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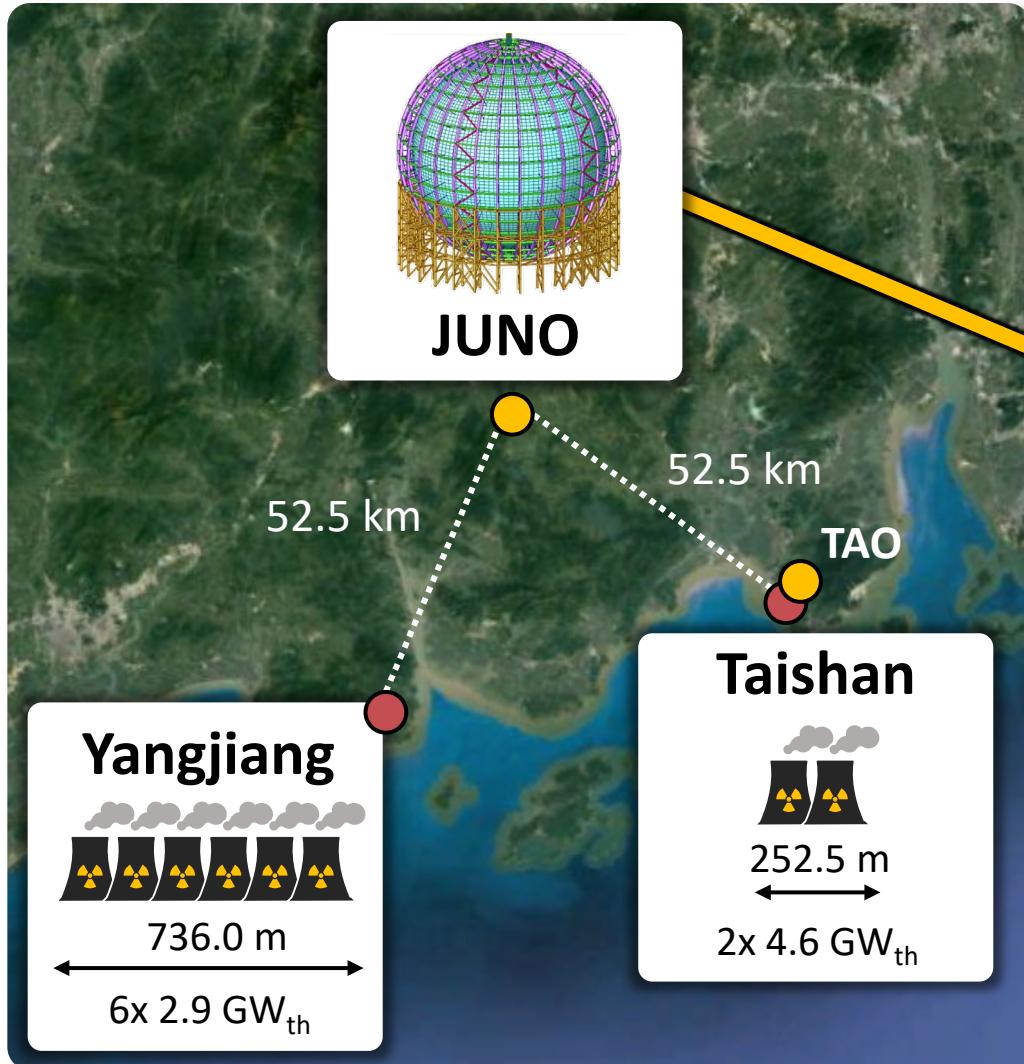
The role of JUNO in leptonic unitarity testing

Inaugurating a new high precision era in the neutrino oscillation field

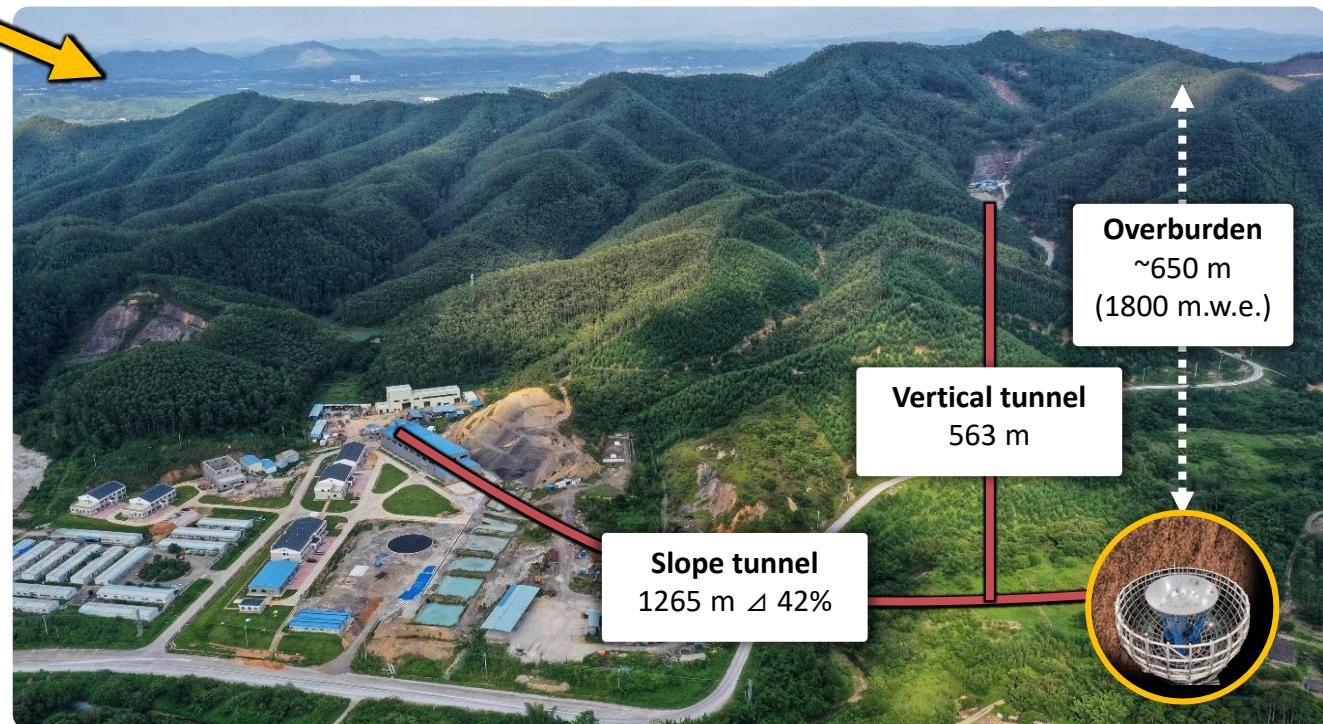
27 June 2022

Andrea Serafini on behalf of the JUNO collaboration
andrea.serafini@infn.it

The Jiangmen Underground Neutrino Observatory



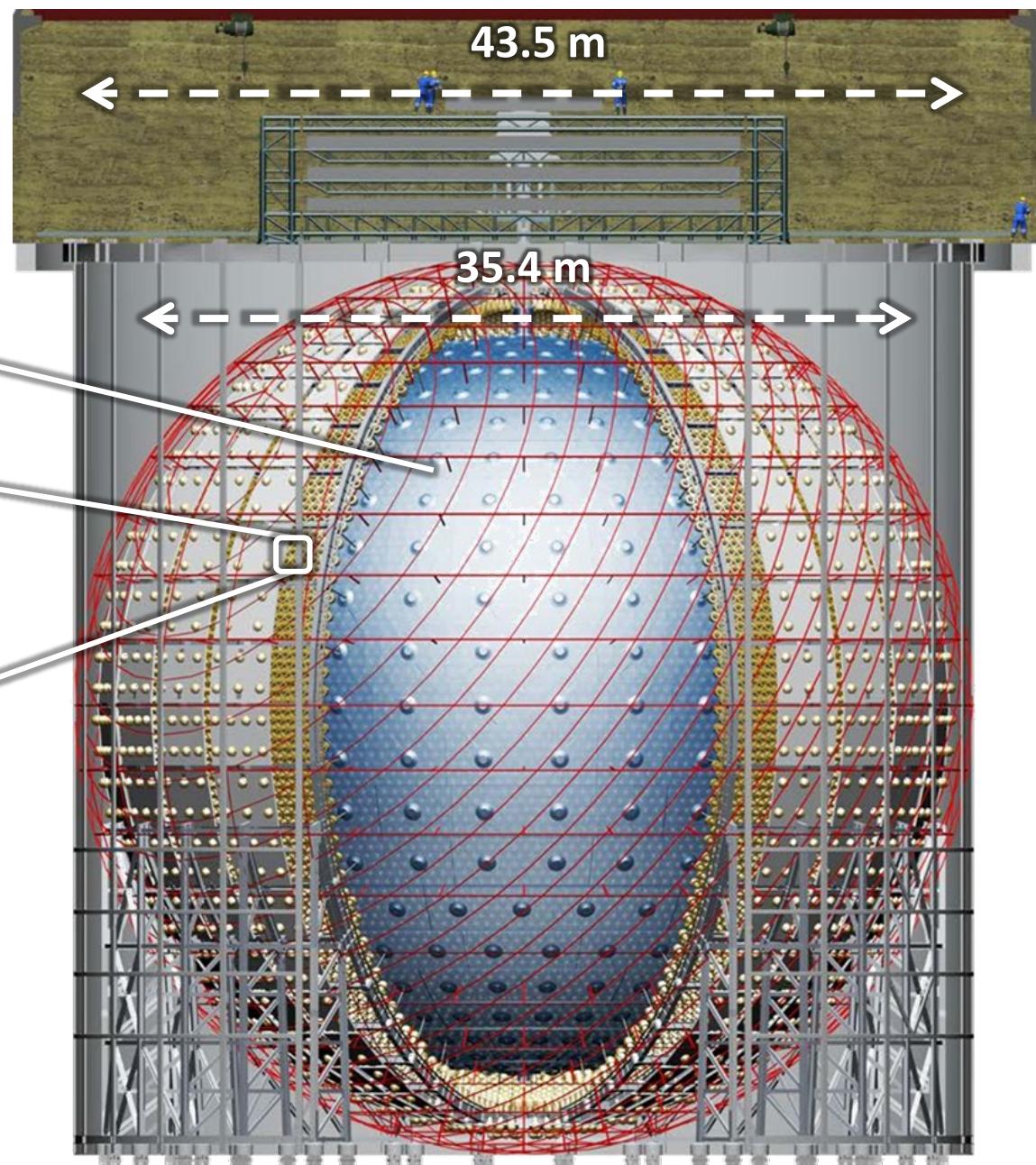
JUNO is a **20 kton** multi-purpose underground **liquid scintillator** detector currently under construction at a baseline of about **52.5 km** from eight **nuclear reactors** in the Guangdong Province of South China.



The JUNO detector

Main requirements:

- **high statistics**
→ 20 kton of liquid scintillator acrylic sphere
- **<3% energy resolution @ 1 MeV**
→ photocoverage > 75%
- **energy-scale systematics below 1%**
→ 17612 20" Large-PMT
→ 25600 3" Small-PMT



	Target mass [kton]	Energy resolution	Light yield [PE/MeV]
Daya Bay	0.02	8%/VE	160
Borexino	0.3	5%/VE	500
KamLAND	1	6%/VE	250
JUNO	20	3%/VE	>1300

The JUNO detection process

JUNO will measure the **antineutrinos** ($\bar{\nu}$) generated in the fissions occurring in 8 nuclear cores at 52.5 km

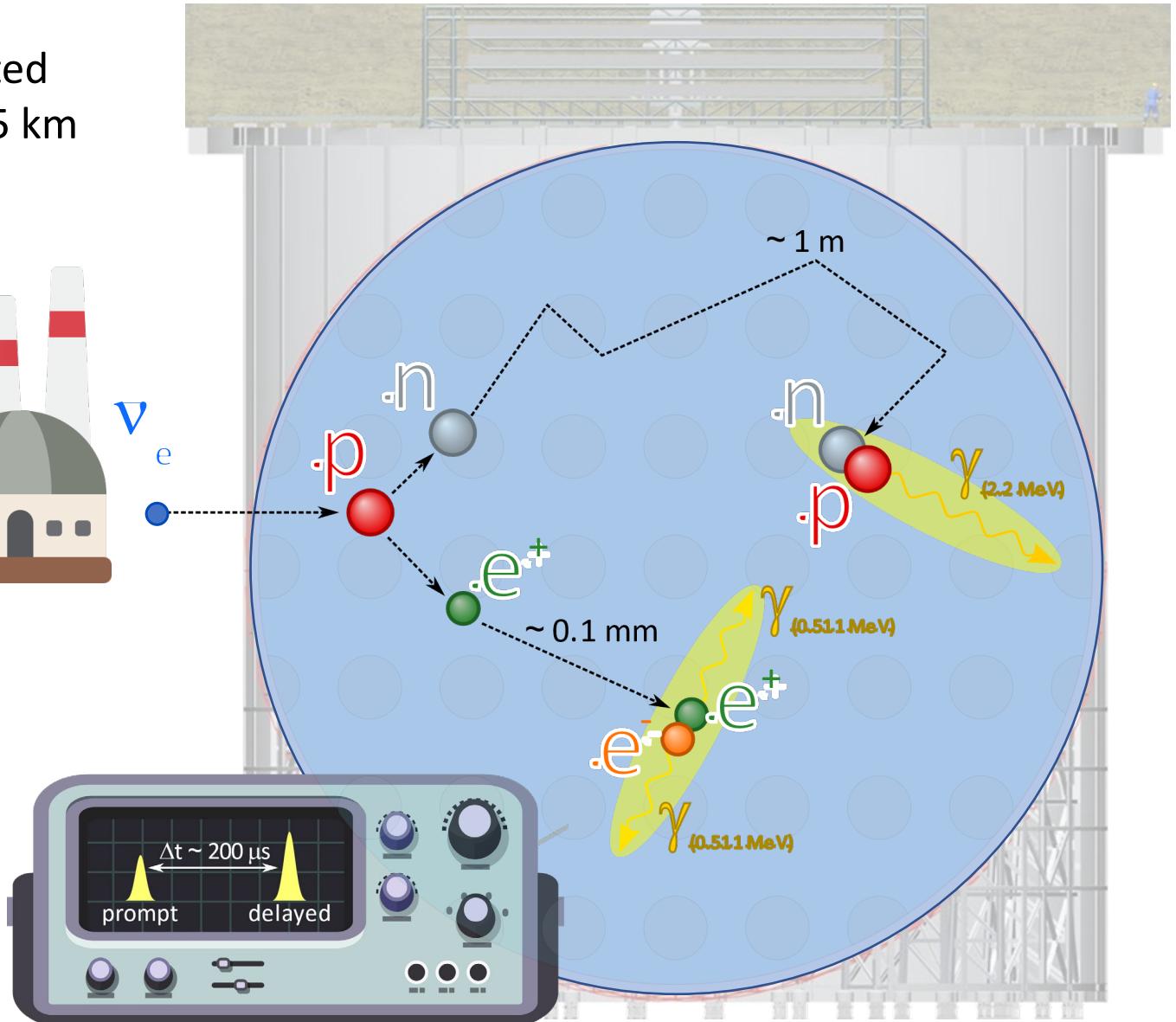
The **detection** is based on a charged current interaction named Inverse Beta Decay (**IBD**) on protons (p)
→ **sensitive only to electron** $\bar{\nu}_e$



Detection relies on a **double coincidence**:

- **prompt** signal: positron (e^+) annihilation
- **delayed** signal: neutron (n) capture

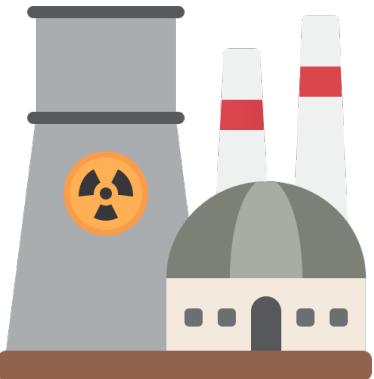
→ **strong handle against most backgrounds**



The JUNO physics program

JUNO can detect neutrinos and antineutrinos coming from several sources:

Reactor



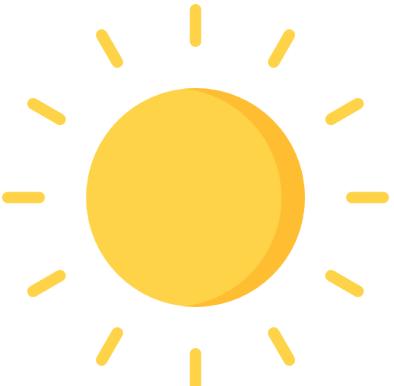
~45/day

Atmosphere



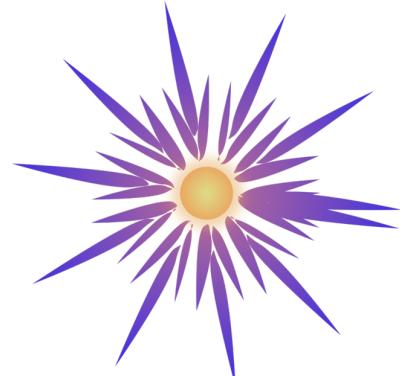
>100/year

Sun



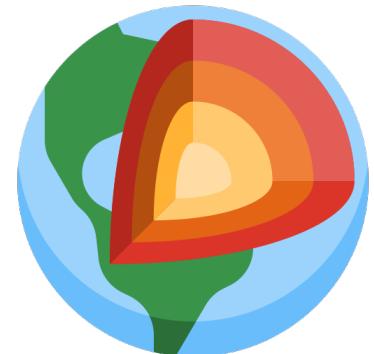
>100/day

Supernovae



$\sim 10^4$ /10 s @ 10kpc

Earth



~400/year

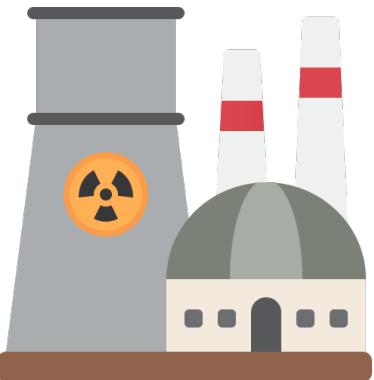
Neutrino oscillation properties

Neutrinos as a probe

The JUNO physics program

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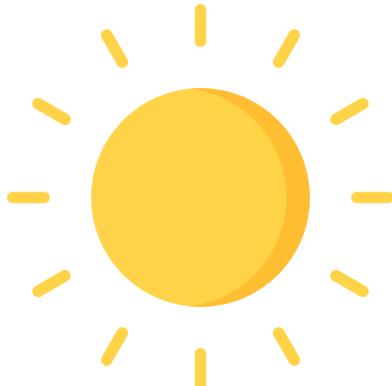
~45/day

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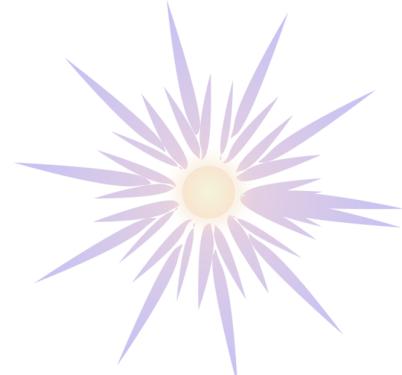
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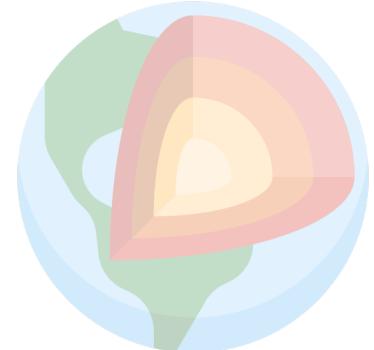
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Neutrino oscillation properties

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A recap of neutrino oscillations

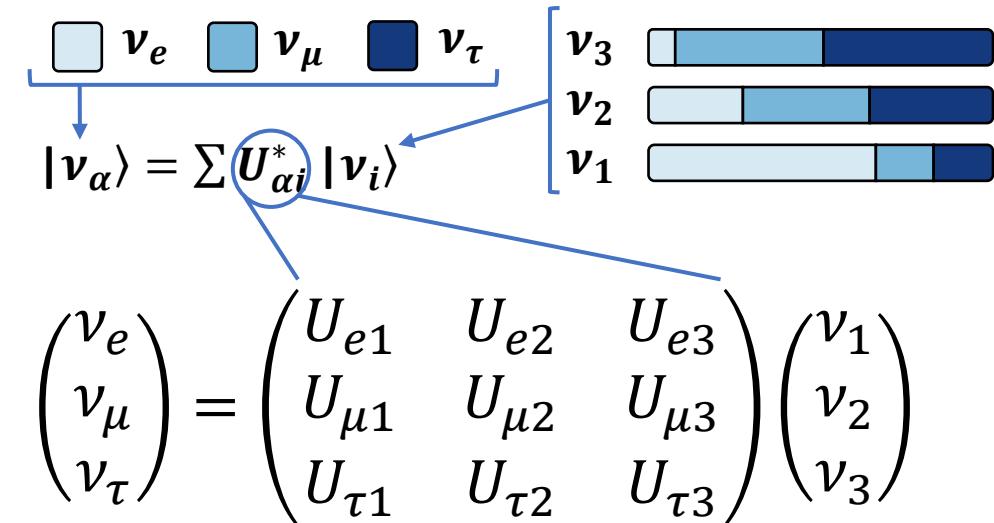
mass eigenstates (ν_i) \neq flavor (ν_α) eigenstates

→ flavor “oscillates” during propagation!

→ phase depends on the splitting between mass states

The matrix connecting mass and flavor states is the Pontecorvo-Maki-Nakagawa-Sakata **(PMNS) mixing matrix.**

→ presumably unitary (?)



A recap of neutrino oscillations

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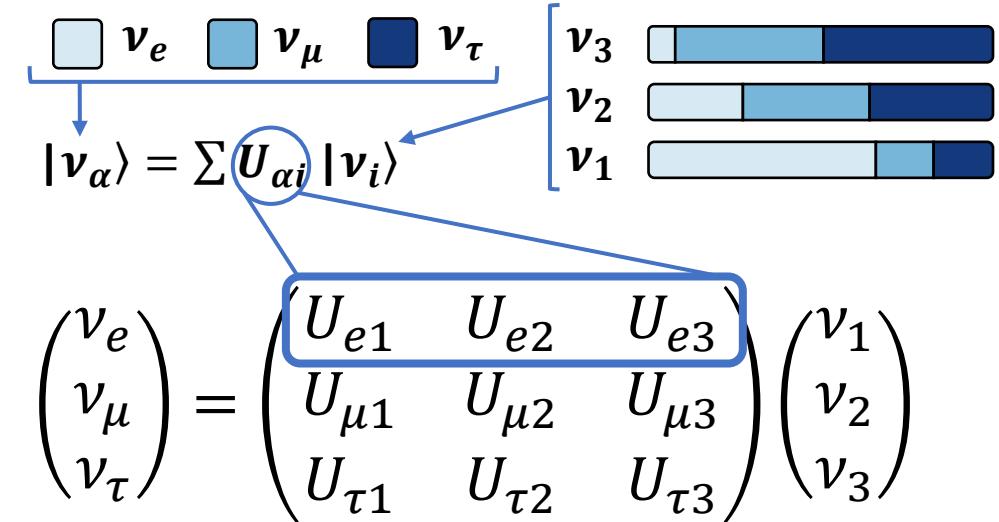
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Unitarity can be **tested directly** via the electron neutrino/antineutrino channel by testing the **Electron Row Unitarity (ERU)**



$$|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2 = 1$$

3 unknowns → 3 constraints required

Testing 3x3 PMNS unitarity

The precision measurement of neutrino mixing parameters is a very powerful tool to test the standard 3-flavor neutrino model via ERU: **3 unknowns → 3 independent constraints required**

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1) Solar → $|U_{e2}|^2 \cdot (|U_{e1}|^2 + |U_{e2}|^2) + |U_{e3}|^4 \quad (\sin^2 \theta_{12} \cos^4 \theta_{13} + \sin^4 \theta_{13}) \quad \sim 5\% \quad [10.5281/zenodo.4134680]$
(SNO + Super-K)

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Indeed, in the Minimal Unitarity Violation (MNU) scheme, the $\bar{\nu}_e$ survival probability:

$$\begin{aligned} P_{\bar{e}\bar{e}}^{MNU} &= (|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2) \cdot P_{\bar{e}\bar{e}} = \\ &= (|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2 \cdot \left(1 - \frac{4|U_{e1}|^2|U_{e2}|^2}{(|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2} \sin^2 \Delta_{21} - \frac{4|U_{e3}|^2(|U_{e1}|^2 \sin^2 \Delta_{31} + |U_{e2}|^2 \sin^2 \Delta_{32})}{(|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2}\right) \end{aligned}$$

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Oscillation physics with JUNO

JUNO aims at determining the neutrino mass ordering @ $>3\sigma$ in 6 years, and will be the first experiment to:

- simultaneously observe fast and slow oscillations (Δm_{21}^2 , θ_{12} , Δm_{31}^2 and θ_{13})
- place <1% precision on Δm_{21}^2 , θ_{12} , Δm_{31}^2

$$\Delta_{ij} = \frac{\Delta m_{ij}^2 L}{4E}$$

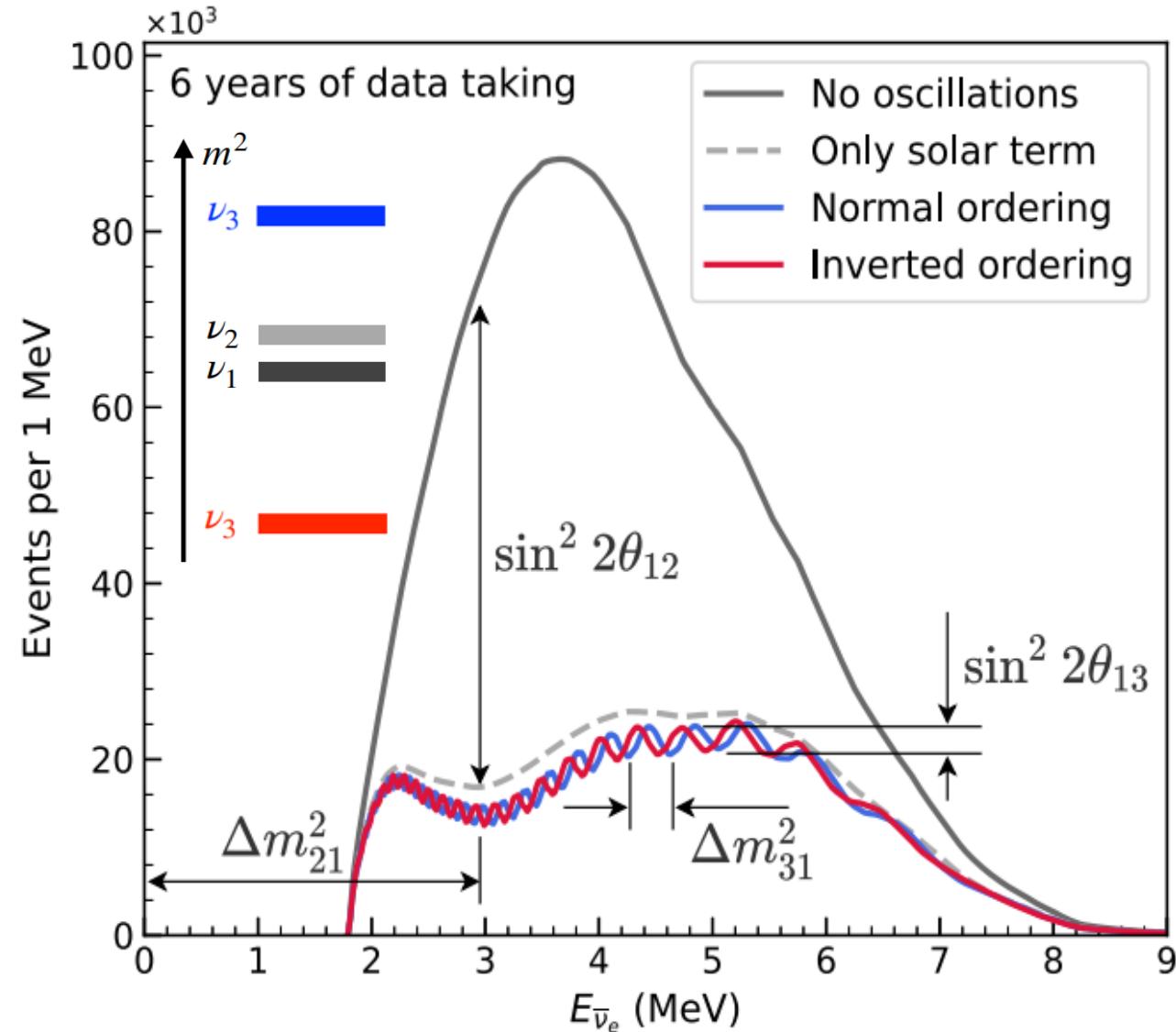
$$P_{\bar{e}\bar{e}} = 1 - P_{21} - P_{31} - P_{32}$$

SLOW $P_{21} = \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \Delta_{21}$

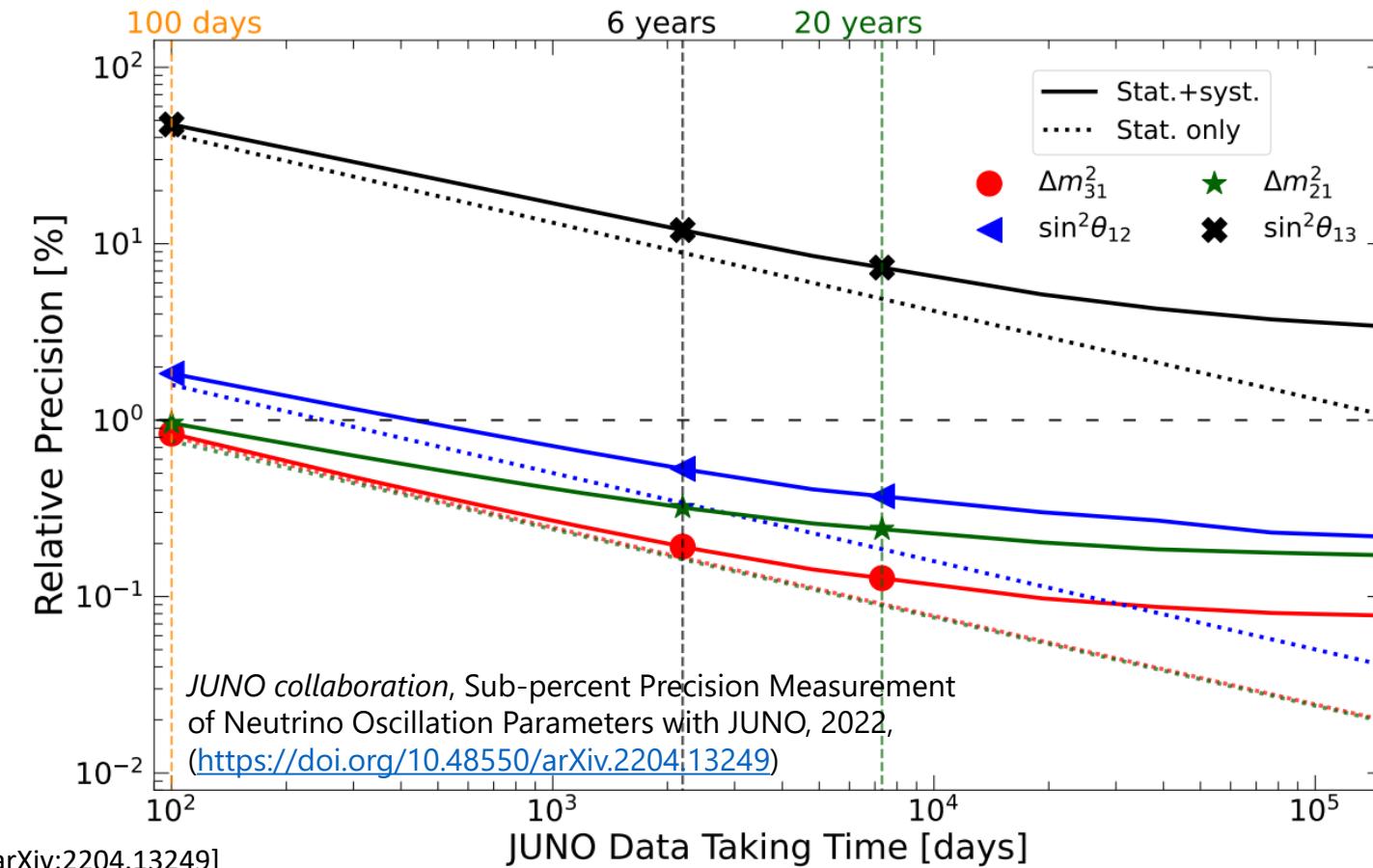
$$P_{31} = \cos^2 \theta_{12} \sin^2 2\theta_{13} \sin^2 \Delta_{31}$$

FAST $P_{32} = \sin^2 \theta_{12} \sin^2 2\theta_{13} \sin^2 \Delta_{32}$

→ probability does not depend on δ_{CP} and θ_{23}



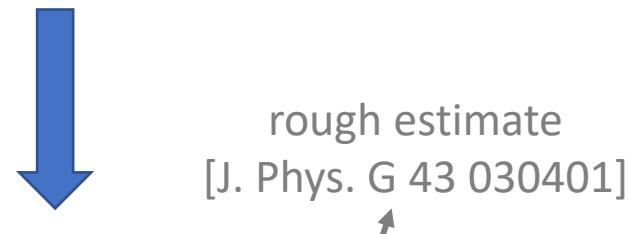
Subpercent measurement of oscillation parameters



	$\sin^2 \theta_{12}$	Δm_{21}^2	$\sin^2 \theta_{13}$	Δm_{31}^2
PDG2020	4.2%	2.4%	3.2%	1.4%
Nufit5.1	4.0%	2.8%	2.8%	1.1%
JUNO 6 years	0.5%	0.3%	12%	0.2%

In <2 years, JUNO will improve the precision on Δm_{21}^2 , θ_{12} , Δm_{31}^2
 → unprecedented <1% level.

In 6 years, the precision on θ_{12} , Δm_{21}^2 and Δm_{31}^2 will reach 0.5%, 0.3% and 0.2%



- constrain PMNS unitarity @ 2.5% level (limited by SNO solar constraint)
- narrow down the parameter space of $0\nu\beta\beta$ effective mass
- constrain the neutrino mass sum rule $\Delta m_{13}^2 + \Delta m_{21}^2 + \Delta m_{32}^2 = 0$

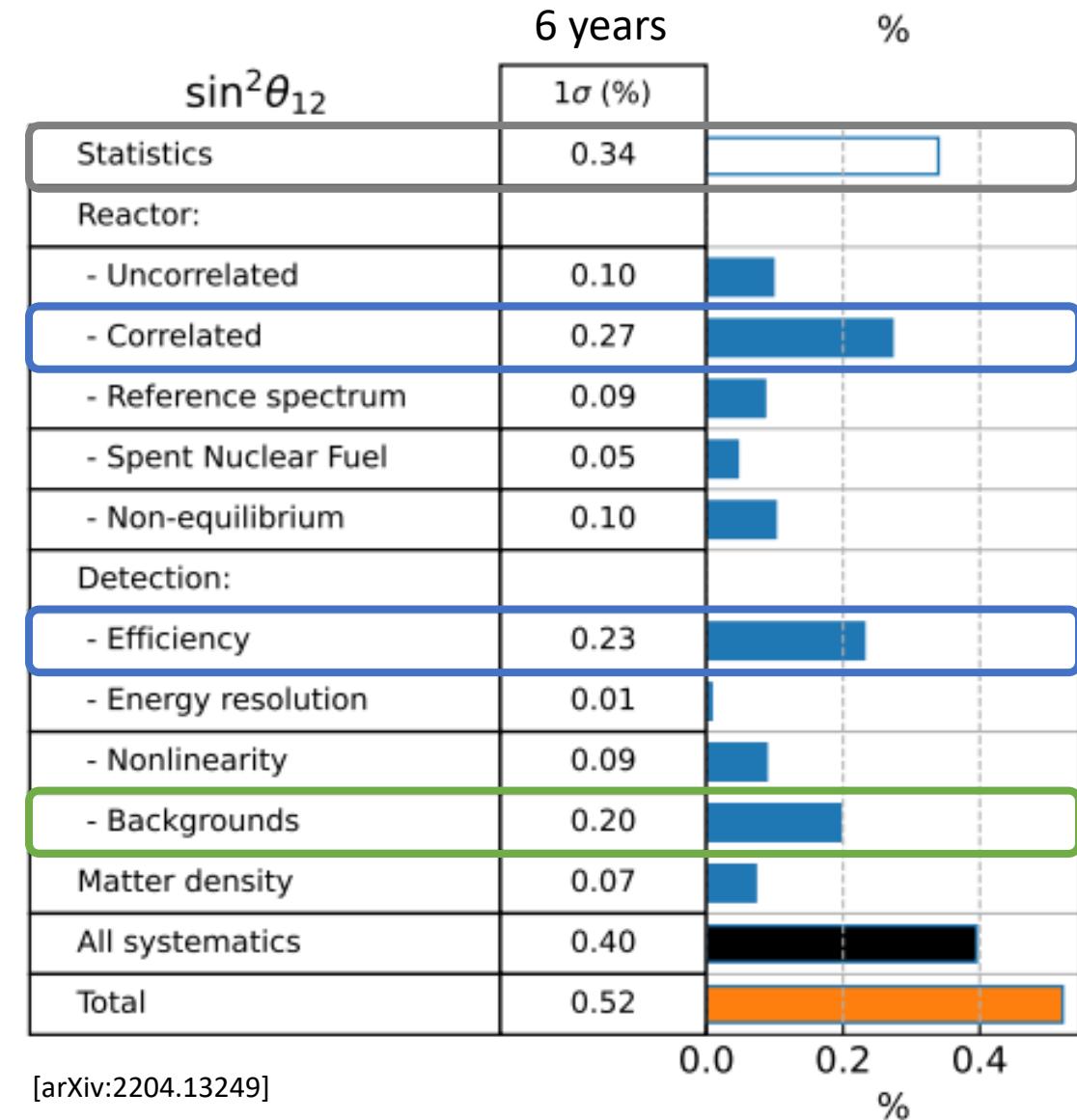
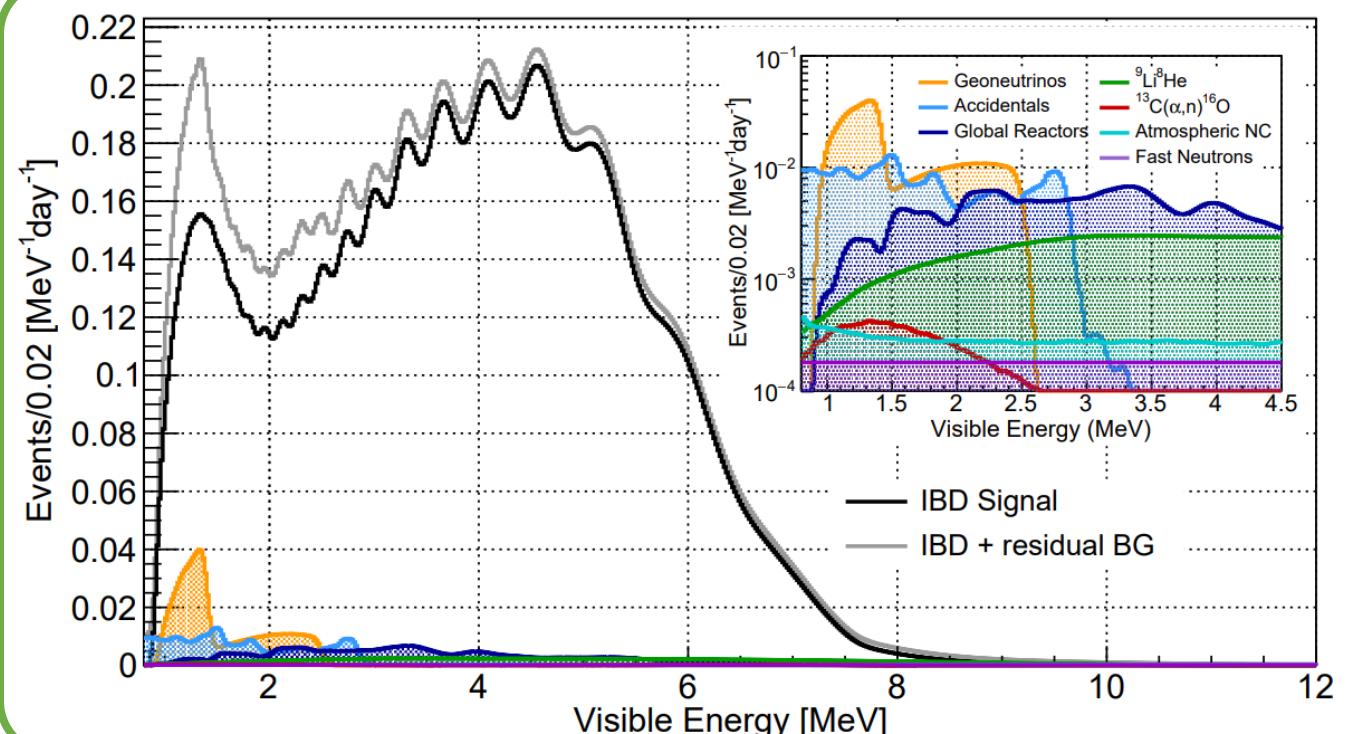
Challenges to θ_{12} measurement

Main systematic uncertainties:

→ rate/normalization:

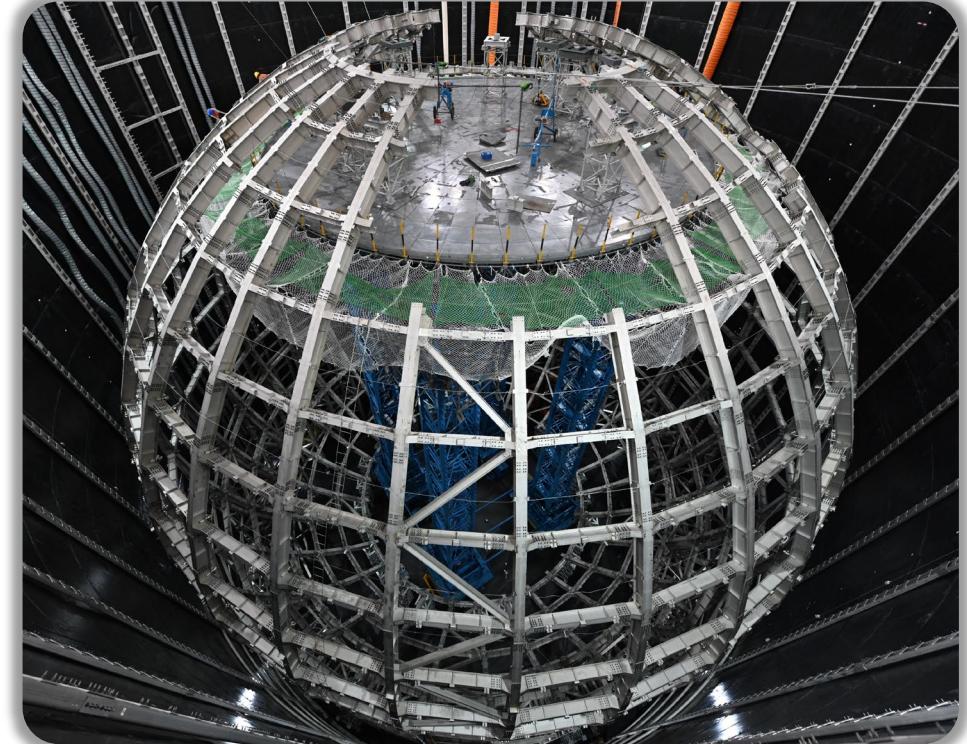
- reactor flux normalization
- detector efficiency

→ backgrounds

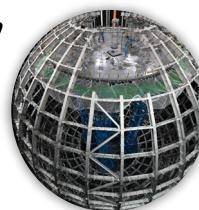


Final remarks

- **JUNO** will inaugurate a **high precision** era in the neutrino oscillation field. In 6 years:
 - **neutrino mass hierarchy** at $>3\sigma$
 - **<1% precision** on θ_{12} , Δm_{21}^2 and Δm_{31}^2
- The 0.5% precision on θ_{12} will provide a constrain to **directly test leptonic unitarity** via ERU at 2.5% (together with solar and reactor constraints)



Thank you !



Back up

PMNS terms

Table 1. Quantities to which each experiment is sensitive: using the PMNS parameterization when unitarity is assumed (center column), using the MP parametrization when unitarity is not assumed (right column).

Experiment	PMNS Quantity	LMM Quantity
Solar Neutral Current	1	$(U_{e1} ^2 + U_{e2} ^2)N_2^2 + U_{e3} ^2 N_3^2$
Solar Charged Current	$\sin^2 \theta_{12} \cos^4 \theta_{13} + \sin^4 \theta_{13}$	$ U_{e2} ^2 (U_{e1} ^2 + U_{e2} ^2) + U_{e3} ^4$
KamLAND	$\cos^4 \theta_{13} \sin^2 (2\theta_{12})$	$4 U_{e1} ^2 U_{e2} ^2$
Daya Bay	$\sin^2 (2\theta_{13})$	$4 U_{e3} ^2 (U_{e1} ^2 + U_{e2} ^2)/N_e^2$
Sterile Neutrino $P_{\alpha\beta}$ ($\alpha \neq \beta$)	0	$ t_{\alpha\beta} ^2$
OPERA	$\cos^4 \theta_{13} \sin^2 (2\theta_{23})$	$4 U_{\mu 3} ^2 U_{\tau 3} ^2 /N_\mu^2$
Long-baseline $P_{\mu e}$ (T2K, NOvA, DUNE, T2HK)	$\sin^2 \theta_{23} \sin^2 (2\theta_{13})$	$4 U_{e3} ^2 U_{\mu 3} ^2 /N_\mu^2$
Long-baseline $P_{\mu\mu}$ (T2K, NOvA, DUNE, T2HK)	$4 \cos^2 \theta_{13} \sin^2 \theta_{23} (1 - \cos^2 \theta_{13} \sin^2 \theta_{23})$	$4 U_{\mu 3} ^2 (U_{\mu 1} ^2 + U_{\mu 2} ^2)/N_\mu^2$

State of the art and prospects

	knowledge as of 2020		expected knowledge beyond 2020		
	dominant	precision (%)	precision (%)	dominant	technique
θ_{12}	SNO	2.3	≤ 1.0	JUNO	<i>reactor</i>
θ_{23}	NOvA	2.0	~ 1.0	DUNE+HK	<i>beam</i>
θ_{13}	DYB	3.3	3.3	DC+DYB+RENO	<i>reactor</i>
δm^2	KL	2.3	≤ 1.0	JUNO	<i>reactor</i>
$ \Delta m^2 $	DYB+T2K	1.3	≤ 1.0	JUNO+DUNE+HK	<i>reactor+beam</i>
$\pm \Delta m^2$	SK	unknown	measured??	JUNO+DUNE+HK	<i>reactor+beam</i>
δ_{CP}	T2K	unknown	measured??	DUNE+HK	<i>beam</i>

A recap of neutrino oscillations

For neutrinos, mass (ν_i) and flavor (ν_α) eigenstates do not correspond.

The Pontecorvo-Maki-Nakagawa-Sakata (PMNS) mixing matrix can be parametrized by:

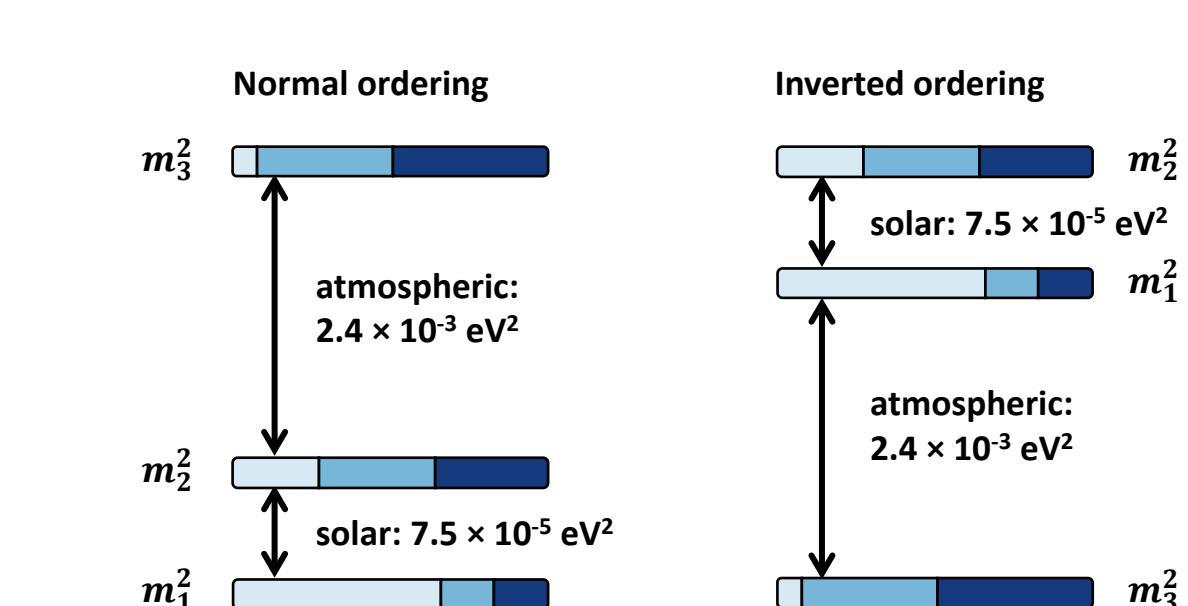
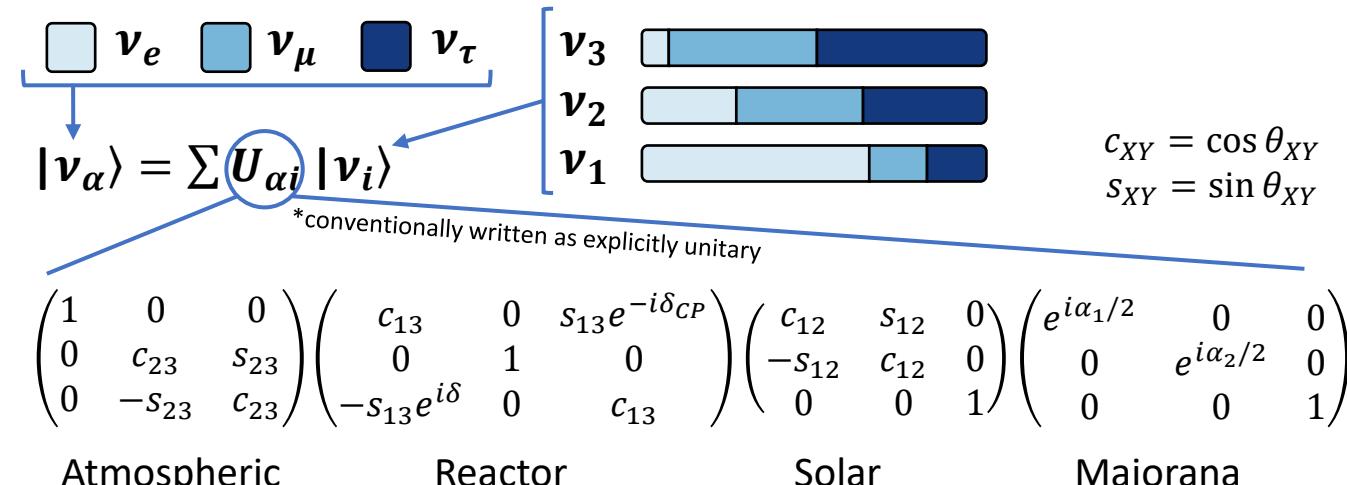
- 3 mixing angles ($\theta_{12}, \theta_{13}, \theta_{23}$)
- 1 CP-violating phase (δ_{CP})

The non-correspondence between mass and flavor eigenstates causes the flavor to “oscillate” during propagation.

The phase of this oscillation depends on the splitting between the mass states:

- Δm_{21}^2
- Δm_{31}^2

With $\Delta m_{21}^2 > 0$ and Δm_{31}^2 and Δm_{32}^2 possibly positive (normal ordering) or negative (inverted).



Testing 3x3 PMNS unitarity

The precision measurement of neutrino mixing parameters is a very powerful tool to test the standard 3-flavor neutrino model. In particular, PMNS unitarity can be tested directly via **Electron Row Unitarity (ERU)**:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad \rightarrow \quad |U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2 = 1$$

3 unknowns → 3 constraints required

In the Minimal Unitarity Violation (MNU) scheme:

$$\begin{aligned} P_{\bar{e}\bar{e}}^{MNU} &= (|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2) \cdot P_{\bar{e}\bar{e}} = \\ &= (|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2 \cdot \left(1 - \frac{4|U_{e1}|^2|U_{e2}|^2}{(|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2} \sin^2 \Delta_{21} - \frac{4|U_{e3}|^2(|U_{e1}|^2 \sin^2 \Delta_{31} + |U_{e2}|^2 \sin^2 \Delta_{32})}{(|U_{e1}|^2 + |U_{e2}|^2 + |U_{e3}|^2)^2} \right) \end{aligned}$$

Additional constraint from solar neutrino experiments (SNO):

$$|U_{e2}|^2 \cdot (|U_{e1}|^2 + |U_{e2}|^2) + |U_{e3}|^4$$

Constrained by SNO at $\sim 12\%$ [Phys. Rev. C 72, 055502]

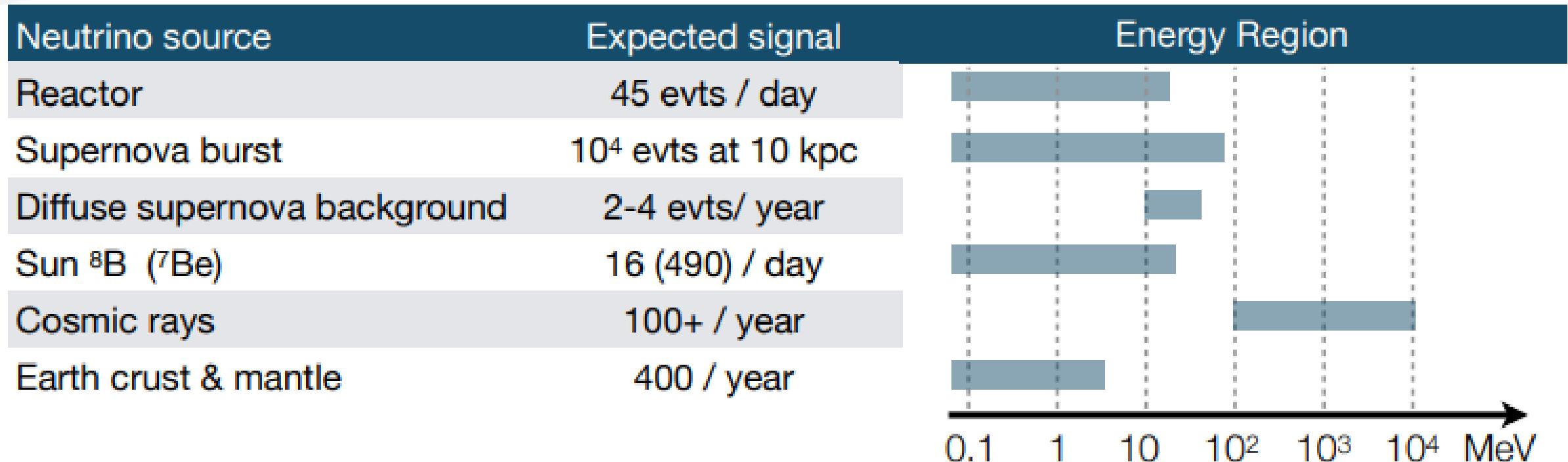
$$\sim \sin^2 2\theta_{12} \cdot \sin^2 \Delta_{21}$$

Measurable by JUNO
with <1% precision

$$\sim \sin^2 2\theta_{13} \cdot \sin^2 \Delta_{3(1,2)}$$

Measurable by JUNO @ $\sim 10\%$ and
constrained by Daya Bay @ <5%
[Phys. Rev. D 93 7, 072011]

The JUNO physics program



The JUNO detector

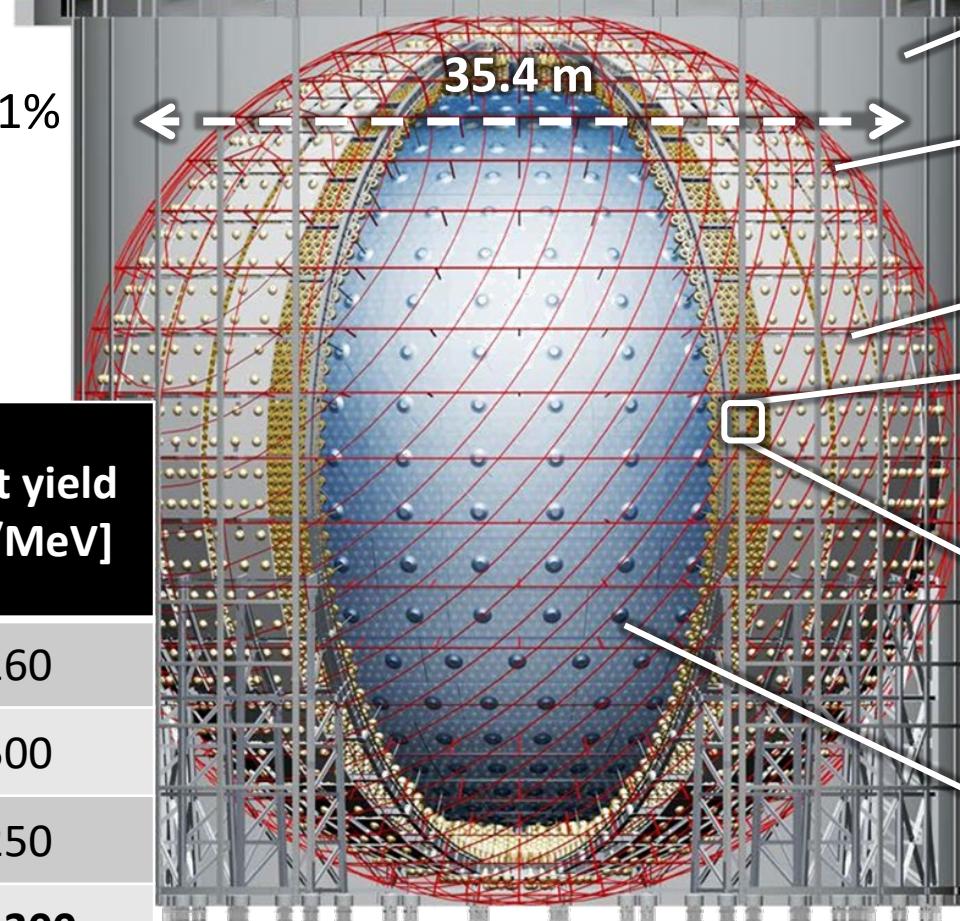
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- energy-scale systematics below 1%
- high statistics

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Top tracker and calibration house



Water pool

Earth magnetic field compensation coils

Veto Large-PMTs



PMTs:
17612 20" Large-PMT
25600 3" Small-PMT
photocoverage > 75%

Acrylic sphere with 20 kton of liquid scintillator (linear alkylbenzene)