



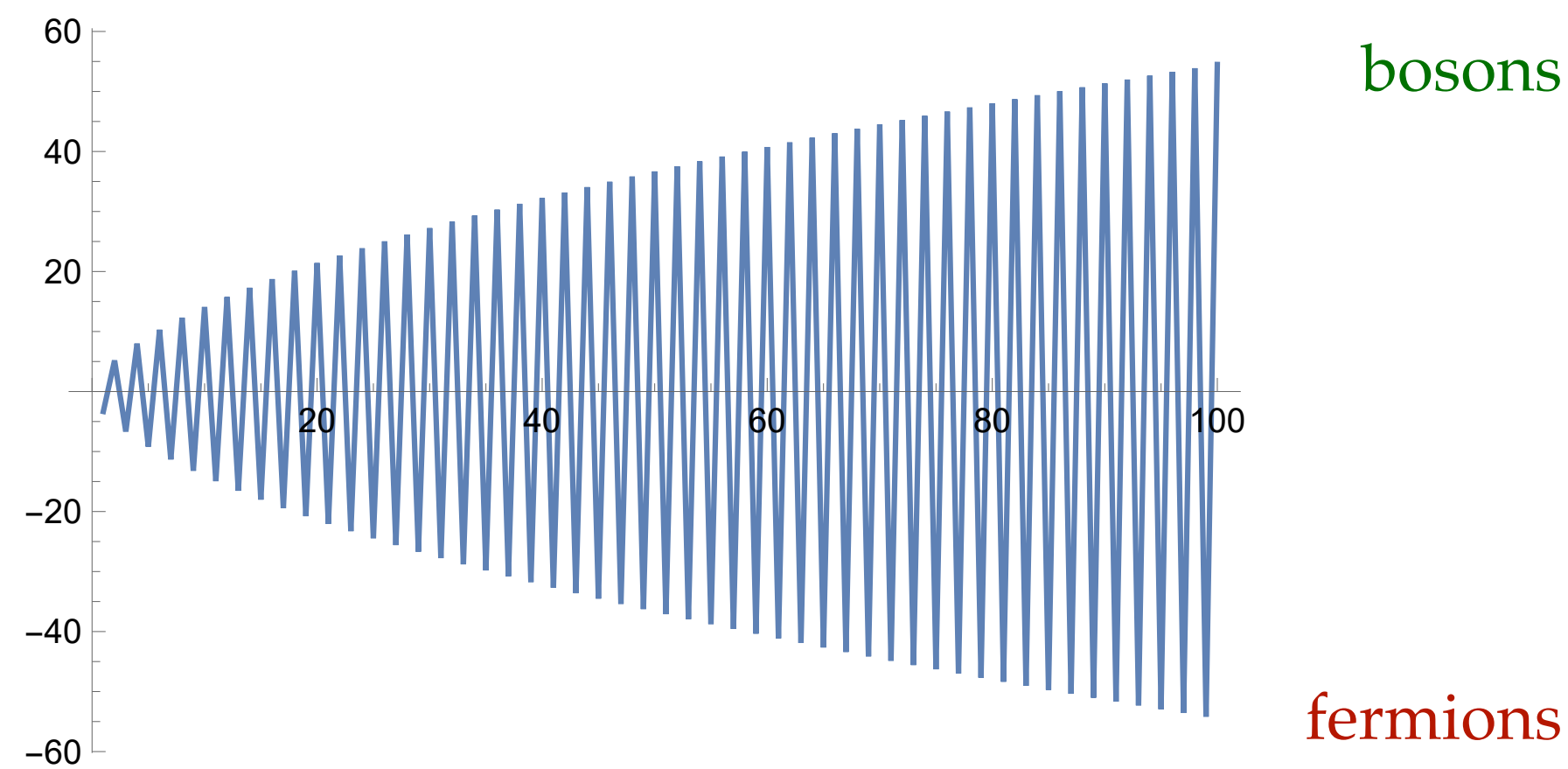
# TACHYONS AND MISALIGNED SUPERSYMMETRY

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Based on work in progress with Giorgio Leone and Ioannis Florakis  
To appear soon

In the mid 90's Keith Dienes conjectured that a hidden *misaligned supersymmetry* is present in the string spectrum of non-tachyonic (non-supersymmetric) vacua



He argued that this *misaligned supersymmetry* be responsible for the finiteness of the one-loop vacuum energy of closed strings

# THE SPECTRUM OF CLOSED STRINGS

$$\begin{aligned}
 \mathcal{Z} &= \sum_s \text{tr}_s P_{\text{GSO}} (-1)^F q^H \bar{q}^{\bar{H}} \\
 &= \tau_2^{1-D/2} \sum_{i,j} \bar{\chi}_i(\bar{q}) \mathcal{N}_{ij} \chi_j(q) \\
 &= \tau_2^{1-D/2} \sum_{m,n} a_{mn} \bar{q}^m q^n
 \end{aligned}$$

Characters of the RCFT

Integer-valued matrix fixed by modular invariance

$a_{nn}$  are (signed) integers and count the number of degrees of freedom at the  $n$ -th mass level

## THE BASIC IDEA OF MISALIGNED SUPERSYMMETRY

$$\bar{\chi}_i(\bar{q}) \chi_j(q) \Rightarrow a_{nn}^{(ij)} \sim A n^{-B} e^{4\pi C_{\text{tot}} \sqrt{n}}$$

The exponential growth is determined by the central charge of the CFT

$$C_{\text{tot}} = C_{\text{left}} + C_{\text{right}} \equiv \sqrt{\frac{c_{\text{left}}}{24}} + \sqrt{\frac{c_{\text{right}}}{24}}$$

## THE BASIC IDEA OF MISALIGNED SUPERSYMMETRY

K. Dienes proved that, when tachyons are absent,

$$\langle a_{nn} \rangle = \sum_{i,j} \mathcal{N}_{ij} a_{nn}^{(ij)} \sim e^{4\pi C_{\text{eff}} \sqrt{n}}$$

with  $C_{\text{eff}} < C_{\text{tot}}$

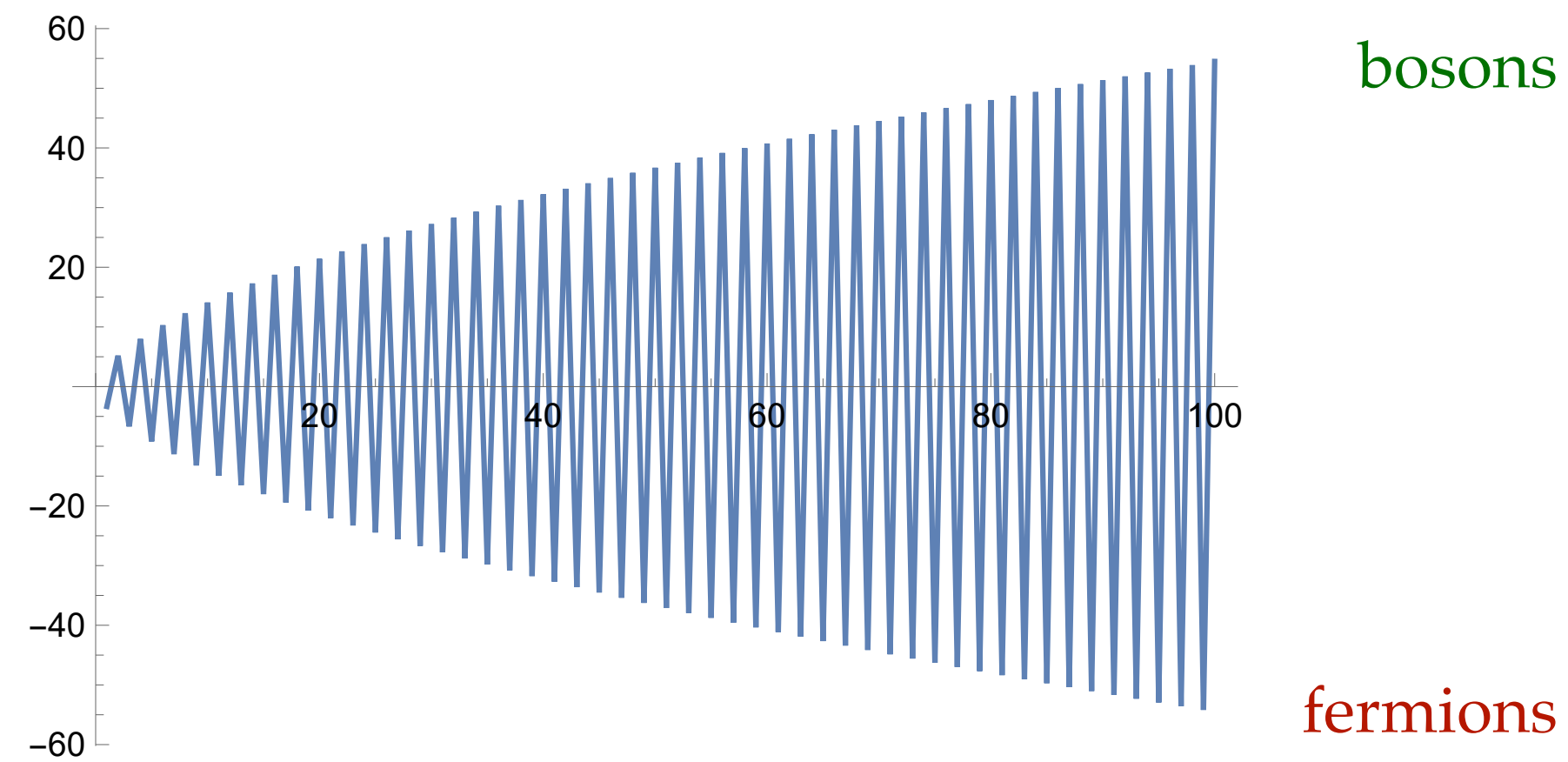
He then conjectured that

$$C_{\text{eff}} = 0$$

[Dienes 1994]

## THE BASIC IDEA OF MISALIGNED SUPERSYMMETRY

He argued that  $C_{\text{eff}} < C_{\text{tot}}$  implies oscillations in the string spectrum



*What is the right trademark for the absence of tachyons?*

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To (try to) answer this question, we have studied all  
non-supersymmetric theories in ten dimensions,  
and their freely-acting deformation in nine dimensions

A clear pattern emerges from this analysis

# NON-SUPERSYMMETRIC STRING VACUA

In D=10 there are 8 non-supersymmetric vacua

$$\begin{aligned} \mathcal{Z}_{16 \times 16} = & V_8(\bar{O}_{16}\bar{O}_{16} + \bar{S}_{16}\bar{S}_{16}) - S_8(\bar{O}_{16}\bar{S}_{16} + \bar{S}_{16}\bar{O}_{16}) \\ & + O_8(\bar{V}_{16}\bar{C}_{16} + \bar{C}_{16}\bar{V}_{16}) - C_8(\bar{V}_{16}\bar{V}_{16} + \bar{C}_{16}\bar{C}_{16}) \end{aligned}$$

tachyon free

$$\mathcal{Z}_{32} = V_8 \bar{O}_{32} - S_8 \bar{S}_{32} + O_8 \bar{V}_{32} - C_8 \bar{C}_{32}$$

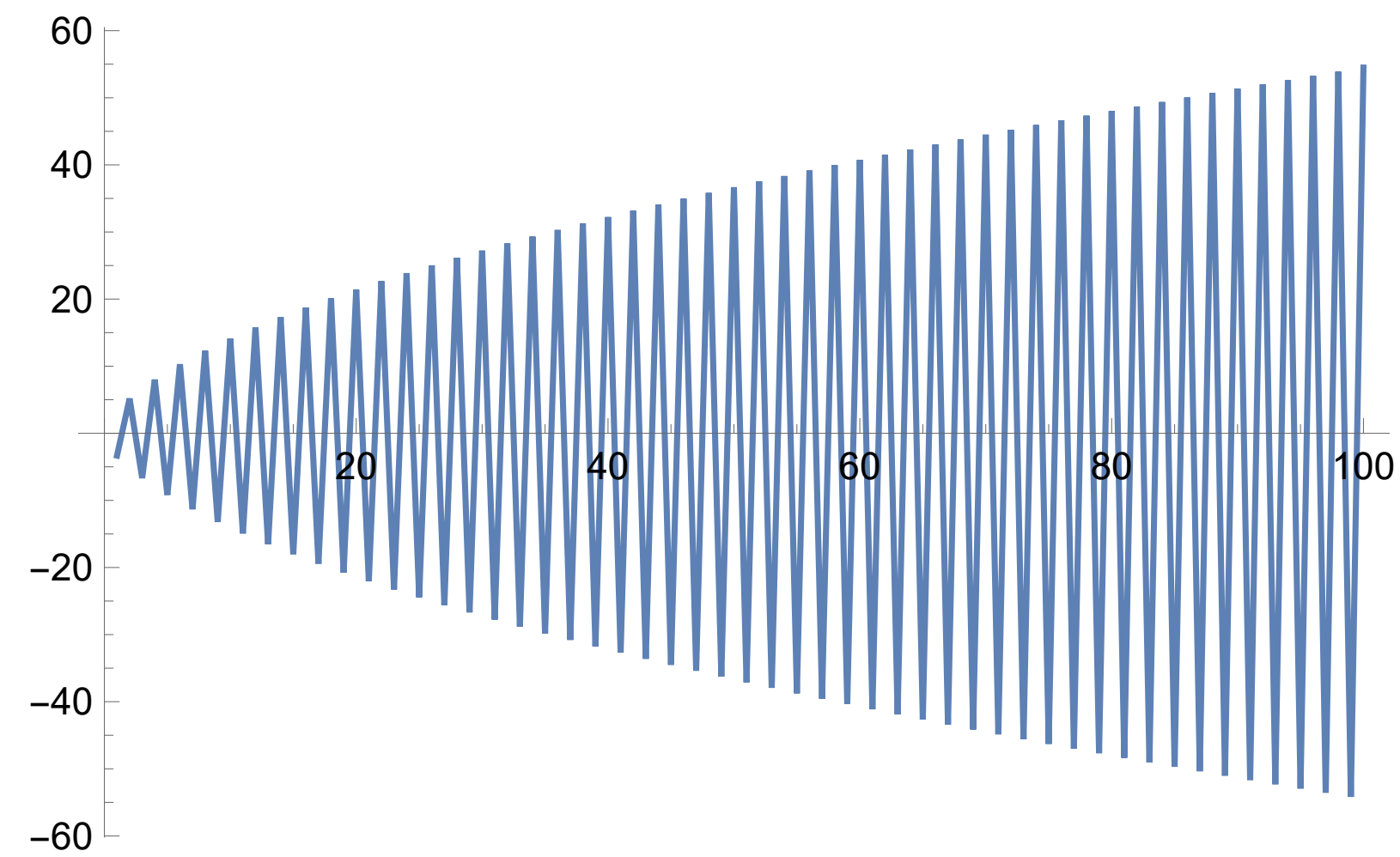
tachyonic

$$SO(24) \times SO(8) \qquad SU(2)^2 \times E_7^2 \qquad SO(16) \times E_8 \qquad SU(16) \times U(1)$$

[0A and 0B theories are purely bosonic]

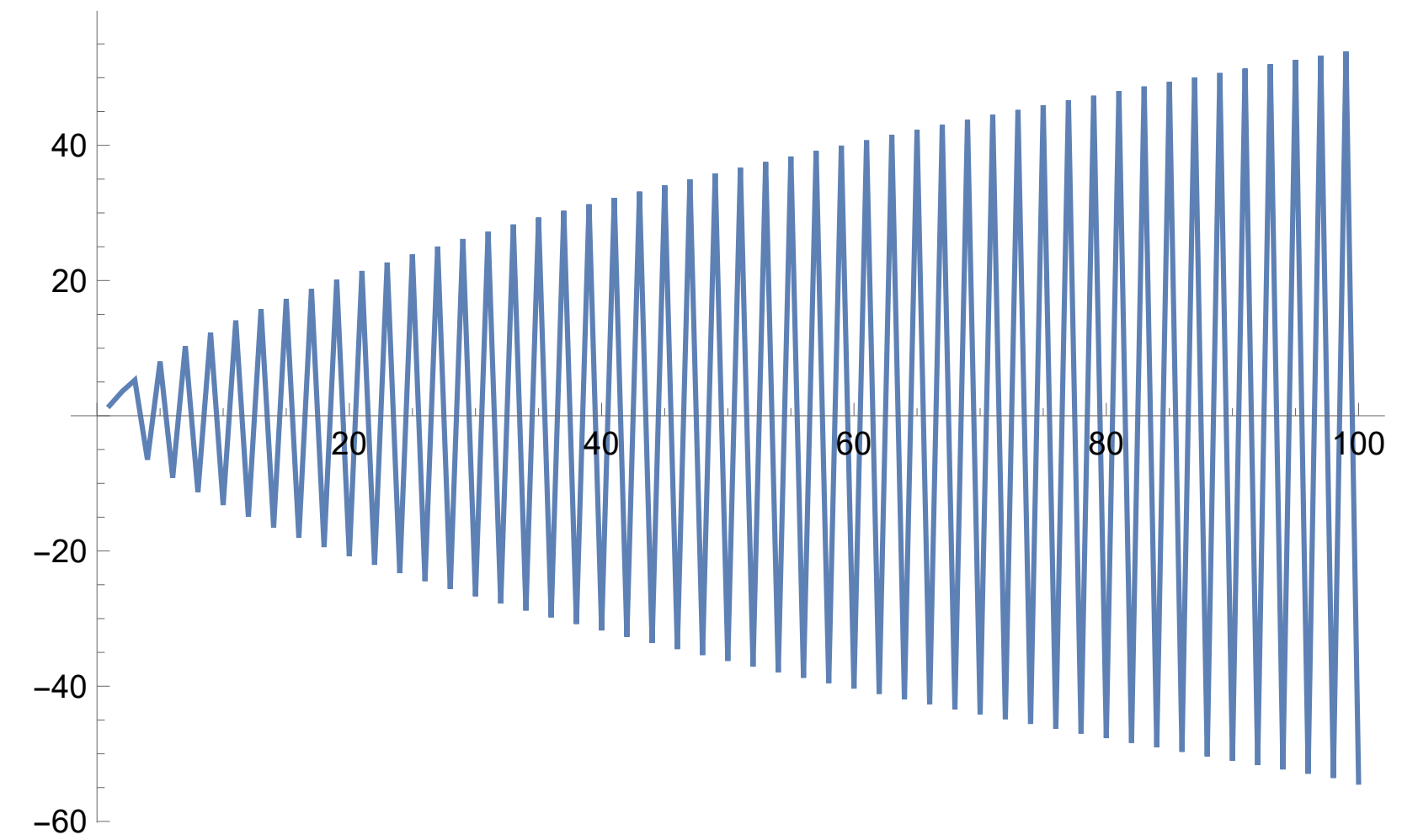


# OSCILLATIONS IN THE STRING SPECTRUM



$O(16) \times O(16)$

[tachyon free]



$SO(32)$

[tachyonic]

## A MORE QUANTITATIVE ANALYSIS OF THE SPECTRUM

$$\chi_i(q) = \sum_n a_n^{(i)} q^{n+H_i} \qquad H_i = h_i - c/24$$

The *circle method* of Hardy-Ramanujan yields the asymptotic growth

$$a_n^{(i)} = i^k \sum_{\ell=1}^{O([\sqrt{n}])} \sum_{\substack{p=0 \\ \gcd(\ell,p)=1}}^{\ell-1} \sum_j \left( \gamma_{\ell,p}^{-1} \right)_{ij} e^{-\frac{2\pi i}{\ell}((p(n+H_i)-p'H_j))} \frac{2\pi a_0^{(j)}}{\ell} \left( \frac{H_j}{n+H_i} \right)^{\frac{1-k}{2}} J_{k-1} \left( \frac{4\pi}{\ell} \sqrt{H_j(n+H_i)} \right)$$

[Kani, Vafa 1990]

Characters with negative  $H_j$  contribute to the exponential growth,  $J_{k-1} \rightarrow I_{k-1}$

A CFT can have many characters with negative  $H_j$   
The most negative one plays the role of the *identity*

# A MORE QUANTITATIVE ANALYSIS OF THE SPECTRUM

Taking into account only the *tachyonic* contributions (this with  $H_j<0$ )

$$\begin{aligned} a_n^{(i)} \sim & e^{4\pi\sqrt{|H_{\min}|}n} + \sum_{j\neq\min} Q_{ij} \, e^{4\pi\sqrt{|H_j|}n} \\ & + \sum_{\ell} P_i(\ell;n) \, e^{\frac{4\pi}{\ell}\sqrt{|H_{\min}|}n} + \sum_{j\neq\min} P_{ij}(\ell;n) \, e^{\frac{4\pi}{\ell}\sqrt{|H_j|}n} \end{aligned}$$

$$\begin{aligned} P(\ell;n+\ell) &= P(\ell;n) \\ \sum_{n=0}^{\ell-1} P(\ell;n) &= 0 \end{aligned}$$

The contribution of the *identity* is *universal*

## THE GROWTH OF THE DEGREES OF FREEDOM

$$\sum_{i,j} \bar{\chi}_i(\bar{q}) \mathcal{N}_{ij} \chi_j(q) \quad \Rightarrow \quad \langle a_{nn} \rangle \sim \sum_{ij} e^{4\pi \sqrt{|H_{\min}^L|n}} \mathcal{N}_{ij} e^{4\pi \sqrt{|H_{\min}^R|n}}$$

Therefore, whenever  $\sum_{i,j} \mathcal{N}_{ij} = 0$

$$C_{\text{eff}} < C_{\text{tot}}$$

This occurs for *all* non-supersymmetric heterotic theories in D=10

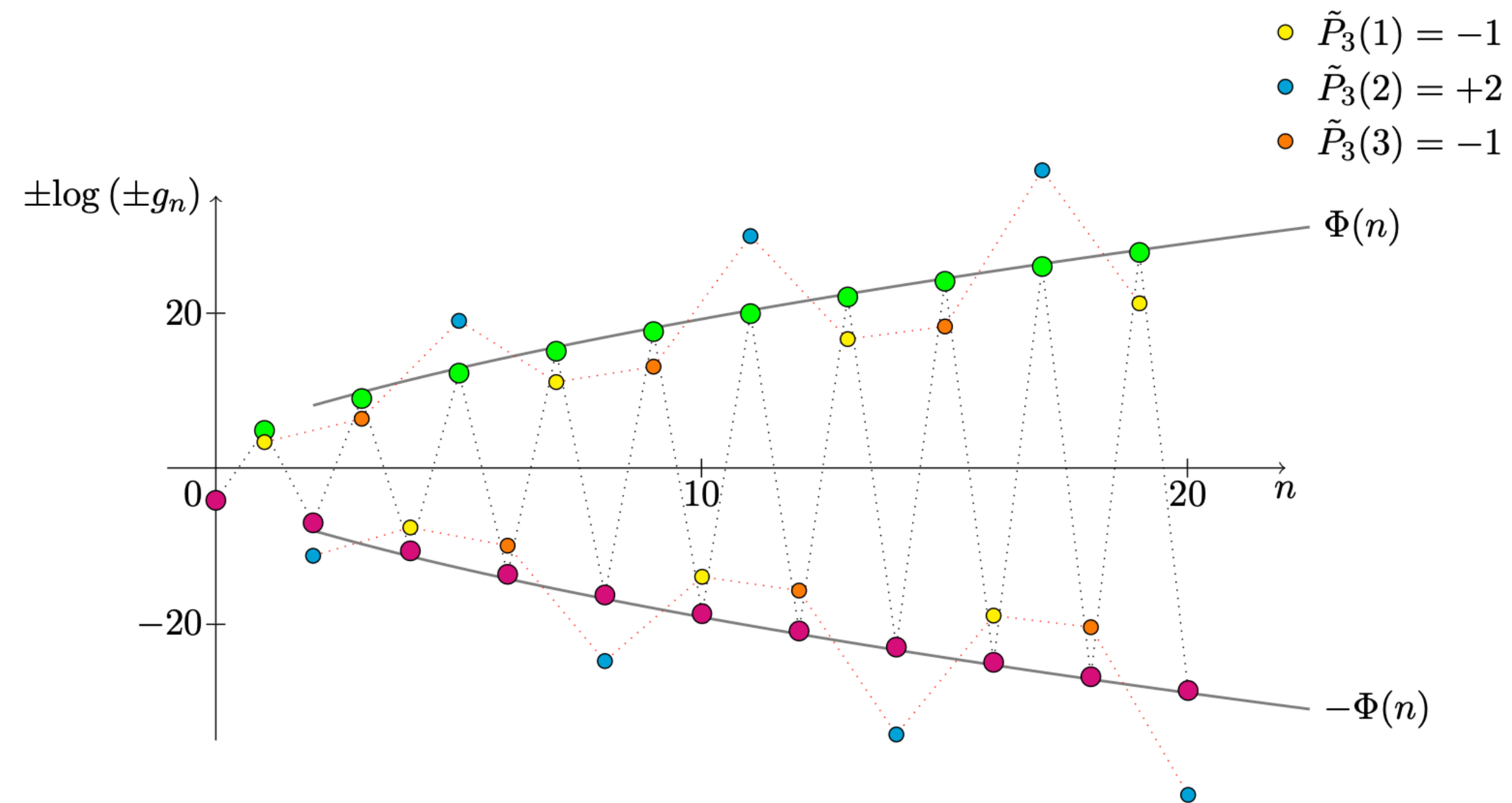
What about the sub-leading exponentials?

# THE GROWTH OF THE DEGREES OF FREEDOM

For the  $O(16) \times O(16)$  heterotic string in D=10 it was “shown” that

$$C_{\text{eff}} = 0$$

[Cribiori et al 2021]



[plot from Cribiori et al 2021]

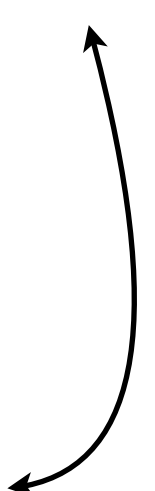
# THE GROWTH OF THE DEGREES OF FREEDOM

For the  $SO(32)$  heterotic string in D=10 instead

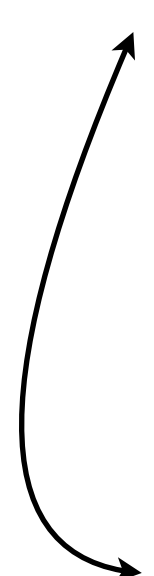
$$\langle a_{nn} \rangle \sim 2 e^{4\pi\sqrt{n/2}} e^{4\pi\sqrt{n/2}} + e^{\frac{4\pi}{2}\sqrt{n/2}} e^{\frac{4\pi}{2}\sqrt{n/2}} + 4 e^{\frac{4\pi}{3}\sqrt{n/2}} e^{\frac{4\pi}{3}\sqrt{n/2}}$$

$\ell = (1, 1)$ 
 $\ell = (2, 2)$ 
 $\ell = (3, 3)$

leading R tachyonic  
character ( $O_8$ )



sub-leading L tachyonic  
character ( $V_8$ )



and thus  $C_{\text{eff}} = 2\sqrt{1/2} < 1 + \sqrt{1/2} = C_{\text{tot}}$

## NINE-DIMENSIONAL FREELY ACTING CONSTRUCTIONS

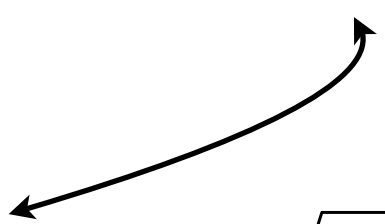
All non-supersymmetric vacua can be built from supersymmetric ones via an involution which contains fermion number

Combining this involution with a shift along a compact direction implies that supersymmetry is broken spontaneously

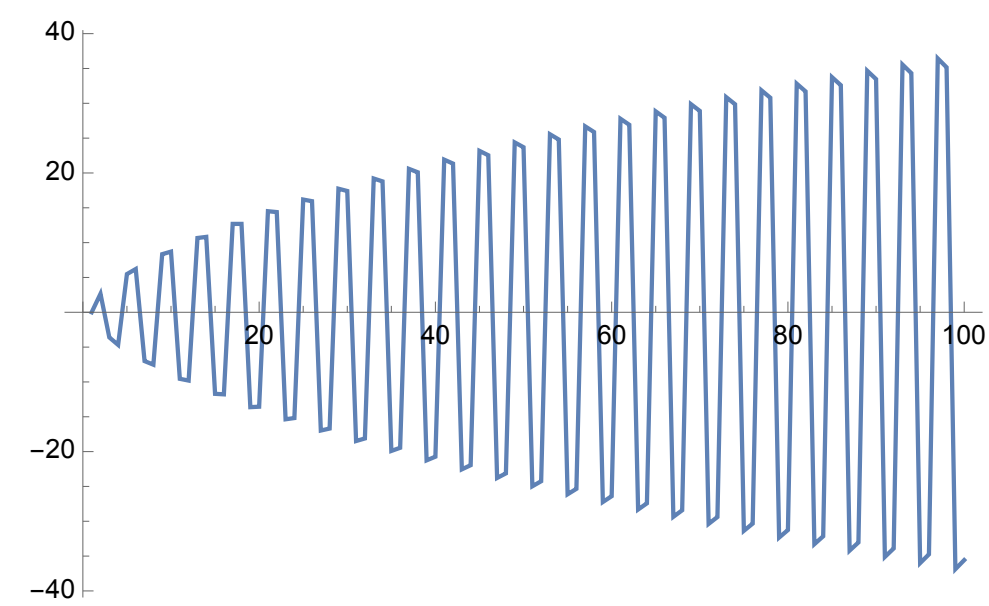
For instance,

$$\mathcal{Z} = \left[ V_8 \Lambda_{2m,n} - S_8 \Lambda_{2m+1,n} + O_8 \Lambda_{2m+1,n+\frac{1}{2}} - C_8 \Lambda_{2m,n+\frac{1}{2}} \right] (\bar{O}_{32} + \bar{S}_{32})$$

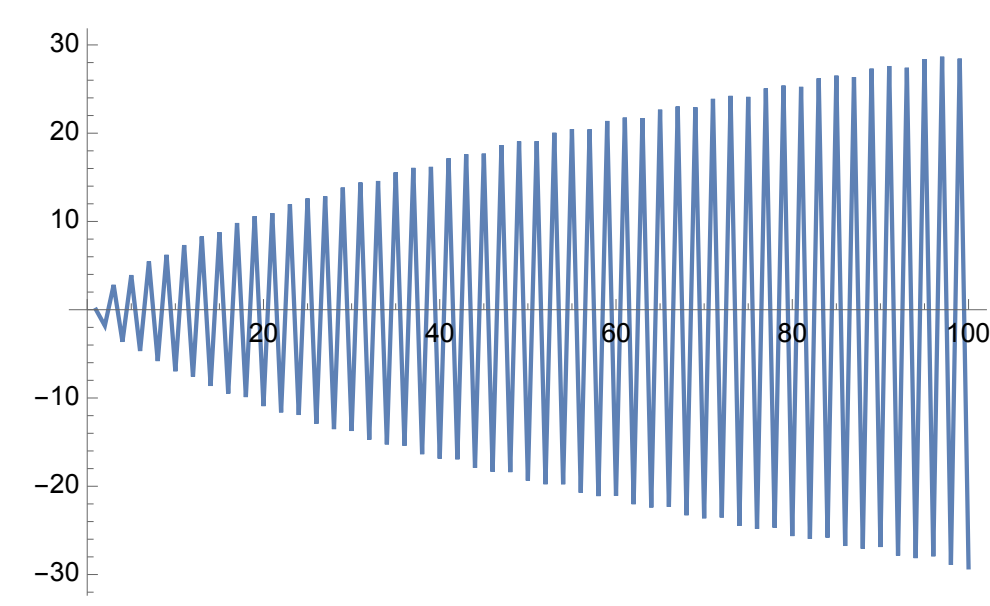
tachyonic for  $\sqrt{2\alpha'}\sqrt{3-2\sqrt{2}} < R < \sqrt{2\alpha'}\sqrt{3+2\sqrt{2}}$



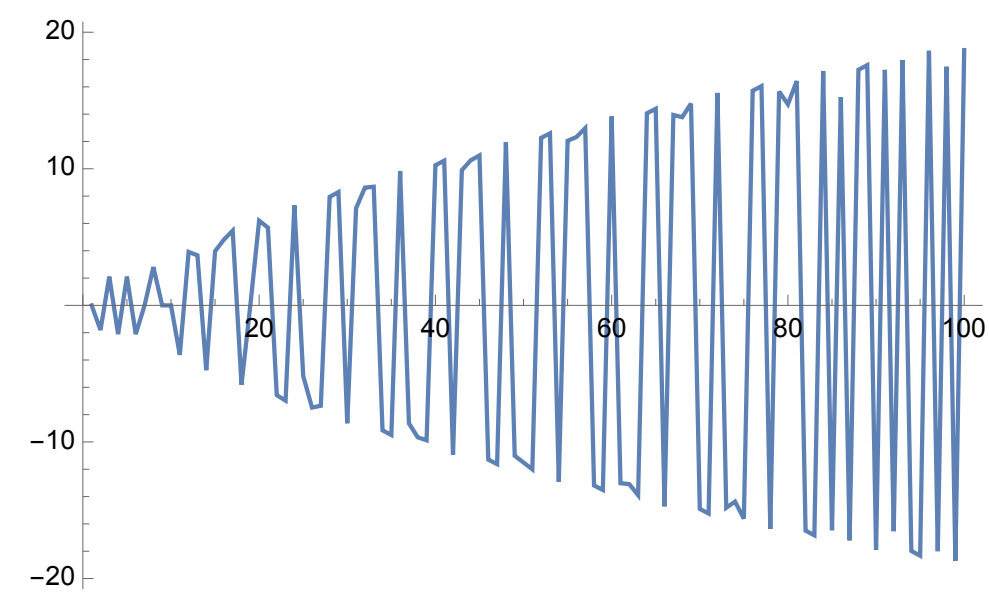
NINE-DIMENSIONAL FREELY-ACTING REALISATION  $O(16) \times O(16)$



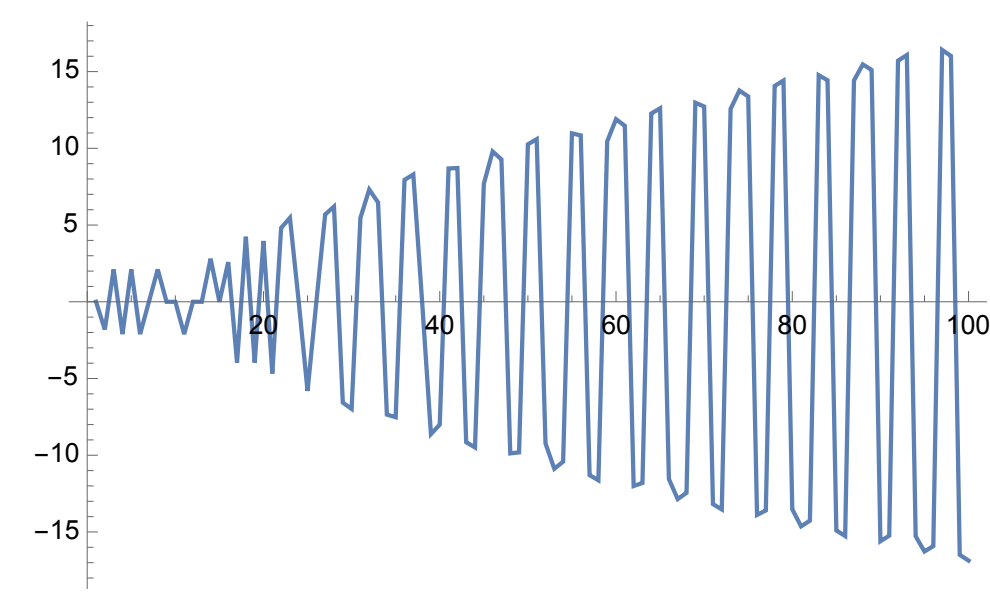
$$R = \sqrt{2\alpha'}$$



non-tachyonic region



$$R = \sqrt{12\alpha'}$$

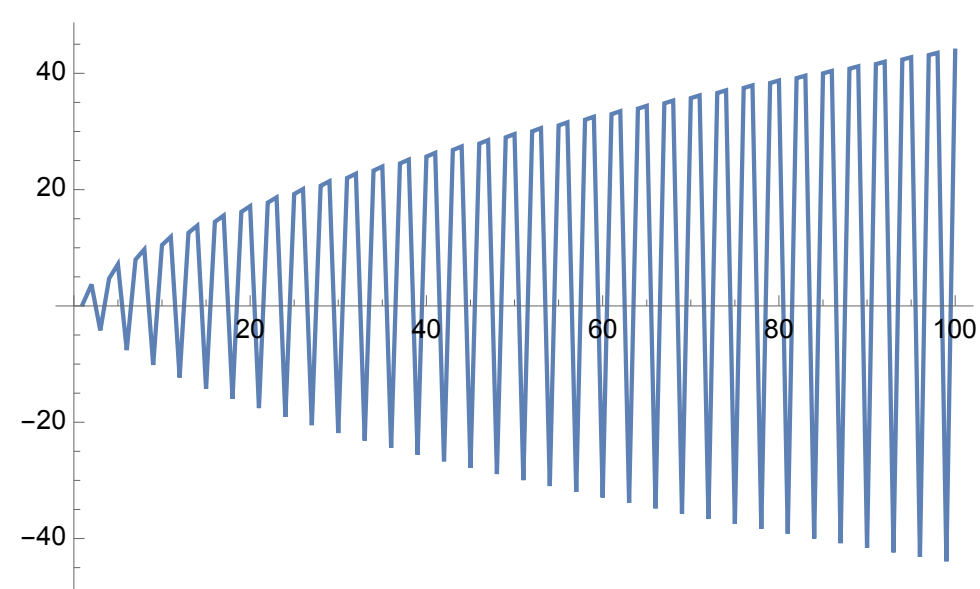


non-tachyonic region

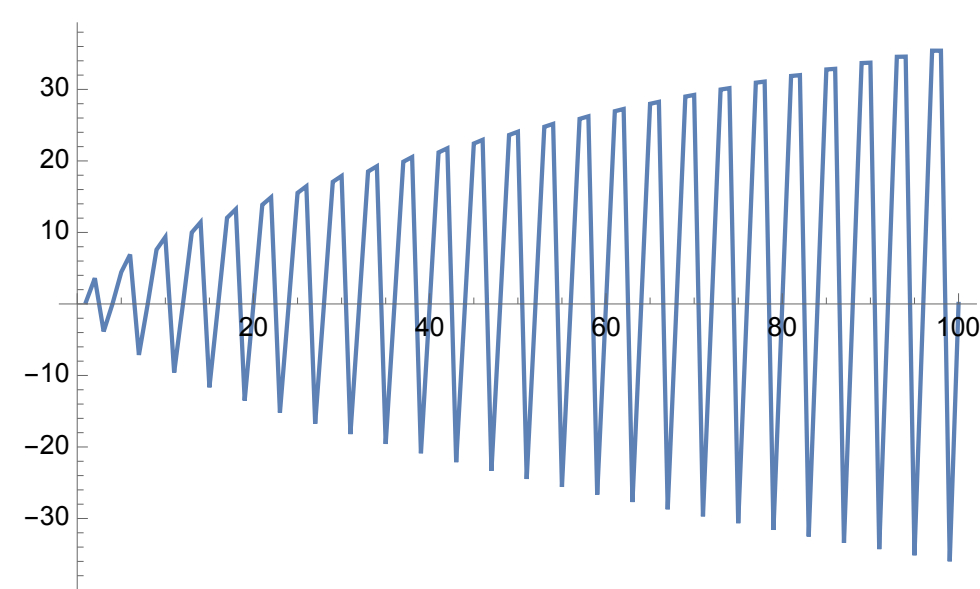
$$R = \sqrt{16\alpha'}$$



NINE-DIMENSIONAL FREELY-ACTING REALISATION  $SO(32)$

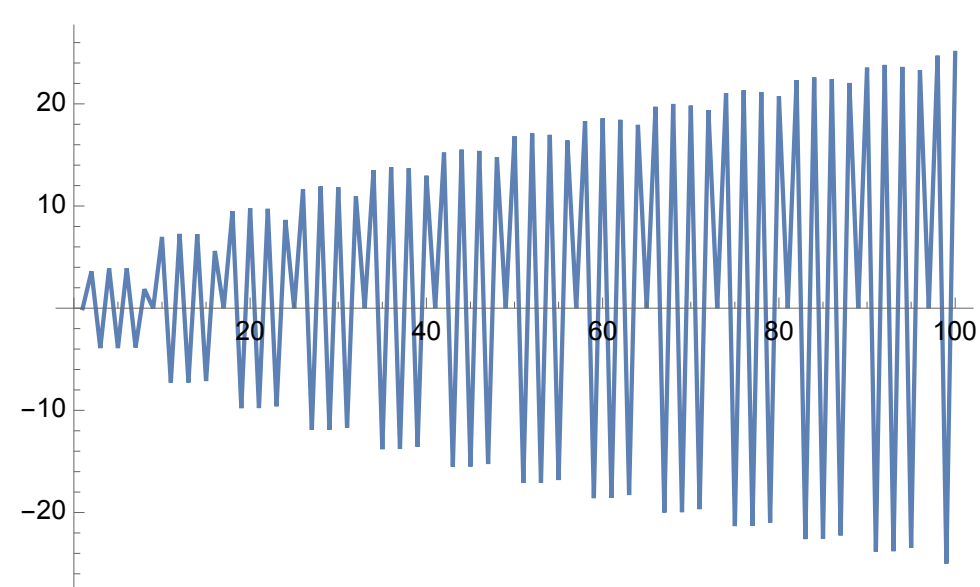


$$R = \sqrt{2\alpha'}$$

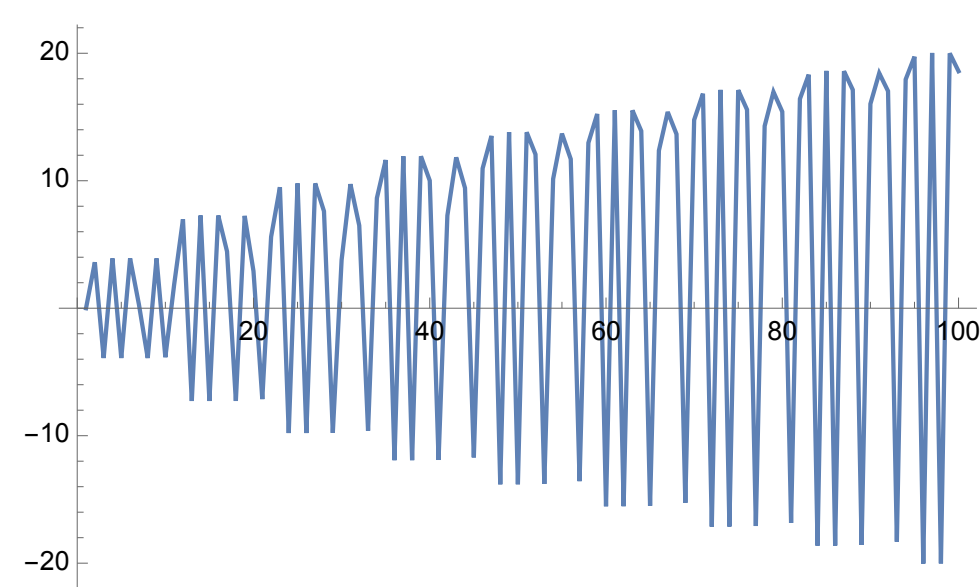


$$R = \sqrt{4\alpha'}$$

tachyonic region



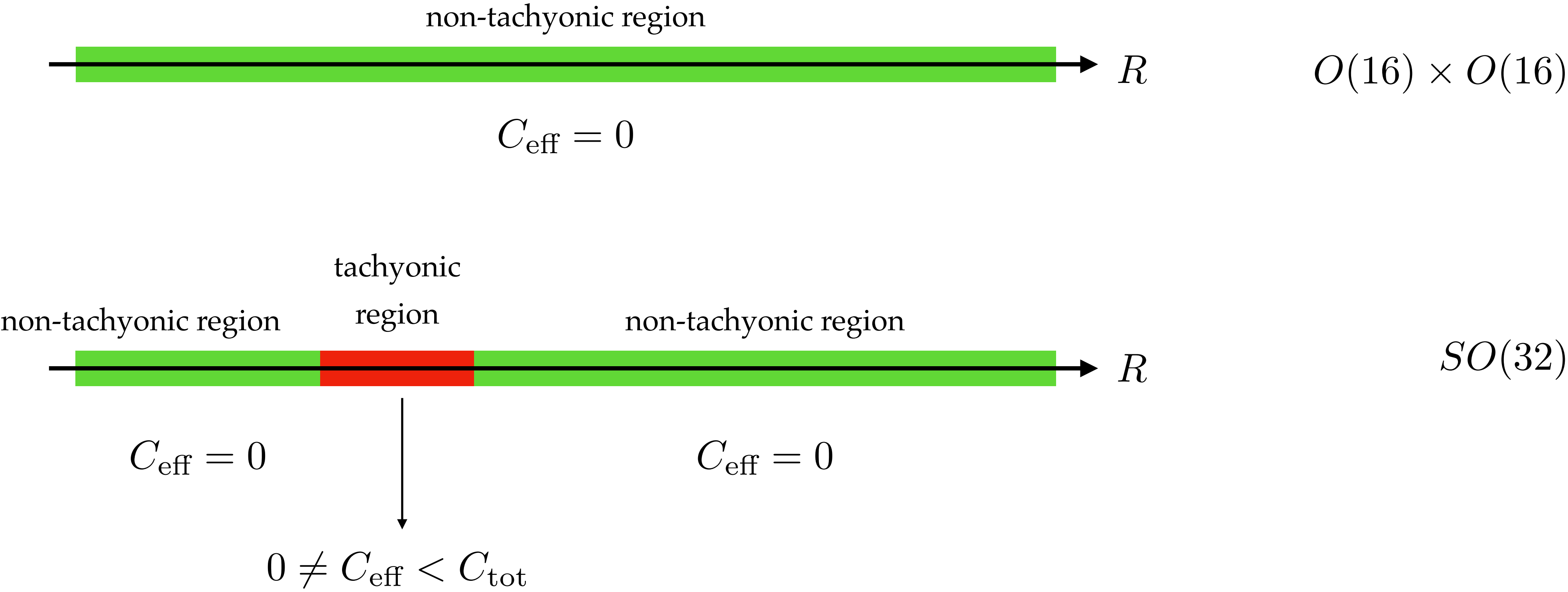
$$R = \sqrt{12\alpha'}$$



$$R = \sqrt{16\alpha'}$$

non-tachyonic region

NINE-DIMENSIONAL FREELY-ACTING REALISATION



# NINE-DIMENSIONAL FREELY-ACTING REALISATION

The analysis is much more complicated because of the presence of many tachyonic characters in the RCFT

For instance, at the point  $R = \sqrt{16\alpha'}$  there are 32 characters

Left sector	Right sector
$H_{12} = -\frac{7}{64}$ $H_{14} = -\frac{23}{64}$	$H_1 = -1$ $H_2 = -\frac{63}{64}$ $H_3 = -\frac{15}{16}$ $H_4 = -\frac{55}{64}$ $H_5 = -\frac{3}{4}$
$H_{16} = -\frac{31}{64}$ $H_{18} = -\frac{31}{64}$	$H_6 = -\frac{39}{64}$ $H_7 = -\frac{7}{16}$ $H_8 = -\frac{15}{64}$ $H_{26} = -\frac{15}{64}$ $H_{27} = -\frac{7}{16}$
$H_{20} = -\frac{23}{64}$ $H_{22} = -\frac{7}{64}$	$H_{28} = -\frac{39}{64}$ $H_{29} = -\frac{3}{4}$ $H_{30} = -\frac{55}{64}$ $H_{31} = -\frac{15}{16}$ $H_{32} = -\frac{63}{64}$

## CONCLUSIONS AND OUTLOOK

Oscillations are a universal feature of generic string vacua

Absence of tachyons (classical stability) seems to imply  $C_{\text{eff}} = 0$

*What is the connection with the results of Kutasov and Seiberg  
and C.A., Cardella, Elitzur and Rabinovici?*

[Kutasov, Seiberg 1991]

[C.A., Cardella, Elitzur,  
Rabinovici 1991]

## CONCLUSIONS AND OUTLOOK

*What is the connection with the results of Kutasov and Seiberg  
and C.A., Cardella, Elitzur and Rabinovici?*

$$\sum_{\{m^2\}} d(m^2) e^{-4\pi m^2 / \Lambda^2} = \frac{3}{\pi} \Lambda^{2-D} \Omega + \sum_{\zeta^*(\rho)=0} C_\rho \Lambda^{\rho-D}$$

degrees of freedom of  
the physical spectrum  
at mass  $m^2$

(finite) vacuum energy

oscillations

$$C_m \Lambda^{\beta_m - D} \cos(\gamma_m \log \Lambda + \phi_m)$$

$$\rho_m = \beta_m + i\gamma_m$$

[C.A., Cardella, Elitzur,  
Rabinovici 1991]

**THANK YOU**