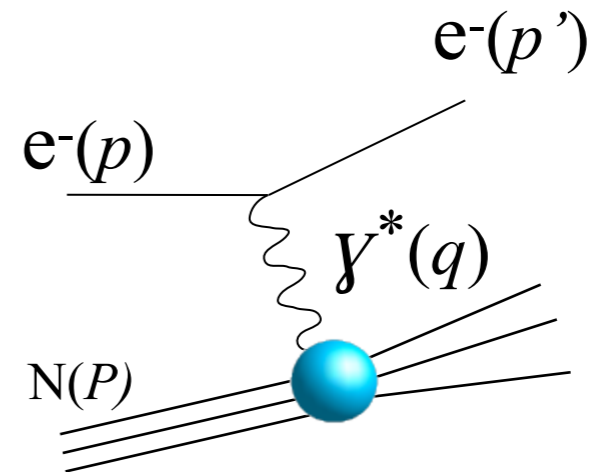


Measurements of α_s from spin structure functions

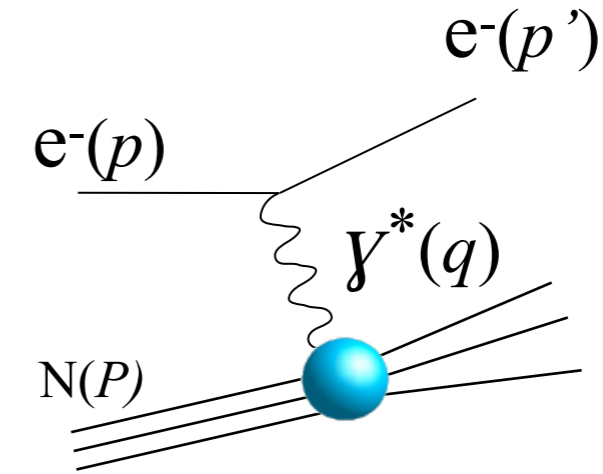
A. Deur
Jefferson Lab

Reminder: inclusive lepton-nucleon scattering

- ◆ $p=(E,\mathbf{p})$, $p'=(E-\nu,\mathbf{p}-\mathbf{q})$, $q=(\nu,\mathbf{q})$
- ◆ γ^* virtual photons: q^2
- ◆ Since $q^2 < 0$ here, we use $Q^2 = -q^2$.
- ◆ Inclusive experiments: only the scattered electrons are detected: target or target fragments are ignored.
- ◆ At high energy, Bjorken scaling variable $x = Q^2 / 2M\nu$ is more convenient than ν .



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Cross section: $\sigma = \sigma_{\text{Mott}} [\alpha F_1(x, Q^2) + \beta F_2(x, Q^2) + \gamma g_1(x, Q^2) + \varpi g_2(x, Q^2)]$

pointlike scattering × (spin independent + spin dependent)

F_1, F_2, g_1 and g_2 : **structure functions**

F_1 and F_2 , are obtained with **unpolarized** beam and target and varying kinematic factors α and β .

g_1 and g_2 are obtained with beam and target both **polarized**, measuring beam spin asymmetries and varying the target spin direction.

Considering the nucleon inclusive spin structure, α_s can be extracted from:

- Q^2 -evolution of $g_1(x, Q^2)$. Complex: involves DGLAP global fit, non-perturbative inputs: quark and gluon distributions, possibly higher-twists for low- Q^2 / large- x data.
- Q^2 -evolution of moment $\int_0^1 g_1(x, Q^2) dx$. Simpler: no x -dependence, non-perturbative inputs: more-or-less well measured axial charges a_0 , a_3 and a_8 (+ possibly higher-twists for low- Q^2 data). Issues: unmeasurable low- x contribution, a_0 is Q^2 -dependent and may have contribution from gluon ΔG pdf (but not the case in \overline{MS}).
- Q^2 -evolution of isovector moment $\int_0^1 g_1^{p-n}(x, Q^2) dx$, i.e. [Bjorken sum](#). Simplest. Axial charge $a_3 = g_A$ precisely measured ($g_A = 1.2762 \pm 0.0005$). DGLAP-evolution known to higher order than single nucleon case. No gluon contribution. But low- x issue and demands measurement on polarized p and n.

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Bjorken sum rule

$$\int g_1^{p-n} dx \equiv \Gamma_1^{p-n} = \frac{1}{6} g_A \left[1 - \frac{\alpha_s}{\pi} - 3.58 \left(\frac{\alpha_s}{\pi} \right)^2 - 20.21 \left(\frac{\alpha_s}{\pi} \right)^3 - 175.7 \left(\frac{\alpha_s}{\pi} \right)^4 - \sim 893 \left(\frac{\alpha_s}{\pi} \right)^5 \right] + \frac{M^2}{Q^2} \left[a_2(\alpha_s) + 4d_2(\alpha_s) + 4f_2(\alpha_s) \right] + \dots$$

↑
 Nucleon's First spin structure function

↑
 Nucleon axial charge. (Value of $\Gamma_1^{p-n}(Q^2)$ in the $Q^2 \rightarrow \infty$ limit)

↑
 pQCD radiative corrections (\overline{MS} Scheme.)

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 Non-perturbative $1/Q^{2n}$ power corrections. (+rad. corr.)

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$\int g_1^{p-n} dx \equiv \Gamma_1^{p-n}$ → Nucleon's First spin structure function
 $\frac{1}{6} g_A$ → Nucleon axial charge. (Value of $\Gamma_1^{p-n}(Q^2)$ in the $Q^2 \rightarrow \infty$ limit)
 $- 3.58 \left(\frac{\alpha_s}{\pi} \right)^2 - 20.21 \left(\frac{\alpha_s}{\pi} \right)^3 - 175.7 \left(\frac{\alpha_s}{\pi} \right)^4 - \sim 893 \left(\frac{\alpha_s}{\pi} \right)^5$ → pQCD radiative corrections (\overline{MS} Scheme.)
 $+ \frac{M^2}{Q^2} \left[a_2(\alpha_s) + 4d_2(\alpha_s) + 4f_2(\alpha_s) \right] + \dots$ → Non-perturbative $1/Q^{2n}$ power corrections. (+rad. corr.)

⇒ Two possibilities to extract $\alpha_s(M_Z)$:

- Do an absolute measurement of $\Gamma_1^{p-n}(Q^2)$ and solve the Bj SR for $\alpha_s(Q^2)$.
 - One α_s per Γ_1^{p-n} experimental data point.
 - Poor systematic accuracy, typically $\Delta\alpha_s/\alpha_s \sim 10\% \Rightarrow$ Not competitive.
- Measurement of Q^2 -dependence of $\Gamma_1^{p-n}(Q^2)$.
 - Need several Γ_1^{p-n} points. Only one (or a few) value of α_s .
 - Good accuracy: 1990's CERN/SLAC data yielded: $\alpha_s(M_Z) = 0.120 \pm 0.009$

Altarelli, Ball, Forte, Ridolfi, Nucl.Phys. B496 337 (1997)

α_s from $\Gamma_1^{p-n}(Q^2)$ measurement in JLab Hall B

AD *et al.* PRD 90, 012009 (2014)

EG1dvcs experiment:

- 6 GeV era of JLab.
- CLAS (CEBAF Large Acceptance Spectr.): 18-48° polar coverage, ~full azimuthal coverage.
- Polarized NH₃ and ND₃ fixed targets.
- Polarized beam.
- **High inclusive statistics** (DVCS process meas.): 6 months, 7 nA $\Rightarrow 2 \times 10^{17}$ e⁻ on target.

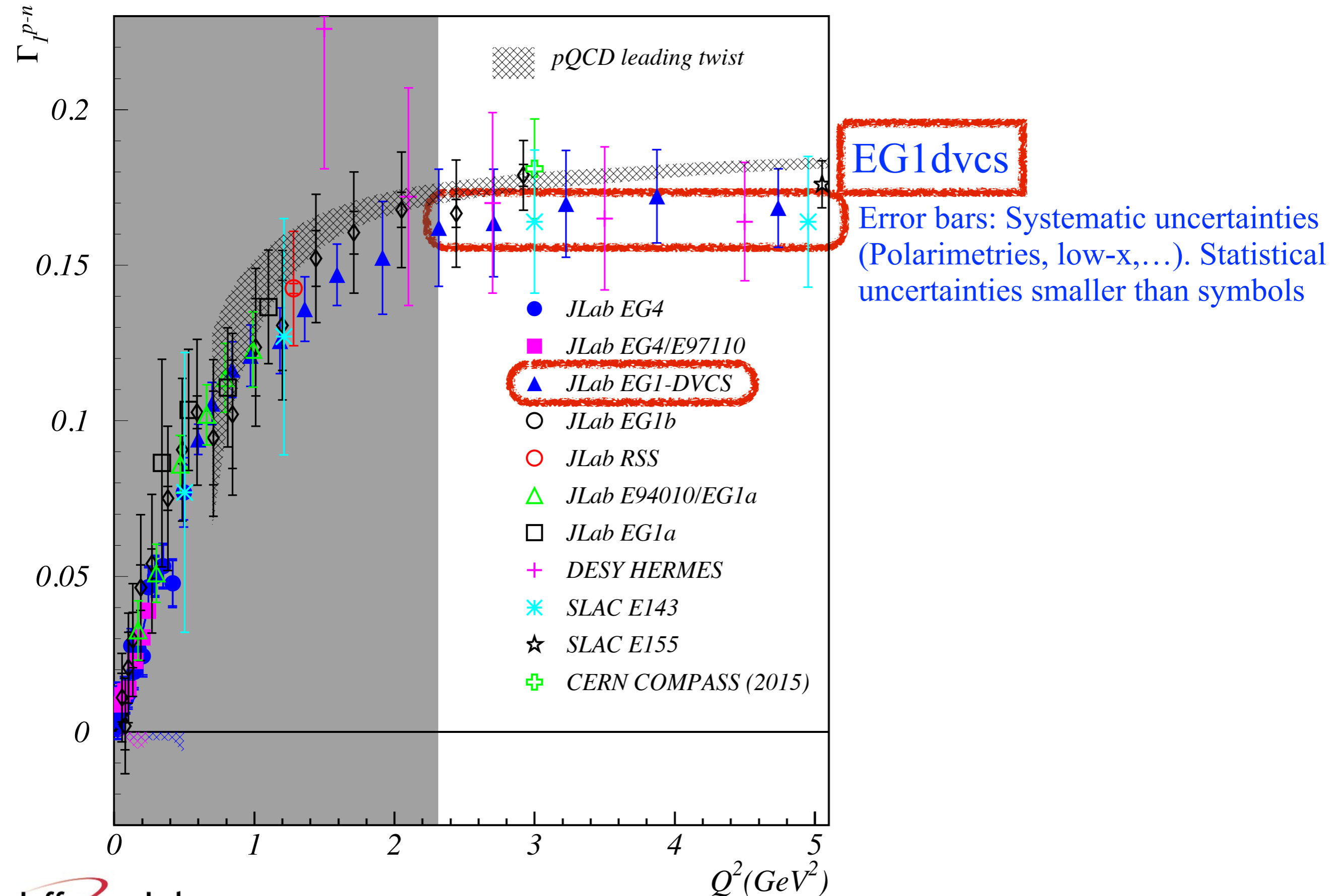
Used EG1dvcs data¹ to avoid uncorrelated systematics between experiments.

- Point-to-point correlated systematics (e.g. polarimetries, nuclear corrections) have minimal impact on uncertainties.
- EG1dvcs data largely dominate world data for **statistics**.
- Restricted Q² range: 2.32 < Q² < 4.74 GeV² rather than 2. < Q² < 10 GeV².

¹ Not fully true: important missing low-x contribution estimated from models fitting world data.

Q ² (GeV ²)	x-range (p)	x-range (d)	$\Gamma_{1,meas}^{p-n}$	$\Gamma_{1,meas+hi.x}^{p-n}$	σ_{meas}^{syst}	$\sigma_{hi.x}^{syst}$	$\Gamma_{1,tot}^{p-n}$	σ^{syst}	σ^{stat}	$\Gamma_{1,meas+hi.x}^{p-n} / \Gamma_{1,tot}^{p-n}$
2.316	0.263-0.864	0.271-0.798	0.0523	0.0515	0.0177	0.0001	0.1621	0.0188	0.0008	0.317
2.707	0.304-0.825	0.326-0.769	0.0398	0.0388	0.0157	0.0008	0.1636	0.0173	0.0006	0.237
3.223	0.362-0.901	0.385-0.799	0.0322	0.0311	0.0152	0.0000	0.1697	0.0171	0.0005	0.183
3.871	0.438-0.893	0.463-0.762	0.0227	0.0206	0.0121	0.0002	0.1721	0.0150	0.0004	0.120
4.739	0.531-0.909	0.663-0.738	0.0145	0.0113	0.0081	0.0002	0.1684	0.0126	0.0002	0.067

Measurement at 6 GeV in JLab Hall B



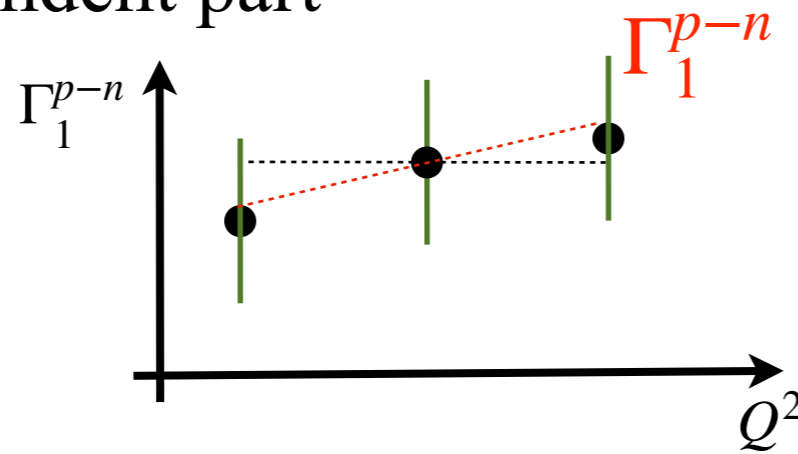
Uncertainties

Experimental systematic uncertainty: Separate point-to-point correlated and point-to-point uncorrelated parts:

- Fit data with expected pQCD form.
- Use *Unbiased estimate* (i.e. force $\chi^2=1$) to assess point-to-point uncorrelated uncert.
- Point-to-point correlated determined so that correlated \oplus uncorrelated = full experimental syst. uncertainties.

Low- x systematic uncertainties: separate Q^2 -dependent and independent parts.

Assume Q^2 -dependent part = $\frac{\text{low-}x \text{ contribution}}{dQ^2} \frac{d\Gamma_1^{p-n}}{dQ^2} \times Q^2\text{-bin size:}$



Add the Q^2 -dependent low- x uncertainty to the point-to-point correlated experimental uncertainty. Use this sum for χ^2 minimization of the fit that yields α_s .

\Rightarrow large parts of the low- x and exp. syst. uncertainties are suppressed.

Uncertainties

Leading uncertainties:

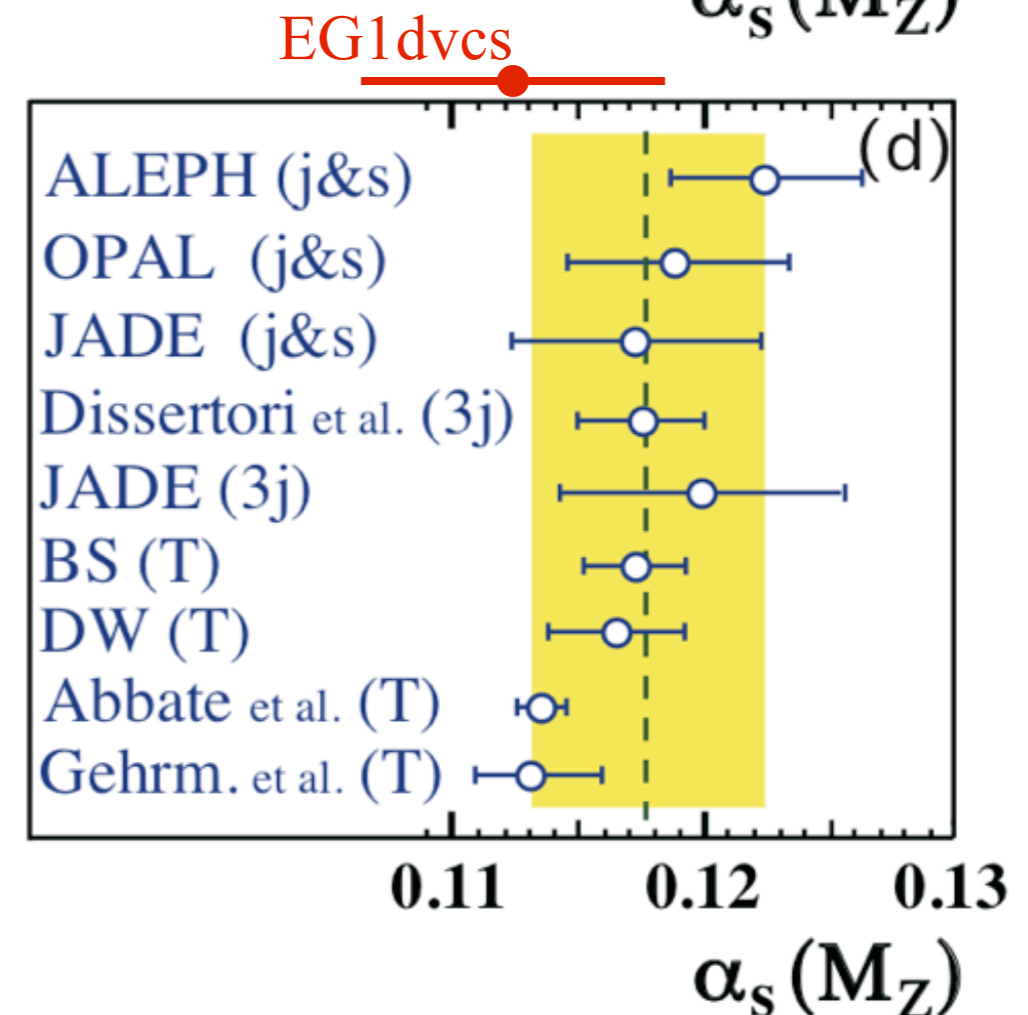
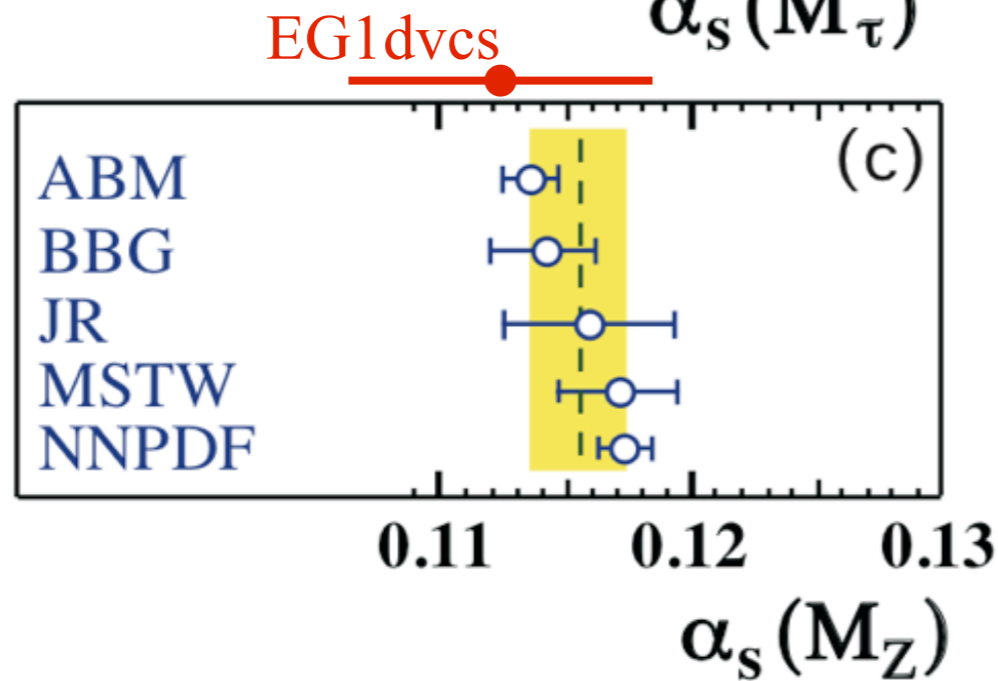
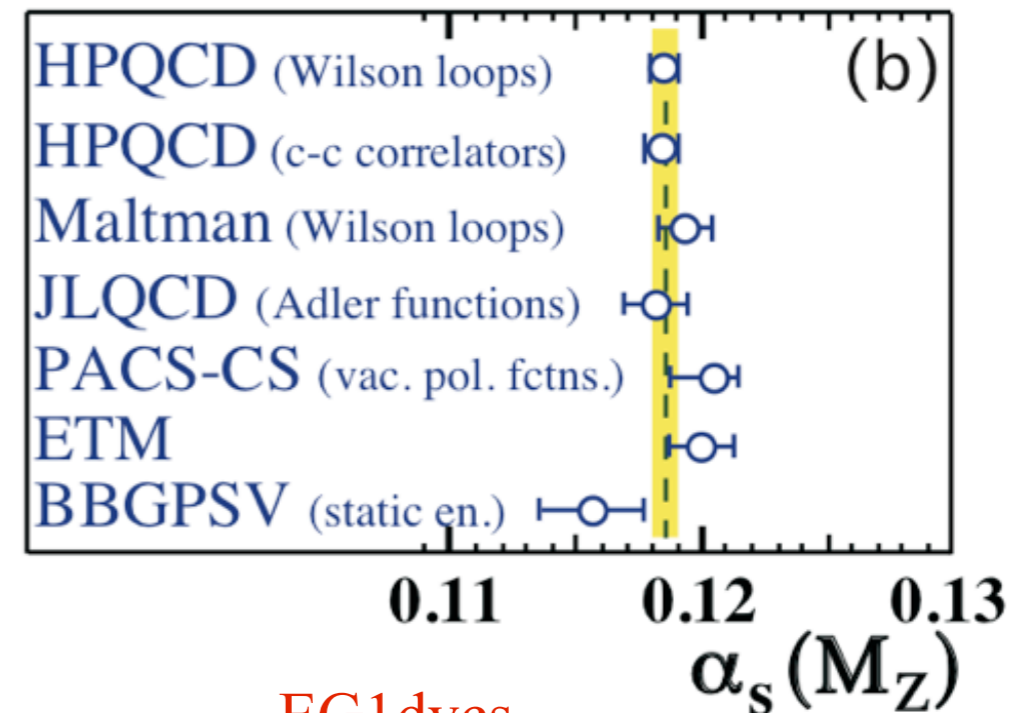
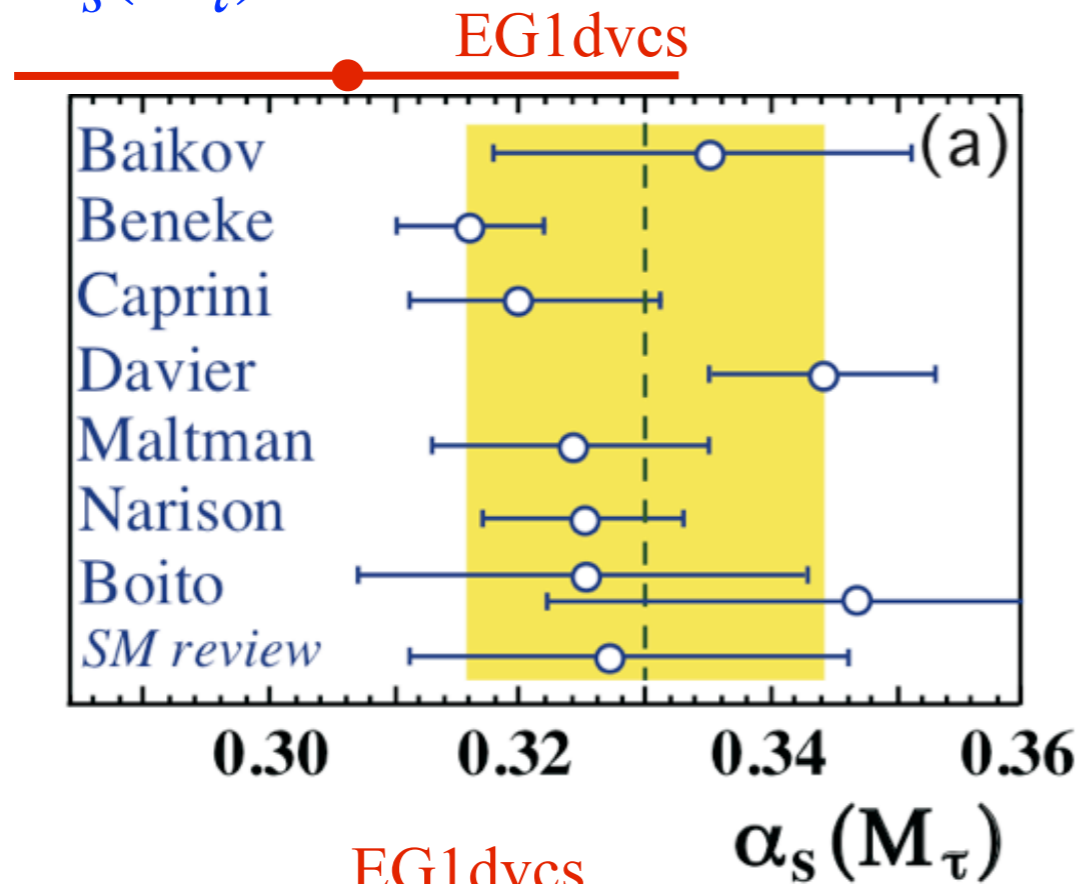
- Point-to-point uncorrelated uncertainties: 4.4%
 - Point-to-point correlated uncertainties: 3.3%
- } **5.5%**

Negligible uncertainties:

- Statistics.
- Twist-4 contributions: $\frac{M^2}{Q^2} [a_2(\alpha_s) + 4d_2(\alpha_s) + 4f_2(\alpha_s)]$
 a_2 (PDF fits) d_2 (meas.) f_2 (Sidorov-Weiss model, with 50% uncertainty). Twist > 4 neglected.
Uncertainties: < 0.1%.
- Bjorken twist-2 series truncation: 1.3%
- β -series truncation (needed it to evolve α_s to M_Z): 0.1%

Result: $\alpha_s(M_Z)=0.1123\pm 0.0061$
 $\alpha_s(M_\tau)=0.306\pm 0.053$

Compared to (then) best world data
(PDG 2014):



Possible future extractions of α_s from $\Gamma_1^{p-n}(Q^2)$

- JLab: approved experiment EG12 using CLAS12 with 11 GeV. Expected precision similar to EG1dvcs, but with Q^2 up to 6 GeV². Low-x issue reduced.
- Electron Ion Collider (EIC)

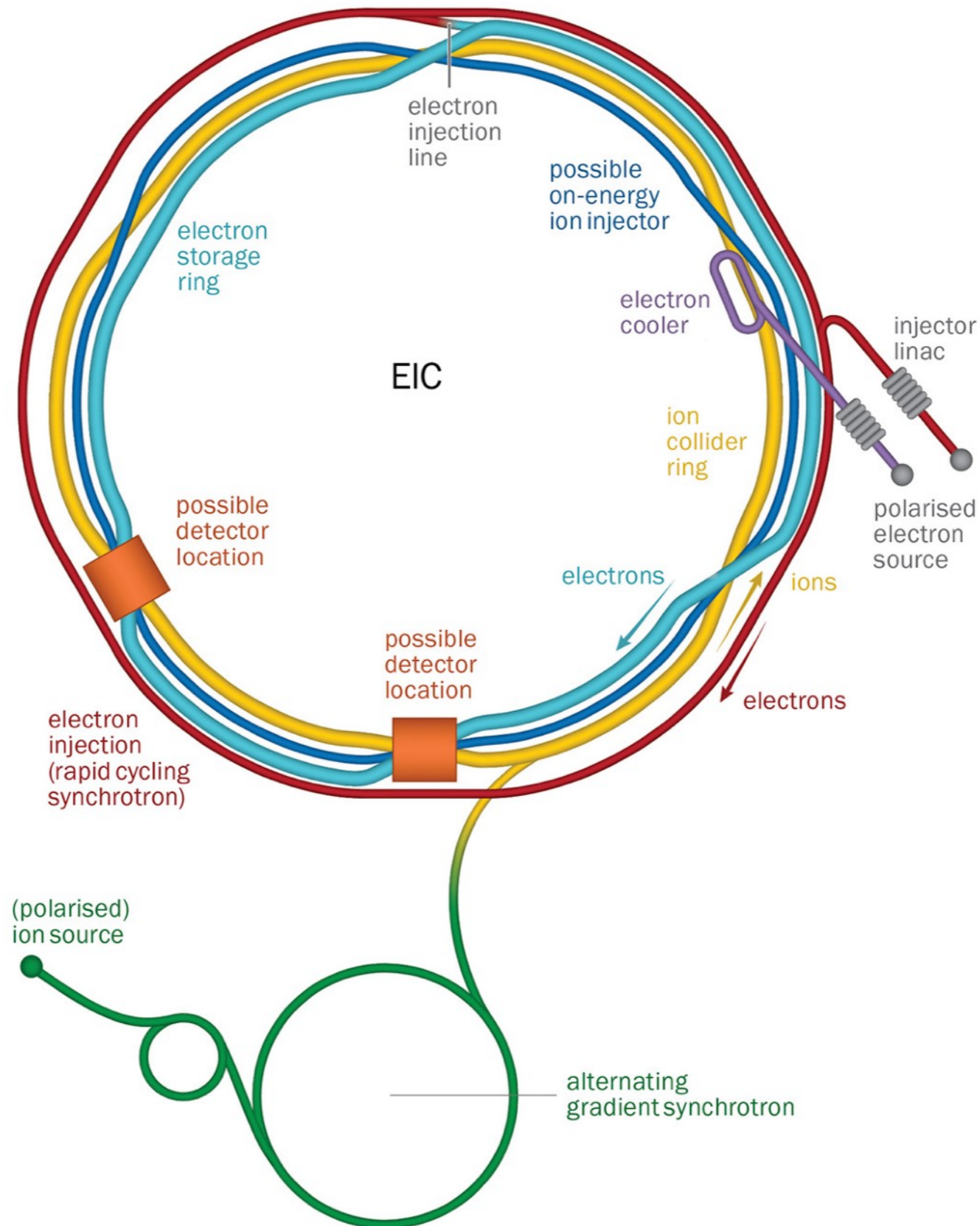
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α_s from $\Gamma_1^{p-n}(Q^2)$ measured at the EIC

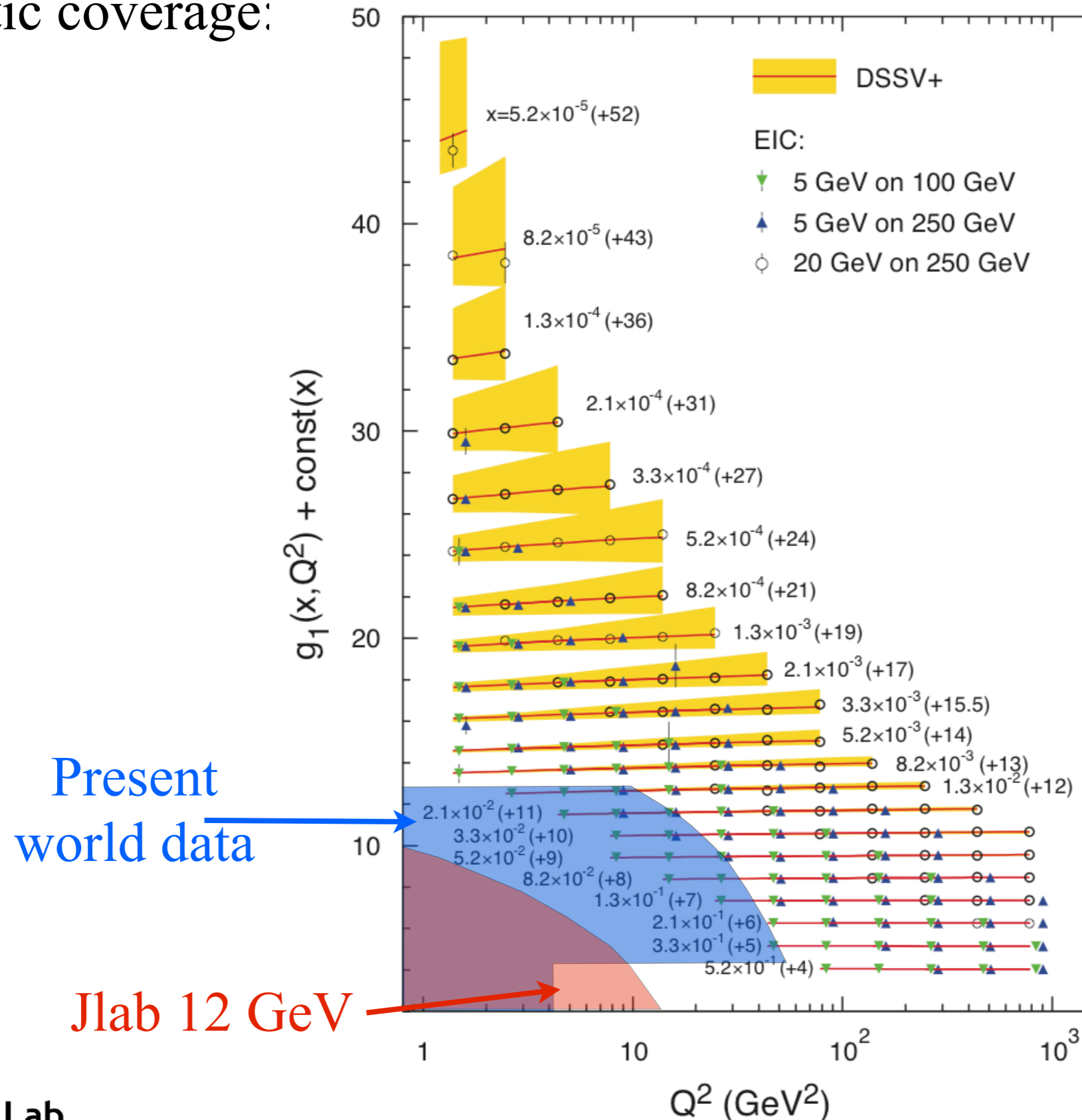
(just a quick exploratory look to see if it is worth pursuing)

2030s:



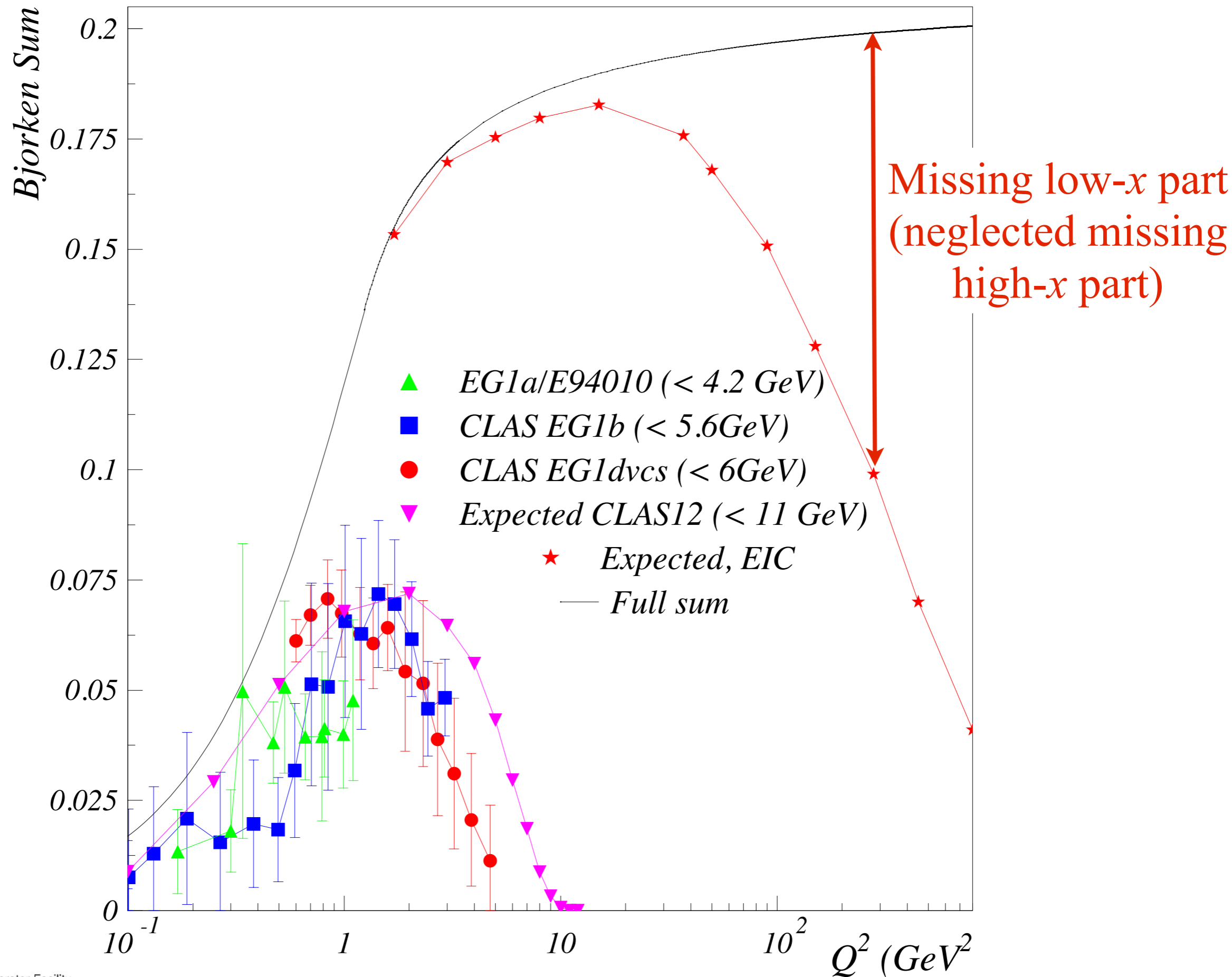
EIC

Use the 5 GeV on 100 GeV, 5 GeV on 250 GeV and 20 GeV on 250 GeV kinematic coverage:

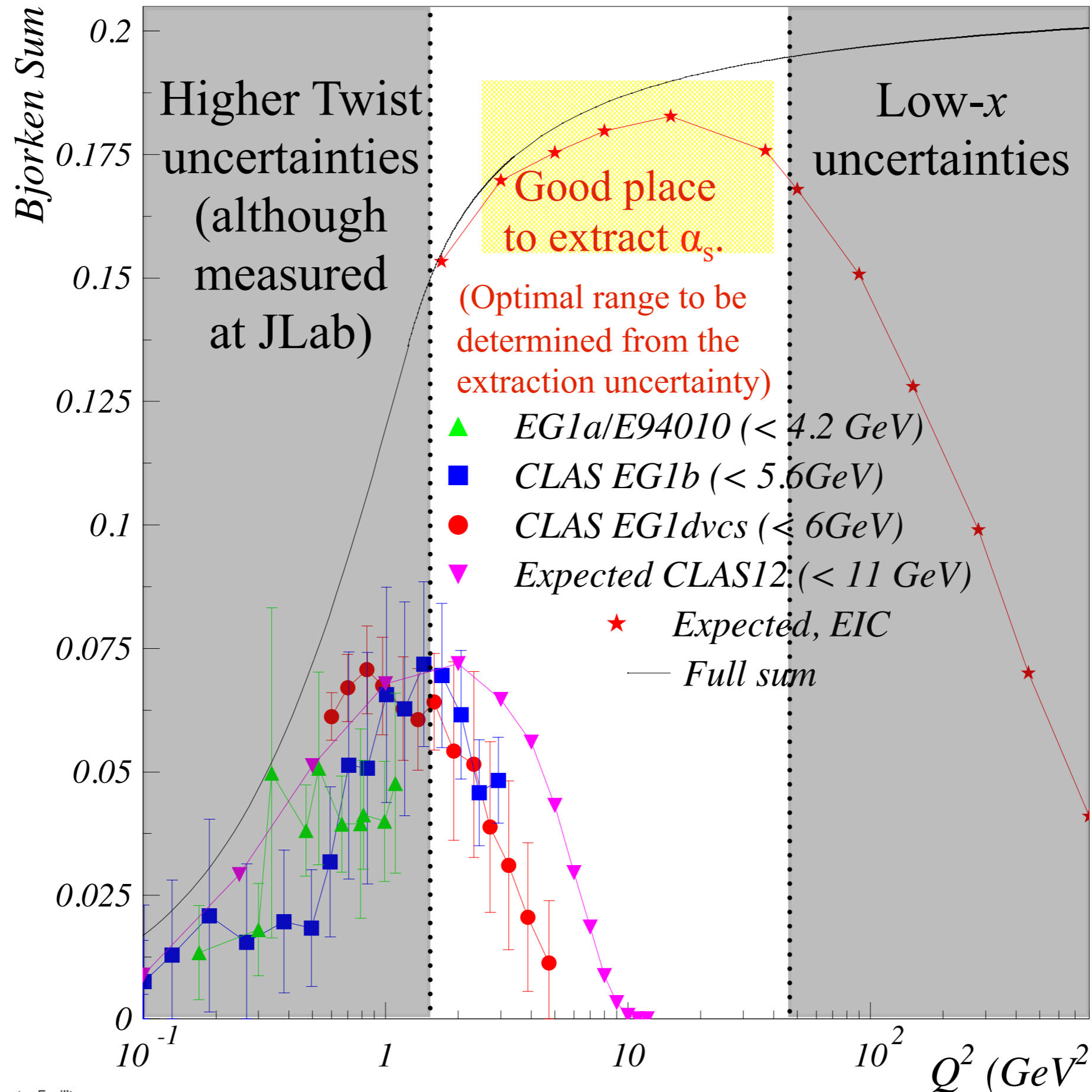


From the EIC white paper,
arXiv:1212.1701

Measured fraction of the Bjorken sum



Measured fraction of the Bjorken sum



Uncertainty budget

Statistics:

- Assume $\Delta\Gamma_1^{p-n} = 0.5\%$ ($Q^2 = 3 \text{ GeV}^2$) to $\Delta\Gamma_1^{p-n} = 0.05\%$ ($Q^2 = 15 \text{ GeV}^2$), not counting other world data (JLab@6&12 GeV, SLAC, CERN, DESY)

EIC:

- Luminosity: $2 \times 10^{33}/\text{s}$,
- $P_e P_N$: 0.5-0.6
- Dilution factor: none
- Duration: a year.
- Q^2 range for α_s fit: $1.5 < Q^2 < 15 \text{ GeV}^2$

Statistics assumed twice better than those of **CLAS EG1b experiment:**

- Luminosity: $10^{34}/\text{s}$,
- $P_b P_t$: 0.2-0.6
- Dilution factor: $\sim 80\%$
- Duration: a few months.
- Q^2 range for α_s fit: $1 < Q^2 < 3 \text{ GeV}^2$

Uncertainty budget

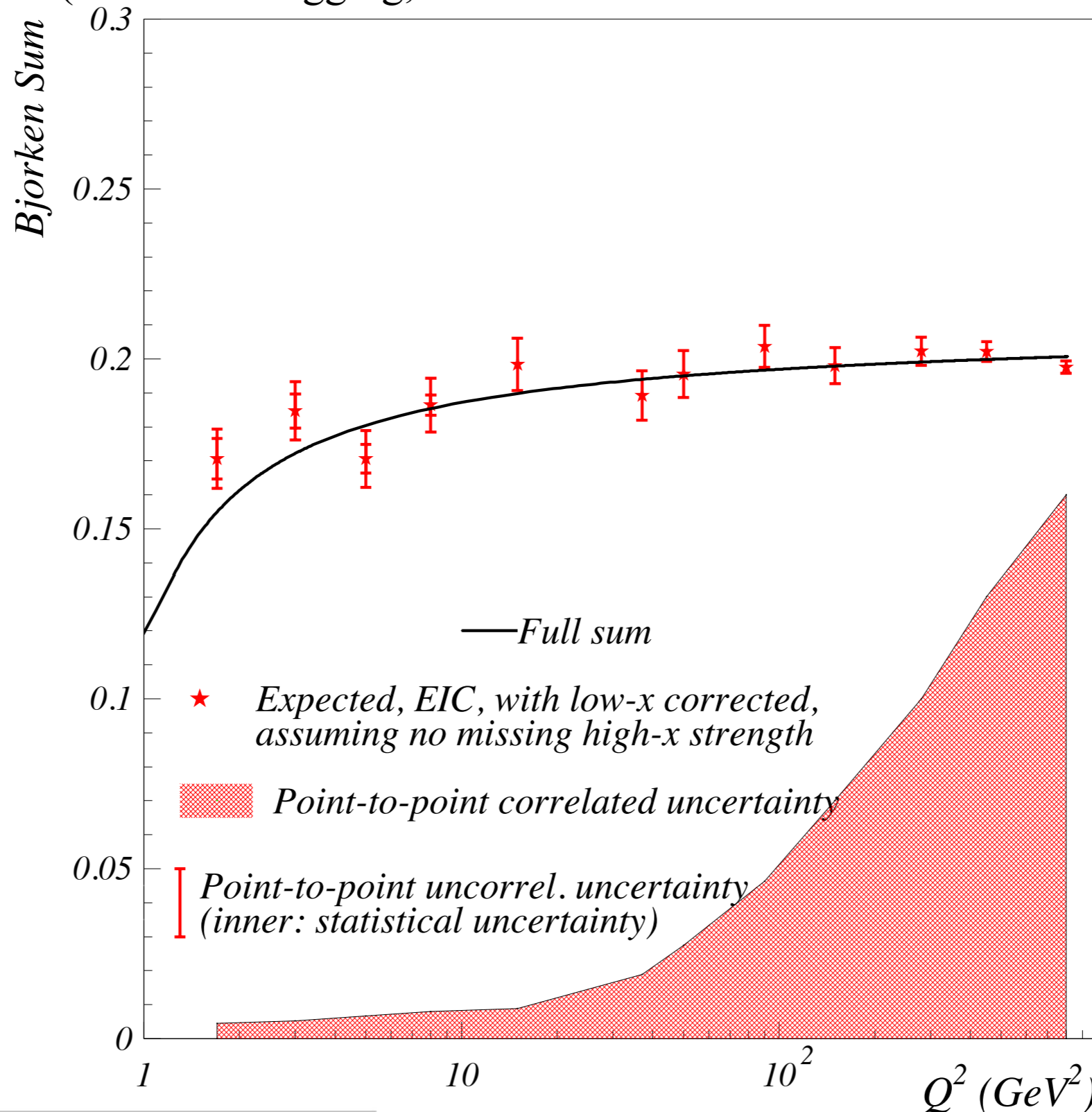
Assumed systematics:

- **Nuclear corrections (D):** 4% or negligible if we can tag the \sim spectator proton.
- **Missing low- x part:** Assume 100% uncertainty on it.
- **Polarimetries:** Assume $\Delta P_e - \Delta P_N = 3\%$.
- **Radiative corrections:** Lower energy data exist. Assume 6%.
- **F_1 to form g_1 from A_1 :** 2.5% (assumed F_2 : 2% for proton and neutron. R : 10%.)
- **g_2 contribution to longitudinal asym:** Measured with transverse pol. ion beam.
- **Dilution/purity:** 0
- **Contamination from particle miss-identification:** Assumed negligible.
- **Detector/trigger efficiencies, acceptance, beam currents:** Neglected (asym).

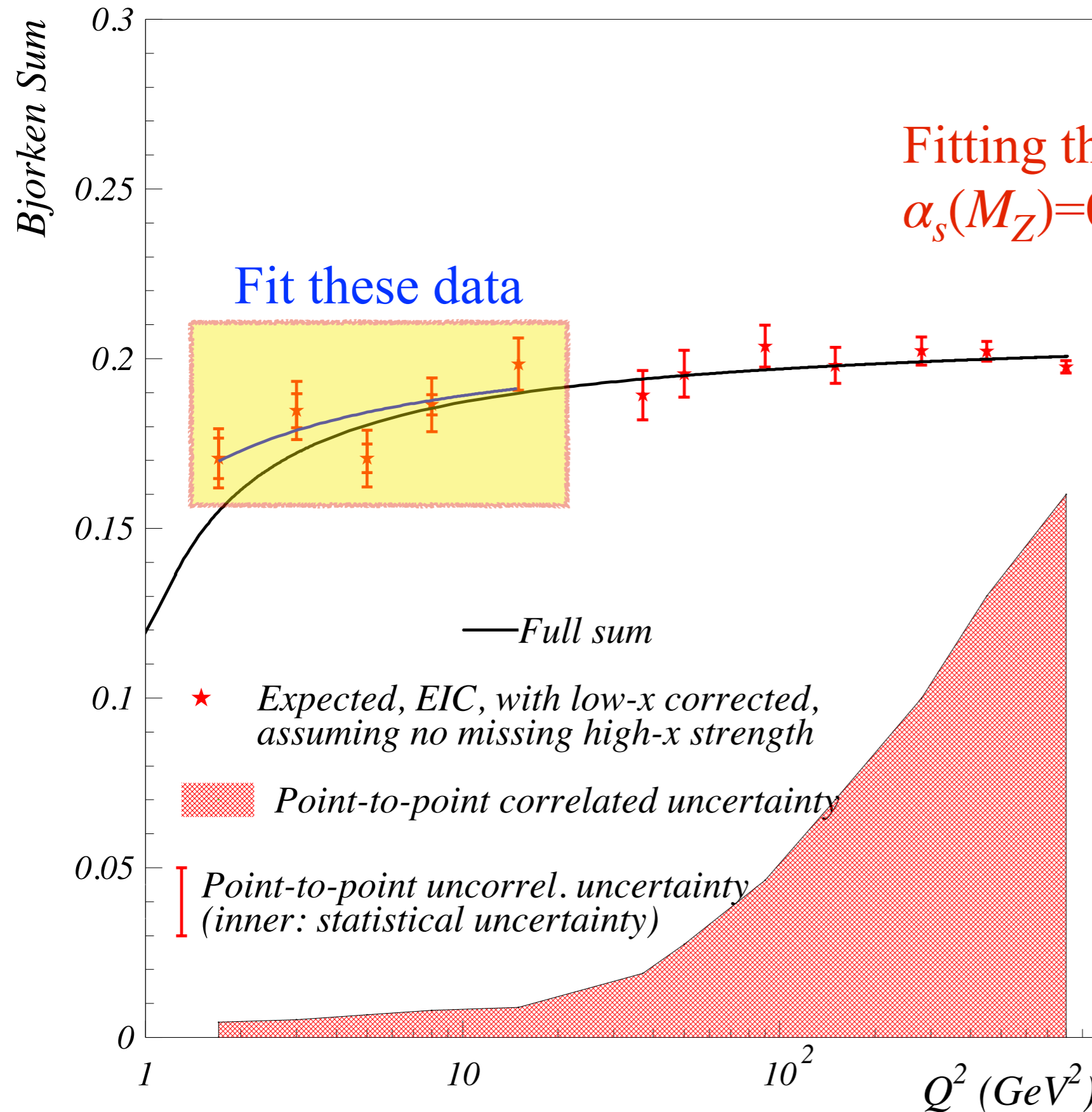
Extraction of $\alpha_s(M_Z)$

For CLAS EG1dvcs, 60% of syst. is point-to-point uncorrelated (excluding the low- x error) \Rightarrow add to stat. uncert.

\Rightarrow data may look like (assume no tagging, i.e. include nuclear correction uncertainty):



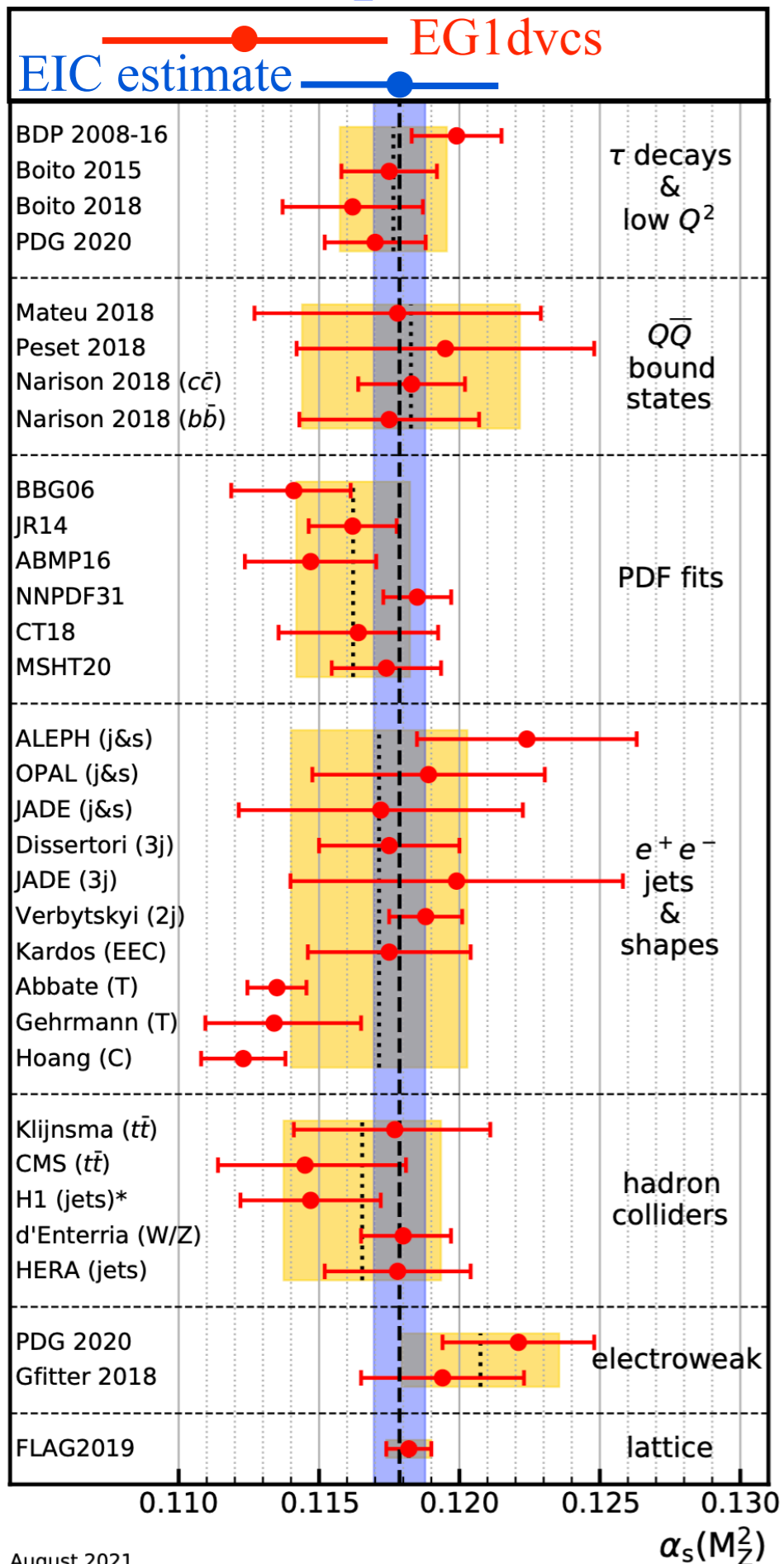
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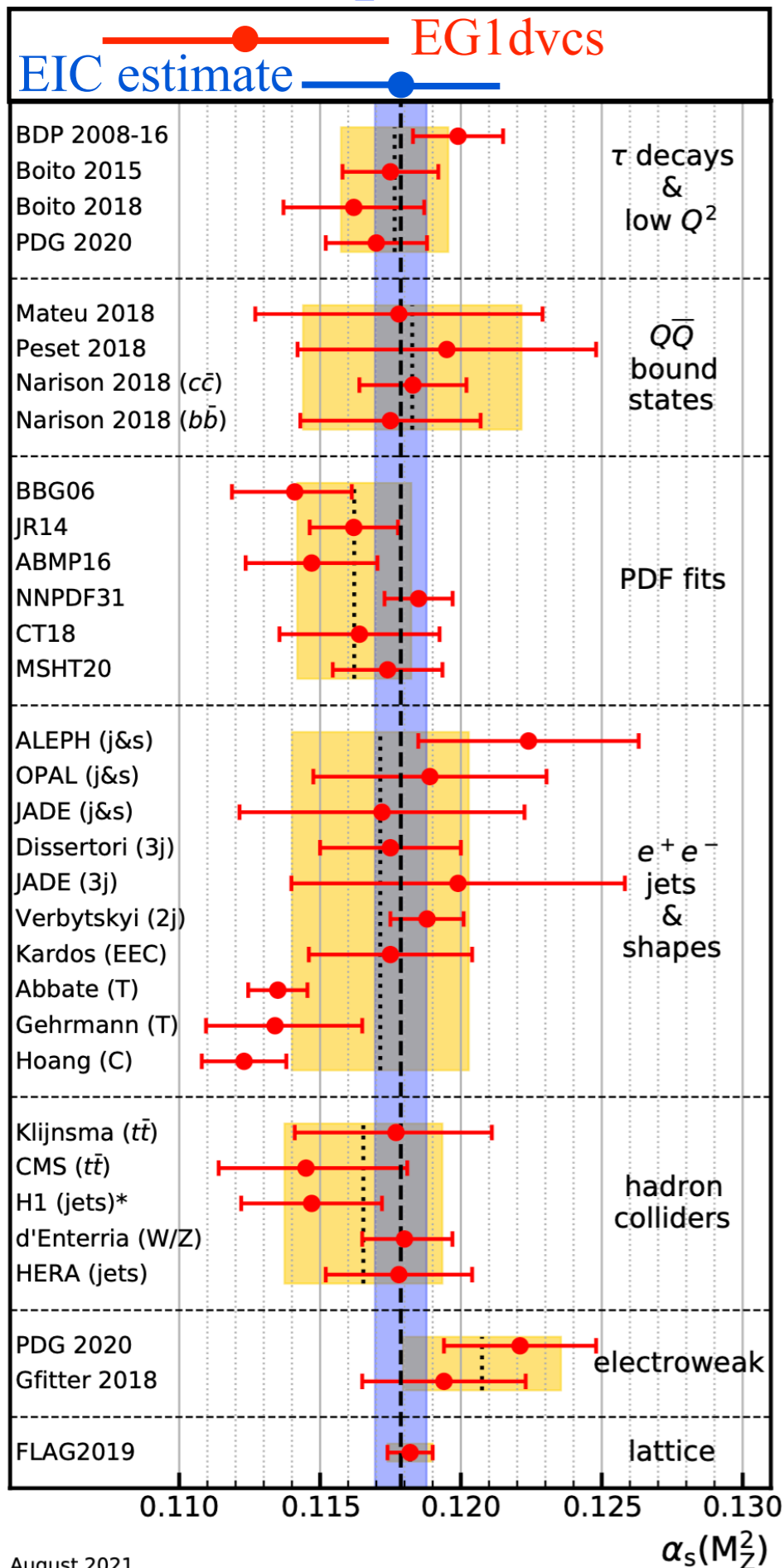
Fitting the data yields:
 $\alpha_s(M_Z) = 0.1175 \pm 0.0033 \pm 0.0005$

(Adding next data point yields:
 $\alpha_s(M_Z) = 0.1162 \pm 0.0030 \pm 0.0021$
 \Rightarrow Not worth including more point).

Compared to EG1dvcs and best world data (PDG 2021):



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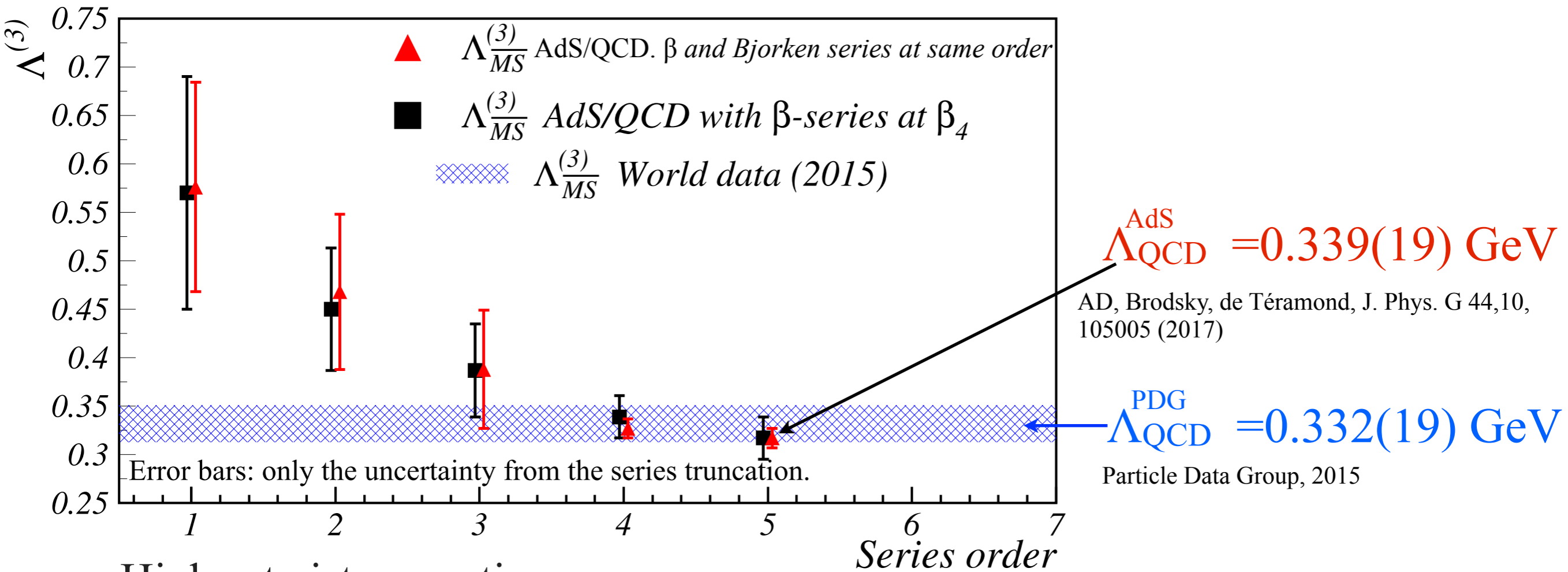
Conclusion:

- Under reasonable assumptions, EIC can yield an accurate measurement of moderate precision.
- Assumed statistics similar to a typical CLAS experiment aiming at measuring inclusive spin structure functions.
- Increasing statistics by factor 10 would yield: $\Delta\alpha_s(M_Z) = \pm 0.0021 \pm 0.0003$.
- If proton tagging is added with D data taking: $\Delta\alpha_s(M_Z) = \pm 0.0016 \pm 0.0003$. **Competitive measurement.**

Summary and perspective

- We discussed the extraction of $\alpha_s(M_Z)$ from the Bjorken sum $\Gamma_1^{p-n}(Q^2) = \int g_1^{p-n}(x, Q^2) dx$.
- The JLab 6 GeV experiment EG1dvcs yielded $\alpha_s(M_Z) = 0.1123 \pm 0.0061$ (PRD, 2014). On the low side, but agreeing with the PDG-2021 world average.
- EG12: JLab CLAS12 with 11 GeV beam. Expected precision similar to EG1dvcs, but better accuracy and with Q^2 up to 6 GeV².
- EIC: exploratory estimate suggests an accurate measurement of acceptable precision, even without dedicated experiment \Rightarrow it seems worth to seriously investigate this possibility.
- Not discussed here:
 - Improving extraction from $\Gamma_1^{p-n}(Q^2)$ with lower Q^2 points, using either
 - Higher-twist corrections.
 - Non-perturbative *model* of $\alpha_s(Q^2)$. Matching between AdS/QCD and pQCD in their overlap domain has yield competitive estimate.
 - $\alpha_s(M_Z)$ from Q^2 -evolution of
 - $g_1(x, Q^2)$;
 - Proton's $\int g_1^p(x, Q^2) dx$.

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