^{t4}$	$C_{v,3}^{t4}$	$C_{v,1}^{t6}$	$C_{v,2}^{t6}$	χ_r^2
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—	-0.090(52)	-0.112(77)	0.274(99)	0.011(39)	0.397(84)	45.68(1.72)

Visualize scale dependence in reduced pseudo-ITD via mock pseudo-PDF fit

$$\Re \mathcal{M}_{\text{fit}}(\nu, z^2) = \int_0^1 dx \cos(x\nu) \Re \mathcal{P}(x, z^2; \alpha, 3)$$

$$\frac{\Gamma(5+\alpha)}{\Gamma(1+\alpha)\Gamma(4)} x^\alpha (1-x)^3$$



Short-Distance Tension

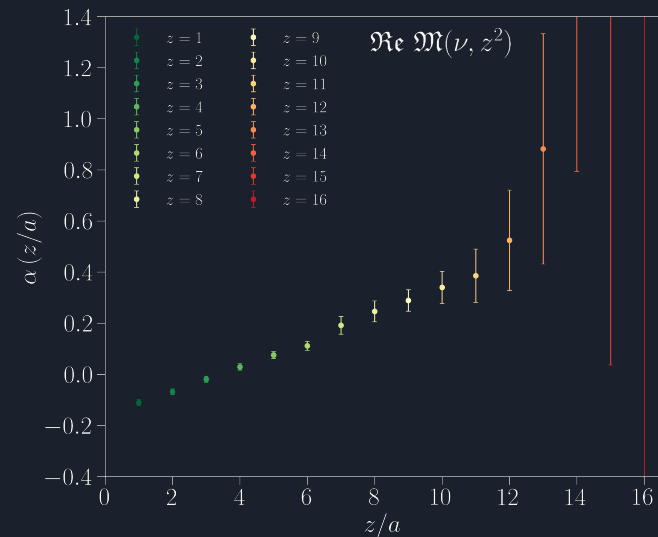
Dramatic effect of a discretization correction

$\{n_{lt}, n_{az}, n_{t4}, n_{t6}\}_{v/+}$	α	β	$C_{v,0}^{lt}$	$C_{v,1}^{lt}$	$C_{v,2}^{lt}$	$C_{v,3}^{lt}$
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-0.279(48)	0.052(53)	-0.371(106)	-0.407(122)	-0.045(37)	0.228(52)	2.620(345)
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$$\frac{\Gamma(5+\alpha)}{\Gamma(1+\alpha)\Gamma(4)} x^\alpha (1-x)^3$$





Short-Distance Tension

Dramatic effect of a discretization correction

$\{n_{lt}, n_{az}, n_{t4}, n_{t6}\}_{v/+}$	α	β	$C_{v,0}^{lt}$	$C_{v,1}^{lt}$	$C_{v,2}^{lt}$	$C_{v,3}^{lt}$
$\{4, 1, 3, 2\}_v$	-0.209(147)	1.330(415)	1.606(257)	0.427(752)	-0.880(409)	-0.675(122)
$\{4, 0, 3, 2\}_v$	-0.376(37)	2.032(496)	1.340(165)	0.335(261)	-0.125(100)	-0.651(140)
$C_{v,1}^{az}$	$C_{v,1}^{t4}$	$C_{v,2}^{t4}$	$C_{v,3}^{t4}$	$C_{v,1}^{t6}$	$C_{v,2}^{t6}$	χ_r^2
-0.279(48)	0.052(53)	-0.371(106)	-0.407(122)	-0.045(37)	0.228(52)	2.620(345)
—	-0.090(52)	-0.112(77)	0.274(99)	0.011(39)	0.397(84)	45.68(1.72)

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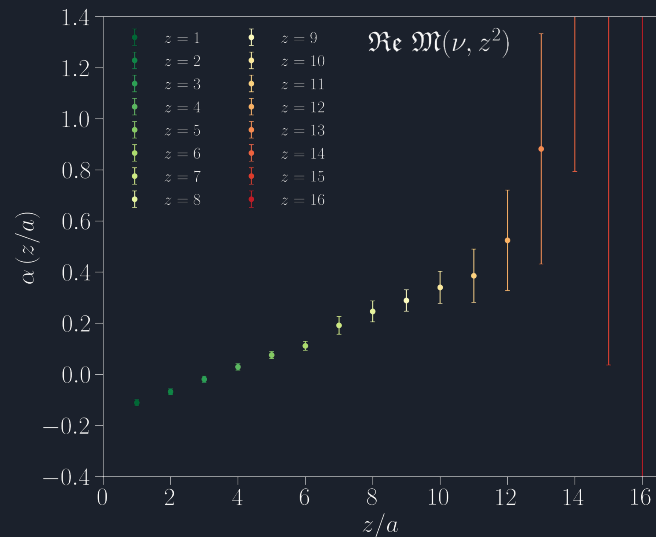
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Evolution/matching with pseudo-PDF fit

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➤ redo two-parameter fits to matched ITD





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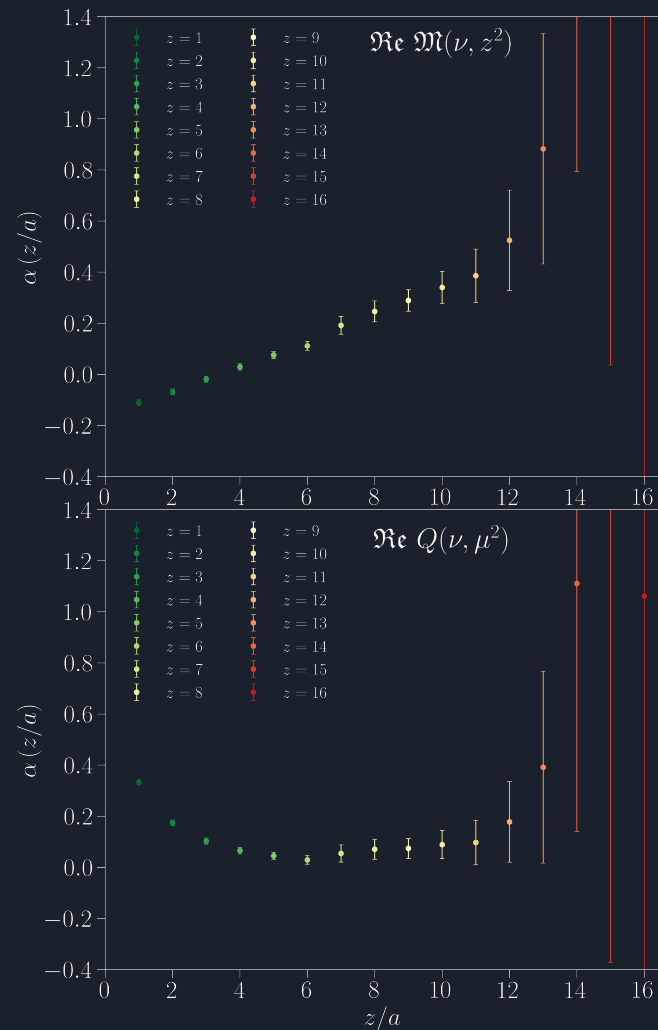
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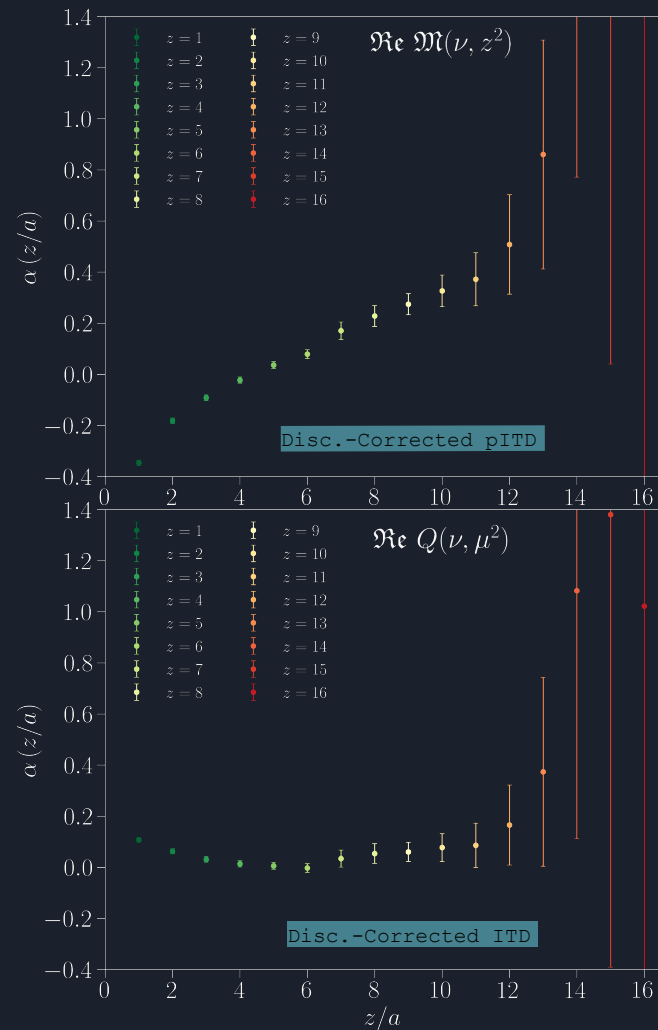
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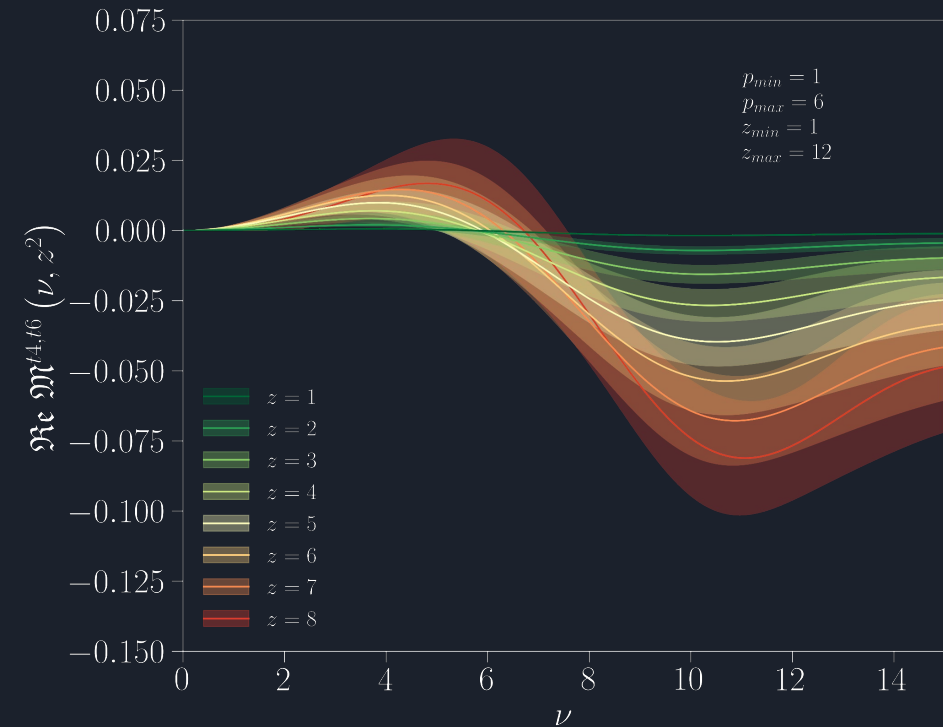
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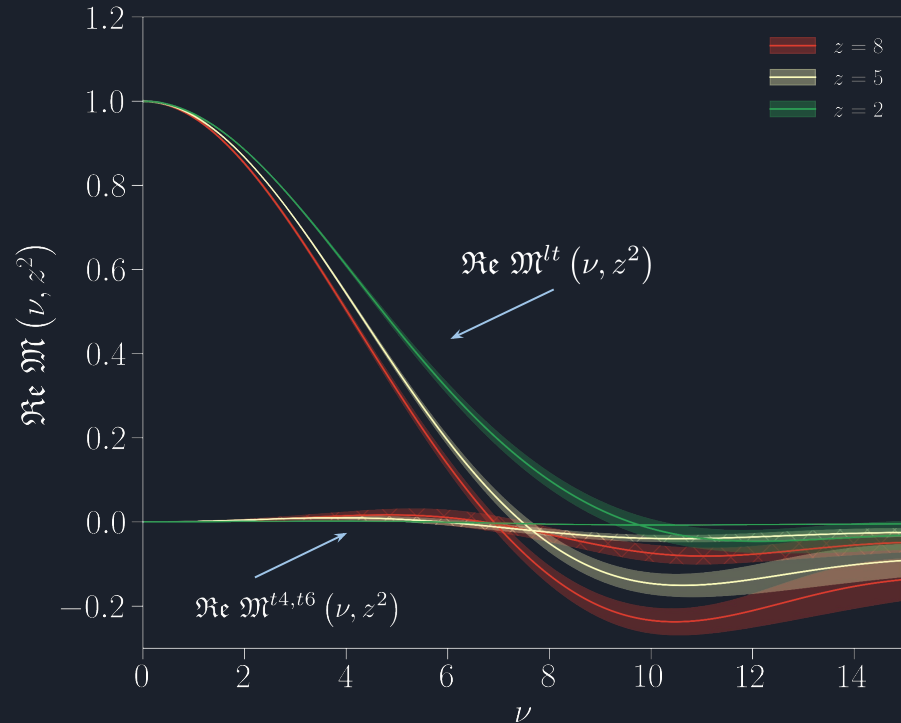
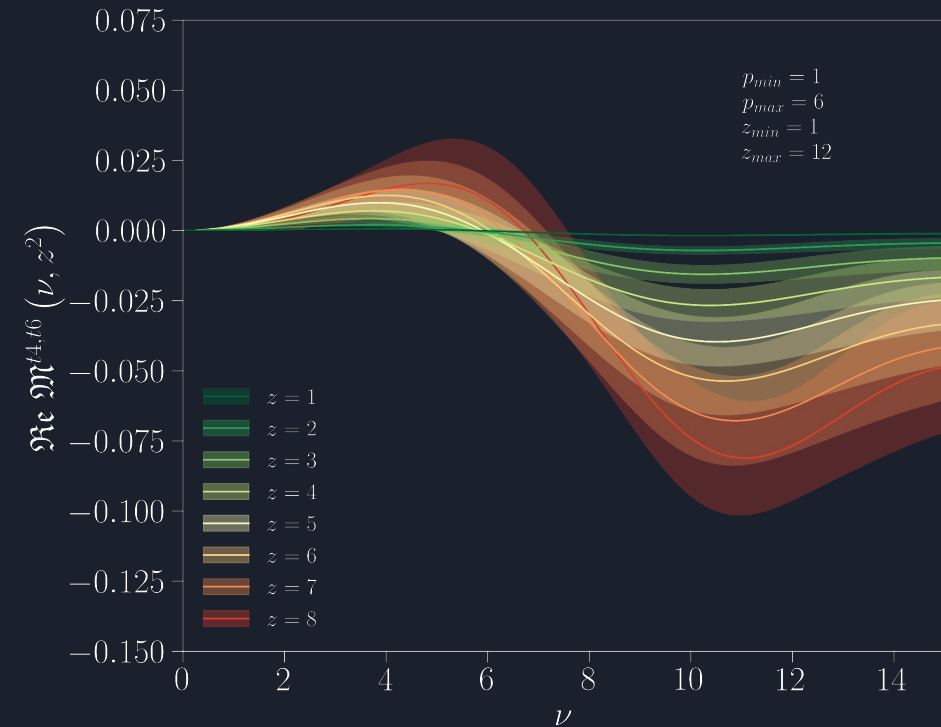
Parameterized Higher-Twist Distribution



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



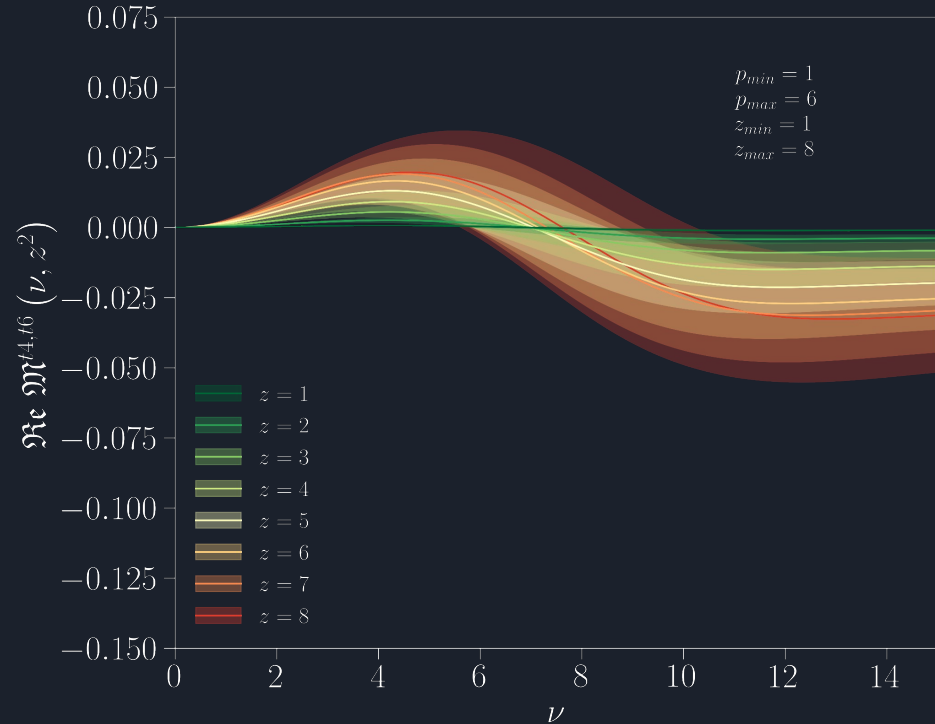
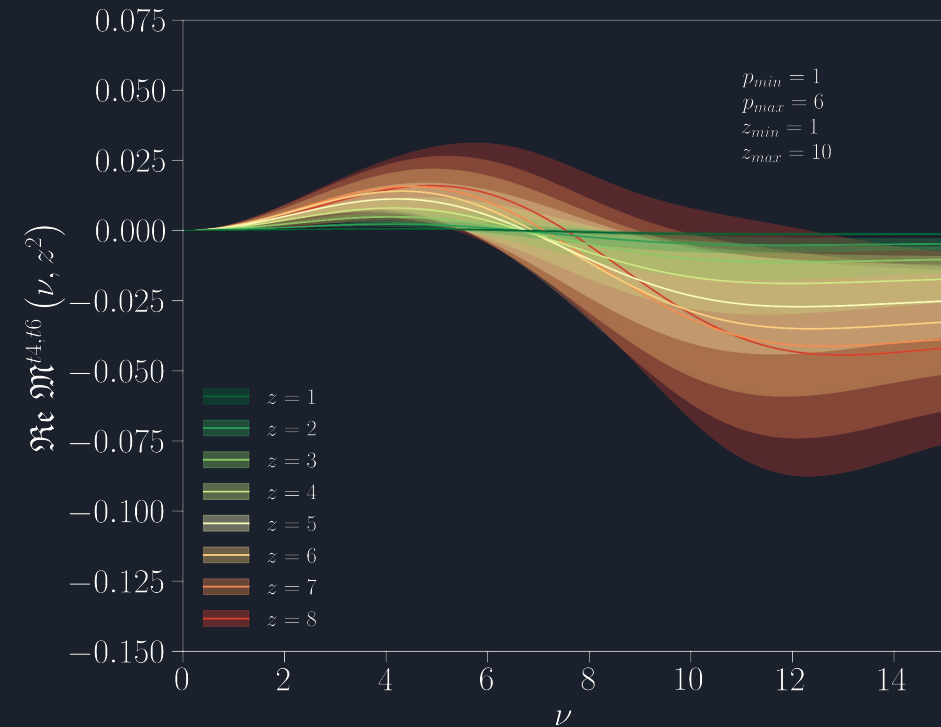
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Higher-Twist Variability - Cuts on Wilson Line



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



Next Steps - Expanded Interpolator Basis

Excited-state contamination

- optimize operator/state overlaps - saturate correlation functions at early temporal separations



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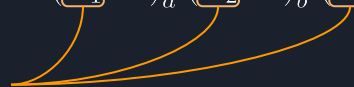
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Generic light-quark nucleon interpolator smeared with distillation

$$\mathcal{O}_i(t) = \epsilon^{abc} (\mathcal{D}_1 \not{u})_a^\alpha (\mathcal{D}_2 \not{d})_b^\beta (\mathcal{D}_3 \not{u})_c^\gamma (t) S_i^{\alpha\beta\gamma}$$

Dirac structure/covariant derivatives





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Discretized continuum-like interpolators of definite permutational symmetries

$$\mathcal{O}_B = (\mathcal{F}_{\mathcal{P}(\mathbf{F})} \otimes \mathcal{S}_{\mathcal{P}(\mathbf{S})} \otimes \mathcal{D}_{\mathcal{P}(\mathbf{D})}) \{q_1 q_2 q_3\} \quad (N_M \otimes (\tfrac{1}{2}^+)_M^1 \otimes D_{L=1,A}^{[2]})^{J^P=\frac{1}{2}^+} \equiv N^2 P_A \tfrac{1}{2}^+$$



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(Generally) Continuum spins reducible under octahedral group

Canonical subductions

- spinors/derivatives combined into object of definite J^P

$$\mathcal{O}_{n_{\Lambda},r}^{\{J\}} = \sum_m S_{n_{\Lambda},r}^{J,m} \mathcal{O}^{\{J,m\}}$$

R. Edwards, et. al., Phys. Rev. D84, 074508 (2011)

J. Dudek and R. Edwards, Phys. Rev. D85, 054016 (2012)

Helicity subductions

C. Thomas, et al., Phys. Rev. D85, 014507 (2012)

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- boost breaks \mathcal{O}_h^P symmetry to little groups

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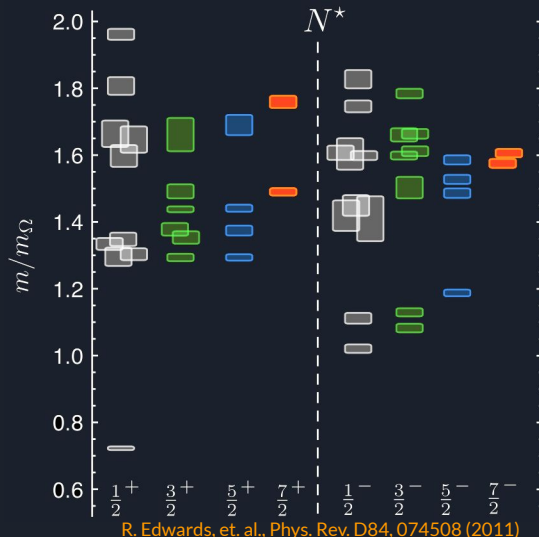
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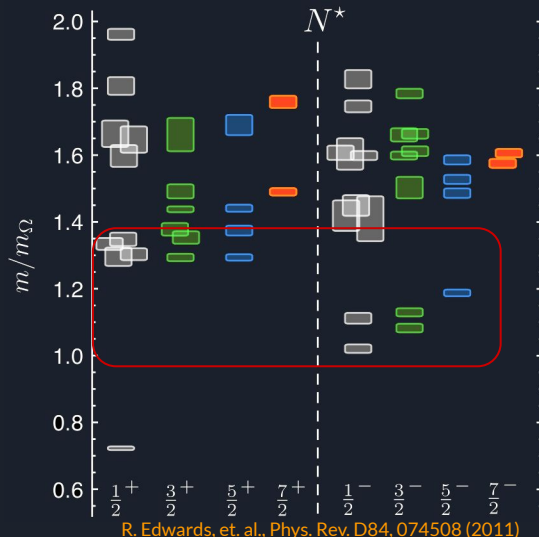
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Summary and Outlook

Hadronic structure accessible from certain lattice calculable matrix elements

- short-distance factorization

Nucleon valence (plus) quark PDF

- distillation (+phasing) - first use in structure studies
 - precise pseudo-ITDs & PDFs
- systematic effects can be reliably addressed

Fidelity of PDFs extracted from Lattice QCD hinges on control of systematic effects

- observed deviation from expected DGLAP evolution of pseudo-PDF at small distances
- neglect of data correlations/Ioffe-time cuts
 - useful, but can yield erroneous results

Repeat calculations at lighter pion masses & finer lattice spacings underway

HadStruc Collaboration



Robert Edwards, **CE**, Nikhil Karthik,
Jianwei Qiu, David Richards, Eloy Romero, Frank Winter ^[1]

Balint Joó^[2]

Carl Carlson, Tanjib Khan, Christopher Monahan,
Kostas Orginos, Raza Sufian^[3]

Wayne Morris, Anatoly Radyushkin ^[4]

Joe Karpie^[5]

Savvas Zafeiropoulos^[6]

Yan-Qing Ma^[7]

Jefferson Lab ^[1], Oak Ridge ^[2], William and Mary ^[3], Old Dominion University ^[4], Columbia University ^[5], Aix Marseille University ^[6], Peking University ^[7]



Supplements





An Ill-Posed Inverse

(pseudo-)ITD/PDF matching relations are ill-posed

$$\mathfrak{M}(\nu, z^2) = \int_{-1}^1 dx \mathcal{K}(x\nu, z^2\mu^2) f_{q/h}(x, \mu^2) + \sum_{k=1}^{\infty} \mathcal{B}_k(\nu) (z^2)^k$$

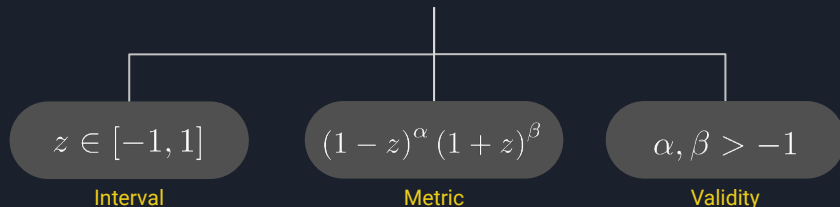
$$\mathcal{Q}(\nu, \mu^2) = \int_{-1}^1 dx e^{i\nu x} f_{q/h}(x, \mu^2)$$

➤ commonly: matched ITD fit by F.T. of parameterized PDF

$$f_{q/h}(x, \mu^2) = N x^\alpha (1-x)^\beta P(x)$$

Jacobi (hypergeometric) polynomials

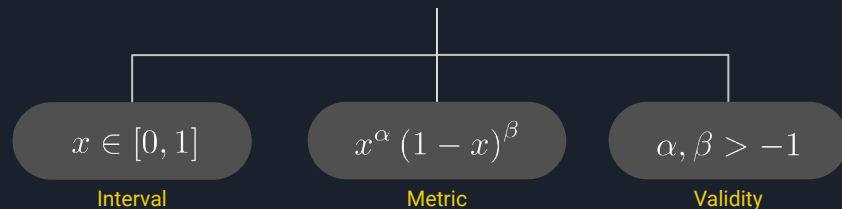
$$P_n^{(\alpha, \beta)}(z) = \frac{\Gamma(\alpha + n + 1)}{n! \Gamma(\alpha + \beta + n + 1)} \sum_{j=0}^n \binom{n}{j} \frac{\Gamma(\alpha + \beta + n + j + 1)}{\Gamma(\alpha + j + 1)} \left(\frac{z-1}{2}\right)^j$$



$$\int_{-1}^1 dz (1-z)^\alpha (1+z)^\beta P_n^{(\alpha, \beta)}(z) P_m^{(\alpha, \beta)}(z) = \delta_{n,m} h_n(\alpha, \beta)$$

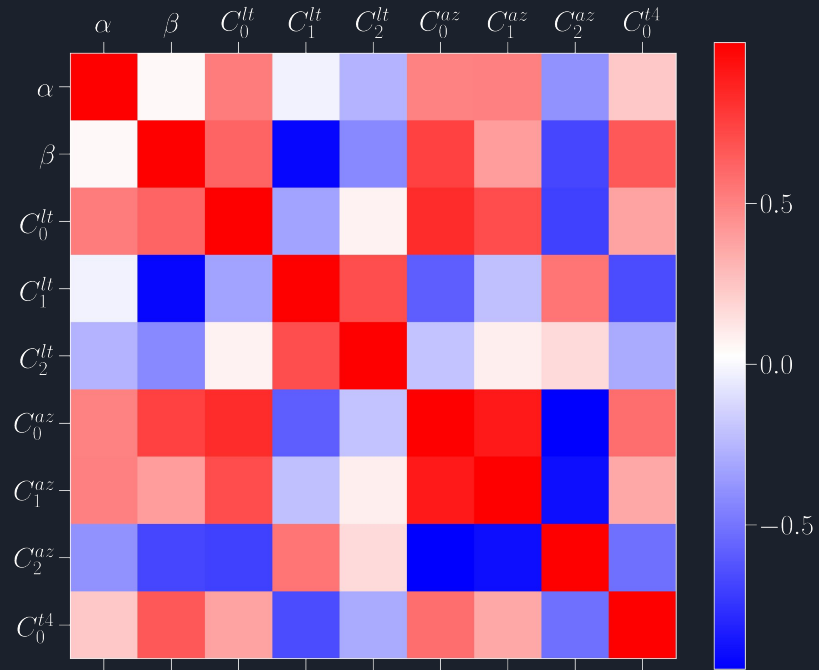
A convenient change of variables: $z \mapsto 1 - 2x$

$$\Omega_n^{(\alpha, \beta)}(x) = \sum_{j=0}^n \underbrace{\frac{\Gamma(\alpha + n + 1)}{n! \Gamma(\alpha + \beta + n + 1)} \binom{n}{j} \frac{(-1)^j \Gamma(\alpha + \beta + n + j + 1)}{\Gamma(\alpha + j + 1)}}_{\omega_{n,j}^{(\alpha, \beta)}} x^j$$



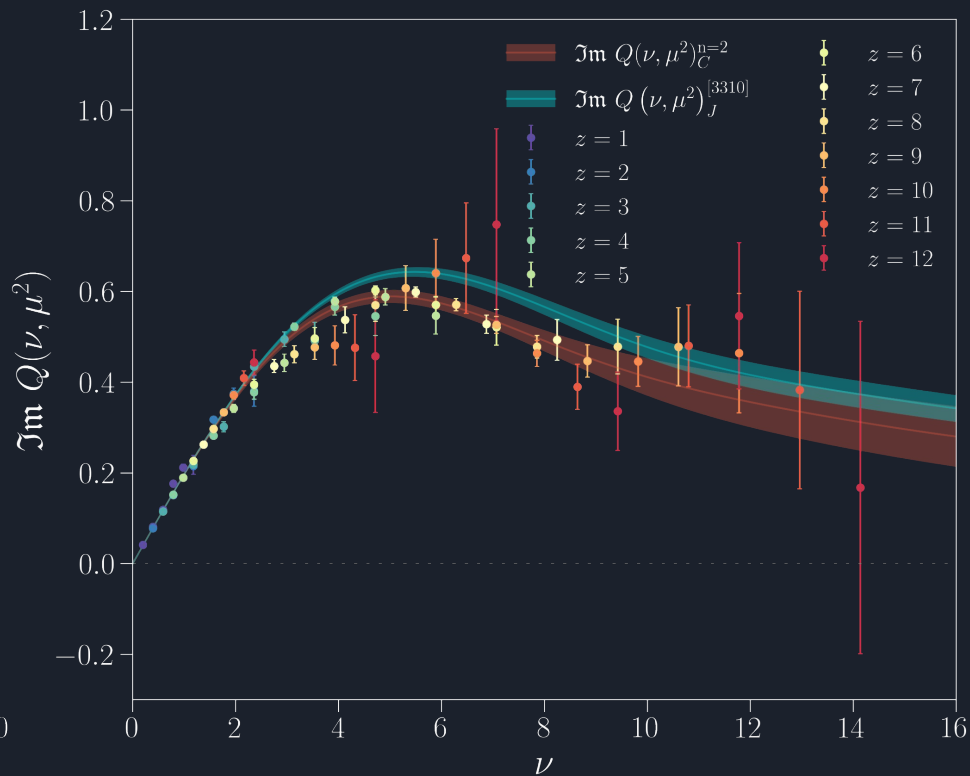
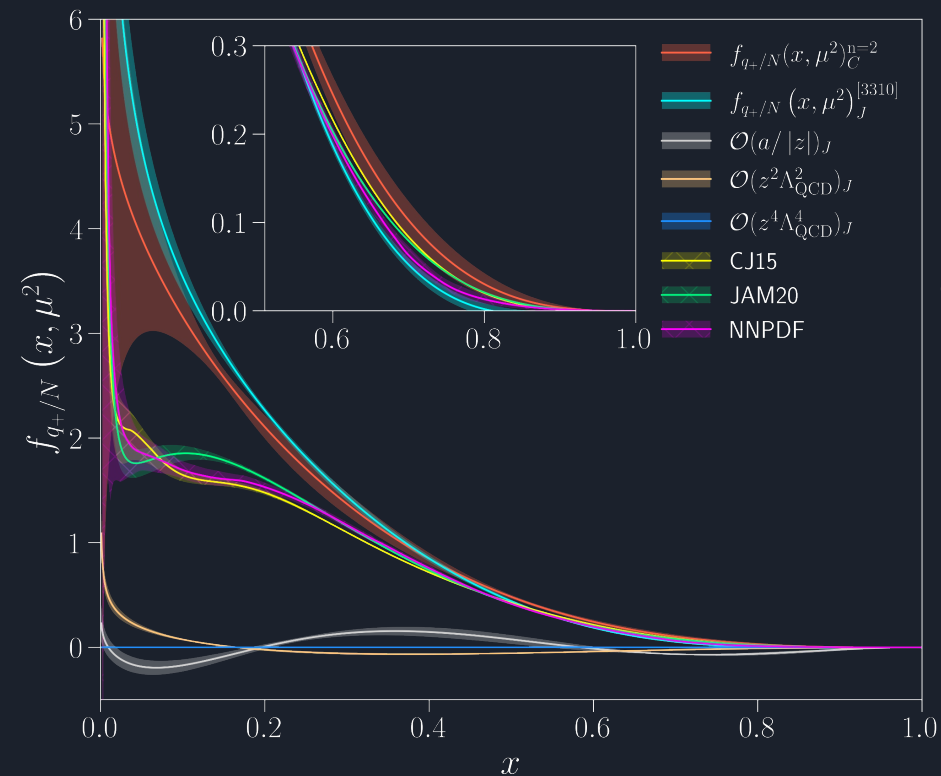
Flexibility of PDF functional form captured without bias via $\{\Omega_n^{(\alpha, \beta)}\}$

$$f_{q/h}(x) = x^\alpha (1-x)^\beta \sum_{n=0}^{\infty} C_{q,n}^{(\alpha, \beta)} \Omega_n^{(\alpha, \beta)}(x)$$





Plus Quark PDF and Leading-Twist ITD





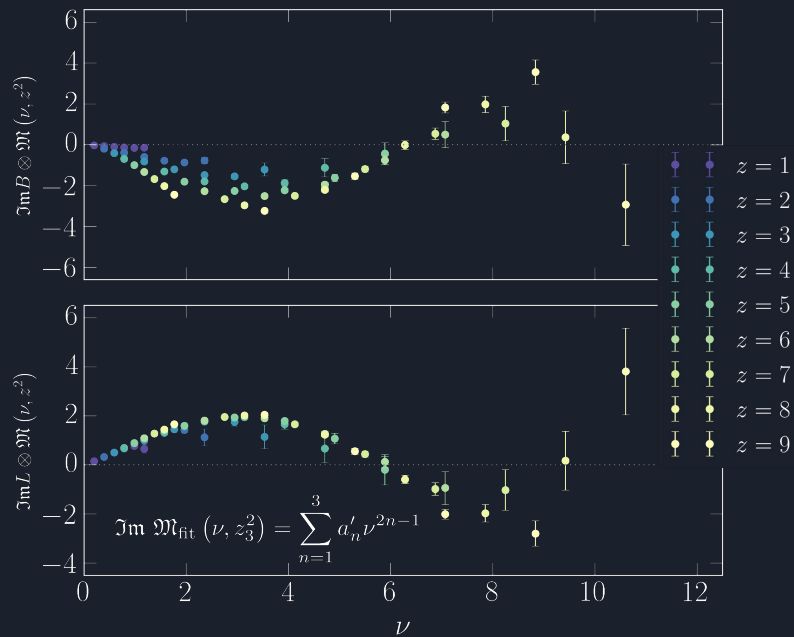
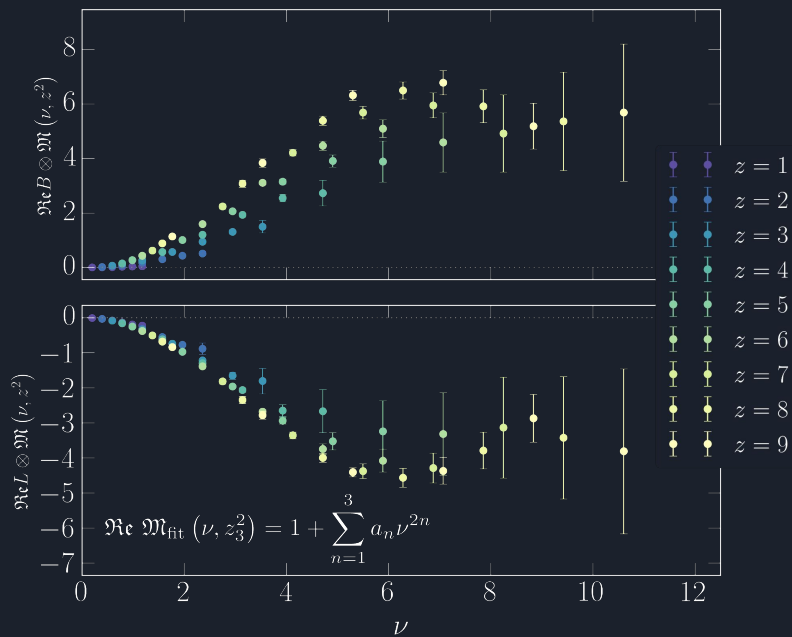
Evolution and Scheme Conversion

Matching reduced pseudo-ITD to ITD requires a continuous description

$$\mathcal{Q}(\nu, \mu^2) = \mathfrak{M}(\nu, z^2) + \frac{\alpha_s C_F}{2\pi} \int_0^1 du \left[\ln \left(\frac{e^{2\gamma_E + 1} z^2 \mu^2}{4} \right) B(u) + L(u) \right] \mathfrak{M}(u\nu, z^2)$$

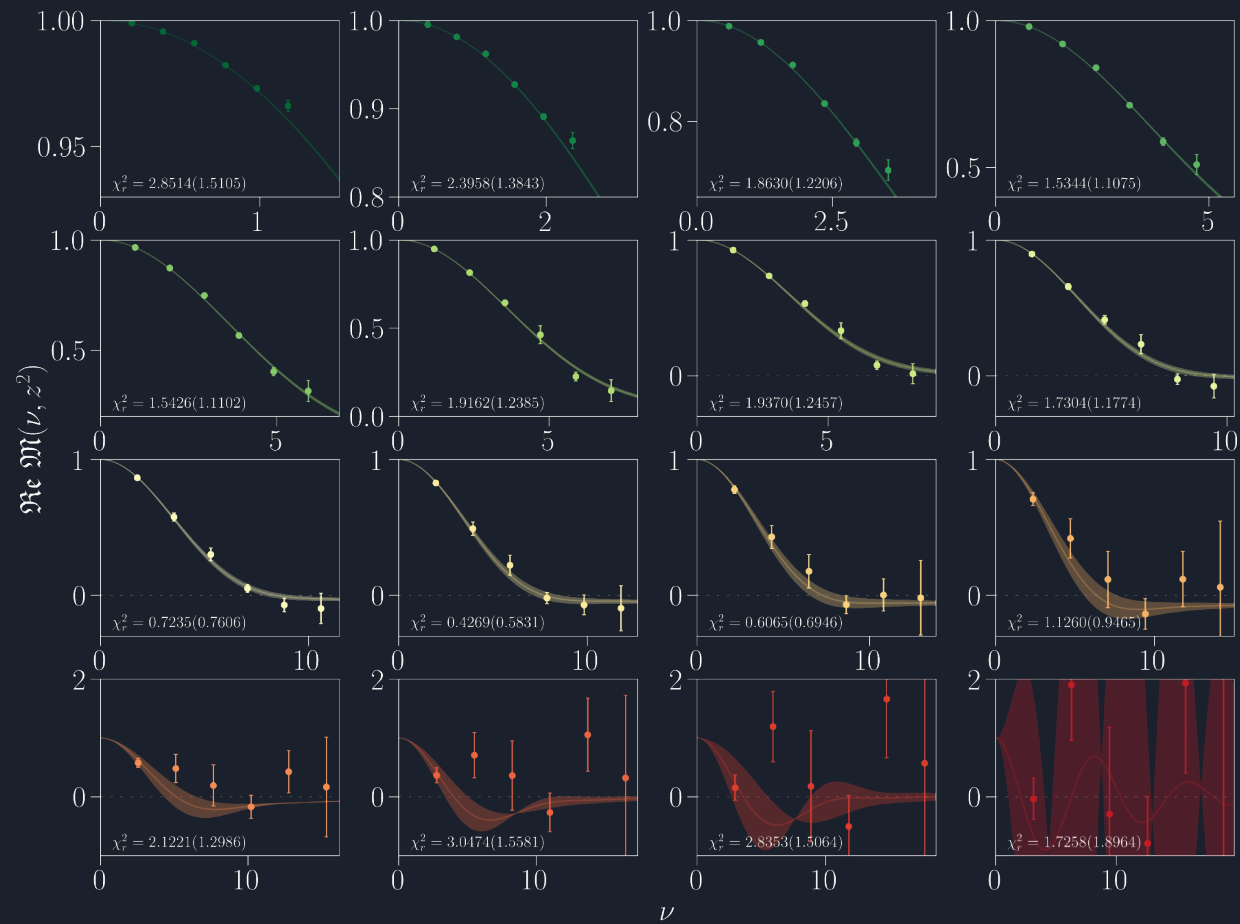
$$B(u) = \left(\frac{1+u^2}{1-u} \right)_+ \quad L(u) = \left[4 \frac{\ln(1-u)}{1-u} - 2(1-u) \right]_+$$

T. Izubuchi, et al., Phys.Rev. D98 (2018) no.5, 056004
 A. Radyushkin, Phys.Lett. B781 (2018) 433-442
 A. Radyushkin, Phys. Rev. D 98 (2018) no.1, 014019
 J.-H. Zhang, et al., Phys.Rev. D97 (2018) no.7, 074508



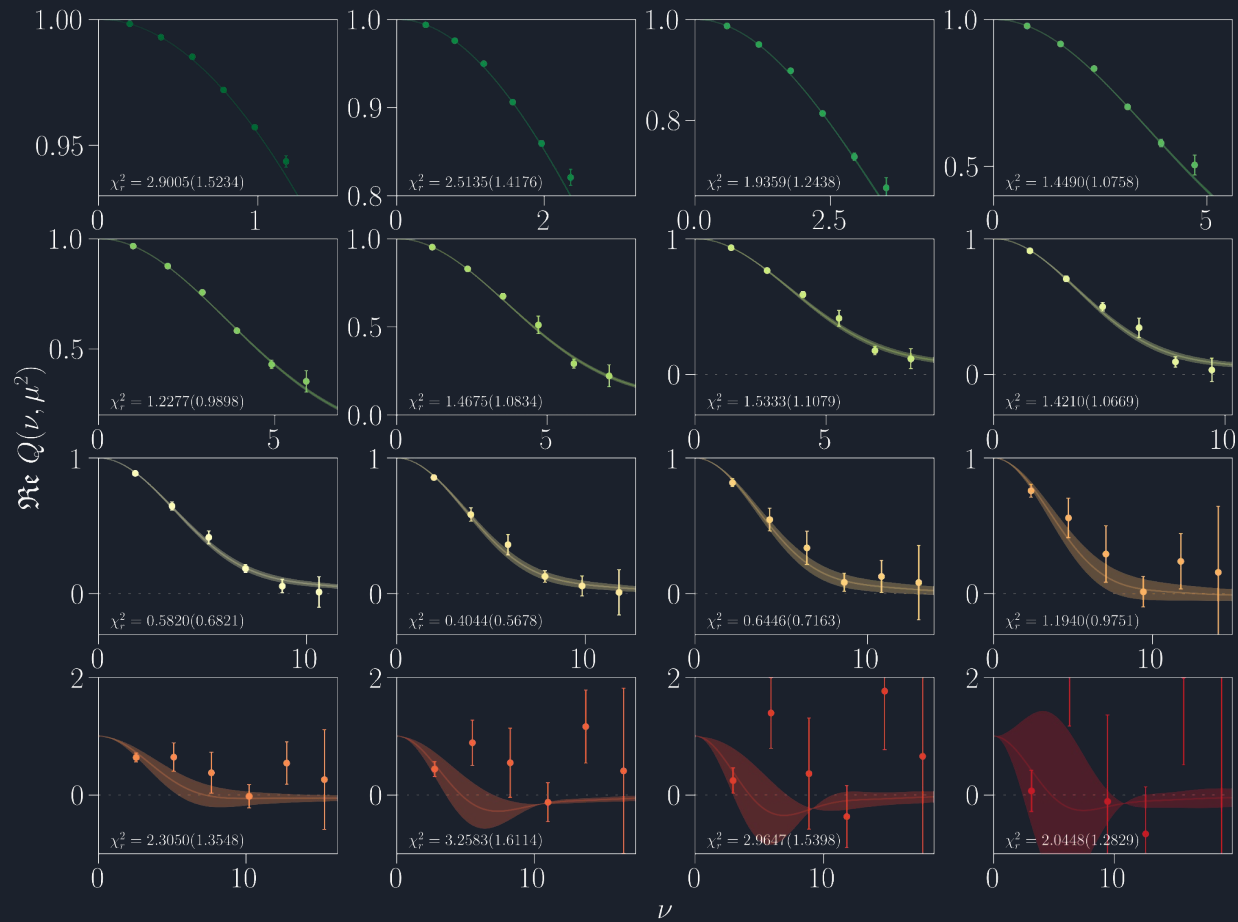


Pseudo-PDF Fit of Reduced Pseudo-ITD

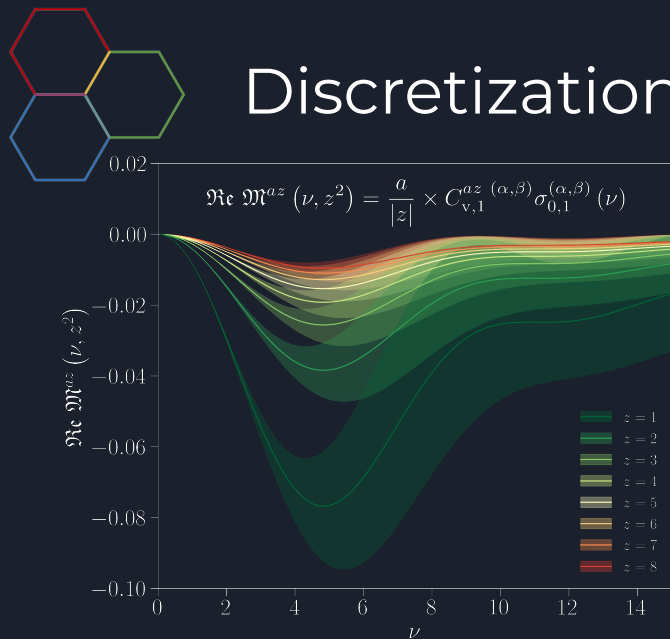




ITD Fit - Pseudo-PDF Form for Convolution

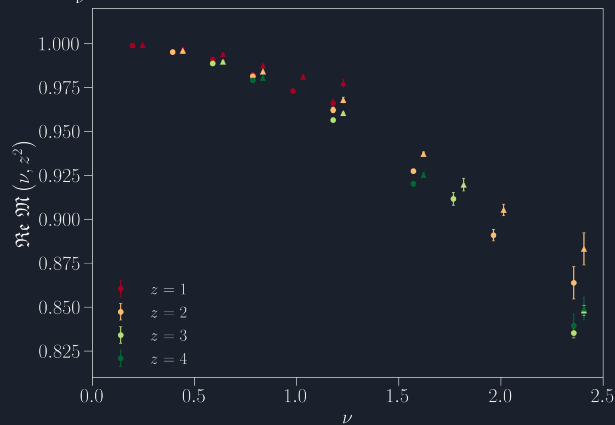


Discretization Effect



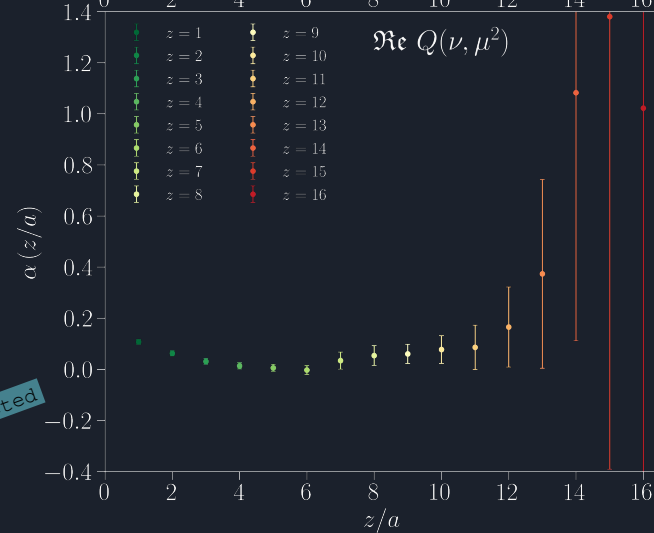
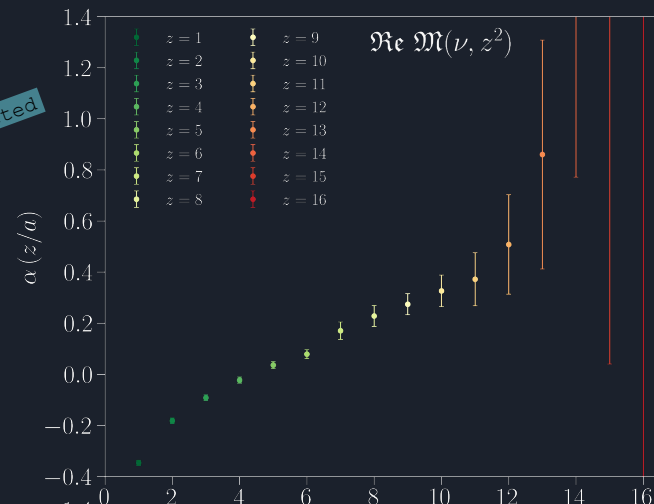
Removal of parameterized discretization effect

- (numerically) slight vertical shift of small- z data
- effect most pronounced for $\nu \lesssim 4.5$



Disc.-Corrected
pITD

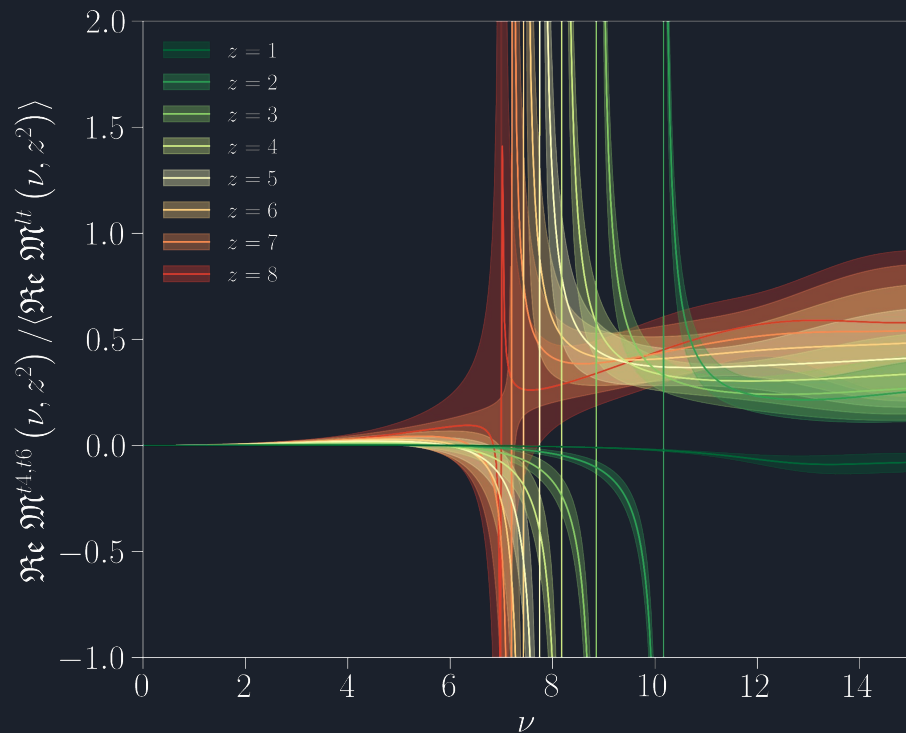
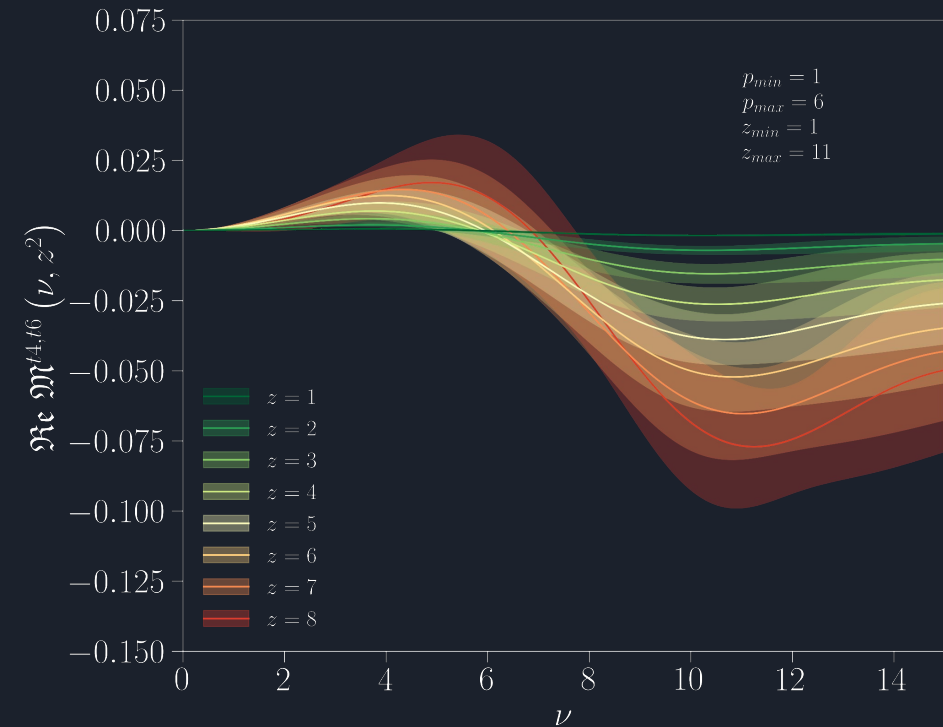
Repeat DGLAP check in (pseudo-)ITD



Disc.-Corrected
ITD



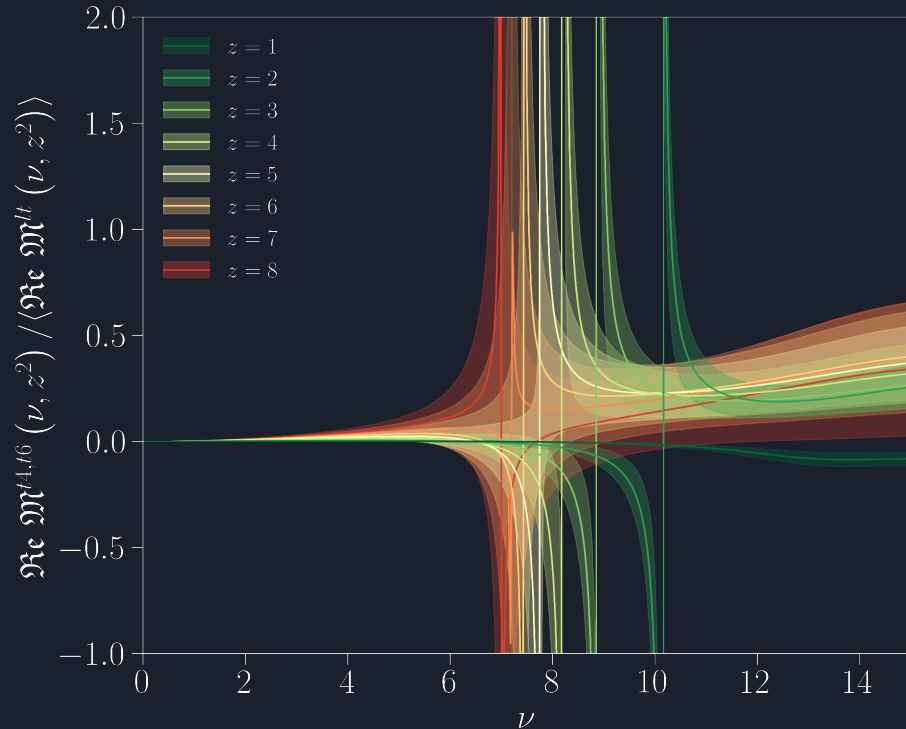
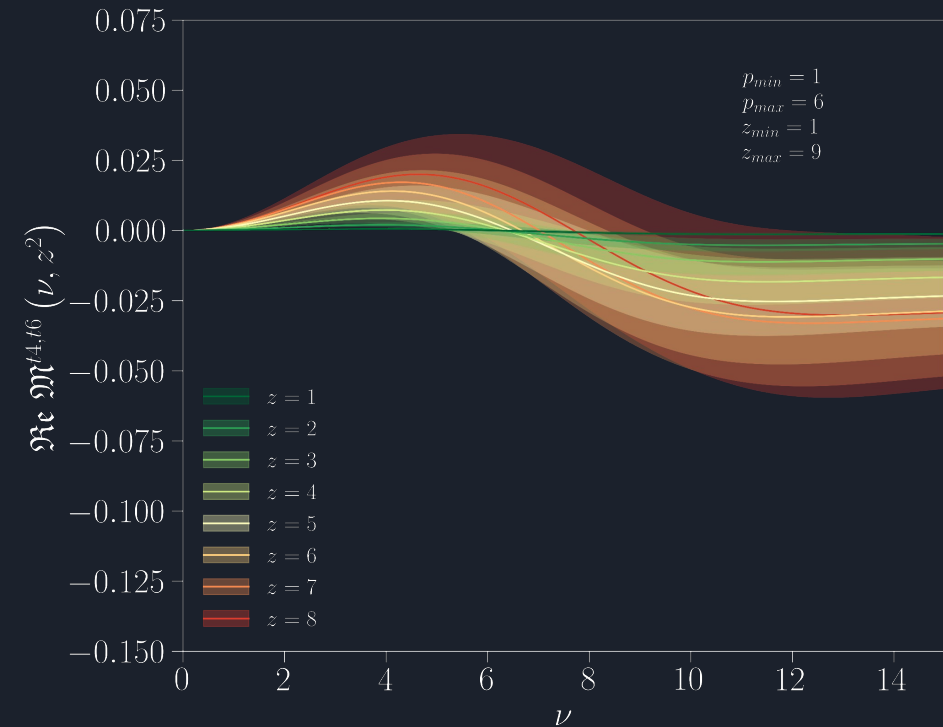
Higher-Twist Variability - Cuts on Wilson Line



$$\text{Re } \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



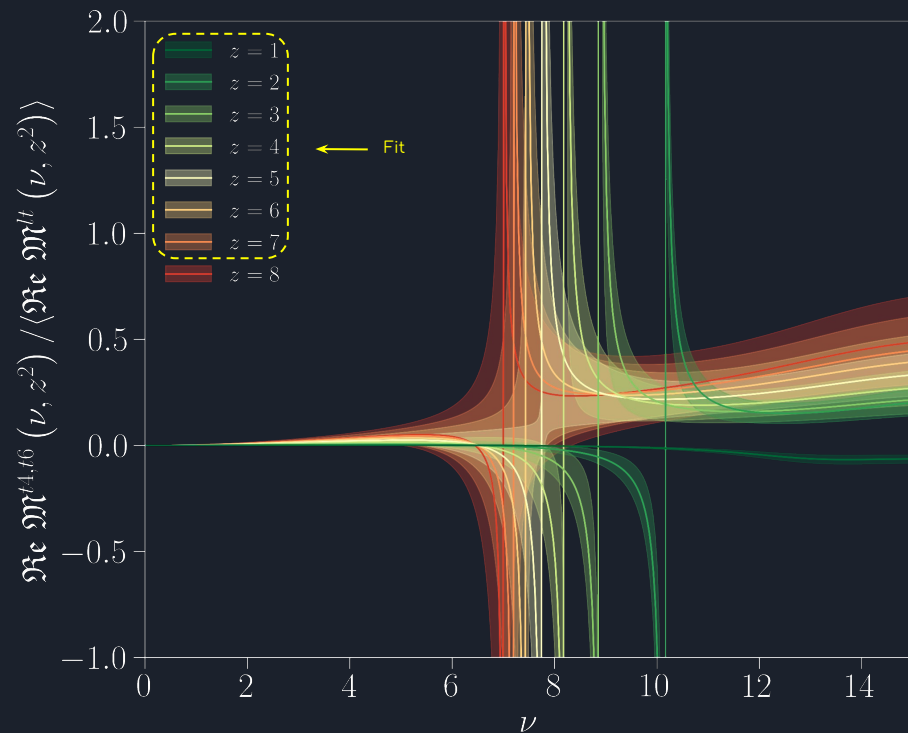
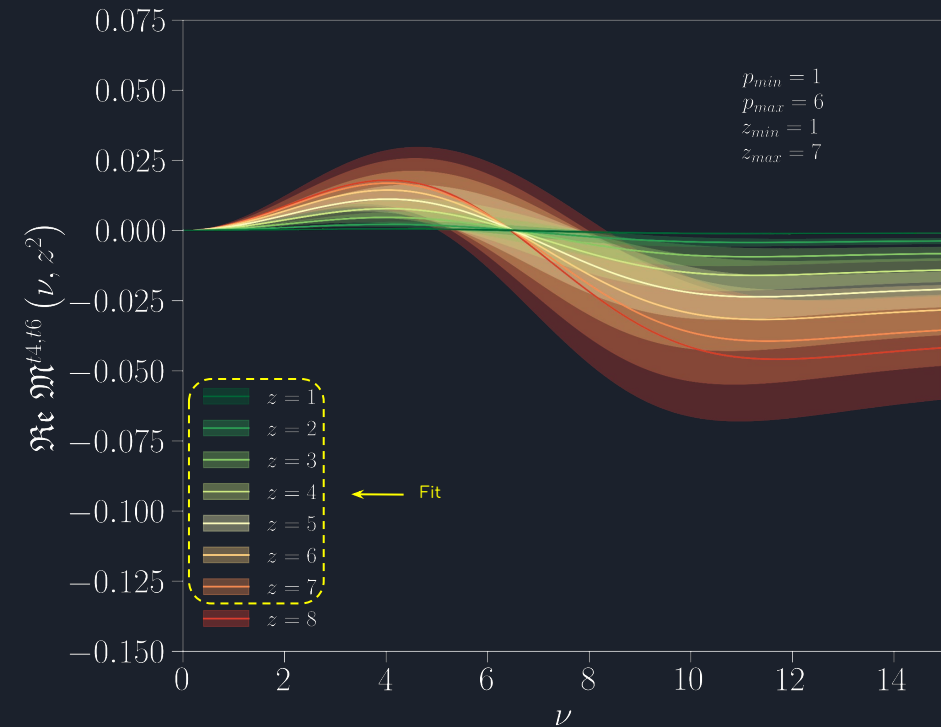
Higher-Twist Variability - Cuts on Wilson Line



$$\text{Re } \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



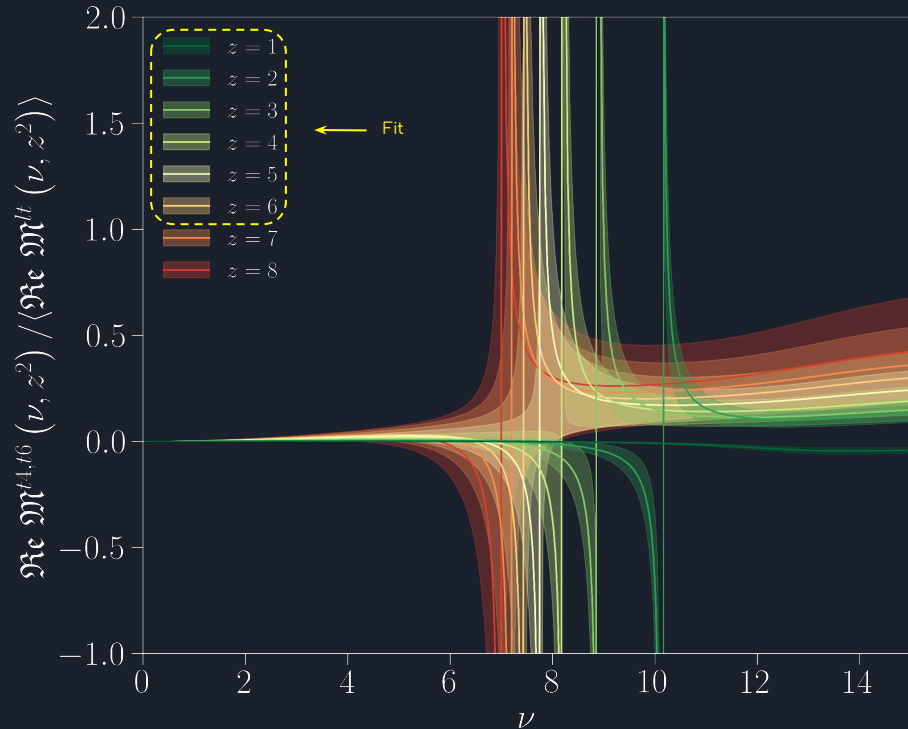
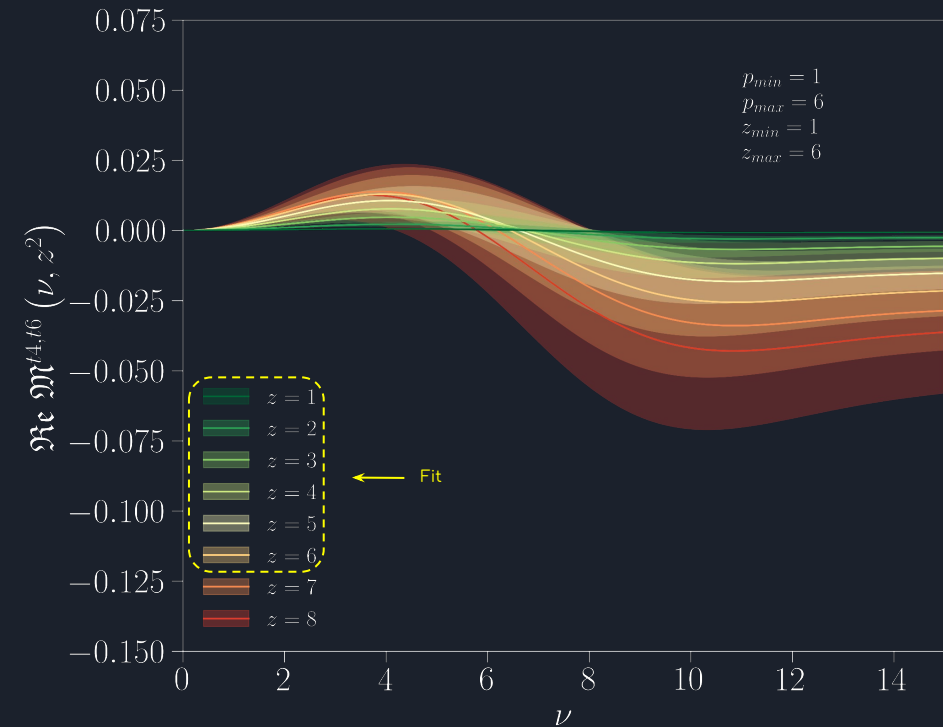
Higher-Twist Variability - Cuts on Wilson Line



$$\text{Re } \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



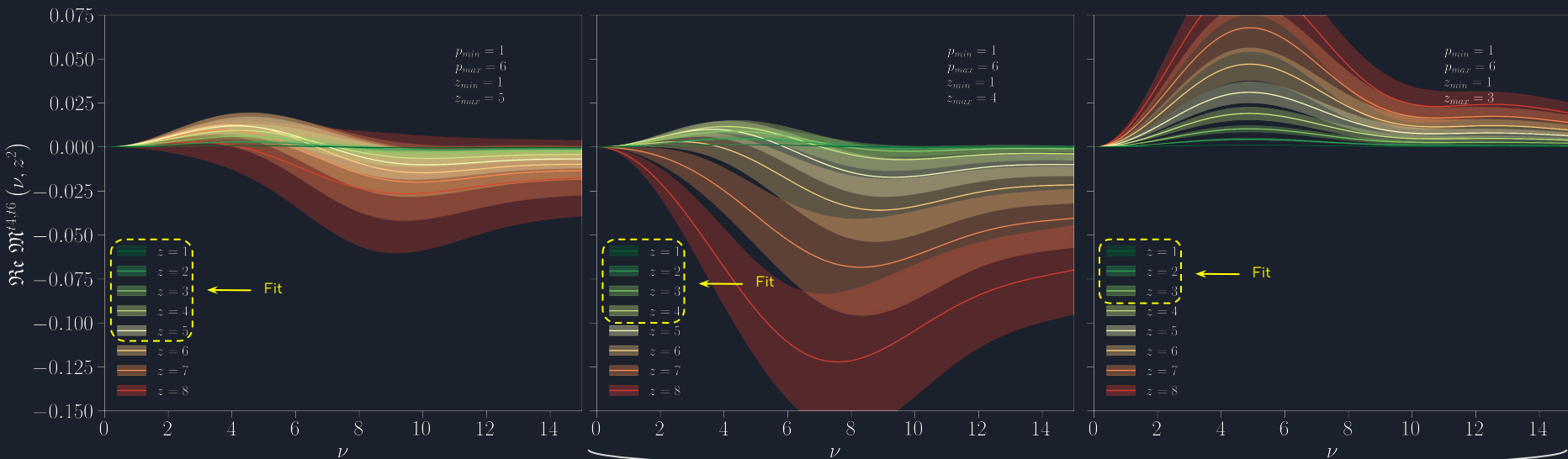
Higher-Twist Variability - Cuts on Wilson Line



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



Higher-Twist Variability - Cuts on Wilson Line

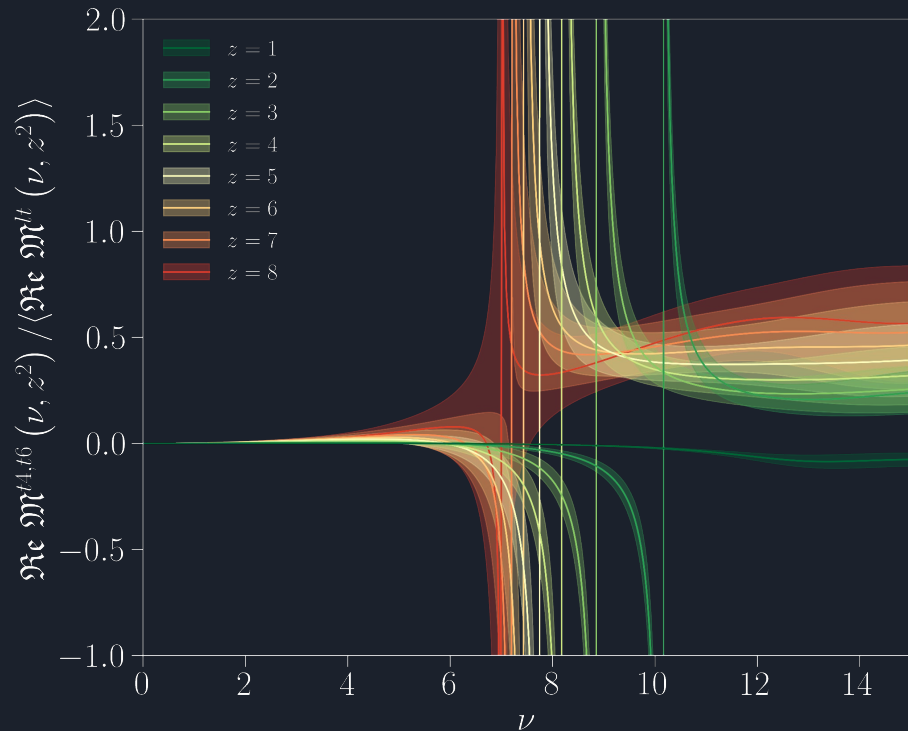
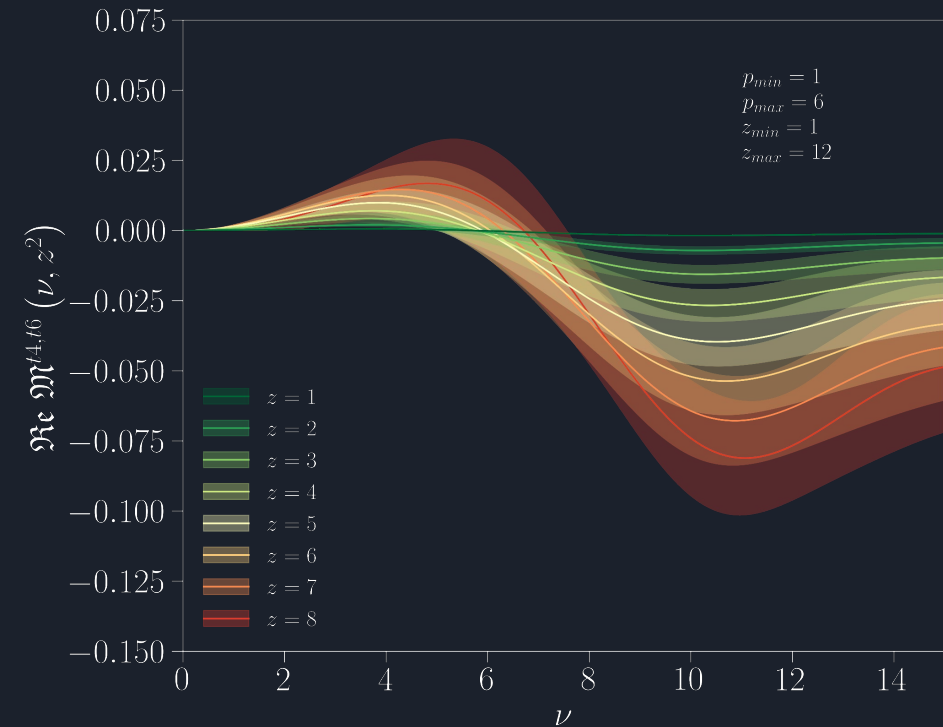


Large variability of parameterized higher-twist effects beyond Wilson line cut
 \Rightarrow effect no longer constrained

$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



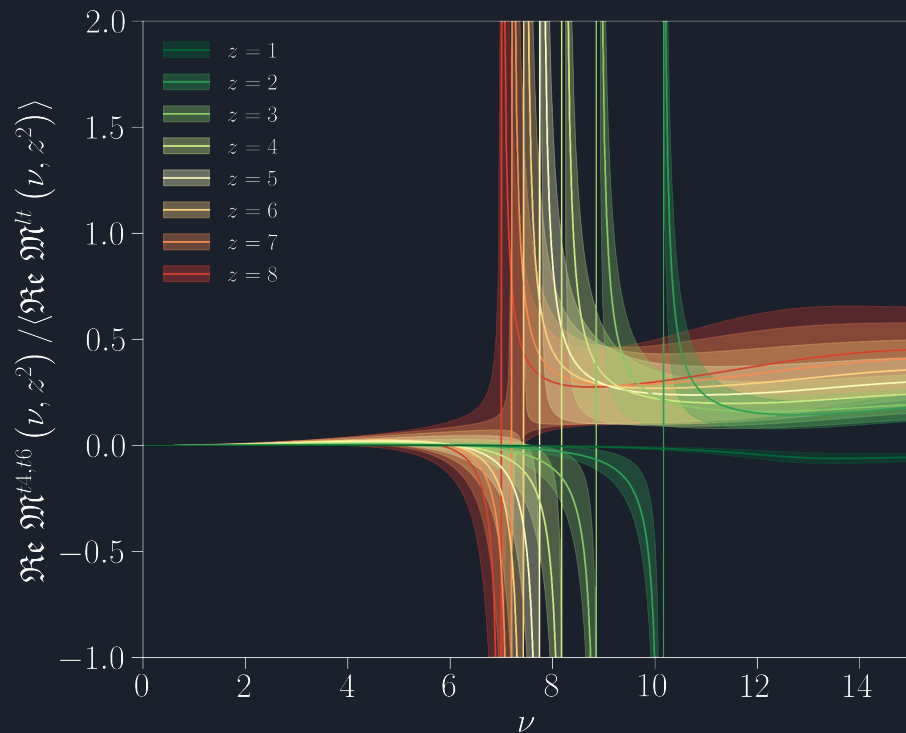
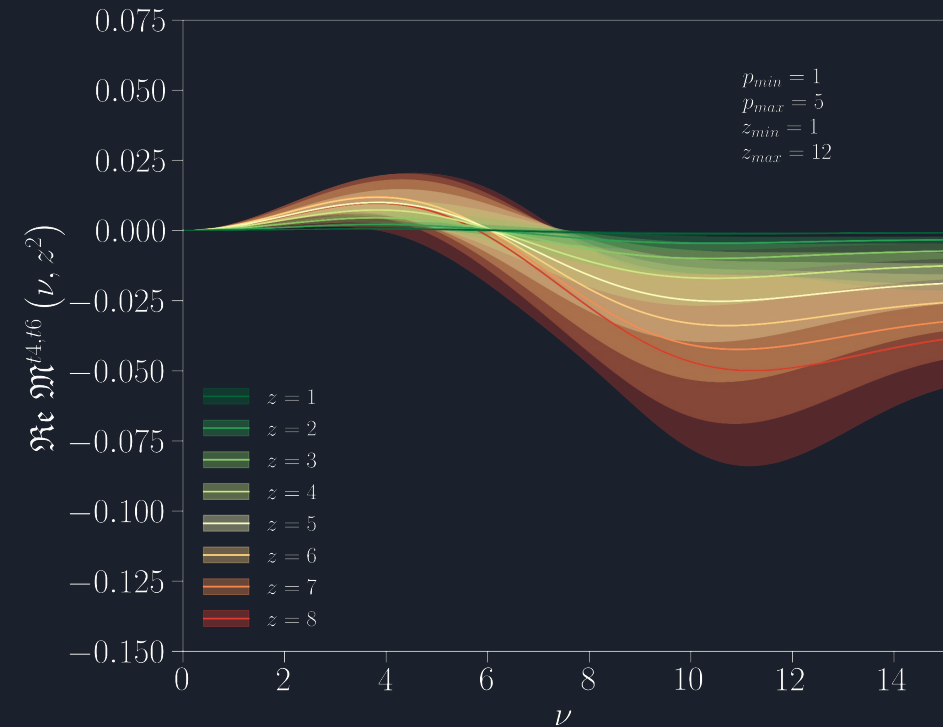
Higher-Twist Variability - Cuts on Momenta



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



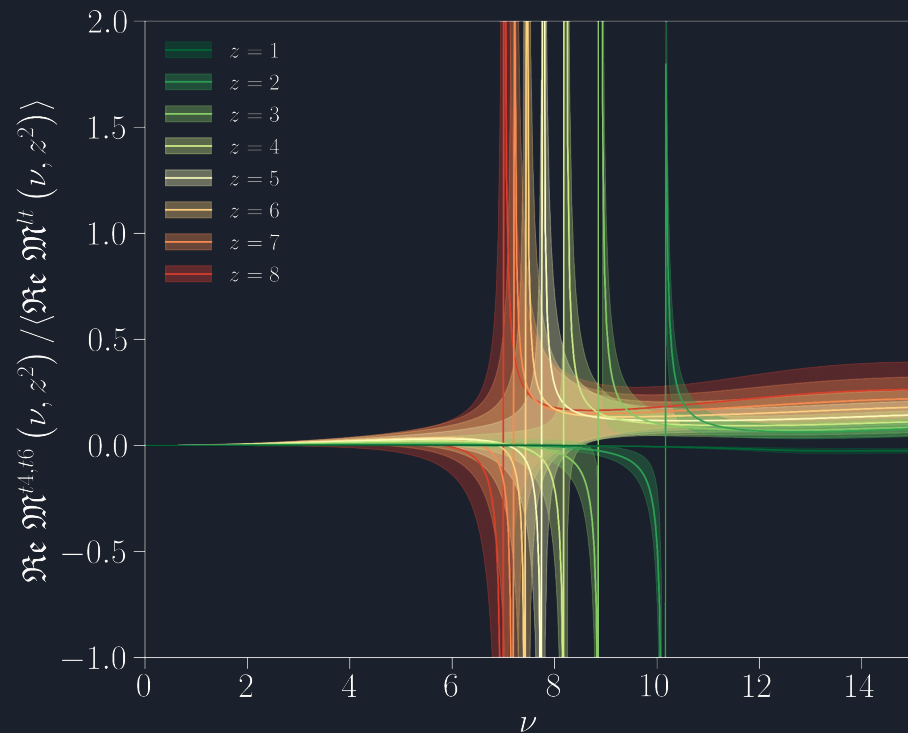
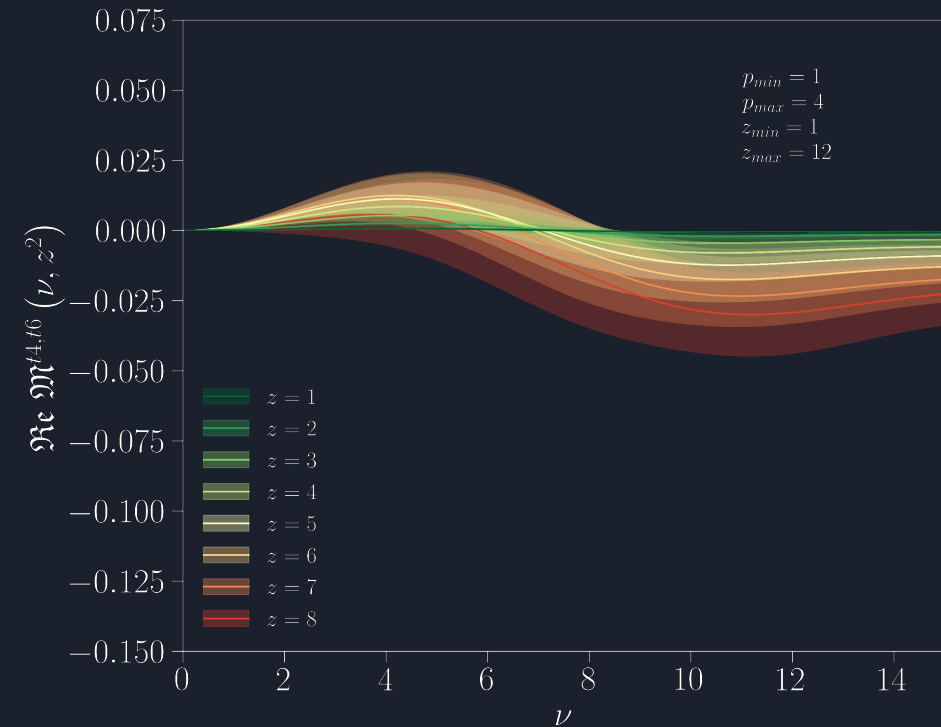
Higher-Twist Variability - Cuts on Momenta



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



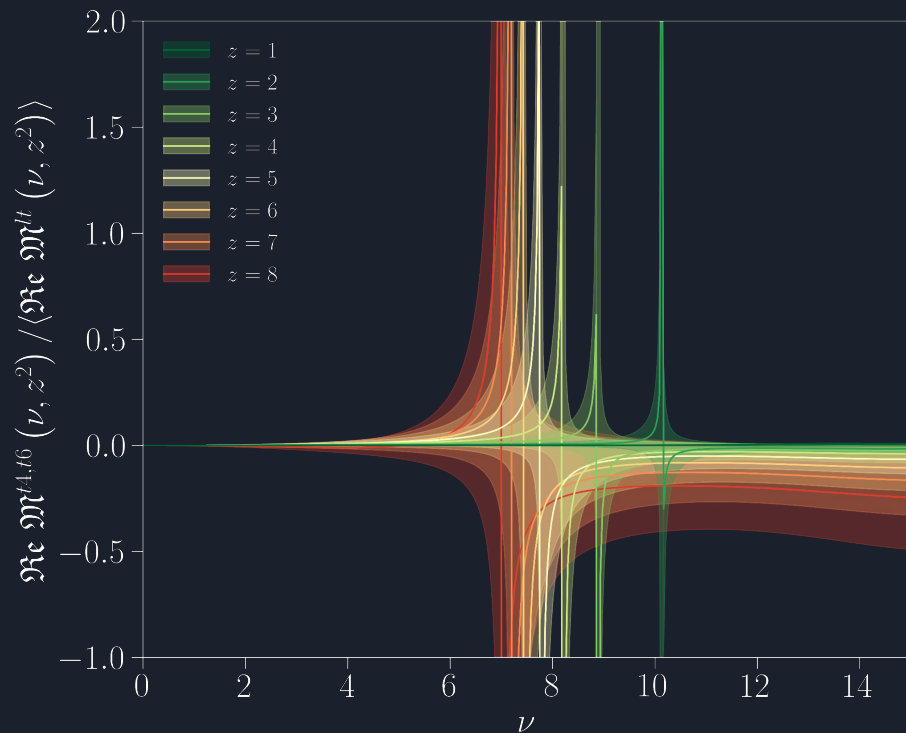
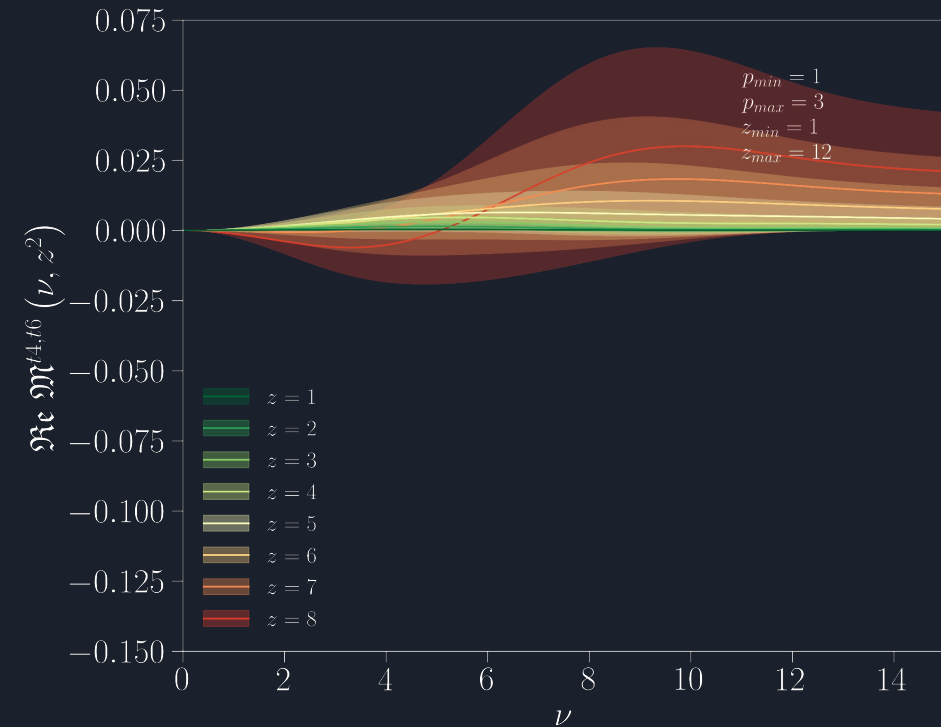
Higher-Twist Variability - Cuts on Momenta



$$\text{Re } \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



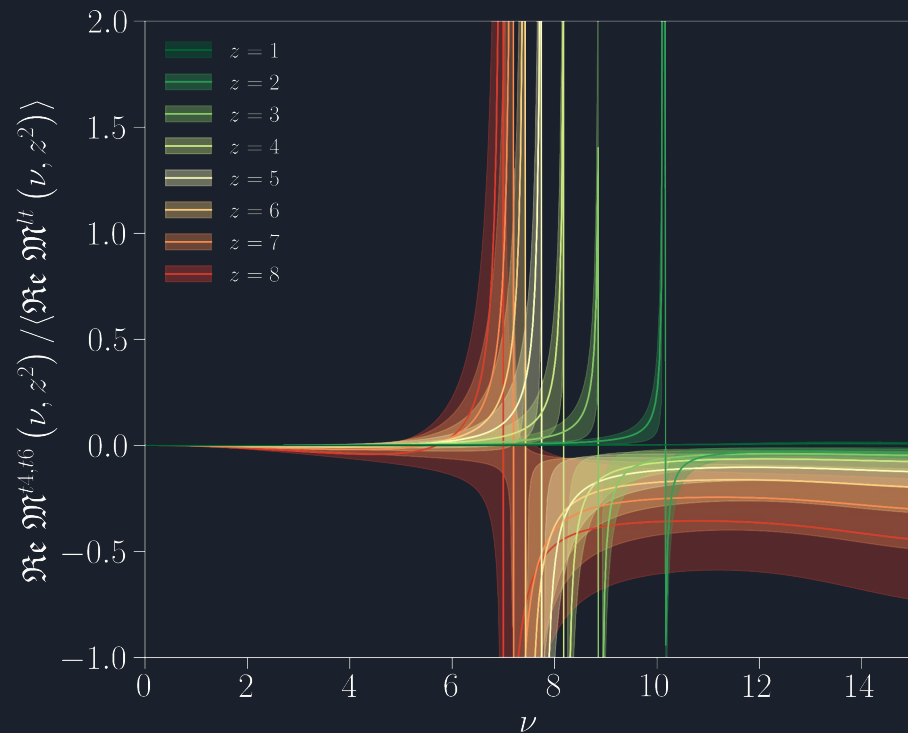
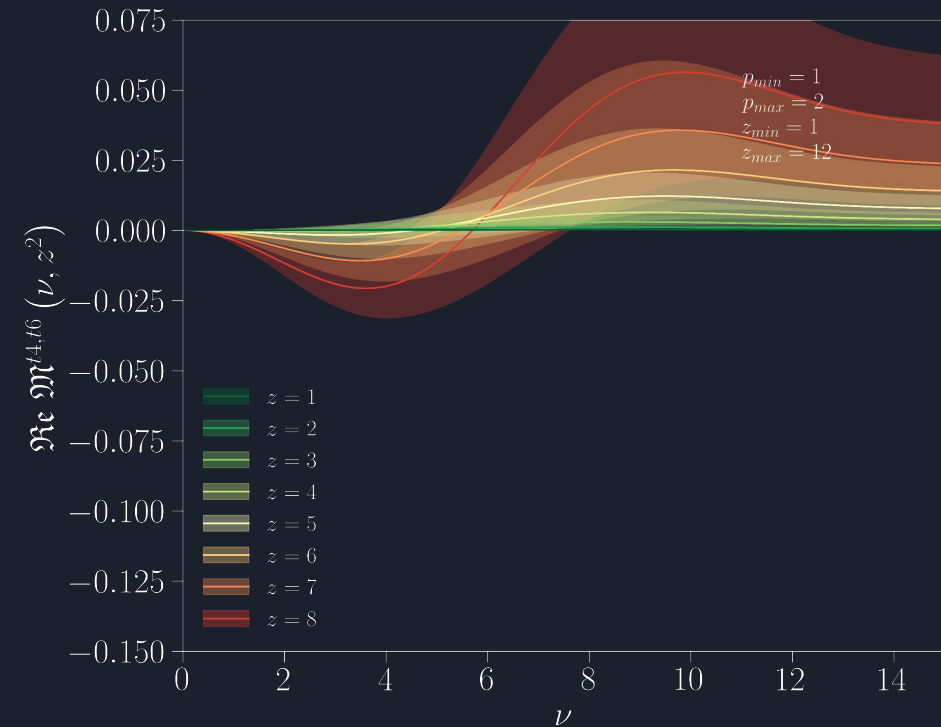
Higher-Twist Variability - Cuts on Momenta



$$\text{Re } \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



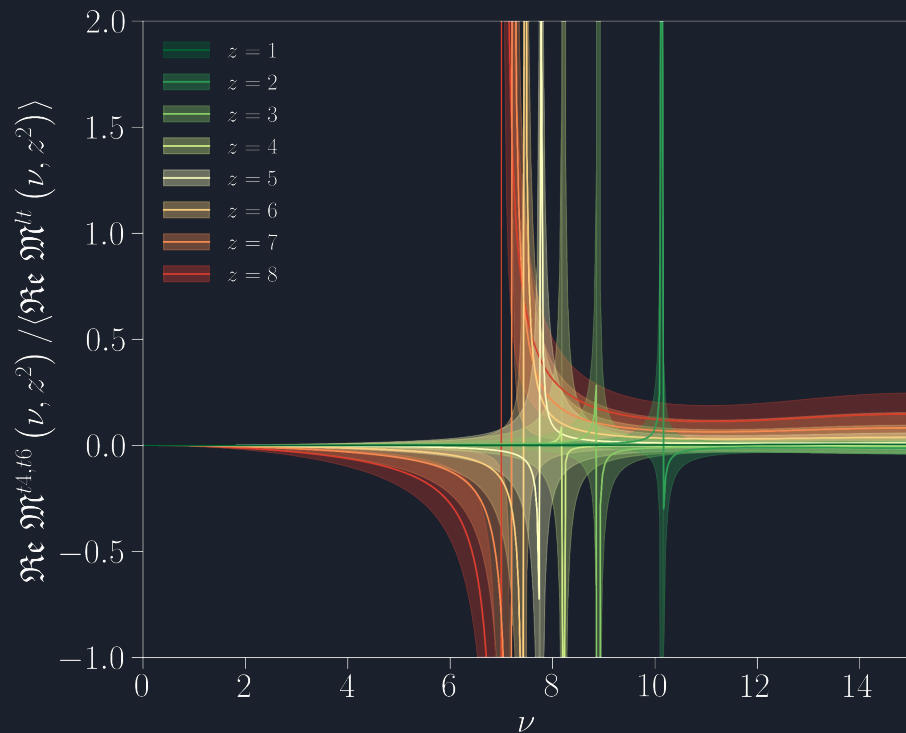
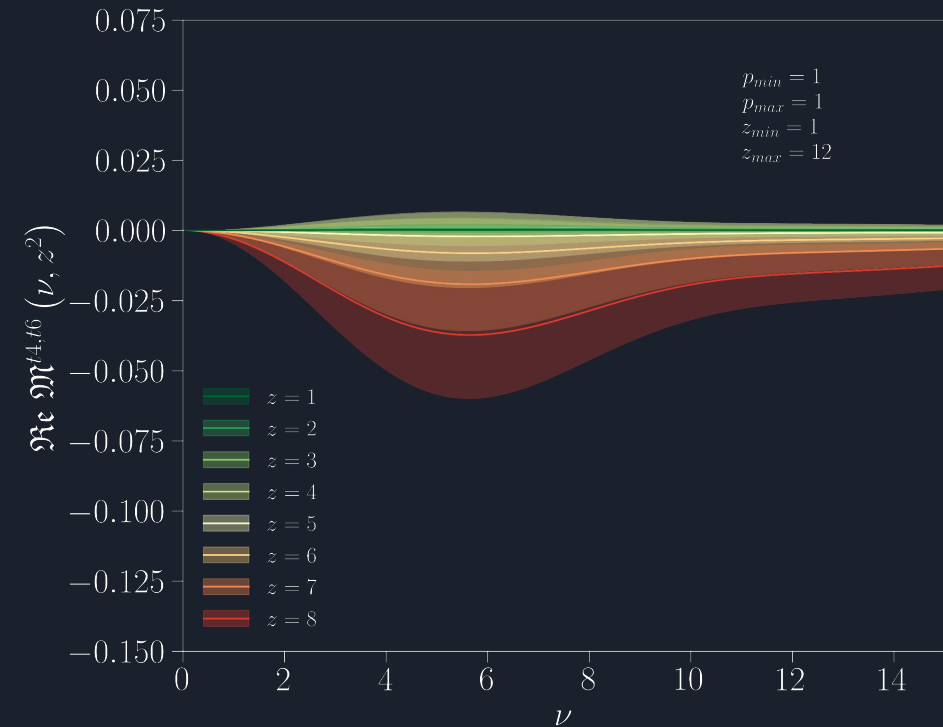
Higher-Twist Variability - Cuts on Momenta



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$



Higher-Twist Variability - Cuts on Momenta



$$\Re \mathcal{M}^{t4,t6}(\nu, z^2) \equiv z^2 \Lambda_{\text{QCD}}^2 \sum_{n=1}^{n_{t4}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t4(\alpha,\beta)} + z^4 \Lambda_{\text{QCD}}^4 \sum_{n=1}^{n_{t6}} \sigma_{0,n}^{(\alpha,\beta)}(\nu) C_{v,n}^{t6(\alpha,\beta)}$$