

Black hole entropy functions and spinning spindles

James Sparks

Mathematical Institute, Oxford

Based on work with Jerome Gauntlett & Dario Martelli

[and Pietro Benetti Genolini, Andrea Boido, Davide Cassani, Chris Couzens, Pietro Ferrero, Juan Pérez Ipiña]

Our programme over the last few years: low-dimensional SUSY AdS/CFT

Highlights:

- New geometric techniques \Rightarrow computation of observables in gravity via extremal problems, without solving Einstein equations
- Allowed to identify dual SCFTs, including for old classes of solutions
- Rich infinite classes of new SCFTs, from branes at geometric singularities/wrapped over cycles, with exact matching of observables
- New supersymmetric compactifications, relation to accelerating black holes, and new perspective on black hole entropy functions

More concretely:

- $\text{AdS}_3 \times \mathbf{Y}$ solutions of IIB string theory, $\text{AdS}_2 \times \mathbf{Y}$ solutions of M-theory
- Generated by a single type of brane (D3, M2), flux \mathbf{G} = differential form
- Supersymmetry with $U(1)_R$ symmetry \Rightarrow Killing vector ξ on \mathbf{Y}

Work around 2007 by [Gauntlett-Kim] *et al* constructed explicit solutions, studied local geometric structure on \mathbf{Y} , but lacking insight into field theory duals.

From 2018: introduced global geometric techniques, exploiting structure on \mathbf{Y} .

We introduce a **supersymmetric action**:

$$\mathbf{S} = \mathbf{S}(\mathbf{Y}, \mathbf{N}_a; \xi),$$

where $\mathbf{N}_a = \int_{\Sigma_a} \mathbf{G}$ are integer quantized fluxes, over cycles $\Sigma_a \subset \mathbf{Y}$.

Key point: depends only on *global* geometric data on \mathbf{Y} .

By construction, supersymmetric AdS solutions *extremize* \mathbf{S} , as a function of

$$\xi = \sum_{i=1}^n b_i \frac{\partial}{\partial \varphi_i} \quad \Rightarrow \quad \text{extremize over } \vec{\mathbf{b}} \in \mathbb{R}^n \quad \Rightarrow \quad \xi_*$$

- We introduced various geometric methods for computing \mathbf{S} in closed form, e.g. developing some new toric geometry [explicit example later]
- Extremal $\mathbf{S}_* = \mathbf{S}(\mathbf{Y}, \mathbf{N}_a; \xi_*)$ = central charge for AdS₃ solutions, entropy for AdS₂ solutions
- ξ_* = superconformal R-symmetry
- Other observables computed similarly, e.g. dimensions of BPS operators

Global description of $\mathbf{Y} \Rightarrow$ identification of SCFT duals \Rightarrow exact matching

Brief overview of AdS_3 results:

- $AdS_3 \times Y_7$ solutions in type IIB, D3-brane flux G_5
- Various explicit solutions, e.g. [Gauntlett-MacConamhna-Mateos-Waldram]
 - ▶ New methods \Rightarrow interpretation as D3-branes wrapped over a **spindle**
 - ▶ New way to preserve supersymmetry (not via a topological twist)
- Infinite classes of D3-branes at Calabi-Yau 3-fold singularities, further wrapped over (orbifold) Riemann surface (no explicit solutions in general)
- $S_* =$ central charge c , with general proof that the gravity result matches field theory central charge (via anomaly polynomials & c -extremization)

Rest of this talk: $\text{AdS}_2 \times \mathbf{Y}_9$ solutions \Rightarrow near horizon of extremal black holes.

- We consider \mathbf{Y}_9 of fibred form

$$\mathbf{X}_7 \hookrightarrow \mathbf{Y}_9 \rightarrow \Sigma,$$

where $\Sigma =$ (orbifold) Riemann surface

- Near horizon of extremal black hole in $\text{AdS}_4 \times \mathbf{X}_7$, horizon Σ
- Fluxes \mathbf{N}_a determine magnetic charges of black hole
- $\mathcal{S} = \mathcal{S}(\mathbf{Y}_9, \mathbf{N}_a; \xi) =$ **entropy function**, similar in spirit to Sen, $\mathcal{S}_* =$ entropy

$\text{AdS}_4 \times \mathbf{X}_7$ black holes:

- Reduction on $\mathbf{X}_7 \Rightarrow$ massless $U(1)$ gauge fields \mathbf{A}_a in AdS_4
- Different $\mathbf{X}_7 \Rightarrow$ different SCFT_3 , with black holes in the bulk
- Expect general families of supersymmetric black holes with magnetic & electric charges \mathbf{p}_a , \mathbf{q}_a , angular momentum \mathbf{J} , and **acceleration**
- Acceleration is a new ingredient \Rightarrow different horizon topologies (spindles)

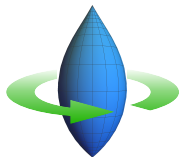
What is known?

- For $\mathbf{X}_7 = \mathbf{S}^7$ and horizon $\Sigma = \mathbf{S}^2$, general dyonic rotating black holes may be constructed in STU gauged supergravity, entropy matched to a dual microstate counting [[Benini-Hristov-Zaffaroni](#)]
- For general \mathbf{X}_7 , class of rotating, accelerating black holes with single electric & magnetic charge constructed in [[Ferrero-Gauntlett-Ipiña-Martelli-JFS](#)]

It seems hopeless to construct general black hole solutions explicitly.

But can we compute their entropy, and understand dual microstate counting?

Well-known in GR that acceleration \Rightarrow conical deficit angles along horizon Σ .



Provided these are *quantized* ($2\pi/n_{\pm}$ at each “pole”), so Σ = orbifold known as a *spindle*, $\text{AdS}_4 \times \mathbf{X}_7$ solution has no conical deficit angles!

This has to do with supersymmetry relating acceleration to magnetic charge, and the latter twists \mathbf{X}_7 over the black hole, removing the deficit angles.

General lesson: lower-dimensional singular solutions can have non-singular uplifts!

Our supersymmetric action $\mathcal{S}(\mathbf{Y}_9, \mathbf{N}_a; \vec{\mathbf{b}})$ allows to compute black hole entropy:

- With zero electric charge & angular momentum (although \exists conjecture for how to include these)
- Magnetic charges $\mathbf{p}_a = \mathbf{N}_a / \mathbf{N}$, where $\mathbf{N} = \int_{\mathbf{X}_7} \mathbf{G}_7 =$ number of M2-branes
- Recall \mathbf{Y}_9 is \mathbf{X}_7 fibred over horizon Σ . We have a general “gravitational block” formula for \mathcal{S} , with contributions coming from each “pole” of Σ
- This remarkably agrees with guesses for entropy functions e.g. for STU black holes, but is both *derived*, and much more general!

Example:

$$\mathbf{S} = \frac{1}{b_0} \left(\frac{1}{\sqrt{\text{Vol}(\mathbf{X}_7)(\vec{b}_+)}} - \frac{\sigma}{\sqrt{\text{Vol}(\mathbf{X}_7)(\vec{b}_-)}} \right) \frac{8\pi^3 \mathbf{N}^{3/2}}{3\sqrt{6}}$$

where $\sigma = \pm 1$ and

- b_0 = component of R-symmetry vector \vec{b} rotating horizon Σ
- $\vec{b}_\pm = \vec{b} \mp \frac{b_0}{n_\pm} \vec{v}_\pm$ (\vec{v}_\pm determined by magnetic charges \mathbf{p}_a /fibration)

Can also be written as

$$\mathbf{S} = \frac{4}{b_0} \left[\mathcal{F}_{S^3}(\vec{b}_+) - \sigma \mathcal{F}_{S^3}(\vec{b}_-) \right]$$

Relation to other approaches to black hole entropy & thermodynamics?

Problem: Euclidean on-shell action of extremal black holes not well-defined.

Solution: look at *complex* family of supersymmetric but non-extremal black hole solutions, turning on chemical potential ω for rotation.

Remarkably, for the families of explicit black hole solutions we find

$$\mathcal{S} |_{b_0 = \omega/2\pi i} = -\text{Euclidean action of black hole!}$$

- Relates near horizon (IR) to bulk and boundary (UV)
- Extremizing right hand side gives Legendre transform \Rightarrow entropy

We have a conjecture for how this should work more generally.

Summary:

- New geometric approach to computing observables in gravity
- Infinite classes of AdS_3 & AdS_2 solutions \Rightarrow new 2d SCFTs with exact matching, near horizon limits of extremal black holes
- New compactifications on spindles, and relation to acceleration in GR

Some questions:

- Black strings for which the AdS_3 solutions are near horizon limits?
- Field theories on spindles/general singular spaces? Microstate counting?
- Entropy function \Leftrightarrow Euclidean black hole action via fixed point theorem?