

Non-local S-matrix in a solvable model

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28 June 2022

SUSY 2022,
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Upcoming article:
DK, A. Pilaftsis

1

Introduction

- Why non-local S-matrix?
- local S-matrix
- Introducing non-locality
- Simplifications
- Physical representation

2

Analytical results

- First analytical result
- Some observations
- Further simplification: Recover conservation of momentum

3

Features

- Far-field— $|\vec{l}| \gg |\vec{k}| \delta l^2$
- Near-field— $|\vec{l}| \ll |\vec{k}| \delta l^2$
- Pattern in space

4

Makes sense check-list

5

Summary—Future

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Why non-local S-matrix?

Non-locality is connected to finite volume effects.

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- Displaced vertex searches. Deeper understanding of how a mediator travels.
- Other possibilities: QFT picture of diffraction, effective interactions (e.g. non-local chiral ET³ or jets), etc.

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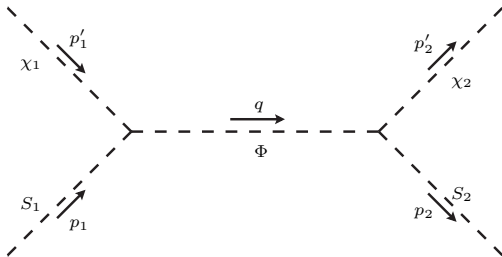
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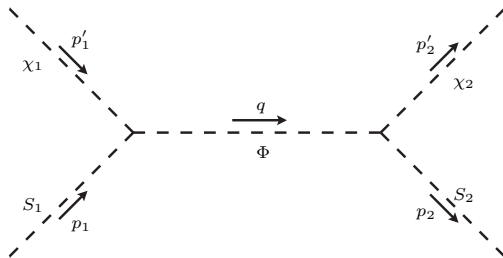
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$$T(p, k) = \lambda g \int \frac{d^4 q}{(2\pi)^4} \frac{1}{q^2 - m_\Phi^2 + i\epsilon} \overbrace{\int d^4 x d^4 y e^{-i(p-q)\cdot x} e^{i(k-q)\cdot y}}^{\delta^{(4)}(p-q)\delta^{(4)}(k-q)} ,$$

with $p = p_1 + p'_1$, $k = p_2 + p'_2$

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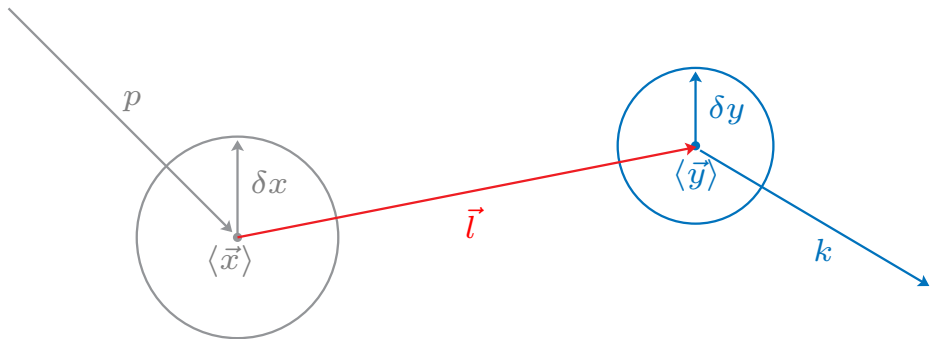
We can simplify the calculations by assuming *steady state* and *spherical symmetry*; $\delta p^0, \delta k^0 \rightarrow 0$, $\delta p^i = 1/\delta x$, and $\delta k^i = 1/\delta y$.

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$$\begin{aligned} T(p, k; \vec{l}, \delta x, \delta y) = & (2\pi) \delta(p^0 - k^0) \lambda g \\ & \int d^3 \vec{x} e^{i \vec{p} \cdot \vec{x} - \vec{x}^2 / \delta x^2} \\ & \int d^3 \vec{y} e^{-i \vec{k} \cdot \vec{y} - \vec{y}^2 / \delta y^2} \\ & \int \frac{d^3 \vec{q}}{(2\pi)^3} \frac{e^{-i \vec{q} \cdot (\vec{x} - \vec{y} - \vec{l})}}{\tilde{q}^2 - |\vec{q}|^2 + i\epsilon}, \end{aligned}$$

with $\vec{l} = \langle \vec{y} \rangle - \langle \vec{x} \rangle$ and $\tilde{q}^2 = p^{02} - m_\Phi^2$.

Physical representation



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Summary-Future

After some algebra, the amplitude becomes

$$T(p, k; \vec{l}, \delta x, \delta y) \sim \delta(p^0 - k^0) \frac{\delta x^3 \delta y^3}{|\vec{L}|} e^{-[(|\vec{p}|^2 + \tilde{q}^2)\delta x^2 + (|\vec{k}|^2 + \tilde{q}^2)\delta y^2]/4} \\ \left[e^{i\tilde{q}|\vec{L}|} \text{Erfc}(z_-) - e^{-i\tilde{q}|\vec{L}|} \text{Erfc}(z_+) \right],$$

where $z_{\pm} = -\frac{i}{2}\tilde{q}\delta l \pm \frac{|\vec{L}|}{\delta l}$, $\vec{L} = \vec{l} - \frac{i}{2}(\vec{p}\delta x^2 + \vec{k}\delta y^2)$,
 $\delta l^2 = \delta x^2 + \delta y^2$, and $|\vec{L}| = \sqrt{\vec{L} \cdot \vec{L}}$.

Some observations

From this form we observe that:

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From this form we observe that:

- Amplitude is finite in the physical region.
- Momentum is not conserved, and the distribution of \vec{k} depends on $\delta x, \delta y$.
- In the limit $\delta x, \delta y \rightarrow 0$, $T \sim \delta x^3 \delta y^3 e^{i\tilde{q}|\vec{l}|}/|\vec{l}| \delta(p^0 - k^0)$.⁴

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Further simplification: Recover conservation of momentum

In an experiment, we expect negligible violation of momentum conservation ($|\vec{p}|\delta x, |\vec{k}|\delta y \gg 1$). So, we introduce a profile that can help us work out the general behaviour of the amplitude:

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Note:

Momenta at each vertex suffer from uncertainties.

Introduce matrix element as usual

$$M \sim \frac{\delta l^3}{|\vec{l}|} e^{-\delta l^2(|\vec{k}|^2 + \tilde{q}^2)/2} \left[e^{i\tilde{q}|\vec{l}|} \text{Erfc}(z_-) - e^{-i\tilde{q}|\vec{l}|} \text{Erfc}(z_+) \right].$$

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In the far-field (*Fraunhofer*) region we recover the inverse-square law (similar to $\delta l \rightarrow 0$). The matrix element becomes:

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Observations:

- Inverse square law.
- Suppressed backwards direction.
- Finite.
- Oscillations of mixed mediators.
- Off-shell mediator may be the most probable outcome.

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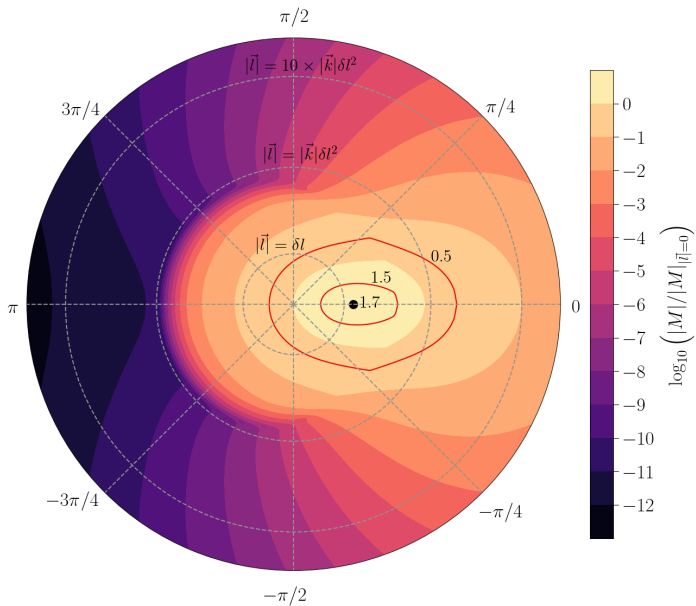
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Observations:

- $\delta l \rightarrow \infty \Rightarrow M \sim i\pi \delta_+ (|\vec{k}|^2 - \tilde{q}^2) + \mathcal{P} \left\{ \frac{1}{|\vec{k}|^2 - \tilde{q}^2} \right\}.$
- Finite.
- Oscillations of mixed mediators.
- Slightly shifted maximum.

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- ✓ Distinct features, *i.e.* falsifiable.

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- **Experiments.**

Thank you!