

# **Thermal Misalignment of Scalar Dark Matter**

**Generating scalar misalignment through Planck suppressed  
interactions**

**Akshay Ghalsasi, Brian Batell, arxiv:2109.04476**



# Roller Coaster Dark Matter

Generating scalar misalignment through Planck suppressed interactions

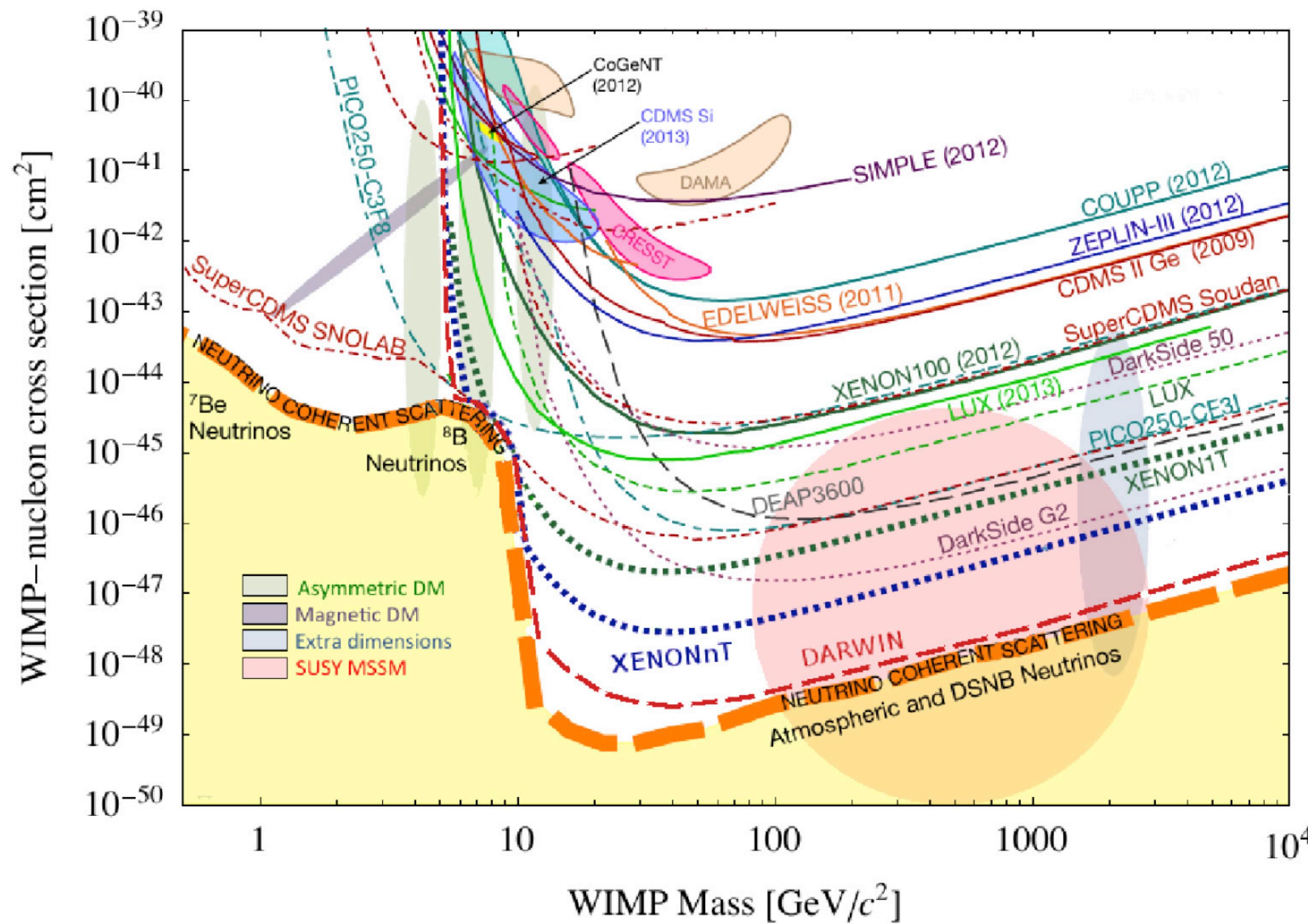
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# Overview

- Introduction to the model
- Visualizing Potential and Analytical Understanding
- Parameter space exploration
- Constraints
- Conclusions and Future Work

# Motivation

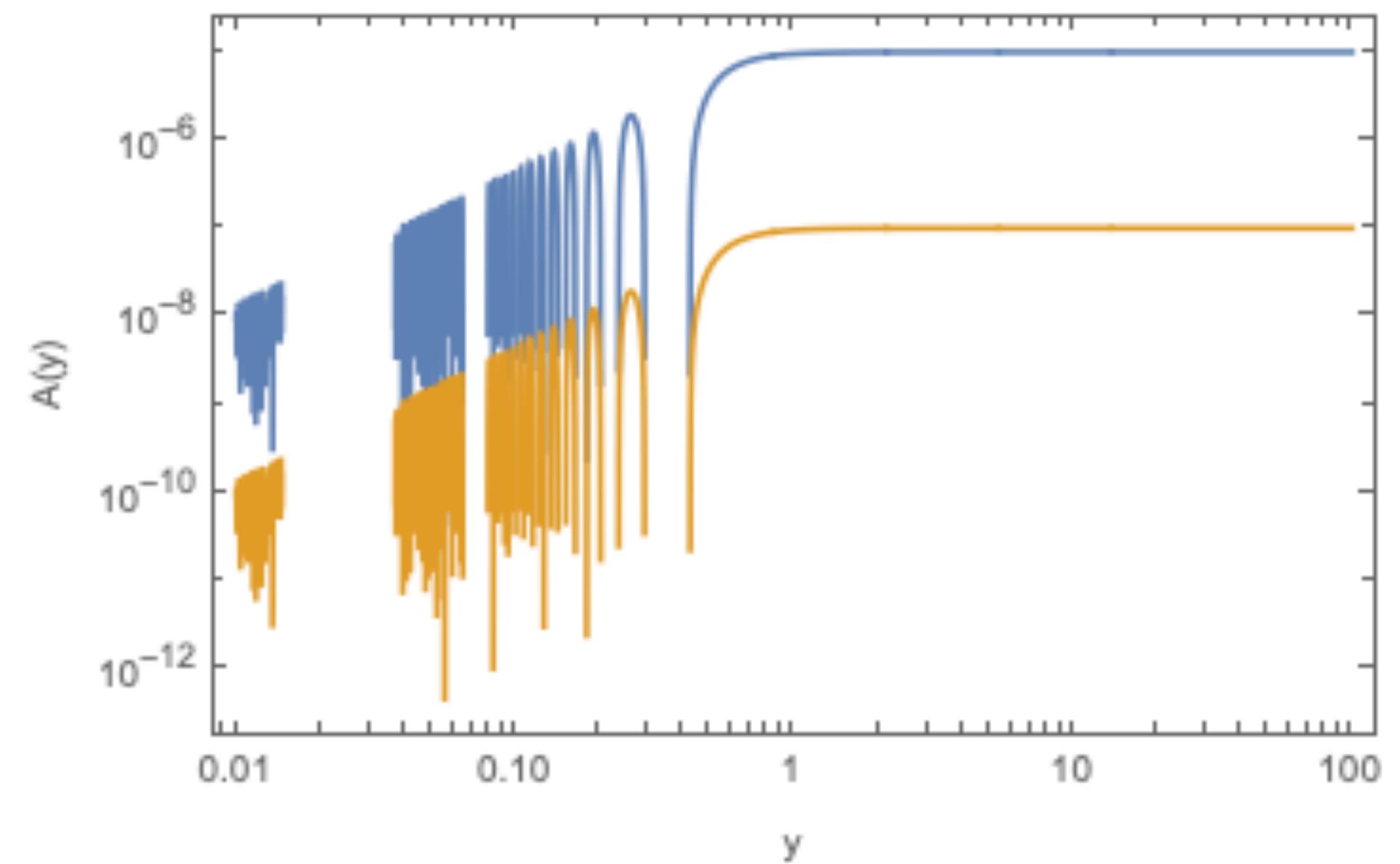
## WIMPs under pressure



# Motivation

## Standard misalignment

- Coherent oscillations of a scalar field can constitute dark matter
- In standard scenario
  - Inflation generates misalignment
  - Scalar oscillates when Hubble equals mass
  - Final yield sensitive to initial conditions
- Worst case scenario - no coupling to SM, impossible to detect



# Introduction to the Model

## The Lagrangian

- Scalar couples to fermions through Planck suppressed 5-dim operator

$$\mathcal{L} \supset - \left[ m_f \left( 1 - \frac{\beta\phi}{M_{\text{pl}}} \right) f_L f_R - \text{h.c.} \right] + \frac{1}{2} m_\phi^2 \phi^2$$

- Large  $\beta$  corresponds to a UV scale smaller than  $M_{\text{pl}}$  (through a Higgs portal)
- Same Lagrangian explored in Weiner and Zurek (2006) for other purposes

# Introduction to the Model

## The Potential

- Free energy of fermions gives rise to effective potential on the scalar

$$\delta V(\phi) = \frac{-g}{2\pi^2} T^4 \int_0^\infty dx x^2 \log \left[ 1 + e^{-\sqrt{x^2 + \frac{1}{y^2}} \left( 1 - \frac{\beta\phi}{M_{\text{pl}}} \right)^2} \right]$$

- Standard equation of motion

$$\ddot{\phi} + 3H\dot{\phi} + \frac{\partial V}{\partial \phi} = 0 \quad \xrightarrow{\hspace{1cm}} \quad \frac{d^2\phi}{dy^2} + \frac{M_{\text{pl}}^2}{m_f^4 \gamma^2 y^6} \frac{\partial V}{\partial \phi} = 0$$

$$y = \frac{T}{m_f}$$

$$\downarrow \quad \kappa = \frac{m_\phi M_{\text{pl}}}{m_f^2}$$

$$\frac{d^2\phi}{dy^2} + \frac{\kappa^2}{\gamma^2 y^6} \left( \phi - \frac{2}{\pi^2} \frac{y^2 M_{\text{pl}}}{\kappa^2} \beta \left( 1 - \frac{\beta\phi}{M_{\text{pl}}} \right) \int_0^\infty dx \frac{x^2}{(1 + e^\xi) \xi} \right) = 0$$

$$\xi = \sqrt{x^2 + \frac{1}{y^2} \left( 1 - \frac{\beta\phi}{M_{\text{pl}}} \right)^2}$$

# Analytical Understanding

- High Temperature regime ( $T \gg m_\psi, y \gg 1$ )

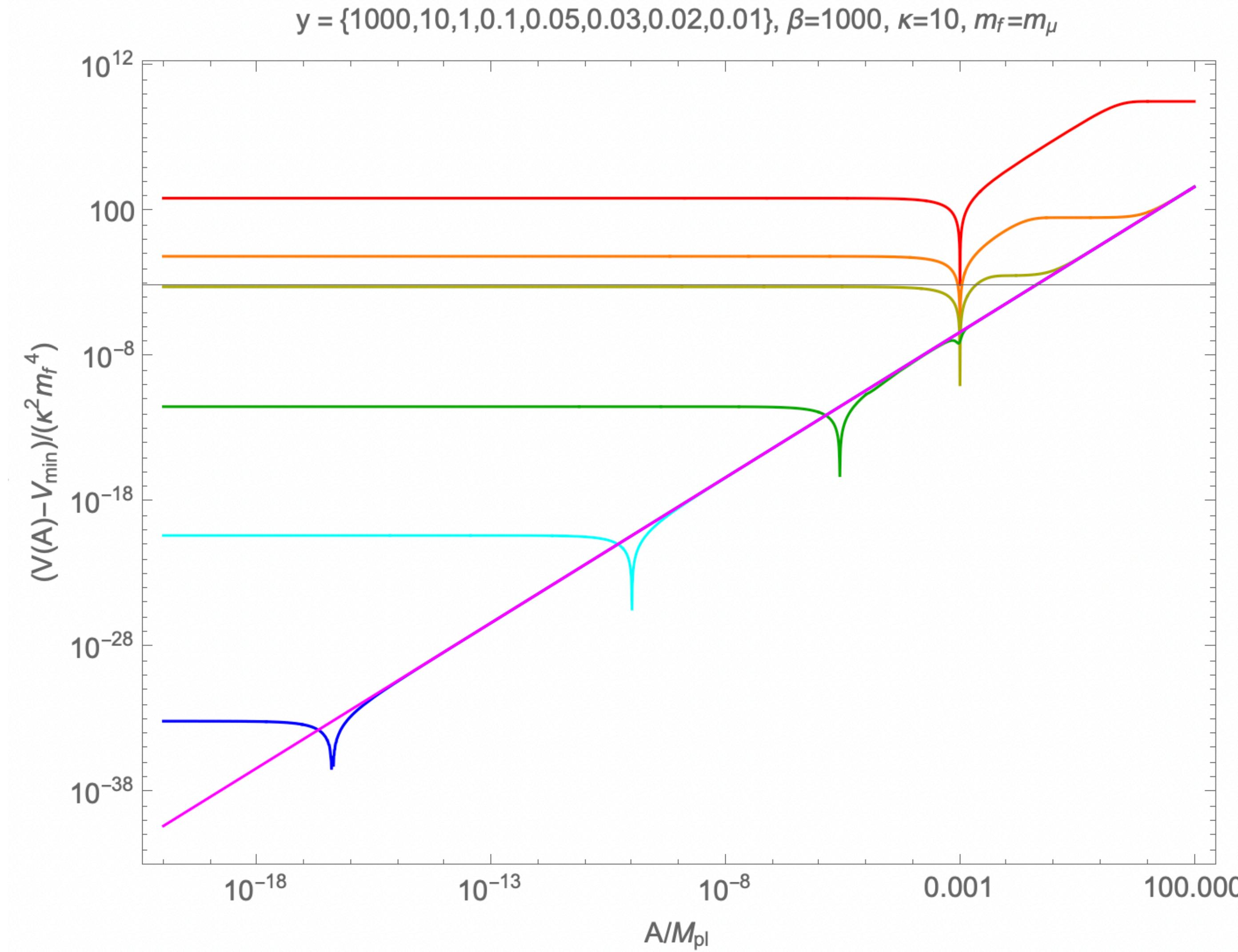
$$\delta V_\phi = \frac{-g}{2\pi^2} T^4 \int_0^\infty dx x^2 \log \left[ 1 + e^{-\sqrt{x^2 + \frac{1}{y^2} \left( 1 - \frac{\beta\phi}{M_{\text{pl}}} \right)^2}} \right] = m_f^4 \left( \frac{y^2}{12} (1 - \beta\phi)^2 + \frac{1}{2} \kappa^2 \phi^2 \right) \rightarrow$$

$$\phi_{\min} = \frac{y^2 \beta}{y^2 \beta^2 + 6\kappa^2}$$

- Low Temperature regime ( $T \ll m_\psi, y \ll 1$ )

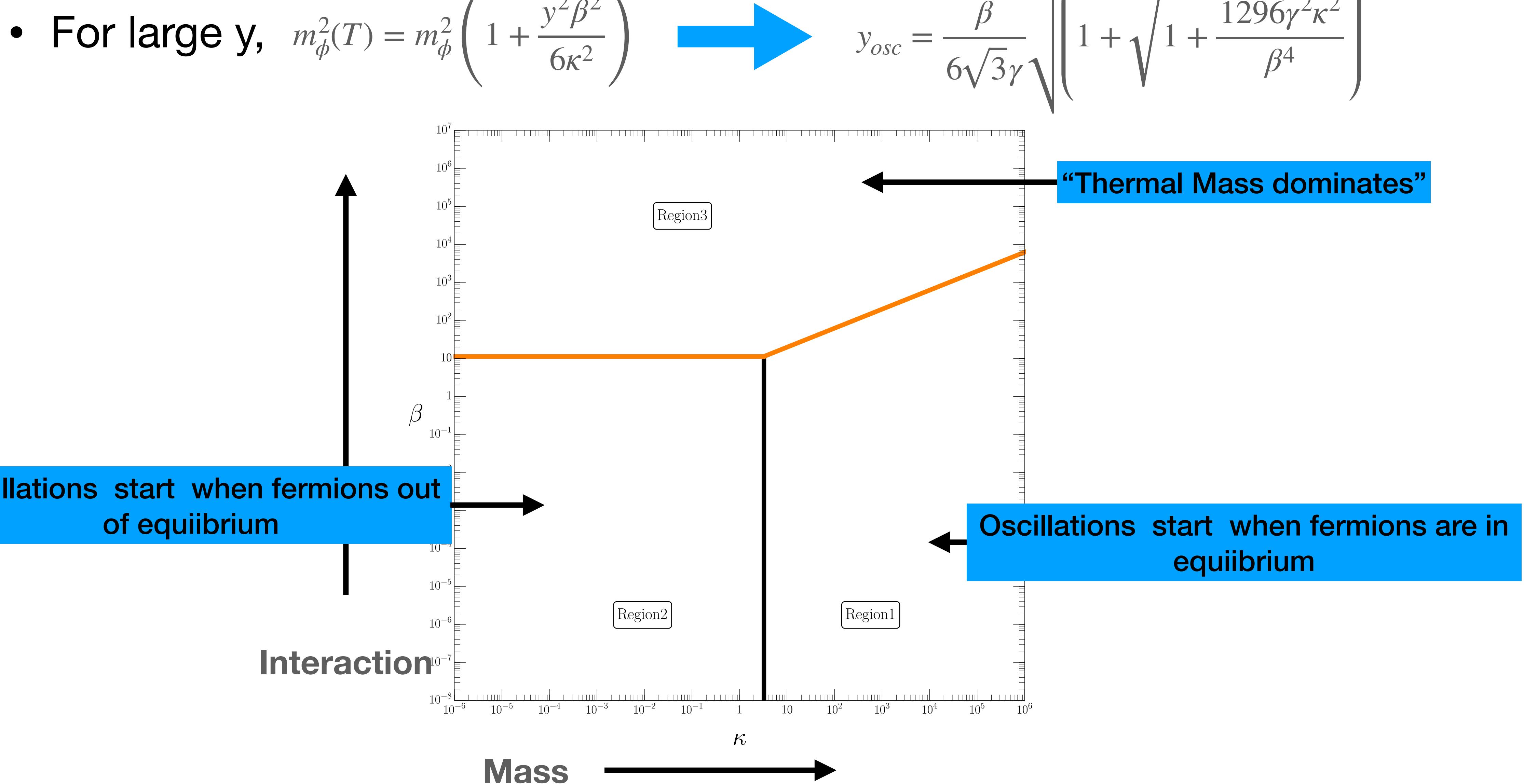
$$\delta V_\phi = m_f^4 \left( y^{5/2} e^{-1/y} + \frac{1}{2} \kappa^2 \phi^2 \right)$$

# Visualizing the Potential



# Parameter space exploration

## Overview

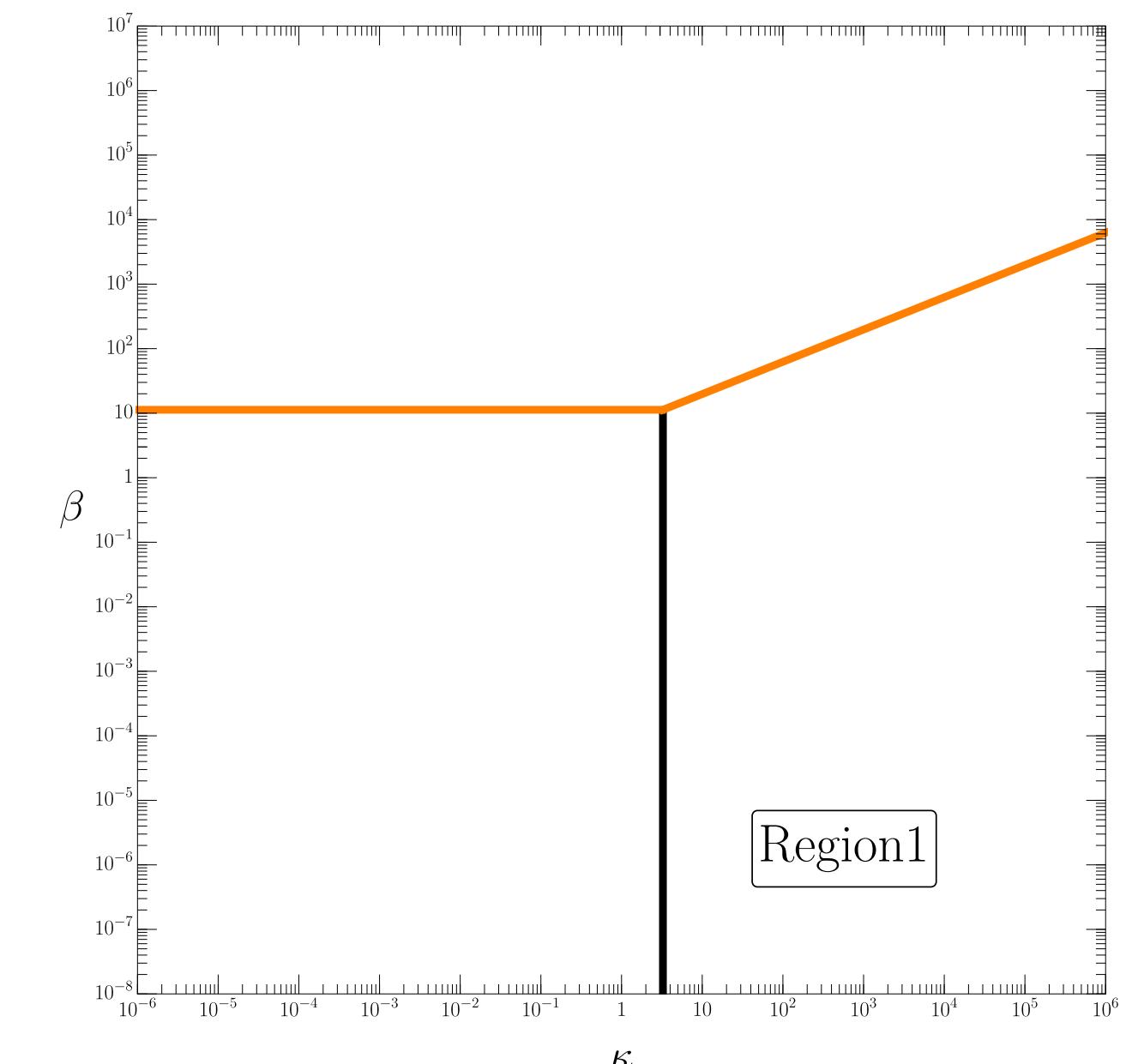


# Parameter space exploration

## Large $\kappa$ and small $\beta$ , Region 1

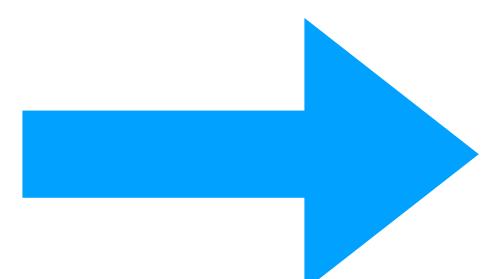
- Large  $\kappa$ , small  $\beta$ , mass of scalar initiates the oscillation

$$y_{osc} \simeq \sqrt{\frac{\kappa}{3\gamma}}$$



- However minima still controlled by  $\beta$ , since  $T \gg m$  potential applies

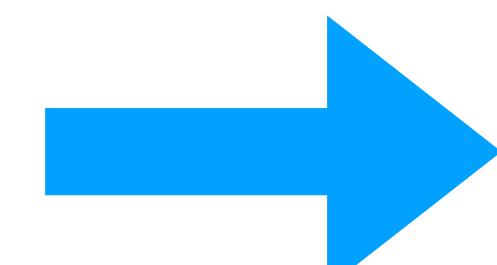
$$\ddot{\phi} + 3H\dot{\phi} + \frac{\partial V}{\partial \phi} = 0$$



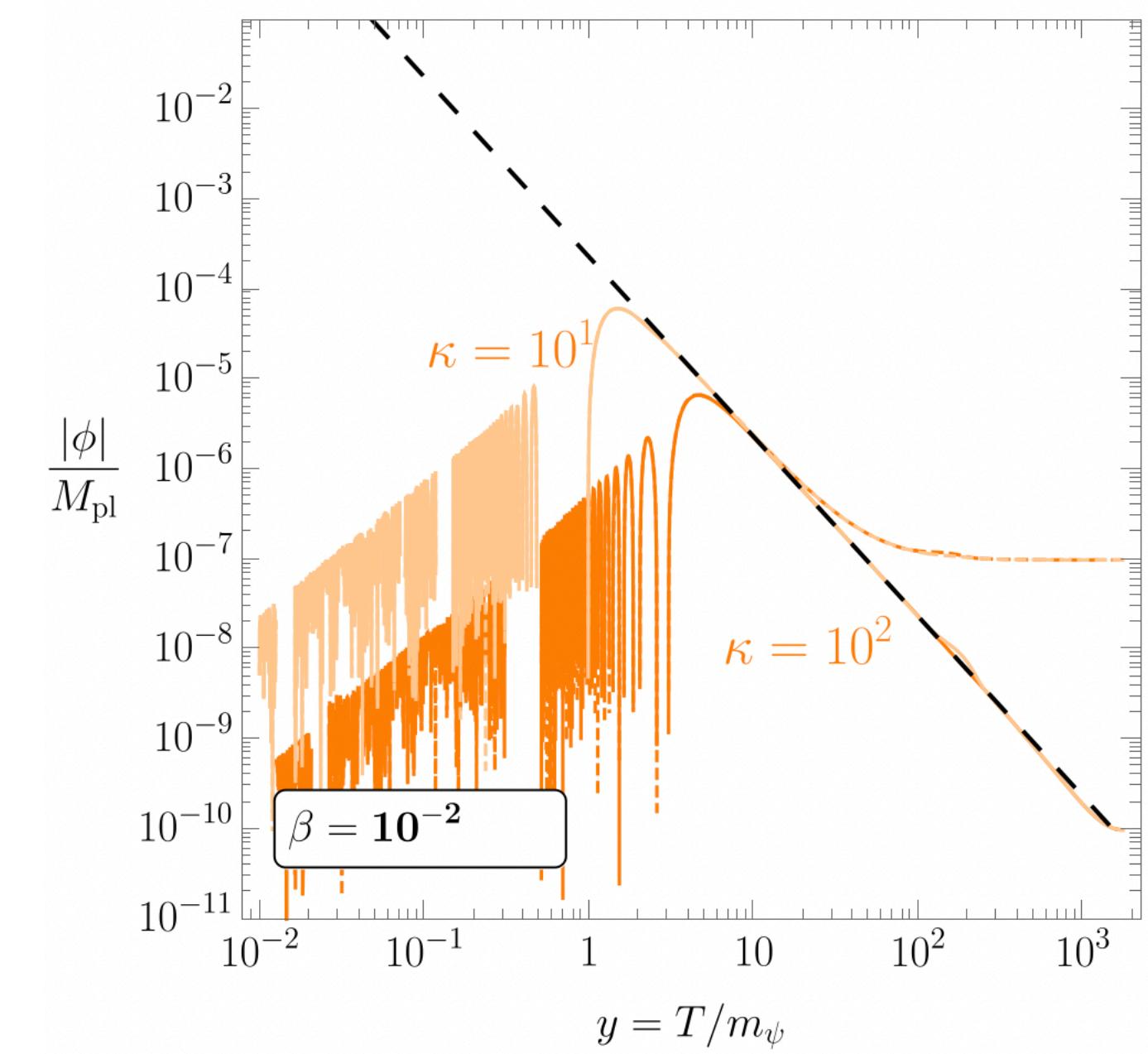
$$\dot{\phi} = \frac{\beta m_\psi^2}{18\gamma} = const.$$

- Thus it can be shown

$$\phi(y) \simeq \frac{\beta M_{pl}}{36\gamma^2 y^2}$$



$$\phi(y_{osc}) \simeq \frac{\beta}{12\gamma\kappa}$$

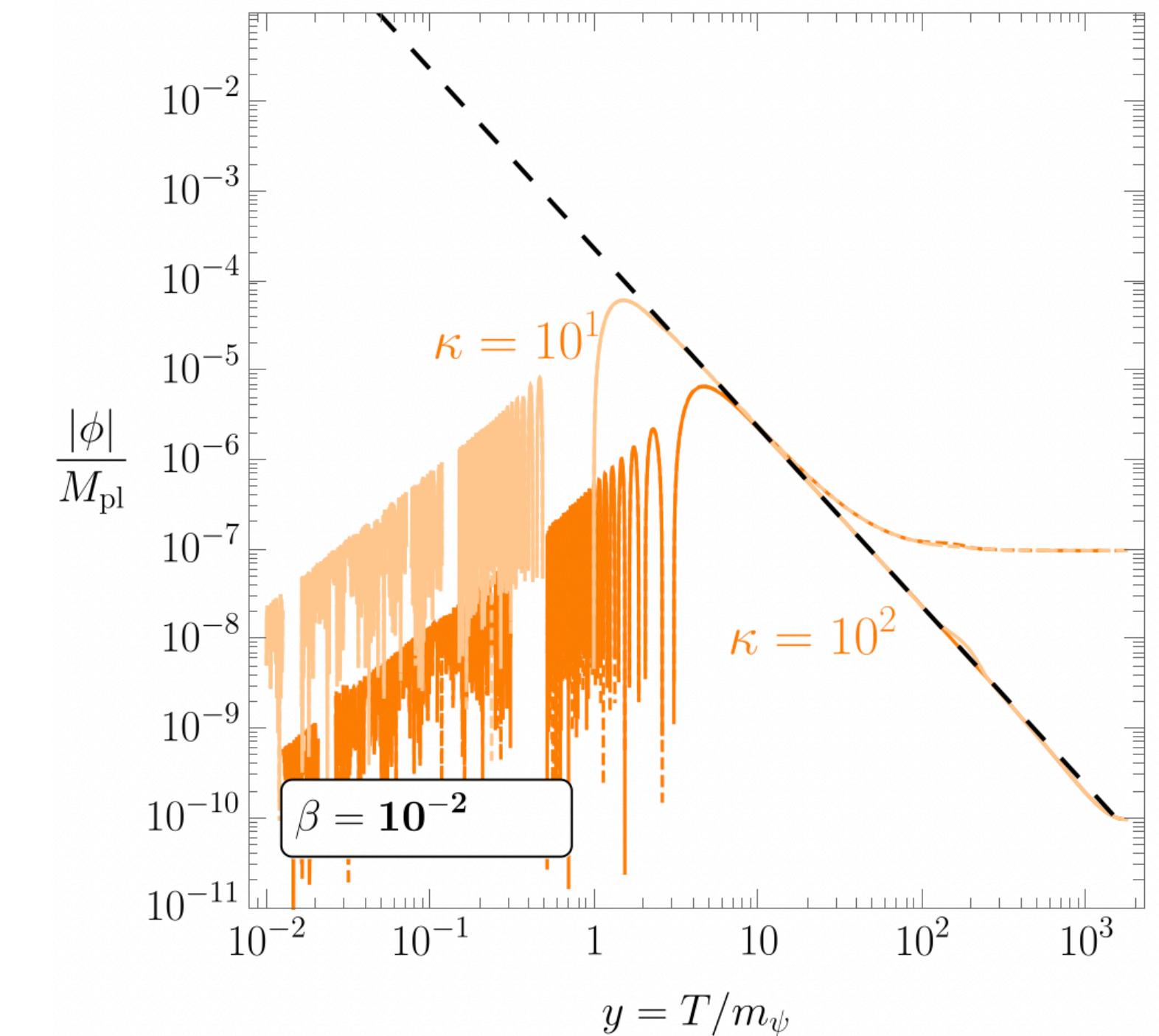


# Parameter space exploration

## Region 1, DM density

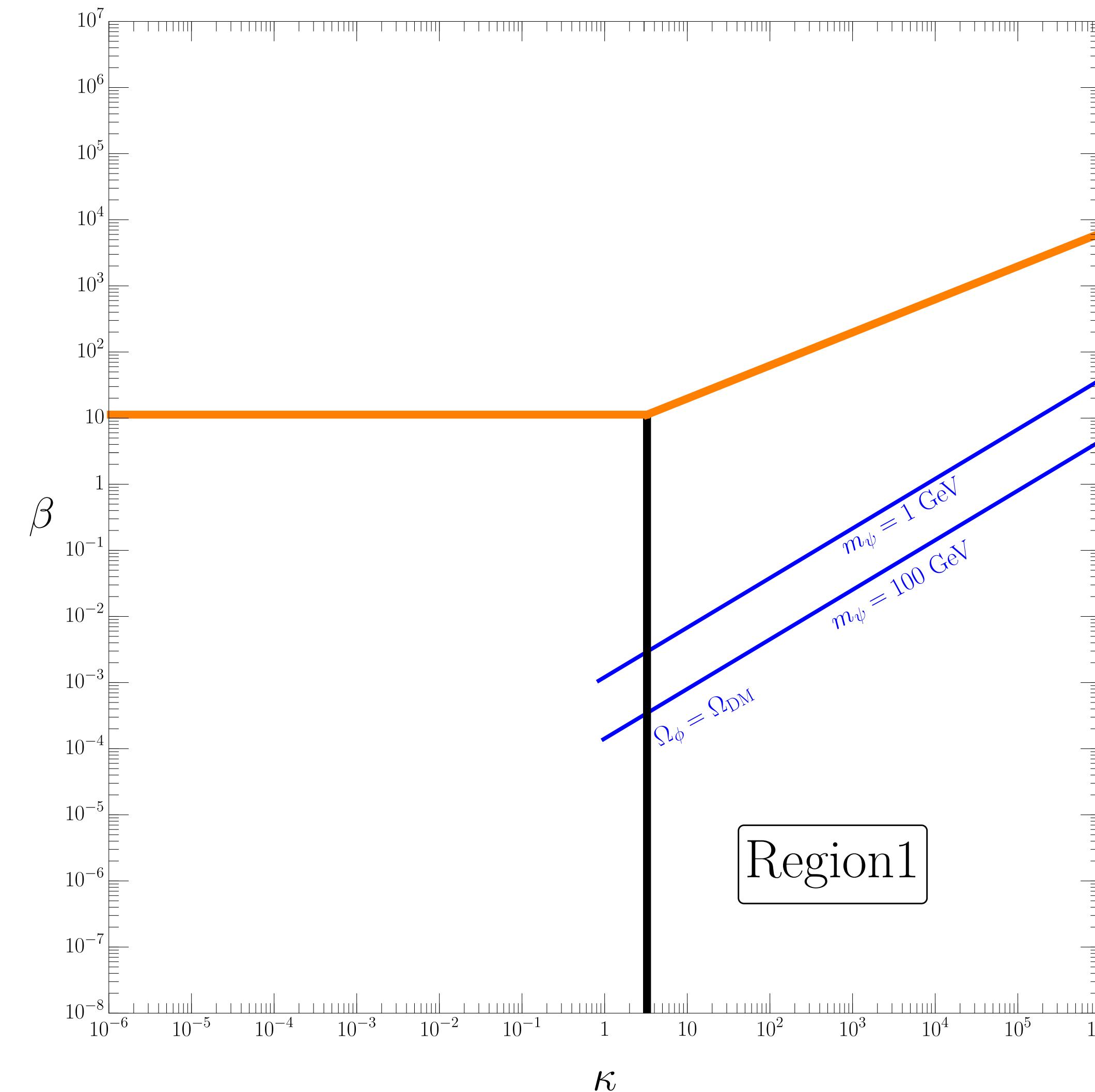
- Initial condition irrelevant as long as  $\phi(y_i) < \phi(y_{osc})$
- DM density can then be written as

$$\Omega_\phi \simeq \Omega_{\text{DM}} \left( \frac{m_\psi}{0.1 \text{ GeV}} \right) \left( \frac{\beta}{0.1} \right)^2 \left( \frac{400}{\kappa} \right)^{3/2} \left( \frac{10.75}{g_*^{\text{osc}}} \right)^{5/4}$$



# Parameter space exploration

## Region 1, Plot



# Parameter space exploration

## Small $\kappa$ and small $\beta$ , Region 2

- For  $y \gg 1$  scalar follows similar path as Region 1

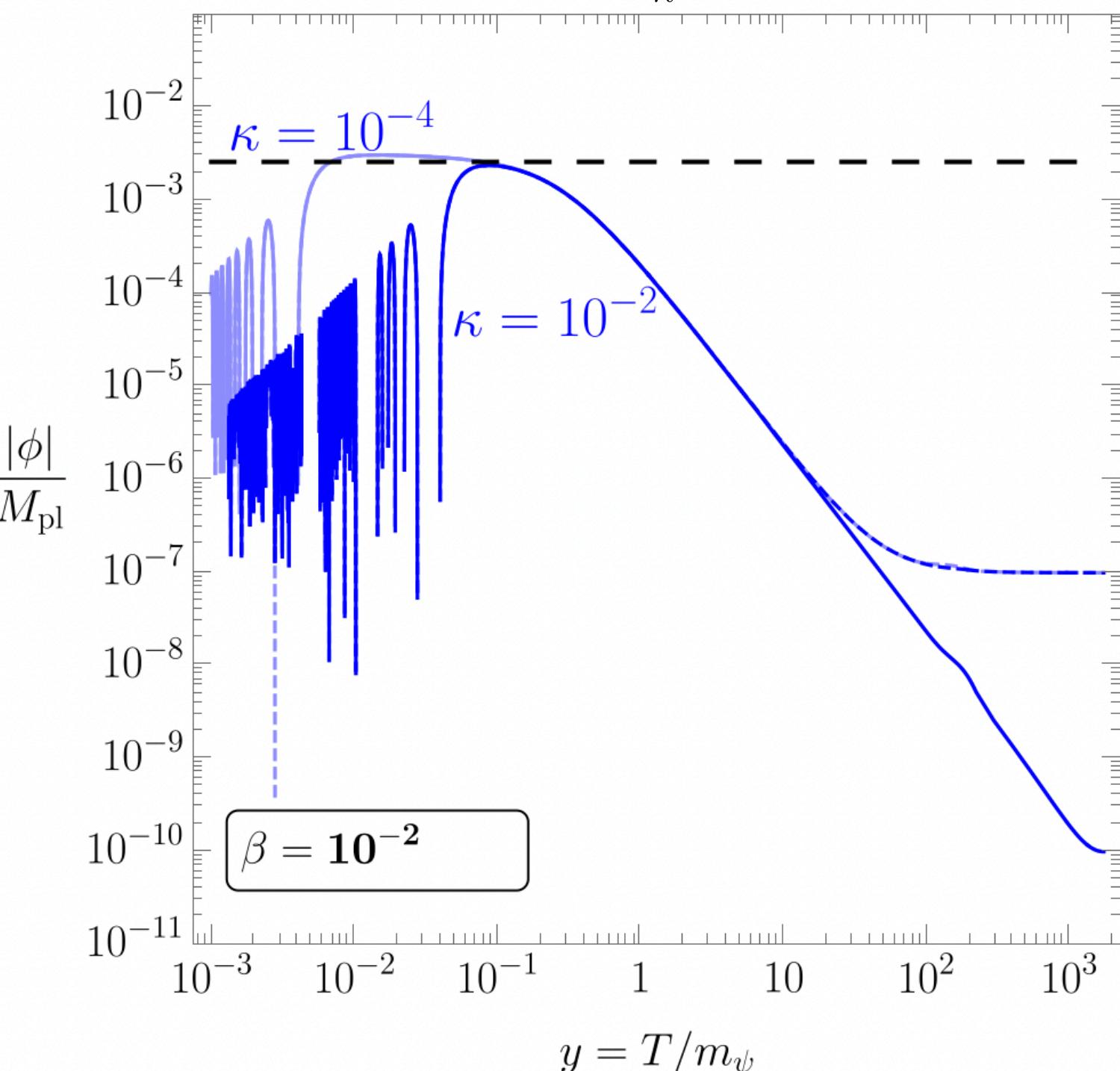
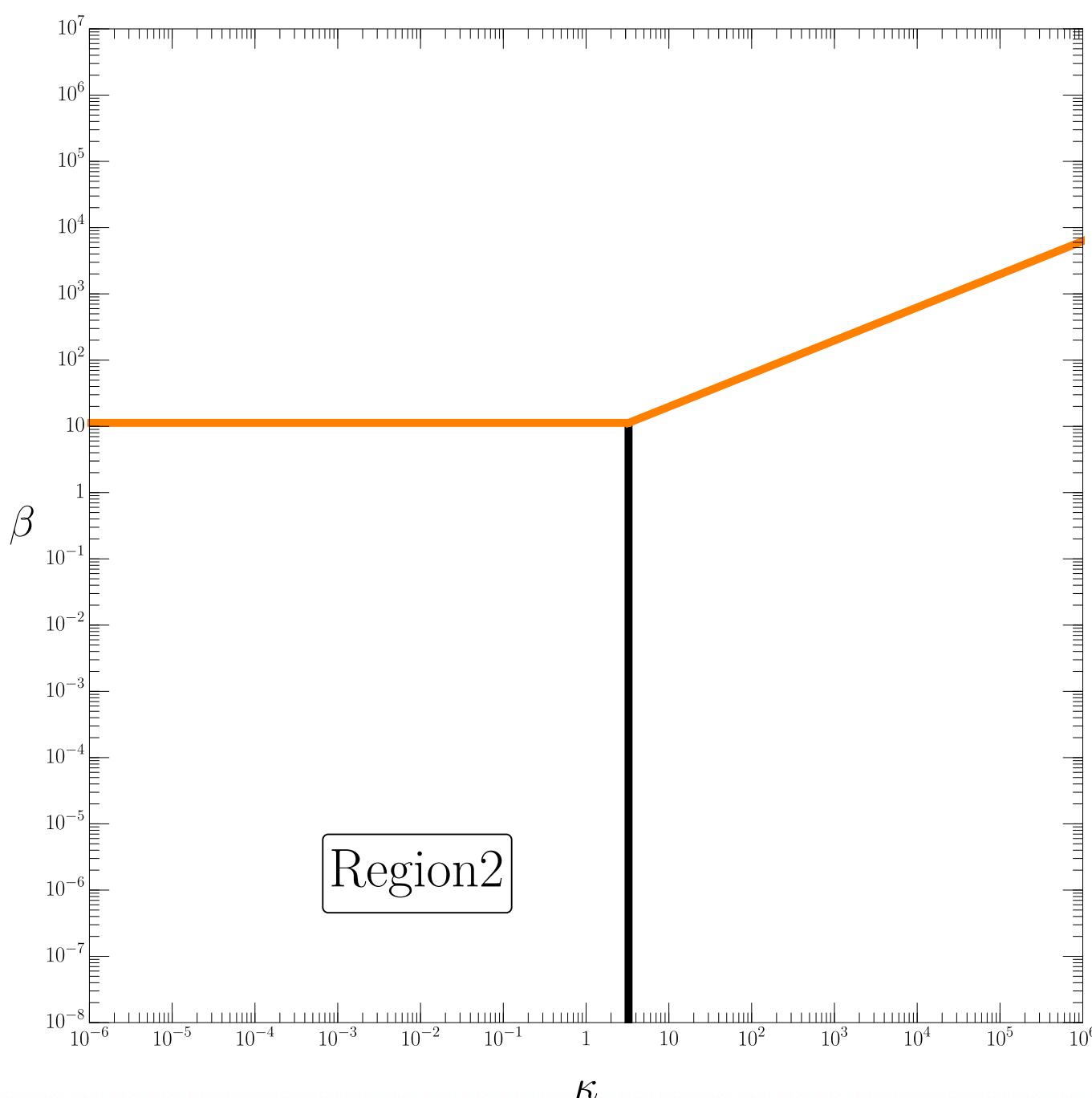
$$\phi(y) \simeq \frac{\beta M_{\text{pl}}}{36\gamma^2 y^2}$$

- For small  $\kappa$  oscillations start for  $y < 1$

$$y_{\text{osc}} \simeq \sqrt{\frac{\kappa}{3\gamma}}$$

- Fermion feedback switches off scalar asymptotes to

$$\phi(y_{\text{osc}}) \simeq \frac{0.27\beta}{\gamma^2}$$

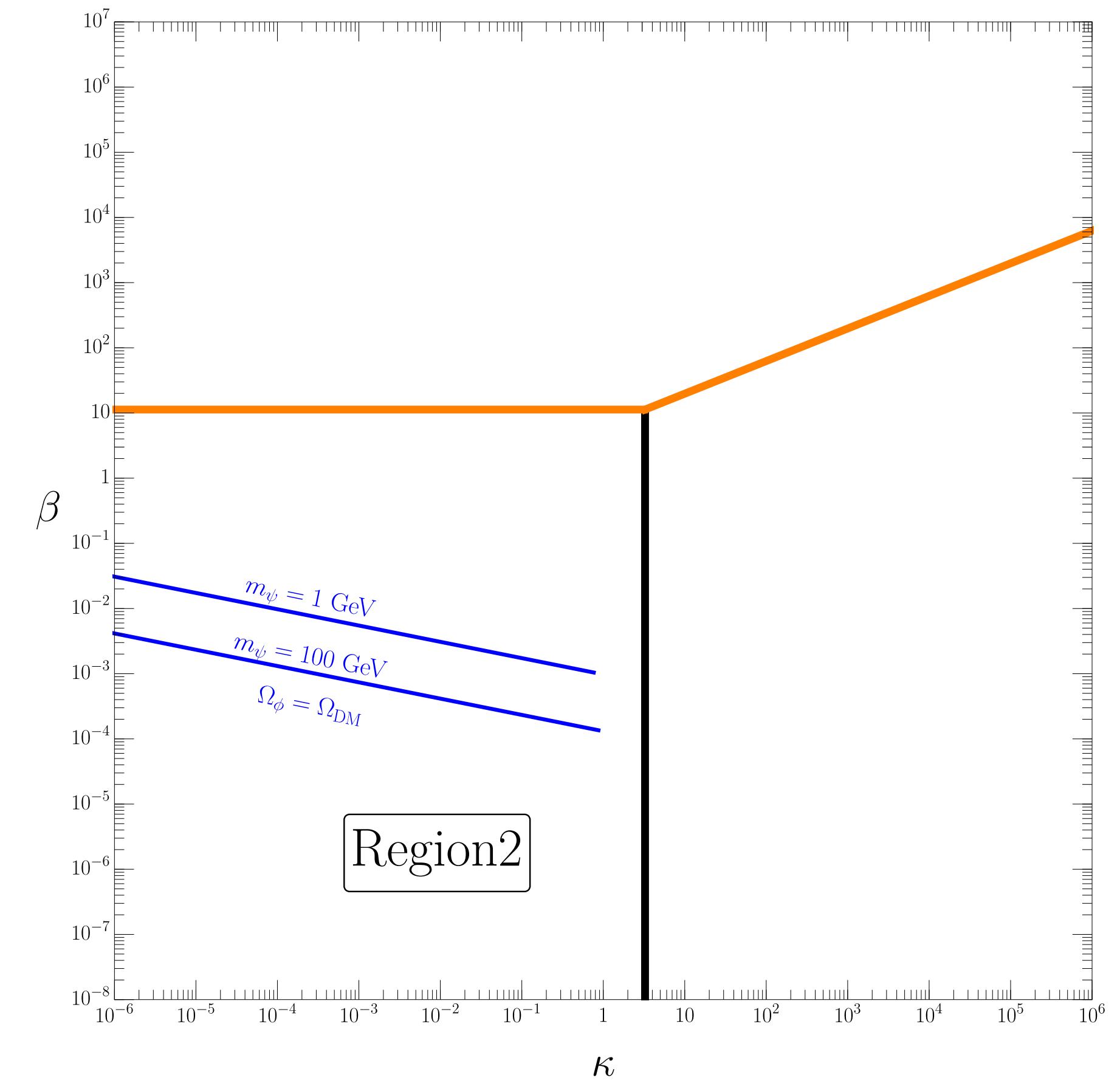


# Parameter space exploration

## Small $\kappa$ and small $\beta$ , DM density, Plots

- Thus the energy density at matter radiation equality is given by

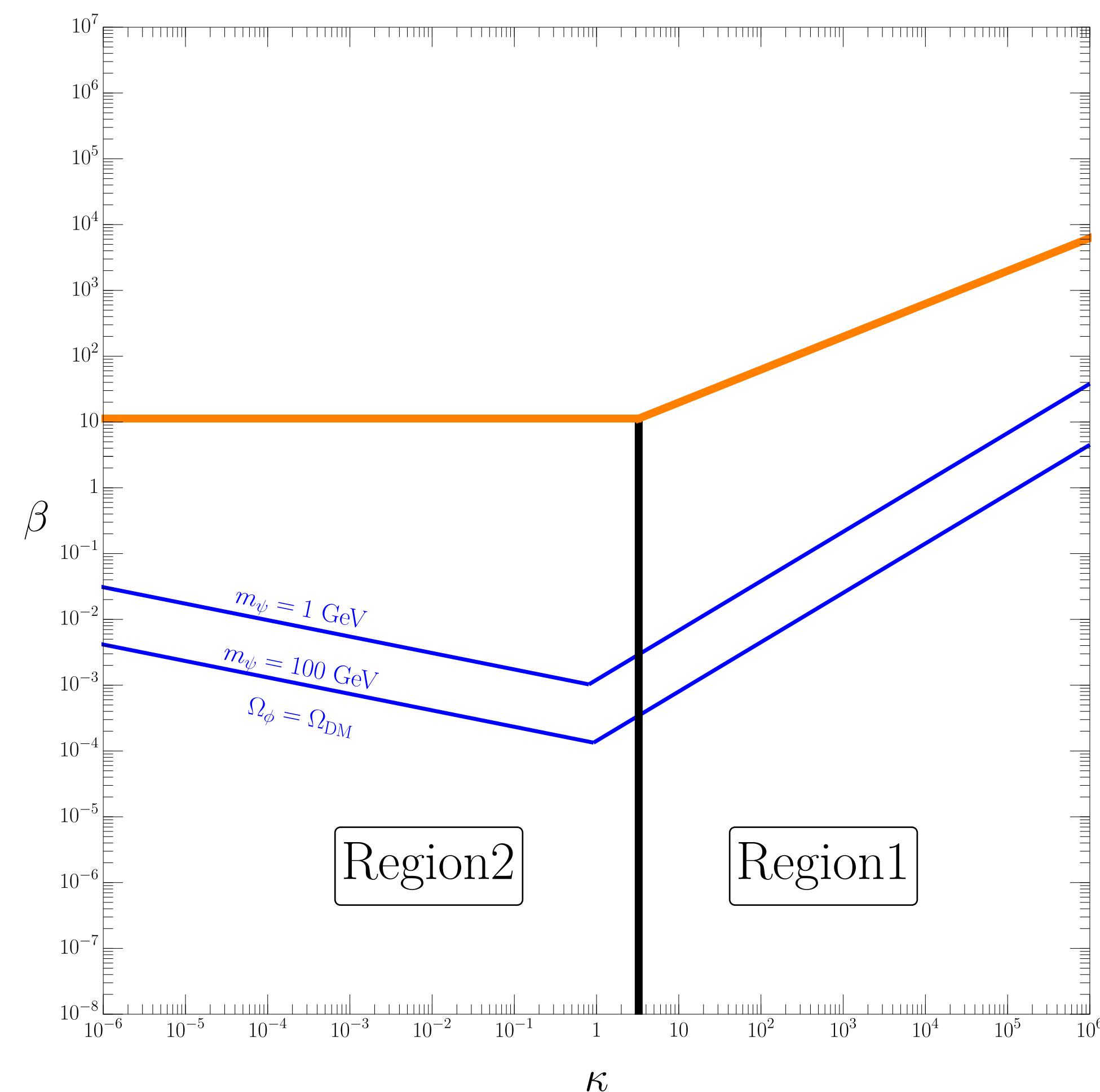
$$\Omega_\phi \simeq \Omega_{\text{DM}} \left( \frac{m_\psi}{0.1 \text{ GeV}} \right) \left( \frac{\beta}{10^{-3}} \right)^2 \left( \frac{\kappa}{0.01} \right)^{1/2} \left( \frac{10.75}{g_*^{\text{osc}}} \right)^{9/4}$$



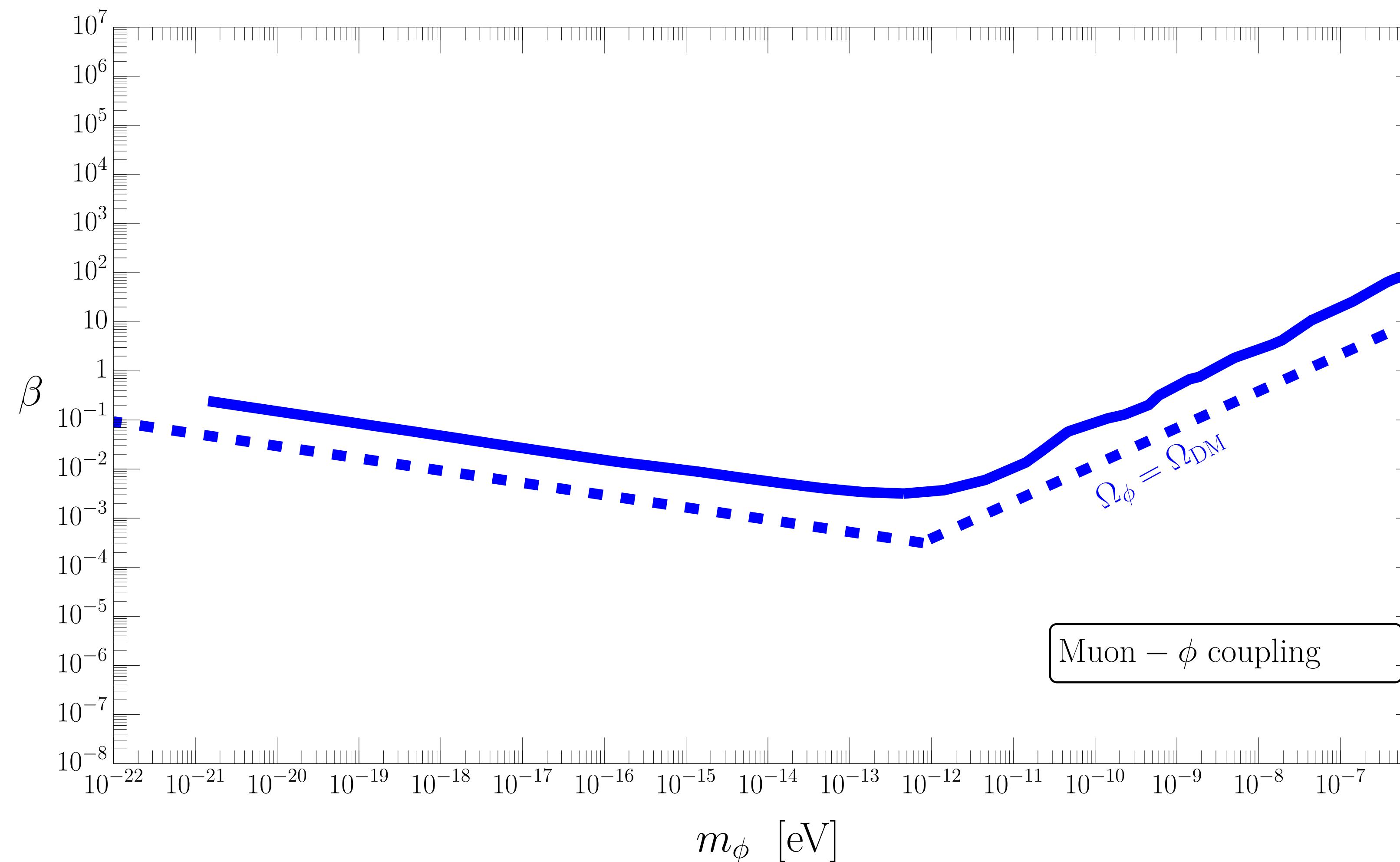
# Parameter space exploration

## Small $\kappa$ and small $\beta$ , Plots

- Combined Plot

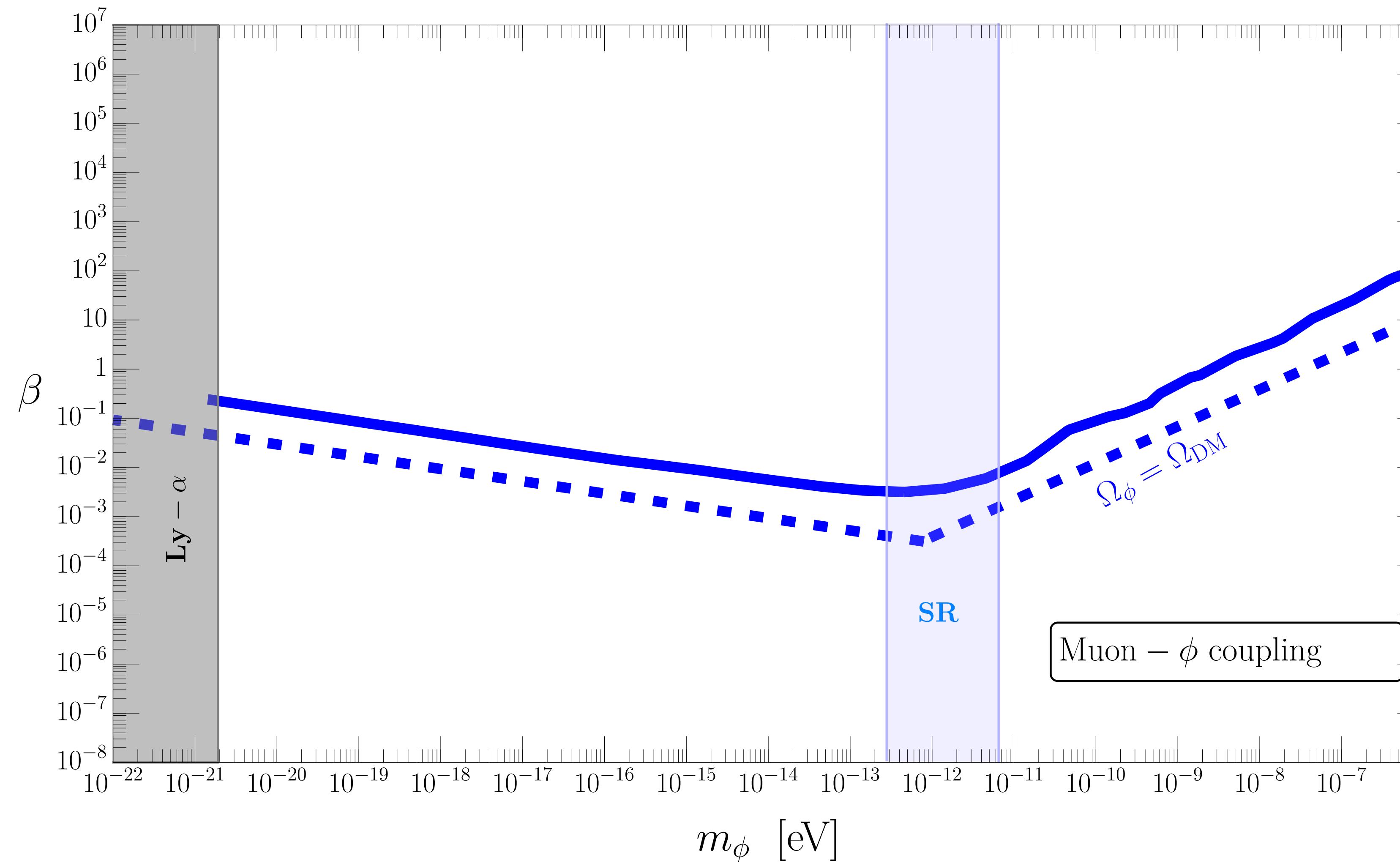


# Muon Case



# Constraints

## Ly- $\alpha$ and Superradiace



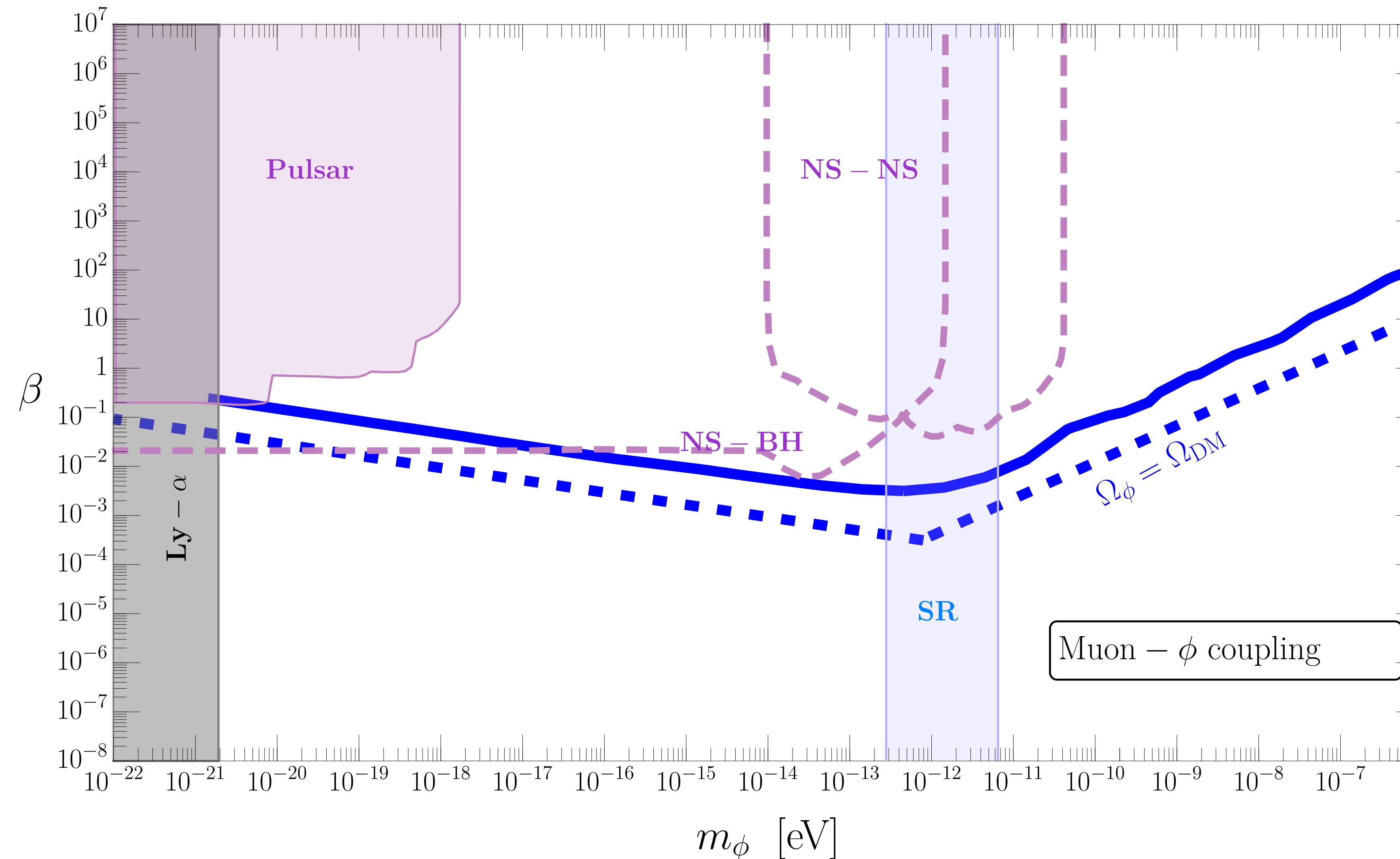
Irsic et. al. : 1703.04683

Baryakhtar et. al. : 2011.11646

# Constraints

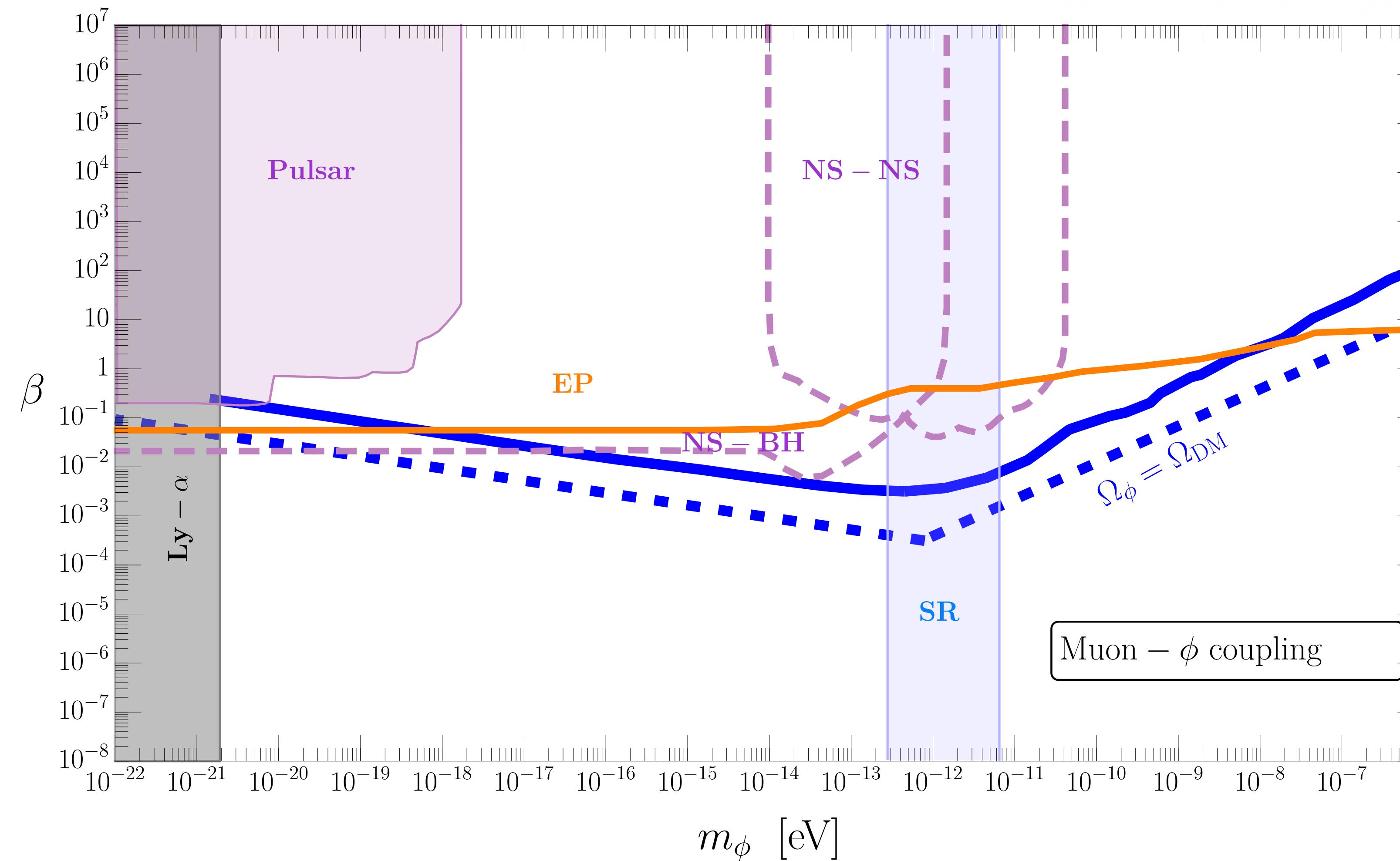
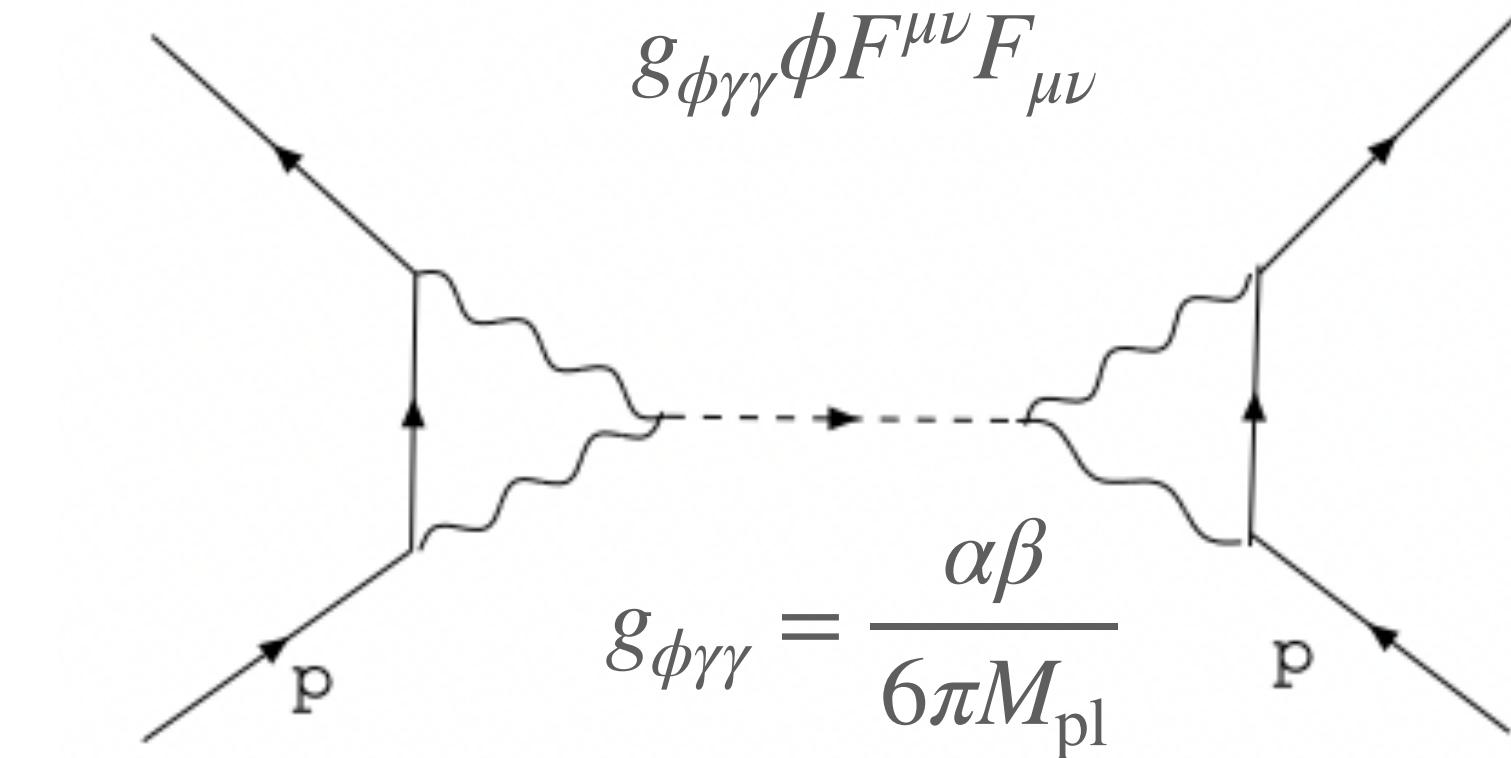
## Astrophysical

Dror et. al. : 1909.12845



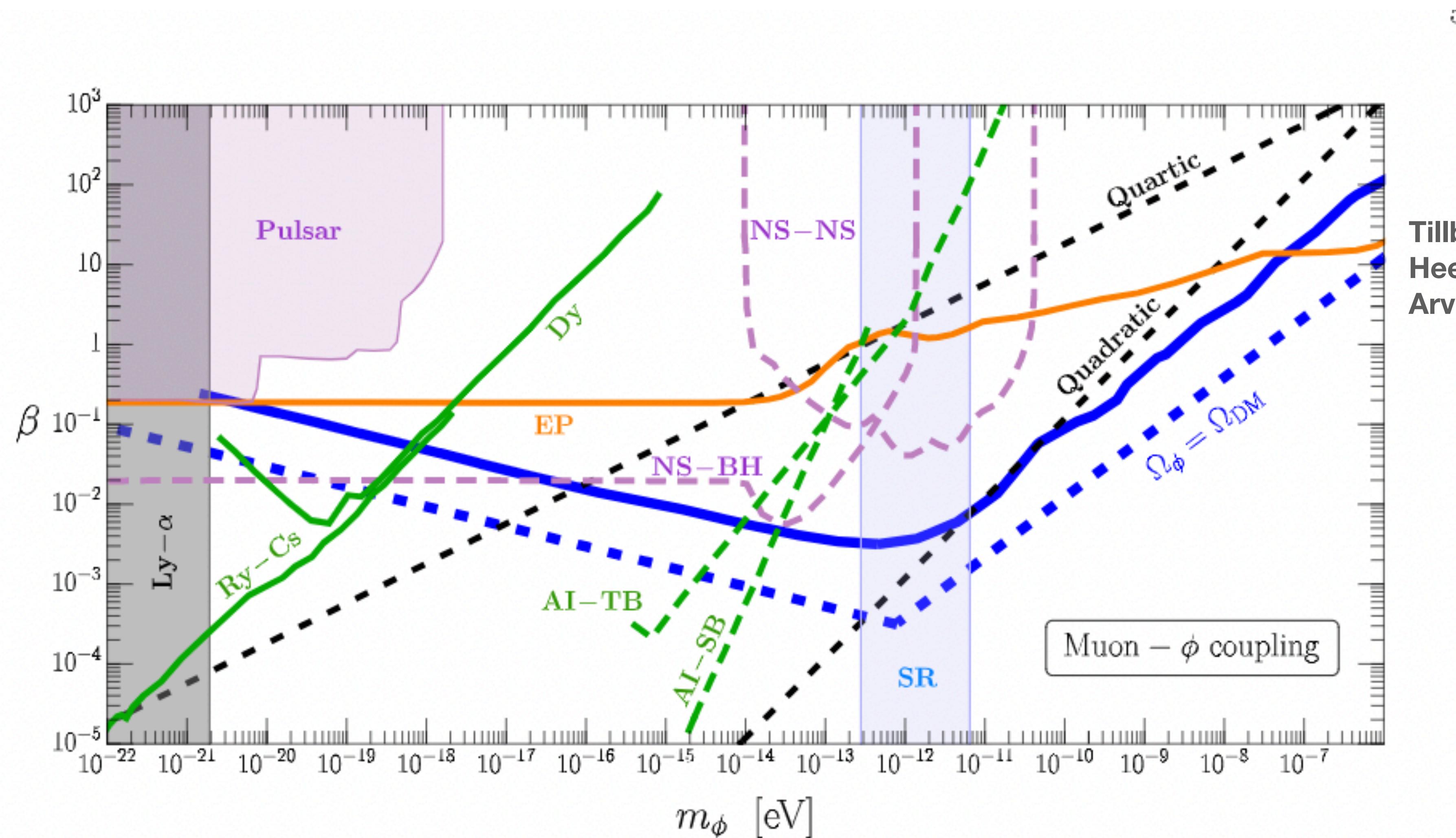
# Constraints

## From coupling to photon, EP



# Constraints

## Atomic Clocks, Interferometers

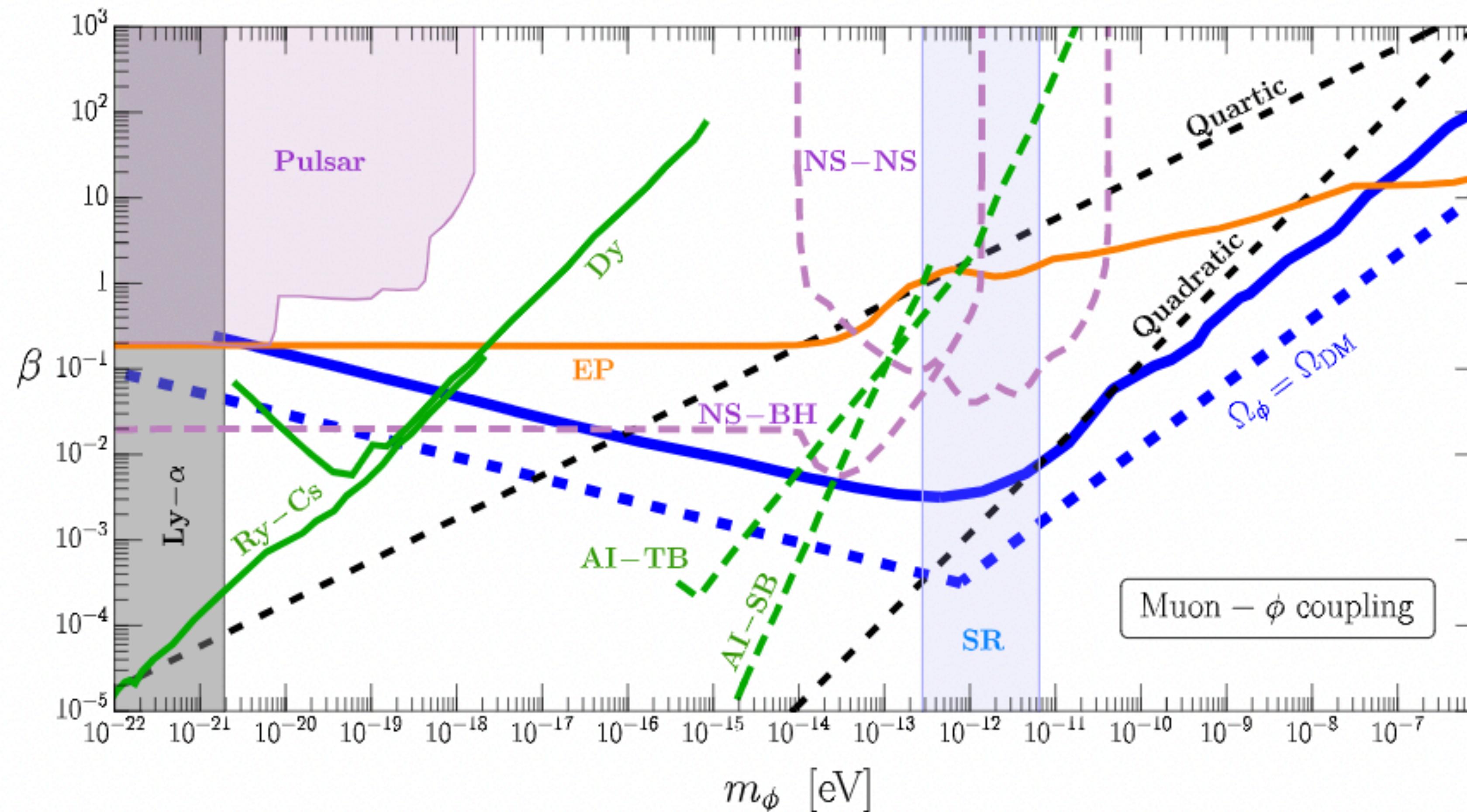


# Constraints Naturalness

Quartic  $\lambda A_{max}^4 > m_A^2 A_{max}^2$

Quadrartic  $\delta m_A^2 = \frac{1}{16\pi^2} \left( \frac{\beta m_f}{M_{pl}} \right)^2 \Lambda_{EW}^2 < m_A^2$

5



# **Constraints**

## **Other constraints that can be avoided**

- Fine-tuning of initial conditions
- Isocurvature constraints
- Decay, Freeze-in and thermalization

# Conclusion

## Things learned so far

- Generic couplings to fermions can generate a misalignment for scalar fairly independently of initial conditions
- Allows us to relate scalar mass and couplings to get required DM density
- Mechanism within reach of being confirmed/ruled out for muons
- Assuming standard cosmology provides strongest constraints on scalar muon coupling in much of the parameter space

# Ongoing Work

## Things to investigate

- Explore Region 3 in depth
- Explore thermal feedback from coupling to photons
- Explore electron and tau lepton parameter space.
- Higgs portal
- Neutrino portal
- UV completion of model
- Others ...

# Questions?

# Backup Slides

# Parameter space exploration

## Large $\beta$ and $\kappa$ , Region 3

- Oscillations start controlled by  $\beta$ , scalar oscillates around

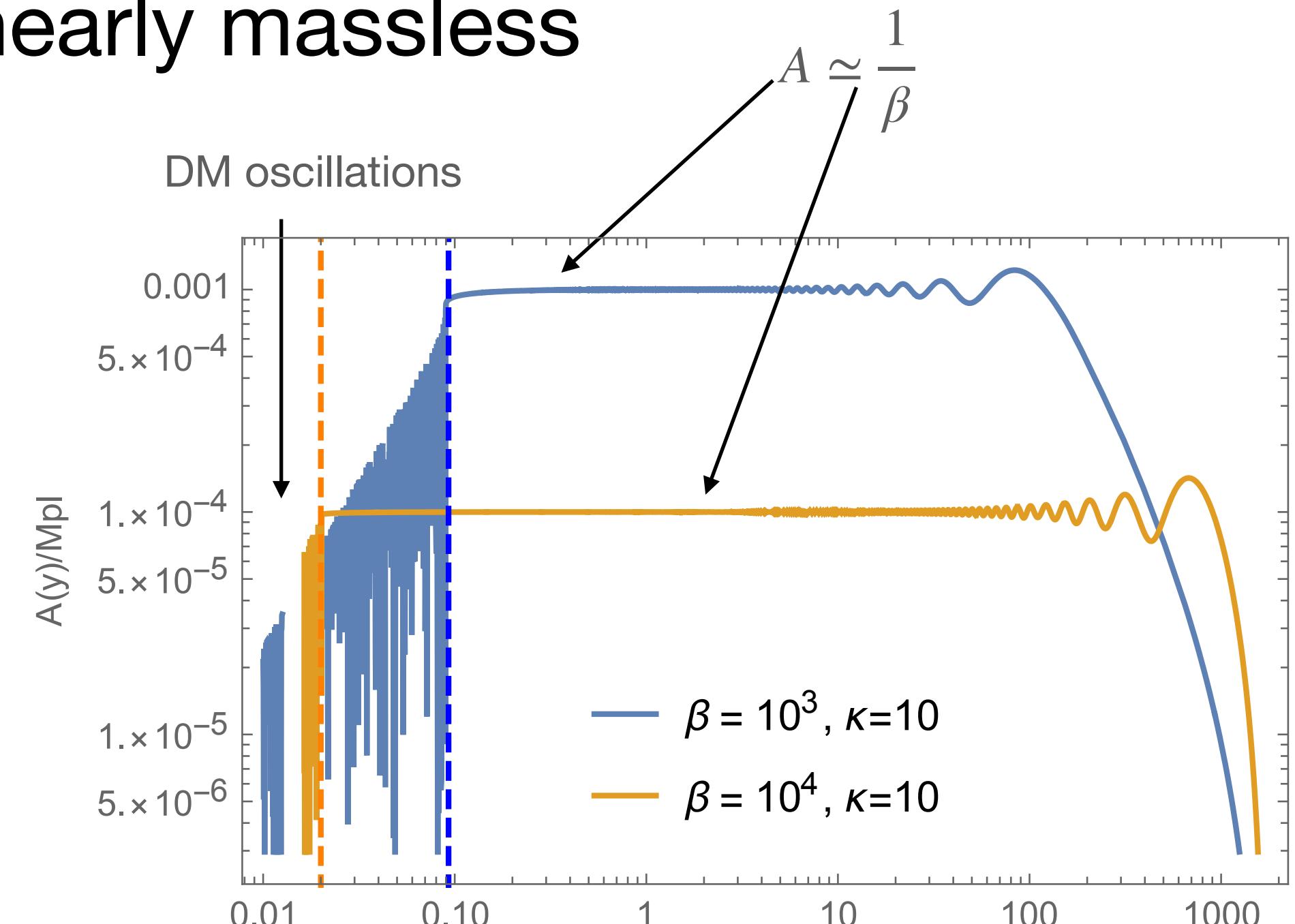
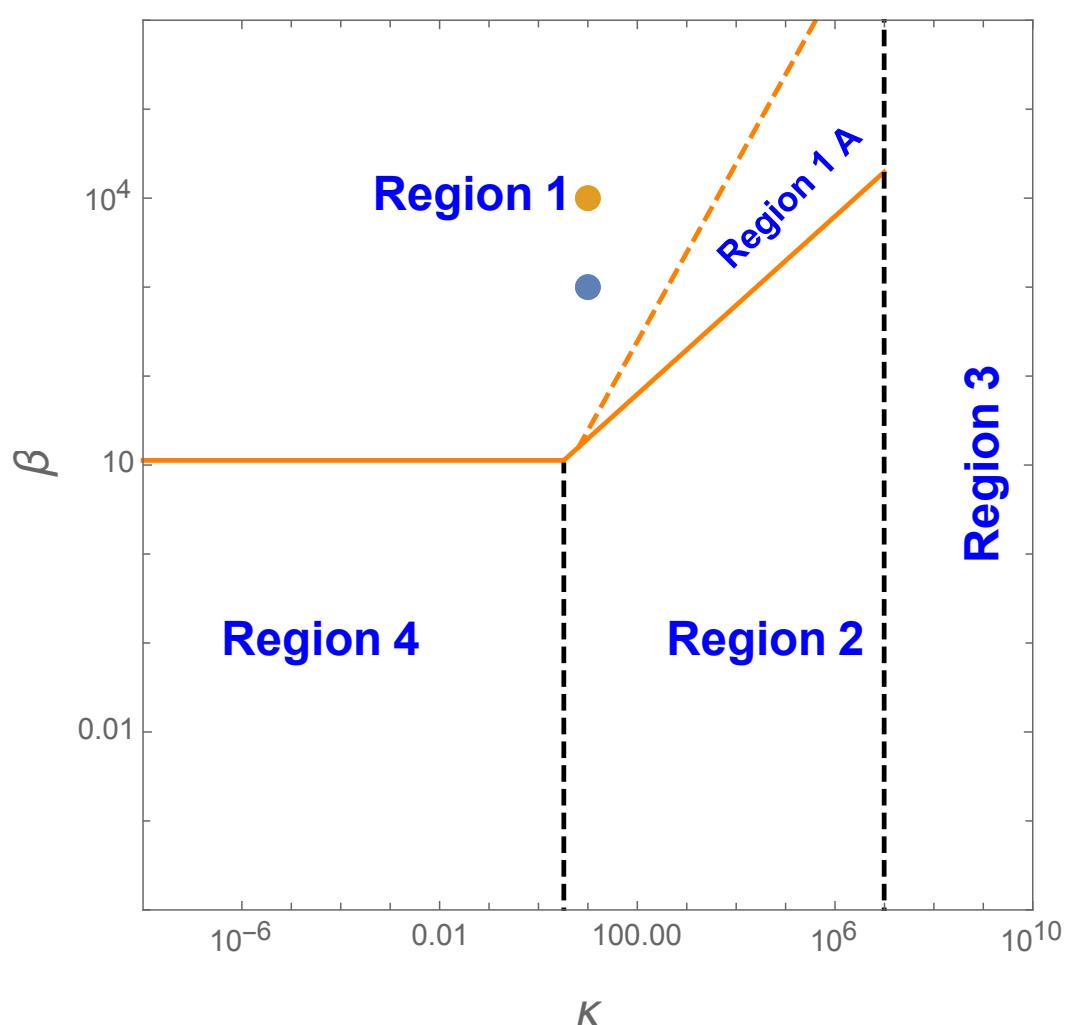
$$A_{min}(y) = \frac{y^2\beta}{y^2\beta^2 + 6\kappa^2} \Big|_{y \gg 1} \simeq \frac{1}{\beta}$$

- Scalar oscillates around  $1/\beta$ , making fermions nearly massless

$$y_{eff} = \frac{T}{m \left( 1 - \frac{\beta A}{M_{pl}} \right)} \simeq \frac{y(y^2\beta^2 + 6\kappa^2)}{6\kappa^2}$$

- Once  $y_{eff} < 1$ , the feedback shuts off at

$$y_d \simeq 2 \times \left( \frac{\kappa}{\beta} \right)^{2/3} + \mathcal{O} \left( \frac{\kappa}{\beta} \right)^{4/3}$$



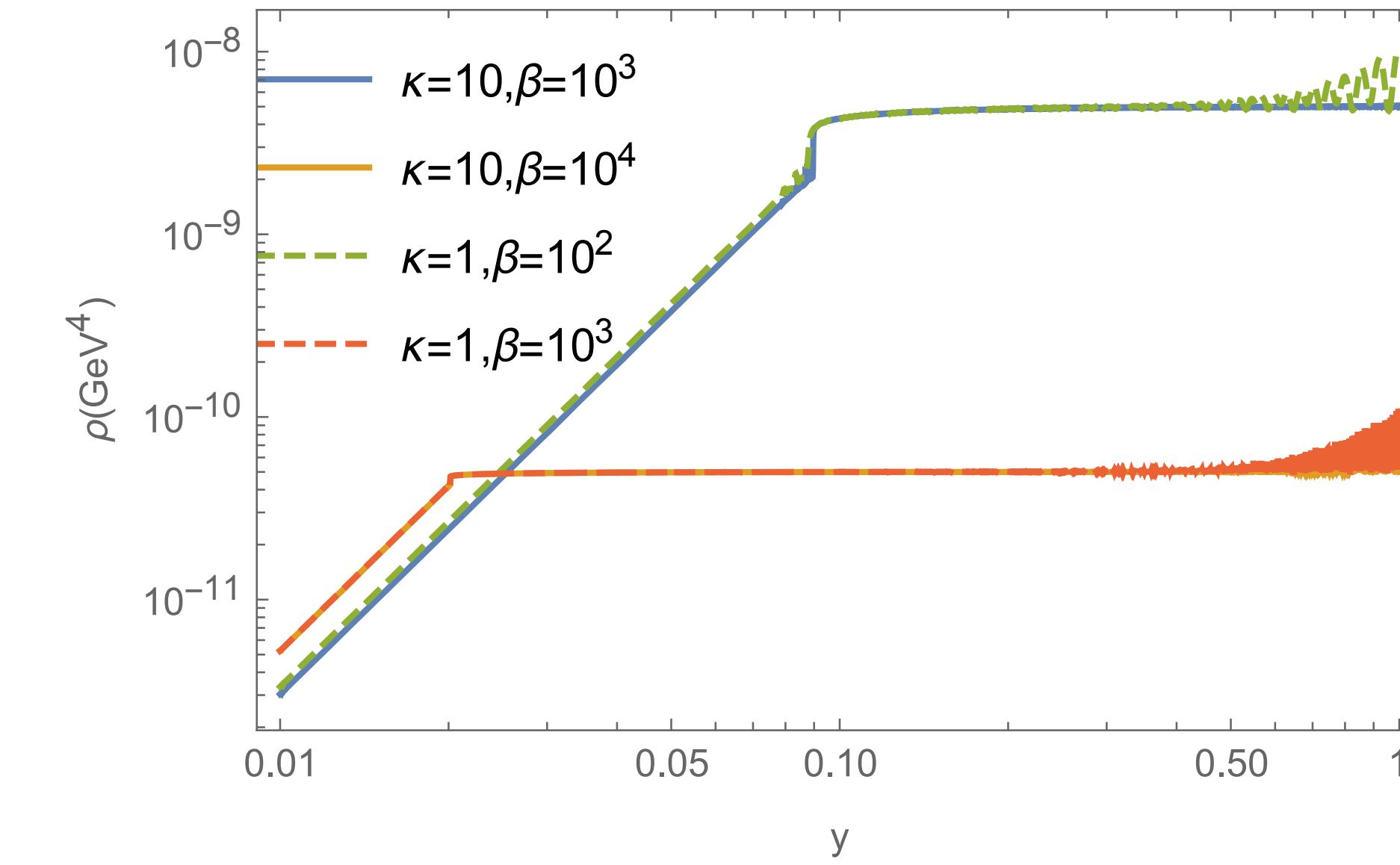
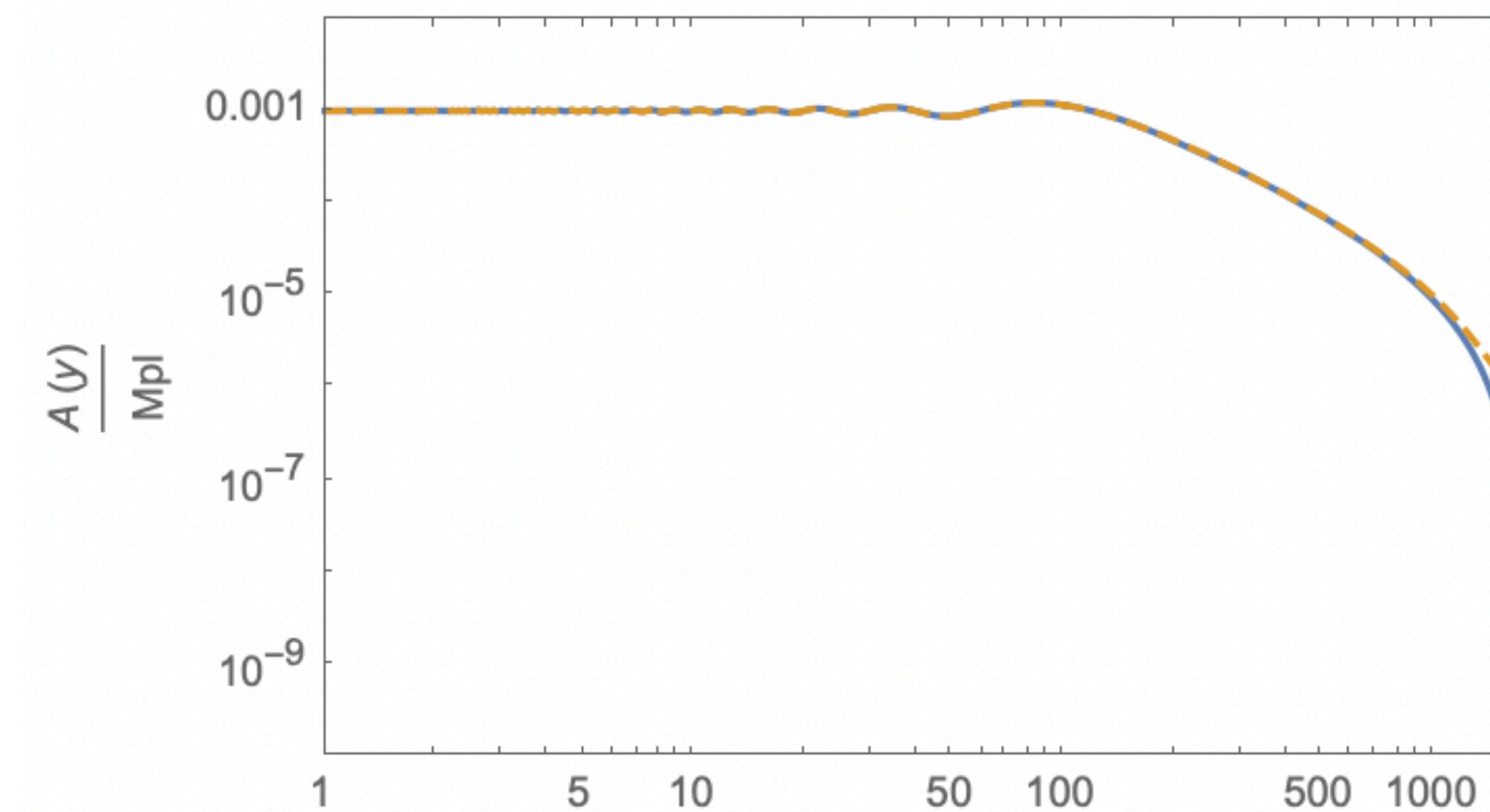
# Parameter space exploration

## Large $\beta$ and $\kappa$ , DM density, Plots

- Expected energy density is independent of  $\kappa, \beta$ . Too large even for electron mass

$$\Omega_{eq} = \xi \frac{\frac{1}{2} m_A^2 A_{min}^2(y_d)}{\frac{\pi^2 g(y_d)}{30} y_d^4 m_f^4} \frac{y_d}{y_{eq}} = 2.8 \times 10^5 \left( \frac{10.75}{g(y_d)} \right) \left( \frac{\xi}{0.1} \right) \left( \frac{m}{0.1 \text{GeV}} \right)$$

$\frac{1}{\beta^2}$



# Parameter Space exploration

## Large $\beta$ and $\kappa$ , Region 1 A

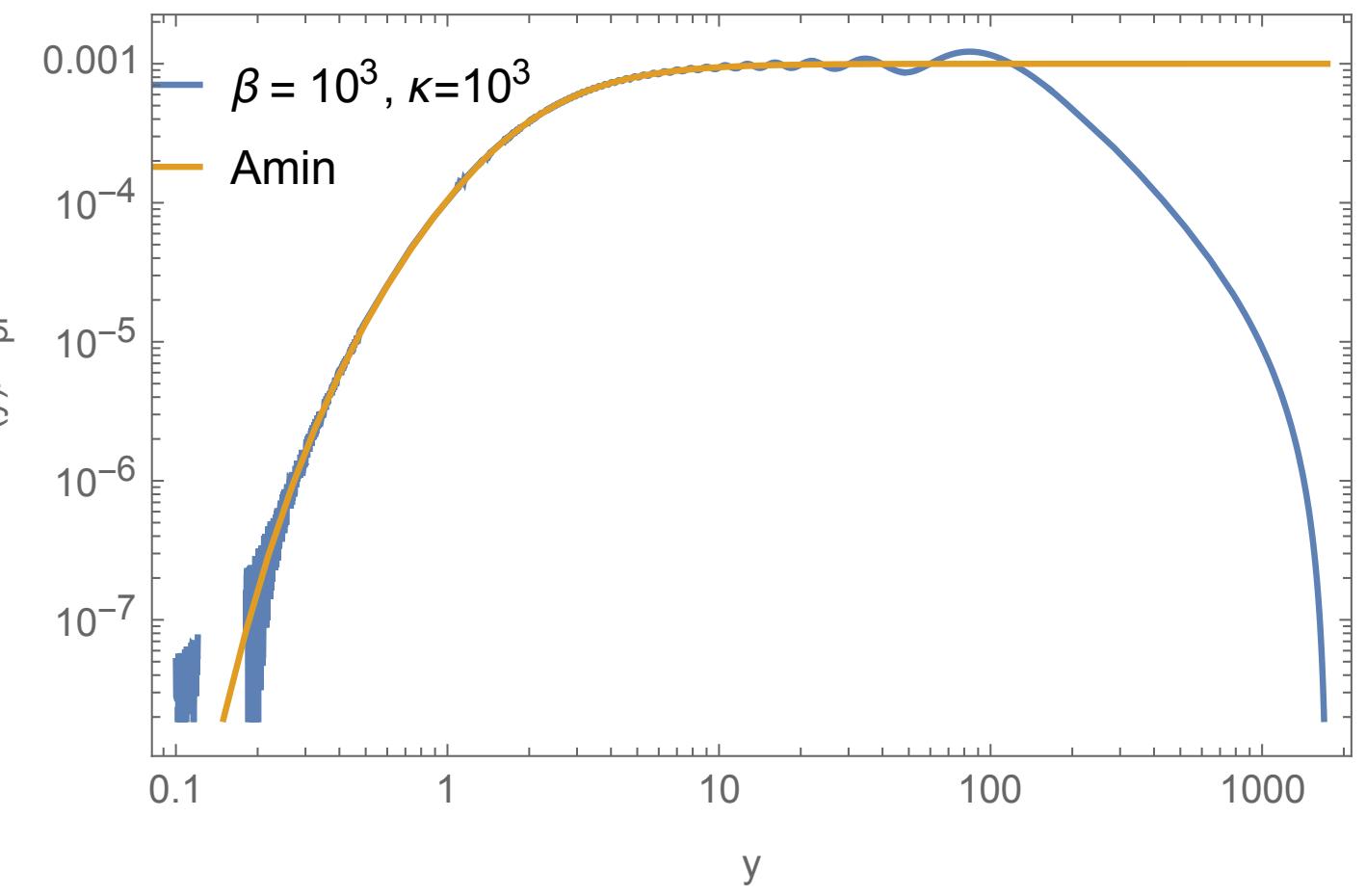
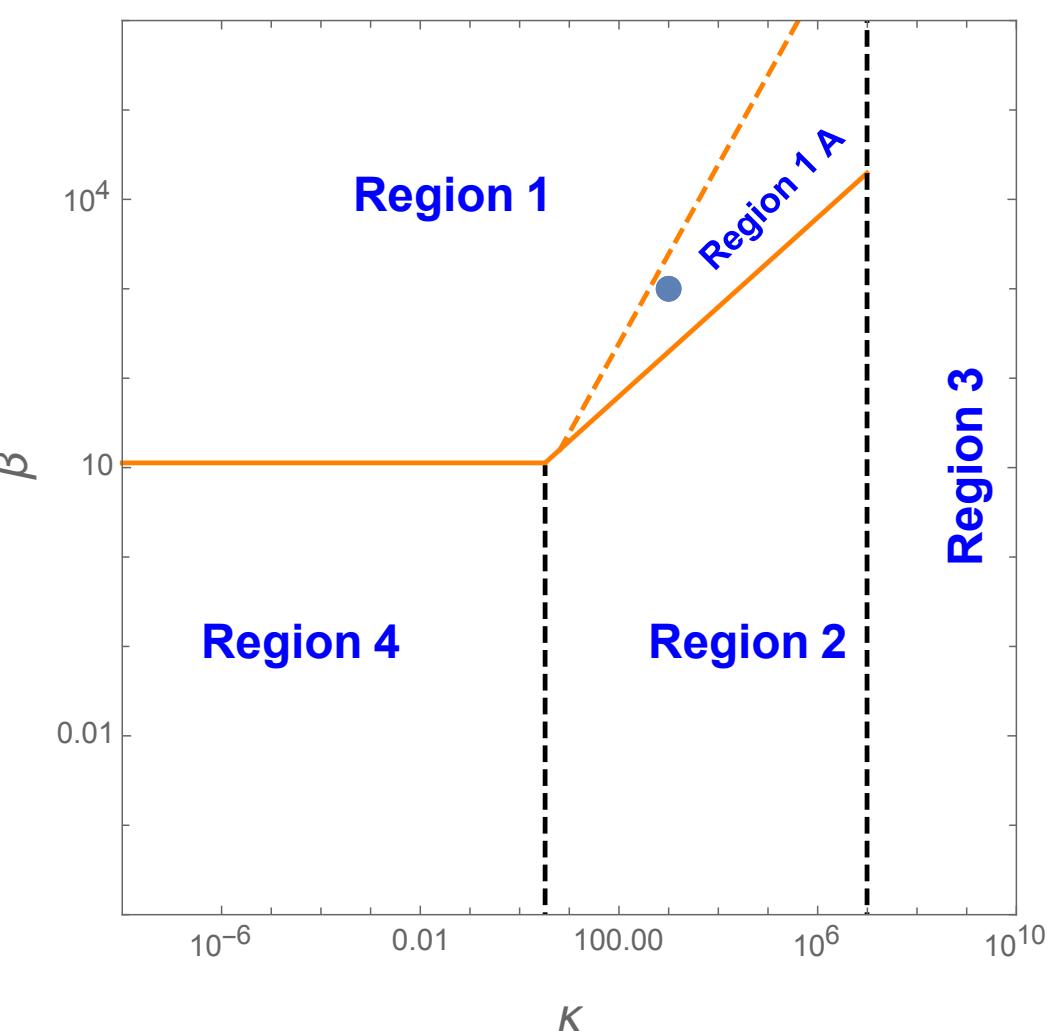
- Oscillations start controlled by  $\beta$ , scalar oscillates around

$$A_{min}(y) = \frac{y^2\beta}{y^2\beta^2 + 6\kappa^2} \Big|_{y \gg 1} \simeq \frac{1}{\beta}$$

- However since  $\kappa \sim \beta$  in this region  $y_d \simeq 2 \times \left(\frac{\kappa}{\beta}\right)^{2/3}$  does not hold
- Instead scalar starts to oscillate like matter at  $y_{osc,DM} \simeq \sqrt{6} \frac{\kappa}{\beta}$
- DM density is given by

$$\Omega_{eq} = \frac{\frac{1}{2}m_A^2 \left( A(y_{osc,DM}) - A_{min}(y_{osc,DM}) \right)^2}{\frac{\pi^2 g_*}{30} y_{osc,DM}^4 m_f^4} \frac{y_{osc,DM}}{y_{eq}} \simeq 8.6 \times 10^8 \left( \frac{\kappa}{\beta^3} \right) \left( \frac{m}{0.1 \text{GeV}} \right)^{A(y)/M_{pl}}$$

- Correct relic density not obtained for muons



# Parameter space exploration

## Large (Huge) $\kappa$ and small $\beta$ , Analytical understanding

- Our fermion feedback starts at  $y = y_i = \frac{T_{EW}}{m_f}$
- For huge  $\kappa$  oscillations start before EW transition, hence before feedback

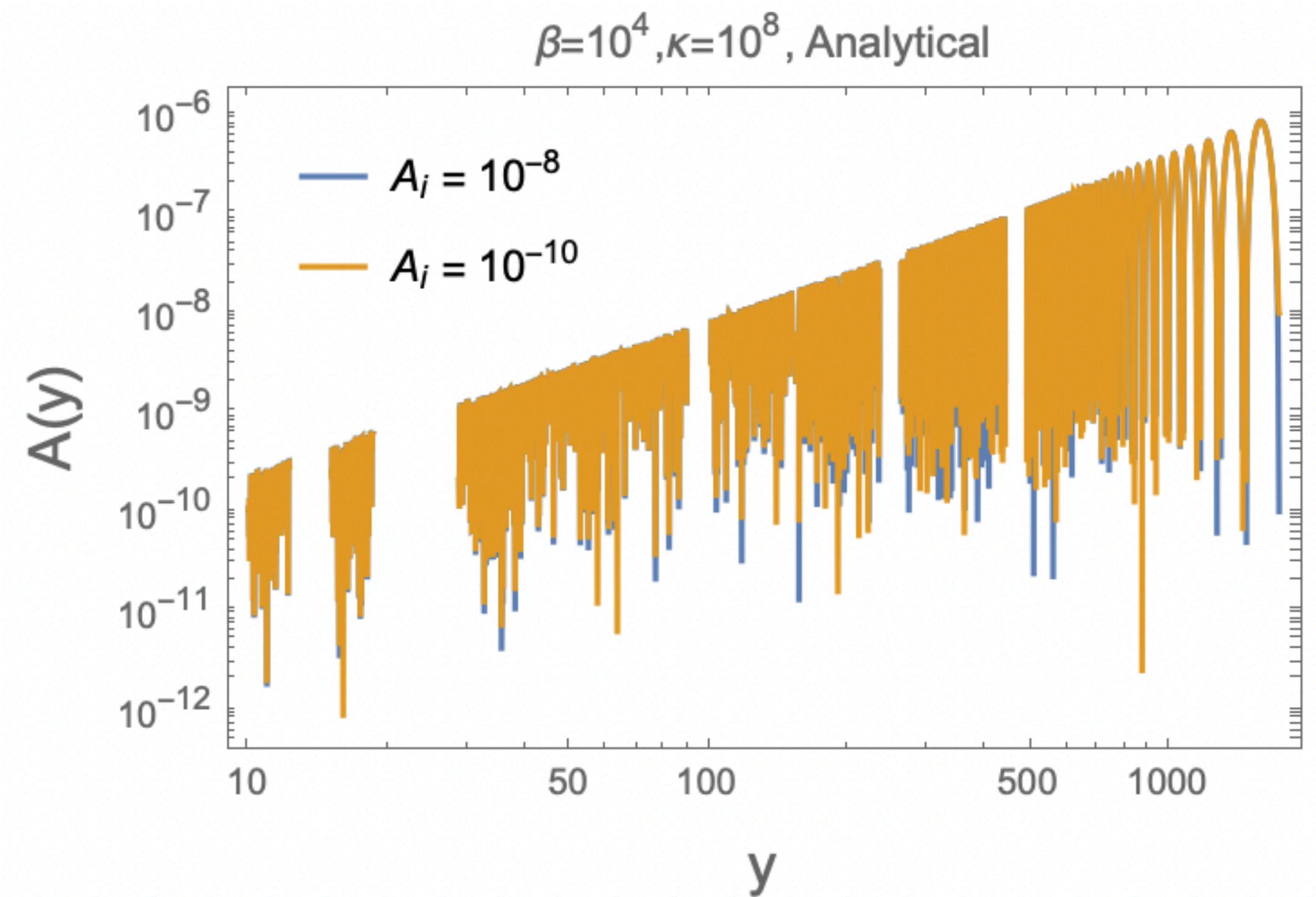
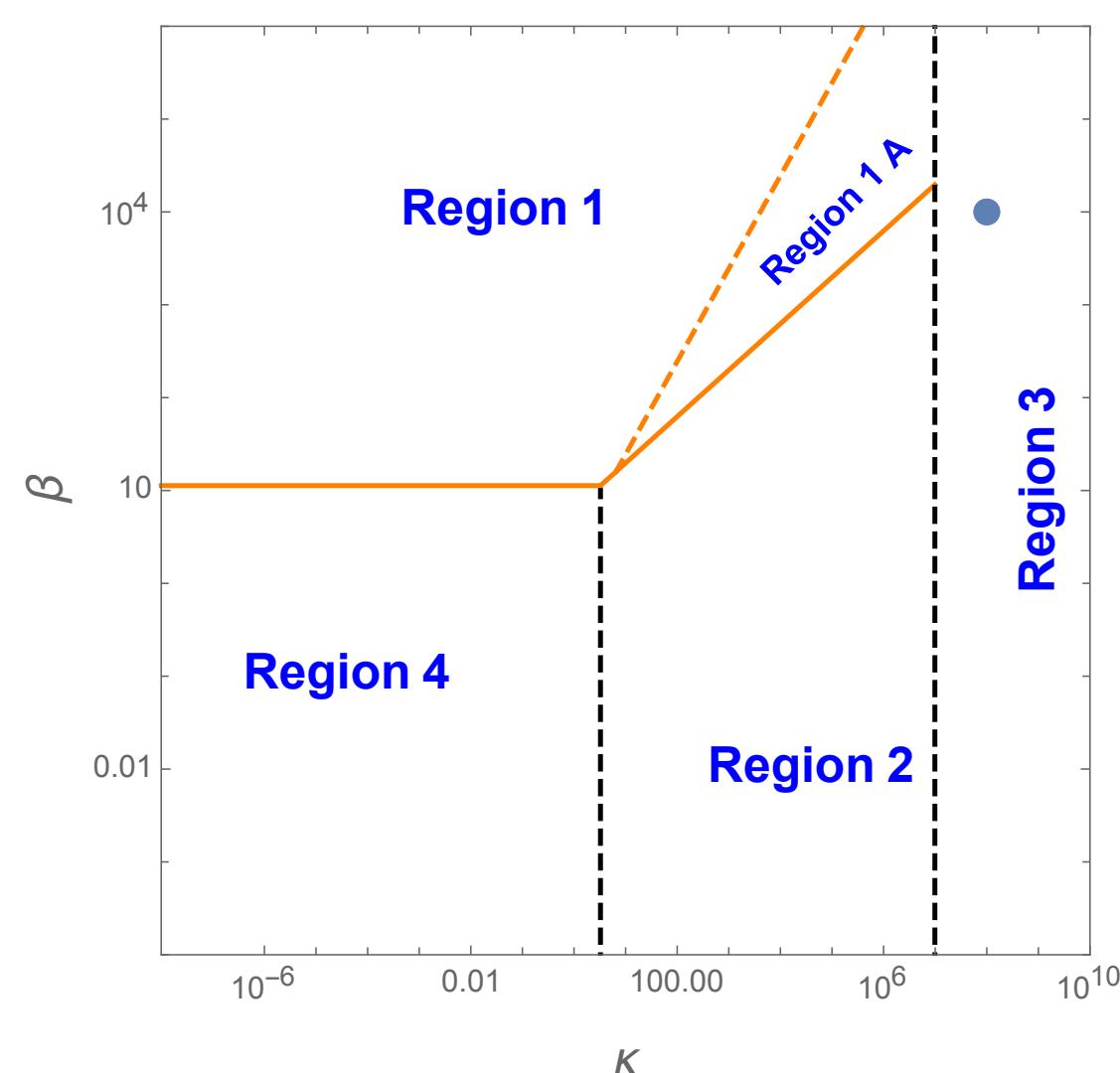
$$y_{osc} > y_i \rightarrow \kappa > 3\gamma y_i^2$$

- Can analytically solve in this regime

$$A(y < y_i) = \frac{y^2 \beta}{6\kappa^2} + \left(\frac{y}{y_i}\right)^{3/2} \frac{(-\beta y_i^2 + 6\kappa^2 A_i)}{6\kappa^2} \cos(\zeta)$$

Minima

Amplitude of oscillation



# Parameter space exploration

## Large (Huge) $\kappa$ and small $\beta$ , DM density, Plots

- Initial condition needs to be small enough, i.e.  $A_i \ll \frac{\beta y_i^2}{6\kappa^2}$
- DM density given by

$$\Omega_{eq} = \frac{\frac{1}{2}m_A^2 A^2(y_i)}{\frac{\pi^2 g(y_i)}{30} y_i^4 m_f^4} \frac{y_i}{y_{eq}} = \frac{5y_i\beta^2}{12\pi^2 g_* y_{eq} \kappa^2} = \frac{8.6 \times 10^7 \beta^2}{\kappa^2} \left( \frac{100}{g(y_i)} \right)$$

